

A New Approach in Coal Mine Exploration Using Cosmic Ray Muons

Reza DARIJANI¹, Ali NEGARESTANI², Mohammad Reza REZAIE²,
Syed Jalil FATEMI³, and Ahmad AKHOND¹

¹Department of Physics, Payame Noor University, Tehran, Iran

²Department of Physics, Kerman Graduate University of Technology, Kerman, Iran
e-mail: mr.rezaie@kgut.ac.ir (corresponding author)

³Department of Physics, Shahid Bahonar University, Kerman, Iran

Abstract

Muon radiography is a technique that uses cosmic ray muons to image the interior of large scale geological structures. The muon absorption in matter is the most important parameter in cosmic ray muon radiography. Cosmic ray muon radiography is similar to X-ray radiography. The main aim in this survey is the simulation of the muon radiography for exploration of mines. So, the production source, tracking, and detection of cosmic ray muons were simulated by MCNPX code. For this purpose, the input data of the source card in MCNPX code were extracted from the muon energy spectrum at sea level. In addition, the other input data such as average density and thickness of layers that were used in this code are the measured data from Pabdana (Kerman, Iran) coal mines. The average thickness and density of these layers in the coal mines are from 2 to 4 m and 1.3 gr/cm³, respectively. To increase the spatial resolution, a detector was placed inside the mountain. The results indicated that using this approach, the layers with minimum thickness about 2.5 m can be identified.

Key words: muon radiography, coal mines, MCNPX.

1. INTRODUCTION

When the primary cosmic ray particles, protons and α particles, penetrate the atmosphere they collide with air nuclei and new particles, mostly pions and kaons are produced. These secondary particles are short lived, and a fraction of them decay into muons instead of interacting with nuclei deeper in the atmosphere. Because these cosmic ray muons weakly interact with air nuclei, and their mean lifespan is long (about 2.2×10^{-6} s), they are the dominant charged particles at the Earth's surface. They arrive at angles ranging from vertical to horizontal (Thompson and Whalley 1975), and their energy spectrum has been measured as a function of the zenith angle θ (Allkofer *et al.* 1981). It is also well known that muons arriving in the horizontal direction have a higher intensity at energies above 100 GeV. Because muons become highly penetrative as their energy becomes higher, these horizontal muons can be used for the purpose of probing the internal structure of gigantic objects. For example, typical horizontally arriving cosmic ray muons with energy of 1 TeV penetrate 2.6 km of water (Tanaka *et al.* 2007). Since muons penetrate through matter their energy decreases. The energy loss depends on the kinetic energy of the muon and, at very high energies, also on the chemical composition of the matter (Malmqvist *et al.* 1979). Some of muons that do not have enough energy to escape from the matter of interest are absorbed in the matter, and as a result the muon flux is reduced. Absorption of muons in matter is independent of geophysical models and directly measures the density length. Geophysical models can be used to predict a host of geophysical measurements and are important ways of calibrating regions in the absence of direct measurements. Meanwhile, geophysical models explain geophysical observations and determine the subsurface Earth properties. By measuring cosmic ray muons intensity in different depths of matter, very useful information regarding the density of matter can be obtained. This information is used for many purposes such as exploration of mines, and prediction of volcanic eruptions and earthquakes. The aim of this research is to obtain a two-dimensional profile of the average density inside a mountain by simulation. To this end, production, tracking, and detection of muons were performed by the MCNPX code. The input data such as the average density and thickness of layers that have been used in MCNPX code are the measured data from Pabdana (Kerman, Iran) coal mines.

2. THEORY

The potential use of high-energy near-horizontal cosmic ray muons for exploring the internal structure of large scale objects such as Egyptian pyramids has been recognized (Alvarez *et al.* 1970). The effective use of such near-horizontal cosmic ray muons for muon radiography can be explained through the following steps.

2.1 The cosmic-ray muon energy spectrum and its dependence on zenith angle

Pions and kaons are produced through nuclear interactions between primary cosmic ray protons and the atmospheric air, some of which decay and produce cosmic ray muons. The energy spectrum of these cosmic ray muons has been measured at different polar angles (Thompson and Whalley 1975). Furthermore, the flux of incident cosmic muons is of critical importance for geophysical imaging since it is used to determine the attenuation produced by the geological target (Lesparre *et al.* 2010). There are two ways to derive the differential flux of muons. The first approach is the use of a full simulation. These Monte Carlo computations can be accomplished by simulation codes like CORSIKA which allows to take into account the geomagnetic and altitude dependence. The second approach to obtain flux models is the use of fitting. The choice of a particular parameterization of the fitting curves may either be inspired by the physics involved in the production of muons from their parents or be guessed to provide a better fit regardless of the physical meaning of the parameters (Gaisser 1990). Some of fitting models only consider the production of muons from pions and kaons decay and assume a primary proton flux of the form: $P_0 e^{-\gamma}$ with $P_0 = 1.8 \text{ (cm}^{-2} \text{ sr}^{-1} \text{ GeV}^{\gamma-1}\text{)}$, and $\gamma = 2.7$. This approach yields the analytical form of the muon spectrum initially proposed by Bugaev *et al.* (1970) and popularized by Gaisser in the form of Eq. 1 (Lesparre *et al.* 2010).

$$\varphi_G(E_0, \theta) = A_G E_0^{-\gamma} \left(\frac{1}{1 + \frac{E_0 \cos \theta}{E_{0,\pi}^{\text{cr}}}} + \frac{B_G}{1 + \frac{E_0 \cos \theta}{E_{0,\kappa}^{\text{cr}}}} \right), \quad (1)$$

where A_G , γ , and B_G are the scale factor, the power index, and the balance factor, respectively. The balance factor depends on the ratio of muons produced by the kaons and pions. $E_{0,\pi}^{\text{cr}}$ and $E_{0,\kappa}^{\text{cr}}$ are critical energies of pions and kaons for the vertical incidence ($\theta = 0$).

2.2 Range of cosmic ray muons through standard rock

The energy loss of a charged particle with some energy E (TeV) on passage through matter with a density length (density \times path length) of X hectogram per cm^2 , $X \text{ (hg/cm}^2 = 100 \text{ gr/cm}^2\text{)}$, can be written as (Tanaka *et al.* 2007, Nagamine *et al.* 1995).

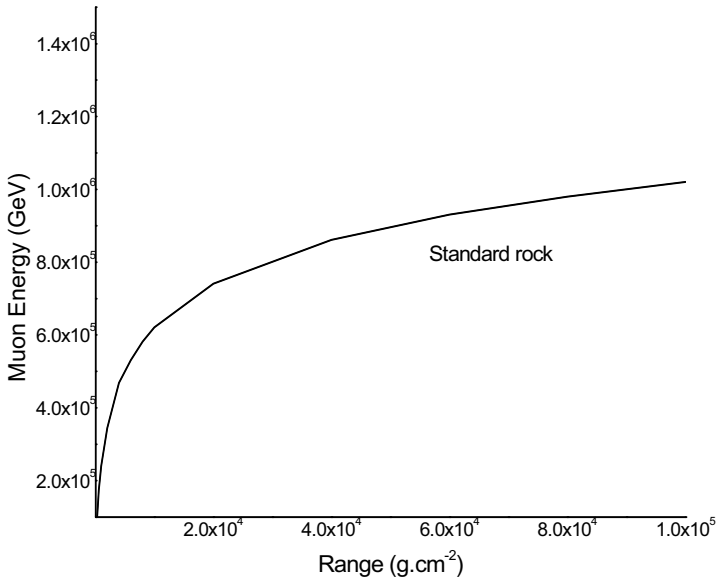


Fig. 1. The muon energy *versus* range in a standard rock.

$$\frac{dE}{dX} = \left[1.88 + 0.077 \text{Ln} \left(\frac{E}{m_\mu} \right) + 3.9E \right] \times 10^{-6} \text{ [TeVcm}^2\text{/gr]} , \tag{2}$$

where the first two terms represent ionization loss and the third term represents stochastic processes mainly due to bremsstrahlung. Choosing $E = 1 \text{ TeV}$ to yield a relatively small contribution of the logarithmic term, the mean rang (density length) X can be obtained by integrating the energy loss formula as Eq. 3 (Nagamine *et al.* 1995, Nagamine 2003).

$$X = 2.5 \times 10^3 \text{Ln}(1.56E + 1) \text{ [hg/cm}^2\text{]} . \tag{3}$$

This equation gives the relationship between range and energy in a standard rock. Standard rock is defined by $z/A = 0.5$, $z^2/A = 5.5$, $Z = 11$, and $\rho = 2.56 \text{ gr/cm}^3$, where Z , A , and ρ are the mean atomic number, the mean mass number, and density of the rock, respectively. Figure 1 shows the muon energy *versus* range in a standard rock.

2.3 Intensity of cosmic ray muons penetrating though rock with density length X

Measurement of muons intensity at various depths is important from different aspects. For instance, it provides information on the electromagnetic processes that reduce the flux (Bektasoglu *et al.* 2012). The measured muons

intensity at depth x in the Earth is equal to the intensity of muons at sea level with an energy greater than $E_c(x)$; disregarding the small fraction of muons which decay in the passage from the ground to the depth x , $E_c(x)$ is the minimum energy necessary for a muon to penetrate to the depth x (Malmqvist *et al.* 1979). Thus a unique relationship exists between density length X ($X = \rho x$) and the intensity of penetrating cosmic ray muons $N_\mu(E_c(x), \theta_z)$. Once X is given, the minimum energy $E_c(x)$ is determined through the X - E relation (Eq. 3), and the integrated flux $N_\mu(E_c(x), \theta_z)$ is given as Eq. 4.

$$N_\mu(E_c, \theta = 90 - \theta_z) = \int_{E_c}^{\infty} \frac{dN}{dE} dE . \quad (4)$$

Conversely, for a substance with an unknown density length X , the measurement of the muons flux, $N_\mu(\theta_z)$, penetrating through the substance with a zenith angel, θ_z , uniquely determines its density length in hg/cm^2 . Small changes in X due to the existence of a region of lower or higher density inside the broadly uniform object lead to differences in $N_\mu(\theta_z)$. The change in $N_\mu(\theta_z)$ informs us of change in X (Nagamine 2003). Figure 2 shows the penetrating muons flux as a function of rock thickness at different zenith angles.

Using Fig. 2, the thickness of rock (in km) is calculated by dividing rock thickness (in km) water equivalent by rock density.

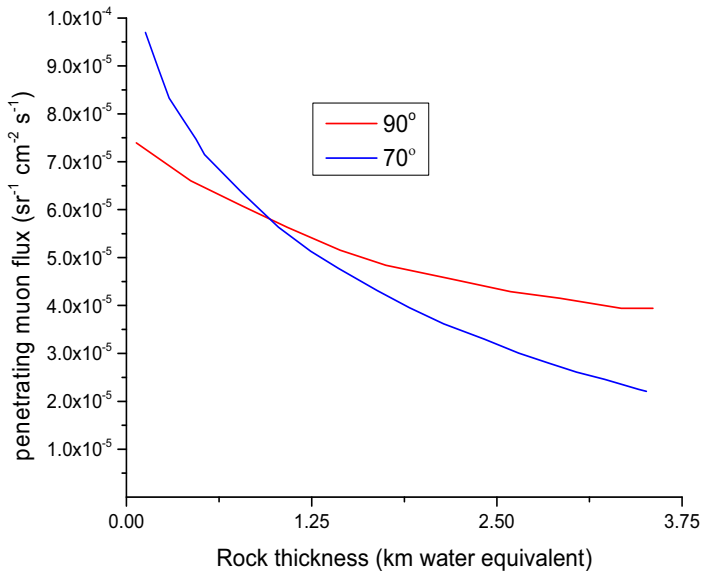


Fig. 2. Penetrating muon flux *versus* rock thickness (Nagamine 2003).

3. DETERMINATION COSMIC RAY MUON PATH THROUGH THE MATTER

For determination of the path of the muon penetrating through gigantic objects, there are two types of practical detection systems that are explained in the following.

3.1 Telescope of position sensitive detector

In this case, the straight line connecting the points where a cosmic ray muon passes through two (or three) counters determines the particle path (Fig. 3).

There are two representative examples of position determination that are explained in the following:

- The passing points can be determined by the difference in the arrival time of scintillation light from a set of at least three photomultipliers attached to the edges of each scintillator.
- The passing points can be determined more straightly by using segmented counter arrays in both horizontal and vertical directions (Fig. 3b).

In the case of counter telescopes of the position sensitive detectors, the spatial resolution of the detection system at the mountain ($\Delta X, \Delta Y$) can be determined from the resolution of the intersection points with each counter ($\Delta x, \Delta y$), the distance between the counters comprising the telescope (l), and the distance between the object (mountain) and the detector (L). This spatial resolution is calculated by the following Eqs. 5 and 6.

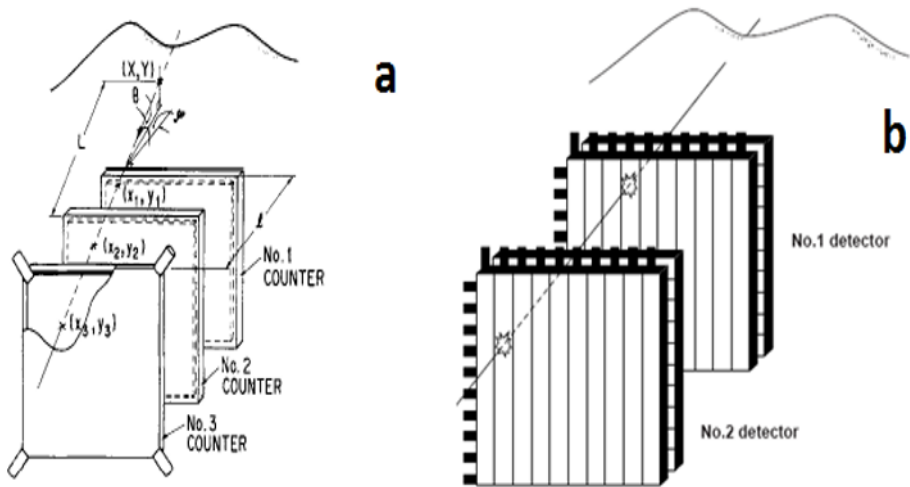


Fig. 3: (a) The counter telescope with three plastic scintillators, and (b) the detection system comprising two segmented plastic counters (Nagamine *et al.* 1995).

$$\Delta X = \frac{L}{l} \Delta x , \quad (5)$$

$$\Delta Y = \frac{L}{l} \Delta y . \quad (6)$$

3.2 Cerenkov light detector

In this case, the passage of the relativistic muon leaves a track of Cerenkov radiation in a selected gas or liquid which can be collected using a concave reflective mirror focused on an array of photomultipliers (Nagamine *et al.* 1995, Tanaka *et al.* 2003).

4. KERMAN COAL MINES

Kerman coal mines are situated in the east of central Iran at the distance between 50 and 220 km north west of Kerman. At these areas there are 9 coal ores that the Main Pabdana and the South Pabdana are active and important mines.

4.1 Main Pabdana coal mine

This coal mine is around 9 km long and 1.1 km wide, and its area is about 10 km². It consists of 14 layers of coal which are located at different depths and have different thicknesses. Their maximum thickness is of about 4 m, while for some others the thickness is about 2 m. The average density of coal and that of rock in this mine are about 1.3 and 2.65 gr/cm³, respectively.

4.2 The South Pabdana mine

This mine like the Main Pabdana was explored from 1966 to 1977. Its area is about 2 km². It consists of 13 layers, each one having an average thickness between 1 and 3.15 m and the average density of 1.3 gr/cm³.

5. MCNPX CODE

MCNPX is a general-purpose Fortran 90 Monte Carlo radiation transport code that transports more than 34 particle types including neutrons, protons, electrons, muons, *etc.* MCNPX is the product of computer codes that are used for modeling, the interaction of particle with matter, radiography, exploration of mine and oil, medical physics, and shielding. MCNPX has many capabilities that enhance physics, sources, tallies, and graphics (Pelowitz 2008) making the results of simulations more accurate. The simulation performed in the present study has many stages that are explained in the following.

5.1 Geometry

In this study geometry is defined as a cone that is 200 m high, parallel to the z-axis, and the opening angle of the cone is 45 degrees. This cone is inter-

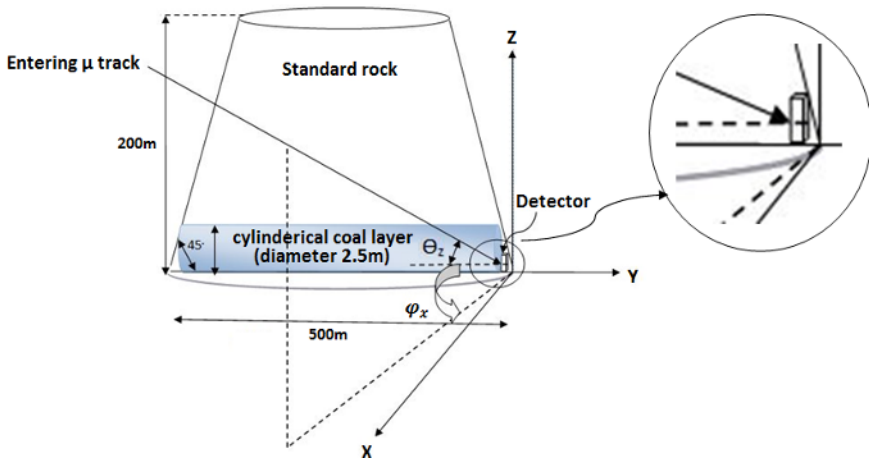


Fig. 4. Simplified model for mountain-detector and polar coordinate system.

sected by two planes parallel to xy plane at $z = 0$ and $z = 200$ m, respectively. Inside this cone, there is a cylindrical layer of coal with diameter of 2.5 m which is located at a depth of 200 m from the mountain top (cone) with a width of 500 m; the remaining volume of the cone is standard rock. The coal was replaced by carbon that is defined by $z/A = 0.5$, $z^2/A = 3$, $Z = 6$, and $\rho = 1.3 \text{ gr/cm}^3$, where Z , A , and ρ are the mean atomic number, the mean mass number, and density of the coal, respectively. A detection system with an area of about 1 m^2 , perpendicular to the y axis, was located at the foot of cone. The geometry and detector that were used in MCNPX code have been shown in Fig. 4.

5.2 Source card

The source card of MCNPX code includes all necessary data for generation of particles (positive and negative muons) with energy spectrum at sea level. The energy distribution and life time of positive and negative muons are equal. But due to the fact that the primary cosmic rays, protons (95%) and He (5%), are essentially positively charged, the secondary component in the atmosphere has a positive charge excess (between 20 and 30%), too. These data with energies ranging from 1 to 10^3 GeV were extracted using Eq. 1. In this equation, values of parameters A_G , γ , and B_G are $0.14 \text{ (cm}^{-2}\text{sr}^{-1}\text{s}^{-1}\text{GeV}^{\gamma-1}\text{)}$, 2.7, and 0.054, respectively. In addition, the values of critical energies of pions and kaons, $E_{0,\pi}^{\text{cr}}$ and $E_{0,\kappa}^{\text{cr}}$, are 115/1.1 and 850/1.1, respectively. Meanwhile, E_0 and θ are, respectively, the energy and zenith angle of incident cosmic ray muon. In this study, the source has been defined using general source and SDEF card for positive and negative muons separately. The main

control keywords for the SDEF card are position, type, and energy of the particle.

5.3 Tally (output)

The appropriate MCNPX tally that is of event-by-event nature and fits the study requirement is the PTRAC card. The main control keywords for the PTRAC card are: (i) output file, (ii) event type filter, and (iii) history filter. The PTRAC card of MCNPX code generates an output file which includes all the information needed for tracking the particle. This information includes: (a) the location (x, y, z) of the particle and its related cell and material number, and (b) the ID number, weight, event time, directional cosine, and energy of the particle. The PTRAC output files by positive and negative muons were equal.

5.4 Analysis of the output file (PTRAC)

The simulation of penetrating cosmic ray muons through the mountain by 12×10^9 muons was performed. The PTRAC output file is simply an ASCII file including the xyz (Cartesian coordinates), the uvw (directional cosine), the energy, and the event time of the particle (muon) that is being simulated (Zareie 2010 and Tajik *et al.* 2013). Some information that has been extracted from output file (PTRAC) is presented in Table 1. Data in Table 1 have been extracted from Fig. 1A in the Appendix.

Table 1

A portion of typical data which have been extracted from PTRAC output file

Extracted information from PTRAC output file		
Particle number: 374512	Entrance to geometry	Exit of geometry
Location of particle (x, y, z) [cm]	(-7207.9, -19849, 23378)	(45.868, -0.225, -95.656)
Directional cosine	(0.22962, 0.62851, -0.74314)	(0.2323, 0.6253, -0.7449)
Energy of particle [MeV]	23000	20372
Time of particle [10^{-8} s]	0.20534	105.56
Cell number	20	5
Surface number	50	5
Material number	1	2

Using the output file (PTRAC), the distance that muon travels through matter can be calculated. Practically, this distance is calculated by determination of the muon path through matter using the topographic data. Since the PTRAC file gives the directional cosines, the direction of the entrance of muon to detector and the number of muons as a function of (θ, φ) , $N_\mu(\theta, \varphi)$,

in every interval of zenith and azimuthal angles with reference to a line perpendicular to the detector plane can be obtained.

6. CALCULATION OF THE DENSITY OF MATTER INSIDE THE MOUNTAIN

There are two possibilities to obtain a two dimensional profile of the density of matter inside the mountain that are mentioned as follows.

6.1 Using the energy loss of muon inside matter

A unique relationship exists between the range (in gr/cm^3) and energy of the penetrating cosmic ray muons (Eq. 3). Thus, for both the input and the output muon energy, the range can be calculated, and from the difference between these calculated ranges, the density length is determined (Darijani *et al.* 2014).

6.2 The use of relative flux of muons

The relative intensity of cosmic ray muons transmitted through matter with reference to those directly transmitted through the air is called the relative flux of muons that is introduced as Eq. 7.

$$n(\theta, \varphi) = \frac{N_{\mu}(\theta, \varphi)}{N_0(\theta, \varphi)}, \quad (7)$$

Equation 7 is a unique relation between density length (X) and the intensity of penetrating cosmic ray muons through matter. Figure 5 indicates

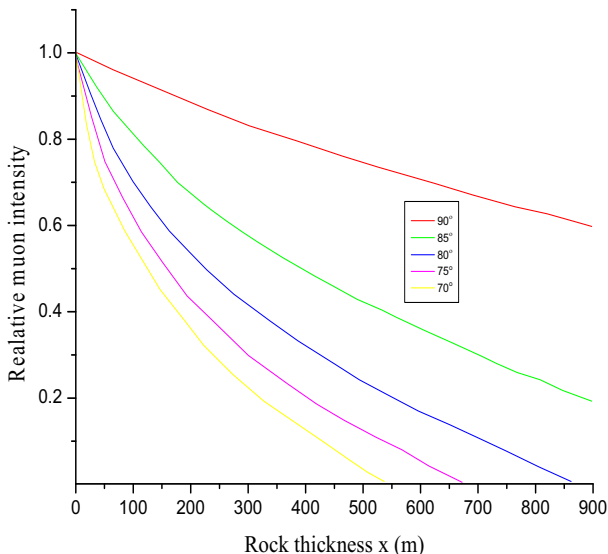


Fig. 5. Relative muon intensity versus rock thickness (Nagamine 2003).

relative intensity of muons as a function of standard rock thickness. So, for a substance with an unknown density length, the measurement of the muon flux penetrating through the substance with zenith and azimuthal angles uniquely determines its density length (in gr/cm^3) (Nagamine 2003). Using the density length, the substance density is determined.

7. RESULTS

The simulation was performed with 12×10^9 particles. Because of the symmetry, only 6×10^9 of particles are toward the geometry (Figs. 4 and 6), and the number of muons that were detected after passing through the geometry are about 45720. Figure 6 shows a two-dimensional picture of production and tracking of muons in the simulation code.

Since cosmic ray muons are constantly irradiating every substance on the Earth at rate of about 70 muon per m^2 per min, the performed simulation is equal naturally to the detection of muons for about 35 hours. The limitation of time is due to the limitation of the number of input particles (2×10^9) in MCNPX code.

Because the size of the detector plane (100×100 cm) is much smaller than the object of interest, the path of each muon can be determined in polar system, as shown in Fig. 4. In this figure the detector plane (xz plane) is

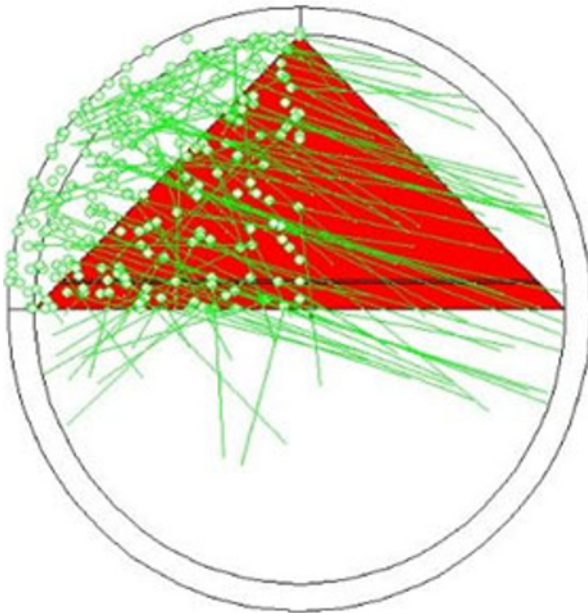


Fig. 6. Two dimensional yz profile of production and tracking of muons. The green dots show muon source points and green lines show muon tracks.

perpendicular to the y axes, so the zenith and the azimuthal angles are obtained with reference to the y axes. The angle between muon path and z axis is θ , and θ_z (zenith angle) is $90 - \theta$ that changes from 0 to 90 degrees (for horizontal and vertical entrance, respectively). The angle between muon path in xy plan and y axis is φ_x (azimuthal angle) that changes from -90 to 90 degrees. These angles were divided into steps, $\Delta\theta$ and $\Delta\varphi$, and in every interval the number of muons was recorded. A portion of the typical file in which the number of muons has been recorded in every interval has been shown in Table 2.

Table 2

A portion of the counted muons in angle intervals

The counted muons in angle intervals $\Delta\theta^\circ$ and $\Delta\Phi^\circ$										
$\Delta\theta^\circ \backslash \Delta\Phi^\circ$...	[0, 3)	[3, 6)	[6, 9)	[9, 12)	[12, 15)	[15, 18)	[18, 21)	[21, 24)	...
...
[-69, -66)	...	2	14	7	12	11	7	8	0	...
[-66, -63)	...	0	7	4	7	21	18	5	7	...
[-63, -60)	...	3	12	2	14	16	13	1	9	...
[-60, -57)	...	7	4	13	12	29	19	2	13	...
[-57, -54)	...	9	3	15	13	18	31	0	15	...
[-54, -51)	...	12	1	2	10	15	49	13	2	...
[-51, -48)	...	14	4	9	6	22	18	8	2	...
[-48, -45)	...	6	2	3	10	15	18	16	9	...
[-45, -42)	...	11	6	5	9	6	23	25	8	...
...

Using the output file (PTRAC), two dimensional profiles of the number of muons and density of matter that muons penetrate through it were plotted and have been shown in Figs. 7-10. Figures 7 and 8 indicate the two dimensional profiles of the number of penetrating particles through the mountain without and with a layer of coal, respectively. The two dimensional profiles of the density of the mountain without and with a layer of coal, that muons penetrate through it, have been shown in Figs. 9 and 10. A comparison of Figs. 9, 10a, and 10b shows the existence of the region with lower density (the layer of coal) in the mountain. Figure 10 shows that around $\varphi_x = 0$ (parallel to the y axis) and for θ_z between 0 and 5 degrees, there is a matter that its density is less than the density of the standard rock. By this approach, the layers with minimum thickness of about 2.5 m are identified.

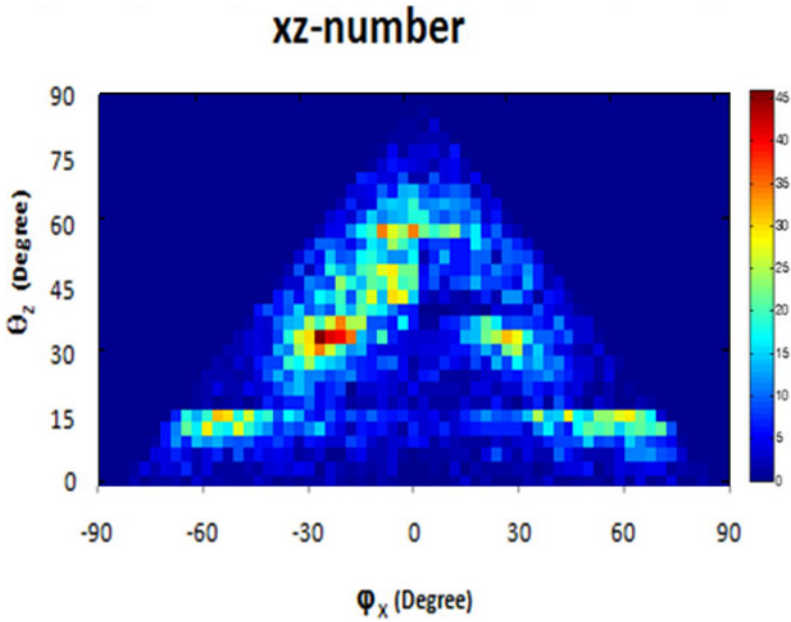


Fig. 7. Two dimensional profile from number of particles in the xz plan without any layer of coal.

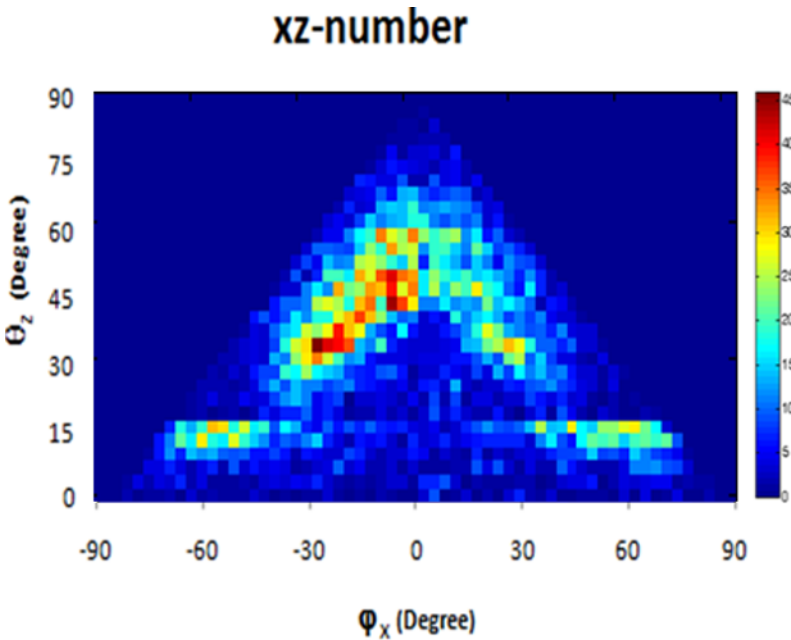


Fig. 8. Two dimensional profile from number of particles in the xz plan with a layer of coal.

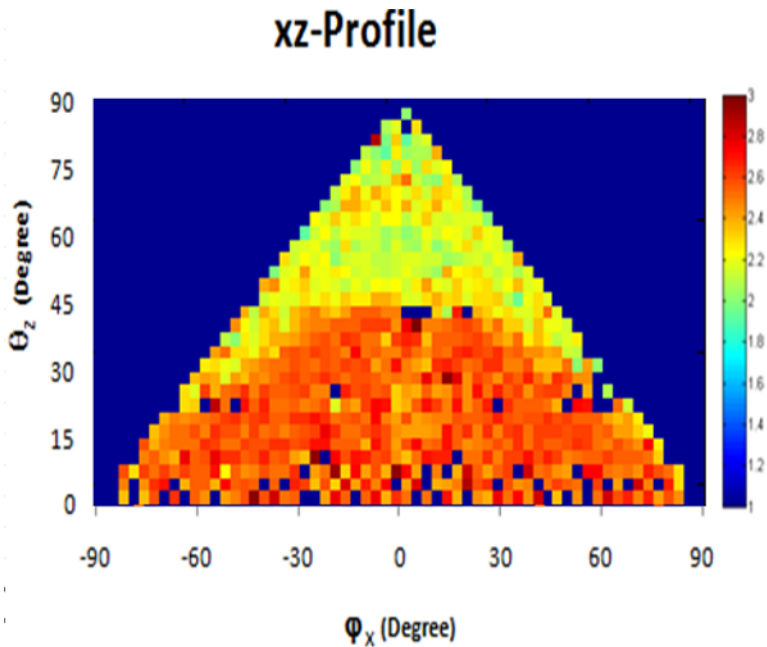


Fig. 9. Two dimensional profile from density of mountain in the xz plan without any layer of coal.

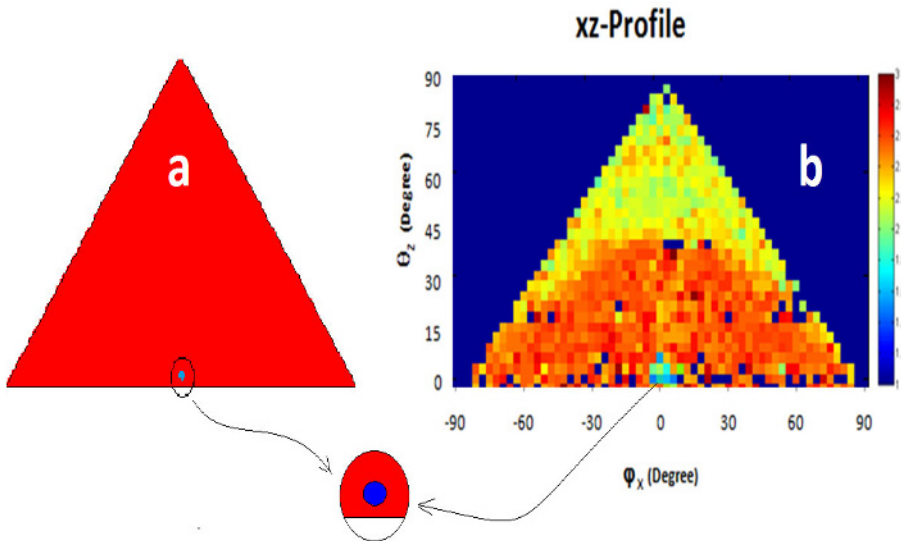


Fig. 10: (a) Two dimensional xz profile of geometry that the code provides (the blue circle shows the xz profile of cylindrical coal layer), and (b) two dimensional profile from density of mountain in xz plan with a layer of coal due to MCNPX simulation results.

8. CONCLUSION

Very high energy muons are most suitable for measuring the density profile of a large scale substance like a mountain. These muons also are used for exploration of mines. The spatial resolution to determine the internal intensity of a mountain depends on not only the spatial resolution of detector but also the distance of the detector from the center of mountain. In this study, the detector was placed inside the mountain (Fig. 4) that increased the spatial resolution and decreased the interest area for imaging. Using many detectors at different altitudes of a mountain, radiography can be done from more parts of mountain, the spatial resolution can be increased, and the exact location of coal layers in a mountain can be determined. Figures 7 to 10 are the results of using one detector.

Appendix

Interpretation of data in PTRAC output file

PTRAC output file is an ASCII file which includes all information required for tracking a particle. This information includes: (a) the location (x, y, z) of the particle and its related cell, and material number (b) the ID number, weight, event time, directional cosine, and energy of the particle. A portion of typical PTRAC output has been shown in Fig. 1A.

```

374512 3000 30 5
3000 2 5 129 20 0 1
-0.72079E+04 -0.19849E+05 0.23378E+05 0.22962E+00 0.62851E+00 -0.74314E+00 0.23000E+05 0.10000E+01
0.20534E+00
3000 3 1 94 10 1 2
-0.13960E+03 -0.50173E+03 0.50133E+03 0.22962E+00 0.62851E+00 -0.74314E+00 0.23000E+05 0.10000E+01
0.10289E+03
3000 4 7 139 101 2 5
-0.8578E+02 -0.35536E+03 0.3279E+03 0.23029E+00 0.62611E+00 -0.7449E+00 0.21862E+05 0.10000E+01
0.10367E+03
3000 5 23 51 10 1 6
-0.83934E+02 -0.35040E+03 0.32205E+03 0.23073E+00 0.62580E+00 -0.7450E+00 0.21836E+05 0.10000E+01
0.10369E+03
3000 6 16 20 20 0 10
-0.2843E+02 -0.20017E+03 0.14257E+03 0.23238E+00 0.62530E+00 -0.74499E+00 0.20372E+05 0.10000E+01
0.10450E+03
3000 7 12 138 30 0 11
-0.10981E+02 -0.15320E+03 0.86600E+02 0.23238E+00 0.62530E+00 -0.74499E+00 0.20372E+05 0.10000E+01
0.10475E+03
3000 8 11 51 20 0 12
0.87891E+01 -0.10000E+03 0.23218E+02 0.23238E+00 0.62530E+00 -0.74499E+00 0.20372E+05 0.10000E+01
0.10503E+03
9000 9 5 51 50 0 13
0.45868E+02 -0.22508E+00 -0.95656E+02 0.23238E+00 0.62530E+00 -0.74499E+00 0.20372E+05 0.10000E+01
0.10556E+03
706380 3000 30 5

```

Fig. 1A. A portion of typical PTRAC output.

The first number of the first line (*i.e.*, 374512) is an ID number for the particle, number 3000 indicates the entrance of the particle to the geometry, number 30 is the number of the cell that has been filtered, and number 5 determines the entrance surface of the source particles in the geometry. In the second line, numbers 5 and 20 show that the particle has been entered from surface 5 into cell 20, number 0 is the material number, number 2 is the line number, and number 129 is the angle between the particle path and the line perpendicular to the surface 5. In the third line, the first six numbers are Cartesian coordinates (x, y, z) and direction cosines (u, v, w) of the entered particle, respectively. Numbers 0.23000E+05, 0.10000E+01, and 0.20534E+00 in the third line are the energy, weight, and event time of the particle, respectively. In line 16, number 9000 characterizes the exit of the particle from the geometry.

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