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Fast multipole method for singular integral equations of second kind

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Abstract

This paper explores the numerical methods for a singular integral equation (SIE), which arise in the study of various problems of mathematical physics and engineering. The idea behind the boundary element method (BEM) is used to discretize the SIE. The fast multipole method (FMM), which is a very efficient and popular algorithm for the rapid solution of boundary value problems, is used to accelerate the BEM solutions of SIE. The effectiveness and accuracy of the proposed method are tested by numerical examples.

Keywords: singular integral equation; boundary element method; fast multipole method; Cauchy principal value; finite-part integral

1 Introduction

Various problems of mathematical physics and engineering can be described by differential equations which can often be reformulated as an equivalent integral equation. For example, the integral equations with singular kernel often arise in practical applications such as potential problems, Dirichlet problem, and radiative equilibrium [1], *etc.*, which in general can be described as

$$au(x) + \frac{b}{\pi} \int_{-1}^1 \frac{u(y)}{y-x} dy + \lambda \int_{-1}^1 k(y,x)u(y) dy = f(x), \quad x \in (-1, 1), \quad (1)$$

where a, b are two real constants, $f(x)$ and $k(y, x)$ are given functions, the first integral in (1) is defined in the sense of Cauchy principal value, and λ is not an eigenvalue. In last few years, many numerical methods have been developed to solve SIE (1), among which collocation methods, Galerkin methods, spectral methods, *etc.* [2–11] have been widely used for solving these kinds of problems for many years. Recently, Xiang and Brunner [12, 13] introduced collocation and discontinuous Galerkin methods for Volterra integral equations with highly oscillatory Bessel kernels, and they concluded that the collocation methods are much more easily implemented and can get higher accuracy than discontinuous Galerkin methods under the same piecewise polynomials space. Volterra integral equations with oscillatory trigonometric kernels are given in [14, 15], which shows that the convergence order of collocation on graded meshes is not necessarily better than that on uniform meshes when the kernels are oscillatory trigonometric. Numerical solutions for the Fredholm integral equations are discussed in [16, 17].

In general, it will give rise to a standard linear system of equations when approximating the solution of SIE by numerical methods. Specially, collocation methods would lead to systems of equations with dense and non-symmetrical coefficient matrices where $O(N^2)$ elements need to be stored, with N being the number of degrees of freedom. Solving the systems of equations directly will need $O(N^3)$ arithmetic operations. Fortunately, Rokhlin and Greengard innovated the fast multipole method (FMM) which has been widely used for solving large scale engineering problems such as potential, elastostatic, Stokes flow, and acoustic wave problems. For one dimension (SIE (1)), the interval $[-1, 1]$ is not a closed curve; however, we can also utilize BEM to solve SIE and accelerate BEM by FMM when the kernel $k(y, x)$ has multipole expansion or $k(y, x) = 0$, for details, see [18–20].

In this paper, we are concerned with the evaluation of SIE

$$u(x) + \lambda \int_{-1}^1 \frac{u(y)}{(y-x)^m} dy = f(x), \quad x \in (-1, 1), \tag{2}$$

where $m \in \mathbb{Z}, m \geq 1, u(x)$ is an unknown function and $f(x)$ is a given function. The integral in (2) is defined in the sense of Hadamard finite-part integral for $m > 1$. For simplicity, we denote SIE (2) as

$$(I + \lambda K)u = f.$$

We approximate the solution of SIE by the collocation methods and utilize the FMM to improve the efficiency of algorithm. The paper is organized as follows. In Section 2 we give a brief description of the FMM, where the multipole expansion theory is introduced and also moment to moment translation (M2M), moment to local translation (M2L), and local to local translation (L2L). In Section 3, we give the convergence analysis of the proposed method. In Section 4 we give preliminary numerical examples to illustrate the effectiveness and accuracy of the proposed method.

2 Fast multipole boundary element method for the solution of (2)

In this section, we recall some basic formulations for the fast multipole boundary element method. In order to solve a SIE numerically for the unknown function, we need to discretize the SIE, firstly. We divide the interval $(-1, 1)$ into several segments $(x_{j-1}, x_j), j = 1, \dots, N$, and use the piecewise constant collocation method [16, 21], then SIE (2) becomes

$$u(\tilde{x}_i) + \lambda \sum_{j=1}^N \int_{x_{j-1}}^{x_j} \frac{u(y)}{(y-\tilde{x}_i)^m} dy = f(\tilde{x}_i), \quad \tilde{x}_i \in (-1, 1). \tag{3}$$

Denote $A = (a_{ij})$ with $a_{ij} = \lambda \int_{x_{j-1}}^{x_j} \frac{1}{(y-\tilde{x}_i)^m} dy + \delta_{ij}$, and $\mathbf{b} = [f(\tilde{x}_1), \dots, f(\tilde{x}_N)]^T, \mathbf{u} = [u(\tilde{x}_1), \dots, u(\tilde{x}_N)]^T$, we obtain a standard linear system of equations

$$A\mathbf{u} = \mathbf{b}.$$

Due to matrix-vector multiplication, solving this system by iterative solvers such as the generalized minimum residue (GMRES) method needs $O(N^2)$ operations, and even worse by direct solvers such as Gauss elimination. Based on the multipole expansion of the kernel, the FMM can be used to accelerate the matrix-vector multiplication.

2.1 Multipole expansion of the kernel

To derive the multipole expansion of kernel $G(y, x) := \frac{1}{(y-x)^m}$, we need the following formulation.

Lemma 1 ([22]) *Let a be any real or complex constant and $|z| < 1$, then we have*

$$(1 + z)^a = 1 + \frac{a}{1!}z + \frac{a(a-1)}{2!}z^2 + \frac{a(a-1)(a-2)}{3!}z^3 + \dots \tag{4}$$

Applying Lemma 1, set $z = \frac{y-y_c}{y_c-x}$, then the kernel $G(y, x)$ can be expressed as the multipole expansion in the following, if $|y - y_c| < |y_c - x|$,

$$\begin{aligned} G(y, x) &= (y_c - x)^{-m} \left(1 + \frac{y - y_c}{y_c - x} \right)^{-m} \\ &= (y_c - x)^{-m} \left(1 + \frac{-m}{1!}z + \frac{-m(-m-1)}{2!}z^2 + \frac{-m(-m-1)(-m-2)}{3!}z^3 + \dots \right) \\ &= \sum_{k=0}^{\infty} O_k(y_c - x) I_k(y - y_c), \end{aligned} \tag{5}$$

where

$$I_k(x) = \frac{x^k}{k!}, \quad k \geq 0, \tag{6}$$

$$O_k(x) = \begin{cases} x^{-m}, & k = 0, \\ x^{-m-k}(-m)(-m-1)\dots(-m-k+1), & k \geq 1. \end{cases} \tag{7}$$

In addition, we have the following two results:

$$I_k(x_1 + x_2) = \sum_{l=0}^k I_k(x_1)I_{k-l}(x_2) = \sum_{l=0}^k I_k(x_2)I_{k-l}(x_1), \tag{8}$$

$$O_k(x_1 + x_2) = \sum_{l=0}^{\infty} O_{k+l}(x_1)I_l(x_2), \quad |x_1| > |x_2|. \tag{9}$$

The integral in (3) is now evaluated as follows:

$$\int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i)u(y) dy = \sum_{k=0}^{\infty} O_k(y_c - \tilde{x}_i)M_k(y_c), \tag{10}$$

where $M_k(y_c) = \int_{x_{j-1}}^{x_j} I_k(y - y_c)u(y) dy$ are called moments about y_c , which are independent of the collocation point \tilde{x}_i .

If the expansion point y_c is moved to a new location $y_{c'}$, we obtain this translation by (6) and

$$M_k(y_{c'}) = \int_{x_{j-1}}^{x_j} I_k(y - y_{c'})u(y) dy,$$

which leads to

$$M_k(y_{c'}) = \sum_{l=0}^k I_{k-l}(y_c - y_{c'}) M_l(y_c). \tag{11}$$

We call this formula moment-to-moment translation (M2M).

On the other hand, if $|y_c - x_l| > |x_l - \tilde{x}_i|$, then (10) can be rewritten as

$$\begin{aligned} \int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i) u(y) dy &= \sum_{k=0}^{\infty} O_k(y_c - \tilde{x}_i) M_k(y_c) \\ &= \sum_{k=0}^{\infty} O_k(y_c - x_l + x_l - \tilde{x}_i) M_k(y_c) \\ &= \sum_{k=0}^{\infty} L_k(x_l) I_k(x_l - \tilde{x}_i) \end{aligned} \tag{12}$$

with

$$L_k(x_l) = \sum_{l=0}^{\infty} O_{k+l}(y_c - x_l) M_l(y_c), \tag{13}$$

where $L_k(x_l)$ denotes the local expansion coefficients about x_l . We call the formula (13) moment-to-local translation (M2L).

If the point for local expansion is moved from x_l to $x_{l'}$, using a local expansion with p terms, we obtain this translation by

$$\begin{aligned} \int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i) u(y) dy &= \sum_{k=0}^p L_k(x_l) I_k(x_l - \tilde{x}_i) \\ &= \sum_{k=0}^p L_k(x_l) I_k(x_l - x_{l'} + x_{l'} - \tilde{x}_i) \\ &= \sum_{k=0}^p L_k(x_l) \sum_{j=0}^k I_k(x_l - x_{l'}) I_{k-j}(x_{l'} - \tilde{x}_i), \end{aligned}$$

which leads to

$$L_k(x_{l'}) = \sum_{j=0}^k L_k(x_l) I_{k-j}(x_{l'} - \tilde{x}_i). \tag{14}$$

We call this formula local-to-local translation (L2L).

2.2 Evaluation of the integrals

The multipole expansion is used to evaluate the integrals that are far away from the collocation point, whereas the direct approach is applied to evaluate the integrals on the remaining segments that are close to the collocation point. In the following, we are con-

cerned with the piecewise constant collocation method for (2), where we can evaluate the moments analytically by

$$\begin{aligned}
 M_k(y_c) &= \int_{x_{j-1}}^{x_j} I_k(y - y_c)u(y) dy \\
 &= u_j \int_{x_{j-1}}^{x_j} \frac{(y - y_c)^k}{k!} dy \\
 &= u_j(I_{k+1}(x_j - y_c) - I_{k+1}(x_{j-1} - y_c)).
 \end{aligned}
 \tag{15}$$

When the integrating interval (x_{j-1}, x_j) is close to the collocation point \tilde{x}_i but not coincidence with the interval (x_{i-1}, x_i) , the integral is not singular, we can evaluate it analytically by

$$\int_{x_{j-1}}^{x_j} \frac{u_j}{(y - \tilde{x}_i)^m} dy = \frac{u_j}{(-m + 1)}(y - \tilde{x}_i)^{-m+1} \Big|_{x_{j-1}}^{x_j}.
 \tag{16}$$

For the integrating interval (x_{i-1}, x_i) , where $\tilde{x}_i \in (x_{i-1}, x_i)$, we can evaluate the integral analytically by

$$\int_{x_{i-1}}^{x_i} \frac{u_i}{(y - c)^m} dy = \begin{cases} \ln(\frac{1-c}{1+c})u_i, & m = 1, \\ u_j \frac{1}{(m-1)!} \frac{d^{(m-1)}}{dc^{(m-1)}} \int_{x_{i-1}}^{x_i} \frac{1}{y-c} dy, & m > 1, \end{cases}
 \tag{17}$$

where $x_{i-1} < c < x_i$.

3 Convergence analysis

In this section, we derive the error bound for (10) when $m = 1$.

Lemma 2 ([23]) *The Cauchy integral operator $K : C^{(0,\alpha)(-1,1)} \rightarrow C^{(0,\alpha)(-1,1)}$ defined by*

$$(Ku)(z) := \frac{1}{\pi i} \int_{-1}^1 \frac{u(y)}{y - z} dy, \quad z \in (-1, 1)
 \tag{18}$$

is bounded, where $C^{(0,\alpha)(-1,1)}$ denotes the Hölder continuous function on $(-1, 1)$.

Lemma 3 ([24]) *Suppose that $f(z)$ is analytic in $|z| < R$, then the Taylor series of $f(z)$ is converging absolutely in $|z| < R$, and we have*

$$\begin{aligned}
 |R_n| &\leq \frac{1}{2\pi} \oint_c \frac{|f(t)||dt|}{|t|(1 - \frac{|z|}{|t|})} \Big| \frac{z}{t} \Big|^n \\
 &\leq M_1 \left(\frac{r}{R} \right)^n,
 \end{aligned}$$

where $r = |z|$, c is a circle centered at origin with radius R and R_n is the remainder of Taylor series.

Theorem 1 *If we apply a multipole expansion with p terms and suppose that $|y_c - \tilde{x}_i|/h \geq 2$, we have the following error bound:*

$$\begin{aligned}
 E_M^p &= \left| \int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i) u(y) dy - \sum_{k=0}^p O_k(y_c - \tilde{x}_i) M_k(y_c) \right| \\
 &\leq C \frac{1}{h} \left(\frac{1}{2} \right)^{p+1},
 \end{aligned} \tag{19}$$

where $h = \max_{y \in (x_{j-1}, x_j)} |y - y_c|$, $C = \int_{x_{j-1}}^{x_j} |u(y)| dy$.

Proof

$$\begin{aligned}
 E_M^p &= \left| \int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i) u(y) dy - \sum_{k=0}^p O_k(y_c - \tilde{x}_i) M_k(y_c) \right| \\
 &= \left| \sum_{k=p+1}^{\infty} O_k(y_c - \tilde{x}_i) M_k(y_c) \right| \\
 &\leq C \sum_{k=p+1}^{\infty} |O_k(y_c - \tilde{x}_i)| \frac{h^k}{k!} \\
 &\leq C \sum_{k=p+1}^{\infty} C_{m+k-1}^k \frac{h^k}{(y_c - \tilde{x}_i)^{m+k}},
 \end{aligned} \tag{20}$$

where $|y - y_c| \leq h$, C_{m+k-1}^k is the binomial coefficients.

For $m = 1$, we have $C_{m+k-1}^k = 1$, then

$$E_M^p \leq C \frac{h^{(p+1)}}{|y_c - \tilde{x}_i|^{(p+2)}} \frac{1}{1 - h/|y_c - \tilde{x}_i|}. \tag{21}$$

Let $\eta = |y_c - \tilde{x}_i|/h$, then the preceding error bound can be written as

$$E_M^p \leq C \frac{1}{h(\eta - 1)} \left(\frac{1}{\eta} \right)^{p+1}, \tag{22}$$

which establishes the desired error bound. □

Theorem 2 *If we apply multipole expansion and local expansion with p terms and suppose that $|y_c - \tilde{x}_i|/h \geq 2$, we have the following error bound:*

$$\begin{aligned}
 E_L^p &= \left| \int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i) u(y) dy - \sum_{k=0}^p L_k(x_l) I_k(x_l - \tilde{x}_i) \right| \\
 &\leq CM \left(\frac{1}{2} \right)^{p+1},
 \end{aligned} \tag{23}$$

where C, M are some constants.

Proof From (12) and (13), we have

$$\begin{aligned}
 E_L^p &= \left| \int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i) u(y) dy - \sum_{k=0}^p L_k(x_l) I_k(x_l - \tilde{x}_i) \right| \\
 &\leq \left| \int_{x_{j-1}}^{x_j} G(y, \tilde{x}_i) u(y) dy - \sum_{k=0}^p O_k(y_c - \tilde{x}_i) M_k(y_c) \right| \\
 &\quad + \left| \sum_{k=0}^p (O_k(y_c - \tilde{x}_i) - \tilde{O}_k(y_c - \tilde{x}_i)) M_k(y_c) \right|,
 \end{aligned} \tag{24}$$

where $\tilde{O}_k(y_c - \tilde{x}_i) = \sum_{n=0}^p O_{k+n}(y_c - x_l) I_n(x_l - \tilde{x}_i)$.

Due to

$$O_k(y_c - \tilde{x}_i) = \frac{k!}{(y_c - \tilde{x}_i)^{k+1}} = \frac{k!}{(y_c - x_l)^{k+1}} \left(1 + \frac{x_l - \tilde{x}_i}{y_c - x_l} \right)^{-k-1}, \tag{25}$$

it follows that

$$|R_{p+1}| = |O_k(y_c - \tilde{x}_i) - \tilde{O}_k(y_c - \tilde{x}_i)| = \left| \sum_{n=p+1}^{\infty} O_{k+n}(y_c - x_l) I_n(x_l - \tilde{x}_i) \right|, \tag{26}$$

which is the remainder of Taylor series expansion of (25).

Define $g(z) = (1 + z)^{-k-1}$, then $g(z)$ is analytic in $|z| < 1$. Since x_l and y_c are well-separated points, we could set $|\frac{x_l - \tilde{x}_i}{y_c - x_l}| < 1/2$. By Lemma 3 we can estimate the remainder of Taylor series R_{p+1} as follows:

$$|R_{p+1}| \leq M_1 \left| \frac{k!}{|(y_c - x_l)^{k+1}} \right| \left(\frac{1}{2} \right)^{p+1}, \tag{27}$$

where M_1 is constant determined by a function $g(z)$.

By Theorem 1, (27) and (24), we have

$$E_L^p \leq C \frac{1}{h} \left(\frac{1}{2} \right)^{p+1} + (p + 1) C M_1 C_1 \left(\frac{1}{2} \right)^{p+1}, \tag{28}$$

where $C_1 = \max\{\frac{h^k}{|(y_c - x_l)^{k+1}}, k = 0, \dots, p\}$.

The desired error bound is established by setting $M = \frac{1}{h} + (p + 1) M_1 C_1$. □

Remark 1 In the FMM, x_l and y_c are well-separated points, we could obtain $|y_c - x_l| > 2h$, which leads to C_1 is bounded.

When we use the FMM to solve SIE (3), the integral operator K is approximated by multipole expansion; if we define K_p as an approximate operator used in the collocation method, then Theorem 1 and Theorem 2 indicate that

$$\lim_{p \rightarrow \infty} \|K_p - K\| = 0. \tag{29}$$

It follows from Lemma 2 and (29) that the integral operator $((I + \lambda K_p)u)(x)$ is bounded.

4 Numerical examples

We illustrate the efficiency and accuracy of the methods described in this paper by numerical examples. Here \hat{u}_N denotes the piecewise constant collocation method, N is the number of collocation points. We choose the uniform mesh, and \tilde{x}_i is the middle point of (x_{i-1}, x_i) . Denote by $I_h^N = \{\tilde{x}_i, i = 1, \dots, N\}$ the set of collocation points.

Example 4.1 We consider

$$u(x) - \frac{1}{2\pi} \int_{-1}^1 \frac{u(y)}{(y-x)^m} dy = f(x), \quad |x| < 1,$$

where

$$f(x) = 1 - \frac{1}{2\pi} \int_{-1}^1 \frac{1}{(y-x)^m} dy,$$

and for $m = 1$ or $m = 2$,

$$u(x) = 1$$

is the exact solution of equation.

Tables 1-4 illustrate the efficiency and accuracy of the methods.

From Tables 1-4, it is easy to see that the proposed method is effective. It might also be noted that $\tilde{x}_i \in I_h^{10}$ but $\tilde{x}_i \notin I_h^{100}$ and $\tilde{x}_i \notin I_h^{1000}$, i.e., the points $\{-0.9, -0.5, -0.1, 0.1, 0.5, 0.9\} \subset I_h^{10}$, but it is not a subset of I_h^{100} or I_h^{1000} .

Table 1 Approximations at $x = -0.9, -0.5, -0.1$ for $u(x) - \frac{1}{2\pi} \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x)$

x	-0.9	-0.5	-0.1
\hat{u}_{10}	1.00000000008664	0.999999999966656	1.000000002837690
\hat{u}_{100}	0.999999988989405	0.999999996983473	0.99999998141215
\hat{u}_{1000}	1.00000006379357	0.99999997488601	0.99999998811889
u	1	1	1

Table 2 Approximations at $x = 0.1, 0.5, 0.9$ for $u(x) - \frac{1}{2\pi} \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x)$

x	0.1	0.5	0.9
\hat{u}_{10}	0.99999996078344	1.00000000199640	1.00000000023095
\hat{u}_{100}	1.00000000689915	1.000000003352249	1.000000033188974
\hat{u}_{1000}	1.00000000150622	1.000000002673433	0.99999995294152
u	1	1	1

Table 3 Approximations at $x = -0.9, -0.5, -0.1$ for $u(x) - \frac{1}{2\pi} \int_{-1}^1 \frac{u(y)}{(y-x)^2} dy = f(x)$

x	-0.9	-0.5	-0.1
\hat{u}_{10}	0.99999993535333	0.99999981428145	0.999999917999549
\hat{u}_{100}	0.99999993845970	0.99999988715871	0.99999990115998
\hat{u}_{1000}	0.99999995690078	0.99999991935489	0.99999993153335
u	1	1	1

Table 4 Approximations at $x = 0.1, 0.5, 0.9$ for $u(x) - \frac{1}{2\pi} \int_{-1}^1 \frac{u(y)}{(y-x)^2} dy = f(x)$

x	0.1	0.5	0.9
\hat{u}_{10}	0.999999917999549	0.999999981428145	0.999999993535332
\hat{u}_{100}	0.999999990036579	0.999999988582430	0.999999992194477
\hat{u}_{1000}	0.999999993375294	0.999999992061794	0.999999995726361
u	1	1	1

Table 5 Approximations at $x = -0.9, -0.5, -0.1$ for $u(x) - \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x)$

x	-0.9	-0.5	-0.1
\hat{u}_{10}	-1.194519628297830	-0.601189820467772	-0.156767445549832
\hat{u}_{100}	-0.902665658196492	-0.496591018161909	-0.094652542268498
\hat{u}_{1000}	-0.900369341742219	-0.499656299326138	-0.099458627117756
u	-0.9	-0.5	-0.1

Table 6 Approximations at $x = 0.1, 0.5, 0.9$ for $u(x) - \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x)$

x	0.1	0.5	0.9
\hat{u}_{10}	0.057786709933708	0.465859007671879	0.872182894009698
\hat{u}_{100}	0.106011730395414	0.507142753891225	0.908324232818361
\hat{u}_{1000}	0.100609029282889	0.500725569118200	0.900861403735125
u	0.1	0.5	0.9

Example 4.2 Let us consider

$$u(x) - \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x), \quad |x| < 1,$$

where

$$f(x) = x - (2 - x \ln(x + 1) + x \ln(1 - x)),$$

then

$$u(x) = x,$$

is the exact solution of equation.

Tables 5-6 also show that the proposed method is effective and the convergence order is $O(1/N)$, which coincides with the classical theory of collocation methods.

Example 4.3 For the more general case, we consider

$$u(x)(1 + x^2) - (1 + x^2) \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x), \quad |x| < 1,$$

where

$$f(x) = x^3 - \left(\frac{4 - \pi}{2} + 2x^2 + x^3 \log[-1 + 2/(1 + x)] \right)$$

Table 7 Approximations at $x = -0.9, -0.5, -0.1$ for $u(x) - \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x)$

x	-0.9	-0.5	-0.1
\hat{u}_{10}	-0.608896542845208	-0.143953205178288	-0.019345651255813
\hat{u}_{100}	-0.400907820963934	-0.096826636160551	-0.002238964950304
\hat{u}_{1000}	-0.402641823825626	-0.099675717661578	-0.001102497056280
u	-0.402762430939227	-0.100000000000000	-0.000990099009901

Table 8 Approximations at $x = 0.1, 0.5, 0.9$ for $u(x) - \int_{-1}^1 \frac{u(y)}{(y-x)} dy = f(x)$

x	0.1	0.5	0.9
\hat{u}_{10}	-0.025111748216482	0.052110055810136	0.363240467300408
\hat{u}_{100}	-0.001202605523378	0.101502077326766	0.409787561038363
\hat{u}_{1000}	0.000783503911385	0.100161733460852	0.403492522553747
u	0.000990099009901	0.1	0.402762430939227

and the exact solution is

$$u(x) = \frac{x^3}{1+x^2}.$$

Tables 7-8 also illustrate the efficiency and accuracy of the methods.

5 Conclusion

In this paper, we explore collocation methods for a singular Fredholm integral equation of the second kind and utilize the FMM to improve the efficiency of algorithm. Based on the multipole expansion of kernel, we demonstrate that the approximate operator used in the collocation equation converges to the initial operator. Numerical examples demonstrate the performance of the proposed algorithm.

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

All authors derived the proof. All authors conceived of the study and participated in its design and coordination. All authors read and approved the final manuscript.

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