

NANO EXPRESS

Open Access



# A Theoretical Simulation of the Radiation Responses of Si, Ge, and Si/Ge Superlattice to Low-Energy Irradiation

Ming Jiang<sup>1</sup>, Haiyan Xiao<sup>1\*</sup> , Shuming Peng<sup>2</sup>, Guixia Yang<sup>2</sup>, Zijiang Liu<sup>3</sup>, Liang Qiao<sup>1</sup> and Xiaotao Zu<sup>1</sup>

## Abstract

In this study, the low-energy radiation responses of Si, Ge, and Si/Ge superlattice are investigated by an ab initio molecular dynamics method and the origins of their different radiation behaviors are explored. It is found that the radiation resistance of the Ge atoms that are around the interface of Si/Ge superlattice is comparable to bulk Ge, whereas the Si atoms around the interface are more difficult to be displaced than the bulk Si, showing enhanced radiation tolerance as compared with the bulk Si. The mechanisms for defect generation in the bulk and superlattice structures show somewhat different character, and the associated defects in the superlattice are more complex. Defect formation and migration calculations show that in the superlattice structure, the point defects are more difficult to form and the vacancies are less mobile. The enhanced radiation tolerance of the Si/Ge superlattice will benefit for its applications as electronic and optoelectronic devices under radiation environment.

**Keywords:** Superlattice, Si and Ge, Radiation, Defect formation and migration

## Background

During the past decades, the Si/Ge superlattice (SL) has attracted much attention in semiconductor research due to its potential contribution to the development of new electronic and optoelectronic devices [1–6]. For example, the study of photoconductivity of Si/Ge SL is of remarkable importance for photodiodes as emitter and receiver for fast optical communication [5]. In its applications like the space electronic component, the micro-electronic component, the solar cell and the space-based electronics [1, 4, 6], the optical and electronic properties of Si/Ge SL may be altered due to the bombardment of high-energy ions from space environment, resulting in performance degradation of the electronic devices. Therefore, it is necessary to investigate the radiation responses of this semiconductor material under extreme working conditions.

Recently, a lot of researchers have studied the radiation damage effects of Si/Ge superlattice [7–16]. Sobolev et al. investigated the influences of electron irradiation on the

photoluminescence (PL) of Si/Ge SL containing mono-layer of pure Ge, and enhanced radiation resistance of the SL structure was found as compared with bulk silicon [12]. Fonseca et al. irradiated the Si/Ge SL with embedded Ge quantum dots (QDs) employing the 2.0 MeV proton irradiation and found an extraordinary high radiation resistance of the QD-in-SL structure [13]. Similar results were obtained by Leitão et al., who reported that the Ge quantum wells (QWs) deposited on a diode structure containing a Si/Ge multilayer structure were more resistant to the proton irradiation as compared with the single Ge QWs [14]. As the promising thermoelectric materials, the thermoelectric characteristic of Si/Ge system may be also affected under the radiation environment [11, 15]. Zheng et al. irradiated the multiple periodic layers of  $\text{Si}_{1-x}\text{Ge}_x/\text{Si}$  employing 5 MeV Si ions, and they found that the thermo-electric figure of merit increases with increasing Si ions fluencies [11]. The defects and structural disorder reduce the cross plane thermal conductivity by absorbing and dissipating phonon along the lattice, and the electronic density of states in the miniband of the QD structure increases the electrical conductivity and the Seebeck coefficient, which all contribute to the increase of figure of merit [11].

\* Correspondence: [hyxiao@uestc.edu.cn](mailto:hyxiao@uestc.edu.cn)

<sup>1</sup>School of Physics, University of Electronic Science and Technology of China, Chengdu 610054, China

Full list of author information is available at the end of the article

Theoretically, Sayed and Windl both investigated the atomic displacements of bulk Si employing the classical molecular dynamics (MD) method [17, 18]. They found that the threshold displacement energies ( $E_d$ s) depend on the knock-on direction and the damaged states are mainly Frenkel pair (FP) defects [17, 18]. Caturla et al. studied the effects of ion mass and energy on the radiation damage of bulk Si employing the MD method [19]. They reported that the production of amorphization as well as isolated point defects and small clusters have a strong dependence on ion mass and a weak relationship to ion energy [19]. Holmström et al. calculated the  $E_d$ s for germanium using the MD method and found that the stable defects are FP defects [20]. Shaw et al. applied an ab initio method to study the effects of antimony and germanium defects on the electronic structure of Si/Ge heterostructures and found that these defects interact with the Si/Ge interfaces, resulting in interface-related localized resonances and large local perturbations to the electronic structure [21]. Despite of these mentioned investigations, no theoretical simulations of dynamic process of radiation damage of Si/Ge SL have been reported in the literature thus far. There still lacks an atomic-level understanding of the micro-structural evolution and the underlying mechanism for defect generation in the semiconductor superlattices.

The ab initio molecular dynamics (AIMD) method has been demonstrated to be an important tool for shedding light on the radiation damage processes and has indeed been successful in simulating the recoil events of a series of semiconductor and ceramic materials [22–27]. As compared with the classical MD method, the interatomic potentials are obtained from electronic structure calculations rather than empirical fitting of experimental results. Consequently, a lot of physical parameters like  $E_d$ s can be determined with ab initio accuracy. In this study, the AIMD method is employed to compare the response behaviors of bulk Si, Ge, and Si/Ge SL under low-energy irradiation. The threshold displacement energies have been determined, and the defect distribution and the pathway for defect generation have been provided. The possible origin for the discrepancy in radiation tolerance between bulk Si (Ge) and Si/Ge SL is also explored. The presented results provide a fundamental insight into the microscopic mechanism of displacement events in bulk Si, Ge, and Si/Ge SL and advance the understanding of the radiation responses of these materials under radiation environment.

## Methods

The low-energy displacement events of bulk Si, Ge, and Si/Ge SL are simulated by the Spanish Initiative for Electronic Simulations with Thousands of Atoms (SIESTA) code. The norm-conserving Troullier-Matrisins pseudopotentials [28]

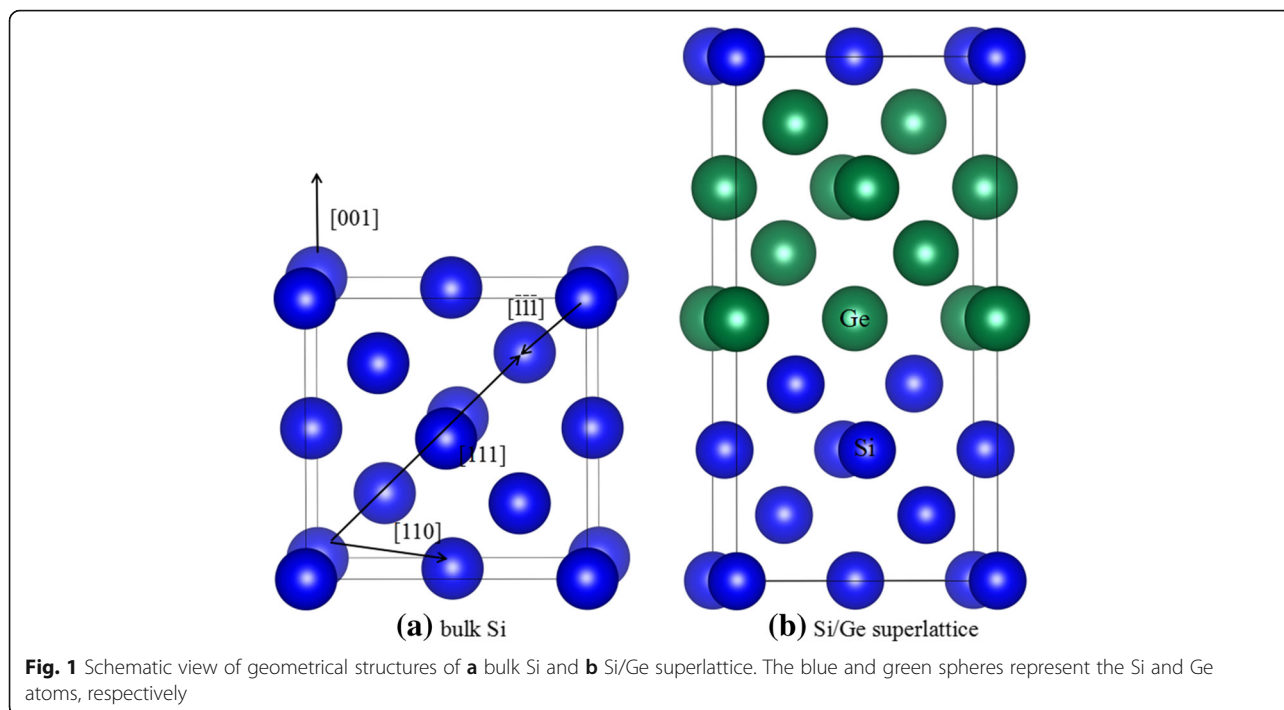
are employed to determine the interaction between ions and electrons, and the exchange-correlation potential is described by the local-density approximation (LDA) in Ceperly-Alder parameterization [29]. The valence wave functions are expanded by a basis set of localized atomic orbitals, and single- $\zeta$  basis sets plus polarization orbital (SZP) are employed, with a K-point sampling of  $1 \times 1 \times 1$  in the Brillouin zone and a cut-off energy of 60 Ry. In the present study, a Si<sub>2</sub>/Ge<sub>2</sub> SL, which consists of two layers of Si alternating with two layers of Ge and totally 288 atoms, is considered. Figure 1 illustrates the geometrical configuration for bulk Si and Si/Ge SL. A specific atom is selected as the primary knock-on atom (PKA), and it is given a kinetic energy to initiate a recoil event. If the PKA returns to its original position at the end of the displacement event, the simulation is restarted at higher recoil energy with an energy increment of 5 eV. Once the PKA is permanently displaced from its lattice site, additional runs are performed to improve the precision to 0.5 eV. For each atom type, four and five principal incidence directions are taken into account for bulk Si (Ge) and Si/Ge SL, respectively. The simulations are conducted with an NVE ensemble and the maximum duration of each run is 1.2 ps to avoid the instability of the system.

## Results and Discussion

### The Displacement Events in Bulk Silicon and Germanium

The lattice constant of bulk Si is determined to be 5.50 Å, which agrees well with the theoretical result of 5.48 Å [30] and the experimental result of 5.43 Å [31]. As compared with bulk Si, the lattice constant of bulk Ge is larger, i.e., 5.71 Å, which is consistent with the calculated result of 5.65 Å [30] and the experimental value of 5.77 Å [31]. Our calculated threshold displacement energies for bulk Si and Ge are summarized in Table 1, along with the associated defects after the displacement events. The configurations for the damage end states of Si and Ge recoils are plotted in Figs. 2 and 3, respectively.

For bulk Si, the  $E_d$  values are slightly smaller than the experimental results of 21 eV for [001] [32], ~47.6 eV for [110] [33], and ~12.9 eV for [111] [34] directions, and both the experiment and our calculations reveal that the damaged end states are Frenkel pair (FP) defect. It is also noted that  $E_d$  values in the present study are generally comparable with the MD results reported by Windl et al. [18], except the case of [110], for which our calculated value of 47 eV is much larger than the MD result of 24 eV. Previous AIMD simulation of ion-solid interactions in SiC revealed that the displacement event is actually a charge-transfer process and the charge transfer to and from recoiling atoms can alter the energy barriers and dynamics for stable defect formation [35]. The lower values of  $E_d$  found by AIMD compared to those determined by classical MD may be due to the fact that



charge transfer that occurs during the recoil events is taken into account by the AIMD method, while in the classical MD simulation, the charge of atoms is fixed. In the study of Windl et al., the kinetic energy is transferred to the PKA to generate one silicon vacancy ( $V_{Si}$ ) and one silicon interstitial ( $Si_{int}$ ) defects [18]. By contrast, in our study, the associated defects for Si[110] displacement event contain two  $V_{Si}$  and two  $Si_{int}$  defects, resulting in higher energies for the formation of the damaged states. The  $E_d$  values for Si[111] and Si $[\bar{1}\bar{1}\bar{1}]$  are very close to each other, i.e., 9.5 and 10 eV, respectively. In both cases, the created defects are  $V_{Si}$  and  $Si_{int}$  (see Fig. 2c, d), whereas the mechanisms of defect generation show different character. In the case of Si[111], the Si

PKA moves along the  $[11\bar{1}]$  direction due to the repulsive interactions and collides with its neighboring Si atom. The Si PKA then scatters away to occupy an interstitial site ( $Si_{int}$ ), and the replaced Si moves back to the lattice site of PKA. The associated defects are one  $V_{Si}$  and one  $Si_{int}$  defects. As for Si $[\bar{1}\bar{1}\bar{1}]$ , the displacement event is relatively simpler, i.e., the Si PKA moves 4.69 Å away from its lattice site to form a  $Si_{int}$  defect. In the cases of Si[001] and Si[110], the  $E_d$ s are determined to be 20 and 47 eV, respectively, indicating that the Si atoms are more difficult to be displaced along the [110] direction. The damage end states for Si[001] and Si[110] are somewhat different. In the case of Si[001], the PKA receives kinetic energy and moves along the [001] direction to collide with its neighboring atoms. The replaced Si atom keeps moving and occupies an interstitial site, as shown in Fig. 2a. As for Si[110], the PKA scatters toward the  $[11\bar{1}]$  direction due to the repulsive interactions between the PKA and its neighboring atoms and hits one neighboring Si atom ( $Si_1$ ). Then, the Si PKA rebounds toward the [111] direction to replace another Si atom ( $Si_2$ ), and the  $Si_2$  atom occupies an interstitial site in the end. The  $Si_1$  atom receives sufficient energy to move along the [110] direction and replaces its neighboring Si atom ( $Si_3$ ), which forms an interstitial defect. In the end, the associated defects are two  $V_{Si}$  and two  $Si_{int}$  defects, as shown in Fig. 2b.

**Table 1** The calculated threshold displacement energies and associated defects after the recoil events in bulk Si and Ge.  $V_X$ : X vacancy ( $X = Si$  or Ge);  $X_{int}$ : X interstitial ( $X = Si$  or Ge)

Direction	Bulk Si		Bulk Ge	
	$E_d$ (eV)	Defect type	$E_d$ (eV)	Defect type
[001]	20, 17.4 <sup>a</sup> , 21 <sup>c</sup>	$V_{Si} + Si_{int}$	18, 18.5 <sup>b</sup> , ~ 18 <sup>f</sup>	$V_{Ge} + Ge_{int}$
[110]	47, 24 <sup>a</sup> , ~ 47.6 <sup>d</sup>	$2V_{Si} + 2Si_{int}$	28.5	$V_{Ge} + Ge_{int}$
[111]	9.5, 11.3 <sup>a</sup> , ~ 12.9 <sup>e</sup>	$V_{Si} + Si_{int}$	9.5, 12.5 <sup>b</sup> , ~ 15 <sup>c</sup>	$V_{Ge} + Ge_{int}$
$[\bar{1}\bar{1}\bar{1}]$	10	$V_{Si} + Si_{int}$	9.5, 10.5 <sup>b</sup>	$V_{Ge} + Ge_{int}$

<sup>a</sup>Ref. [17]

<sup>b</sup>Ref. [20]

<sup>c</sup>Ref. [32]

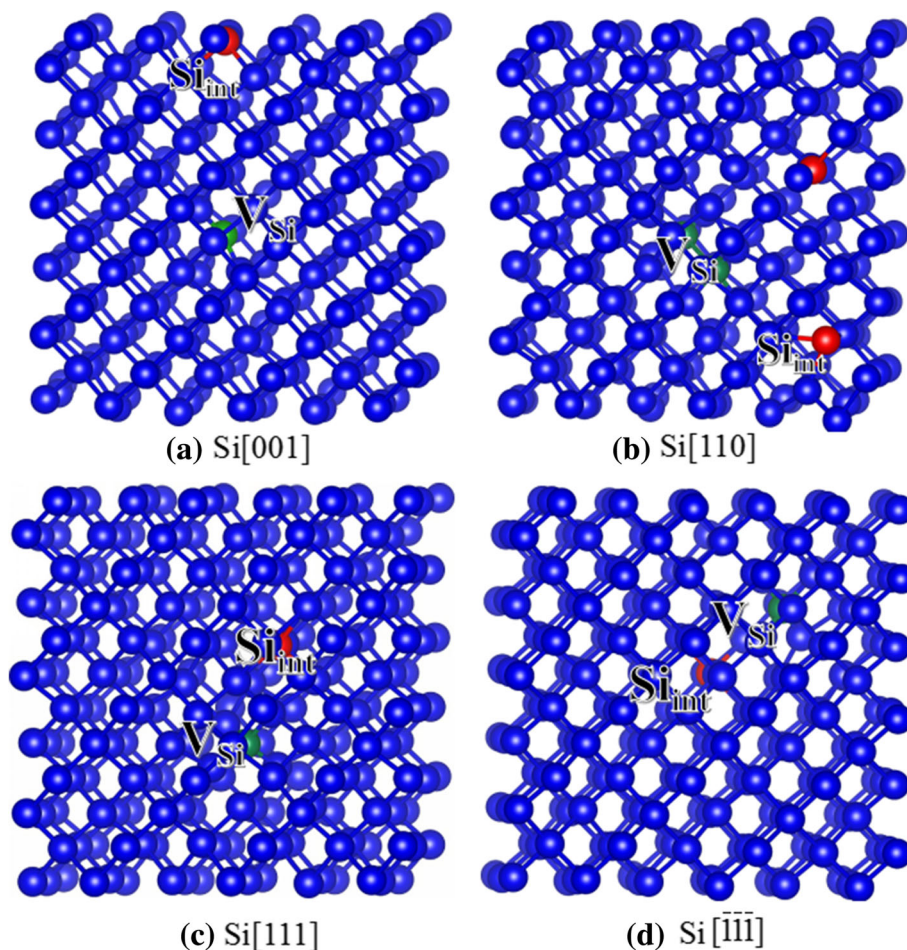
<sup>d</sup>Ref. [33]

<sup>e</sup>Ref. [34]

<sup>f</sup>Ref. [36]

For bulk Ge, the values of  $E_d$  are in good agreement with the experimental value of ~ 18 eV [36] and the theoretical value of 18.5 eV [20] for [001] direction. It is





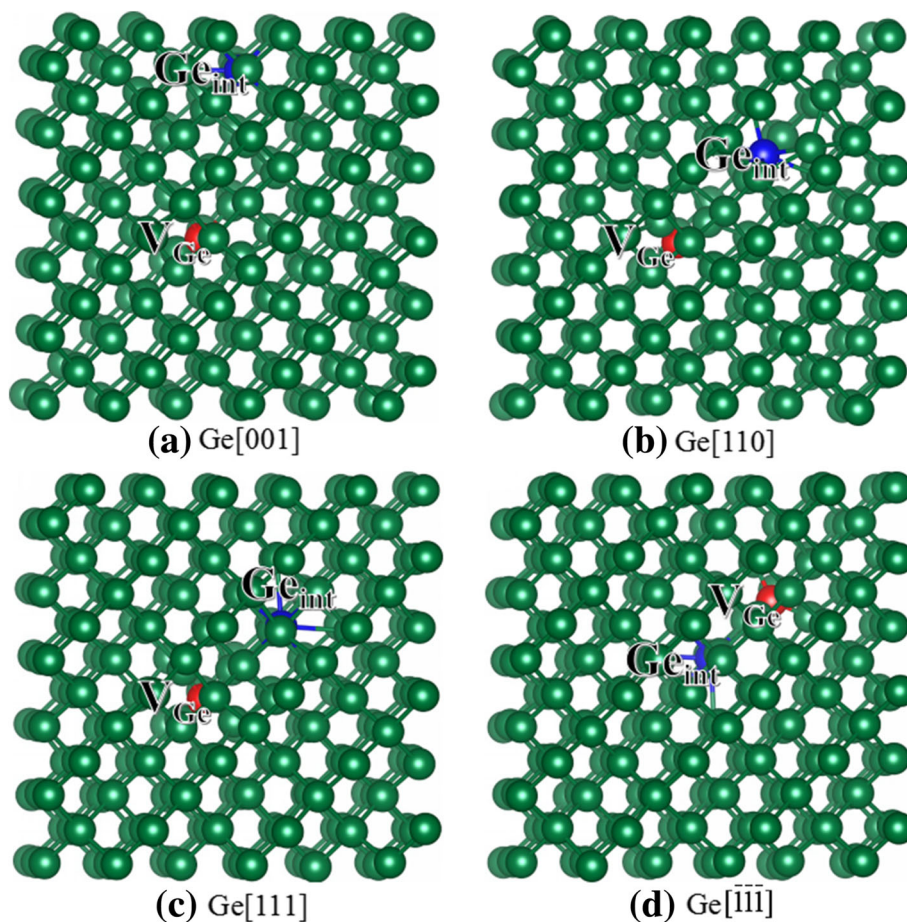
**Fig. 2 a–d** Schematic view of geometrical structures of damage Si after recoil events. The green and red spheres represent the vacancy and interstitial defects, respectively.  $V_{Si}$ : silicon vacancy;  $Si_{int}$ : silicon interstitial

noted that the present value of 9.5 eV is comparable to the Holmström’s result of 12.5 eV [20] for [111] direction, which are smaller than the experimental value of ~ 15 eV [36]. For Ge[111] and Ge  $\bar{1}\bar{1}\bar{1}$ , the determined  $E_d$  values are as small as 9.5 eV, indicating that the Ge atoms are easily to be displaced along these two directions. In both cases, the associated defects are germanium vacancy and germanium interstitial (see Fig. 3c, d). For Ge  $\bar{1}\bar{1}\bar{1}$ , the Ge PKA does not follow a straight path, but gets strongly deflected by one of its nearest neighbors to occupy an interstitial site ( $Ge_{int}$ ). By contrast, in the case of Ge[111], the Ge PKA moves 4.92 Å along the [111] direction to form an interstitial defect ( $Ge_{int}$ ). As compared with the  $E_d$  of Ge[001], the value of Ge[110] is 10 eV larger, indicating that the Ge atom is more difficult to be displaced along the [110] direction. Although the associated defects for Ge[001] and Ge[110] are similar, the mechanisms for defect generation are somewhat different. The Ge PKA receives kinetic energy and moves along the

[001] direction to collide with its neighboring atoms. The replaced Ge atom keeps moving and occupies an interstitial site, as shown in Fig. 3a. As for Ge[110], the Ge recoil collides with its first neighboring Ge atom ( $Ge_1$ ) along the [110] direction and rebounds along the [111] direction, resulting in the formation of  $Ge_{int}$ . The  $Ge_1$  atom leaves its lattice site and replaces its neighboring Ge atom ( $Ge_2$ ). Subsequently, the  $Ge_2$  atom moves back to the lattice site of  $Ge_1$  and eventually only one  $V_{Ge}$  and one  $Ge_{int}$  defects are formed, as shown in Fig. 3b. These results suggest that in bulk Si and Ge, the  $E_d$ s are strongly dependent on the crystallographic direction, and the atoms are more difficult to be displaced along the [110] direction. The radiation damage end states in bulk Si and Ge are mainly FP defects, i.e., vacancy and interstitial defects.

#### The Displacement Events in Si/Ge Superlattice

In this study, the displacement events of  $Si_2/Ge_2$  SL, which contains two layers of Si alternating with two



**Fig. 3 a–d** Schematic view of geometrical structures of damage Ge after recoil events. The red and blue spheres represent the vacancy and interstitial defects, respectively.  $V_{Ge}$ : germanium vacancy;  $Ge_{int}$ : germanium interstitial

layers of Ge (see Fig. 1b), are considered. The Si and Ge atoms that are adjacent to the Si/Ge interface are selected as the PKA. The  $E_d$ s for Si and Ge recoils and the associated defects are listed in Table 2. The defect configurations for Si and Ge recoils are illustrated in Figs. 4 and 5, respectively. It is noted that in the case of Si[111], no defects are created even at energies up to 100 eV. Due to the computational restrictions, we did not

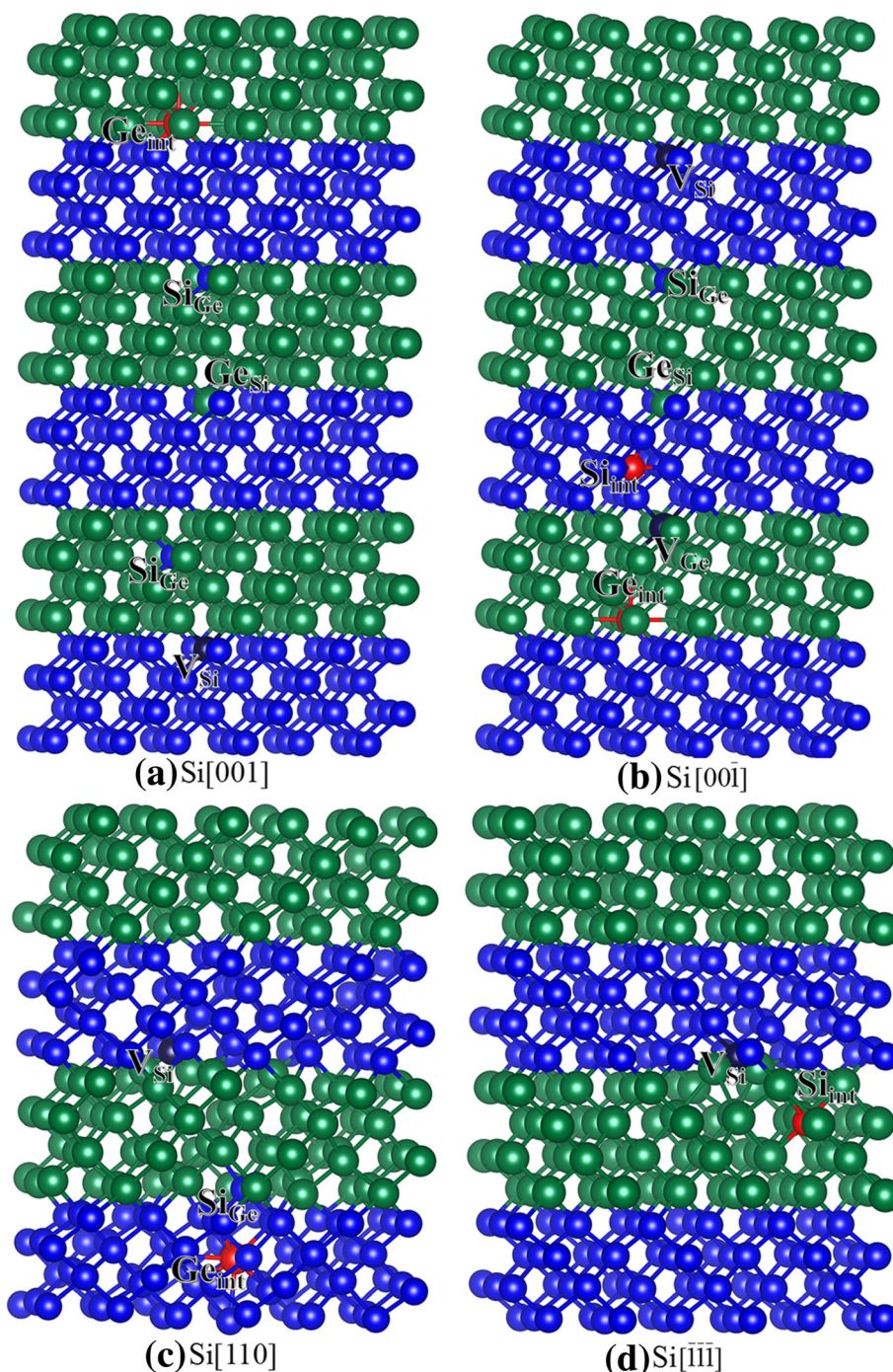
perform further simulations of recoil events at energies higher than 100 eV, and the exact  $E_d$  value for Si[111] is not determined.

In the Si/Ge SL structure, the Si PKA is found to be easily displaced along the  $[111]$  direction, as indicated by the small  $E_d$  value of 10 eV. The pathway for defect generation is very simple, i.e., the Si PKA moves 4.61 Å away from its lattice site and forms a  $Si_{int}$  defect. For

**Table 2** The calculated threshold displacement energies and associated defects after the recoil events in Si/Ge superlattice.  $V_X$ : X vacancy ( $X = Si$  or  $Ge$ );  $X_{int}$ : X interstitial ( $X = Si$  or  $Ge$ );  $X_Y$ : X occupying the Y lattice site ( $X$  and  $Y = Si$  or  $Ge$ )

Direction	Si recoils		Ge recoils	
	$E_d$ (eV)	Defect type	$E_d$ (eV)	Defect type
[001]	46.5	$V_{Si} + Si_{Ge} + Ge_{Si} + Si_{Ge} + Ge_{int}$	16	$Si_{Ge} + Ge_{Si}$
[001̄]	42.5	$V_{Si} + Si_{Ge} + Ge_{Si} + Si_{int} + Ge_{int} + V_{Ge}$	17.5	$V_{Ge} + Ge_{Si} + Si_{int}$
[110]	38.5	$V_{Si} + Si_{Ge} + Ge_{int}$	20	$V_{Ge} + Ge_{Si} + Si_{int}$
[111]	> 100	–	10	$V_{Ge} + Ge_{int}$
[111̄]	10	$V_{Si} + Si_{int}$	13.5	$V_{Ge} + Ge_{int}$

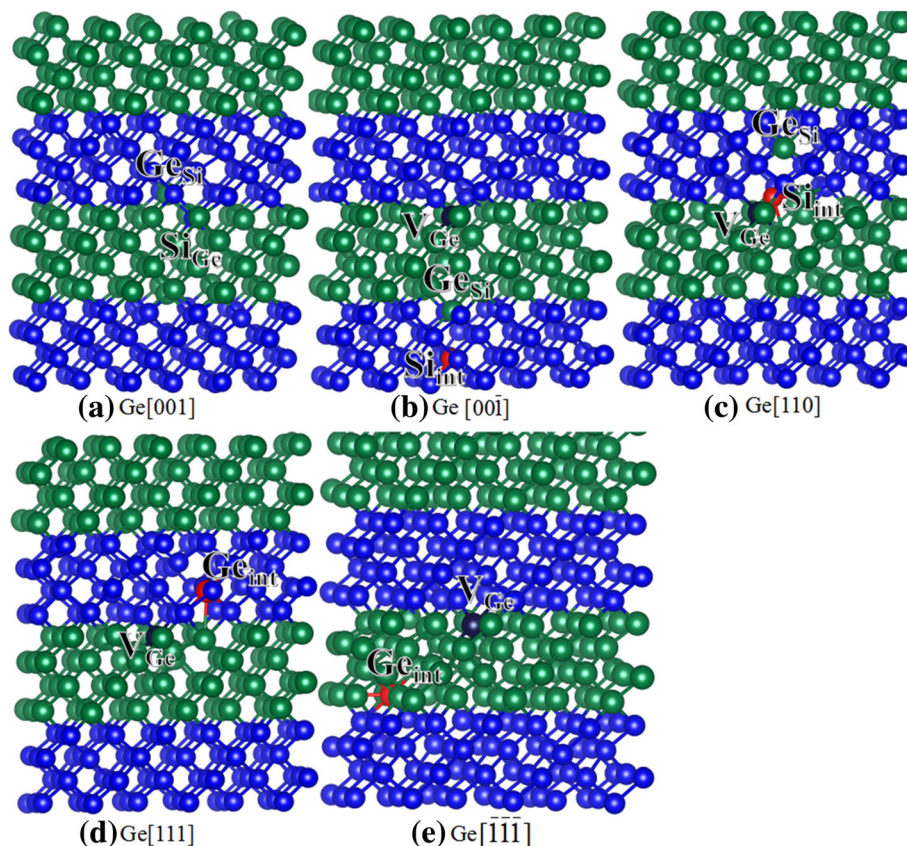




**Fig. 4 a-d** Schematic view of geometrical structures of damage Si/Ge superlattice after Si recoil events. The blue and green spheres represent the Si and Ge atoms, respectively.  $V_X$ : X vacancy (X=Si or Ge);  $X_{int}$ : X interstitial (X=Si or Ge);  $X_Y$ : X occupying the Y lattice site (X and Y=Si or Ge). The purple and red spheres represent the vacancy and interstitial defects, respectively

Si[001] and Si[00 $\bar{1}$ ], the  $E_{ds}$  are determined to be 46.5 and 42.5 eV, respectively, and the damaged defects are different as expected. In the case of Si[001], the Si PKA moves along the [001] direction to replace its neighboring Ge atom ( $Si_{Ge}$ ), and the replaced Ge atom collides

with its adjacent Si atom and occupies its lattice site, forming a  $Ge_{Si}$  antisite defect. The replaced Si atom receives sufficient energy and further replaces another Ge atom ( $Si_{Ge}$ ), which finally occupies an interstitial site. Eventually, the associated defects are one  $V_{Si}$ , one  $Ge_{int}$ ,



**Fig. 5 a–e** Schematic view of geometrical structures of damage Si/Ge superlattice after Ge recoil events. The blue and green spheres represent the Si and Ge atoms, respectively.  $V_X$ : X vacancy (X = Si or Ge);  $X_{int}$ : X interstitial (X = Si or Ge);  $X_Y$ : X occupying the Y lattice site (X and Y = Si or Ge). The purple and red spheres represent the vacancy and interstitial defects, respectively

and three antisite defects. As for Si[00 $\bar{1}$ ], two neighboring Ge atoms and one neighboring Si atom are also involved in the displacement event, and the damaged states contain two vacancies, two interstitials, and two antisite defects, as shown in Fig. 4b. In the case of Si [110], the Si atom moves to hit its neighboring Si atom and scatters toward the [11 $\bar{1}$ ] direction. Then, the Si PKA replaces one neighboring Ge atom, which occupies an interstitial site in the end. After the displacement events, the associated defects contain one  $V_{Si}$ , one  $Si_{Ge}$ , and one  $Ge_{int}$  defects. As compared with the bulk Si, the Si atoms in Si/Ge SL are generally more difficult to be displaced except the case of [110] and the mechanisms of defect generation are more complex, indicating that the bulk Si and Si/Ge SL show different radiation responses to irradiation. Our results are consistent with the experiments carried out by Fonseca et al. and Leitão et al. [13, 14], who also found that the radiation resistance of the SL structure was enhanced as compared with the bulk silicon.

For Ge recoils in Si/Ge SL, the Ge atoms are easily to be displaced along the [111] and [ $\bar{1}\bar{1}\bar{1}$ ] directions, which

are similar to the Ge recoil events in bulk Ge. Although the radiation damage end states for Ge[111] and Ge[ $\bar{1}\bar{1}\bar{1}$ ] are very similar, i.e., Ge FP defects, the mechanisms of defect generation are different. In the case of Ge[111], the Ge PKA moves 4.77 Å away from its lattice site and forms a  $Ge_{int}$  defect. For the Ge [ $\bar{1}\bar{1}\bar{1}$ ], the Ge atom moves along the [ $\bar{1}\bar{1}\bar{1}$ ] direction to replace its neighboring Ge atom. The collided Ge atom moves along this direction and occupies an interstitial site in the end. It is noted that the  $E_d$  values of 16 eV for Ge[001] and 17.5 eV for Ge[00 $\bar{1}$ ] are comparable with the value of 18 eV for Ge[001] in bulk Ge, whereas the associated defects show different character. In the case of Ge[001], the Ge PKA receives sufficient energy but scatters along the [111] direction to replace its neighboring Si atom, forming a  $Ge_{Si}$  antisite defect. Then, the replaced Si atom occupies the Ge PKA lattice site and forms an antisite defect ( $Si_{Ge}$ ). In the case of Ge [00 $\bar{1}$ ], the Ge PKA moves 5.63 Å away to replace its neighboring Si atom. The Si atom moves along this direction and forms a  $Si_{int}$  defect. As compared with the Ge[110] in bulk Ge, the  $E_d$  for Ge[110] in Si/Ge SL is



8.5 eV smaller, and the associated defects are more complex, as indicated by one  $V_{Ge}$ , one  $Ge_{Si}$ , and one  $Si_{int}$  defects. Comparing the Ge recoil events in bulk Ge and SL, we find that the Ge atoms in Si/Ge SL are more resistant along the [110] direction. For other displacement events, the  $E_{ds}$  are generally comparable with those for bulk states. However, the radiation damage end states in bulk Ge and Si/Ge SL are different, and some antisite defects are created in Si/Ge SL structure. These results suggest that the Ge recoils in Si/Ge SL structure show different radiation responses to irradiation. Comparing the Si and Ge recoils in SL structure, we find that the displacement events of Si atoms are much more affected than Ge, i.e., the  $E_{ds}$  for Si atoms in SL structure are generally increased, which may lead to enhanced radiation resistance of Si/Ge SL. Sobolev et al. have found that the Si/Ge SLs show extraordinarily high radiation hardness as compared with bulk Si [12], which is consistent with our results.

**The Defect Formation Energy and Migration Barrier in Bulk Si, Ge, and Si/Ge Superlattice**

In bulk Si and Ge, the damaged states are mainly vacancy and interstitial defects. As for Si/Ge SL, the associated defects contain vacancy, interstitial, and antisite defects and the mechanisms of defect generation are generally more complex. The discrepancy in the resistance to defect formation between bulk component materials and Si/Ge SL may result in their different radiation tolerances. To further investigate the origin of the different radiation responses of these semiconductor materials, we calculate the formation energies of vacancy, interstitial and antisite defects in bulk states and SL structures and the migration barrier of the most favorable defects employing density functional theory method. The computations are based on a supercell

**Table 3** The defect formation energies in bulk Si, Ge, and Si/Ge superlattice.  $V_X$ : X vacancy (X = Si or Ge);  $X_{int}$ : X interstitial (X = Si or Ge); FP defect: Frenkel pair defect

Defect type	Defect formation energies (eV)		
	Si/Ge SL	Bulk Ge	Bulk Si
$V_{Si}$	2.85	–	3.60, 3.61 <sup>a</sup> , 3.56 <sup>b</sup>
$V_{Ge}$	2.73	2.23, 2.09 <sup>a</sup>	–
$Si_{int}$	3.77	–	3.77, 3.75 <sup>c</sup> , 3.29 <sup>d</sup>
$Ge_{int}$	3.52	2.97, 2.92 <sup>e</sup>	–
Si FP	5.19	–	4.62, 4.26 <sup>b</sup>
Ge FP	5.01	4.15	–

<sup>a</sup>Ref. [39]

<sup>b</sup>Ref. [38]

<sup>c</sup>Ref. [40]

<sup>d</sup>Ref. [37]

<sup>e</sup>Ref. [42]

consisting of 64 atoms, with a  $6 \times 6 \times 6$  k-point sampling in real space and a cutoff energy of 500 eV.

The defect formation energies in bulk Si, Ge, and Si/Ge SL are listed in Table 3, along with other calculated results. In bulk Si, the formation energies for  $V_{Si}$ ,  $Si_{int}$ , and Si FP defects are calculated to be 3.60, 3.77, and 4.62 eV, respectively, which are in reasonable agreement with other calculations [37–40]. Our results indicate that the  $V_{Si}$  defect is easier to be created in bulk Si. Similarly, the  $V_{Ge}$  defect in bulk Ge is energetically more favorable than the  $Ge_{int}$  and Ge FP defects, as indicated by the smallest defect formation energy of 2.23 eV, which compares well with the theoretical value of 2.09 eV [39]. As for the Si/Ge SL, the formation energy of  $V_{Ge}$  is determined to be 2.73 eV, which is smaller than the formation energies of other defects. The next favorable defect is the  $V_{Si}$  defect, and the formation energy is determined to be 2.85 eV. It is noted that the value of 3.52 eV for  $Ge_{int}$  is smaller than the value of 3.77 eV for  $Si_{int}$  defect. As for FP defect, the formation energy is obviously larger, i.e., 5.19 eV for Si FP and 5.01 eV for Ge FP, suggesting that the FP defects are difficult to be created. As compared with the bulk states, the defect formation energies for Si/Ge SL structure are generally larger except for the defects of  $V_{Si}$  and  $Si_{int}$ , indicating that in SL structure, the point defects are generally more difficult to form. Such discrepancy in the resistance to defect formation between bulk states and Si/Ge SL structure may result in their different responses to irradiation.

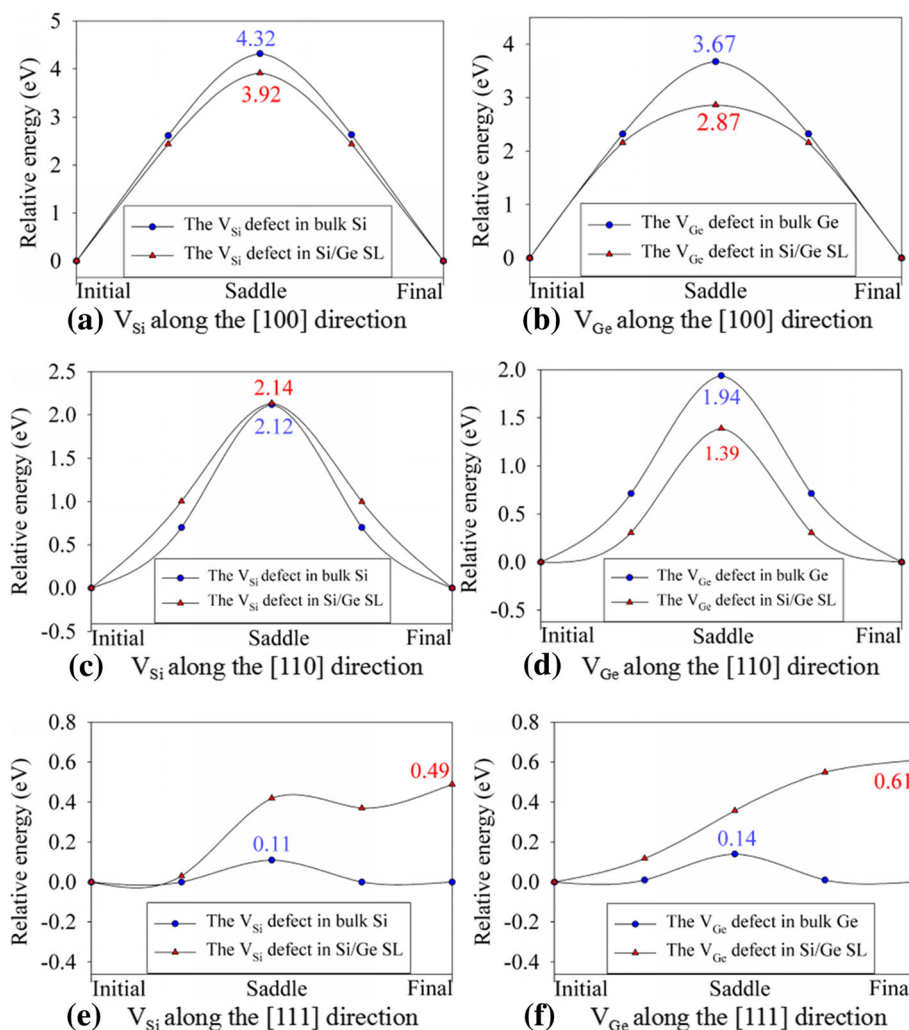
Based on the optimized structures, the migration behaviors of the  $V_{Ge}$  and  $V_{Si}$  defects that are the most favorable defects in bulk and Si/Ge SL structures are further investigated. The  $V_{Ge}$  and  $V_{Si}$  defects which are adjacent to the Si/Ge interface are taken into account, and the migration barriers are summarized in Table 4. It is noted that the migration barriers along the [100] and [110] directions for  $V_{Ge}$  defects are smaller than those for  $V_{Si}$  defects, and the energy barrier for  $V_{Ge}$  migration along the [111] direction is slightly larger than that for  $V_{Si}$  migration, which are consistent with the results reported by Cowern et al. [41].

**Table 4** The defect migration barrier in bulk Si, Ge, and Si/Ge superlattice.  $V_X$ : X vacancy (X = Si or Ge)

Defect type	Direction	Migration barrier (eV)		
		Si/Ge SL	Bulk Ge	Bulk Si
$V_{Si}$	[100]	3.92	–	4.32
	[110]	2.14	–	2.12, 2.85 <sup>a</sup>
	[111]	0.49	–	0.11
$V_{Ge}$	[100]	2.87	3.67	–
	[110]	1.39	1.94, 2.1 <sup>a</sup>	–
	[111]	0.61	0.14	–

<sup>a</sup>Ref. [41]





**Fig. 6** The migration barrier of silicon vacancy ( $V_{Si}$ ) and germanium vacancy ( $V_{Ge}$ ) defects obtained by a cluster nudged elastic band method. **a**  $V_{Si}$  along the [100] direction; **b**  $V_{Ge}$  along the [100] direction; **c**  $V_{Si}$  along the [110] direction; **d**  $V_{Ge}$  along the [110] direction; **e**  $V_{Si}$  along the [111] direction; **f**  $V_{Ge}$  along the [111] direction

The energy landscapes of defect migration along the [100], [110], and [111] directions are plotted in Fig. 6. In Fig. 6a, the migration barriers of the  $V_{Si}$  defect along the [100] direction are determined to be 4.32 and 3.92 eV in bulk Si and Si/Ge SL, respectively. As for the [110] direction, the migration barrier of 2.14 eV for  $V_{Si}$  in the Si/Ge SL structure is very close to the value of 2.12 eV in bulk Si. Comparing the migration barrier along each direction, we find that the [111] direction is the most favorable migration direction for both Si and Ge vacancies, as indicated by the significantly smaller migration barriers. Especially, the  $V_{Si}$  defects migrate more easily along the [111] direction in bulk Si than Si/Ge SL, since the energy barrier of 0.11 eV in the bulk state is much smaller (see Fig. 6e). As for the  $V_{Ge}$  defects, the migration barriers along the [100] direction are calculated to be 3.67 eV in bulk Ge and 2.87 eV in Si/Ge SL. In the

case of [110] direction, the energy barriers are determined to be 1.94 and 1.39 eV in the bulk and SL structures, respectively. Similar to the case of Si vacancy migration, the  $V_{Ge}$  defects are easier to migrate along the [111] direction. Also, the migration occurs more easily in bulk Ge than Si/Ge SL, as shown in Fig. 6f. Our calculations suggest that both Si and Ge vacancies are more mobile in the bulk states than SL structure, which may result in void formation and even volume swelling. This may contribute to different responses to irradiation for the bulk and SL structures.

**Conclusions**

In summary, low-energy displacement events in bulk Si, Ge, and Si/Ge superlattice (SL) have been investigated by an ab initio molecular dynamics method. In bulk Si and Ge, the threshold displacement energies are shown

to be dependent on the crystallographic direction and the atoms are more difficult to be displaced along the [110] direction. The damaged states in bulk states are mainly vacancy and interstitial defects. In the Si/Ge SL structure, the Si atoms are more resistant along the [111] direction, while the Ge atoms are more difficult to be displaced along the [110] direction. Our calculations show that the energies for the Ge recoils in the SL structure are generally comparable to those in the bulk Ge, whereas the energies for the Si recoils in the SL structure are generally much larger than those in bulk Si, indicative of enhanced radiation resistance of the Si/Ge SL. Defect formation energy calculations show that the point defects in the Si/Ge SL generally have higher formation energies, indicating that in the SL structure the point defects are generally more difficult to form. It is also found that the [111] direction is the most favorable migration path for both Si and Ge vacancies, and both vacancies are more mobile in the bulk states than in SL structure. Our calculations suggest that the enhanced radiation resistance of Si/Ge SL is beneficial to its application as electronic and optoelectronic devices under extreme working conditions like radiation.

#### Abbreviations

AIMD: Ab initio molecular dynamics;  $E_d$ : Threshold displacement energy; FP: Frenkel pair; Ge: Germanium;  $Ge_{int}$ : Germanium interstitial;  $Ge_{Si}$ : Germanium occupying the silicon lattice site; LDA: Local-density approximation; MD: Molecular dynamics; NVE: Microcanonical ensemble; PKA: Primary knock-on atom; PL: Photoluminescence; QD: Quantum dot; QW: Quantum well; Si: Silicon; SIESTA: Spanish Initiative for Electronic Simulations with Thousands of Atoms;  $Si_{Ge}$ : Silicon occupying the germanium lattice site;  $Si_{int}$ : Silicon interstitial; SL: Superlattice; SZP: Single- $\zeta$  basis sets plus polarization orbital;  $V_{Ge}$ : Germanium vacancy;  $V_{Si}$ : Silicon vacancy

#### Acknowledgements

The theoretical calculations were performed using the supercomputer resources at TianHe-1 located at National Supercomputer Center in Tianjin.

#### Funding

Haiyan Xiao was supported by the NSAF Joint Foundation of China (Grant No. U1530129). Zijiang Liu was supported by the National Natural Science Foundation of China (Grant No. 11464025) and the New Century Excellent Talents in University under Grant No. NECT-11-0906.

#### Availability of Data and Materials

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

#### Authors' Contributions

HX and XZ designed the calculations. MJ conducted the calculations and wrote the manuscript. SP, GY, ZL, and LQ contributed the discussion and interpretation of the results. All authors read and approved the final manuscript.

#### Competing Interests

The authors declare that they have no competing interests.

#### Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

#### Author details

<sup>1</sup>School of Physics, University of Electronic Science and Technology of China, Chengdu 610054, China. <sup>2</sup>Institute of Nuclear Physics and Chemistry, Chinese Academy of Engineering Physics, Mianyang 621900, China. <sup>3</sup>Department of Physics, Lanzhou City University, Lanzhou 730070, China.

Received: 22 February 2018 Accepted: 18 April 2018

Published online: 02 May 2018

#### References

- Garg J, Bonini N, Marzari N (2011) High thermal conductivity in short-period superlattices. *Nano Lett* 11:5135–5141
- Froyen S, Wood DM, Zunger A (1988) Structural and electronic properties of epitaxial thin-layer SiGe superlattices. *Phys Rev B* 37:6893–6907
- Hossain MZ, Johnson HT (2012) Electron-dependent thermoelectric properties in Si/Si<sub>1-x</sub>Ge<sub>x</sub> heterostructures and Si<sub>1-x</sub>Ge<sub>x</sub> alloys from first-principles. *Appl Phys Lett* 100:253901
- Menzel D, Koschinski W, Dettmer K, Schoenes J (1999) Photoconductivity of Si/Ge buffers, superlattices, and multiple quantum wells. *Thin Solid Films* 342:312–316
- Tair F, Sekkal N, Amrani B, Adli W, Boudaoud L (2007) First principles calculation of the opto-electronic properties of (110) growth axis SiGe superlattices. *Superlattices Microst* 41:44–55
- Thomas IO, Srivastava GP (2016) Lattice thermal conduction in ultra-thin nanocomposites. *J Appl Phys* 119:244309
- Sobolev NA, Ivlev GD, Gatskevich EI, Leitão JP, Fonseca A, Carmo MC, Lopes AB, Sharaev DN, Kibbel H, Presting H (2003) Pulsed laser annealing of Si–Ge superlattices. *Mater Sci Eng C-Mater* 23:19–22
- Yamaguchi M (1995) Radiation resistance of compound semiconductor solar cells. *J Appl Phys* 78:1476–1480
- Mamor M, Pipeleers B, Auret FD, Vantomme A (2011) Defect production in strained p-type Si<sub>1-x</sub>Ge<sub>x</sub> by Er implantation. *J Appl Phys* 109:013715
- Novikov AV, Yablonskiy AN, Platonov VV, Obolenskiy SV, Lobanov DN, Krasilnik ZF (2010) Effect of irradiation on the luminescence properties of low-dimensional SiGe/Si(001) heterostructures. *Semiconductors* 44:329–334
- Zheng B, Budak S, Muntele C, Xiao Z, Celaschi S, Mutele I, Chhay B, Zimmerman RL, Holland LR, Ila D (2011) Improvement on thermoelectric characteristics of layered nanostructure by ion beam bombardment. *MRS Proc* 929:013715
- Sobolev NA, Korshunov FP, Sauer R, Thonke K, König U, Presting H (1996) Influence of electron irradiation and annealing on the photoluminescence of Si/Ge superlattices and Si/Ge quantum wells. *J Cryst Growth* 167:502–507
- Fonseca A, Sobolev NA, Leitão JP, Alves E, Carmo MC, Zakharov ND, Werner P, Tonkikh AA, Cirlin GE (2006) Influence of defects on the optical and structural properties of Ge dots embedded in an Si/Ge superlattice. *J Lumin* 121:417–420
- Leitão JP, Santos NM, Sobolev NA, Correia MR, Stepina NP, Carmo MC, Magalhães S, Alves E, Novikov AV, Shaleev MV, Lobanov DN, Krasilnik ZF (2008) Radiation hardness of GeSi heterostructures with thin Ge layers. *Mater Sci Eng B-Solid* 147:191–194
- Budak S, Guner S, Smith C, Minamisawa RA, Zheng B, Muntele C, Ila D (2009) Surface modification of Si/Ge multi-layers by MeV Si ion bombardment. *Surf Coat Tech* 203:2418–2421
- Nylandsted Larsen A, O'Raifeartaigh C, Barklie RC, Holm B, Priolo F, Franzo G, Lulli G, Bianconi M, Nipoti R, Lindner JKN, Mesli A, Grob JJ, Cristiano F, Hemment PLF (1997) MeV ion implantation induced damage in relaxed Si<sub>1-x</sub>Ge<sub>x</sub>. *J Appl Phys* 81:2208–2218
- Sayed M, Jefferson JH, Walker AB, Cullis AG (1995) Computer simulation of atomic displacements in Si, GaAs, and AlAs. *Nucl Instrum Meth Phys Res Sect B-Beam Interact Mater Atoms* 102:232–235
- Windl W, Lenosky TJ, Kress JD, Voter AF (1998) First-principles investigation of radiation induced defects in Si and SiC. *Nucl. Instrum. Meth. Phys. Res. Sect. B-Beam Interact. Mater. Atoms* 141:61–65
- Caturla MJ, Díaz de la Rubia T, Marqués LA, Gilmer GH (1996) Ion-beam processing of silicon at keV energies: a molecular-dynamics study. *Phys Rev B* 54:16683–16695
- Holmström E, Nordlund K, Kuronen A (2010) Threshold defect production in germanium determined by density functional theory molecular dynamics simulations. *Phys Scripta* 81:035601
- Shaw MJ, Briddon PR, Jaros M (1996) Electronic structure of imperfect Si/Ge heterostructures. *Phys Rev B* 54:16781–16785

22. Xiao HY, Gao F, Zu XT, Weber WJ (2009) Threshold displacement energy in GaN: Ab initio molecular dynamics study. *J Appl Phys* 105:123527
23. Liu B, Xiao HY, Zhang Y, Aidhy DS, Weber WJ (2013) Ab initio molecular dynamics simulations of threshold displacement energies in SrTiO<sub>3</sub>. *J Phys-Condens Mat* 25:485003
24. Jiang M, Xiao HY, Zhang HB, Peng SM, Xu CH, Liu ZJ, Zu XT (2016) A comparative study of low energy radiation responses of SiC, TiC and ZrC. *Acta Mater* 110:192–199
25. Lucas G, Pizzagalli L (2005) Ab initio. *Phys Rev B* 72:161202
26. Xue SW, Zu XT, Zhou WL, Deng H, Xiang X, Zhang L, Deng H (2008) Effects of post-thermal annealing on the optical constants of ZnO thin film. *J Alloy Compd* 448:21–26
27. Zu XT, Yang L, Gao F, Peng SM, Heinisch HL, Long XG, Kurtz RJ (2009) Properties of helium defects in bcc and fcc metals investigated with density functional theory. *Phys Rev B* 80:1993–2006
28. Troullier N, Martins JL (1991) Efficient pseudopotentials for plane-wave calculations. *Phys Rev B* 43:1993–2006
29. Perdew JP, Burke K, Ernzerhof M (1996) Generalized gradient approximation made simple. *Phys Rev Lett* 77:3865–3868
30. Heyd J, Peralta JE, Scuseria GE, Martin RL (2005) Energy band gaps and lattice parameters evaluated with the Heyd-Scuseria-Ernzerhof screened hybrid functional. *J Chem Phys* 123:174101
31. Claudio T, Stein N, Petermann N, Stroppa DG, Koza MM, Wiggers H, Klobes B, Schierning G, Hermann RP (2016) Lattice dynamics and thermoelectric properties of nanocrystalline silicon-germanium alloys. *Phys Status Solidi A-Appl Mat* 213:515–523
32. Loferski JJ, Rappaport P (1958) Radiation damage in Ge and Si detected by carrier lifetime changes: damage thresholds. *Phys Rev* 111:432–439
33. Koike J, Parkin DM, Mitchell TE (1992) Displacement threshold energy for type IIa diamond. *Appl Phys Lett* 60:1450–1452
34. Corbett JW, Watkins GD (1965) Production of divacancies and vacancies by electron irradiation of silicon. *Phys Rev* 138:A555–A560
35. Gao F, Xiao H, Zu X, Posselt M, Weber WJ (2009) Defect-enhanced charge transfer by ion-solid interactions in SiC using large-scale ab initio molecular dynamics simulations. *Phys Rev Lett* 103:027405
36. Chen Y, MacKay JW (1969) Orientation and temperature dependence of electron damage in n-type germanium. *Philos Mag* 19:357–367
37. Rinke P, Janotti A, Scheffler M, Van de Walle CG (2009) Defect formation energies without the band-gap problem: combining density-functional theory and the GW approach for the silicon self-interstitial. *Phys Rev Lett* 102:026402
38. Goedecker S, Deutsch T, Billard L (2002) A fourfold coordinated point defect in silicon. *Phys Rev Lett* 88:235501
39. Ramprasad R, Zhu H, Rinke P, Scheffler M (2012) New perspective on formation energies and energy levels of point defects in nonmetals. *Phys Rev Lett* 108:066404
40. Lenosky TJ, Kress JD, Kwon I, Voter AF, Edwards B, Richards DF, Yang S, Adams JB (1997) Highly optimized tight-binding model of silicon. *Phys Rev B* 55:1528–1544
41. Cowern NE, Simdyankin S, Ahn C, Bennett NS, Goss JP, Hartmann JM, Pakfar A, Hamm S, Valentin J, Napolitani E, De Salvador D, Bruno E, Mirabella S (2013) Extended point defects in crystalline materials: Ge and Si. *Phys Rev Lett* 110:155501
42. Kamiyama E, Sueoka K, Vanhellefont J (2013) Formation energy of intrinsic point defects in Si and Ge and implications for Ge crystal growth. *ECS J Solid State Sci Technol* 2:P104–P109

Submit your manuscript to a SpringerOpen<sup>®</sup> journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

---

Submit your next manuscript at ► [springeropen.com](http://springeropen.com)

---