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# 5-nm LiF as an Efficient Cathode Buffer Layer in Polymer Solar Cells Through Simply Introducing a C<sub>60</sub> Interlayer

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## Abstract

Lithium fluoride (LiF) is an efficient and widely used cathode buffer layer (CBL) in bulk heterojunction polymer solar cells (PSCs). The LiF thickness is normally limited to 1 nm due to its insulting property. Such small thickness is difficult to precise control during thermal deposition, and more importantly, 1-nm-thick LiF cannot provide sufficient protection for the underlying active layer. Herein, we demonstrated the application of a very thick LiF as CBL without sacrificing the device efficiency by simply inserting a C<sub>60</sub> layer between the active layer and LiF layer. The devices with the C<sub>60</sub>/LiF (5 nm) double CBLs exhibit a peak power conversion efficiency (PCE) of 3.65%, which is twofold higher than that (1.79%) of LiF (5 nm)-only device. The superior performance of the C<sub>60</sub>/LiF (5 nm)-based devices is mainly attributed to the good electrical conductivity of the C<sub>60</sub>/LiF (5 nm) bilayer, arising from the intermixing occurred at the C<sub>60</sub>/LiF interface. Besides, the formation of a P3HT/C<sub>60</sub> subcell and the optical spacer effect of C<sub>60</sub> also contribute to the increase in short-circuit current density ( $J_{sc}$ ) of the device. With further increase of LiF thickness to 8 nm, a PCE of 1.10% is attained for the C<sub>60</sub>/LiF-based device, while the negligible photovoltaic performance is observed for the LiF-only device. All in all, our results show that C<sub>60</sub>/LiF bilayer is a promising alternative to LiF single layer due to its high tolerance to the LiF thickness variations.

Keywords: Polymer solar cells, Thick LiF buffer layer, C<sub>60</sub>/LiF bilayer, Mixed morphology

## Background

Solution-processed bulk heterojunction polymer solar cells (PSCs) have received increasing attention in recent decades because of their potential advantages such as low cost, light weight, and possibility to fabricate large-scale, flexible, and semitransparent devices [1–5]. By far, the relatively low power conversion efficiency (PCE) compared to silicon-based solar cells is still a major limitation that hinders their practical application. To achieve commercialization of this promising technology, extensive research efforts have focused on increasing the efficiency of PSCs. Until now, PCEs in the range of 11–13% have been demonstrated, primarily owing to the development of novel conjugated polymer donor and non-fullerene acceptor materials [6–12]. Besides, the introduction of

anode/cathode buffer layer between the active layer and the electrode provides another efficient means to improve the device performance [13–21].

PSCs can be divided into conventional and inverted structures according to whether the indium-tin-oxide (ITO) electrode serves as the anode or the cathode. For the conventional PSCs with ITO as anode, a low work function metal such as Ca is commonly used as cathode buffer layer (CBL) to reduce the work function of the cathode (e.g., Al, Ag). However, Ca is easily oxidized when exposed to air, resulting in the poor stability of the devices. Another widely used CBL in PSCs is lithium fluoride (LiF), which has been demonstrated to enhance the device performance through the formation of an interfacial dipole at the cathode interface [22]. Nevertheless, the thickness of LiF is limited to less than 2 nm (generally  $\sim 1$  nm) due to its insulating property [23, 24]. Such a small thickness is very difficult to be controlled via thermal deposition. Furthermore, 1-nm-thick LiF cannot provide sufficient



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protection for the underlying active layer during the evaporation of hot metal atoms [17, 25].

To address these problems, we have previously reported five stacks of  $C_{60}$ /LiF CBL, which substantially improved the device efficiency and stability of PSCs due to its good electrical conductivity even though a very thick LiF was used [26]. However, the fivestacked C<sub>60</sub>/LiF film was prepared by alternating deposition of C<sub>60</sub> and LiF layers. This preparation process is very complicated and time-consumed, and significantly increases the cost of device fabrication. In this work, we adopted a C<sub>60</sub>/LiF bilayer as CBL to achieve the same effect as five-stacked C<sub>60</sub>/LiF CBL. After depositing a  $C_{60}$  layer prior to the LiF evaporation, a thick LiF is allowed to be used without sacrificing the device efficiency. The PSCs with C<sub>60</sub>/LiF double CBLs maintained a ~3% PCE over a wide range of LiF thickness  $(1 \sim 6 \text{ nm})$ , and showed a PCE of 1.10% even at a very thick LiF, 8 nm. In contrast, the PSCs with LiF single CBL exhibited a rapid decrease of PCE with increasing LiF thickness and had negligible photovoltaic performance at LiF thickness of 8 nm. Besides, the peak efficiency (3.77%) of  $C_{60}$ /LiF-based devices is ~ 23% higher than that (3.06%) of LiF-only device. Taken all together, these results indicate that C<sub>60</sub>/LiF bilayer is a more promising candidate as a CBL compared to single LiF layer.

#### Methods

#### **Fabrication of PSCs**

ITO-coated glass substrates (Delta Technologies, LTD) were cleaned in acetone and isopropyl alcohol (IPA) under sonication for 5 min each and then treated by  $O_2$  plasma for 60 s to generate the hydrophilic surface. The filtered poly(3,4-ethylenedioxythiophene):poly(styrenesulfonate)

(PEDOT:PSS) solution (H. C. Starck, Clevios PH 500) was spin-coated onto the cleaned glass/ITO substrates at a speed of 2000 rpm for 50 s, followed by baking at 110 °C for 20 min under nitrogen atmosphere. Subsequently, the samples were transferred to a N<sub>2</sub>-purged glovebox (< 0.1 ppm O<sub>2</sub> and H<sub>2</sub>O) for spin-coating of photoactive layer.

P3HT (Rieke Metals Inc., 4002-EE, 91–94% regioregularity) and PCBM (American Dye Source, purity > 99.5%) were dissolved in chlorobenzene with a weight ratio of 1:1. The mixed solution was filtered using a 0.45 µm filter and then spin-coated on top of the PEDOT:PSS layer at 1000 rpm for 50 s, followed by thermal annealing at 130 °C for 20 min, which produced a ~ 160-nm-thick active layer measured using a Dektek surface profiler. The  $C_{60}$ , LiF, and Al (75 nm) electrode were sequentially deposited by thermal evaporation at a base pressure of 1 × 10<sup>-6</sup> mbar. The deposition rate and film thickness were monitored with a quartz crystal sensor. A circular-shaped shadow mask of 1 mm diameter was put on the sample to define the active area before the Al deposition.

#### Characterization

The current density-voltage (J-V) characteristics were measured using a Keithley 2400 system under simulated Air Mass 1.5 Global (AM 1.5 G) solar illumination at an intensity of 100 mW/cm<sup>2</sup>, which was calibrated by a power meter (OPHIR, Nova-Oriel) and a reference silicon solar cell. The measurements were carried out with the PSCs inside the glovebox. Atomic force microscope (AFM) images were taken with a Veeco Dimension-Icon AFM operated in tapping mode. Absorption spectra were obtained using a Varian Cary 50 UV/Vis spectrophotometer. Photo-induced charge extraction by linearly increasing voltage (Photo-CELIV) measurements were performed on PSCs under ambient conditions. A pulsed N<sub>2</sub> laser (337.1 nm, 1.4 ns) was used to generate the charge carriers, which were then extracted by a reversebias voltage ramp that was applied after 100 µs delay time. The current transients were recorded using a digital storage oscilloscope (50  $\Omega$  input impedance). During and after illumination, an offset voltage was applied to compensate the built-in potential of the devices, which prevents an initial photocurrent prior to the application of the voltage ramp. The mobility of the carriers can be calculated according to the following equation [27, 28]:

$$\mu = \frac{2d^2}{3At_{\max}^2 \left[1 + 0.36\frac{\Delta j}{j(0)}\right]}$$
(1)

where  $\mu$  is the charge carrier mobility, *d* is the thickness of the active layer, *A* is the voltage rise speed,  $t_{\text{max}}$  is the time when the extraction current reaches the maximum value,  $\Delta j$  is the current extraction peak height, and *j*(0) is the displacement current of the capacitance.

#### **Results and Discussion**

Figure 1 shows the *J*-*V* characteristics, recorded under 100 mW/cm<sup>2</sup> illumination (AM 1.5 G), of the PSCs with and without different thicknesses of  $C_{60}$  sandwiched between the active layer and 5-nm-thick LiF layer. The device without the  $C_{60}$  layer shows S-shaped curve, resulting in the low fill factor (FF) and therefore the low PCE, despite the typical short-circuit current density ( $J_{sc}$ ) and open-circuit voltage ( $V_{oc}$ ). The low FF is rationalized in terms of the insulating property of LiF, which blocks the electron injection/extraction when the LiF layer is too thick and thus leads to the large series resistance ( $R_{s}$ ) and small shunt resistance ( $R_{sh}$ ) of the device as shown in Table 1 ( $R_{s}$  and  $R_{sh}$  were calculated from the inverse slope of photo *J*-*V* curve at 0 mA/cm<sup>2</sup> and 0 V, respectively). As for the  $J_{sc}$ , the normal value



 $(9.23 \text{ mA/cm}^2)$  implies that the built-in electric filed inside the device (from work function difference between anode and cathode) is sufficient to promote the electron transport through LiF (5 nm) CBL by tunneling. After introducing 3-nm-thick C<sub>60</sub> layer between P3HT:PCBM and LiF (5 nm) layers, the S-shape disappears and the FF increases significantly from 32.4 to 56.3%. The increased FF arises from the reduced  $R_{s}$ , which implies that the C<sub>60</sub> (3 nm)/LiF (5 nm) bilayer possesses better electrical conductivity than single LiF (5 nm) layer. With the increase of C<sub>60</sub> thickness, the FF first increases, reaching a maximum value of 67% at 8 nm and then decreases slightly with further increasing  $C_{60}$  thickness. Due to the recovery of FF, the  $C_{60}/\text{LiF}$  (5 nm)-based devices show a maximum PCE of 3.65%, which is two times higher than that (1.79%) of LiF (5 nm)-only device. To demonstrate the reproducibility of the results, the average photovoltaic parameters and standard deviations of the studied devices were calculated from a batch of five devices, as shown in Additional file 1: Table S1. For each device, all the parameters including  $J_{sc}$ ,  $V_{oc}$ , FF, and PCE are highly reproducible with little variation, which validates the reliability of the results presented in Table 1.

To find out the reasons leading to the high FF for the  $C_{60}/\text{LiF}$  (5 nm)-based PSCs, AFM measurements were performed to examine the morphology of the LiF layer on the  $C_{60}$  surface. Figure 2 shows the height (top) and phase (bottom) images, recorded by tapping mode AFM, of the P3HT:PCBM films without and with the  $C_{60}$  (35 nm), LiF (5 nm), and  $C_{60}$ (35 nm)/LiF (5 nm) layers deposited on top (image size 500 nm × 500 nm). The pristine P3HT:PCBM film exhibits a very smooth surface with a low rootmean-square (rms) roughness of 0.81 nm (height image) and shows fibrillar crystalline domains of P3HT (phase image) [29]. After depositing 35-nmthick C<sub>60</sub> and 5-nm-thick LiF on top, the rms roughness increases to 1.36 and 1.67 nm, respectively. Although there is no significant difference in rms roughness between the top C<sub>60</sub> and LiF layers, the surface morphologies of these two films are very different. The 35-nm-thick C<sub>60</sub> shows larger aggregates (spherical shape) as compared to 5-nm-thick LiF, which can also be observed in their phase images. When depositing the C<sub>60</sub> (35 nm)/LiF (5 nm) bilayer on the P3HT:PCBM film, both the  $C_{60}$  (large size) and LiF (small size) aggregates are observed, indicating that the underlying  $C_{60}$  layer is not completely covered by 5-nm-thick LiF. Therefore, some intermixing occurs at the C<sub>60</sub>/LiF interface, which results in the good electrical conductivity of C<sub>60</sub>/LiF (5 nm) bilayer considering the percolation path formed by C<sub>60</sub> molecules.

To further investigate the influence of the  $C_{60}/\text{LiF}$  double CBLs on the device performance of PSCs, we fix the  $C_{60}$  thickness at the optimum value of 25 nm while changing the LiF thickness from 0.5 to 8 nm. For comparison, the devices with LiF single CBL were also fabricated. Figure 3 shows the *J*-*V* characteristics, recorded under 100 mW/cm<sup>2</sup> illumination (AM 1.5 G), of the PSCs using LiF single and  $C_{60}/\text{LiF}$  double CBLs with

**Table 1** Photovoltaic parameters for the P3HT:PCBM-based PSCs with and without different thicknesses of  $C_{60}$  inserted between the active layer and 5-nm-thick LiF layer

CBL	J <sub>sc</sub> (mA/cm <sup>2</sup> )	$V_{\rm oc}$ (V)	FF (%)	PCE (%)	$R_{\rm s}~(\Omega~{\rm cm}^2)$	$R_{\rm sh}~(\Omega~{\rm cm}^2)$
LiF	9.23	0.60	32.4	1.79	87.21	232.70
C <sub>60</sub> (3 nm)/LiF	9.28	0.58	56.3	3.03	7.74	293.46
C <sub>60</sub> (5 nm)/LiF	8.74	0.57	66.7	3.32	6.03	732.93
C <sub>60</sub> (8 nm)/LiF	9.13	0.57	67.0	3.48	6.22	768.36
C <sub>60</sub> (12 nm)/LiF	9.34	0.58	65.1	3.53	6.14	655.67
C <sub>60</sub> (15 nm)/LiF	9.49	0.60	64.0	3.65	5.24	351.63
C <sub>60</sub> (25 nm)/LiF	9.97	0.56	65.3	3.65	6.13	726.52
C <sub>60</sub> (35 nm)/LiF	8.97	0.56	61.9	3.11	8.19	751.58



varying thicknesses of LiF. The corresponding photovoltaic parameters of the devices are summarized in Table 2. The devices with LiF single CBL have a maximum PCE of 3.06% at the optimal LiF thickness of 1 nm. Further increasing the thickness leads to a rapid decrease in PCE to 0.79% at 6 nm and 0.06% at 8 nm. In contrast, the devices with C<sub>60</sub> (25 nm)/LiF double CBLs exhibit improved performance with a peak efficiency of 3.77% at the LiF thickness of 1 nm. More importantly, as the thickness increases to 6 and 8 nm, PCEs of 2.65 and 1.10% are attained, respectively, which are significantly higher than those of LiF-only devices. It should be mentioned that the results presented in Table 2 is also highly reproducible, as demonstrated by the very small standard deviations of the device characteristic parameters (Additional file 1: Table S2). For instance, the standard deviation of the device efficiency is less than 0.2% (0.1% for most devices), indicating high reproducibility. Furthermore, the average PCE shows the same trend as observed in Table 2, which implies that the comparison of efficiency among different groups is reliable.

As shown in Table 2, the improvement in PCE for the  $C_{60}$  (25 nm)/LiF-based PSCs mainly arises from the increase in FF and  $J_{\rm sc}$  due to the reduced  $R_{\rm s}$ . To better understand the  $R_{\rm s}$  reduction, we investigate the charge transport properties of the LiF single layer and  $C_{60}$ /LiF bilayer using the photo-CELIV technique [30, 31]. Additional file 1: Figure S1 shows the photo-CELIV current transients, recorded at varying voltage rise speeds, for the PSCs with the LiF single and  $C_{60}$ /LiF double CBLs. In photo-CELIV, the time of extraction current maximum ( $t_{\rm max}$ ) is used for estimating the charge carrier mobility according to Eq. 1 [27]. The calculated



CBL	J <sub>sc</sub> (mA/cm <sup>2</sup> )	$V_{\rm oc}$ (V)	FF (%)	PCE (%)	$R_{\rm s}$ ( $\Omega$ cm <sup>2</sup> )	$R_{\rm sh}~(\Omega~{\rm cm}^2)$
LiF (0.5 nm)	9.15	0.59	53.8	2.91	5.50	208.94
LiF (1 nm)	9.21	0.58	57.2	3.06	6.14	253.56
LiF (6 nm)	7.20	0.58	18.9	0.79	195.21	62.65
LiF (8 nm)	1.05	0.58	10.2	0.06	4749.92	175.15
C <sub>60</sub> /LiF (0.5 nm)	10.03	0.59	60.2	3.57	4.95	323.33
C <sub>60</sub> /LiF (1 nm)	10.10	0.56	66.6	3.77	5.05	626.63
C <sub>60</sub> /LiF (6 nm)	8.21	0.54	59.8	2.65	8.98	744.54
C <sub>60</sub> /LiF (8 nm)	6.96	0.54	29.2	1.10	14.08	82.15

**Table 2** Photovoltaic parameters for the P3HT:PCBM-based PSCs using LiF single and  $C_{60}$  (25 nm)/LiF double CBLs with different thicknesses of LiF

mobilities of the LiF (6 nm)-only device are 3.71, 3.40, and 3.59  $\times$  10<sup>-5</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for the voltage slopes of 10, 20, and 30 kV/s, respectively, implying that the mobility is independent on the voltage rise speed. In contrast, the estimated mobilities of the  $C_{60}$  (25 nm)/LiF (6 nm)-based device are 3.81, 3.56, and  $3.09 \times 10^{-4}$  cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for the voltage slopes of 10, 20, and 30 kV/s, respectively, which are one order of magnitude higher than those of the LiF (6 nm)-only device. The increased mobility after introducing a C<sub>60</sub> layer can be attributed to the improved electrical conductivity arising from the intermixing occurred at the  $C_{60}$ /LiF interface. In addition, it is noted that the photo-CELIV peak for the LiF (6 nm)-only device is broader than that for the  $C_{60}$ (25 nm)/LiF (6 nm)-based device, which indicates a more dispersive charge transport resulting from the larger imbalance between the electron and hole mobilities [32, 33]. This imbalance is attributed to the extremely low electron mobility for the LiF (6 nm)-only device considering that the extraction of electrons is blocked by the thick LiF layer. The accumulated electrons at the P3HT:PCBM/LiF interface screen the applied electric field and thereby decrease the rate of charge extraction in the device. In contrast, the narrow peak for the C<sub>60</sub> (25 nm)/LiF (6 nm)-based device implies the balanced electron and hole mobilities as well as the improved electron extraction owing to the good conductivity of the  $C_{60}$ (25 nm)/LiF (6 nm) bilayer.

Besides the significant improvement in FF, the  $J_{\rm sc}$  is slightly enhanced after incorporation of  $C_{60}$  (25 nm) layer. Considering that the spin-coated P3HT:PCBM blend film consists of a P3HT-rich region near the top surface [34, 35], we speculate that the excitons generated in this region can be dissociated at the P3HT/ $C_{60}$  interface for  $C_{60}$  (25 nm)/LiF-based devices, which leads to the increased  $J_{\rm sc}$  compared to the devices without the  $C_{60}$  interlayer. To verify this speculation, we fabricated the PSCs with a device structure of ITO/PEDOT:PSS/  $P3HT/C_{60}$  (25 nm)/LiF/Al, where the thickness of P3HT is varied from 5 to 100 nm. Figure 4 shows the J-V characteristics of these devices under 100 mW/cm<sup>2</sup> illumination (AM 1.5 G), and the corresponding photovoltaic parameters are summarized in Additional file 1: Table S3. It is found that the  $J_{sc}$  of P3HT/C<sub>60</sub>-based solar cells increases as the P3HT thickness decreases, which is rationalized in terms of the limited exciton diffusion length in P3HT (~ 10 nm). The  $J_{sc}$  reaches a maximum value of 1.34 mA/cm<sup>2</sup> at the P3HT thickness of 10 nm and then drops with further decease of the thickness to 5 nm due to the insufficient absorption. As mentioned above, such P3HT/C<sub>60</sub> subcell is most likely formed after depositing 25-nm-thick C<sub>60</sub> on top of the P3HT:PCBM active layer, which results in a 1.34 mA/cm<sup>2</sup> increase in  $J_{\rm sc}$  under ideal conditions for C<sub>60</sub>/LiF-based devices [36]. By comparing the  $J_{\rm sc}$  values of the devices with and without the  $C_{60}$  (25 nm) interlayer, the enhancement in  $J_{sc}$  is around 1 mA/cm<sup>2</sup> (except for the LiF (8 nm)-based devices), which is consistent with our speculation.

After introducing a  $C_{60}$  layer between the P3HT:PCBM and LiF layers, the optical field distribution within the



TO/PEDOT:PSS/P3HT (x nm)/C<sub>60</sub> (25 nm)/LiF (1 nm)/Al using varying thicknesses of P3HT

solar cell is most likely altered, which will cause the variation in  $J_{sc}$  [26, 37]. To investigate this effect, we first simulated the electric field intensity inside the P3HT:PCBM active layer for the devices with and without the  $C_{60}$ interlayer. As shown in Additional file 1: Figure S2a, the integrated field intensity for the devices incorporating a  $C_{60}$  layer is weaker in the short-wavelength region and stronger in the long-wavelength region as compared to the device without  $C_{60}$  interlayer. This trend becomes more remarkable, and simultaneously, a red shift is observed with increasing the  $C_{60}$  thickness. Additional file 1: Figure S2b shows the absorption spectra of the pristine C<sub>60</sub> film, and the P3HT:PCBM films with and without different CBLs deposited on top. Comparing the absorption spectra of P3HT:PCBM/C<sub>60</sub> (25 nm) films with and without 8-nm-thick LiF, the two curves overlap almost completely, indicating that LiF does not absorb visible light. On the other hand, the P3HT:PCBM/C<sub>60</sub> films have higher absorption in the wavelength ranges of 400~510 nm and 580~680 nm when compared to the pristine P3HT:PCBM film. This absorption enhancement becomes more pronounced with increasing  $C_{60}$  thickness. Intuitively, the absorption enhancement in the 400~510 nm wavelength range arises from the  $C_{60}$  absorption (400~550 nm). Additional file 1: Figure S2c shows the incident photon-tocurrent conversion efficiency (IPCE) spectra of the PSCs with LiF (5 nm) single and C<sub>60</sub> (25 nm)/LiF (5 nm) double CBLs. Compared to the LiF-only device, the device with  $C_{60}$ /LiF double CBLs has a lower IPCE at the short wavelengths due to the parasitic absorption in the  $C_{60}$  film, and shows a higher IPCE at long wavelengths, owing to the optical spacer effect as well as the contribution of  $P3HT/C_{60}$ subcell.

From Table 2, it is noticed that the  $C_{60}$  (25 nm)/ LiF (8 nm)-based device exhibits a low PCE of 1.10% although this efficiency is still much higher than that (0.06%) of the LiF (8 nm)-only device. The low PCE is the result of the small  $J_{\rm sc}$  and FF, which is caused by the large  $R_{\rm s}$ . As discussed above, the C<sub>60</sub> (35 nm)/ LiF (5 nm) film has good electrical conductivity due to the formation of the mixed morphology between C<sub>60</sub> and LiF layers (see Fig. 2). To find the reason for the high resistance of the  $C_{60}$  (25 nm)/LiF (8 nm) film, AFM measurements were performed on P3HT:PCBM films without and with the  $C_{60}$  (25 nm), LiF (8 nm), and C<sub>60</sub> (25 nm)/LiF (8 nm) layers deposited on top. As shown in Additional file 1: Figure S3, large spherical aggregates are formed in the  $C_{60}$ (25 nm) film while relatively small aggregates are formed in the LiF (8 nm) film, which is similar to the observation in Fig. 2. When depositing 8-nm-thick LiF on top of the  $C_{60}$  (25 nm) layer, the morphology (small aggregates) is very similar to that of the pristine LiF film, indicating that the underlying  $C_{60}$  aggregates are completely covered by 8-nm-thick LiF. Therefore, we speculate that a thick LiF accumulates at the top of the  $C_{60}$  (25 nm)/LiF (8 nm) bilayer film, which hinders the electron extraction and therefore leads to the high  $R_s$  of the device.

#### Conclusions

In summary, we have demonstrated that a thick LiF can be used as CBL in P3HT:PCBM-based PSCs by simply introducing a C<sub>60</sub> layer between the active layer and the LiF layer. The devices with the  $C_{60}/LiF$ (5 nm) double CBLs exhibit a peak efficiency of 3.65%, while the LiF (5 nm)-only device shows a two times lower PCE of 1.79%. The improved device performance mainly results from the high FF due to the good electrical conductivity of the  $C_{60}/\text{LiF}$  bilayer. In addition, the  $J_{\rm sc}$  is also improved after introducing a  $C_{60}$  interlayer, which can be attributed to the contribution of P3HT/C<sub>60</sub> subcell as well as the optical spacer effect of C<sub>60</sub>. Further increasing the LiF thickness to 8 nm leads to the rapid decrease of PCE to 1.10 and 0.06% for the C<sub>60</sub>/LiF-based device and LiFonly device, respectively. The decline in PCE of the device with  $C_{60}/\text{LiF}$  (8 nm) double CBLs is caused by the impeded electron transport, owing to the accumulated LiF at the top of the C<sub>60</sub> (25 nm)/LiF (8 nm) bilayer. All in all, these results indicate that the  $C_{60}/LiF$ bilayer is a more promising CBL as compared to LiF single layer for fabricating highly efficient and largescale PSCs.

#### Additional file

Additional file 1: Supporting information. Table S1. Average photovoltaic performance parameters for the P3HT:PCBM-based PSCs with and without different thicknesses of C<sub>60</sub> inserted between the active layer and 5-nm-thick LiF layer. Table S2. Average photovoltaic performance parameters for the P3HT:PCBM-based PSCs using LiF single and C<sub>60</sub> (25 nm)/LiF double CBLs with different thicknesses of LiF. Figure S1. Photo-CELIV curves for the devices with (a) the LiF (6 nm) single and (b) C<sub>60</sub> (25 nm)/LiF (6 nm) double CBLs. Table S3. Photovoltaic parameters of the P3HT/C<sub>60</sub> (25 nm)-based PSCs with the P3HT thickness varied from 5 to 100 nm. Figure S2. (a) Simulated electric field intensity within the active layer versus the thickness of C<sub>60</sub> interlayer for the PSCs having the following structure: ITO (150 nm)/PEDOT:PSS (45 nm)/ P3HT:PCBM (180 nm)/C<sub>60</sub> (x nm)/Al (120 nm). (b) Absorption spectra of the pristine C<sub>60</sub> film and the P3HT:PCBM blend films with and without different CBLs deposited on top. (c) Incident photon-to-current conversion efficiency (IPCE) spectra for the devices with and without the  $C_{60}$  interlayer. Figure S3. AFM height (top) and phase (bottom) images of  $C_{60}$  (25 nm), LiF (8 nm), and  $C_{60}$ (25 nm)/LiF (8 nm) layers deposited on P3HT:PCBM blend films. (DOC 1663 kb)

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#### Authors' Contributions

XDL carried out the experiments, analyzed the data, and wrote the manuscript. LJG and YHZ provided helpful suggestions. All authors read and approved the final manuscript.

#### **Competing Interests**

The authors declare that they have no competing interests.

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