

Research Article

Littlewood-Paley g -Functions and Multipliers for the Laguerre Hypergroup

Jizheng Huang^{1,2}

¹ College of Sciences, North China University of Technology, Beijing 100144, China

² CEMA, Central University of Finance and Economics, Beijing 100081, China

Correspondence should be addressed to Jizheng Huang, hjzheng@163.com

Received 4 November 2010; Accepted 13 January 2011

Academic Editor: Shusen Ding

Copyright © 2011 Jizheng Huang. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

Let $L = -(\partial^2/\partial x^2 + (2\alpha + 1/x)(\partial/\partial x) + x^2(\partial^2/\partial t^2))$; $(x, t) \in (0, +\infty) \times \mathbb{R}$, where $\alpha \geq 0$. Then L can generate a hypergroup which is called Laguerre hypergroup, and we denote this hypergroup by \mathbf{K} . In this paper, we will consider the Littlewood-Paley g -functions on \mathbf{K} and then we use it to prove the Hörmander multipliers on \mathbf{K} .

1. Introduction and Preliminaries

In [1], the authors investigated Littlewood-Paley g -functions for the Laguerre semigroup. Let

$$\mathcal{L}_\alpha = \sum_{i=1}^d x_i \frac{\partial^2}{\partial x_i^2} + (\alpha_i + 1 - x_i) \frac{\partial}{\partial x_i}, \quad (1.1)$$

where $\alpha = (\alpha_1, \dots, \alpha_d)$, $x_i > 0$, then define the following Littlewood-Paley function \mathcal{G}_α by

$$\mathcal{G}_\alpha f(x) = \left(\int_0^\infty |t \nabla_\alpha P_t^\alpha f(x)|^2 \frac{dt}{t} \right)^{1/2}, \quad (1.2)$$

where $\nabla_\alpha = (\partial_t, \sqrt{x_1} \partial_{x_1}, \dots, \sqrt{x_d} \partial_{x_d})$ and P_t^α is the Poisson semigroup associated to \mathcal{L}_α . In [1], the authors prove that \mathcal{G}_α is bounded on $L^p(\mu_\alpha)$ for $1 < p < \infty$. In this paper, we consider the following differential operator

$$L = - \left(\frac{\partial^2}{\partial x^2} + \frac{2\alpha + 1}{x} \frac{\partial}{\partial x} + x^2 \frac{\partial^2}{\partial t^2} \right); \quad (x, t) \in (0, +\infty) \times \mathbb{R}, \quad (1.3)$$

where $\alpha \geq 0$. It is well known that it can generate a hypergroup (cf. [2, 3] or [4]). We will define Littlewood-Paley g -functions associated to L and prove that they are bounded on $L^p(\mathbf{K})$ for $1 < p < \infty$. As an application, we use it to prove the Hörmander multiplier theorem on \mathbf{K} .

Let $\mathbf{K} = [0, \infty) \times \mathbf{R}$ equipped with the measure

$$dm_\alpha(x, t) = \frac{1}{\pi\Gamma(\alpha + 1)} x^{2\alpha+1} dx dt, \quad \alpha \geq 0. \quad (1.4)$$

We denote by $L^p_\alpha(\mathbf{K})$ the spaces of measurable functions on \mathbf{K} such that $\|f\|_{\alpha,p} < +\infty$, where

$$\begin{aligned} \|f\|_{\alpha,p} &= \left(\int_{\mathbf{K}} |f(x, t)|^p dm_\alpha(x, t) \right)^{1/p}, \quad 1 \leq p < \infty, \\ \|f\|_{\alpha,\infty} &= \operatorname{esssup}_{(x,t) \in \mathbf{K}} |f(x, t)|. \end{aligned} \quad (1.5)$$

For $(x, t) \in \mathbf{K}$, the generalized translation operators $T_{(x,t)}^{(\alpha)}$ are defined by

$$\begin{aligned} &T_{(x,t)}^{(\alpha)} f(y, s) \\ &= \begin{cases} \frac{1}{2\pi} \int_0^{2\pi} f\left(\sqrt{x^2 + y^2 + 2xy \cos \theta}, s + t + xy \sin \theta\right) d\theta, & \text{if } \alpha = 0, \\ \frac{\alpha}{\pi} \int_0^{2\pi} \int_0^1 f\left(\sqrt{x^2 + y^2 + 2xyr \cos \theta}, s + t + xy r \sin \theta\right) r(1-r^2)^{\alpha-1} dr d\theta, & \text{if } \alpha > 0. \end{cases} \end{aligned} \quad (1.6)$$

It is known that $T_{(x,t)}^{(\alpha)}$ satisfies

$$\|T_{(x,t)}^{(\alpha)} f\|_{\alpha,p} \leq \|f\|_{\alpha,p}. \quad (1.7)$$

Let $M_b(\mathbf{K})$ denote the space of bounded Radon measures on \mathbf{K} . The convolution on $M_b(\mathbf{K})$ is defined by

$$(\mu * \nu)(f) = \int_{\mathbf{K} \times \mathbf{K}} T_{(x,t)}^{(\alpha)} f(y, s) d\mu(x, t) d\nu(y, s). \quad (1.8)$$

It is easy to see that $\mu * \nu = \nu * \mu$. If $f, g \in L^1_\alpha(\mathbf{K})$ and $\mu = f m_\alpha$, $\nu = g m_\alpha$, then $\mu * \nu = (f * g) m_\alpha$, where $f * g$ is the convolution of functions f and g defined by

$$(f * g)(x, t) = \int_{\mathbf{K}} T_{(x,t)}^{(\alpha)} f(y, s) g(y, -s) dm_\alpha(y, s). \quad (1.9)$$

The following lemma follows from (1.7).

Lemma 1.1. *Let $f \in L^1_\alpha(\mathbf{K})$ and $g \in L^p_\alpha(\mathbf{K})$, $1 \leq p \leq \infty$. Then*

$$\|f * g\|_{\alpha,p} \leq \|f\|_{\alpha,1} \|g\|_{\alpha,p}. \tag{1.10}$$

$(\mathbf{K}, *, i)$ is a hypergroup in the sense of Jewett (cf. [5, 6]), where i denotes the involution defined by $i(x, t) = (x, -t)$. If $\alpha = n - 1$ is a nonnegative integer, then the Laguerre hypergroup \mathbf{K} can be identified with the hypergroup of radial functions on the Heisenberg group \mathbf{H}^n .

The dilations on \mathbf{K} are defined by

$$\delta_r(x, t) = (rx, r^2t), \quad r > 0. \tag{1.11}$$

It is clear that the dilations are consistent with the structure of hypergroup. Let

$$f_r(x, t) = r^{-(2\alpha+4)} f\left(\frac{x}{r}, \frac{t}{r^2}\right). \tag{1.12}$$

Then we have

$$\|f_r\|_{\alpha,1} = \|f\|_{\alpha,1}. \tag{1.13}$$

We also introduce a homogeneous norm defined by $\|(x, t)\| = (x^4 + 4t^2)^{1/4}$ (cf. [7]). Then we can define the ball centered at $(0, 0)$ of radius r , that is, the set $B_r = \{(x, t) \in \mathbf{K} : \|(x, t)\| < r\}$.

Let $f \in L^1_\alpha(\mathbf{K})$. Set $x = \rho(\cos \theta)^{1/2}$, $t = 1/2\rho^2 \sin \theta$. We get

$$\int_{\mathbf{K}} f(x, t) dm_\alpha(x, t) = \frac{1}{2\pi\Gamma(\alpha + 1)} \int_{-\pi/2}^{\pi/2} \int_0^\infty f\left(\rho(\cos \theta)^{1/2}, \frac{1}{2}\rho^2 \sin \theta\right) \rho^{2\alpha+3} (\cos \theta)^\alpha d\rho d\theta. \tag{1.14}$$

If f is radial, that is, there is a function ψ on $[0, \infty)$ such that $f(x, t) = \psi(\|(x, t)\|)$, then

$$\begin{aligned} \int_{\mathbf{K}} f(x, t) dm_\alpha(x, t) &= \frac{1}{2\pi\Gamma(\alpha + 1)} \int_{-\pi/2}^{\pi/2} (\cos \theta)^\alpha d\theta \int_0^\infty \psi(\rho) \rho^{2\alpha+3} d\rho \\ &= \frac{\Gamma((\alpha + 1)/2)}{2\sqrt{\pi}\Gamma(\alpha + 1)\Gamma(\alpha/2 + 1)} \int_0^\infty \psi(\rho) \rho^{2\alpha+3} d\rho. \end{aligned} \tag{1.15}$$

Specifically,

$$m_\alpha(B_r) = \frac{\Gamma((\alpha + 1)/2)}{4\sqrt{\pi}(\alpha + 2)\Gamma(\alpha + 1)\Gamma(\alpha/2 + 1)} r^{2\alpha+4}. \tag{1.16}$$

We consider the partial differential operator

$$L = -\left(\frac{\partial^2}{\partial x^2} + \frac{2\alpha + 1}{x} \frac{\partial}{\partial x} + x^2 \frac{\partial^2}{\partial t^2}\right). \tag{1.17}$$

L is positive and symmetric in $L^2_\alpha(\mathbf{K})$, and is homogeneous of degree 2 with respect to the dilations defined above. When $\alpha = n - 1$, L is the radial part of the sublaplacian on the Heisenberg group \mathbf{H}^n . We call L the generalized sublaplacian.

Let $L_m^{(\alpha)}$ be the Laguerre polynomial of degree m and order α defined in terms of the generating function by

$$\sum_{m=0}^{\infty} s^m L_m^{(\alpha)}(x) = \frac{1}{(1-s)^{\alpha+1}} \exp\left(-\frac{xs}{1-s}\right). \quad (1.18)$$

For $(\lambda, m) \in \mathbf{R} \times \mathbf{N}$, we put

$$\varphi_{(\lambda, m)}(x, t) = \frac{m! \Gamma(\alpha + 1)}{\Gamma(m + \alpha + 1)} e^{i\lambda t} e^{-(1/2)|\lambda|x^2} L_m^{(\alpha)}(|\lambda|x^2). \quad (1.19)$$

The following proposition summarizes some basic properties of functions $\varphi_{(\lambda, m)}$.

Proposition 1.2. *The function $\varphi_{(\lambda, m)}$ satisfies that*

- (a) $\|\varphi_{(\lambda, m)}\|_{\alpha, \infty} = \varphi_{(\lambda, m)}(0, 0) = 1$,
- (b) $\varphi_{(\lambda, m)}(x, t) \varphi_{(\lambda, m)}(y, s) = T_{(x, t)}^{(\alpha)} \varphi_{(\lambda, m)}(y, s)$,
- (c) $L\varphi_{(\lambda, m)} = |\lambda|(4m + 2\alpha + 2)\varphi_{(\lambda, m)}$.

Let $f \in L^1_\alpha(\mathbf{K})$, the generalized Fourier transform of f is defined by

$$\widehat{f}(\lambda, m) = \int_{\mathbf{K}} f(x, t) \varphi_{(-\lambda, m)}(x, t) dm_\alpha(x, t). \quad (1.20)$$

It is easy to show that

$$\begin{aligned} (f * g)^\wedge(\lambda, m) &= \widehat{f}(\lambda, m) \widehat{g}(\lambda, m), \\ \widehat{f}_r(\lambda, m) &= \widehat{f}(r^2\lambda, m). \end{aligned} \quad (1.21)$$

Let $d\gamma_\alpha$ be the positive measure defined on $\mathbf{R} \times \mathbf{N}$ by

$$\int_{\mathbf{R} \times \mathbf{N}} g(\lambda, m) d\gamma_\alpha(\lambda, m) = \sum_{m=0}^{\infty} \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} \int_{\mathbf{R}} g(\lambda, m) |\lambda|^{\alpha+1} d\lambda. \quad (1.22)$$

Write $L^p_\alpha(\widehat{\mathbf{K}})$ instead of $L^p(\mathbf{R} \times \mathbf{N}, d\gamma_\alpha)$. We have the following Plancherel formula:

$$\|f\|_{\alpha, 2} = \|\widehat{f}\|_{L^2_\alpha(\widehat{\mathbf{K}})}, \quad f \in L^1_\alpha(\mathbf{K}) \cap L^2_\alpha(\mathbf{K}). \quad (1.23)$$

Then the generalized Fourier transform can be extended to the tempered distributions. We also have the inverse formula of the generalized Fourier transform.

$$f(x, t) = \int_{\mathbf{R} \times \mathbf{N}} \widehat{f}(\lambda, m) \varphi_{(\lambda, m)}(x, t) d\gamma_{\alpha}(\lambda, m) \tag{1.24}$$

provided $\widehat{f} \in L^1_{\alpha}(\widehat{\mathbf{K}})$.

In the following, we give some basic notes about the heat and Poisson kernel whose proofs can be found in [8]. Let $\{H^s\} = \{e^{-sL}\}$ be the heat semigroup generated by L . There is a unique smooth function $h((x, t), s) = h_s(x, t)$ on $\mathbf{K} \times (0, +\infty)$ such that

$$H^s f(x, t) = f * h_s(x, t). \tag{1.25}$$

We call h_s is the heat kernel associated to L . We have

$$h_s(x, t) = \int_{\mathbf{R}} \left(\frac{\lambda}{2 \sinh(2\lambda s)} \right)^{\alpha+1} e^{-(1/2)\lambda \coth(2\lambda s)x^2} e^{i\lambda t} d\lambda, \tag{1.26}$$

$$h_s(x, t) \leq C s^{-\alpha-2} e^{-(A/s)\|(x,t)\|^2}.$$

Let $\{P^s\} = \{e^{-s\sqrt{L}}\}$ be the Poisson semigroup. There is a unique smooth function $p((x, t), s) = p_s(x, t)$ on $\mathbf{K} \times (0, +\infty)$, which is called the Poisson kernel, such that

$$P^s f(x, t) = f * p_s(x, t). \tag{1.27}$$

The Poisson kernel can be calculated by the subordination. In fact, we have

$$p_s(x, t) = \frac{4s}{\sqrt{\pi}} \Gamma\left(\alpha + \frac{5}{2}\right) \int_0^{\infty} \left(\frac{\lambda}{\sinh \lambda} \right)^{\alpha+1} \left((s^2 + x^2 \lambda \coth \lambda)^2 + (2\lambda t)^2 \right)^{-(2\alpha+5)/4} \times \cos\left(\left(\alpha + \frac{5}{2}\right) \arctan\left(\frac{2\lambda t}{s^2 + x^2 \lambda \coth \lambda} \right) \right) d\lambda, \tag{1.28}$$

$$p_s(x, t) \leq C s \left(s^2 + \|(x, t)\|^2 \right)^{-(\alpha+5/2)}.$$

The heat maximal function M_H is defined by

$$M_H f(x, t) = \sup_{s>0} |H^s f(x, t)| = \sup_{s>0} |(f * h_s)(x, t)|. \tag{1.29}$$

The Poisson maximal function M_P is defined by

$$M_P f(x, t) = \sup_{s>0} |P^s f(x, t)| = \sup_{s>0} |(f * p_s)(x, t)|. \tag{1.30}$$

The Hardy-Littlewood maximal function is defined by

$$M_B f(x, t) = \sup_{r>0} \frac{1}{m_\alpha(B_r)} \int_{B_r} T_{(x,t)}^{(\alpha)}(|f|)(y, s) dm_\alpha(y, s) = \sup_{r>0} (|f| * b_r)(x, t), \quad (1.31)$$

where $b(x, t) = (1/(m_\alpha(B_1)))\chi_{B_1}(x, t)$.

The following proposition is the main result of [8].

Proposition 1.3. M_B , M_P and M_B are operators on \mathbf{K} of weak type $(1, 1)$ and strong type (p, p) for $1 < p \leq \infty$.

The paper is organized as follows. In the second section, we prove that Littlewood-Paley g -functions are bounded operators on $L^p_\alpha(\mathbf{K})$. As an application, we prove the Hörmander multiplier theorem on \mathbf{K} in the last section.

Throughout the paper, we will use C to denote the positive constant, which is not necessarily same at each occurrence.

2. Littlewood-Paley g -Function on \mathbf{K}

Let $k \in \mathbb{N}$, then we define the following G -function and g_k^* -function

$$\begin{aligned} g_k(f)^2(x, t) &= \int_0^\infty \left| \partial_s^k P^s f(x, t) \right|^2 s^{2k-1} ds, \\ g_k^*(f)^2(x, t) &= \int_0^\infty \left(\int_{\mathbf{K}} s^{-(\alpha+1)} \left(1 + s^{-2} \|(y, r)\|^4 \right)^{-k} \left| \partial_s P^s T_{(y,r)}^{(\alpha)} f(x, t) \right|^2 dm_\alpha(y, r) \right) ds. \end{aligned} \quad (2.1)$$

Then, we can prove

Theorem 2.1. (a) For $k \in \mathbb{N}$ and $f \in L^2(\mathbf{K})$, there exists $C_k > 0$ such that

$$\|g_k(f)\|_{\alpha,2} = C_k \|f\|_{\alpha,2}. \quad (2.2)$$

(b) For $1 < p < \infty$ and $f \in L^p(\mathbf{K})$, there exist positive constants C_1 and C_2 , such that

$$C_1 \|f\|_{\alpha,p} \leq \|g_k(f)\|_{\alpha,p} \leq C_2 \|f\|_{\alpha,p}. \quad (2.3)$$

(c) If $k > (\alpha + 2)/2$ and $f \in L^p(\mathbf{K})$, $p > 2$, then there exists a constant $C > 0$ such that

$$\|g_k^*(f)\|_{\alpha,p} \leq C \|f\|_{\alpha,p}. \quad (2.4)$$

Proof. (a) When $k \in \mathbb{N}$, by the Plancherel theorem for the Fourier transform on \mathbf{K} ,

$$\begin{aligned} \|g_k(f)\|_{\alpha,2}^2 &= \int_{\mathbf{K}} \left(\int_0^\infty \left| \partial_s^k P^s f(x,t) \right|^2 s^{2k-1} ds \right) dm_\alpha(x,t) \\ &= \int_0^\infty \left(\int_{\mathbf{R} \times \mathbf{N}} \left| (\partial_s^k P^s f)^\wedge(\lambda, m) \right|^2 d\gamma_\alpha(\lambda, m) \right) s^{2k-1} ds \\ &= \int_0^\infty \left(\int_{\mathbf{R}} \sum_{m=0}^\infty \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} \left| (\partial_s^k P^s f)^\wedge(\lambda, m) \right|^2 |\lambda|^{\alpha+1} d\lambda \right) s^{2k-1} ds. \end{aligned} \tag{2.5}$$

Since

$$\left(\partial_s^k P^s f \right)^\wedge(\lambda, m) = \widehat{f}(\lambda, m) \left(-\sqrt{(4m + 2\alpha + 2)|\lambda|} \right)^k e^{-s\sqrt{(4m+2\alpha+2)|\lambda|}}, \tag{2.6}$$

we get

$$\begin{aligned} \|g_k(f)\|_{\alpha,2}^2 &= \int_0^\infty \left(\int_{\mathbf{R}} \sum_{m=0}^\infty \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} \left| \widehat{f}(\lambda, m) \right|^2 \left((4m + 2\alpha + 2)|\lambda| \right)^k e^{-2s\sqrt{(4m+2\alpha+2)|\lambda|}} |\lambda|^{\alpha+1} d\lambda \right) s^{2k-1} ds. \end{aligned} \tag{2.7}$$

By

$$\int_0^\infty e^{-2s\sqrt{(4m+2\alpha+2)|\lambda|}} s^{2k-1} ds = C_k \left((4m + 2\alpha + 2)|\lambda| \right)^{-k}, \tag{2.8}$$

we have

$$\|g_k(f)\|_{\alpha,2}^2 = C_k \int_{\mathbf{R}} \sum_{m=0}^\infty \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} \left| \widehat{f}(\lambda, m) \right|^2 |\lambda|^{\alpha+1} d\lambda = C_k \|f\|_{\alpha,2}^2. \tag{2.9}$$

Therefore

$$\|g_k(f)\|_{\alpha,2} = C_k \|f\|_{\alpha,2}. \tag{2.10}$$

(b) As $\{P^s\}$ is a contraction semigroup (cf. Proposition 5.1 in [3]), we can get $\|g_k(f)\|_{\alpha,p} \leq C_2 \|f\|_{\alpha,p}$ (cf. [9]). For the reverse, we can prove by polarization to the identity and (a) (cf. [10]).

(c) We first prove

$$\int_{\mathbf{K}} g_k^*(f)^2(x,t) \psi(x,t) dm_\alpha(x,t) \leq C \int_{\mathbf{K}} g_1(f)^2(x,t) M_B \psi(x,t) dm_\alpha(x,t), \tag{2.11}$$

where $0 \leq \psi \in L^q_\alpha(\mathbf{K})$ and $\|\psi\|_{\alpha,q} \leq 1, 1/q + 2/p = 1$.

Since $k > (\alpha + 2)/2$, we know

$$\int_{\mathbf{K}} (1 + \|(y, r)\|^4)^{-k} dm_{\alpha}(y, r) < \infty. \quad (2.12)$$

By Proposition 1.3,

$$\begin{aligned} & \int_{\mathbf{K}} g_k^*(f)^2(x, t) \psi(x, t) dm_{\alpha}(x, t) \\ &= \int_{\mathbf{K}} \left(\int_0^{\infty} \int_{\mathbf{K}} s^{-(\alpha+1)} (1 + s^{-2}\|(y, r)\|^4)^{-k} \left| \partial_s P^s T_{(y,r)}^{(\alpha)} f(x, t) \right|^2 dm_{\alpha}(y, r) ds \right) \psi(x, t) dm_{\alpha}(x, t) \\ &= \int_0^{\infty} \int_{\mathbf{K}} s^{-(\alpha+1)} \left| \partial_s P^s f(y, r) \right|^2 \left(\int_{\mathbf{K}} T_{(x,t)}^{(\alpha)} (1 + s^{-2}\|(y, r)\|^4)^{-k} \psi(x, t) dm_{\alpha}(x, t) \right) dm_{\alpha}(y, r) ds \\ &\leq C \int_{\mathbf{K}} g_1(f)^2(y, r) M_B \psi(y, r) dm_{\alpha}(y, r) \\ &\leq C \|g_1(f)\|_{\alpha, p}^2 \|M_B \psi\|_{\alpha, q} \leq C \|f\|_{\alpha, p}^2. \end{aligned} \quad (2.13)$$

Therefore $\|g_k^*(f)\|_{\alpha, p} \leq C \|f\|_{\alpha, p}$. This gives the proof of Theorem 2.1. \square

We can also consider the Littlewood-Paley g -function that is defined by the heat semigroup as follows: let $k \in \mathbb{N}$, we define

$$\begin{aligned} \mathcal{G}_k^H(f)^2(x, t) &= \int_0^{\infty} \left| \partial_s^k H^s f(x, t) \right|^2 s^{2k-1} ds, \\ \mathcal{G}_k^{H,*}(f)^2(x, t) &= \int_0^{\infty} \left(\int_{\mathbf{K}} s^{-(\alpha+1)} (1 + s^{-2}\|(y, r)\|^4)^{-k} \left| \partial_s H^s T_{(y,r)}^{(\alpha)} f(x, t) \right|^2 dm_{\alpha}(y, r) \right) ds. \end{aligned} \quad (2.14)$$

Similar to the proof of Theorem 2.1, we can prove

Theorem 2.2. (a) For $k \in \mathbb{N}$ and $f \in L^2(\mathbf{K})$, there exists $C_k > 0$ such that

$$\|\mathcal{G}_k^H(f)\|_{\alpha, 2} = C_k \|f\|_{\alpha, 2}. \quad (2.15)$$

(b) For $1 < p < \infty$ and $f \in L^p(\mathbf{K})$, there exist constants C_1 and C_2 , such that

$$C_1 \|f\|_{\alpha, p} \leq \|\mathcal{G}_k^H(f)\|_{\alpha, p} \leq C_2 \|f\|_{\alpha, p}. \quad (2.16)$$

(c) If $k > (\alpha + 2)/2$ and $f \in L^p(\mathbf{K})$, $p > 2$, then $\|\mathcal{G}_k^{H,*}(f)\|_{\alpha, p} \leq C \|f\|_{\alpha, p}$.

By Theorem 2.2, we can get (cf. [10])

Corollary 2.3. *Let $k \in \mathbb{N}$ and $f \in L^2(\mathbf{K})$, if $\mathcal{G}_k^H(f) \in L^p(\mathbf{K})$, $1 < p < \infty$, then $f \in L^p(\mathbf{K})$ and there exists $C > 0$ such that*

$$C\|f\|_{\alpha,p} \leq \|\mathcal{G}_k^H(f)\|_{\alpha,p}. \tag{2.17}$$

3. Hörmander Multiplier Theorem on \mathbf{K}

In this section, we prove the Hörmander multiplier theorem on \mathbf{K} . The main tool we use is the Littlewood-Paley theory that we have proved.

We first introduce some notations. Assume Ψ is a function defined on $\mathbf{R} \times \mathbf{N}$, then let $\Delta_- \Psi(\lambda, 0) = \Psi(\lambda, 0)$ and for $m \geq 1$,

$$\begin{aligned} \Delta_- \Psi(\lambda, m) &= \Psi(\lambda, m) - \Psi(\lambda, m - 1), \\ \Delta_+ \Psi(\lambda, m) &= \Psi(\lambda, m + 1) - \Psi(\lambda, m). \end{aligned} \tag{3.1}$$

Then we define the following differential operators:

$$\begin{aligned} \Lambda_1 \Psi(\lambda, m) &= \frac{1}{|\lambda|} (m \Delta_- \Psi(\lambda, m) + (\alpha + 1) \Delta_+ \Psi(\lambda, m)), \\ \Lambda_2 \Psi(\lambda, m) &= \frac{-1}{2\lambda} ((\alpha + m + 1) \Delta_+ \Psi(\lambda, m) + m \Delta_- \Psi(\lambda, m)). \end{aligned} \tag{3.2}$$

We have the following lemma.

Lemma 3.1. *Let $g(\lambda, m) = ((4m + 2\alpha + 2)|\lambda|)e^{-(4m+2\alpha+2)|\lambda|^s}h(\lambda, m)$, where $k \in \mathbb{N}$, $h(\lambda, m)$ is a $([(\alpha + 1)/2] + 1)$ times differentiable function on \mathbb{R}^2 and satisfies*

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right)^j h(\lambda, m) \right| \leq C_j ((4m + 2\alpha + 2)|\lambda|)^{-j} \tag{3.3}$$

for $j = 0, 1, 2, \dots, [(\alpha + 1)/2] + 1$. Then one has

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right) g(\lambda, m) \right| \leq C \max \left\{ \frac{1}{|\lambda|^s}, 1 + \frac{m}{|\lambda|^s} \right\} e^{-\epsilon(4m+2\alpha+2)|\lambda|^s}, \tag{3.4}$$

where $0 < \epsilon < 1$ and $s > 0$.

Proof. Without loss of the generality, we can assume that $\lambda > 0$. when $m = 0$, we have

$$\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) = 2 \frac{\partial}{\partial \lambda}. \tag{3.5}$$

It is easy to calculate

$$\left| \frac{\partial}{\partial \lambda} g(\lambda, 0) \right| \leq C \frac{1}{\lambda s} e^{-\epsilon(4m+2\alpha+2)\lambda s}. \quad (3.6)$$

When $m \geq 1$, we have

$$\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) = 2 \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right). \quad (3.7)$$

Since

$$\begin{aligned} \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right) g(\lambda, m) &= ((4m+2\alpha+2)|\lambda|) e^{-(4m+2\alpha+2)|\lambda|s} \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right) h(\lambda, m) \\ &+ \frac{\partial}{\partial \lambda} \left\{ ((4m+2\alpha+2)|\lambda|) e^{-(4m+2\alpha+2)|\lambda|s} \right\} h(\lambda, m) \\ &- \frac{m}{\lambda} \Delta_{-1} f(m) g(m-1), \end{aligned} \quad (3.8)$$

we get

$$\left| \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right) g(\lambda, m) \right| \leq C \left(1 + \frac{m}{\lambda s} \right) e^{-\epsilon(4m+2\alpha+2)\lambda s}. \quad (3.9)$$

Then Lemma 3.1 is proved. \square

Then we can prove Hörmander multiplier theorem on the Laguerre hypergroup \mathbf{K} .

Theorem 3.2. *Let $h(\lambda, m)$ be a $([(\alpha+1)/2] + 1)$ times differentiable function on \mathbb{R}^2 and satisfies*

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right)^j h(\lambda, m) \right| \leq C_j ((4m+2\alpha+2)|\lambda|)^{-j} \quad (3.10)$$

for $j = 0, 1, 2, \dots, [(\alpha+1)/2] + 1$ and T is an operator which is defined by $\widehat{Tf}(\lambda, m) = h(\lambda, m) \widehat{f}(\lambda, m)$, then T is bounded on $L^p_\alpha(\mathbf{K})$, where $1 < p < \infty$.

Proof. We just prove the theorem for $2 < p < \infty$, for $1 < p < 2$; we can get the result by the dual theorem. By Theorem 2.2, Corollary 2.3 and the note that $Tf \in L^2(\mathbf{K})$, it is sufficient to prove the following:

$$\mathcal{G}_2^H(Tf)(x, t) \leq C \mathcal{G}_1^{H,*}(f)(x, t), \quad (x, t) \in \mathbf{K}. \quad (3.11)$$

Let $u_s = H^s f$ and $U^s = H^s(Tf)$, then we can get

$$U^{s+t} = G_t * u_s(x, t), \quad (3.12)$$

where $\widehat{G}_t(\lambda, m) = e^{-2(2m+\alpha+1)|\lambda|t} h(\lambda, m)$.

Differentiating (3.12) with respect to t and s , then assuming that $t = s$, we can get

$$\partial_s^2 H^{2s}(Tf) = F_s * \partial_s H^s f, \quad (3.13)$$

where

$$\widehat{F}_s(\lambda, m) = -((4m + 2\alpha + 2)|\lambda|) e^{-(4m+2\alpha+2)|\lambda|s} h(\lambda, m). \quad (3.14)$$

Therefore

$$\left| \partial_s^2 H^{2s}(Tf)(x, t) \right| \leq \int_{\mathbf{K}} F_s(y, r) \left| T_{(x,t)}^{(\alpha)} \partial_s H^s f(y, r) \right| dm_\alpha(y, r). \quad (3.15)$$

By the Cauchy-Schwartz inequality,

$$\left| \partial_s^2 H^{2s}(Tf)(x, t) \right|^2 \leq A(s) \int_{\mathbf{K}} \left(1 + s^{-2} \|(y, r)\|^4\right)^{-1} \left| T_{(x,t)}^{(\alpha)} \partial_s H^s f(y, r) \right|^2 dm_\alpha(y, r), \quad (3.16)$$

where

$$A(s) = \int_{\mathbf{K}} \left(1 + s^{-2} \|(x, t)\|^4\right) |F_s(x, t)|^2 dm_\alpha(x, t). \quad (3.17)$$

In the following, we prove

$$A(s) \leq C s^{-\alpha-3}. \quad (3.18)$$

We write

$$\begin{aligned} A(s) &= \int_{\|(x,t)\| \leq \sqrt{s}} \left(1 + s^{-2} \|(x, t)\|^4\right) |F_s(x, t)|^2 dm_\alpha(x, t) \\ &\quad + \int_{\|(x,t)\| > \sqrt{s}} \left(1 + s^{-2} \|(x, t)\|^4\right) |F_s(x, t)|^2 dm_\alpha(x, t) \\ &= A_1(s) + A_2(s). \end{aligned} \quad (3.19)$$

For $A_1(s)$, we can easily get

$$\begin{aligned}
 A_1(s) &\leq C \int_{\mathbf{K}} |F_s(x, t)|^2 dm_\alpha(x, t) = C \int_{\mathbf{R} \times \mathbf{N}} \left| \widehat{F}_s(\lambda, m) \right|^2 d\gamma_\alpha(\lambda, m) \\
 &= C \int_{\mathbf{R} \times \mathbf{N}} ((4m + 2\alpha + 2)|\lambda|)^2 e^{-(8m+4\alpha+4)|\lambda|s} H^2(\lambda, m) d\gamma_\alpha(\lambda, m) \\
 &\leq C \int_{\mathbf{R}} \sum_{m=0}^{\infty} \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} ((4m + 2\alpha + 2)|\lambda|)^2 e^{-(8m+4\alpha+4)|\lambda|s} |\lambda|^{\alpha+1} d\lambda \\
 &= Cs^{-\alpha-4} \int_{\mathbf{R}} \sum_{m=0}^{\infty} \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} ((4m + 2\alpha + 2)|\lambda|)^2 e^{-(8m+4\alpha+4)|\lambda|s} |\lambda|^{\alpha+1} d\lambda \\
 &\leq Cs^{-\alpha-4} \sum_{m=0}^{\infty} (4m + 2\alpha + 2)^{-2} \leq Cs^{-\alpha-4}.
 \end{aligned} \tag{3.20}$$

For $A_2(s)$, we have

$$\begin{aligned}
 A_2(s) &\leq Cs^{-2} \int_{\mathbf{K}} (4t^2 + x^4) |F_s(x, t)|^2 dm_\alpha(x, t) \\
 &= Cs^{-2} \int_{\mathbf{K}} \left| (2it - |x|^2) F_s(x, t) \right|^2 dm_\alpha(x, t) \\
 &= Cs^{-2} \int_{\mathbf{R} \times \mathbf{N}} \left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right) \widehat{F}_s(\lambda, m) \right|^2 d\gamma_\alpha(\lambda, m).
 \end{aligned} \tag{3.21}$$

By Lemma 3.1,

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right) \widehat{F}_s(\lambda, m) \right| \leq C \max \left\{ \frac{1}{|\lambda|s}, 1 + \frac{m}{|\lambda|s} \right\} e^{-\epsilon(4m+2\alpha+2)|\lambda|s}, \tag{3.22}$$

where $0 < \epsilon < 1$.

So

$$\begin{aligned}
 A_2(s) &\leq Cs^{-2} \int_{\mathbf{R} \times \mathbf{N}} e^{-\epsilon(8m+4\alpha+4)|\lambda|s} d\gamma_\alpha(\lambda, m) \\
 &= Cs^{-\alpha-4} \int_{\mathbf{R} \times \mathbf{N}} e^{-\epsilon(8m+4\alpha+4)|\lambda|s} d\gamma_\alpha(\lambda, m) \\
 &\leq Cs^{-\alpha-4}.
 \end{aligned} \tag{3.23}$$

Therefore (3.18) holds. Then

$$\left| \partial_s^2 H^{2s}(Tf)(x, t) \right|^2 \leq Cs^{-\alpha-4} \int_{\mathbf{K}} \left(1 + s^{-2} \|(y, r)\|^4 \right)^{-1} \left| T_{(x,t)}^{(\alpha)} \partial_s H^s f(y, r) \right|^2 dm_\alpha(y, r). \tag{3.24}$$

Integrating the both sides of the above inequality with $s^3 ds$, we have

$$G_2^H(x, t) \leq CG_1^{H,*}(f)(x, t). \quad (3.25)$$

Then Theorem 3.2 is proved. \square

Acknowledgments

This Papers supported by National Natural Science Foundation of China under Grant no. 11001002 and the Beijing Foundation Program under Grants no. 201010009009, no. 2010D005002000002.

References

- [1] C. E. Gutiérrez, A. Incognito, and J. L. Torrea, "Riesz transforms, g -functions, and multipliers for the Laguerre semigroup," *Houston Journal of Mathematics*, vol. 27, no. 3, pp. 579–592, 2001.
- [2] M. M. Nessibi and K. Trimèche, "Inversion of the Radon transform on the Laguerre hypergroup by using generalized wavelets," *Journal of Mathematical Analysis and Applications*, vol. 208, no. 2, pp. 337–363, 1997.
- [3] K. Stempak, "An algebra associated with the generalized sub-Laplacian," *Polska Akademia Nauk. Instytut Matematyczny. Studia Mathematica*, vol. 88, no. 3, pp. 245–256, 1988.
- [4] K. Trimèche, *Generalized Wavelets and Hypergroups*, Gordon and Breach Science Publishers, Amsterdam, The Netherlands, 1997.
- [5] R. I. Jewett, "Spaces with an abstract convolution of measures," *Advances in Mathematics*, vol. 18, no. 1, pp. 1–101, 1975.
- [6] W. R. Bloom and H. Heyer, *Harmonic Analysis of Probability Measures on Hypergroups*, vol. 20 of *de Gruyter Studies in Mathematics*, Walter de Gruyter & Co., Berlin, Germany, 1995.
- [7] K. Stempak, "Mean summability methods for Laguerre series," *Transactions of the American Mathematical Society*, vol. 322, no. 2, pp. 671–690, 1990.
- [8] J. Huang and H. Liu, "The weak type $(1, 1)$ estimates of maximal functions on the Laguerre hypergroup," *Canadian Mathematical Bulletin*, vol. 53, no. 3, pp. 491–502, 2010.
- [9] E. M. Stein, *Topics in Harmonic Analysis Related to the Littlewood-Paley Theory*, Princeton University Press, Princeton, NJ, USA, 1970.
- [10] E. M. Stein, *Singular Integrals and Differentiability Properties of Functions*, Princeton University Press, Princeton, NJ, USA, 1970.