



Study on the anomalous quartic $W^+W^-\gamma\gamma$ couplings of electroweak bosons in e^-p collisions at the LHeC and the FCC-he

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Abstract In this paper, a study is carried out on the $e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ production to probe quartic $W^+W^-\gamma\gamma$ couplings at the Large Hadron electron Collider (LHeC) with $\sqrt{s} = 1.30, 1.98$ TeV and at the Future Circular Collider-hadron electron (FCC-he) with $\sqrt{s} = 3.46, 5.29$ TeV. Production cross-sections are determined for both at leptonic and hadronic decay channels of the W -boson. With the data from future e^-p colliders, it is possible to obtain sensitivity measures at 95% C.L. on the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings which are competitive with the limits obtained by the LHC, as well as with others limits reported in the literature. The production mode $e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ in e^-p collisions offers a window for study the quartic $W^+W^-\gamma\gamma$ electroweak bosons couplings at the LHeC and the FCC-he, which provides a much cleaner collision environment than the LHC.

Keywords Models beyond the standard model · W bosons · Quartic gauge boson couplings

1 Introduction

A property of the weak interaction is that its gauge bosons W^\pm and Z can couple to each other in certain combinations and also to γ . The gauge bosons W^\pm , Z , and γ through mixing with each other represent some of the Standard Model (SM) [1–3] particles most strongly coupled to Electroweak Symmetry Breaking (EWSB). Due to the non-Abelian nature

of the SM electroweak theory, gauge bosons interact with each other and the SM predicts the existence of the Triple Gauge Couplings (TGC) and the Quartic Gauge Couplings (QGC). While the anomalous TGC (aTGC) and the anomalous QGC (aQGC) are deviations from the SM. Therefore, it is important to measure both aTGC and aQGC to further test the SM or have indications of new physics Beyond the Standard Model (BSM). Study of anomalous $WW\gamma\gamma$ couplings sensitivity is the main topic in this article. For this purpose, we use the effective Lagrangian formalism which has been utilized extensively for parameterizing new physics BSM in many processes of particle physics. This technique provides a model-independent parameterization of any new physics characterized by higher-order operators.

Studies for the aQGC $WW\gamma\gamma$ have been theoretically carried out at lepton-lepton colliders with the processes $e^+e^- \rightarrow VVV$ [4–11], $e^+e^- \rightarrow VVFF$ [12,13], $e\gamma \rightarrow VVF$ [14,15], $\gamma\gamma \rightarrow VVV$ [16,17], $\gamma\gamma \rightarrow VV$ [18], $e^+e^- \rightarrow e^+\gamma^*e^- \rightarrow VVFF$ [19] and at hadron-hadron colliders with the processes $pp \rightarrow VVV$ [20–26], $pp \rightarrow VVFF$ [27–29], $pp \rightarrow p\gamma^*p \rightarrow pVVF$ [29] and $pp \rightarrow p\gamma^*\gamma^*p \rightarrow pVVp$ [30–34], and $pp \rightarrow p\gamma^*\gamma^*p \rightarrow pVVVp$ [35], at lepton-hadron colliders with the process $ep \rightarrow VVFF$ where $V = W^\pm, Z, \gamma$ and $F = e, j, \nu$. Searches for processes containing the aQGC have been performed through the process $e^+e^- \rightarrow WW\gamma$ by the L3, DELPHI and OPAL Collaborations at the large electron positron (LEP) collider [36–39], the process $p\bar{p} \rightarrow pW^+W^-\bar{p} \rightarrow pe^+\nu_e\bar{\nu}_p$ by the D0 Collaboration at the Tevatron [40], the processes $pp \rightarrow p\gamma^*\gamma^*p \rightarrow pWWp$ and $pp \rightarrow W\gamma jj$ by the CMS Collaboration [41,42] and the process $pp \rightarrow pW^+W^-p \rightarrow pe^\pm\nu\mu^\mp\nu p$ by the ATLAS Collaboration at the LHC [43]. In the post-LHC era the present and future colliders contemplate in their physics programs the study of the aQGC: the High-

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Luminosity Large Hadron Collider (HL-LHC), the High-Energy Large Hadron Collider (HE-LHC) [44], the Large Hadron electron Collider (LHeC) [45–52], the Future Circular Collider-hadron electron (FCC-he) [51, 52], the International Linear Collider (ILC) [53], the compact linear collider (CLIC) [54], the circular electron positron collider (CEPC) [55] and the future circular collider e^+e^- (FCC-ee) [56].

The LHC may not provide high precision measurements due to strong interactions of pp collisions. An ep collider may be a very good option to complement the LHC physics program. Since ep colliders have high luminosity and high energy, the effects of new physics BSM may appear by probing the interaction of W -boson with the photon which requires measuring $WW\gamma\gamma$ couplings precisely. The LHeC and the FCC-he are planned to produce ep collisions at energies from 1.30 to 5.29 TeV [45–50]. The LHeC is a suggested deep inelastic electron-nucleon scattering machine which has been planned to collide electrons with energy from 60 GeV to possibly 140 GeV, with protons with an energy of 7 TeV. In addition, the FCC-he is designed to collide electrons with energy from 60 to 140 GeV, with protons with an energy of 50 TeV.

In this paper, we present our results in a model-independent way for the total cross-section of the process $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$ at the $e^- \gamma^*$ mode, as well as limits on the aQGC $WW\gamma\gamma$ at the LHeC with $\sqrt{s}=1.30, 1.98$ TeV and $\mathcal{L} = 10, 100 \text{ fb}^{-1}$ and at the FCC-he with $\sqrt{s} = 3.46, 5.29$ TeV and $\mathcal{L} = 100, 1000 \text{ fb}^{-1}$. For our study, we use an effective Lagrangian approach which provides a generic platform for introducing the effect of new physics BSM by adding additional terms in the Lagrangian of the SM. Specifically, we consider a scenario where the electroweak theory is realized linearly and the lowest-order aQGC are given by dimension-eight operators, and with a focus on the so-called genuine aQGC operators, that is, operators that generate the aQGC but do not have any aTGC associated with them [60].

The paper is organized as follows: In Section II, we give the general expressions for the effective Lagrangian. In Section III, we give a motivation on photon-induced process at future ep colliders. In Section IV, we evaluate the total cross-section of the reaction $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$. In Section V, we derive the 95% C.L. allowed sensitivity measures on the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings at the LHeC and the FCC-he. In Section VI, we summarize our conclusions.

2 Dimension-8 operators set relevant for the process

$$e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$$

A suitable and relatively modern approach to observe the effects of new physics BSM in a model-independent formalism is to use an effective Lagrangian description of the SM.

Starting from our present theoretical, phenomenological and experimental understanding, treating the SM in an effective Lagrangian approach is a well-motivated starting point since we have no present evidence of BSM physics. In practice, this means defining a scale, Λ , of new physics higher than the energy scale being probed in the experiment and using the fields of the SM to write higher dimension operators in addition to dimension-4 operators of the SM. Following the context of Refs. [57–61], effective field theory in which the SM is extended by higher-dimensional operators composing by all possible combinations of the SM fields is given by:

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_j \frac{c_j^{(8)}}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots, \quad (1)$$

Here, only even-dimension operators can contribute if we require lepton and baryon number conservation. For this reason, the leading effective operators which give contribution to vertices including multi-bosons are expected from dimension-6 operators. Gauge boson operators have been described by either linear or non-linear effective Lagrangians. In the nonlinear approach, the SM gauge symmetry is conserved and is realized by using the chiral Lagrangian parameterization [6, 62]. The aTGC and aQGC in this approach appear as dimension-6 operators. However, the SM gauge symmetry in the linear approach is broken by means of Higgs scalar doublet [6, 15]. Generally, dimension-6 operators used to examine the QGC provide great convenience for comparing LEP results [5]. Therefore, C and P conserving non-linear effective Lagrangian for $WW\gamma\gamma$ aQGC that define dimension-6 operators is described as follows:

$$\mathcal{L}_0 = -\frac{e^2}{16\pi\Lambda^2} a_0 F_{\mu\nu} F^{\mu\nu} \mathbf{W}^\alpha \mathbf{W}_\alpha, \quad (2)$$

$$\mathcal{L}_c = -\frac{e^2}{16\pi\Lambda^2} a_c F_{\mu\alpha} F^{\mu\beta} \mathbf{W}^\alpha \mathbf{W}_\beta, \quad (3)$$

where $F_{\mu\nu}$ and \mathbf{W}^α are defined in the usual way as in SM.

In our study in order to separate the effects of the aQGC, we shall consider effective operators that lead to the aQGC without an aTGC associated to them. The lowest dimension operators that leads to quartic interactions are of dimension-8. Therefore, genuine quartic vertices are of dimension-8. These operators are three classes of genuine aQGC operators and they are given in Table 1. This type of genuine aQGC operators in which we are interested do not have an aTGC counterpart. In the set of genuine aQGC operators given in Table 1, Φ stands for the Higgs doublet, and the covariant derivatives of the Higgs field is given by $D_\mu \Phi = (\partial_\mu + igW_\mu^j \frac{\sigma^j}{2} + \frac{i}{2} g' B_\mu) \Phi$, and σ^j ($j = 1, 2, 3$) represent the Pauli matrices, while $W^{\mu\nu}$ and $B^{\mu\nu}$ are the gauge field strength tensors for $SU(2)_L$ and $U(1)_Y$, respectively.

Table 1 Set of genuine aQGC operators of dimension-8 for $WW\gamma\gamma$ vertex [60]

Operator name	Operator
S-type operators	
$O_{S,0}$	$[(D_\mu \Phi)^\dagger (D_\nu \Phi)] \times [(D^\mu \Phi)^\dagger (D^\nu \Phi)]$
$O_{S,1}$	$[(D_\mu \Phi)^\dagger (D^\mu \Phi)] \times [(D_\nu \Phi)^\dagger (D^\nu \Phi)]$
M-type operators	
$O_{M,0}$	$Tr[W_{\mu\nu} W^{\mu\nu}] \times [(D_\beta \Phi)^\dagger (D^\beta \Phi)]$
$O_{M,1}$	$Tr[W_{\mu\nu} W^{\nu\beta}] \times [(D_\beta \Phi)^\dagger (D^\mu \Phi)]$
$O_{M,2}$	$[B_{\mu\nu} B^{\mu\nu}] \times [(D_\beta \Phi)^\dagger (D^\beta \Phi)]$
$O_{M,3}$	$[B_{\mu\nu} B^{\nu\beta}] \times [(D_\beta \Phi)^\dagger (D^\mu \Phi)]$
$O_{M,4}$	$[(D_\mu \Phi)^\dagger W_{\beta\nu} (D^\mu \Phi)] \times B^{\beta\nu}$
$O_{M,5}$	$[(D_\mu \Phi)^\dagger W_{\beta\nu} (D^\nu \Phi)] \times B^{\beta\mu}$
$O_{M,6}$	$[(D_\mu \Phi)^\dagger W_{\beta\nu} W^{\beta\nu} (D^\mu \Phi)]$
$O_{M,7}$	$[(D_\mu \Phi)^\dagger W_{\beta\nu} W^{\beta\mu} (D^\nu \Phi)]$
T-type operators	
$O_{T,0}$	$Tr[W_{\mu\nu} W^{\mu\nu}] \times Tr[W_{\alpha\beta} W^{\alpha\beta}]$
$O_{T,1}$	$Tr[W_{\alpha\nu} W^{\mu\beta}] \times Tr[W_{\mu\beta} W^{\alpha\nu}]$
$O_{T,2}$	$Tr[W_{\alpha\mu} W^{\mu\beta}] \times Tr[W_{\beta\nu} W^{\nu\alpha}]$
$O_{T,5}$	$Tr[W_{\mu\nu} W^{\mu\nu}] \times B_{\alpha\beta} B^{\alpha\beta}$
$O_{T,6}$	$Tr[W_{\alpha\nu} W^{\mu\beta}] \times B_{\mu\beta} B^{\alpha\nu}$
$O_{T,7}$	$Tr[W_{\alpha\mu} W^{\mu\beta}] \times B_{\beta\nu} B^{\nu\alpha}$
$O_{T,8}$	$B_{\mu\nu} B^{\mu\nu} B_{\alpha\beta} B^{\alpha\beta}$
$O_{T,9}$	$B_{\alpha\mu} B^{\mu\beta} B_{\beta\nu} B^{\nu\alpha}$

It is worth mentioning that the LEP2 constraints on the $WW\gamma\gamma$ vertex [63] described in terms of the anomalous a_0/Λ^2 and a_c/Λ^2 couplings can be translated into bounds on $f_{M,i}^2$ with $i = 0 - 7$. The following expressions show the relations between the $\frac{f_{M,i}}{\Lambda^4}$ couplings for the $WW\gamma\gamma$ vertex and the $\frac{a_0}{\Lambda^2}$ and $\frac{a_c}{\Lambda^2}$ couplings [5,60,64]:

$$\frac{f_{M,0}}{\Lambda^2} = \frac{a_0}{\Lambda^2} \frac{1}{g^2 v^2}, \tag{4}$$

$$\frac{f_{M,1}}{\Lambda^2} = -\frac{a_c}{\Lambda^2} \frac{1}{g^2 v^2}, \tag{5}$$

$$\frac{f_{M,0}}{\Lambda^2} = \frac{f_{M,2}}{2} = \frac{f_{M,6}}{2}, \tag{6}$$

Table 2 The aQGC altered with dimension-8 operators are shown with X

	WWWW	WWZZ	ZZZZ	WWγZ	WWγγ	ZZZγ	ZZγγ	Zγγγ	γγγγ
O_{S0}, O_{S1}	X	X	X						
$O_{M0}, O_{M1}, O_{M6}, O_{M7}$	X	X	X	X	X	X	X		
$O_{M2}, O_{M3}, O_{M4}, O_{M5}$		X	X	X	X	X	X		
O_{T0}, O_{T1}, O_{T2}	X	X	X	X	X	X	X	X	X
O_{T5}, O_{T6}, O_{T7}		X	X	X	X	X	X	X	X
O_{T8}, O_{T9}			X			X	X	X	X

$$\frac{f_{M,1}}{\Lambda^2} = \frac{f_{M,3}}{2} = -\frac{f_{M,5}}{2} = \frac{f_{M,7}}{2}. \tag{7}$$

To complement this section, in Table 2, all the aQGC altered with dimension-8 operators are presented.

3 Photoproduction at the LHeC and the FCC-he

Photon interactions have been extensively studied at HERA [70], LEP [71], Tevatron [40] and LHC [72]. In a similar manner, a significant fraction of lepton-hadron collisions at the LHeC and the FCC-he will involve quasi-real photon interactions. The LHeC and the FCC-he can to some extent be considered as a high energy $e\gamma^*$, γp , $\gamma^* p$ and $\gamma^* \gamma^*$ collisions. On this topic, the future lepton-hadron colliders offer excellent new opportunities for the study of high energy particle collisions, thus significantly extending the physics capabilities of a lepton-hadron collider. With these options, a large number of new and exciting measurements become accessible with $e\gamma^*$, γp , $\gamma^* p$ and $\gamma^* \gamma^*$ collisions. Because the photons couple directly to all fundamental fields carrying the electromagnetic currents as leptons, quarks, W 's, etc.. High energy $e\gamma^*$, γp , $\gamma^* p$ and $\gamma^* \gamma^*$ collisions will provide a comprehensive laboratory for exploring virtually every aspect of the SM and BSM physics. A review of the studies made on $e\gamma^*$, γp , $\gamma^* p$ and $\gamma^* \gamma^*$ collisions physics on future colliders is made in Refs. [57–59,73–86].

It is appropriate to mention that the studies of photon interactions at the LHC are possible due to experimental signatures of events involving photon exchanges such as the presence of very forward scattered protons and of large rapidity gaps in forward directions. However, to tag efficiently photon-induced processes and to keep backgrounds under control, some processes require very forward proton detectors [87]. The photon-induced processes have been measured in $p\bar{p}$ collisions at Tevatron-Fermilab using the large rapidity gap signature. The exclusive two-photon production of lepton pairs and the diffractive photoproduction of J/ψ mesons were studied in Refs. [88–90], respectively. In both cases, clear signals were obtained with low backgrounds.

As we mentioned above, scenarios like the LHeC and the FCC-he offer an unique opportunity to build ep collider, which can also be operated in γp collisions [73–78]. These conversions are made by converting the incoming electrons or protons into an intense beam of high energy photons. In addition, the ep colliders also provide the opportunity to examine $\gamma^*\gamma^*$, γ^*e and γ^*p modes with quasi-real photons through the equivalent photon approximation (EPA) [87,91,92].

The phenomenological investigations at lepton-hadron colliders generally contain usual deep inelastic scattering reactions where the colliding hadron dissociates into partons. These reactions have been extensively studied in the literature, while the processes elastic and semi-elastic, such as $e\gamma^*$, $\gamma^*\gamma^*$ and γ^*p have been much less studied. These processes have simpler final states with respect to lepton-hadron processes. In this case, these processes compensate for the advantages of lepton-hadron processes such as high luminosity and high center-of-mass energy. In addition, $e\gamma^*$ have effective luminosity and much higher energy compared to the process $\gamma^*\gamma^*$ collision. This may be significant because of the high energy dependencies of the cross-section containing the new physics parameters. For all the aforementioned, it is expected that the γ^*e collisions to have a high sensitivity to the aQGC $WW\gamma\gamma$.

Regarding $e\gamma^*$ collisions, these can be discerned from usual deep inelastic scattering collisions by means of two experimental signatures [93]. The first signature is the forward large rapidity gap [89,90,94,95]. Quasi-real photons have a low virtuality and scatter with small angles from the beam pipe. As the transverse momentum carried by a quasi-real photon is small, photon-emitting protons should also be scattered with small angles and exit the central detector without being detected. This causes a decreased energy deposit in the corresponding forward region. As a result, one of the forward regions of the central detector has a significant lack of energy. This defines the forward large-rapidity gap, and usual ep deep inelastic collisions can be rejected by applying a selection cut on this quantity. The second experimental signature is provided by the forward detectors [45,96,97] which are capable to detect particles with a large pseudorapidity. When a photon emitted by a proton is scattered with a large pseudorapidity, it exceeds the pseudorapidity coverage of the central detectors. In these processes, the proton can be detected by the forward detectors provides a distinctive signal for $e\gamma^*$ collisions. In this regard, the LHeC Collaboration has a program of forward physics with extra detectors located in a region between a few tens up to several hundreds of metres from the interaction point [45].

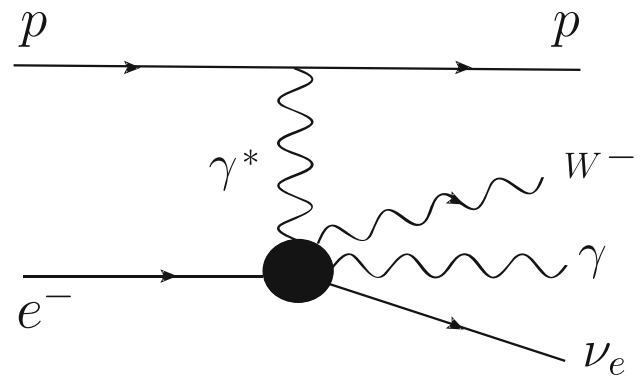


Fig. 1 Feynman diagram for the signal process $e^-p \rightarrow e^-\gamma^*p \rightarrow pW^-\gamma\nu_e$. New physics (represented by a black circle) in the electroweak sector can modify the quartic gauge couplings

4 The total cross-section for one exchanged quasi-real photon

γ^* photons emitted from proton beams collide with the incoming electron, and $e\gamma^*$ collisions are generated. The process $e^-\gamma^* \rightarrow W^-\gamma\nu_e$ participates as a subprocess in the process $e^-p \rightarrow e^-\gamma^*p \rightarrow pW^-\gamma\nu_e$. In addition, the diagram of the process $e^-p \rightarrow e^-\gamma^*p \rightarrow pW^-\gamma\nu_e$ is given in Fig. 1. The Feynman diagrams for the subprocess $e^-\gamma^* \rightarrow W^-\gamma\nu_e$ are shown in Fig. 2. Therefore, we find the total cross-section of the main process $e^-p \rightarrow e^-\gamma^*p \rightarrow pW^-\gamma\nu_e$ by integrating the cross-section for the subprocess $e^-\gamma^* \rightarrow W^-\gamma\nu_e$. The total cross-section of this process can be written as:

$$\sigma(e^-p \rightarrow pW^-\gamma\nu_e) = \int f_{\gamma^*}(x) \hat{\sigma}(e^-\gamma^* \rightarrow W^-\gamma\nu_e) dx. \tag{8}$$

Here, the spectrum of EPA photons $f_{\gamma^*}(x)$ is defined as follows [91,98]:

$$f_{\gamma^*}(x) = \frac{\alpha}{\pi E_p} \left\{ [1-x] \left[\varphi\left(\frac{Q_{max}^2}{Q_0^2}\right) - \varphi\left(\frac{Q_{min}^2}{Q_0^2}\right) \right] \right\} \tag{9}$$

with $x = E_\gamma/E_p$, $Q_{max}^2 = 2 GeV^2$ is the maximum virtuality of the photon and Q_{min}^2 is:

$$Q_{min}^2 = \frac{m_p^2 x^2}{1-x}. \tag{10}$$

In addition, the explicit form of function φ contained in Eq. (9) is:

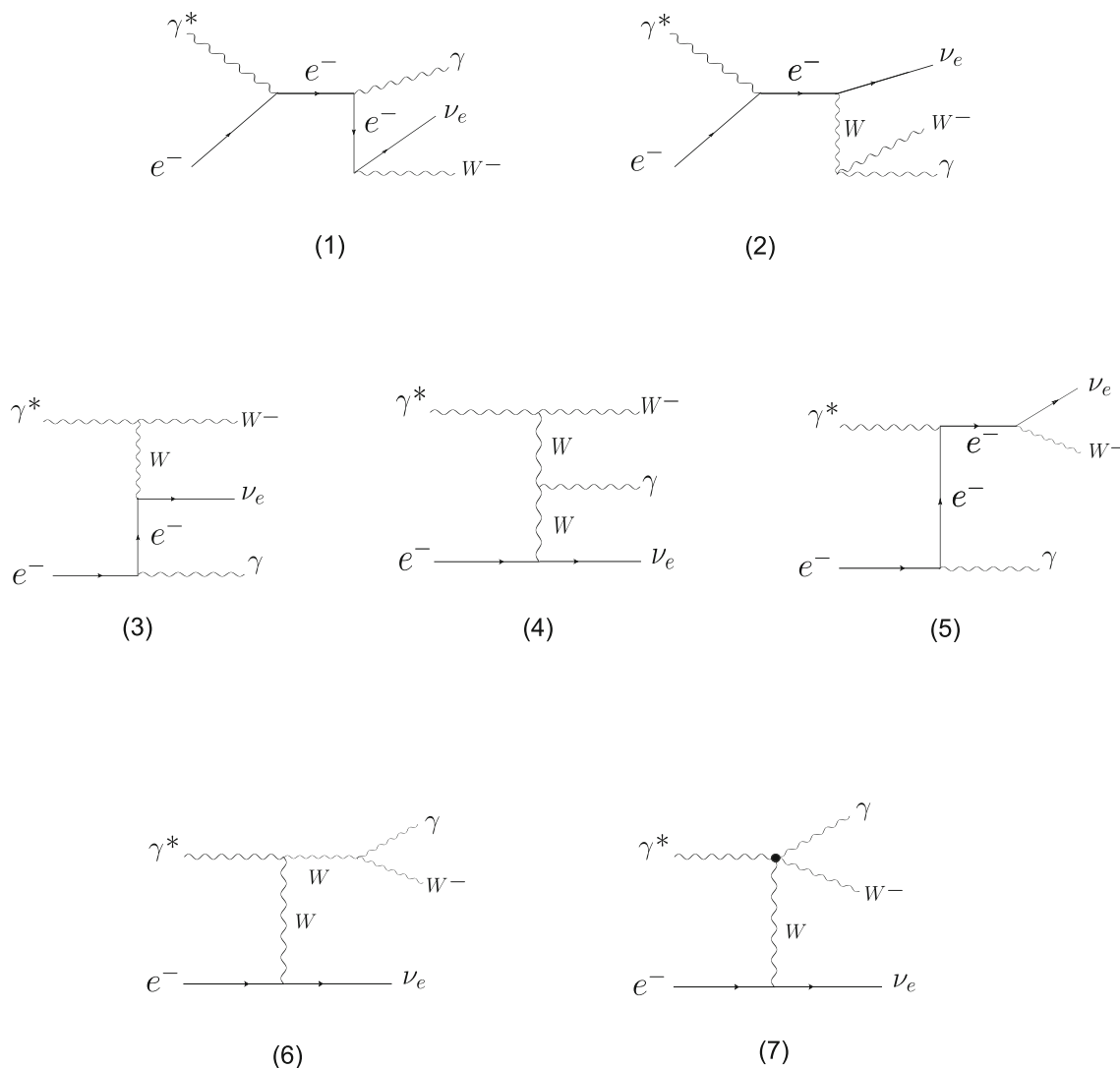


Fig. 2 Representative Feynman diagrams contributing to the subprocess $e^- \gamma^* \rightarrow W^- \gamma \nu_e$

$$\begin{aligned} \varphi(\theta) = & (1 + ay) \left[-\ln \left(1 + \frac{1}{\theta} \right) + \sum_{k=1}^3 \frac{1}{k(1 + \theta)^k} \right] + \frac{y(1 - b)}{4\theta(1 + \theta)^3} \\ & + c \left(1 + \frac{y}{4} \right) \left[\ln \left(\frac{1 - b + \theta}{1 + \theta} \right) + \sum_{k=1}^3 \frac{b^k}{k(1 + \theta)^k} \right], \end{aligned} \tag{11}$$

where explicitly a, b, c and y are given by:

$$y = \frac{x^2}{(1 - x)}, \tag{12}$$

$$a = \frac{1 + \mu_p^2}{4} + \frac{4m_p^2}{Q_0^2} \approx 7.16, \tag{13}$$

$$b = 1 - \frac{4m_p^2}{Q_0^2} \approx -3.96, \tag{14}$$

$$c = \frac{\mu_p^2 - 1}{b^4} \approx 0.028. \tag{15}$$

In our calculations, we analyze signals and backgrounds of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ through the expression given by Eq. (8) and using the MadGraph5_aMC@NLO [99] package in which the aQGC are implemented through FeynRules [100] package through dimension-8 effective Lagrangians related to the anomalous quartic $WW\gamma\gamma$ couplings.

In order to make our numerical computation more realistic, the kinematic study of the $W^- \gamma \nu_e$ production starts with the usual detection and isolation cuts on the final state leptons and quarks. The SM process with final state should be accepted as a background for the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$. Additionally, we have considered the following background processes $\nu_e \gamma l \nu_l, \nu_e \gamma jj, \nu_e \gamma \gamma jj$ and $\nu_e \gamma jjj$ for both leptonic and hadronic decay channel of the W-boson. We

Table 3 Definition of the fiducial regions of the fully leptonic and hadronic $W^- \gamma \nu_e$ analyses

Fiducial Requirements		
Selected cuts of $f_{M,i}$ for the hadronic decay of the W -boson		
$p_T^j > 30 \text{ GeV}$,	$p_T^\gamma > 150 \text{ GeV}$,	$p_T^\nu > 20 \text{ GeV}$
$100 \text{ GeV} > M_{jj} > 60 \text{ GeV}$,	$ \eta_j < 5$,	$ \eta_\gamma < 2.5$,
$\Delta R_{jj(\min)} = 0.4$		
$\Delta R_{\gamma j(\min)} = 0.4$		
Selected cuts of $f_{T,i}$ for the hadronic decay of the W -boson		
$p_T^j > 30 \text{ GeV}$,	$p_T^\gamma > 170 \text{ GeV}$,	$p_T^\nu > 20 \text{ GeV}$
$100 \text{ GeV} > M_{jj} > 60 \text{ GeV}$,	$ \eta_j < 5$,	$ \eta_\gamma < 2.5$,
$\Delta R_{jj(\min)} = 0.4$		
$\Delta R_{\gamma j(\min)} = 0.4$		
Selected cuts of $f_{M,i}$ for the leptonic decay of the W -boson		
$p_T^l > 20 \text{ GeV}$,	$p_T^\gamma > 150 \text{ GeV}$,	$p_T^\nu > 20 \text{ GeV}$,
$ \eta_l < 2.5$,	$ \eta_\gamma < 2.5$,	
$\Delta R_{ll(\min)} = 0.4$		
$\Delta R_{\gamma l(\min)} = 0.4$		
Selected cuts of $f_{T,i}$ for the leptonic decay of the W -boson		
$p_T^l > 20 \text{ GeV}$,	$p_T^\gamma > 170 \text{ GeV}$,	$p_T^\nu > 20 \text{ GeV}$,
$ \eta_l < 2.5$,	$ \eta_\gamma < 2.5$,	
$\Delta R_{ll(\min)} = 0.4$		
$\Delta R_{\gamma l(\min)} = 0.4$		

know that the high dimensional operators could affect the p_T^γ photon transverse momentum, especially at the region with a large p_T^γ values, which can be very useful to distinguish signal and background events (see Section V). By applying a cut in the missing energy, we reduce the background to consider. Therefore, a set of kinematic cuts used for the analysis of signal and background processes are summarized in Table 3.

In Tables 4 and 5, we present the values of σ_{SM} and σ_{Tot} of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ for the center-of-mass energies of 1.30, 1.98, 3.46 and 5.29 TeV. Here, σ_{Tot} is given as follows,

$$\sigma_{Tot} = \sigma_{SM} + \sigma_{int} + \sigma_{BSM}, \tag{16}$$

where σ_{SM} is the SM cross-section, σ_{int} is the interference term between SM and the new physics contribution and σ_{BSM} is the contribution due to BSM physics, respectively. Values of the total cross-section predictions for the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ after applying the cuts described in the text, show that the total cross-section increases with increasing of the center-of-mass energy, as well as with the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ contribution. For instance, from Tables 4 and 5, we obtain $\sigma(\sqrt{s}, \frac{f_{T,6}}{\Lambda^4}) = (3.18 \times 10^7)\sigma_{SM}$ and $\sigma(\sqrt{s}, \frac{f_{T,5}}{\Lambda^4}) = (2.66 \times 10^7)\sigma_{SM}$ for the leptonic and hadronic channels, with $\sqrt{s} = 5.29 \text{ TeV}$ at the FCC-he. In general, the total cross-section predicted for the $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ signal and for the center-of-mass

energies of $\sqrt{s} = 1.30, 1.98, 3.46, 5.29 \text{ TeV}$ is increased by $\mathcal{O}(10^1 - 10^7)$ orders of magnitude with respect to the total cross-section of the SM, that is to say $\sigma(\sqrt{s}, \frac{f_{M,i}}{\Lambda^4}) \approx \mathcal{O}(10^1 - 10^5)\sigma_{SM}$ and $\sigma(\sqrt{s}, \frac{f_{T,i}}{\Lambda^4}) \approx \mathcal{O}(10^4 - 10^7)\sigma_{SM}$, with $i = 1 - 5, 7$ and $j = 0 - 2, 5, 6, 7$.

For the aQGC $f_{M,0-5,7}/\Lambda^4$ and $f_{T,0-2,5,6,7}/\Lambda^4$ taking one at a time, we get the results for the total cross-section (σ_{Tot}) as shown in Figs. 3, 4, 5, 6, 7, 8, 9 and 10 and Tables 4 and 5 at the LHeC with $\sqrt{s} = 1.30, 1.98 \text{ TeV}$ and the FCC-he with $\sqrt{s} = 3.46, 5.29 \text{ TeV}$, respectively. The color lines in Figs. 3, 4, 5, 6, 7, 8, 9 and 10 show the deviations from the total SM value of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ as a function of $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$. In these figures, we consider the leptonic and hadronic decays of the W -boson in the final state of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$, where $W \rightarrow \nu_l l$, $W \rightarrow jj'$ with $\nu_l = \nu_e, \nu_\mu$, $l = e^-, \mu$ and $j = u, c, \bar{d}, \bar{s}$, $j' = d, s, \bar{u}, \bar{c}$. Figures 8 and 10 illustrate more clear effect of the dimension-8 operators on the total cross-section of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ with the leptonic and hadronic decay of the W -boson, and for the FCC-he with $\sqrt{s} = 5.29 \text{ TeV}$. The highest cross-section in value is obtained for $\sigma(\sqrt{s}, f_{T,5}/\Lambda^4) = 2.10 \times 10^5 \text{ pb}$ followed by $\sigma(\sqrt{s}, f_{T,6}/\Lambda^4) = 1.05 \times 10^5 \text{ pb}$ and $\sigma(\sqrt{s}, f_{T,7}/\Lambda^4) = 2.54 \times 10^4 \text{ pb}$ for the hadronic channel as shown in Table 5 as well as by Fig. 10. A direct comparison of the results shown in Tables 4 and 5 for the total cross-section of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ projected by

Table 4 Total cross-sections of the process $e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ at the LHeC with $\sqrt{s} = 1.30, 1.98$ TeV and at the FCC-he with $\sqrt{s} = 3.46, 5.29$ depending on 13 anomalous couplings obtained by dimension-8 operators. Also, all anomalous couplings for the LHeC and the FCC-he are taken as equal to 1×10^{-8} and 5×10^{-9} GeV⁻⁴, respectively. Here, we consider that only one of the anomalous couplings deviates from the SM at any given time

	LHeC		FCC-he	
	Leptonic decay		Leptonic decay	
$\sigma(e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e)$ (pb)				
SM	7.55×10^{-6}	4.04×10^{-5}	3.93×10^{-5}	1.63×10^{-4}
Couplings	$\sqrt{s} = 1.30$ TeV	$\sqrt{s} = 1.98$ TeV	$\sqrt{s} = 3.46$ TeV	$\sqrt{s} = 5.29$ TeV
f_{M0}/Λ^4	2.99×10^{-3}	9.35×10^{-2}	1.09×10^{-2}	7.77×10^{-1}
f_{M1}/Λ^4	2.04×10^{-3}	5.70×10^{-2}	7.43×10^{-3}	6.00×10^{-1}
f_{M2}/Λ^4	1.29×10^{-1}	4.02	4.73×10^{-1}	3.25×10^1
f_{M3}/Λ^4	8.81×10^{-2}	2.45	3.23×10^{-1}	2.58×10^1
f_{M4}/Λ^4	9.77×10^{-3}	3.06×10^{-1}	3.56×10^{-2}	2.54
f_{M5}/Λ^4	6.77×10^{-3}	1.87×10^{-1}	2.52×10^{-2}	1.96
f_{M7}/Λ^4	5.25×10^{-4}	1.43×10^{-2}	1.96×10^{-3}	1.46×10^{-1}
f_{T0}/Λ^4	2.85×10^{-1}	10.00	3.03	2.14×10^2
f_{T1}/Λ^4	7.15×10^{-1}	2.34×10^1	5.79	4.82×10^2
f_{T2}/Λ^4	9.60×10^{-2}	3.16	8.18×10^{-1}	7.75×10^1
f_{T5}/Λ^4	3.07	1.08×10^2	3.26×10^1	2.26×10^3
f_{T6}/Λ^4	7.68	2.52×10^2	6.25×10^1	5.19×10^3
f_{T7}/Λ^4	1.04	3.40×10^1	8.79	7.63×10^2

Table 5 Total cross-sections of the process $e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ at the LHeC with $\sqrt{s} = 1.30, 1.98$ TeV and at the FCC-he with $\sqrt{s} = 3.46, 5.29$ TeV depending on 13 anomalous couplings obtained by dimension-8 operators. Also, all anomalous couplings for the LHeC and the FCC-he are taken as equal to 1×10^{-8} and 5×10^{-9} GeV⁻⁴, respectively. Here, we consider that only one of the anomalous couplings deviates from the SM at any given time

	LHeC		FCC-he	
	Hadronic decay		Hadronic decay	
$\sigma(e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e)$ (pb)				
SM	1.98×10^{-5}	8.95×10^{-5}	2.37×10^{-4}	6.38×10^{-4}
Couplings	$\sqrt{s} = 1.30$ TeV	$\sqrt{s} = 1.98$ TeV	$\sqrt{s} = 3.46$ TeV	$\sqrt{s} = 5.29$ TeV
f_{M0}/Λ^4	1.85×10^{-2}	3.17×10^{-1}	1.06	8.33
f_{M1}/Λ^4	5.66×10^{-3}	1.03×10^{-1}	9.16×10^{-2}	8.04×10^{-1}
f_{M2}/Λ^4	7.94×10^{-1}	1.36×10^1	4.50×10^1	3.58×10^2
f_{M3}/Λ^4	2.45×10^{-1}	4.43	3.90	3.48×10^1
f_{M4}/Λ^4	6.05×10^{-2}	1.04	3.43	2.90×10^1
f_{M5}/Λ^4	1.88×10^{-2}	3.39×10^{-1}	2.97×10^{-1}	2.63
f_{M7}/Λ^4	1.46×10^{-3}	2.58×10^{-2}	2.33×10^{-2}	2.02×10^{-1}
f_{T0}/Λ^4	1.30	2.89×10^1	1.48×10^2	1.55×10^3
f_{T1}/Λ^4	1.93	4.19×10^1	7.10×10^1	7.47×10^2
f_{T2}/Λ^4	2.79×10^{-1}	6.03	1.77×10^1	1.87×10^2
f_{T5}/Λ^4	1.40×10^1	3.10×10^2	1.59×10^3	1.70×10^4
f_{T6}/Λ^4	2.08×10^1	4.51×10^2	7.61×10^2	8.07×10^3
f_{T7}/Λ^4	3.01	6.52×10^1	1.94×10^2	2.02×10^3

the LHeC and the FCC-he for both leptonic and hadronic channels of the W-boson, indicate a difference of one and up to two orders of magnitude of the FCC-he with respect to the LHeC. Similar behavior is observed in Figs. 3, 4, 5, 6, 7, 8, 9 and 10.

To close this section, it is worth mentioning that our results show that a nonzero aQGC enhances the production cross-section at large energies of the e^-p system with respect to the SM prediction as can be seen in Figs. 3, 4, 5, 6, 7, 8, 9 and 10.

5 Projections on the aQGC $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ at the LHeC and the FCC-he

The presence of new physics characterized by the parameters $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ may be quantified by a simple χ^2 method that varies the parameters and is based on:

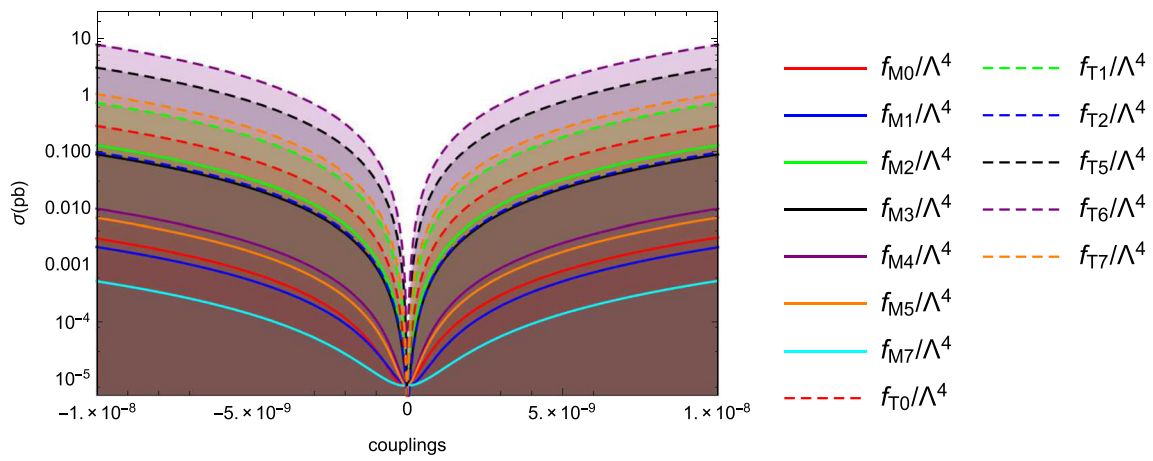


Fig. 3 For leptonic channel, the total cross-sections of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ as a function of the anomalous couplings at the LHeC with $\sqrt{s} = 1.30$ TeV

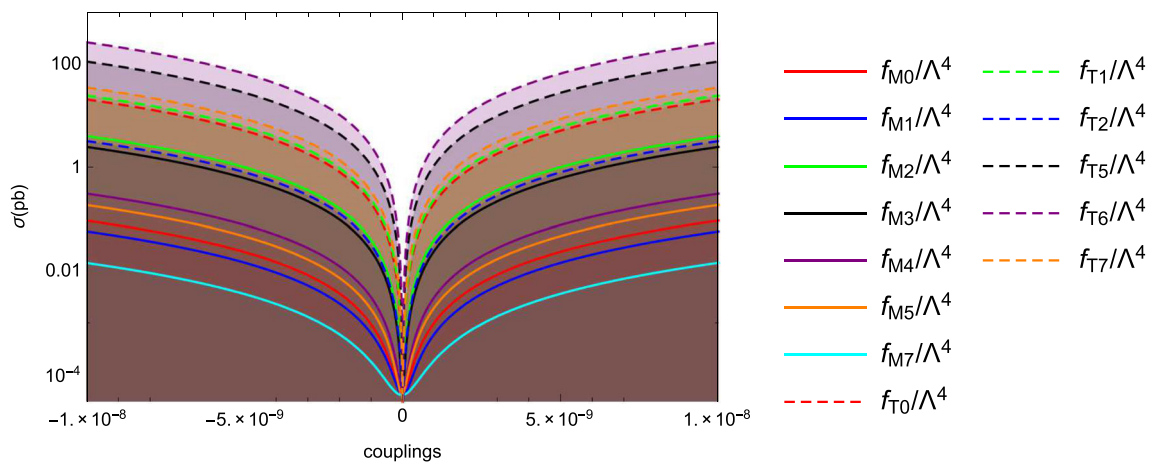


Fig. 4 Same as in Fig. 3, but for $\sqrt{s} = 1.98$ TeV at the LHeC

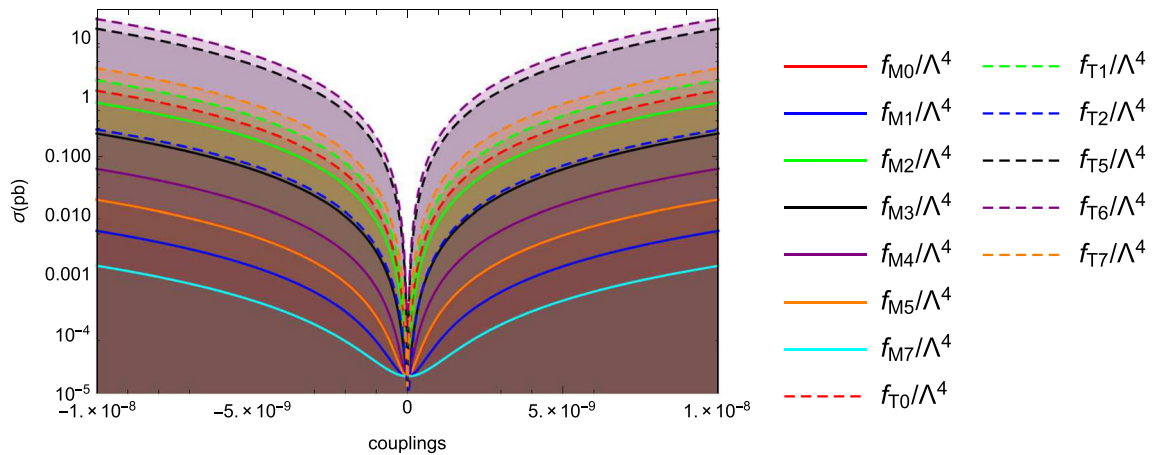


Fig. 5 Same as in Fig. 3, but for hadronic decay

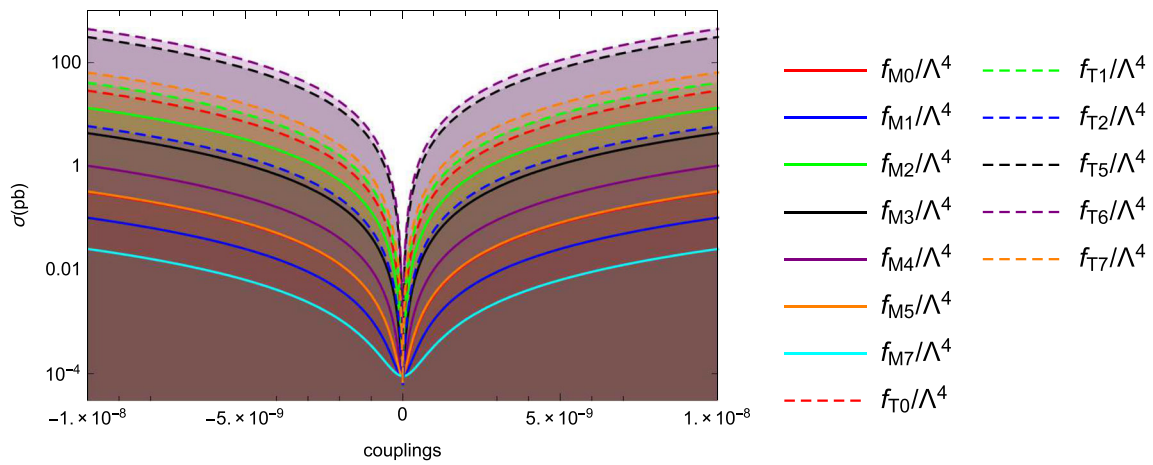


Fig. 6 Same as in Fig. 4, but for hadronic decay

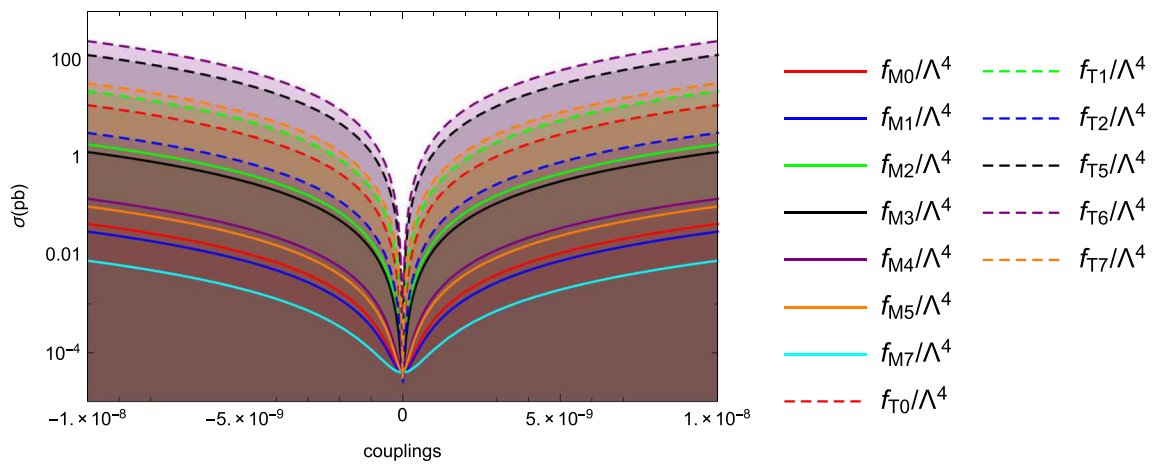


Fig. 7 For leptonic channel, the total cross-sections of the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ as a function of the anomalous couplings at the FCC-he with $\sqrt{s} = 3.46$ TeV

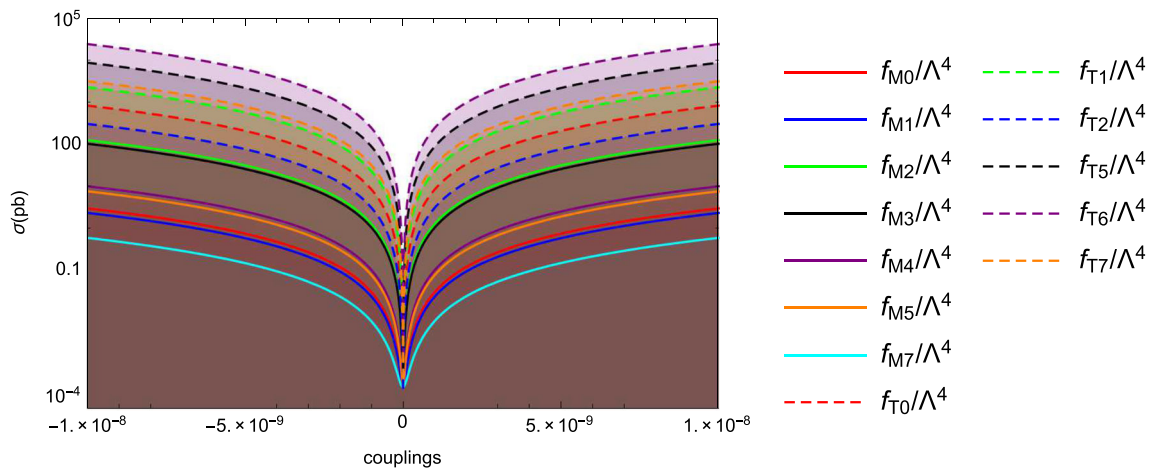


Fig. 8 Same as in Fig. 7, but for $\sqrt{s} = 5.29$ TeV at the FCC-he

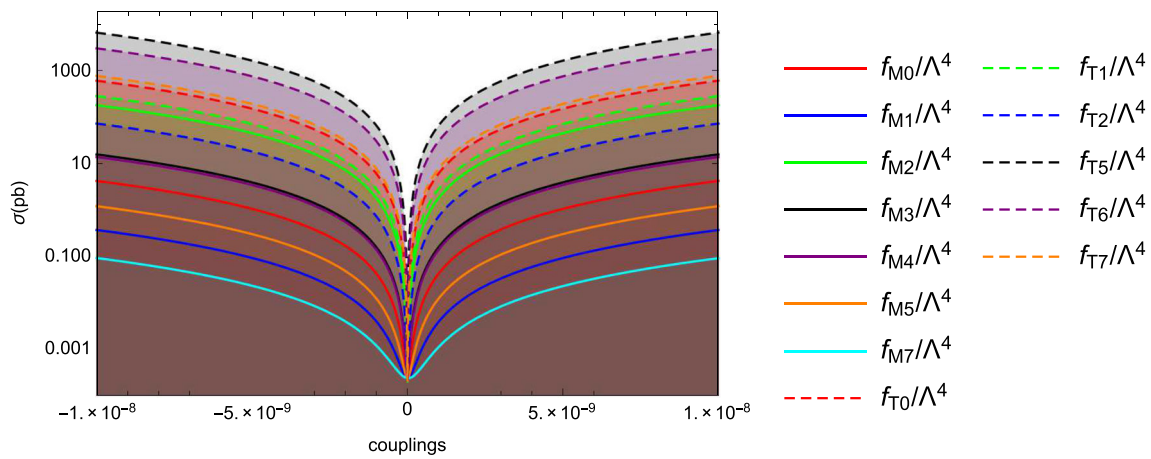


Fig. 9 Same as in Fig. 7, but for hadronic decay

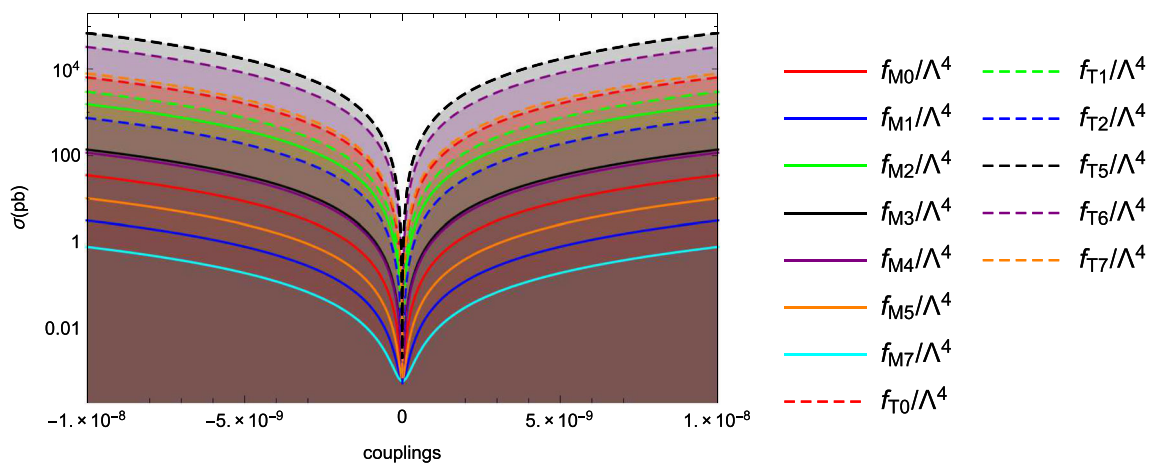


Fig. 10 Same as in Fig. 8, but for hadronic decay

$$\chi^2(f_{M,i}/\Lambda^4, f_{T,j}/\Lambda^4) = \left(\frac{\sigma_{SM}(\sqrt{s}) - \sigma_{Tot}(\sqrt{s}, f_{M,i}/\Lambda^4, f_{T,j}/\Lambda^4)}{\sigma_{SM}(\sqrt{s})\delta_{st}} \right)^2. \quad (17)$$

where $\delta_{st} = \frac{1}{\sqrt{N_{SM}}}$ is the statistical error and N_{SM} is the total number of events only coming from SM backgrounds:

$$N_{SM} = \mathcal{L} \times \sigma_{SM}. \quad (18)$$

In order to quantify the expected limits on $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$, advantage has been taken in this analysis of the fact that the aQGC enhance the total cross-section at high energies (see Tables 4 and 5). To get an idea of the LHeC and FCC-he constraining power, in Tables 6 and 7, we show the expected bounds on the aQGC $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ from the $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$ production. We present in the rows of Tables 6 and 7 the expected LHeC and FCC-he limits. In these tables, attainable sensitivity on $f_{M,0-5,7}/\Lambda^4$ and $f_{T,0-2,5-7}/\Lambda^4$ at the LHeC and the FCC-he runs is already higher than the present direct limits stemming from LEP [36–

39] and Tevatron [40] and our limits are competitive with the limits reported by the LHC [41–43]. At the FCC-he, these aQGC can be tested in the $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$ production mode with $\sqrt{s} = 5.29$ TeV and $\mathcal{L} = 1000$ fb $^{-1}$. Our limits stronger on the Wilson coefficients of the operators $\mathcal{O}_{T,j}$ are listed below for the FCC-he with $\sqrt{s} = 5.29$ TeV and $\mathcal{L} = 1000$ fb $^{-1}$ at 95% C.L., for the hadronic decay channel of the W-boson in the final state:

$$\frac{f_{T5}}{\Lambda^4} = [-0.237; 0.270] \text{ TeV}^{-4}, \quad (19)$$

$$\frac{f_{T6}}{\Lambda^4} = [-0.330; 0.420] \text{ TeV}^{-4}, \quad (20)$$

$$\frac{f_{T7}}{\Lambda^4} = [-1.000; 0.550] \text{ TeV}^{-4}. \quad (21)$$

For the operators $\mathcal{O}_{M,i}$, the constraints on the Wilson coefficients do not degrade substantially and become in a factor $\mathcal{O}(2 - 5)$ weaker than those obtained in Refs. [103, 107] through the $Z\gamma jj$ production.

Table 6 Sensitivity measures on the aQGC at the 95% C. L. via the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ at the LHeC with $\sqrt{s} = 1.30, 1.98$ TeV. Here, we consider that only one of the anomalous couplings deviates from the SM at any given time

Couplings (TeV ⁻⁴)	Leptonic channel		Hadronic channel	
	10 fb ⁻¹	100 fb ⁻¹	10 fb ⁻¹	100 fb ⁻¹
LHeC, $\sqrt{s} = 1.30$ TeV				
f_{M0}/Λ^4	$[-1.45; 1.25] \times 10^3$	$[-0.86; 0.66] \times 10^3$	$[-6.75; 6.92] \times 10^2$	$[-3.76; 3.93] \times 10^2$
f_{M1}/Λ^4	$[-2.02; 2.10] \times 10^3$	$[-0.88; 0.94] \times 10^3$	$[-1.24; 1.26] \times 10^3$	$[-6.96; 7.10] \times 10^2$
f_{M2}/Λ^4	$[-2.02; 2.10] \times 10^2$	$[-1.12; 1.20] \times 10^2$	$[-1.02; 1.06] \times 10^2$	$[-5.62; 6.09] \times 10^1$
f_{M3}/Λ^4	$[-2.39; 2.55] \times 10^2$	$[-1.31; 1.47] \times 10^2$	$[-1.86; 1.93] \times 10^2$	$[-1.03; 1.10] \times 10^2$
f_{M4}/Λ^4	$[-7.38; 7.51] \times 10^2$	$[-4.13; 4.25] \times 10^2$	$[-3.91; 3.62] \times 10^2$	$[-2.27; 1.97] \times 10^2$
f_{M5}/Λ^4	$[-9.07; 8.73] \times 10^2$	$[-5.17; 4.84] \times 10^2$	$[-6.98; 6.73] \times 10^2$	$[-3.98; 3.73] \times 10^2$
f_{M7}/Λ^4	$[-3.27; 3.19] \times 10^3$	$[-1.85; 1.77] \times 10^3$	$[-2.46; 2.49] \times 10^3$	$[-1.38; 1.41] \times 10^3$
f_{T0}/Λ^4	$[-1.26; 1.20] \times 10^2$	$[-0.73; 0.66] \times 10^2$	$[-7.46; 7.37] \times 10^1$	$[-4.21; 4.12] \times 10^1$
f_{T1}/Λ^4	$[-7.70; 7.78] \times 10^1$	$[-4.31; 4.39] \times 10^1$	$[-6.19; 5.90] \times 10^1$	$[-3.55; 3.26] \times 10^1$
f_{T2}/Λ^4	$[-1.92; 2.32] \times 10^2$	$[-1.01; 1.41] \times 10^2$	$[-1.34; 1.86] \times 10^2$	$[-0.67; 1.18] \times 10^2$
f_{T5}/Λ^4	$[-3.65; 3.92] \times 10^1$	$[-2.00; 2.27] \times 10^1$	$[-2.10; 2.36] \times 10^1$	$[-1.13; 1.39] \times 10^1$
f_{T6}/Λ^4	$[-2.15; 2.61] \times 10^1$	$[-1.12; 1.59] \times 10^1$	$[-1.67; 2.03] \times 10^1$	$[-0.87; 1.23] \times 10^1$
f_{T7}/Λ^4	$[-0.71; 0.58] \times 10^2$	$[-0.43; 0.30] \times 10^2$	$[-4.87; 4.73] \times 10^1$	$[-2.77; 2.63] \times 10^1$
LHeC, $\sqrt{s} = 1.98$ TeV				
f_{M0}/Λ^4	$[-3.84; 3.51] \times 10^2$	$[-2.24; 1.91] \times 10^2$	$[-2.40; 2.45] \times 10^2$	$[-1.34; 1.39] \times 10^2$
f_{M1}/Λ^4	$[-4.57; 4.77] \times 10^2$	$[-2.53; 2.73] \times 10^2$	$[-4.28; 4.30] \times 10^2$	$[-2.40; 2.42] \times 10^2$
f_{M2}/Λ^4	$[-5.42; 5.84] \times 10^1$	$[-2.96; 3.38] \times 10^1$	$[-3.67; 3.80] \times 10^1$	$[-2.04; 2.16] \times 10^1$
f_{M3}/Λ^4	$[-7.10; 7.24] \times 10^1$	$[-3.96; 4.10] \times 10^1$	$[-6.42; 6.69] \times 10^1$	$[-3.55; 3.83] \times 10^1$
f_{M4}/Λ^4	$[-2.01; 2.03] \times 10^2$	$[-1.13; 1.15] \times 10^2$	$[-1.37; 1.33] \times 10^2$	$[-7.78; 7.41] \times 10^1$
f_{M5}/Λ^4	$[-2.77; 2.43] \times 10^2$	$[-1.64; 1.30] \times 10^2$	$[-2.51; 2.28] \times 10^2$	$[-1.46; 1.24] \times 10^2$
f_{M7}/Λ^4	$[-9.56; 9.18] \times 10^2$	$[-5.46; 5.08] \times 10^2$	$[-8.48; 8.67] \times 10^2$	$[-4.73; 4.92] \times 10^2$
f_{T0}/Λ^4	$[-2.22; 2.35] \times 10^1$	$[-1.22; 1.36] \times 10^1$	$[-2.42; 2.26] \times 10^1$	$[-1.40; 1.24] \times 10^1$
f_{T1}/Λ^4	$[-2.02; 2.18] \times 10^1$	$[-1.10; 1.26] \times 10^1$	$[-1.99; 1.94] \times 10^1$	$[-1.13; 1.08] \times 10^1$
f_{T2}/Λ^4	$[-5.46; 6.07] \times 10^1$	$[-2.95; 3.56] \times 10^1$	$[-0.45; 0.59] \times 10^2$	$[-0.23; 0.37] \times 10^2$
f_{T5}/Λ^4	$[-9.64; 10.09]$	$[-5.33; 5.78]$	$[-7.12; 7.13]$	$[-4.00; 4.01]$
f_{T6}/Λ^4	$[-6.22; 6.64]$	$[-3.41; 3.83]$	$[-6.07; 5.83]$	$[-3.46; 3.23]$
f_{T7}/Λ^4	$[-1.95; 1.59] \times 10^1$	$[-1.18; 0.83] \times 10^1$	$[-1.55; 1.58] \times 10^1$	$[-8.67; 8.94]$

To complement our results obtained on the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings which are summarized in Tables 6 and 7, we calculated the sensitivity on the aQGCs at the 95% C. L. via the process $e^- p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ at the FCC-he with $\sqrt{s} = 3.46, 5.29$ TeV with the combined data for the leptonic and hadronic decays of W -boson given in Table 3. Our results are illustrated in Table 8. In this case the sensitivities on the aQGCs for the combined data get better with respect to the leptonic and hadronic decays of W -boson given by Tables 6 and 7.

In the post-LHC era, the FCC-he is one of the proposed colliders in the new energy frontier at the LHC and would provide proton beam energies up to 50 TeV and electron beam energies from 60 to 140 GeV. In this case the expected bounds on the $f_{M,0-5,7}/\Lambda^4$ and $f_{T,0-2,5-7}/\Lambda^4$ Wilson coef-

ficients can reach a sensitivity of approximately one order of magnitude stronger than our present limits given in Tables 6, 7 and 8.

In Ref. [107], a study phenomenological through the two-to-two scattering of electroweak gauge bosons is carried out to determined the partial-wave unitarity constraints on the lowest-dimension effective operators which generate aQGC. Quantitatively, its results are summarized in Tables 1, 2, 3 and 4 of Ref. [107]. Our results on the $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ Wilson coefficients show that our limits are competitive with the results reported in Ref. [107], and in some cases our limits are stronger by one order of magnitude as shown in Tables 7 and 8 of our manuscript, as well as in Tables 2 and 3 of Ref. [107].

Table 7 Sensitivity measures on the aQGC at the 95% C. L. via the process $e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ at the FCC-he with $\sqrt{s} = 3.46, 5.29$ TeV. Here, we consider that only one of the anomalous couplings deviates from the SM at any given time

Couplings (TeV ⁻⁴)	Leptonic channel		Hadronic channel	
	100 fb ⁻¹	1000 fb ⁻¹	100 fb ⁻¹	1000 fb ⁻¹
FCC-he, $\sqrt{s} = 3.46$ TeV				
f_{M0}/Λ^4	$[-2.98; 2.97] \times 10^2$	$[-1.68; 1.67] \times 10^2$	$[-4.69; 4.88] \times 10^1$	$[-2.60; 2.78] \times 10^1$
f_{M1}/Λ^4	$[-3.56; 3.60] \times 10^2$	$[-2.00; 2.03] \times 10^2$	$[-1.56; 1.67] \times 10^2$	$[-0.86; 0.97] \times 10^2$
f_{M2}/Λ^4	$[-4.54; 4.58] \times 10^1$	$[-2.54; 2.58] \times 10^1$	$[-7.10; 7.55]$	$[-3.90; 4.35]$
f_{M3}/Λ^4	$[-5.35; 5.57] \times 10^1$	$[-2.96; 3.18] \times 10^1$	$[-2.57; 2.40] \times 10^1$	$[-1.48; 1.31] \times 10^1$
f_{M4}/Λ^4	$[-1.60; 1.69] \times 10^2$	$[-0.88; 0.97] \times 10^2$	$[-2.69; 2.61] \times 10^1$	$[-1.53; 1.45] \times 10^1$
f_{M5}/Λ^4	$[-2.05; 1.91] \times 10^2$	$[-1.18; 1.04] \times 10^2$	$[-9.03; 8.86] \times 10^1$	$[-5.11; 4.95] \times 10^1$
f_{M7}/Λ^4	$[-7.13; 7.12] \times 10^2$	$[-4.01; 4.00] \times 10^2$	$[-3.03; 3.48] \times 10^2$	$[-1.61; 2.07] \times 10^2$
f_{T0}/Λ^4	$[-1.66; 1.65] \times 10^1$	$[-9.34; 9.23]$	$[-3.40; 4.03]$	$[-1.79; 2.42]$
f_{T1}/Λ^4	$[-1.05; 1.34] \times 10^1$	$[-0.54; 0.83] \times 10^1$	$[-0.61; 0.48] \times 10^1$	$[-0.38; 0.25] \times 10^1$
f_{T2}/Λ^4	$[-0.27; 0.38] \times 10^2$	$[-0.13; 0.24] \times 10^2$	$[-0.91; 1.27] \times 10^1$	$[-0.45; 0.81] \times 10^1$
f_{T5}/Λ^4	$[-5.02; 5.09]$	$[-2.81; 2.88]$	$[-1.01; 1.26]$	$[-0.52; 0.77]$
f_{T6}/Λ^4	$[-3.57; 3.70]$	$[-1.98; 2.11]$	$[-1.37; 2.01]$	$[-0.67; 1.31]$
f_{T7}/Λ^4	$[-1.05; 0.92] \times 10^1$	$[-0.62; 0.49] \times 10^1$	$[-3.36; 3.21]$	$[-1.92; 1.78]$
FCC-he, $\sqrt{s} = 5.29$ TeV				
f_{M0}/Λ^4	$[-5.27; 5.14] \times 10^1$	$[-2.99; 2.86] \times 10^1$	$[-2.08; 2.13] \times 10^1$	$[-1.16; 1.21] \times 10^1$
f_{M1}/Λ^4	$[-5.86; 5.94] \times 10^1$	$[-3.28; 3.36] \times 10^1$	$[-6.91; 7.09] \times 10^1$	$[-3.86; 4.02] \times 10^1$
f_{M2}/Λ^4	$[-7.92; 8.10]$	$[-4.42; 4.60]$	$[-3.15; 3.25]$	$[-1.75; 1.85]$
f_{M3}/Λ^4	$[-8.67; 8.99]$	$[-4.81; 5.13]$	$[-1.01; 1.13] \times 10^1$	$[-0.55; 0.66] \times 10^1$
f_{M4}/Λ^4	$[-2.75; 2.92] \times 10^1$	$[-1.51; 1.68] \times 10^1$	$[-1.20; 1.12] \times 10^1$	$[-6.96; 6.12]$
f_{M5}/Λ^4	$[-3.32; 3.19] \times 10^1$	$[-1.89; 1.77] \times 10^1$	$[-3.95; 3.84] \times 10^1$	$[-2.24; 2.14] \times 10^1$
f_{M7}/Λ^4	$[-1.23; 1.11] \times 10^2$	$[-0.72; 0.60] \times 10^2$	$[-1.40; 1.42] \times 10^2$	$[-7.83; 8.01] \times 10^1$
f_{T0}/Λ^4	$[-2.90; 2.88]$	$[-1.64; 1.62]$	$[-1.51; 1.43]$	$[-0.86; 0.79]$
f_{T1}/Λ^4	$[-1.69; 1.82]$	$[-0.92; 1.05]$	$[-2.21; 2.10]$	$[-1.27; 1.16]$
f_{T2}/Λ^4	$[-4.56; 5.00]$	$[-2.47; 2.92]$	$[-3.92; 4.76]$	$[-2.04; 2.88]$
f_{T5}/Λ^4	$[-0.83; 0.96]$	$[-0.44; 0.57]$	$[-4.33; 4.66] \times 10^{-1}$	$[-2.37; 2.70] \times 10^{-1}$
f_{T6}/Λ^4	$[-0.49; 0.58]$	$[-0.26; 0.35]$	$[-0.61; 0.70]$	$[-0.33; 0.42]$
f_{T7}/Λ^4	$[-1.63; 1.35]$	$[-0.98; 0.71]$	$[-1.57; 1.12]$	$[-1.00; 0.55]$

Our results for the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings are competitive with those reported in Ref. [103] through the $Z\gamma jj$ production at $\sqrt{s} = 13$ TeV and integrated luminosity of 35.9 fb^{-1} at the LHC. A direct comparison of the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings given in Ref. [103] with our results reported in Tables 6, 7 and 8, shows that in some cases our bounds for $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ are more stringent than those reported in Table 4 of Ref. [103] for the CMS Collaboration at the LHC. In Ref. [104], a search at 95% C.L. for the aQGC $f_{M,0,1,6,7}/\Lambda^4$ and $f_{T,0,1,2}/\Lambda^4$ through electroweak production of WW , WZ , and ZZ boson pairs in association with two jets in proton-proton collisions at the LHC with $\sqrt{s} = 13$ TeV is reported. Our results are competitive with those reported by Ref. [104].

Other experimental results on the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings reported by the CMS and ATLAS Collaborations are the followings. With $\sqrt{s} = 8$ TeV and integrated luminosity of 19.7 fb^{-1} the CMS experiment searching for exclusive or quasi-exclusive WW production via the signal topology $pp \rightarrow p\gamma^*\gamma^*p \rightarrow pW^+W^-p$ [41,42]. Their research are translated into limits on the aQGC $f_{M,0,1,2,3}/\Lambda^4$. In addition, the CMS experiment measure the electroweak-induced production of W and two jets, where the W -boson decays leptonically, and experimental limits on the aQGC $f_{M,0-7}/\Lambda^4$ and $f_{T,0-2,5-7}/\Lambda^4$ are set at 95% C.L. [41]. In another investigation with $\sqrt{s} = 8$ TeV and $\mathcal{L} = 20.2 \text{ fb}^{-1}$ of pp collisions the ATLAS experiment, was studied the production of $WV\gamma$ events in $e\nu\mu\nu\gamma$, $evjj\gamma$ and $\mu\nu jj\gamma$ final states [69]. The results reported in

Table 8 Sensitivity measures on the aQGC at the 95% C. L. via the process $e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ at the FCC-he with $\sqrt{s} = 3.46, 5.29$ TeV and 1000 fb^{-1} for the combined data. Here, we consider that only one of the anomalous couplings deviates from the SM at any given time

Couplings	$\sqrt{s} = 3.46$ TeV	$\sqrt{s} = 5.29$ TeV
f_{M0}/Λ^4	$[-2.25; 2.43] \times 10^1$	$[-1.00; 1.03] \times 10^1$
f_{M1}/Λ^4	$[-0.71; 0.82] \times 10^2$	$[-2.61; 2.75] \times 10^1$
f_{M2}/Λ^4	$[-3.36; 3.81]$	$[-1.49; 1.59]$
f_{M3}/Λ^4	$[-1.24; 1.10] \times 10^1$	$[-3.68; 4.46]$
f_{M4}/Λ^4	$[-1.33; 1.26] \times 10^1$	$[-5.91; 5.27]$
f_{M5}/Λ^4	$[-4.36; 4.09] \times 10^1$	$[-1.54; 1.43] \times 10^1$
f_{M7}/Λ^4	$[-1.34; 1.76] \times 10^2$	$[-5.56; 5.15] \times 10^1$
f_{T0}/Λ^4	$[-2.13; 1.52]$	$[-7.26; 6.59] \times 10^{-1}$
f_{T1}/Λ^4	$[-3.12; 2.09] \times 10^1$	$[-8.15; 8.10] \times 10^{-1}$
f_{T2}/Λ^4	$[-0.35; 0.75] \times 10^1$	$[-1.49; 2.22]$
f_{T5}/Λ^4	$[-0.43; 0.69]$	$[-1.91; 2.35] \times 10^{-1}$
f_{T6}/Λ^4	$[-0.54; 1.14]$	$[-0.21; 0.29]$
f_{T7}/Λ^4	$[-1.68; 1.48]$	$[-0.79; 0.39]$

these studies are weaker than those reported in our present article.

To be consistent and put in perspective the current limits with a small luminosity compared to the end run of the HL-LHC, as well as with the projections of limits at future electron-proton colliders with full luminosity, we present the expected bounds on the anomalous quartic $WW\gamma\gamma$ couplings $f_{M,0-1}/\Lambda^4$ and $f_{T,0-1}/\Lambda^4$ at the HL-LHC with $\sqrt{s} = 14$ TeV and $\mathcal{L} = 3000 \text{ fb}^{-1}$:

$$\frac{f_{M0}}{\Lambda^4} = [-4.0; 4.0] \text{ TeV}^{-4}, \tag{22}$$

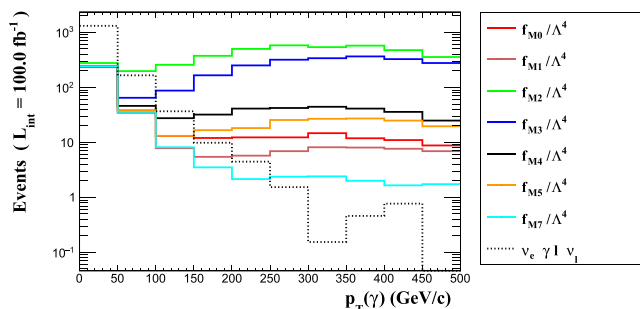
$$\frac{f_{M1}}{\Lambda^4} = [-12.0; 12.0] \text{ TeV}^{-4}, \tag{23}$$

$$\frac{f_{T0}}{\Lambda^4} = [-0.6; 0.6] \text{ TeV}^{-4}, \tag{24}$$

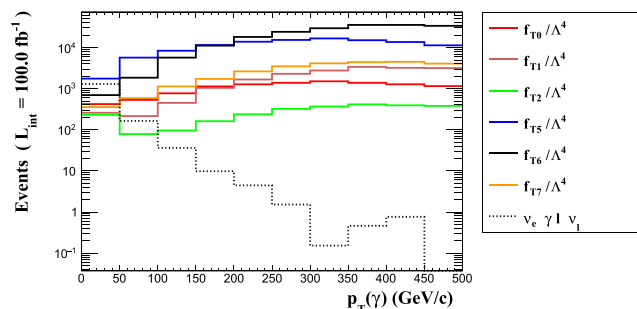
$$\frac{f_{T1}}{\Lambda^4} = [-0.4; 0.4] \text{ TeV}^{-4}. \tag{25}$$

These results show that even in the case of the HL-LHC, our limits on the aQGC given in Tables 7 and 8 are competitive and in some cases are better with respect to the limits given by Eqs. (22)–(25).

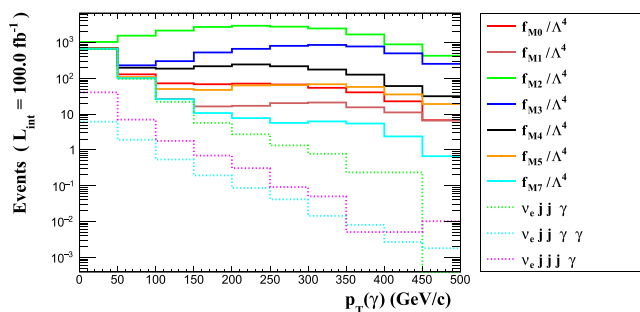
Phenomenological results on the aQGC $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ at the LHeC with $\sqrt{s} = 1.30, 1.98$ TeV and the FCC-he with $\sqrt{s} = 3.46, 5.29$ TeV are presented in Refs. [57–59]. In Ref. [59], the $ep \rightarrow e^- \gamma^* p \rightarrow eW\gamma q'X \rightarrow e\nu_l q'X$



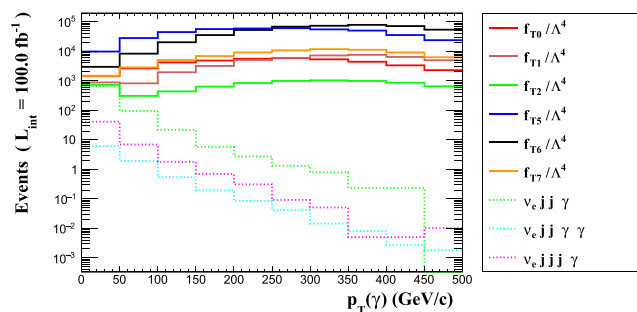
(a) Leptonic decay



(b) Leptonic decay



(c) Hadronic decay



(d) Hadronic decay

Fig. 11 The number of expected events as a function of the p_T^γ photon transverse momentum for the $e^-p \rightarrow e^- \gamma^* p \rightarrow p W^- \gamma \nu_e$ signal and backgrounds at the LHeC with $\sqrt{s} = 1.98$ TeV. The distributions are for

$f_{M,i}/\Lambda^4$ with $i = 0, 1, 2, 3, 4, 5, 7$, $f_{T,j}/\Lambda^4$ with $j = 0, 1, 2, 5, 6, 7$ and various backgrounds for both leptonic and hadronic decay channel of the W -boson

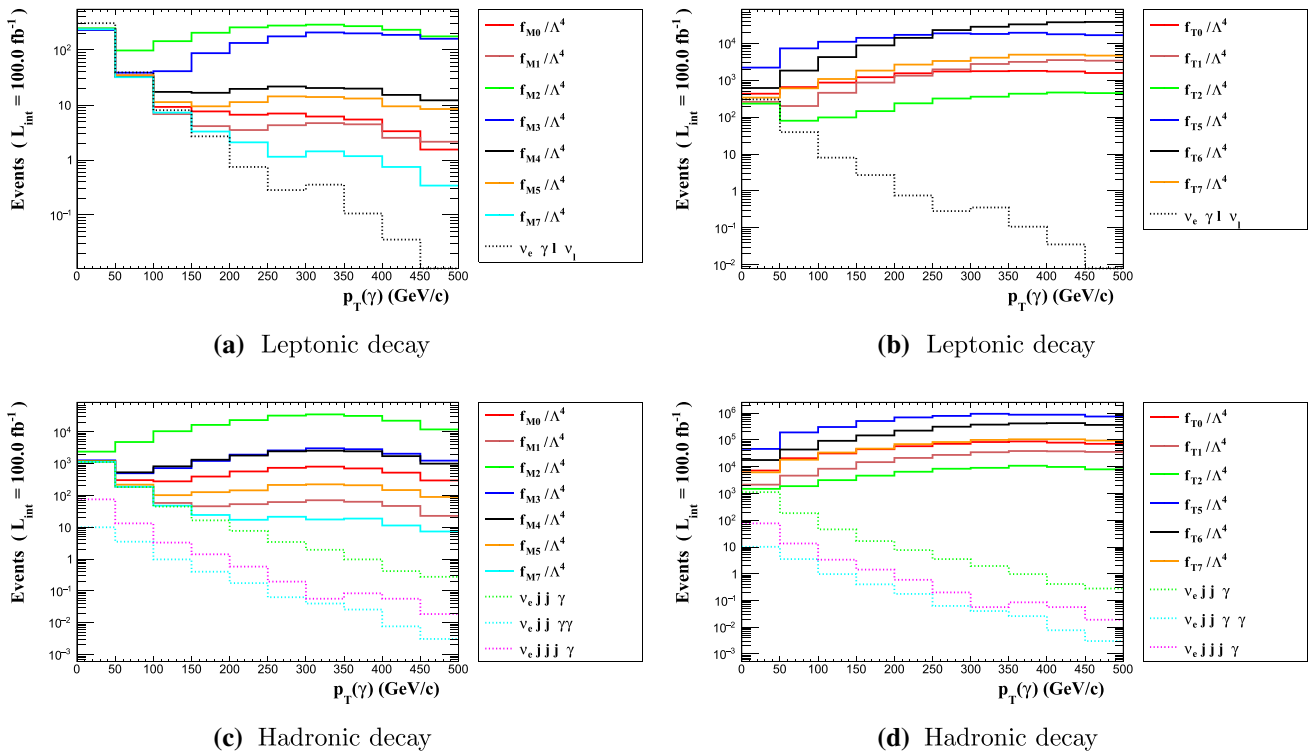


Fig. 12 The number of expected events as a function of the p_T^γ photon transverse momentum for the $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$ signal and backgrounds at the FCC-he with $\sqrt{s} = 3.46$ TeV. The dis-

tributions are for $f_{M,i}/\Lambda^4$ with $i = 0, 1, 2, 3, 4, 5, 7$, $f_{T,j}/\Lambda^4$ with $j = 0, 1, 2, 5, 6, 7$ and various backgrounds for both leptonic and hadronic decay channel of the W -boson

channel gets sensitivity measures of the order of 10^{-1} for some anomalous $f_{T,j}/\Lambda^4$ couplings. In another study, sensitivity measures on the aQGC of the order of 10^1 are reported by Refs. [57,58] via the process $e^-p \rightarrow e^- \gamma^* \gamma^* p \rightarrow e^- W^+ W^- p$ with the subprocess $\gamma^* \gamma^* \rightarrow W^+ W^-$. In the case of our present article, in Tables 6, 7 and 8, we summarize all of the sensitivity measures on the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings obtained at $\sqrt{s} = 1.30, 1.98$ TeV and $\sqrt{s} = 3.46, 5.29$ TeV with the production mode $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$. Our results on the aQGC $f_{M,0-5,7}/\Lambda^4$ and $f_{T,0-2,5-7}/\Lambda^4$ for the different energy stages above mentioned provided sensitivity measures of the order of 10^{-1} , which is similar to those sensitivity measures report by [57–59] at the LHeC and the FCC-he, with other channels. For other reviews experimental and phenomenological, the reader can check Refs. [41,42,57,58,69,105–107].

Photon in the final state of the process $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$ at the LHeC and FCC-he have the advantage of being identifiable with high purity and efficiency. Thus, the single-photon and diphoton channels are especially sensitive for new physics BSM in terms of modest backgrounds, excellent mass resolution and the clean experimental signature. Furthermore, as we mentioned above, in order to quantify the expected limits on $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$, advantage has

been taken in this analysis of the fact that the aQGC enhance the cross-sections at high energies. In Figs. 11 and 12, the number of expected events as a function of the $p_T(\gamma)$ transverse momentum for the $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$ signal and backgrounds at the LHeC and the FCC-he. The distributions are for $\frac{f_{M,i}}{\Lambda^4}$ with $i = 0, 1, 2, 3, 4, 5, 7$, $\frac{f_{T,j}}{\Lambda^4}$ with $j = 0, 1, 2, 5, 6, 7$ and various backgrounds for both leptonic and hadronic decay channel of the W -boson. In these figures, the solid histograms correspond to σ_{Tot} cross-section, and the dashed distribution corresponds to the other SM background sources. In addition, this distribution clearly shows great sensitivity with respect to the anomalous $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ couplings for both cases leptonic and hadronic. The analysis of these distributions is important to be able to discriminate the basic acceptance cuts for $W^- \gamma \nu_e$ events at the LHeC and the FCC-he.

6 Summary and conclusions

In this work, in the effective Lagrange approach, we study the $e^-p \rightarrow e^- \gamma^* p \rightarrow pW^- \gamma \nu_e$ channel at the LHeC and the FCC-he as a way to perform sensitivity measures on the total cross-section and on the anomalous $f_{M,i}/\Lambda^4$

and $f_{T,j}/\Lambda^4$ couplings defining dimension-8 effective operators. Since the aQGC $WW\gamma\gamma$ described through effective Lagrangian have dimension-8, they have very strong energy dependence. Therefore, the anomalous cross-section containing the $WW\gamma\gamma$ vertex has a higher energy than the SM cross-section. In addition, the future ep collider will possibly generate a final state with two or more gauge bosons. Hence, it will have a great potential to investigate the aQGC. High energy accelerated e^- and p beams at these colliders radiate quasi-real photons, and thus $e\gamma^*$, γ^*p and $\gamma^*\gamma^*$ collisions are produced from the e^-p process itself. Therefore, ep colliders will provide an important opportunity to probe $e\gamma^*$, γ^*p and $\gamma^*\gamma^*$ collisions at high energies. For the new physics searches at ep colliders have a very clean experimental environment, since they have no interference with weak and strong interactions.

Regarding the comparison with present experimental limits, we find that the $f_{M,i}/\Lambda^4$ and $f_{T,j}/\Lambda^4$ constraints are significantly competitive with the ones achievable at the CMS and ATLAS Collaborations at the LHC through the $Z\gamma jj$ production [103]. Specifically, our results provide stringent limits on the parameters $f_{T,5}/\Lambda^4$, $f_{T,6}/\Lambda^4$ and $f_{T,7}/\Lambda^4$ by a factor of $\mathcal{O}(1.9-3.7)$ than those reported in Ref. [103]. Our limits stronger at 95% C.L. on the Wilson coefficients are $\frac{f_{r5}}{\Lambda^4} = [-0.237; 0.270] \text{ TeV}^{-4}$, $\frac{f_{r6}}{\Lambda^4} = [-0.330; 0.420] \text{ TeV}^{-4}$ and $\frac{f_{r7}}{\Lambda^4} = [-1.000; 0.550] \text{ TeV}^{-4}$ for $\sqrt{s} = 5.29 \text{ TeV}$ and $\mathcal{L} = 1000 \text{ fb}^{-1}$ at 95% C.L. for the FCC-he.

We conclude by mentioning that our projections at the LHeC and the FCC-he are interpreted in the approach of dimension-8 effective field theory operators through the $e^-p \rightarrow e^-\gamma^*p \rightarrow pW^-\gamma\nu_e$ channel. Our results indicate that the $e^-p \rightarrow e^-\gamma^*p \rightarrow pW^-\gamma\nu_e$ production is convincing for searching for the dimension-8 operators $\mathcal{O}_{M,0-5,7}$ and $\mathcal{O}_{T,0-2,5,6,7}$, and as a consequence of the Wilson coefficients $f_{M,0-5,7}/\Lambda^4$ and $f_{T,0-2,5,6,7}/\Lambda^4$ with clean environments, as well as with good sensitivity.

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