

Letter

η , η' mixing angle and η' gluonium content extraction from the KLOE R_ϕ measurement^{*}

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Received: 11 April 2008 / Revised: 6 August 2008

Published online: 7 October 2008 – © Società Italiana di Fisica / Springer-Verlag 2008

Communicated by Haiyan Gao

Abstract. In this report the extraction of the η , η' mixing angle and of the η' gluonium content from the $R_\phi = Br(\phi(1020) \rightarrow \eta'\gamma)/Br(\phi(1020) \rightarrow \eta\gamma)$ is updated. The η' gluonium content is estimated by fitting R_ϕ , together, with other decay branching ratios. The extracted parameters are: $Z_G^2 = 0.12 \pm 0.04$ and $\varphi_P = (40.4 \pm 0.9)^\circ$.

PACS. 14.40.Aq π , K , and η mesons – 13.20.Jf Decays of other mesons – 12.39.Mk Glueball and non-standard multi-quark/gluon states

1 Introduction

The η' -meson, being a pure $SU(3)$ singlet, has been considered for years the meson within which a gluon condensate contribution can show up. In this paper we try to extract the gluon condensate and the η , η' mixing angle in the constituent quark model using the approach from [1] and the wave function spatial overlapping parameters introduced by ref. [2]. In particular the same method of ref. [3] will be used but in addition the $\pi^0 \rightarrow \gamma\gamma$ and $\eta \rightarrow \gamma\gamma$ branching ratios are fitted according to the prescriptions from [4]. This method is chosen because it relates our measurement of $Br(\phi \rightarrow \eta'\gamma)/Br(\phi \rightarrow \eta\gamma)$ [5] to

the η' gluonium content and the η , η' mixing angle. The η and η' mixing angle and the presence of a gluonium component in the η' -meson have been mostly investigated in the past, but are still without a definitive conclusion [6]. Following the approach from [1,3] the η and η' wave functions can be decomposed in three terms: the u , d quark wave function $|q\bar{q}\rangle = \frac{1}{\sqrt{2}}(|u\bar{u}\rangle + |d\bar{d}\rangle)$, the strange component $|s\bar{s}\rangle$ and the $|glue\rangle$. The wave functions are written as follows:

$$\begin{aligned} |\eta'\rangle &= \cos(\varphi_G) \sin(\varphi_P) |q\bar{q}\rangle + \cos(\varphi_G) \cos(\varphi_P) |s\bar{s}\rangle \\ &\quad + \sin(\varphi_G) |glue\rangle \\ |\eta\rangle &= \cos(\varphi_P) |q\bar{q}\rangle - \sin(\varphi_P) |s\bar{s}\rangle, \end{aligned}$$

where φ_P is the η , η' mixing angle and $Z_G^2 = \sin^2 \phi_G$ the gluon contribution. The ratio of the two branching ratios: $R_{\phi(1020)} = Br(\phi(1020) \rightarrow \eta'\gamma)/Br(\phi(1020) \rightarrow \eta\gamma)$ is related to this decomposition by the formula

$$\begin{aligned} R_{\phi(1020)} &= \cot^2(\varphi_P) \cos^2(\varphi_G) \\ &\quad \times \left(1 - \frac{m_s}{\bar{m}} \frac{Z_{NS}}{Z_S} \frac{\tan \varphi_V}{\sin 2\varphi_P} \right)^2 \left(\frac{p_{\eta'}}{p_\eta} \right)^3. \quad (1) \end{aligned}$$

In this formula $p_{\eta'}$ and p_η are the momenta of the η' - and η -meson, respectively, $m_s/\bar{m} = 2m_s/(m_u + m_d)$ is the constituent quark masses ratio, Z_{NS} describes the spatial wave function overlapping between the $q\bar{q}$ component of the ω -meson and η -meson, and Z_S the one between the $s\bar{s}$ component of the η - and $\phi(1020)$ -meson, ϕ_V is the ω ,

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$\phi(1020)$ mixing angle. The parameters Z_S , Z_{NS} , ϕ_V and m_s/\bar{m} are taken from [7] in which the $Br(\phi(1020) \rightarrow \eta'\gamma)$ and $Br(\phi(1020) \rightarrow \eta\gamma)$ are fitted together with other $V \rightarrow P\gamma$ decays (V indicates the vector mesons ρ , ω , $\phi(1020)$ and P the pseudoscalars π^0 , η , η').

As in the KLOE [8] paper [5] we fit the ratio $R_{\phi(1020)}$ from the KLOE measurement

$$\begin{aligned} R_{\phi(1020)} &= \frac{Br(\phi(1020) \rightarrow \eta'\gamma)}{Br(\phi(1020) \rightarrow \eta\gamma)} \\ &= 4.77 \pm 0.09_{\text{stat}} \pm 0.19_{\text{syst}} \times 10^{-3} \end{aligned}$$

together with the available data [9] on $\Gamma(\eta' \rightarrow \gamma\gamma)/\Gamma(\pi^0 \rightarrow \gamma\gamma)$, $\Gamma(\eta' \rightarrow \rho\gamma)/\Gamma(\omega \rightarrow \pi^0\gamma)$ and $\Gamma(\eta' \rightarrow \omega\gamma)/\Gamma(\omega \rightarrow \pi^0\gamma)$. The dependence of these ratios on the mixing angle φ_P and the gluonium content ϕ_G is given by the following equations:

$$\begin{aligned} \frac{\Gamma(\eta' \rightarrow \gamma\gamma)}{\Gamma(\pi^0 \rightarrow \gamma\gamma)} &= \frac{1}{9} \left(\frac{m_{\eta'}}{m_{\pi^0}} \right)^3 \\ &\times \left(5 \cos \phi_G \sin \varphi_P + \sqrt{2} \frac{f_q}{f_s} \cos \phi_G \cos \varphi_P \right)^2, \quad (2) \end{aligned}$$

$$\frac{\Gamma(\eta' \rightarrow \rho\gamma)}{\Gamma(\omega \rightarrow \pi^0\gamma)} = 3 \frac{Z_{NS}}{\cos(\phi_V)} \left(\frac{m_{\eta'}^2 - m_\rho^2}{m_\omega^2 - m_\pi^2} \cdot \frac{m_\omega}{m_{\eta'}} \right) X_{\eta'}^2, \quad (3)$$

$$\begin{aligned} \frac{\Gamma(\eta' \rightarrow \omega\gamma)}{\Gamma(\omega \rightarrow \pi^0\gamma)} &= \frac{1}{3} \left(\frac{m_{\eta'}^2 - m_\omega^2}{m_\omega^2 - m_\pi^2} \cdot \frac{m_\omega}{m_{\eta'}} \right)^3 \\ &\times \left[Z_{NS} X_{\eta'} + 2 \frac{m_s}{\bar{m}} Z_S \cdot \tan \phi_V Y_{\eta'} \right]^2. \quad (4) \end{aligned}$$

Using the value of Z_{NS} and Z_S from [7], we obtain $\varphi_P = (39.7 \pm 0.7)^\circ$ and $Z_G^2 = \sin^2 \varphi_G = 0.14 \pm 0.04$, $P(\chi^2) = 49\%$. Imposing $\varphi_G = 0$ the χ^2 probability of the fit decreases to 1%. The ratio of the Γ 's is obtained using the branching fractions of the decay and the total decay widths Γ_ω , Γ_{π^0} from the PDG 2006 [9]. All the correlations amongst the measurements of the several branching ratios are taken into account. The correlations are due to the choice of normalising all decay widths to the $\Gamma(\omega \rightarrow \pi^0\gamma)$ and to the use of a constrained fit technique in the PDG 2006 in order to obtain more accurate estimates.

The parameter m_s/\bar{m} is determined mainly by $K^{*+} \rightarrow K^+\gamma$ while ϕ_V is given by the $V \rightarrow \pi^0\gamma$ transitions, giving negligible correlations to the φ_P and Z_G^2 parameters. On the other hand, the parameters Z_S , Z_{NS} are strongly correlated to the mixing angle parameter, φ_P , in eq. (1). The constraint $\Gamma(\eta' \rightarrow \gamma\gamma)/\Gamma(\pi^0 \rightarrow \gamma\gamma)$ is, instead, independent of the parameters Z_{NS} and Z_S .

In ref. [3] a similar procedure to the one of [7] was followed taking into account also the possibility of having a gluonium content. They find $Z_G^2 = 0.04 \pm 0.09$ that deviates of 1 σ from our result but with a larger error.

In [3] and [10] this difference was attributed to the use of overlapping parameters obtained by a fit which assumes no gluonium content. In order to check out this possibility, we have performed several tests on the fit procedure. We first performed a new fit using the overlapping parameter Z_S and Z_{NS} extracted by the fit in ref. [3], where the

Table 1. Summary of the results obtained using new values for Z_{NS} and Z_S from [3].

Used inputs	φ_P	Z_G^2
1, 2, 3, 4 $\Gamma(\omega \rightarrow \pi^0\gamma)$ from PDG 2006 [9]	$(40.1 \pm 0.7)^\circ$	0.12 ± 0.04
1, 3, 4 $\Gamma(\omega \rightarrow \pi^0\gamma)$ from PDG 2006 [9]	$(40.4 \pm 0.9)^\circ$	0.12 ± 0.05
1, 2, 3, 4 $\Gamma(\omega \rightarrow \pi^0\gamma)$ from KLOE [11]	$(40.0 \pm 0.7)^\circ$	0.13 ± 0.04

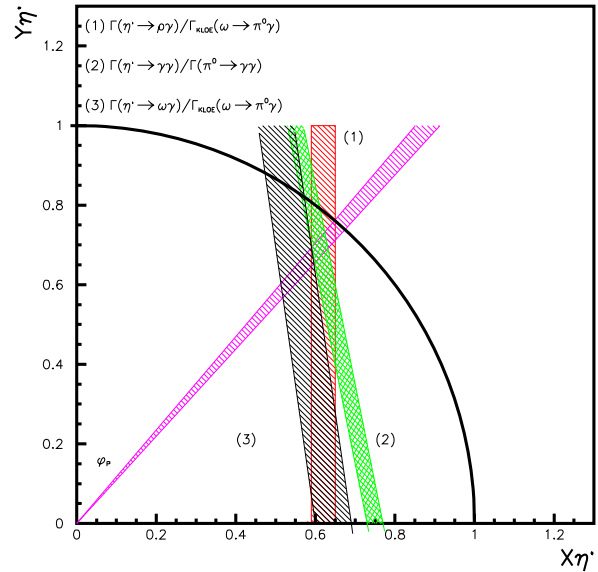


Fig. 1. 68% C.L. region in the plane ($X_{\eta'} = \cos(\varphi_G) \sin(\varphi_P)$, $Y_{\eta'} = \cos(\varphi_G) \cos(\varphi_P)$) for different decay width ratios.

gluonium content was left free, together with all the other parameters: $Z_{NS} = 0.86 \pm 0.03$, $Z_S = 0.79 \pm 0.05$, $\phi_V = (3.2 \pm 0.1)^\circ$, $\frac{m_s}{\bar{m}} = 1.24 \pm 0.07$. We obtained a result in perfect agreement with our previous determination: the errors remain unchanged while the central values move to $\varphi_P = 40.1$, $Z_G^2 = 0.12$.

The value of the fit has been also repeated for different values of Z_{NS} and Z_S in the range 0.5–1.3, and the resulting Z_G^2 varied between 0.07 and 0.18, showing small sensitivity to the used parameters Z_{NS} and Z_S that cannot cause the different result obtained by ref. [3]. Excluding the $P \rightarrow \gamma\gamma$ constraint from the fit we obtain $\varphi_P = (40.4 \pm 0.9)^\circ$ and $Z_G^2 = 0.12 \pm 0.05$, showing that this constraint improves the sensitivity for the gluonium content. All results are summarized in table 1.

A global fit to all the $V \rightarrow P\gamma$ ratios of the branching fractions is in progress. This will allow the overlapping parameters to be left free as in the approach of ref. [3], that is quite different than ours in both fit procedure and input values.

2 New value using KLOE $\Gamma(\omega \rightarrow \pi^0\gamma)$ measurement

We have recently published [11] a new preliminary measurement of the $Br(\omega \rightarrow \pi^0\gamma) = (8.40 \pm 0.19)\%$. Using this

value we obtain $Z_G^2 = 0.13 \pm 0.04$ and $\varphi_P = (40.0 \pm 0.7)^\circ$. The allowed regions in the plane ($X_{\eta'} = \cos(\varphi_G) \sin(\varphi_P)$, $Y_{\eta'} = \cos(\varphi_G) \cos(\varphi_P)$), corresponding to the constraints in eqs. (2)–(4) are shown in fig. 1. Theoretical parameters are taken from [7]. All the allowed regions do not overlap on the no-gluonium line $X_{\eta'}^2 + Y_{\eta'}^2 = 1$, suggesting that the no-gluonium picture is wrong.

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