

Design, performance, and calibration of CMS hadron-barrel calorimeter wedges

The CMS-HCAL Collaboration

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Received: 13 June 2007 / Revised version: 25 February 2008 /
 Published online: 9 April 2008 – © Springer-Verlag / Società Italiana di Fisica 2008

Abstract. Extensive measurements have been made with pions, electrons and muons on four production wedges of the compact muon solenoid (CMS) hadron barrel (HB) calorimeter in the H2 beam line at CERN with particle momenta varying from 20 to 300 GeV/ c . The time structure of the events was measured with the full chain of preproduction front-end electronics running at 34 MHz. Moving-wire radioactive source data were also collected for all scintillator layers in the HB. The energy dependent time slewing effect was measured and tuned for optimal performance.

1 Introduction

The Compact Muon Solenoid (CMS) is a general purpose experiment designed to study pp collisions at the Large Hadron Collider (LHC) [1, 2]. The CMS detector calorimeter has been designed to detect cleanly the diverse signatures of new physics through the measurement of jets with moderate precision and by measuring missing transverse energy flow. The design goal of the calorimeter is to have an energy resolution of $100\%/\sqrt{E} \oplus 5\%$ where E is in GeV [3]. The CMS experiment has a 4 T superconducting solenoidal magnet of length 13 m and inner diameter 5.9 m. The magnet determines many of the features of the CMS calorimeters because the barrel and end-cap calorimeters are located inside this magnet. The hadron calorimeter (HCAL) plays a fundamental role in most anticipated discoveries at the LHC. The HCAL is used to measure the timing and energy of hadronic showers, as well as their angle and position, needed for the generation of level-1 trigger primitives, the high level trigger, and offline reconstruction of jets and missing transverse energy [4–6]. Figure 1 illustrates the calorimeters in and around the CMS solenoidal magnet. The cylindrically symmetric hadron barrel (HB) calorimeter consisting of alternating layers of brass and plastic scintillator plates surrounds the lead tungstate (PbWO_4) electromagnetic calorimeter (ECAL-EB). The HB design maximizes the number of interaction lengths inside the coil, which requires minimizing the amount of space devoted to the active medium. The scintillator tiles are read out with embedded wavelength shifting (WLS) fibers which achieves the required compact and

hermetic design. This technology was first developed at Protvino and by the UA1 collaboration and has been successfully used in an upgrade of the CDF endcap calorimeter [7–10]. Brass was chosen as the absorber material because it is non-magnetic. This design makes construction relatively simple, lends itself to projective tower geometry, and eliminates uninstrumented gaps.

The CMS HCAL contains 9072 readout channels organized into four subsystems: barrel (HB, 2592 channels), endcap (HE, 2592 channels), outer (HO, 2160 channels) and forward (HF, 1728 channels). This paper addresses the design, performance, and calibration of the HB. The performance of the HE, HO, and HF were also extensively investigated and are reported elsewhere [11–14].

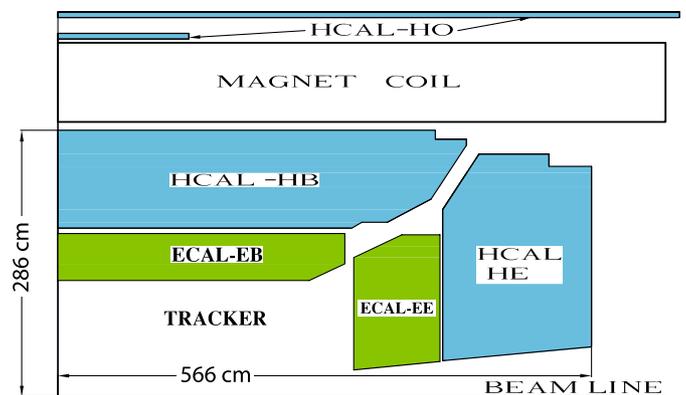


Fig. 1. Location of the ECAL and the HCAL detectors (quarter slice-longitudinal cross section) in and around the CMS magnet

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While the beam tests of the HB prototype and bench tests of the calibration systems were reported previously [15, 16], this paper reports the first measurements of production modules with the complete electronics chain. As discussed in Sect. 6, one of the goals of this study was to relate the radioactive source measurements to beam measurements in order to calibrate the HB wedges which were not exposed to particle beams. After these measurements were performed, the HB wedges were assembled into the final barrel configuration.

This paper is organized as follows. The design details of the hadronic calorimeter are presented in Sect. 2. Topics related to electronics and data acquisition and the test beam setup are in Sects. 3 and 4. Section 5 gives the HB performance in particle beams. Following a summary of the radioactive source calibration in Sect. 6, we draw conclusions and summarize our work in Sect. 7.

2 HCAL barrel design

2.1 Absorber geometry

The HB covers the pseudorapidity range $-1.3 < \eta < 1.3$ and consists of 36 identical azimuthal wedges ($\Delta\phi = 20^\circ$) which form two half-barrels (HB+ and HB-). Each half-barrel is inserted from either end of the cryostat of the superconducting solenoid. Each wedge is further segmented into four azimuthal ($\Delta\phi = 5^\circ$) sectors. The plates are bolted together in a staggered geometry resulting in a configuration that contains no projective passive material for the full radial extent of a wedge (see Fig. 2). The innermost and outermost plates are made of stainless steel to provide structural strength. The scintillator

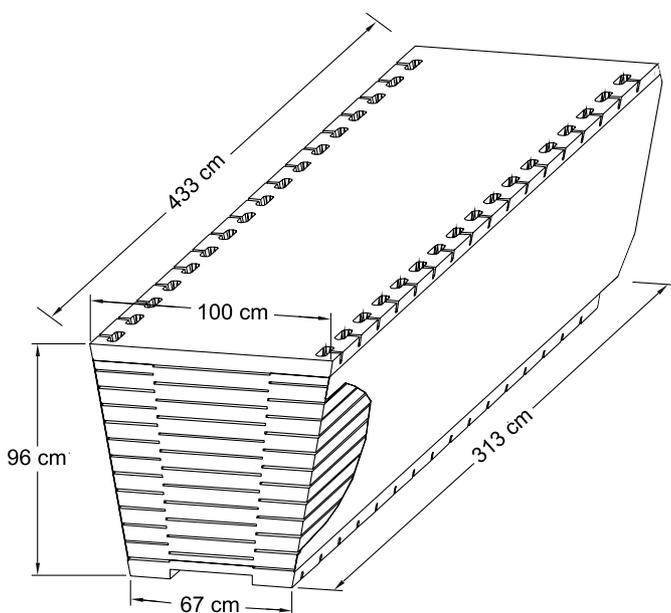


Fig. 2. Isometric view of an HB wedge: the scintillator trays (Fig. 3) are inserted into slots at the end of the wedge

is divided into 16 η sectors, resulting in a segmentation of $(\Delta\eta, \Delta\phi) = (0.087, 0.087)$. The wedges are bolted together and the gap between the wedges is less than 2 mm.

The absorber itself consists of a 40 mm thick front steel plate, followed first by eight 50.5 mm thick brass plates, and then six 56.5 mm thick brass plates, with a final 75 mm thick steel back plate. The total absorber thickness at 90° is 5.82 interaction lengths (λ_I). The HB effective thickness increases with polar angle and is $10.6 \lambda_I$ at $|\eta| = 1.3$. The electromagnetic crystal calorimeter [17] in front of the HB adds $\sim 1.1 \lambda_I$ independent of η .

The brass absorber is commonly known as C26000 (cartridge brass) and composed of 70% Cu and 30% Zn. The density is 8.83 g/cm^3 . The radiation length is $X_0 = 1.49 \text{ cm}$ and the nuclear interaction length is $\lambda_I = 16.42 \text{ cm}$.

2.2 Scintillator

The CMS HCAL active elements consist of about 70 000 scintillator tiles. In order to limit the number of individual physical elements, the tiles of a given azimuthal section and depth layer are grouped into a single scintillator unit, referred to as a tray.

Figure 3 shows the end portion of a typical tray. This design proved to be robust and practical. We tested each scintillator tray and the optical readout chain before installation into the absorber structure. The construction of the absorber structure and the scintillator assemblies were independent.

The first layer of scintillator (layer-0) is located in front of the steel support plate and is made of 9-mm thick Bicon BC408. The last scintillator layer (layer-16) is 9-mm thick Kuraray SCSN81. The others are all 3.7-mm thick Kuraray SCSN81 plates. The active material choice for the HB was Kuraray SCSN81 scintillator because of its long-term stability and acceptable radiation hardness.

A tray is made of individual optically independent scintillators with white painted edges wrapped in Tyvek 1073D sheets. The scintillators are attached to a 0.5-mm thick plastic substrate with plastic rivets. Light from each tile

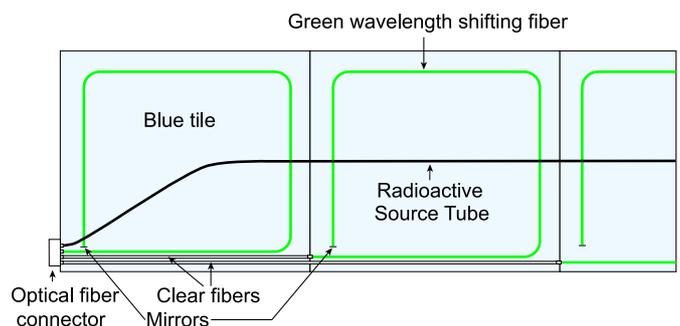


Fig. 3. Schematic of a partial scintillator tray, showing green wavelength shifting fibers, clear fibers, and the radioactive source tube locations

is collected with a 0.94-mm diameter green double-clad wavelength-shifting fiber (Kuraray Y11) which is inserted in a ball-groove machined in the scintillator. The top of the tray is covered with a 2-mm thick polystyrene sheet which carries the optical readout fibers. The end tips of the wavelength shifting fibers embedded inside the scintillator trays are coated by vacuum evaporation technique with a thin layer of aluminum at the tip which serves as a reflector. The average reflectivity is $\sim 83\%$ with a spread of about 6.5% [18].

The wavelength shifting fibers are spliced to clear fibers to minimize optical attenuation in transporting the light from the scintillator plates to the photosensors located a few meters away. A fusion splicer was developed for this task and the light transmission across the splice is 92.6% with an rms of 1.8%. The clear fiber is terminated at a diamond-finished optical connector at the edge of the scintillator tray. An optical cable then further transports light to an optical unit which arranges the fibers into readout towers and guides the light to a hybrid photodiode (HPD) [19].

The completed tray was tested with a collimated ^{137}Cs source. This source illuminates a 4-cm diameter spot on the tray and is positioned with a computer controlled translation stage at selected locations on the tray. We determined the relative light yield of each tile and the uniformity of each tray. The light yield of individual tiles has an rms variation of 4.6%, while the transverse uniformity of a tile has a spread of 4.5%. A uniformly illuminated tile has a light yield variation of about 6%. The overall uniformity from tile to tile as measured for a hadron shower is about 8% per tile. For hadronic showers, the weighted average rms over a tower with 16 tiles is about 3.5% and for muons it is $8\%\sqrt{16} = 2\%$.

For calibration purposes, the top of each tray also houses a 1-mm diameter stainless steel tube, called a radioactive *source tube*, that guides a ^{137}Cs (or ^{60}Co) source welded on the tip of a thin stainless steel wire across the center of each tile in a tray (Fig. 3). During the assembly stage, the scintillator trays are tested by guiding the wire source through the source tubes. The rms of the ratio of the light yields with the collimated source to that of the wire source is 1.3%.

There is considerable material between the active elements of the ECAL-EB and HCAL-HB. Originally, layer-0 was planned to be read independently from the other scintillators to sample early hadronic showers initiated either in the EB or in the inert material. Since the original plan was to read layer-0 independently, a thicker (9 mm) and a brighter scintillator (BC408) was used. It produced about three times more light than the other scintillator layers. For space and other limitations, it was not possible to read layer-0 separately. Since layer-0 was brighter, the optical signals were attenuated by neutral density filters placed between the optical fibers and the HPD housing and added to the light of the other scintillators in the same tower. Transmission filters of 20, 30 and 40% were tested in the beam and the energy resolution had a very broad minimum at 30% and changed by only 1% between 20 and 40%. The 30% transmission neutral density filter was

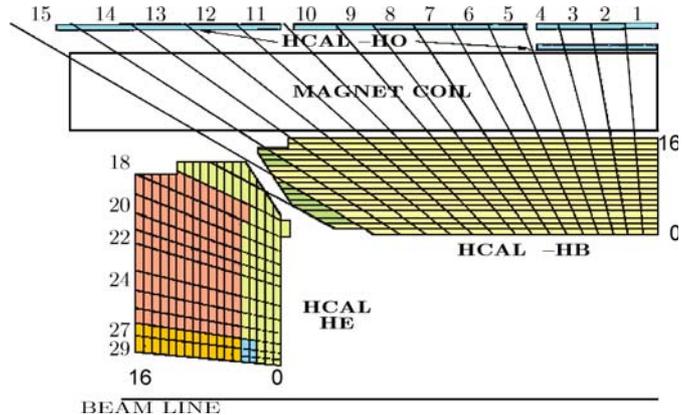


Fig. 4. The HCAL tower segmentation for one-fourth of the HB, HO, and HE detectors is shown above. The numbers on *top* and on the *left* refer to the tower numbers. The numbers on the *right* and on the *bottom* (0–16) indicate the scintillator layers inserted into slots in the absorber. The *shading* represents independent longitudinal readouts in the HB/HE overlap and the small angle regions

chosen so the light level is roughly the same as the other 15 layers.

2.3 Segmentation

The segmentation of the HB is illustrated in Fig. 4 (Table 1). Towers 1 through 14 all have a single longitudinal readout. Towers 15 and 16 are segmented in depth. The front segment of tower 15 contains either 12 or 13 scintillator layers. The rear segment of Tower 15 has three scintillator layers. Tower 16 has five layers in the front

Table 1. The HB tower sizes and depths are summarized below. The thicknesses refer to the center of the tower. Note that tower 16 overlaps with the HCAL endcap (HE) calorimeter. See also Fig. 4

Tower number	$\eta_{\min}-\eta_{\max}$	Thickness (λ_I)
1	0.000–0.087	5.39
2	0.087–0.174	5.43
3	0.174–0.261	5.51
4	0.261–0.348	5.63
5	0.348–0.435	5.80
6	0.435–0.522	6.01
7	0.522–0.609	6.26
8	0.609–0.696	6.57
9	0.696–0.783	6.92
10	0.783–0.870	7.32
11	0.870–0.957	7.79
12	0.957–1.044	8.30
13	1.044–1.131	8.89
14	1.131–1.218	9.54
15	1.218–1.305	10.3
16	1.305–1.392	overlaps HE

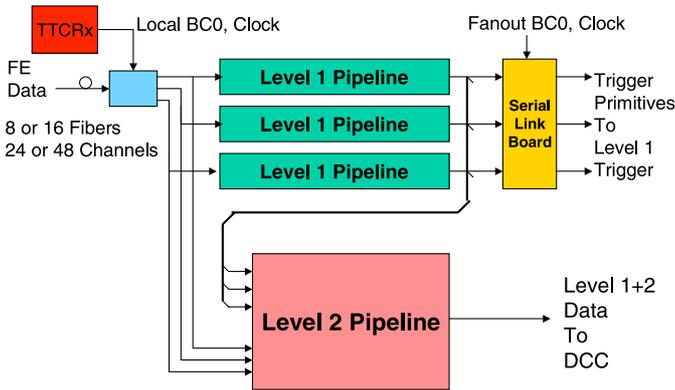


Fig. 6. Block diagram of the HTR electronics. The data from the detector drives two pipelines; one to the CMS trigger and a second to the CMS data acquisition via DCC

The HTR includes two data pipe-lines. The trigger pipe-line assigns the front-end data to a particular LHC bunch crossing and sends them to the CMS trigger. The data acquisition pipe-line stacks the front-end data and sends it to the DCC which performs the data acquisition task.

The raw input data stream in the HTR is deserialized and synchronized to the local clock. A programmable delay of up to a few clock cycles is used to align data from different input fibers. The channel numbers carried on one fiber are demultiplexed. Each channel is then fed to a linearizing look-up table which converts the raw input data to a 16-bit linear energy value. Next, a finite-impulse response (FIR) filter is used to subtract the pedestal and assign all the energy to a single bunch crossing. This performs the same function as a traditional analog shaper, but has the advantage of being easily reprogrammable. Finally, the energy is converted to E_T and compressed to 8 bits according to a non-linear transformation specified by the CMS level-1 calorimeter trigger, and a comparison is performed to check if the signal may repre-

sent a muon. This compressed output plus a muon ID bit is sent to level-1. The final synchronization and serial transmission is performed by a synchronization and link board (SLB).

3.3 Data concentrator card

The LVDS link receiver boards use Channel Link [23] technology from National Semiconductor. Each board contains three independent link receivers which can operate at 20–66 MHz (16-bit words). Buffering for 128 K 32-bit words is provided for each link with provision to discard data if buffer occupancy exceeds a programmable threshold. Event building, protocol checking, event number checking and bit error correction are performed independently for each link.

The DCC logic is designed to operate continuously at the full speed of the two input PCI busses, namely $(33 \text{ MHz}) \times (32 \text{ bits}) \times (2)$. The event builder and output logic circuitry must thus run at an average rate of at least 66 MHz (32-bit words) or 264 MBytes/s. The event builder output is sent in parallel to several destinations.

The outputs are:

1. The data acquisition (DAQ) output: Every event is sent via SLINK-64 to the CMS DAQ. The contents of each event may be controlled by configuration registers.
2. The trigger data output: The trigger information sent to the CMS level-1 trigger is also sent via SLINK-64 to a special trigger DAQ system for monitoring of the trigger performance.
3. The spy output: A selected subset of events is sent to a VME-accessible memory for monitoring and diagnostics.

Error detection and recovery are a primary consideration in a large synchronous system and the DCC contains logic dedicated to this purpose. Figure 7 shows the main DAQ data pipe-line and buffering. Hamming error correction is used for the LVDS links between the HTR and DCC.

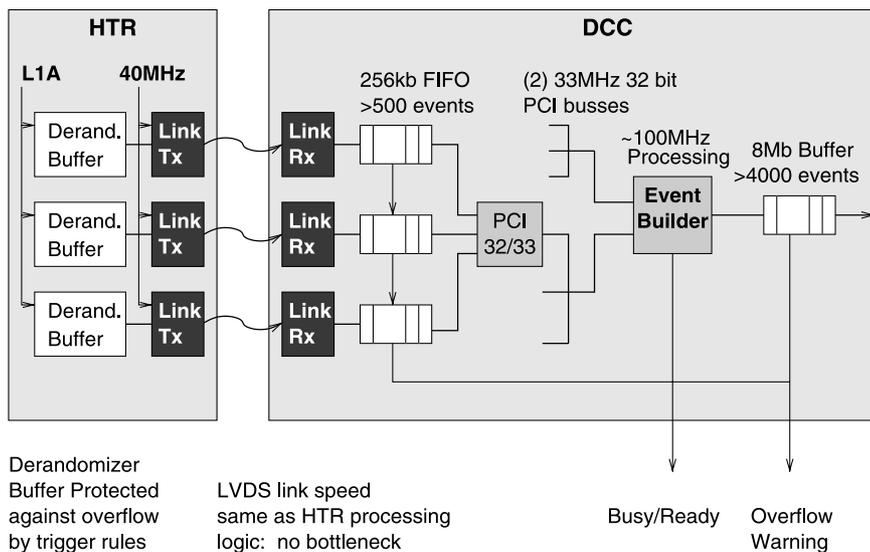


Fig. 7. HCAL DAQ buffering schematic. The HTR sends the received data to the DCC. The DCC stores event fragments in FIFO buffers on the input and output

All single-bit errors are corrected and all double-bit errors are detected by this technique. Event synchronization is checked by means of an event number in the header and trailer of each event, which are checked against the TTC event number.

4 Test beam setup and preliminary measurements

The data were recorded during 2002 at the CERN H2 test beam. A moving platform held two production HB wedges plus a pre-prototype electromagnetic crystal calorimeter which was inserted in front of the HB for some of the measurements. The two-dimensional movement of the platform in ϕ and η directions allowed the beam to be directed onto any tower of the calorimeter. Four scintillation counters were located approximately three meters upstream of the calorimeters and a coincidence between a subset of these counters was used for the trigger. The beam size varied depending on the beam energy and magnet settings but was typically several millimeters in the transverse dimensions.

A pre-prototype electromagnetic calorimeter module was constructed from 49 lead tungstate crystals arranged in a 7×7 array and coupled to individual photomultiplier tubes. The front face dimensions are approximately $2.2 \times 2.2 \text{ cm}^2$, and the length of the crystal is 23 cm or 25.8 radiation lengths. From subsequent test beam data it became clear that production ECAL modules are far superior to this pre-production module. Even though data were taken with this ECAL module, the results will not be discussed further here but in a subsequent paper where EB production modules were available.

4.1 Data sets

Data were taken with electron beam momenta of 20, 30, 50 and 100 GeV/ c , pion beam momenta of 20, 30, 50, 100, and 300 GeV/ c , and a muon beam of momentum 225 GeV/ c . The momentum spread of the pion beam was established by collimators which were typically set to provide momentum spread that is less than $\Delta p/p = \pm 1\%$, and the beam momentum spread contributes negligibly to the measured calorimeter resolution. The test beam is created by extracting the primary proton beam over 1.2 s and impinging on a primary target in the SPS in CERN. The beam is spread uniformly with no particular time structure. The time of each beam particle is determined by a TDC to an accuracy of about 1 ns using the beam counters.

Beams of muons, pions and electrons were directed into the centers of all accessible towers, which included all four 5° sectors of the bottom wedge, but was limited to the bottom two 5° sectors of the top wedge for mechanical reasons. For both wedges, the scan covered the entire η range (16 sectors). These data were taken without the electromagnetic calorimeter module.

Additionally, several special runs were taken with the moving radioactive wire source which illuminated 2092 tiles for the two wedges and with a LED pulser which injected short light pulses to the HPDs.

4.2 Laser calibration and monitoring system

In addition to a LED light injection system, a 300 mJ nitrogen laser is used to monitor the calorimeter and check on the long term relative time calibration. The laser light at 337 nm has two paths. The first path carries the UV light by quartz fibers to a tiny green scintillator inside the RBX housing and lights up all 19 pixels of each HPD. The second path carries the UV light with another quartz fiber path to layer 9 scintillator of each tower. The quartz fiber path length is adjusted so that the timing of the signals in each tower mimics (within 1 ns) an actual event in the center of the CMS detector. The laser system was tested in the test beam and indeed the timing accuracy was within one nanosecond.

4.3 Noise performance

The front-end electronics was operated at 33.79 MHz for all measurements reported in this paper. Figure 8 shows the noise distribution at the highest gain setting of the multi-range ADC for a single QIE channel. The signal is summed over four time slices, and the signal arrives within 2 ns at the start of the second time slice. The least count corresponds to 2080 electrons measured by precise charge injection. The observed noise spread in the pedestal is about 5000 electrons. From the beam measurements, this corresponds to an energy of $\sim 0.2 \text{ GeV}$. After the 2002 test beam, the front-end electronics was upgraded to 40 MHz, the operational mode of the LHC. There is no significant common mode noise.

4.4 Time structure

The time structure of the HB pulse is dominated by the decay time of the wavelength shifting fibers. The front-end electronics, including the photo-detector is fast enough not to distort the pulse shape. An accurate knowledge of the pulse shape is essential to correlate energy deposit with a particular LHC beam crossing. We made two independent measurements of the time structure of the pulse of deposited energy. First, we investigated with a photomultiplier and commercial electronics; and second, we used the HPD and QIE electronics.

4.4.1 Measurements with a photomultiplier tube

We used a single 10-stage, 2-inch photo-multiplier tube (RCA 6342A) and the output was fed into a digital oscilloscope which recorded voltage in 0.4 ns bins. For this study, the output of 3×3 towers were read by the photomultiplier tube, with the beam aimed at the center of

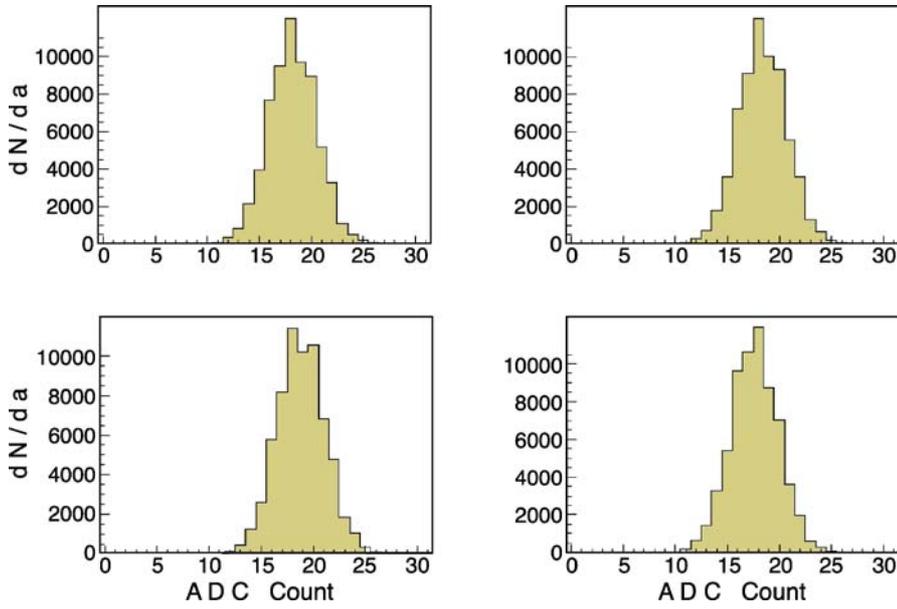


Fig. 8. Measured pedestal distributions (dN/da) in units of raw ADC counts in the most sensitive QIE range. The rms is 2.3 ADC counts

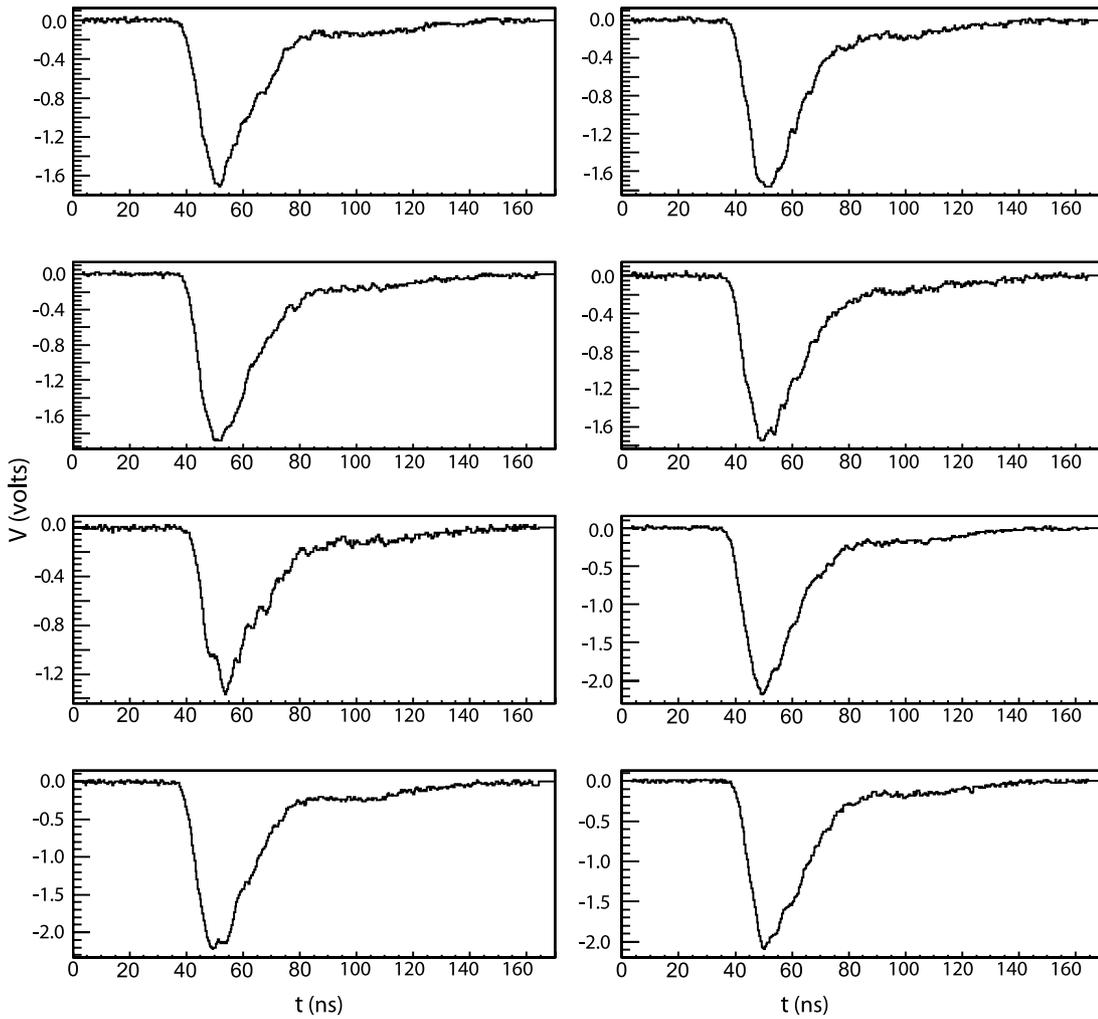


Fig. 9. Calorimeter pulses for eight individual events, voltage vs. time, observed with a photomultiplier for 300 GeV/c pion showers

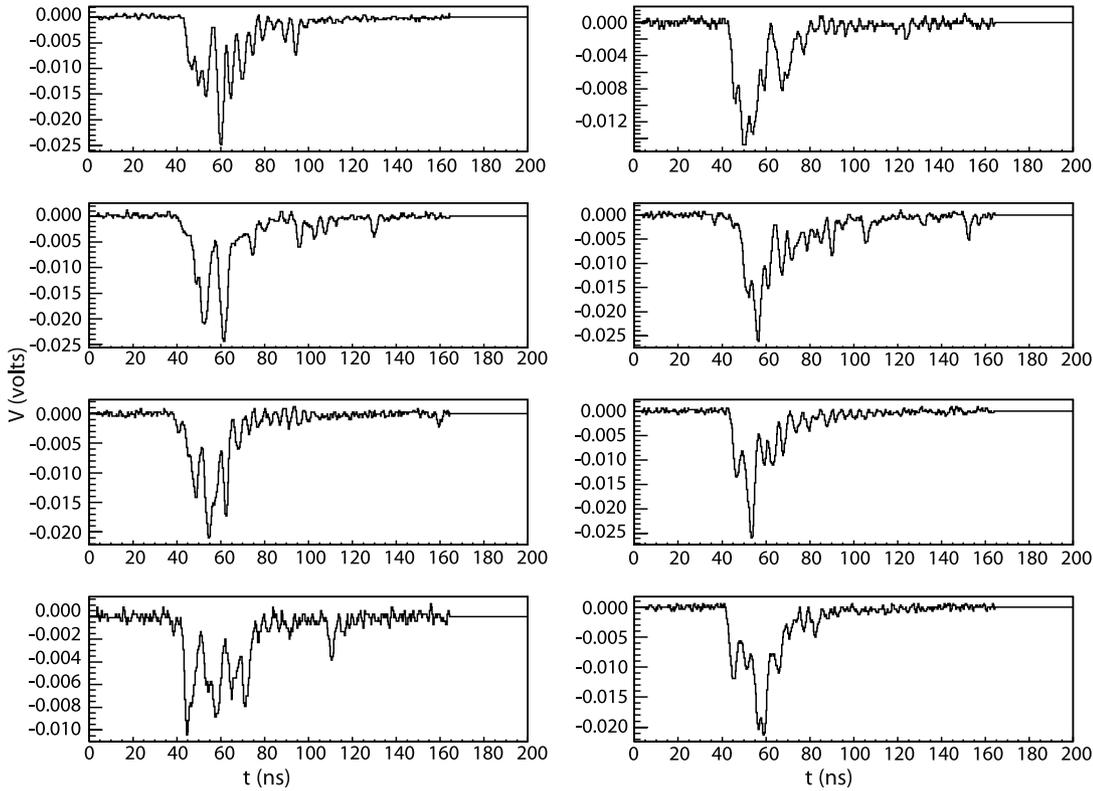


Fig. 10. Calorimeter pulses for eight individual events, voltage vs. time, observed with a photomultiplier for 225 GeV/c incident muons

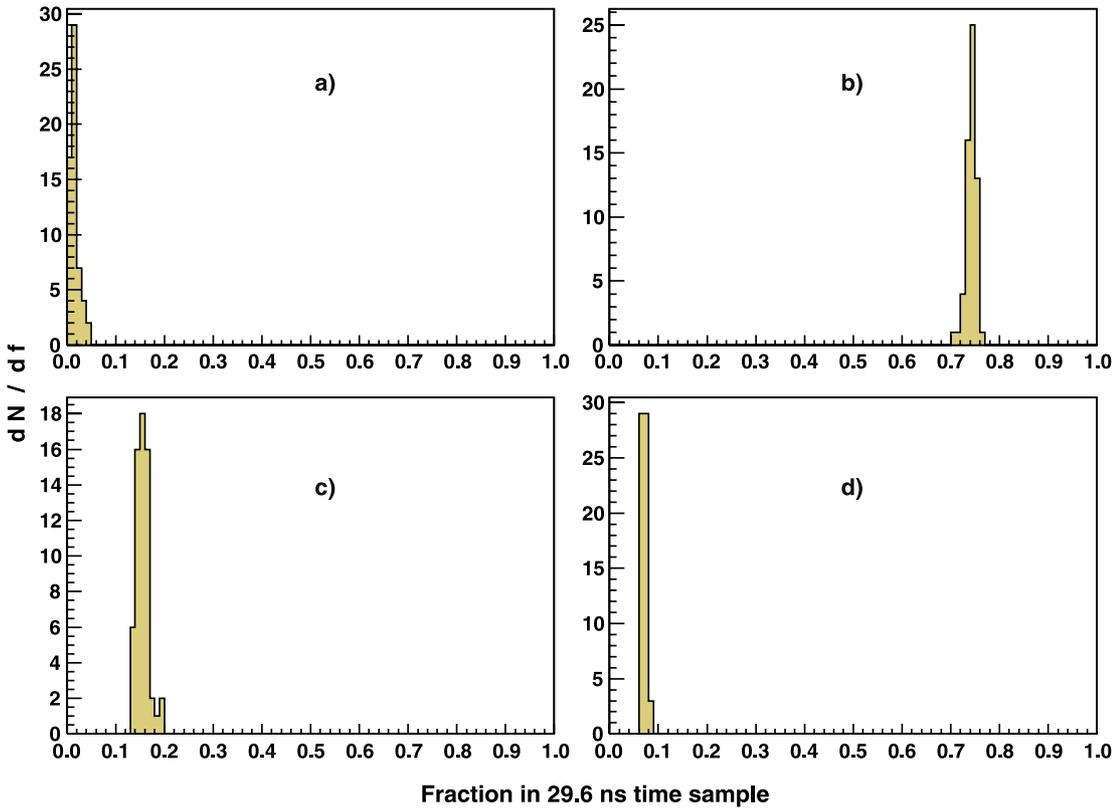


Fig. 11. Fraction of energy (f) observed in 29.6 ns time samples for 200 GeV/c pion showers. **a** First time sample, **b** second time sample, **c** third time sample, and **d** fourth time sample

the central tower. We recorded 300 and 20 GeV/c pion and 100 GeV/c electron showers, as well as signals from high energy muons. Data from 300 GeV/c pions and 225 GeV/c

muons are shown in Figures 9 and 10, respectively. In the case of muons, fluctuations in the number of photoelectrons lead to large event by event fluctuations.

4.4.2 HPD/QIE measurements

Twenty time samples were recorded for each QIE channel per event. The timing was adjusted so that the event occurred near the middle of the time sequence (sample number 10) allowing us to record the noise performance well before and after the energy pulse. At a frequency of 33.79 MHz, each time sample corresponds to 29.6 ns (slightly longer than the 25 ns between LHC bunch crossings). Figure 11 shows the fraction of energy observed in four consecutive time bins of width 29.6 ns. The energy was summed over 3×3 towers, with the beam positioned in the center of the central tower. Using the TDC time information from the beam counters, the events were selected where the beam signal starts in the first 2 ns of the second of four time slices. Most of the signal is collected in two time samples. Note that $\sim 75\%$ of the total energy arrives in a single time sample and this fraction has small dispersion. The reduced clock frequency used in these tests is the result of a design problem in the first procured QIE chips used at the time of data taking. Subsequent beam tests with the upgraded electronics functioning at 40 MHz showed essentially the same results.

Figure 12 shows pulse shapes for 30 GeV/c electrons and 300 GeV/c pions. For this measurement the signal from 3×3 towers around the beam was read by the QIEs. The timing information was obtained by using the TDC information provided by the beam counters. The time distribution of the signal is completely dominated by the wavelength shifting green fibers. The pulse shape was stable for showers for momentum range from 20 to 300 GeV/c and was independent of particle type.

If we select the phase of the test beam particles by requiring that 10%–12% of the observed energy is in the 10th time sample, then all the events have the same temporal shape. We have the freedom to adjust this phase during CMS running using programmable delays provided for each individual channel of the front-

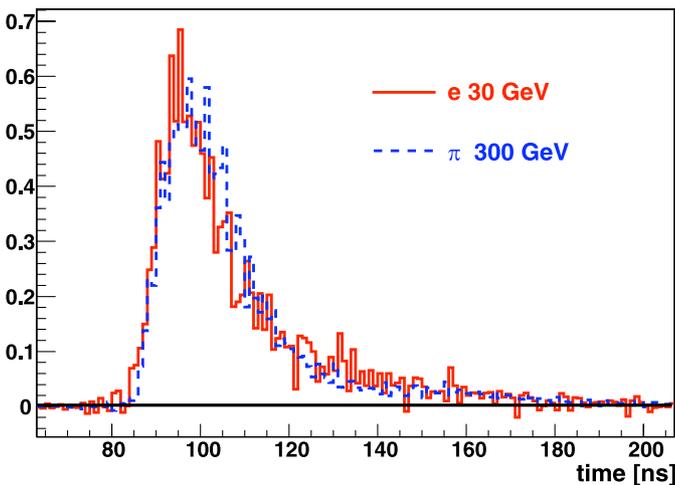


Fig. 12. Pulse shape for 30 GeV/c electrons and 300 GeV/c pions

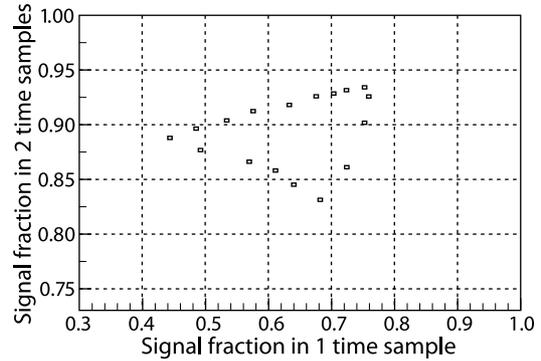


Fig. 13. Fraction of signal observed in two time samples vs. in a single time sample as the phase is changed in 18 intervals spanning a complete time sample. The zero point is suppressed for each axis

end electronics. Figure 13 shows the fraction of energy in two time samples (vertical scale) versus the fraction of energy observed in a single time sample (horizontal scale). The data points represent different choices of the phase. It is possible to have 75% of the energy in a single time sample and more than 90% of the energy in two time samples. Note that the signal fraction measurement thus can provide an arrival measurement of ~ 2 ns.

The QIE introduces a time delay of up to the order of 10 ns depending on pulse height. This is due to the input impedance of the QIE. The time slew can be reduced at the expense of increasing the amount of noise. Figure 14 shows the time slew effect measured in a test setup in the laboratory and test beam measurements. The production modules were constructed with about 5000 e noise and a maximum time slew for very low energy particles of up to 10 ns. Pulse height measurements can be used to correct for this effect.

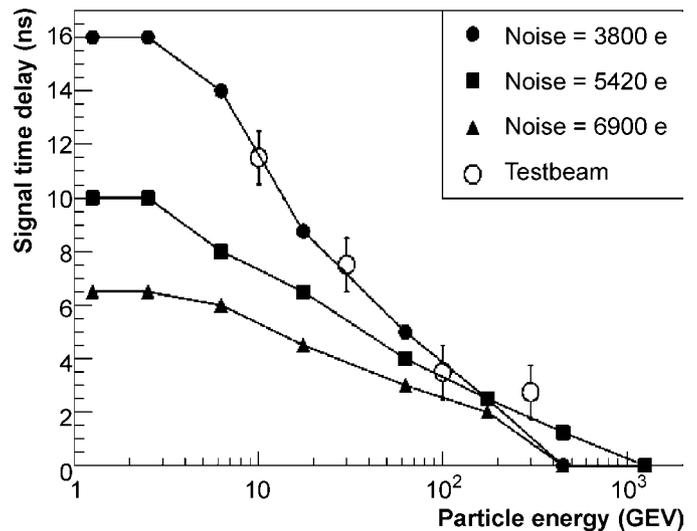


Fig. 14. Time slewing vs. energy for laboratory bench and test beam measurements for three choices of pedestal noise

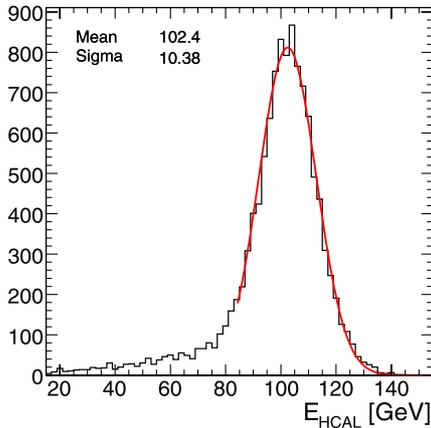


Fig. 15. The energy distribution in the HB (5×5 towers) for 100 GeV/ c pions. The *solid line* is a single Gaussian fit to the histogram data with best fit values for the Gaussian parameters shown in the figure legend

5 The HB response to pions and muons

The absolute energy scale of the HB was defined using a beam of 50 GeV/ c pions directed at the calorimeter. Even though the pre-prototype EB was in the beam, only those pions which deposited less than 2 GeV in the EB were selected for the HB calibration. The calibration factor was extracted as the ratio between the pion beam momentum and the mean of the distribution of summed HB charge deposited in a 5×5 tower array centered on the beam position. Typical calibration constant is 0.2 GeV/fC.

With the 100 GeV/ c pion beam centered on a tower, the fraction of energy observed in the central tower is on average 82%. If the energy in the neighboring eight towers is added, forming a 3×3 array, 98% of the pion energy is observed. The energy response for 100 GeV/ c pions in a 5×5 array is shown in Fig. 15. The beam was centered at tower 4 in η direction, and the energy was summed over four time slices. As stated earlier, only events that deposited less than 2 GeV in the electromagnetic calorimeter were selected, and most of these events deposited less than 1 GeV in the EB. The figure shows the energy in the HB alone. A Gaussian fit gives the energy resolution σ at this energy of about 10%. The low energy tail in Fig. 15 is due to leakage of late developing showers. The outer hadron (HO) calorimeter in CMS is designed to capture this leakage but it was not installed during these tests.

It is important to understand the response of the HB to muons in order to provide redundant calibration to the primary radioactive source method but also for particle identification purposes. A feature bit will be reported for use in higher level triggers, when a muon is identified by the HB. The signal, with 225 GeV/ c incident muons, is about 2.5 GeV on average and clearly visible above the electronic noise (~ 0.2 GeV). Figure 16 shows the muon energy response from a single tower. A Landau fit to this distribution results in 1.64 GeV for the most probable and 2.5 GeV for the mean value. A hadronic shower deposits most of its energy in two or three depth layers with large event

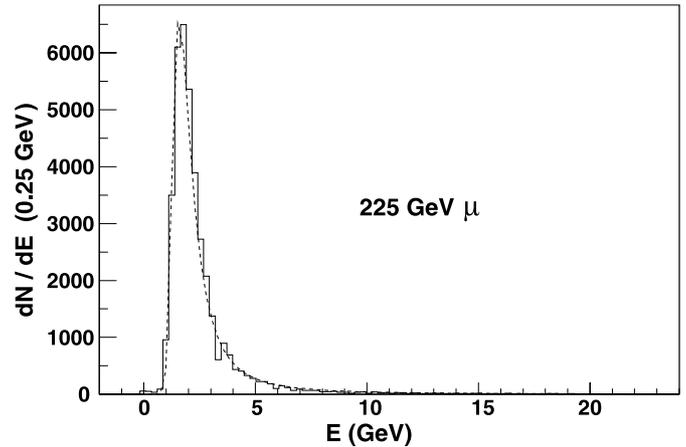


Fig. 16. Observed energy distribution in the HB for 225 GeV/ c incident muons. The most probable value for the energy deposited is 1.64 GeV while the mean is 2.5 GeV. The *dashed curve* is the result of a fit by a Landau distribution

by event fluctuations, so the constant term is substantially larger.

6 Radioactive source calibration

The radioactive source injects the same amount of energy into each tile of each tower. The DC value of the source current is measured by the QIEs at the center of each tile by reading the total charge every 25 ns. The rms of the pedestals is about 0.7 fC. The signal above pedestal is about 1 fC, depending on the strength of the source with an rms value of about 1 fC. During a typical measurement of a tile about 250 000 measurements are recorded, and the rms of the average signal above pedestal is about $1 \text{ fC} / \sqrt{250\,000} = 0.002 \text{ fC}$. Even though an individual recording of a measurement has a error of about 100%, the very large number of measurements reduces the error of the mean to 0.2%. The normalization is performed by comparing the sum of the tile currents (suitably weighted with an average shower shape in depth) to the energy response of the tower in the test beam of known energy. The achieved absolute calibration is better than 2% and reproducible to 1% [16].

The radioactive source calibration by ^{137}Cs (or ^{60}Co) is performed for every scintillator tile. The ratio of the radioactive source signal to the energy response to pions in the beam for each tower provides the initial calibration. This ratio depends on the type and activity of the source. The signal means from 1745 tiles, corrected for the different fiber attenuation, are shown as a histogram in Fig. 17. The rms width of the distribution is 8% and it is consistent with the measurements of tile uniformity made during tray construction.

The dependence of the ratio of radioactive signal to the 100 GeV/ c electron beam signal on the tower number is shown in Fig. 18. The same dependence for the ratio of radioactive signal to the muon signal is presented in Fig. 19.

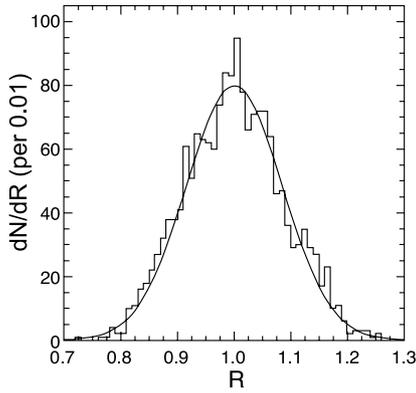


Fig. 17. Distribution of scintillating tile response (R) to the source calibration normalized to the mean of the distribution. The *line* represents a single Gaussian fit to the data. The fitted rms spread of about 8% is consistent with the spread in values measured with the collimated source made during tile assembly

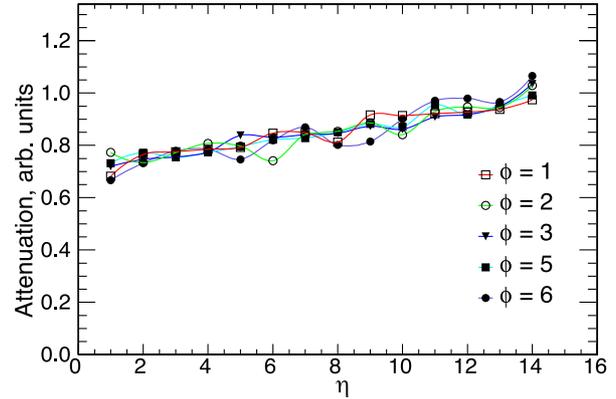


Fig. 20. Measurement of the attenuation vs. η number for five different ϕ numbers (80 towers total). The attenuation occurs in the propagation of the signals in fibers from the towers to the readout box which is located on the $\eta = 14$ end of the wedge

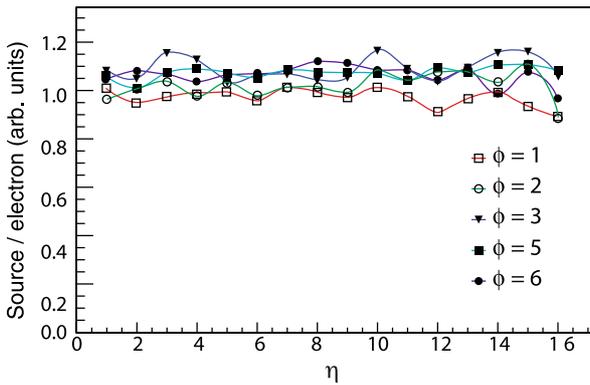


Fig. 18. Ratio of the radioactive source signal (fit to peak value) to the 100 GeV/c incident electron signal in the HB vs. η number of the tower for five different ϕ numbers (80 towers total)

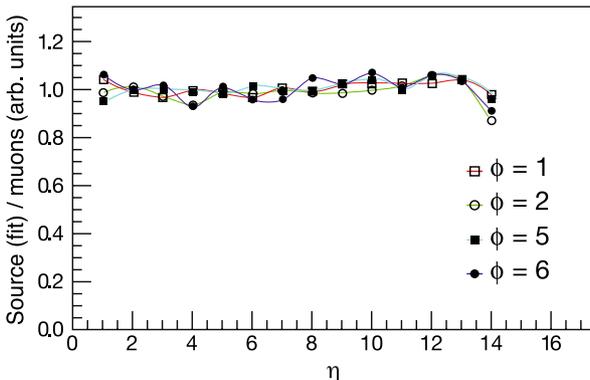


Fig. 19. Ratio of the radioactive source signal to the HB signal due to a 225 GeV/c beam of incident muons vs. η number for four different ϕ numbers (64 towers total)

The electron response agrees with the radioactive source measurements to an rms of 5%, while the muon data agree to 3%. The greater spread in the electron data is due to

the fact that the electron shower is concentrated in the first few scintillator layers, while the source measurement is averaged over all layers with equal weight as is the muon signal. While the muon data are better suited to establish the HB tower-to-tower relative calibration, the data from electrons establish the absolute energy scale. Therefore, we have a cross check of the calibration established by pions.

The length of the optical fibers between the scintillators and the HPDs varies with η . Figure 20 shows the relative response to the radioactive source as a function of η . These data show that the light level at the HPD is about 30% lower for $\eta = 1$ towers when compared to $\eta = 14$ towers. The $\eta = 15$ and 16 towers are not included in these measurements because they are segmented longitudinally.

7 Summary and conclusions

We have described the design and construction of the CMS HCAL barrel calorimeter. The results presented in this paper come from the first data set with the nearly final electronics that was available at the time of the beam tests (2002). The HB electronics noise is HB 0.2 GeV per tower. The pulse shape was determined to be 75% contained in a single 29.6 ns time sample. A radioactive source system was used to record the response of each scintillator tile of HB to the source. The relationship of the HB response to particle beams and sources then established a method to extend the calibration to all those HB modules which were not exposed to particle beams. Thus, these data are used to establish the initial (pre-LHC beam) calibration of the entire HB detector. Since only four HB wedges out of 36 wedges were eventually calibrated in the test beam, it is crucial that the measurements by the radioactive wire source to be used for calibration give the correct energy calibration for all the wedges. The measurements described in this paper show that the consistency between the test beam and the wire source measurements is better than 4% for hadron showers. Thus, the wire source measurements

of every tile of every tower of every wedge will have a calibration constant (initial) to about 4%. The calibration will be improved in the CMS by physics events. The Hadron Barrel calorimeter design and construction meets the requirements for physics as stated in the Technical Design Report and will be ready to take data on day one of LHC operations.

Acknowledgements. This project was carried out with financial support from CERN, the U.S. Department of Energy (DoE), the U.S. National Science Foundation (NSF), RMKI-KFKI (Hungary, OTKA grant T 016823), the Scientific and Technical Research Council of Turkey (TÜBİTAK), Turkish Atomic Energy Agency (TAEK), Bogazici University Research Fund (Grant no: 04B301), the Russian Ministry for Education and Science and the Russian State Committee for Atomic Energy.

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