



3PA-induced optical limiting in pure and barium borate decorated MoS₂ nanocomposites

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Abstract

Pure MoS₂ and barium borate nanorods embedded MoS₂ microspheres were successfully synthesized by adopting an environmental-friendly hydrothermal method. Powder XRD data confirmed the formation of BBO: MoS₂ nanocomposite made of barium borate [$2\theta = 22^\circ$, (018)] along with the MoS₂ [$2\theta = 14^\circ$, (002)]. The presence of few layers MoS₂ was ascertained from the measured difference in characteristic E_{2g}^1 (379 cm⁻¹) and A_{1g} (407 cm⁻¹) Raman modes of MoS₂. FESEM images portrayed the formation of microspheres for pure MoS₂ due to reduction in the surface energy induced by citric acid. In the nanocomposite, barium borate embedded themselves as nanorods upon MoS₂ microsphere forming an urchin-like structure. Ground-state absorption studies revealed a hyperchromic shift in absorption peak due to the overlapping of transition states between highly transparent barium borate and strongly absorbing MoS₂ in the visible region. Femtosecond laser pulses (800 nm, 150 fs, 80 MHz) were employed to perform the open-aperture Z-scan experiments revealing that both pure MoS₂ and BBO: MoS₂ nanocomposites possessed reverse saturable absorption (RSA), attributed to both 2PA and 3PA processes. Interestingly, pure MoS₂ exhibited sequential 3PA (1PA + 2PA) due to the presence of near-resonant energy states, while the dominance of transparent BBO in the nanocomposite induced genuine/instantaneous 3PA process. In closed-aperture Z-scans (for determining the sign and magnitude of nonlinear refraction), the pattern switched from self-focusing (MoS₂) to self-defocusing (BBO: MoS₂) nature. Both pure MoS₂ and BBO: MoS₂ nanocomposites exhibited strong optical limiting with a lower onset limiting threshold (0.11 μJ/cm²) due to synergetic effects of nonlinear absorption (3PA) and nonlinear refraction (self-defocusing). BBO:(0.02 M) MoS₂ nanocomposite possessing a urchin structure depicted strong NLO coefficients with $\gamma_{3PA} = 2.12 \times 10^{-21} \text{ m}^3/\text{W}^2$, $n_2 = -11.1 \times 10^{-17} \text{ m}^2/\text{W}$ and $\chi^{(3)} = 14.0 \times 10^{-19} \text{ m}^2/\text{V}^2$, which were higher than pure MoS₂ and other composites.

Keywords MoS₂ decorated barium borate · Hydrothermal · Two-/three-photon absorption · Z-scan

1 Introduction

A wide variety of inorganic and organic materials are being extensively investigated to achieve efficient nonlinear optics (NLO)-driven devices for applications such as in laser medicine, signal processing, defence and precision measurement. Over the past few years, graphene, two-dimensional transition metal dichalcogenides (TMDC),

topological insulators (TI), metal organic frameworks (MOF) have been investigated extensively as a promising optical limiting material for laser safety devices and saturable absorber for ultrafast fibre lasers [1–3]. Among TMDCs, molybdenum disulphide (MoS₂) has garnered increasing attention, primarily because of its peculiar band structure with exotic properties and represents an abundant, geographically ubiquitous and potentially cheap analogue

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of graphene [4]. It is a typical transition metal sulphide with a layered structure of covalently bonded S–Mo–S, which is bonded to adjacent layers by weak van der Waals (vdW) interactions between neighbouring S–S layers [5]. It possesses layer-dependent optical, electrical and thermal properties, i.e. the bandgap of MoS₂ decreases with increase in the number of layers due to the quantum confinement effects (*direct band gap: 1.9 eV to indirect band gap: 1.2 eV*) [6]. The ultimate crossover from direct to indirect bandgap nature can be accounted for a combined effect of quantum confinement and long-range columbic effects [7]. Further, thermal conductivity of the MoS₂ decreases from bulk material to a lower dimension nanostructures [8]. Earlier works on the investigation of NLO properties of pure MoS₂ revealed a saturable absorption (SA) behaviour, i.e. for higher input powers the absorption becomes nonlinear and the material exhibited high transmittance [9]. An interesting switch-over in the nonlinear absorption behaviour of MoS₂ saturable absorption to reverse saturable absorption was demonstrated through nanocomposites formation [10]. Especially, incorporating ZnO [11], TiO₂ [12] and organic polymers like PMMA [13] yielded reverse saturable absorption (RSA) due to two-photon absorption (TPA) based optical limiting in different MoS₂ nanocomposites. Optical limiting experiments on ZnO and PMMA composites with MoS₂ demonstrated low onset optical limiting threshold of 1.1 J cm⁻², 2.3 J cm⁻², respectively [11]. Impressed by the interesting NLO behaviour of MoS₂ nanocomposites, this article attempts to replace the above mentioned decorative materials with a well-known inorganic NLO material, beta barium borate (BBO). The chosen decorative exhibits many interesting properties such as wide transparency (190–3500 nm), high laser damage threshold, large birefringence and good mechanical properties suitable for NLO devices [14]. Therefore, by incorporating BBO, the major concern of lower linear transmittance in the visible-NIR region can be circumvented in MoS₂, which in turn makes them suitable for optical limiting applications. Often, these nanocomposites are expected to possess strong thermal stabilities which are essential for performing optical limiting action using intense ultrashort pulses. Additionally, advanced photonic materials made up of materials having extreme nonlinear absorption (reverse saturable and saturable absorption) are attractive for laser photonic devices such as optical limiters and mode lockers [15]. Thus, the combination of saturable absorbing MoS₂ with reverse saturable absorbing BBO nanostructures can be a potential system to investigate nonlinear absorption-based optical limiting for widely developing femtosecond IR lasers. These materials can be utilized to protect optical sensors from damage against the intense laser pulses and interrupt the transmission of a powerful laser pulse to eye-safe levels. Thus, the

NLO properties of pure MoS₂ and BBO decorated on different molar concentration of MoS₂ were studied by open-aperture (OA) and closed-aperture (CA) Z-Scan technique using femtosecond laser pulses (800 nm, 150 fs, 80 MHz). Such an investigation helps us to further the understanding of NLO responses at nanoscale and provide avenues to develop new promising optical limiters for highly dangerous ultrafast IR laser pulses.

2 Material preparation

Preparation of pure MoS₂ and decoration of BBO on the surface of MoS₂ structure were achieved by simple one-pot hydrothermal technique. Initially, sodium molybdate and thiourea were taken in ratio of 2:3 and were dissolved in 120 ml distilled water. To control the agglomeration in the MoS₂ system, citric acid was added as a surfactant. The homogeneous white solution obtained on addition of citric acid was transferred to 150 ml Teflon-lined autoclave set-up and maintained at 180 °C for 24 h [16]. Following heat treatment, the solution was brought to room temperature naturally. The acquired black colour suspension in an orange colour solution was washed several times using ethanol and water and filtered repeatedly to remove the excess surfactant from the suspension. Finally, the obtained powders were kept at 200 °C in an oven to remove the adsorbed water molecules. For decoration of BBO on MoS₂, barium chloride, boric acid and sodium hydroxide were dissolved in 80 ml distilled water and kept at stirring for one hour [17]. After 30 min, 0.01 M, 0.02 M and 0.03 M concentration of acquired MoS₂ powders were added to the solution and kept at 120 °C for 24 h. The acquired BBO: MoS₂ particles were repeatedly washed and dried at 400 °C for 3 h resulting in the formation of ash white colour powder.

3 Preliminary identification: XRD

Preliminary confirmation of the material formation was achieved through powder XRD studies employed with Rigaku Ultima IV XRD at the scan speed of 7.19 degrees/min. The recorded XRD pattern of prepared pure and barium borate decorated MoS₂ is as shown in Fig. 1. The formation of pure MoS₂ and BBO: MoS₂ nanocomposites was confirmed through indexing the observed XRD diffraction peaks using JCPDS card Nos. 065-1951 (MoS₂) and 015-0862 (BBO). Identified diffraction crystal planes of MoS₂ are (002), (004), (100), (102), (103), (105), (106), (110), (112) which corresponds to the angle $2\theta = 14^\circ, 29^\circ, 32^\circ, 34^\circ, 39^\circ, 47^\circ, 55^\circ, 58^\circ, 60^\circ$ and beta barium borate are (018), (110), (1010), (202), (205), (119), (1112), (218), (306),

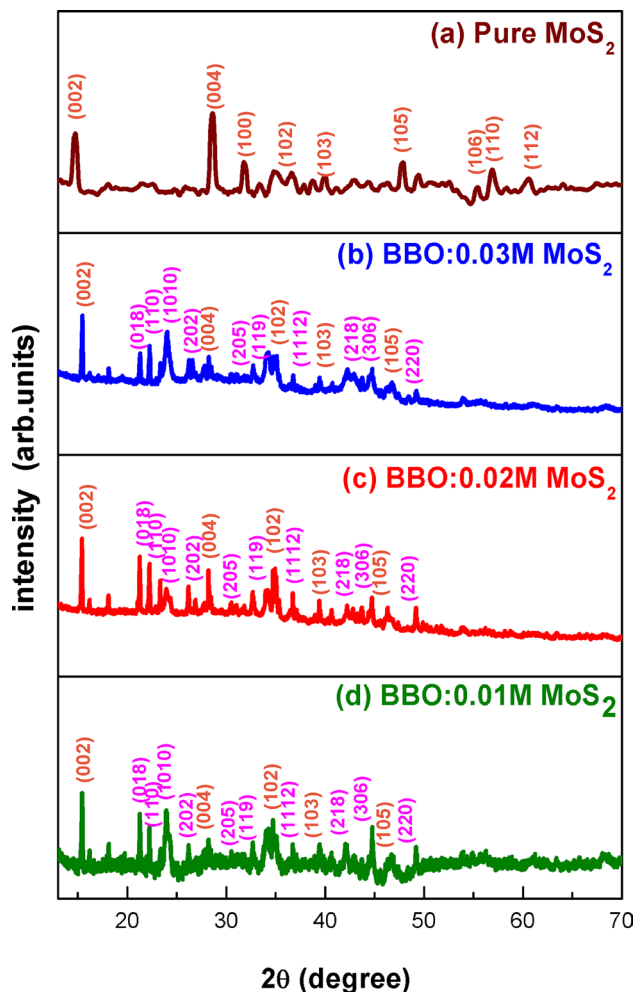


Fig. 1 XRD patterns of (a) pure MoS₂ and barium borate decorated (b) 0.03 M MoS₂ (c) 0.02 M MoS₂ (d) 0.01 M MoS₂ composites

(220) that corresponds to the angle $2\theta = 22^\circ, 24^\circ, 26^\circ, 28^\circ, 32^\circ, 34^\circ, 37^\circ, 42^\circ, 45^\circ, 50^\circ$, respectively. Literature reports show that MoS₂ has a hexagonal shape unit cell with P6₃/mmc space group, i.e. each layer of MoS₂ is composed of S–Mo–S stacks, where a single molybdenum atom is surrounded by six sulphur atoms. Similarly, beta barium borate crystallizes in the R3c space group with rhombohedral structure having lattice constant of $a = b = 7.3 \text{ \AA}$ and $C = 39 \text{ \AA}$. In pure MoS₂, diffraction peaks of (002) and (004) are highly intense compared to the other peaks, which expose the presence of well stacked layered structure of MoS₂ [18]. As the concentration of MoS₂ varied as 0.01 M, 0.02 M and 0.03 M in the BBO: MoS₂ composite, major diffraction peak of MoS₂ (002) suffered a small shift in $2\theta (= 15^\circ)$ position due to the surfactant. However, the diffraction peaks of barium borate remain unchanged in different concentrations of MoS₂ which shows the dominance of crystalline nature of BBO in composite. Absence of any other peaks that corresponds to the by-products,

impurities or other phases of title compound signifies the formation of pure and BBO decorated MoS₂. Thus, preliminary confirmation made using XRD ascertain the formation of BBO: MoS₂ composite without destructing its original crystal structure.

4 Molecular vibrations analysis: Raman spectroscopy

Molecular structural arrangement of MoS₂ and barium borate decorated on different concentration of MoS₂ was studied by Raman spectroscopy in the spectral range of 100–1500 cm⁻¹. Two characteristic peaks were observed at low range for pure MoS₂ at 379 cm⁻¹ and 407 cm⁻¹ which corresponds to in-plane E_{2g}^1 and out-of-plane A_{1g} vibrational modes of the MoS₂. At in-plane vibration mode, sulphur S atoms vibrate in one direction and molybdenum Mo atom vibrates in another direction and at out-of-plane vibration mode, sulphur S atoms only vibrate in a vertical direction. It is known that with increase in number of single layers the A_{1g} mode shifts to higher frequencies and the E_{2g}^1 mode shifts to lower frequencies, i.e. difference between the corresponding peak frequency of A_{1g} and E_{2g}^1 increases as a function of numbers of layer. Earlier report on Raman spectrum of MoS₂ exposes that single layer of MoS₂ shows A_{1g} and E_{2g}^1 peaks at 405 cm⁻¹ and 384 cm⁻¹ [19]. Here in the present case, Raman spectrum of pure MoS₂ prepared by hydrothermal method shows A_{1g} and E_{2g}^1 peaks at 407 cm⁻¹ and 379 cm⁻¹, respectively. Difference in the peak position (Fig. 2) of pure MoS₂ with earlier report [19] and shift in A_{1g} peak towards higher wavenumber and E_{2g}^1 peak towards lower wavenumber can be attributed to the change in layer number obtained in surfactant assisted hydrothermal method. Also in the Raman spectrum of composites, A_{1g} mode encountered a shift towards higher wavenumber 419 cm⁻¹ and it represents the increase in layer thickness during composite formation, while E_{2g}^1 mode exhibited relatively weak vibration due to the thermal exfoliation at 400 °C and it corresponds to the increase in crystal defect states [20]. Further the dominance of BBO, suppressed the E_{2g}^1 and E_g^1 vibrational peaks of MoS₂. Although minimal variations in intensity of BBO were observed in nanocomposites, Raman spectrum was not altered significantly due to dominance of BBO vibrations and lower concentration of MoS₂. The notable characteristic peaks of BBO in nanocomposites are 493 and 602 cm⁻¹ (intra-ring bending vibration of BO₃), 703 cm⁻¹ (symmetric breathing vibration of B–O), 812 cm⁻¹ (symmetric B–O–B bridge bond vibration), 913, 996, 1115 (asymmetric stretching of BO₄) and 1264 cm⁻¹

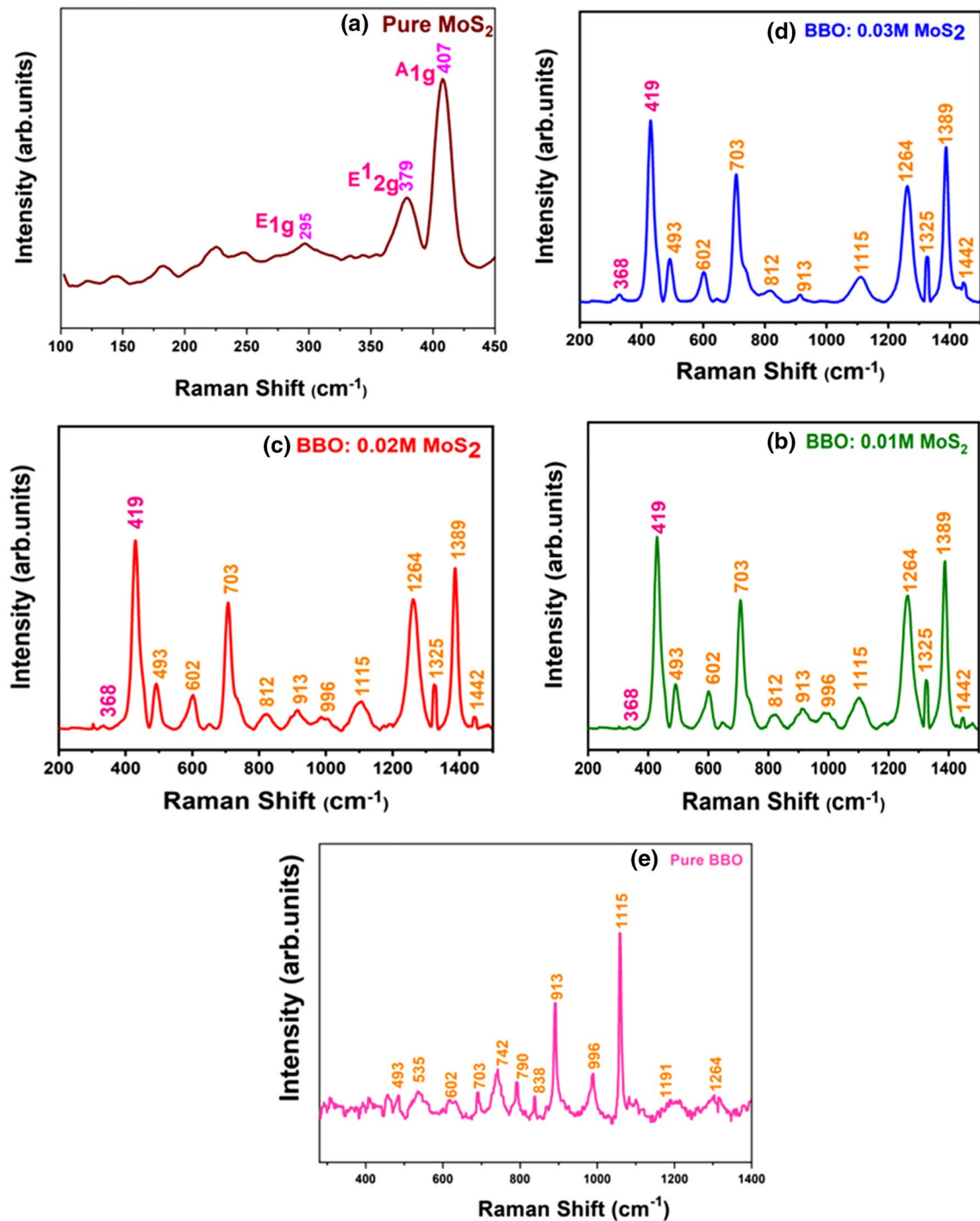
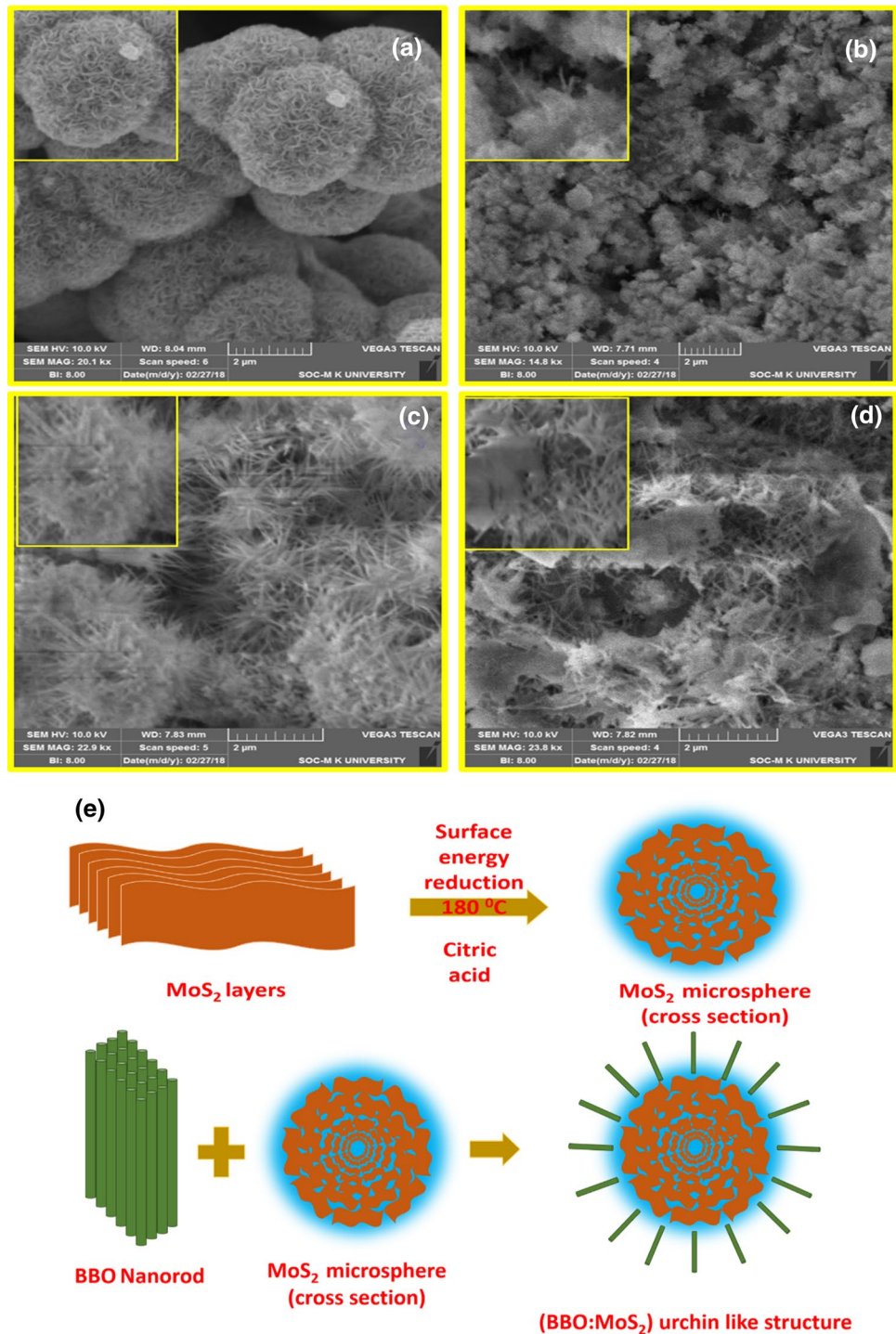


Fig. 2 Raman spectra of (a) pure MoS₂ and barium borate decorated (b) 0.03 M MoS₂ (c) 0.02 M MoS₂ (d) 0.01 M MoS₂ composites (e) pure BBO [17]

(symmetric vibration of B₂-O₃ group), 1325, 1389 and 1442 cm⁻¹ (stretching vibration of the terminal B-O₃). It is clearly inferred that almost all the BBO vibrations are observed in the Raman spectra of nanocomposite along with the characteristic vibrations of MoS₂, thereby

confirming the formation of BBO: MoS₂ composite without any molecular defects.

Fig. 3 Morphological images of (a) Pure MoS_2 and barium borate decorated (b) 0.03 M MoS_2 (c) 0.02 M MoS_2 (d) 0.01 M MoS_2 and (e) schematic of the pictorial growth mechanism of pure MoS_2 and barium borate on MoS_2



5 Morphological analysis: SEM

The morphological images (Fig. 3) of pure and barium borate decorated MoS_2 were recorded using a scanning electron microscope (VEGA3 TESCAN). It can be clearly visualized that MoS_2 appears like a microsphere (Fig. 3a), while BBO decorating on MoS_2 microsphere appears like an urchin structure (Fig. 3b–d). In general, MoS_2 is a 2D

layered structure material in which molybdenum (Mo^+) and sulphur (S^-) atoms are assembled in a same plane to form themselves like nanosheets. In the present case, the possible formation of MoS_2 microsphere is due to the reduction of surface energy induced by the citric acid surfactant. Here, the surfactant (citric acid) act as a reducing agent, which easily reduces Mo^{6+} to Mo^{4+} , and then bond with S^{2-} [16]. Formed MoS_2 nanosheets roll themselves

together and aggregate into marigold flower like structure with average diameter of 2.68 μm through an electrostatic interaction between the nucleates of MoS_2 and citric acid. In the BBO: MoS_2 nanocomposites, with increase in MoS_2 concentration, transformation of microsphere (pure and 0.03 M MoS_2) to urchin (0.02 M MoS_2) to layer (0.01 M MoS_2) like morphology was witnessed. Earlier report shows that BBO grows as elongated nanorod due to the prolonged heating at 120 $^\circ\text{C}$ for 24 h [17]. In the BBO: MoS_2 composite, BBO grow themselves as nanorod with an average length of 0.92 μm (0.03 M MoS_2), 1.21 μm (0.02 M MoS_2), 1.32 μm (0.01 M MoS_2) upon the MoS_2 microspheres. The possible mechanism is that barium nucleate themselves upon the surface of readily available MoS_2 microspheres and form themselves as BBO nucleates which then grow into BBO nanorods under prolonged heating (Fig. 3e). At higher concentration (0.03 M) of MoS_2 , slightly agglomerated MoS_2 microsphere (due to thermal shocking applied at 400 $^\circ\text{C}$ for composite formation) with traces of smaller length BBO rods was observed confirming the transition of microsphere to urchin-like structure. Interestingly at 0.02 M MoS_2 , these BBO nanorods strongly embed on the surface of the MoS_2 microsphere to appear like an urchin structure (0.02 M MoS_2) with a diameter 1.69 μm (MoS_2 microsphere) and length 1.2 μm (barium borate nanorod) (Fig. 3c). At lower concentration (0.01 M) of MoS_2 , due to higher density of BBO nanorods the microspheres structure got destructed to form a sheet-like structure. This mainly arises because interlayer distance of MoS_2 in microsphere morphology was increased due to the incorporation of barium borate. Thus, the formation of peculiar urchin structure BBO: MoS_2 nanocomposite can be advantageous for laser applications due to the availability of spherical MoS_2 (higher stability) along with 1D BBO rods (confined optical excitations).

6 Linear optical studies: absorption spectroscopy

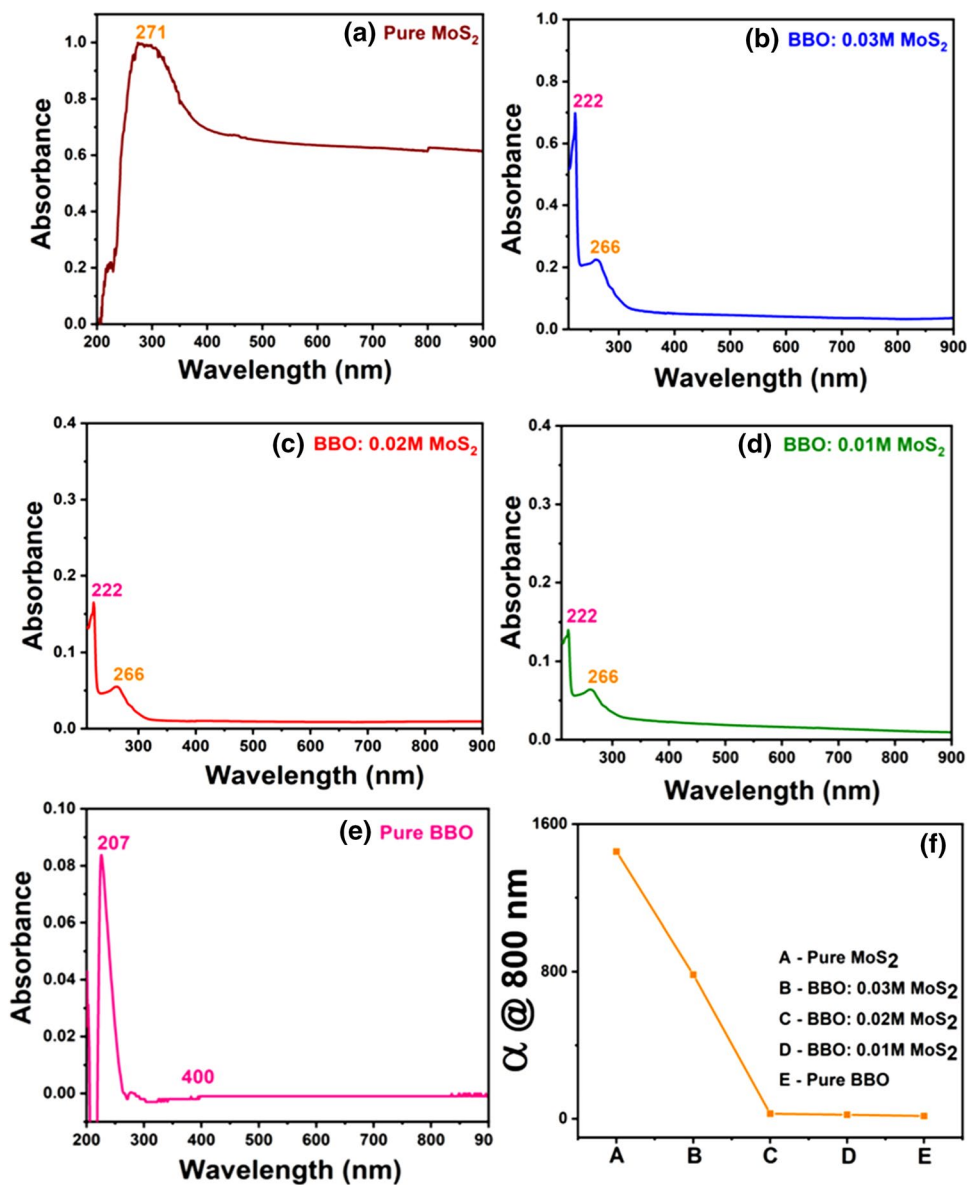
The recorded linear optical absorption spectrum (Fig. 4) of pure and BBO decorated MoS_2 was studied by UV-Visible spectrophotometer (ELICO-SL159) in the range of 190–800 nm. The absorption maximum of MoS_2 was observed at 271 nm with absorption edge at 350 nm, which can be attributed to the excitonic features of MoS_2 microspheres. It has a strong absorption at visible to near infrared (400–800 nm) region due to the direct transition of excitons at K point of Brillouin zone. In BBO: MoS_2 composite, the absorption maximum of MoS_2 was blue shifted to 266 nm from 271 nm when decorating the barium borate nanoparticles upon MoS_2 . Literature reports reveal that absorption maxima of beta barium borate nanorods

were found at 201 nm and exhibit complete transparency throughout the visible region [14]. Here, the absorption peak of barium borate was observed at 222 nm which was strongly red shifted compared to its bulk form due to the influence of 1D morphology and MoS_2 interaction. Also, in the composite the intensity of absorbance increases with increase in the concentration of MoS_2 . The absorbance of BBO: MoS_2 composite is found to be low throughout the visible spectral region when compared with pure MoS_2 . This hypochromic shift in the absorbance of BBO: MoS_2 clearly shows that the high linear transmittance (transparent properties) of composite arises from the contribution of highly transparent barium borate. Hyperchromic shift in absorption peak was due to the overlapping of transition states between highly transparent BBO and strongly absorbing MoS_2 in visible region. Similar observation was made for the BBO: rGO nanocomposites in which the maximum absorption peak of BBO (207 nm) was red shifted due to the restored electronic conjugation within the sheets of rGO and inclusion of barium borate [21]. Importantly, the visible and infrared region absorption of MoS_2 was greatly reduced due to the decorated transparent barium borate. Thus, the change in elemental composition along with morphology has created minor alterations in band structure which open up the possibility to explore different nonlinear absorption processes in the materials. The linear absorption coefficient of nanocomposites @ 800 nm (Fig. 4f) showed a nonlinear dependence at lower concentration of MoS_2 which can be attributed to the dominance of BBO and change in morphology of composite. Here, BBO: (0.01 and 0.02 M) MoS_2 composite possesses high linear transmittance in NIR region (800 nm) imitating the absorbance pattern of BBO, while BBO:0.03 M MoS_2 composite shows high absorbance as like pure MoS_2 . It is also to be admitted that as hydrothermal method was adopted, control of MoS_2 concentration was quiet challenging. Thus due to the inclusion of transparent BBO, the linear transmittance suffered a considerable change and among the samples, BBO: (0.01 M MoS_2) composite possess high linear transmittance in NIR region (800 nm) making them superior for optical limiting action in the NIR spectral region.

7 Nonlinear optical studies: Z-Scan

Z-Scan technique is an effective tool to understand the characteristic features of nonlinear interactions and measure third-order nonlinear optical properties such as nonlinear absorption (NLA) and nonlinear refraction (NLR) coefficients of the materials. Open-aperture and closed-aperture Z-scan patterns were recorded for both pure and BBO: MoS_2 nanocomposites using femtosecond laser pulses (Ti: Sapphire, 800 nm, 150 fs, 80 MHz)

Fig. 4 Absorbance spectra of (a) pure MoS₂ and barium borate decorated (b) 0.03 M MoS₂ (c) 0.02 M MoS₂ (d) 0.01M MoS₂ composites (e) pure barium borate [17] and (f) linear absorbance @800 nm



as an excitation source. Since the experiments were performed using high repetition rate pulses, thermal effects will also contribute to the observed NLO processes. In the typical experiment, the prepared samples were individually dispersed in ethylene glycol having 65% linear transmittance and were taken in a 1-mm-thick cuvette. The transmittance was measured as a function of sample position by moving the sample along the focussed beam. A graph was drawn between normalized transmittance (T_n) and position (Z) of the sample. A theoretical fit for the obtained experimental data was made using a nPA equation as proposed by Sheik-Bahae et al. [22]. The

obtained open-aperture and closed-aperture patterns are shown in Figs. 5 and 7, respectively, in which solid lines represent the theoretical fits and the scattered data (circles) denote the experimental data. The following formulae from Sheik-Bahae formalism are used to estimate the normalized transmittance (nPA) of theoretically and are given by [23],

$$\text{For an open aperture, } T_{nPA} = \frac{1}{\left[1 + (n-1)\beta_n L \left(\frac{I_0}{1 + \left(\frac{z}{z_0}\right)^2} \right)^{n-1} \right]^{\frac{1}{n-1}}}$$

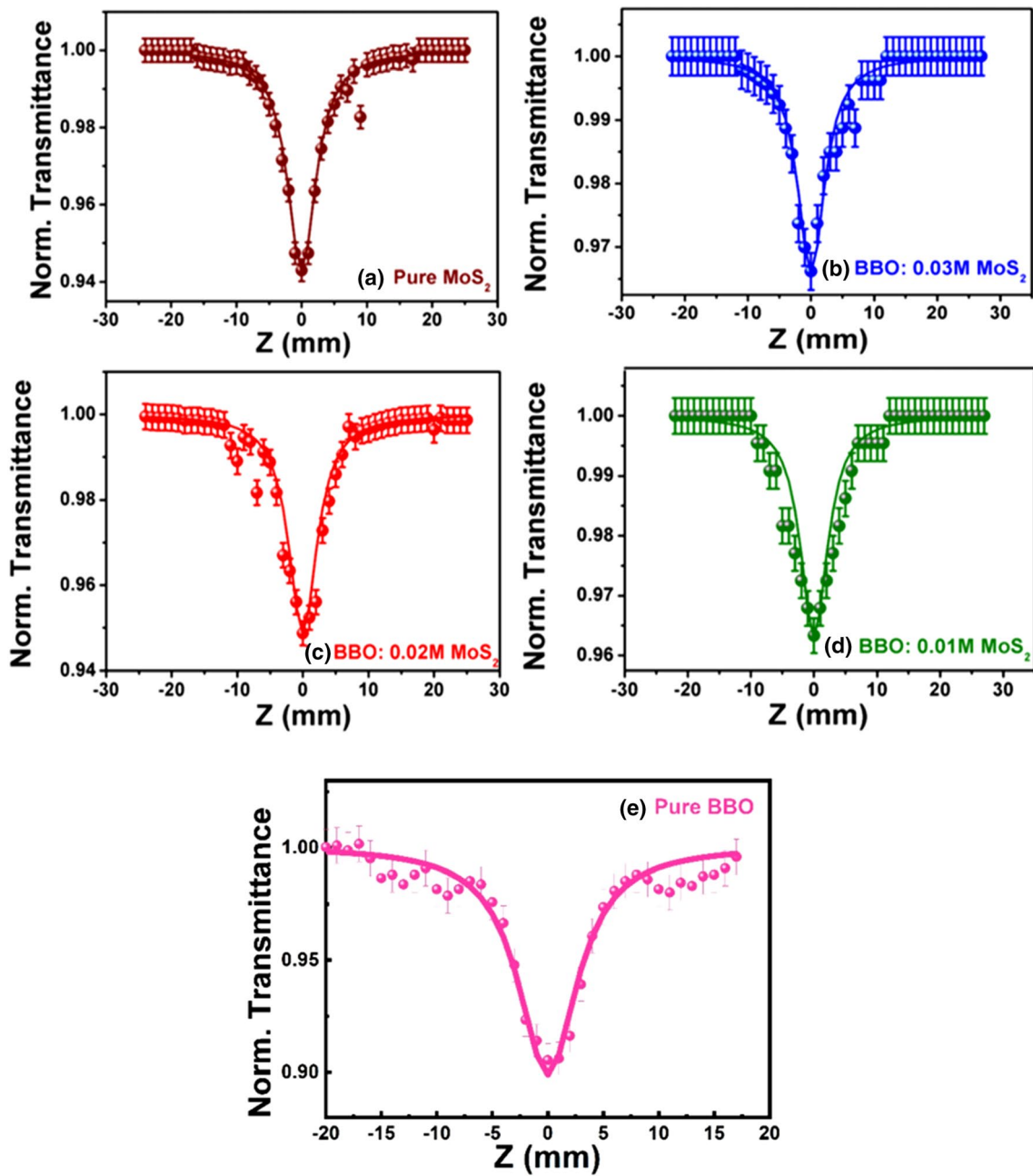


Fig. 5 Open-aperture Z-scan data of (a) pure MoS₂ and barium borate decorated (b) 0.03 M MoS₂ (c) 0.02 M MoS₂ (d) 0.01 M MoS₂ composites, (e) pure barium borate [17]. The excitation was with

femtosecond pulses (800 nm, 150 fs, 80 MHz) with a peak intensity of 295 MW/cm². Symbols are experimental data points, and the solid lines are theoretical fits to the data

For the closed aperture, $T_{CA} = 1 \pm \frac{4\Delta\theta\left(\frac{z}{z_0}\right)}{\left[1+\left(\frac{z}{z_0}\right)^2\right]\left[9+\left(\frac{z}{z_0}\right)^2\right]}$ where β_n is the nonlinear absorption coefficient with n as order of the nonlinear absorption ($n = 1, 2, 3, \dots$), L is the effective thickness of the sample, I_0 is the incident intensity of the laser beam at the focal point, $\left(z_0 = \frac{\pi\omega_0^2}{\lambda}\right)$

is the Rayleigh length, and $\Delta\theta$ is the phase distortion of the transmitted laser beam due to refraction.

Open-aperture Z-scan technique was used to evaluate the nonlinear absorption of the material. Narrow and sharp valley pattern observed for pure and BBO decorated MoS₂ demonstrates the occurrence of reverse saturable absorption (RSA) (Fig. 5), where the transmittance of light decreased with increase in light intensity. Theoretical fits made on the experimental data

revealed the observed reverse saturable absorption can be ascribed due to two-photon absorption (2PA) and three-photon absorption (3PA) processes. It is interesting to observe that the experimental data of both pure MoS₂ and BBO: MoS₂ composite show best fit for both 2PA and 3PA equations. In the UV–Visible pattern, pure MoS₂ possesses absorption in visible and NIR region and thus avails electronic states for 1PA (800 nm, 1.55 eV) and 3PA (271 nm, 4.58 eV) processes directly. The literature shows the indirect bandgap value of MoS₂ changes with the number of layers, and band structure calculated between K and Γ points in the Brillouin zone can be used to display the possible transitions responsible for non-linear absorption [1]. Strong absorption in the wavelength of excitation (800 nm) suggests that the observed 3PA is most likely to be a sequential (1PA + ESA induced genuine 2PA) process rather than simultaneous (genuine 3PA) process. The possible electronic transition involved in observed sequential 3PA of MoS₂ (Fig. 6a) is as follows: under IR excitation, an electron absorbs a photon from the ground state (E_0) and excite themselves to first excited state (E_1 , 1.55 eV—inter band transition of MoS₂ between ground state and $1s_A$ state). Then, it simultaneously absorbs two photons to transit themselves to the third excited state (E_3 , 4.57 eV—quasi continuum state of MoS₂) leading to the possible sequential 3PA through 1PA + genuine 2PA process.

However, in the BBO: MoS₂ (Fig. 6b) composite due to the dominance of BBO, the material shows negligible absorption in NIR region. Earlier reports show that non-linear absorption in pure BBO nanostructures excited with similar laser pulses exhibited 2PA process [17]. Unlike pure MoS₂, due to the non-availability of near-resonant state of excitation, BBO: MoS₂ composite undergoes genuine

3PA process (via virtual states and because sufficient peak intensity is available in the femtosecond pulses used). Here, the possible electronic transition responsible for genuine 3PA is as follows: an electron in ground state (E_0) absorbs three photons simultaneously to transit themselves to available first excited state (E_3 , 4.66 eV—quasi continuum state of MoS₂) through genuine 3PA. In pure barium borate excited under similar laser (ultrashort pulse) excitation, the electrons in the lowest state could possibly be excited to the self-trapped exciton state by simultaneously absorbing two photons leading to genuine 2PA (Fig. 6c) [17]. Therefore, alteration in band structure resulted in sequential 3PA in MoS₂ and genuine 3PA in BBO: MoS₂. It is be mentioned here that to resolve and confirm the postulates of proposed mechanism, intensity-dependent Z-scan studies and transient absorption spectroscopic studies with femtosecond pulses are absolutely necessary. The estimated 2PA and 3PA absorption coefficients (β_{2PA} and γ_{3PA}) of both pure and BBO: MoS₂ composite from the theoretical fits are summarized in Table 1. Considering the possible factors that can affect the experimental data like focal spot size, peak intensities, fitting procedures and calibration of intensity, the error value was fixed as $\sim 5\%$ for the chosen experimental condition. Similar error values are reported elsewhere in the literature and as the error correction is employed uniformly and the change in non-linearity was directly compared. The error bars indicated in Figs. 5, 7 and 8 are indicative of possible inaccuracies in the data collection. We anticipate the maximum error resulting from such factors in the evaluated values of NLO coefficients to be 5%. Similar reports on Z-scan measurement with error values are already available in the literature. It is interesting to observe that BBO: MoS₂ composite possesses higher 3PA and 2PA coefficients than pure MoS₂.

Fig. 6 Energy-level diagram demonstrating sequential 3PA for (a) Pure MoS₂, (b) genuine 2PA-induced ESA for BBO: MoS₂ nanocomposite and (c) genuine 2PA for BBO [17]. Solid lines represent the available energy states and possible electron transitions during laser excitation. And virtual states are represented in dotted lines

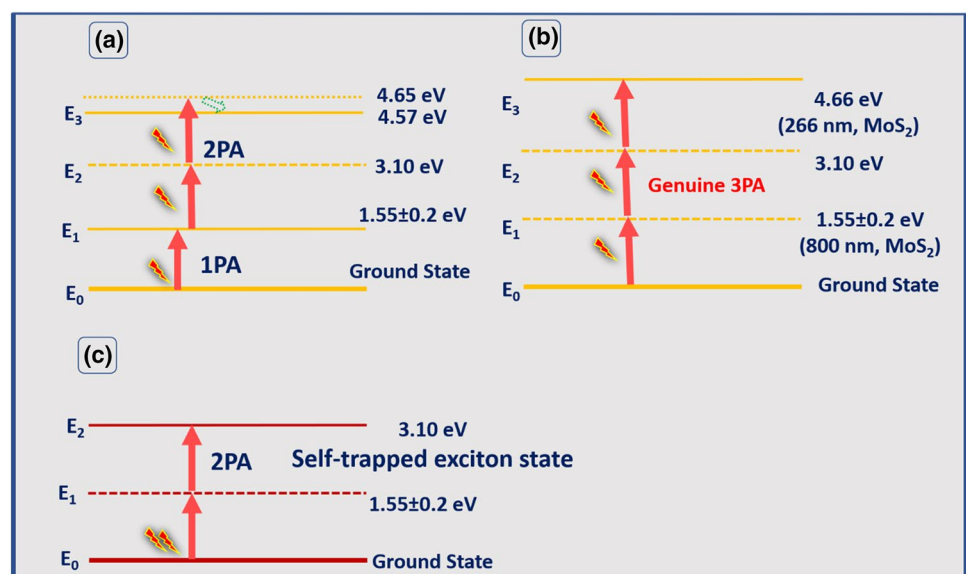


Table 1 2PA & 3PA-induced NLO coefficients of Pure and BBO: MoS₂ composites

S. no.	Sample name	Femtosecond laser (800 nm, 150 fs, 80 MHz)			
		Nonlinear absorption coefficient	Nonlinear refractive index (n_2) $\times 10^{-17}$ m ² /W	NLO susceptibility ($\chi^{(3)}$) $\times 10^{-19}$ m ² /V ²	Onset optical limiting threshold $\mu\text{J}/\text{cm}^2$
1	MoS ₂	$0.15 \pm 0.01 \times 10^{-21}$ m ³ /W ² (3PA) $5.45 \pm 0.27 \times 10^{-10}$ m/W (2PA)	2.28 ± 0.11	2.84 ± 0.14	0.16 ± 0.1
2	BBO:0.03 M MoS ₂	$0.49 \pm 0.02 \times 10^{-21}$ m ³ /W ² (3PA) $5.26 \pm 0.26 \times 10^{-10}$ m/W (2PA)	-8.29 ± 0.41	9.91 ± 0.49	0.25 ± 0.1
3	BBO:0.02 M MoS ₂	$2.12 \pm 0.11 \times 10^{-21}$ m ³ /W ² (3PA) $5.73 \pm 0.28 \times 10^{-10}$ m/W (2PA)	-11.1 ± 0.55	14.0 ± 0.7	0.11 ± 0.1
4	BBO:0.01 M MoS ₂	$0.10 \pm 0.01 \times 10^{-21}$ m ³ /W ² (3PA) $5.01 \pm 0.25 \times 10^{-10}$ m/W (2PA)	-2.05 ± 0.10	3.09 ± 0.15	0.48 ± 0.2
5	BBO [17]	$6.50 \pm 0.32 \times 10^{-10}$ m/W (2PA)	-1.00 ± 0.05	0.98 ± 0.05	0.39 ± 0.2

However, 2PA coefficient of pure BBO is higher than bare MoS₂ and BBO: MoS₂ composite. In the composite, BBO: (0.02 M) MoS₂ composite possesses higher value due to various factors like urchin morphology with microsphere structure of MoS₂, 1D structure of barium borate and overlapping the transition states of barium borate and MoS₂.

Nonlinear refractive index (n_2) of the material was obtained from the closed-aperture Z-scan method. Pure MoS₂ shows the valley followed by a peak pattern (Fig. 7a) indicating self-focusing behaviour with positive nonlinear refractive index. Self-focusing behaviour switches to the self-defocusing (peak followed by a valley as shown in Figs. 7b–d) nature as the sign is reversal and is associated with negative nonlinear refractive index due to the inclusion of barium borate in BBO: MoS₂ composite. It is fascinating to note that the sign of nonlinear refractive index switches from negative to positive. Earlier reports show that pure barium borate nanoparticles exhibit negative nonlinear refraction [17]. In this experiment, thermo-optic effect cannot be completely neglected because high repetition rate of 80 MHz was used in excitation source, and thus, the observed nonlinear refraction can be explained through thermal lens model. Under 800 nm, 150 fs, 80 MHz laser excitation, MoS₂ showed self-focusing effect (positive nonlinear refraction leading to converging of laser), while BBO exhibited self-defocusing effect (negative nonlinear refraction leading to diverging of laser). Thus, the observed self-defocusing effect in BBO: MoS₂ composite clearly exposed the dominance of BBO in the complex system. A change in refractive index as function of temperature induced during laser excitation creates two different lenses in MoS₂ (Convex lens like) and BBO (Concave lens like). Therefore, due to dominance of BBO in BBO: MoS₂ composite, a change in the sign of nonlinear refractive index was observed. The estimated nonlinear

refractive index (n_2) of pure and barium borate decorated MoS₂ is given in Table 1. Nonlinear refractive index of BBO: MoS₂ was higher than the pure MoS₂ and BBO. The superiority of MoS₂ in nonlinear refraction arises predominantly from high thermal conductivity capability of MoS₂ ($2.3 \text{ Wm}^{-1}\text{K}^{-1}$ for few layer MoS₂) [24] with which it can transform the induced local heating during excitation (high repetition rate) resulting in stronger thermal lens model. Also, third-order nonlinear optical susceptibility of the all pure and compositing materials was estimated to be $2.84 \times 10^{-19} \text{ m}^2/\text{V}^2$, $9.91 \times 10^{-19} \text{ m}^2/\text{V}^2$, $14.0 \times 10^{-19} \text{ m}^2/\text{V}^2$ and $3.08 \times 10^{-19} \text{ m}^2/\text{V}^2$ for pure MoS₂, BBO: 0.03 M MoS₂, BBO: 0.02 M MoS₂, BBO: 0.01 M MoS₂, respectively. Among all the samples, BBO decorated on 0.02 M concentration of MoS₂ possesses a higher NLO coefficient such as nonlinear absorption ($\gamma_{3PA} = 2.12 \times 10^{-21} \text{ m}^3/\text{W}^2$, $\beta_{2PA} = 5.73 \times 10^{-10} \text{ m}/\text{W}$), nonlinear refractive index ($n_2 = -11.1 \times 10^{-17} \text{ m}^2/\text{W}$) and third-order NLO susceptibility ($\chi^{(3)} = 14 \times 10^{-19} \text{ m}^2/\text{V}^2$ due to urchin morphology with the availability of spherical MoS₂ (higher thermal stability against high repetition rate femtosecond laser) and 1D BBO rod (confined optical excitations).

Optical limiting curves were extracted from the open-aperture Z-scan data and using a fluence equation ($4\sqrt{\ln 2} E_{in}/\pi^{3/2} \omega(z)^2$) are shown in Fig. 8. The graphs represent a nonlinear pattern in which output transmittance changes nonlinearly with input fluence. And the deviation point at which the nonlinearity begins is termed as onset optical limiting behaviour of the material. Thus, the prepared materials can act as optical limiters that are transmitting a low intensity while blocking the high input intensity. Here, the observed optical limiting arises mainly due to nonlinear absorption (2PA and 3PA) and nonlinear refraction (self-defocusing).

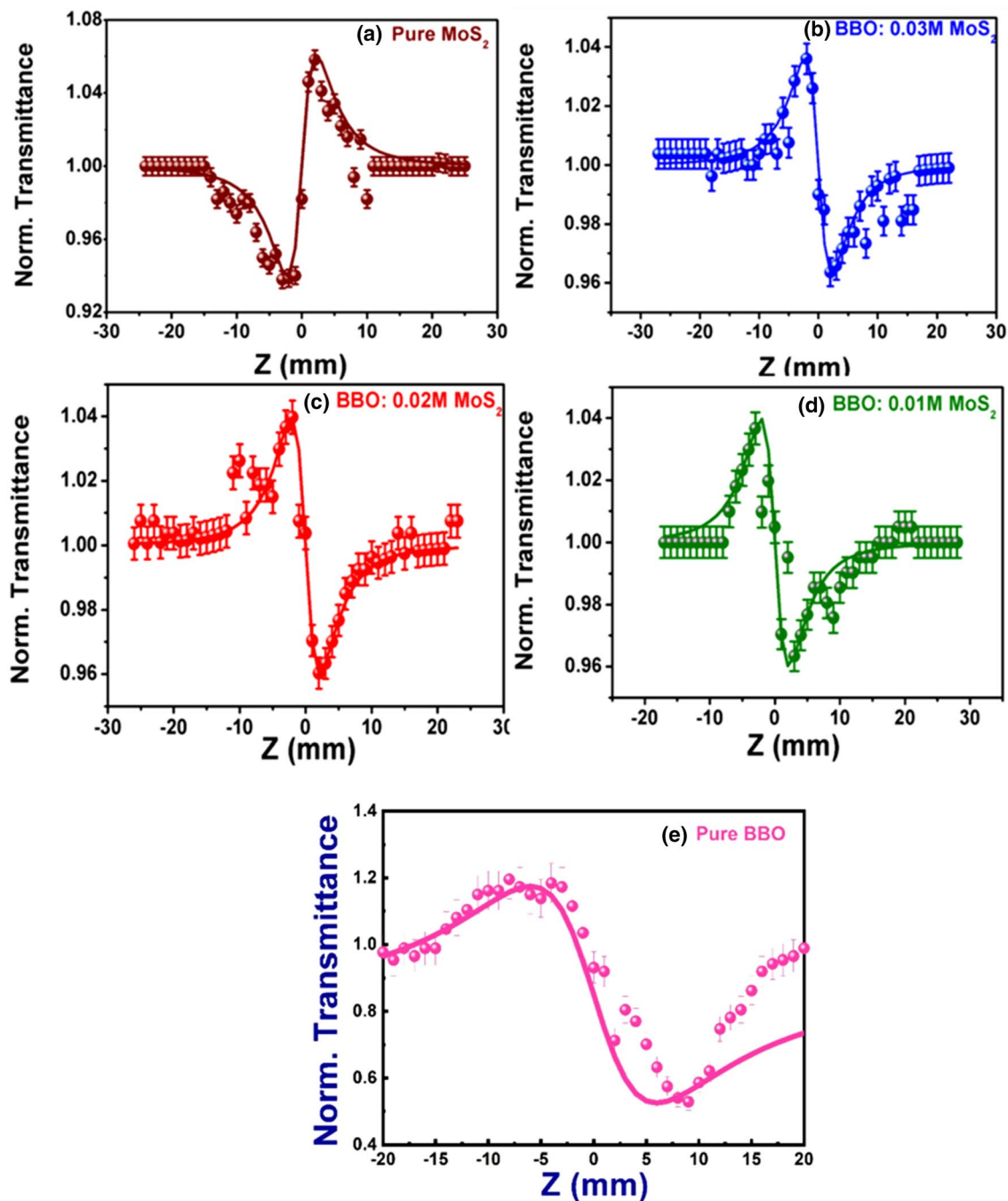


Fig. 7 Closed-aperture pattern of (a) pure MoS_2 and barium borate decorated (b) 0.03 M MoS_2 (c) 0.02 M MoS_2 (d) 0.01 M MoS_2 composites (e) pure barium borate [17]. The excitation was with fem-

tosecond pulses (800 nm, 150 fs, 80 MHz) with a peak intensity of 295 MW/cm^2 . Symbols are experimental data points, and the solid lines are theoretical fits to the data

Among the samples, onset optical limiting threshold of BBO decorate on 0.02 M of MoS_2 ($0.11 \mu\text{J/cm}^2$) and pure MoS_2 ($0.16 \mu\text{J/cm}^2$) is lower than other composites. Further, these values are found to be lower than the reported values of other materials like copper niobite

($0.21 \mu\text{J/cm}^2$) [23], bismuth ($2.16 \mu\text{J/cm}^2$) [25] and zinc oxide ($128 \mu\text{J/cm}^2$) [26] excited with similar laser. As optical limiting threshold is an important parameter in OL device fabrication, an extrapolation of OA data was done and the input fluence at which normalized

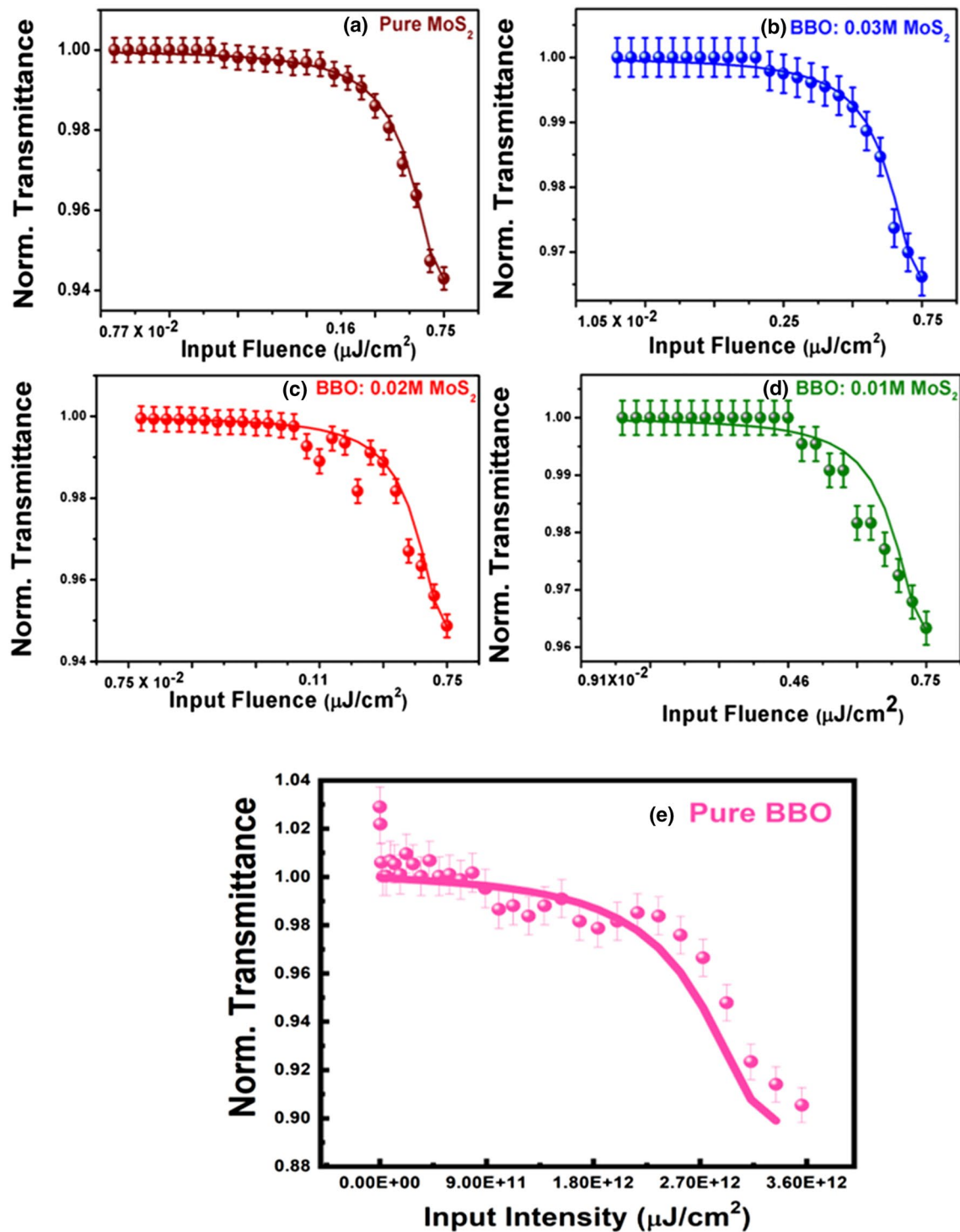


Fig. 8 Optical limiting curves of (a) pure MoS₂ and barium borate decorated (b) 0.03 M MoS₂ (c) 0.02 M MoS₂ (d) 0.01 M MoS₂ composites (e) pure barium borate [17]. The excitation was with fem-

tosecond pulses (800 nm, 150 fs, 80 MHz) with a peak intensity of 295 MW/cm². Symbols are experimental data points, and the solid lines are theoretical fits to the data

transmittance becomes 50% was estimated to be few tens of m J/cm² for all samples. A comparison of NLO coefficients of various composites excited with similar laser is presented in Table 2. Synergetic effects of

2PA/3PA- and self-defocusing-induced optical limiting with lower onset optical limiting threshold and higher NLO coefficients make BBO: MoS₂ nanocomposite

Table 2 Summary of the NLO coefficients of various composites excited with Ti: Sapphire laser pulses (at 800 nm)

Sample name	Nonlinear absorption coefficient (β)	Nonlinear refractive index (n_2)	Onset optical limiting threshold
MoS ₂ -TiO ₂ [12]	1.77×10^{-10} m/W	–	22.3 mJ/cm ²
MoS ₂ -PMMA [13]	1.77×10^{-10} m/W	–	21.5 μ J/cm ²
Fe-Ag [27]	5.40×10^{-15} m/W	–	–
Ag-Polymer [28]	18×10^{-12} m/W	2.8×10^{-11} cm ² /W	38 mJ/cm ²
CdFe ₂ O ₄ -rGO [29]	21×10^{-12} m/W	4.3×10^{-16} m ² /W	0.65 μ J/cm ²

preferable choice for laser safety devices for ultrashort IR laser pulses.

8 Conclusions

In summary, MoS₂ microspheres and barium borate nanorods decorated MoS₂ microspheres with urchin-like structure were prepared by hydrothermal method. Both MoS₂ and BBO: MoS₂ exhibited reverse saturable absorption due to 2PA and 3PA. Sequential and genuine 3PA processes occurred for pure MoS₂ and barium borate decorated MoS₂, respectively, and was confirmed with the aid of available energy states from absorption spectrum. Nonlinear refraction of pure MoS₂ depicted self-focusing behaviour which switched to self-defocusing in BBO: MoS₂ composite. The dynamic processes involved in the observed nonlinearity of MoS₂, BBO and BBO: MoS₂ composites are identified as sequential 3PA [$E_0 \xrightarrow{1PA} E_1$ (interband transition, 1.55 eV) $\xrightarrow{\text{genuine2PA}} E_3$ (continuum state, 4.58 eV)], genuine 2PA [$E_0 E_0 \xrightarrow{\text{genuine2PA}} E_2 E_2$ (self-trapped exciton state of BBO, 3.10 eV)] and genuine 3PA [$E_0 \xrightarrow{\text{genuine3PA}} E_3 E_3$ (continuum state of MoS₂, 4.66 eV)], respectively. The tunability of NLO properties achieved by varying the content of MoS₂ and BBO: (0.02 M) MoS₂ nanocomposite possess high nonlinear absorption coefficient (β), nonlinear refractive index (n_2) and third-order NLO susceptibility ($\chi^{(3)} = 14.0 \times 10^{-19}$ m²/V²). Multiphoton ($nPA, n = 2,3$)-induced optical limiting action in femtosecond IR domain was demonstrated both for pure MoS₂ and BBO: MoS₂ composite and lower onset optical limiting threshold of $\sim \mu$ J/cm² opens the possibility of utilizing BBO:(0.02 M) MoS₂ for laser safety devices and as an energy stabilizer in microjoule lasers based surgery.

Compliance with ethical standards

Conflicts of interest There are no conflicts to declare.

References

- Khudyakov DV, Borodkin AA, Lobach AS, Mazin DD, Vartapevov SK (2019) Optical nonlinear absorption of a few-layer MoS₂ under green femtosecond excitation. *Appl Phys B* 125:73
- Ma C, Wang C, Gao B, Adams J, Wu G, Zhang H (2019) Recent progress in ultrafast lasers based on 2D materials as a saturable absorber. *Appl Phys Rev* 6:041304
- He J, Tao L, Zhang H, Zhou B, Li J (2019) Emerging 2D materials beyond graphene for ultrafast pulse generation in fiber lasers. *Nanoscale* 11:2577–2593
- Kumar NA, Dar MA, Gul R, Baek JB (2015) Graphene and molybdenum disulfide hybrids: synthesis and applications. *Mater Today* 18:286–298
- Cao J, Zhou J, Zhang Y, Liu X (2017) A clean and facile synthesis strategy of MoS₂ nanosheets grown on multi-wall CNTs for enhanced hydrogen evolution reaction performance. *Sci Rep* 7:8825
- Wen X, Gong Z, Li D (2019) Nonlinear optics of two dimensional transition metal dichalcogenides. *Info Mat* 1:317–337
- Edvinsson T (2018) Optical quantum confinement and photocatalytic properties in two-, one- and zero-dimensional nanostructures. *R Soc Open Sci* 5:180387
- Lezama IG, Arora A, Ubaldini A, Barreateau C, Giannini E, Potemski M, Morpurgo AF (2015) Indirect-to-direct band gap crossover in few-layer MoTe₂. *Nano Lett* 15:2336–2342
- He M, Quan C, He C, Huang Y, Zhu L, Yao Z, Xu X (2017) Enhanced nonlinear saturable absorption of MoS₂/graphene nanocomposite films. *J Phys Chem C* 121:27147–27153
- Wei R, Zhan H, Hu Z, Qiao T, He X, Guo Q, Qiu J (2016) Ultra-broadband nonlinear saturable absorption of high-yield MoS₂ nanosheets. *Nanotechnology* 27:305203
- Qu B, Ouyang Q, Yu X, Luo W, Qi L, Chen Y (2015) Nonlinear absorption, nonlinear scattering, and optical limiting properties of MoS₂-ZnO composite-based organic glasses. *Phys Chem Chem Phys* 17:6036–6043
- Wei R, Tian X, Hu Z, Zhang H, Qiao T, He X, Qiu J (2016) Vertically standing layered MoS₂ nanosheets on TiO₂ nanofibers for enhanced nonlinear optical property. *Opt Express* 24:25337–25344
- Liang G, Tao L, Tsang YH, Zeng L, Liu X, Li J, Wen Q (2019) Optical limiting properties of a few-layer MoS₂/PMMA composite under excitation of ultrafast laser pulses. *J Mater Chem C* 7:495–502
- Babeela C, Girisun TS, Vinitha G (2015) Optical limiting behavior of β -BaB₂O₄ nanoparticles in pulsed and continuous wave regime. *J Phys D Appl Phys* 48:065102
- Zhang Y, Lu D, Yu H, Zhang H (2019) Low-Dimensional Saturable Absorbers in the Visible Spectral Region. *Adv Opt Mater* 7:1800886

16. Feng G, Wei A, Zhao Y, Liu J (2015) Synthesis of flower-like MoS₂ nanosheets microspheres by hydrothermal method. *J Mater Sci-Mater Electron* 26:8160–8166
17. Girisun TCS, Somayaji RM, Priyadarshani N, Rao SV (2017) Femtosecond third order optical nonlinearity and optical limiting studies of (γ and β)—barium borate nanostructures. *Mater Res Bull* 87:102–108
18. Lin Hongtao, Xiaoya C, Hongling L, Min Y, Yanxing Q (2010) Hydrothermal synthesis and characterization of MoS₂ nanorods. *Mater Lett* 64:1748–1750
19. Li X, Zhu H (2015) Two-dimensional MoS₂: properties, preparation, and applications. *J Materiomics* 1:33–44
20. Lee Jung E, Jaemin J, Taeg YK, Sujin K, Seong-IK Junghyo N, Sunmin R, Ki TN, Min HL (2017) Catalytic synergy effect of MoS₂/reduced graphene oxide hybrids for a highly efficient hydrogen evolution reaction. *RSC Adv* 7:5480–5487
21. Muruganandi G, Saravanan M, Vinitha G, Raj MJ, Girisun TS (2018) Barium borate nanorod decorated reduced graphene oxide for optical power limiting applications. *Opt Mater* 75:612–618
22. Sheik-Bahae M, Said AA, Van Stryland EW (1989) High-sensitivity, single-beam n_2 measurements. *Opt Lett* 14:955–957
23. Priyadarshani N, Rao SV, Girisun TS (2016) Investigation of the femtosecond optical limiting properties of monoclinic copper niobate. *Appl Phys B* 122:256
24. Gandi AN, Schwingenschlöggl U (2016) Thermal conductivity of bulk and monolayer MoS₂. *EPL* 113:36002
25. Yang C, Shang ZJ, Wang Z, Peng H, Tang XD, Li B, Chen Y (2015) Size control of semimetal bismuth nanoparticles and the UV–visible and IR absorption spectra. *Int J Opt* 44:7–11
26. Shaik UP, Kumar PA, Krishna MG, Rao SV (2014) Morphological manipulation of the nonlinear optical response of ZnO thin films grown by thermal evaporation. *Mater Res Express* 1:046201
27. Sridharan K, Endo T, Cho SG, Kim J, Park TJ, Philip R (2013) Single step synthesis and optical limiting properties of Ni–Ag and Fe–Ag bimetallic nanoparticles. *Opt Mater* 35:860–867
28. Misra N, Rapolu M, Rao SV, Varshney L, Kumar V (2016) Nonlinear optical studies of inorganic nanoparticles–polymer nanocomposite coatings fabricated by electron beam curing. *Opt Laser Technol* 79:24–31
29. Saravanan M, Girisun TCS, Rao SV (2017) Super-paramagnetic and unusual nonlinear absorption switching behaviour of an in situ decorated CdFe₂O₄-rGO nanocomposite. *J Mater Chem C* 5:9929–9942

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