# An Observation About Conformal Points on Surfaces 

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#### Abstract

We study the existence of points on a compact oriented surface at which a symmetric bilinear two-tensor field is conformal to a Riemannian metric. We give applications to the existence of conformal points of surface diffeomorphisms and vector fields.


Keywords Conformal points • Poincaré-Hopf • Line fields

Mathematics Subject Classification Primary 53C18; Secondary 57R22

## 1 Statement of Results

### 1.1 Conformal Points

Let $\Sigma$ be a compact, oriented surface, possibly with non-empty boundary $\partial \Sigma$. Denote by $C_{1}, \ldots, C_{n}$ the boundary components of $\Sigma$ with the induced orientation. Let $\operatorname{Sym}\left(\left(T^{*} \Sigma\right)^{\otimes 2}\right) \rightarrow \Sigma$ be the bundle of symmetric bilinear tensors on $\Sigma$. Fix a Riemannian metric $g$ on $\Sigma$, that is, a positive-definite section of $\operatorname{Sym}\left(\left(T^{*} \Sigma\right)^{\otimes 2}\right) \rightarrow \Sigma$.

Definition 1.1 We say that a section $h$ of $\operatorname{Sym}\left(\left(T^{*} \Sigma\right)^{\otimes 2}\right) \rightarrow \Sigma$ is conformal to $g$ at the point $z \in \Sigma$ if there exists $c \in \mathbb{R}$ such that $h_{z}=c g_{z}$.

Motivated by [1], the goal of this note is to study the set of points

$$
\mathcal{C}(g, h) \subset \Sigma
$$

[^0]at which $h$ is conformal to $g$ (see Theorem 1.2) in order to investigate conformal points of diffeomorphisms $F: \Sigma \rightarrow \Sigma$, in which case $h=F^{*} g$ (see Theorem 1.7 and Corollary 1.8), and of vector fields (see Corollary 1.10).

Our main observation is that $\mathcal{C}(g, h)$ is the zero-set of a section $H^{a}$ in a distinguished vector bundle $E^{a} \rightarrow \Sigma$ over the surface, which we describe now. Let $\operatorname{End}(T \Sigma) \rightarrow \Sigma$ be the bundle of endomorphisms of $T \Sigma$ and let

$$
\begin{equation*}
E^{a} \subset \operatorname{End}(T \Sigma) \tag{1.1}
\end{equation*}
$$

be the subbundle of those endomorphisms which are symmetric with respect to $g$ and have zero trace. For all $z \in \Sigma$, an element of $R \in E_{z}^{a}$ has the matrix expression

$$
\left(\begin{array}{cc}
a & b \\
b & -a
\end{array}\right) \quad a, b \in \mathbb{R}
$$

with respect to a positive, orthonormal basis of $T_{z} \Sigma$. Thus any non-zero element $R \in E_{z}^{a}$ is, up to a positive scalar multiple, a reflection $R: T_{z} \Sigma \rightarrow T_{z} \Sigma$ along a line in $T_{z} \Sigma$. In particular, the $S^{1}$-bundle associated with $E^{a}$ is the bundle of unoriented lines in $T \Sigma$. This $S^{1}$-bundle is doubly covered by the bundle of oriented lines in $T \Sigma$ which, in turn, is the unit-tangent bundle of $\Sigma$, that is, the $S^{1}$-bundle associated with $T \Sigma \rightarrow \Sigma$. The above discussion shows that $E^{a}$ is an oriented plane bundle over $\Sigma$ with Euler number

$$
\begin{equation*}
e\left(E^{a}\right)=2 e(T \Sigma)=2 \chi(\Sigma) \tag{1.2}
\end{equation*}
$$

Given a symmetric bilinear two-tensor field $h$ over $\Sigma$, let $H$ be the section of $\operatorname{End}(T \Sigma)$ representing $h$ with respect to $g$, namely

$$
\begin{equation*}
g_{z}\left(u, H_{z} v\right)=h_{z}(u, v), \quad \forall z \in \Sigma, \forall u, v \in T_{z} \Sigma \tag{1.3}
\end{equation*}
$$

We denote by

$$
\begin{equation*}
H^{a}:=H-\frac{\operatorname{tr} H}{2} I \tag{1.4}
\end{equation*}
$$

the section of $E^{a}$ corresponding to the trace-free part of $H$. Here $I$ is the section of $\operatorname{End}(T \Sigma)$ such that $I_{z}$ is the identity of $T_{z} \Sigma$ for all $z \in \Sigma$.

Thus, we conclude that

$$
z \in \mathcal{C}(g, h) \quad \Longleftrightarrow \quad H_{z}^{a}=0
$$

This relationship implies that given any closed set $C \subset \Sigma$ and any field $R$ of unoriented lines on $\Sigma \backslash C$, there is a section $h$ such that $\mathcal{C}(g, h)=C$ and $R$ is induced by $h$ on $\Sigma \backslash C$, see e.g. [9, Theorem A]. Moreover, we see that, generically, $h$ has only finitely many, non-degenerate (as zeros of $H^{a}$ ) conformal points, and all of them lie in the interior of $\Sigma$. If $h$ has only finitely many (degenerate or non-degenerate) critical points,
and all of them lie in the interior of $\Sigma$, then we can use the Poincaré-Hopf Theorem for unoriented line fields on oriented surfaces with boundary to algebraically count conformal points. To give the precise statement, let us introduce some notation under the assumption that $\mathcal{C}(g, h)$ is finite and $\mathcal{C}(g, h) \subset \Sigma \backslash \partial \Sigma$. For each $z \in \mathcal{C}(g, h)$, we define

$$
\operatorname{ind}_{(g, h)}(z) \in \mathbb{Z}
$$

as the index of $z$ seen as a zero of the section $H^{a}$ of $E^{a} \rightarrow \Sigma$. Thus $\operatorname{ind}_{(g, h)}(z) \in$ $\{-1,+1\}$ if $z$ is non-degenerate. We count the elements in $\mathcal{C}(g, h)$ algebraically via the integer

$$
\begin{equation*}
[\mathcal{C}(g, h)]:=\sum_{z \in \mathcal{C}(g, h)} \operatorname{ind}_{(g, h)}(z) \in \mathbb{Z} \tag{1.5}
\end{equation*}
$$

Moreover, for every boundary component $C_{i}$ of $\Sigma$, with $i=1, \ldots, n$, we define

$$
\begin{equation*}
w_{i}(g, h) \in \mathbb{Z} \tag{1.6}
\end{equation*}
$$

as the winding number of the section $\left.H^{a}\right|_{C_{i}}$ with respect to $R^{i} \in E^{a}$, where $R^{i}(z)$ is the reflection along the line $T_{z} \partial \Sigma \subset T_{z} \Sigma$ for $z \in C_{i}$.

Theorem 1.2 Let $g$ be a Riemannian metric on a compact, oriented surface $\Sigma$. Then the following two statements hold.
(1) For any symmetric bilinear two-tensor field $h$ over $\Sigma$ such that $\mathcal{C}(g, h)$ is finite and $\mathcal{C}(g, h) \subset \Sigma \backslash \partial \Sigma$, the equality

$$
\begin{equation*}
[\mathcal{C}(g, h)]=2 \chi(\Sigma)+\sum_{i=1}^{n} w_{i}(g, h) \tag{1.7}
\end{equation*}
$$

holds, where $\chi(\Sigma)$ denotes the Euler characteristic of $\Sigma$.
(2) Let $\mathcal{C} \subset \Sigma \backslash \partial \Sigma$ be a finite set of points, $\iota: \mathcal{C} \rightarrow \mathbb{Z}$ an arbitrary function, and $w_{1}, \ldots, w_{n} \in \mathbb{Z}$ arbitrary integers satisfying

$$
\begin{equation*}
\sum_{z \in \mathcal{C}} \iota(z)=2 \chi(\Sigma)+\sum_{i=1}^{n} w_{i} . \tag{1.8}
\end{equation*}
$$

Then there exists a symmetric bilinear two-tensor field hover $\Sigma$ such that $\mathcal{C}=$ $\mathcal{C}(g, h), \iota(z)=\operatorname{ind}_{(g, h)}(z)$ for all $z \in \mathcal{C}$ and $w_{i}=w_{i}(g, h)$ for all $i=1, \ldots, n$.

Remark 1.3 For the convenience of the reader, we give a short proof of Theorem 1.2 in Sect. 2 although this can be deduced from the literature. For statement (1), we refer to $[8,14,16]$ which deal with the Poincaré-Hopf Theorem for oriented line fields on surfaces with boundary and to [11, III.2.2], [3, 12, 13] which deal with the PoincaréHopf theorem for unoriented line fields on surfaces without boundary. The version
of statement (1) has been generalized from quadratic forms to symmetric totally real $n$-forms in [6, Section 3]. For statement (2), we refer to the Extension Theorem in [10, p. 145].

Remark 1.4 Passing to the orientation double cover, Theorem 1.2 also holds for nonorientable surfaces.

Remark 1.5 In the literature, there are several formulas for the computation of the indices $\operatorname{ind}_{(g, h)}(z)$ and the winding numbers $w_{i}(g, h)$, for instance [19, Theorem 1], [4, Theorem 2.1].

We discuss now two situations where the set $\mathcal{C}(g, h)$ naturally appears.

### 1.2 Umbilical Points of Immersed Surfaces

First, let us consider a smooth immersion $\rho: \Sigma \rightarrow \mathbb{R}^{3}$. Here we take $g^{\rho}$ and $h^{\rho}$ to be the first and the second fundamental form of the immersion $\rho$, respectively, with respect to the ambient Euclidean metric. The elements of $\mathcal{C}\left(g^{\rho}, h^{\rho}\right)$ are the so-called umbilical points, namely points at which the two principal curvatures of the immersion coincide. In this case, (1.7) yields the well-known result that $\left[\mathcal{C}\left(g^{\rho}, h^{\rho}\right)\right]=2 \chi(\Sigma)$. In general, it is natural to ask which further conditions must the points $z \in \mathcal{C}\left(g^{\rho}, h^{\rho}\right)$ and their indices satisfy besides $\left[\mathcal{C}\left(g^{\rho}, h^{\rho}\right)\right]=2 \chi(\Sigma)$. For instance, Loewner's conjecture [18] asserts that ind $(z) \leq 2$ for all $z \in \mathcal{C}\left(g^{\rho}, h^{\rho}\right)$. This conjecture implies Carathéodory's conjecture [7,17], which asserts that if $\rho$ is a convex embedding (hence $\Sigma=S^{2}$ and $\chi(\Sigma)=2)$, then $\mathcal{C}\left(g^{\rho}, h^{\rho}\right)$ contains at least two points. Examples where $\mathcal{C}\left(g^{\rho}, h^{\rho}\right)$ consists exactly of two points, both having index two, are ellipsoids of revolution, where the two poles are umbilical.

Remark 1.6 In [5] the notion of $k$-roundings of immersed $m$-manifolds in $\mathbb{R}^{n}$ has been introduced so that umbilical points correspond to $k=2, m=2$ and $n=3$. If $k, m, n$ satisfy a certain relation, then $k$-roundings are generically isolated and can algebraically be counted via an index.

### 1.3 Conformal Points of a Diffeomorphism

The second situation in which $\mathcal{C}(g, h)$ naturally appears is when $h=F^{*} g$, where $F: \Sigma \rightarrow \Sigma$ is any orientation-preserving diffeomorphism of $\Sigma$. In this case, $\mathcal{C}\left(g, F^{*} g\right)$ is the set of so-called conformal points of $F$ (with respect to $g$ ). Assuming that $\mathcal{C}\left(g, F^{*} g\right)$ is finite and $\mathcal{C}\left(g, F^{*} g\right) \subset \Sigma \backslash \partial \Sigma$, we are going to give a formula for $w_{i}\left(g, F^{*} g\right)$ in terms of the behavior of $F$ at the boundary. To state the result, for $i=1, \ldots, n$ let $\nu_{i}: C_{i} \rightarrow T \Sigma$ be the outward normal at the boundary component $C_{i}$ and $\tau_{i}: C_{i} \rightarrow T \Sigma$ be the unit vector tangent to $C_{i}$ in the positive direction. The pair $\left(v_{i}, \tau_{i}\right)$ then forms a positive orthonormal frame for $g$ along $C_{i}$. We trivialize $\left.T \Sigma\right|_{\partial \Sigma}=\sqcup_{i} C_{i} \times \mathbb{R}^{2}$ using $\left(v_{i}, \tau_{i}\right)$ at $C_{i}, i=1, \ldots, n$. Since $F$ maps boundary components to boundary components (not necessarily the same) we can express $\mathrm{d} F$
in this trivialization as

$$
\left.\mathrm{d} F\right|_{C_{i}}=: N_{i}=c_{i}\left(\begin{array}{cc}
a_{i} & 0  \tag{1.9}\\
b_{i} & 1
\end{array}\right) .
$$

Here $a_{i}, c_{i}: C_{i} \rightarrow(0, \infty), b_{i}: C_{i} \rightarrow \mathbb{R}$ and $\left(a_{i}, b_{i}\right)$ is never equal to $(1,0)$ since $F$ has no conformal point on $C_{i}$ by assumption.

Theorem 1.7 For all $i=1, \ldots, n$ we have the equality

$$
\begin{equation*}
w_{i}\left(g, F^{*} g\right)=w\left(a_{i}-1, b_{i}\right) \tag{1.10}
\end{equation*}
$$

where $w\left(a_{i}-1, b_{i}\right)$ is the winding number of the curve $\left(a_{i}-1, b_{i}\right): C_{i} \cong S^{1} \rightarrow \mathbb{R}^{2} \backslash\{0\}$ around the origin.

This formula, which will be proved in Sect. 3, allows us to compute $w_{i}\left(g, F^{*} g\right)$ if we understand the behavior of $F$ at points on the boundary sufficiently well. A remarkable example of this phenomenon is illustrated by the next corollary.

Corollary 1.8 If $F: \Sigma \rightarrow \Sigma$ is the identity on the boundary and preserves an area form on $\Sigma$, then

$$
\begin{equation*}
w_{i}\left(g, F^{*} g\right)=0, \quad \forall i=1, \ldots, n \tag{1.11}
\end{equation*}
$$

It follows that for this type of diffeomorphisms

$$
\begin{equation*}
[\mathcal{C}(F)]=2 \chi(\Sigma), \tag{1.12}
\end{equation*}
$$

that is, the number of conformal points of such an $F$ is twice the Euler characteristic.
Proof By (1.10) the assertion is equivalent to showing $w_{i}\left(a_{i}-1, b_{i}\right)=0$. Since $F$ is the identity at the boundary we conclude that $\mathrm{d} F \cdot \tau_{i}=\tau_{i}$ and thus $c_{i}=1$ in (1.9). Since $F$ preserves an area form, it follows that det $N_{i}=1$, which implies that $a_{i}=1$ in (1.9). Therefore, the curve $\left(a_{i}-1, b_{i}\right)=\left(0, b_{i}\right)$ is contained in the $y$-axis and does not cross 0 . We conclude that its winding number around the origin $w\left(a_{i}-1, b_{i}\right)$ vanishes.

Remark 1.9 Equation (1.12) was proved in [1], when $\Sigma=D^{2}$, and $F$ satisfies some additional conditions, which hold, for instance, when $F$ is $C^{1}$-close to the identity,

If we linearize the property of being a conformal point for a diffeomorphism at the identity of $\Sigma$, we get a corresponding condition for conformal points of vector fields on $\Sigma$. This condition is easier phrased after reinterpreting conformality in terms of complex geometry, as we explain next.

### 1.4 Conformal Points and Complex Structures

Let $J$ be the complex structure associated with the Riemannian metric $g$ and the orientation of $\Sigma$. In other words, $J$ yields a section of $\operatorname{End}(T \Sigma)$ such that $v$ and $J_{z} v$ form a positive, orthogonal basis of $T_{z} \Sigma$ for all $z \in \Sigma$ and all $v \in T_{z} \Sigma \backslash\{0\}$. Thus, $J_{z}$ has the matrix expression

$$
\left(\begin{array}{cc}
0 & -1  \tag{1.13}\\
1 & 0
\end{array}\right)
$$

with respect to a positive, orthonormal basis of $T_{z} \Sigma$. An endomorphism $H: T_{z} \Sigma \rightarrow$ $T_{z} \Sigma$ commutes with $J_{z}$ if and only if $H$ has the matrix expression

$$
\left(\begin{array}{cc}
a & -b \\
b & a
\end{array}\right) \quad a, b \in \mathbb{R}
$$

in such a basis. In particular, we deduce that $H$ is, up to a scalar multiple, a rotation matrix. Analogously, $H$ anticommutes with $J_{z}$ if and only if $H$ has the matrix expression

$$
\left(\begin{array}{cc}
a & b \\
b & -a
\end{array}\right) \quad a, b \in \mathbb{R}
$$

in such a basis. In particular, we deduce that $E^{a}$, see (1.1), is exactly the bundle of endomorphisms anticommuting with $J$. Therefore, if we denote by $E^{c} \rightarrow \Sigma$ the bundle of endomorphisms commuting with $J$, we get the splitting

$$
\begin{equation*}
\operatorname{End}(T \Sigma)=E^{c} \oplus E^{a} \tag{1.14}
\end{equation*}
$$

Furthermore, as $J$-complex line bundle we can write

$$
E^{a} \cong T \otimes \overline{T^{*}}, \quad T:=T^{(1,0)} \Sigma
$$

where $T^{(1,0)} \Sigma$ is the holomorphic tangent bundle of $\Sigma$ and $\overline{T^{*}}$ denotes the conjugate of the dual bundle of $T$. With this identification, a local section of $E^{a}$ is given by $\frac{\partial}{\partial z} \otimes \mathrm{~d} \bar{z}$ where $z$ is a local holomorphic coordinate compatible with $J$. Thus, the Euler number of $E^{a}$ as real oriented plane bundle coincides with its Chern number as complex line bundle. Using that $c_{1}(T)=\chi(\Sigma)$, this gives another derivation of (1.2) by computing

$$
c_{1}\left(E^{a}\right)=c_{1}\left(T \otimes \overline{T^{*}}\right)=c_{1}(T)-c_{1}\left(T^{*}\right)=c_{1}(T)+c_{1}(T)=2 c_{1}(T)=2 \chi(\Sigma)
$$

Finally, let us assume that $z$ is a conformal point of an orientation-preserving diffeomorphism $F: \Sigma \rightarrow \Sigma$. Then,

$$
\begin{equation*}
\left(F^{*} g\right)_{z}=c g_{z} \text { for some } c>0 \tag{1.15}
\end{equation*}
$$

If we denote by $M$ the matrix representation of $\mathrm{d}_{z} F$ with respect to positive, orthonormal bases of $T_{z} \Sigma$ and $T_{F(z)} \Sigma$, then (1.15) can be rewritten as

$$
M^{T} M=c I .
$$

This condition is equivalent to saying that $M$ is, up to a scalar multiple, a rotation matrix. Since $J_{z}$ and $J_{F(z)}$ are represented by the matrix (1.13), we see that $\mathrm{d}_{z} F_{J_{z}}=J_{F(z)} \mathrm{d}_{z} F$. We conclude that $z$ is a conformal point of $F$ if and only if $F$ is $j$-holomorphic at $z$ with respect to $J$-holomorphic coordinates around $z$ and $F(z)$.

### 1.5 Conformal Points of Vector Fields

Let $f$ be a vector field on $\Sigma$. Let $F_{t}: \Sigma \rightarrow \Sigma$ be the time- $t$ map of the flow of $f$. Suppose that $z \in \Sigma$ is a point such that $F_{t}(z) \in \mathcal{C}\left(g,\left(F_{t}\right)^{*} g\right)$ for all $t$ close to zero. In particular, $F_{t}$ is $J$-holomorphic at $z$ in a local $J$-holomorphic chart for all small $t$. Taking the derivative in $t$ at $t=0$, we conclude that the vector field $f$ is $\underline{j}$-holomorphic at $z$. In other words $\bar{\partial}_{j} f$ is a section of $E^{a}$ which vanishes at $z$. Here $\bar{\partial}_{j}$ denotes the Cauchy-Riemann operator sending sections of the holomorphic tangent bundle $T=T^{(1,0)} \Sigma$ to sections of $T \otimes \overline{T^{*}} \cong E^{a}$, and $f$ is identified with its image under the isomorphism

$$
\begin{equation*}
T \Sigma \rightarrow T^{\mathbb{C}} \Sigma \cong T^{(1,0)} \Sigma \oplus T^{(0,1)} \Sigma \rightarrow T^{(1,0)} \Sigma \tag{1.16}
\end{equation*}
$$

where $T^{\mathbb{C}} \Sigma$ is the complexification of $T \Sigma$.
Let $\mathcal{C}(J, f)$ be the set of zeros of $\bar{\partial}_{j} f$. If $\mathcal{C}(J, f)$ is finite and $\mathcal{C}(J, f) \subset \Sigma \backslash \partial \Sigma$, then we can associate an index $\operatorname{ind}_{(J, f)}(z)$ to each $z \in \mathcal{C}(J, f)$ and a winding number $w_{i}(J, f)$ representing the relative winding number of $\bar{\partial}_{J} f$ with respect to the canonical section $R_{i}$ along $C_{i}$ for every $i=1, \ldots, n$. Defining the algebraic count

$$
[\mathcal{C}(J, f)]:=\sum_{z \in \mathcal{C}(J, f)} \operatorname{ind}_{(J, f)}(z),
$$

we get the following consequence of Theorem 1.2(1).

Corollary 1.10 Let $J$ be a complex structure on a compact surface $\Sigma$ and $f$ a vector field on $\Sigma$ such that $\mathcal{C}(\jmath, f)$ is finite and $\mathcal{C}(\jmath, f) \subset \Sigma \backslash \partial \Sigma$. Then the equation

$$
[\mathcal{C}(J, f)]=2 \chi(\Sigma)+\sum_{i=1}^{n} w_{i}(J, f)
$$

holds.

### 1.6 An Open Question

Given any Riemannian metric $g$ on $\Sigma$ and diffeomorphism $F: \Sigma \rightarrow \Sigma$, it is interesting to ask which further restrictions must the points $z$ of $\mathcal{C}\left(g, F^{*} g\right)$, their indices $\operatorname{ind}_{\left(g, F^{*} g\right)}(z)$ and the numbers $w_{i}\left(g, F^{*} g\right)$ satisfy besides Eq. (1.7). This question is related to the uniformization theorem for compact surfaces with boundary via Theorem 1.2(2). For instance, given any two metrics $g$ and $h$ on $\Sigma=S^{2}$ or $\Sigma=D^{2}$, we can find a diffeomorphism $F: \Sigma \rightarrow \Sigma$ such that $F^{*} g$ and $h$ are conformal at every point, see [15, Theorem 1]. Thus, $\mathcal{C}\left(g, F^{*} g\right)=\mathcal{C}(g, h), \operatorname{ind}_{\left(g, F^{*} g\right)}(z)=\operatorname{ind}_{(g, h)}(z)$ for every $z$ in this set, and $w_{i}\left(g, F^{*} g\right)=w_{i}(g, h)$ for all $i=1, \ldots, n$. As a consequence of Theorem 1.2(2), there are no further restrictions in this case.

As an instructive example, let us construct a diffeomorphism $F: S^{2} \rightarrow S^{2}$ having exactly one conformal point without resorting to the uniformization theorem. We start by defining a suitable vector field $f$ (a section of $T^{(1,0)} \Sigma$ under the identification (1.16)) and then we will define $F:=F_{t}$, the time- $t$ map of $f$ for $t>0$ small enough. The complex manifold $S^{2}$ is obtained gluing together two charts with holomorphic local coordinates $z, w \in \mathbb{C}$ having the relationship $z w=1$. In the $z$-chart, we define $f$ via $f_{1}(z):=z^{3} \rho\left(|z|^{2}\right)$ where $\rho:[0, \infty) \rightarrow \mathbb{C}$ is any smooth curve such that $\rho^{\prime}(s) \neq 0$ for all $s \in[0, \infty)$ and such that $\rho(s)=s$ for small $s$ and $\rho(s)=-1 / s$ for large $s$. Thus $\bar{\partial} f_{1}(z)=z^{3} \rho^{\prime}\left(|z|^{2}\right) z$ is different from zero if $z \neq 0$. On the other hand, for $z$ near zero, $\bar{\partial} f_{1}(z)=z^{4}$ and thus $z=0$ is a holomorphicity point of index 4. In the $w$-chart, $f_{1}$ gets transformed into

$$
f_{2}(w):=-w^{2} f_{1}\left(\frac{1}{w}\right)=-\frac{1}{w} \rho\left(\frac{1}{|w|^{2}}\right)
$$

which for $w$ small enough becomes $f_{2}(w)=\frac{|w|^{2}}{w}=\bar{w}$. Thus $f_{2}$ smoothly extends at $w=0$ and $\bar{\partial} f_{2}(0)=1 \neq 0$. Defining $f$ as $f_{2}$ in the $w$-chart, we get the desired vector field.

We now claim that the time- $t$ flow map $F_{t}$ of $f$ has exactly one conformal point at $z=0$, provided that $t>0$ is small enough. Indeed, if $z$ is close to 0 , then $f_{1}(z)=z^{4} \bar{z}$ is real analytic and therefore the Cauchy-Kovalevskaya Theorem implies that $F_{t}(z)$ is real analytic and the expansion $F_{t}(z)=z+t z^{4} \bar{z}+O\left(t^{2}|z|^{6}\right)$ holds. Therefore, $\bar{\partial} F_{t}(z)=t z^{4}+O\left(t^{2}|z|^{5}\right)$, which vanishes only for $z=0$ when $t>0$ is small enough. On the other hand, in any compact neighborhood of 0 in the $w$-chart, we have $F_{t}(w)=w+t f_{2}(w)+O\left(t^{2}\right)$ and $\bar{\partial} F_{t}(w)=t \bar{\partial} f_{2}(w)+O\left(t^{2}\right)$ which is nowhere vanishing if $t>0$ is small enough since $\left|\bar{\partial} f_{2}(w)\right|$ is bounded away from zero.

Coming back to the case of an arbitrary surface $\Sigma$, we see that in general there are metrics $g$ and $h$ such that $h$ and $F^{*} g$ are not conformal at all points, no matter how we choose the diffeomorphism $F$. The easiest examples where this happens is when $\Sigma=\mathbb{T}^{2}$, or when $\Sigma=D^{2}$ and we require in addition the diffeomorphism $F$ to be the identity at the boundary. For instance, on $\mathbb{T}^{2}$ conformal classes of metrics $g$ are classified by lattices $\Gamma$ in $\mathbb{C}$, up to Euclidean isometries and homotheties, where $g$ is the Riemannian metric on $\mathbb{T}^{2}=\mathbb{C} / \Gamma$ induced by the Euclidean metric on $\mathbb{C}$. To get an example on the disc, let us identify $D^{2}$ with the unit Euclidean disc in $\mathbb{C}$. Let $g$ be the Euclidean metric on $D^{2}$. Recall that the group of diffeomorphisms
$\varphi: D^{2} \rightarrow D^{2}$ such that $g$ and $\varphi^{*} g$ are conformal at all points consists of the Möbius transformations preserving $D^{2}$. Consider $G: D^{2} \rightarrow D^{2}$ to be any diffeomorphism such that $\left.G\right|_{\partial D^{2}} \neq\left.\varphi\right|_{\partial D^{2}}$ for all $\varphi$. Such a $G$ surely exists since if $\varphi$ is not the identity, then $\varphi$ can have at most two fixed points on the boundary. If we define $h:=G^{*} g$, then there is no diffeomorphism $F: D^{2} \rightarrow D^{2}$ which is identity at the boundary and such that $F^{*} h$ and $g$ are conformal at every point. Indeed, if such an $F$ exists, then $(G \circ F)^{*} g=F^{*} G^{*} g=F^{*} h$ is conformal to $g$ at all points, which means that $F \circ G=\varphi$ for some Möbius transformation $\varphi$ preserving the disc. Since $F$ is the identity at the boundary, this would imply that $G=\varphi$ on the boundary. A contradiction.

Thus, in the case of $\mathbb{T}^{2}$ and of $D^{2}$, it is meaningful to ask if there is a metric $g$ and a diffeomorphism $F$ (being the identity on the boundary in the case of $D^{2}$ ) such that $\mathcal{C}\left(g, F^{*} g\right)$ is empty. If one can find a vector field $f$ (vanishing on the boundary in the case of $D^{2}$ ) such that $\mathcal{C}(J, f)=\varnothing$, then $\mathcal{C}\left(g, F_{t}^{*} g\right)=\varnothing$ for small $t \neq 0$, as well, where $F_{t}$ is the time- $t$ map of the flow of $f$.

In the case of $\Sigma=\mathbb{T}^{2}$, we can readily find such a vector field for all conformal classes of complex structures. Indeed, let $\mathbb{T}^{2}=\mathbb{C} / \Gamma$ where $\Gamma$ is a lattice in $\mathbb{C}$ and let $J$ be the complex structure on $\Sigma$ induced by that on $\mathbb{C}$. Up to Euclidean isometries and homotheties, we can assume that $\Gamma$ is generated by $1, \tau \in \mathbb{C}$, where $\tau=a+i b$ with $b>0$. Consider the vector field which in a global holomorphic trivialization of $T^{(1,0)} \Sigma$ is written as $f(z)=\mathrm{e}^{\frac{2 \pi i}{b} \operatorname{Im} z}$. Notice that $f$ is well-defined since it is invariant under translations by 1 and $\tau$. Moreover,

$$
\bar{\partial}_{j} f(z)=\frac{\partial}{\partial \bar{z}} \mathrm{e}^{\frac{\pi}{b}(z-\bar{z})}=-\frac{\pi}{b} f(z)
$$

which is nowhere vanishing.
However, we do not know if such a vector field $f$ exists on $D^{2}$. Since vector fields on $D^{2}$ correspond to functions in a global trivialization of $T^{(1,0)} D^{2}$, we have the following open question.
Question 1.11 Does there exist a smooth function $f: D^{2} \rightarrow \mathbb{C}$ satisfying the following two conditions?
(1) $\forall z \in D^{2}, \quad \frac{\partial f}{\partial \bar{z}}(z) \neq 0$.
(2) $\forall z \in \partial D^{2}, \quad f(z)=0$.

### 1.7 Plan of the Paper

Theorem 1.2 is proven in Sect. 2. Theorem 1.7 is proven in Sect. 3.

## 2 Proof of Theorem 1.2

We prove Theorem 1.2(1). Let $h$ be a symmetric bilinear two-tensor field over $\Sigma$ such that $\mathcal{C}(g, h)$ is finite and $\mathcal{C}(g, h) \subset \Sigma \backslash \partial \Sigma$. Recall the definition of $H$ and $H^{a}$ from (1.3) and (1.4).

If $\Sigma$ has no boundary, then $[\mathcal{C}(g, h)]=e\left(E^{a}\right)=2 \chi(\Sigma)$ by the Poincaré-Hopf Theorem for oriented plane bundles [2, Theorem 11.17]. If $\Sigma$ has boundary, let $\hat{\Sigma}$ be the closed, oriented surface that we obtain from $\Sigma$ by gluing a disc $D_{1}, \ldots, D_{n}$ along each boundary component $C_{1}, \ldots, C_{n}$. The gluing maps $D^{2} \rightarrow D_{i}$ have the Euclidean disc

$$
D^{2}=\left\{(x, y) \in \mathbb{R}^{2} \mid x^{2}+y^{2} \leq 1\right\}
$$

as domain and send the boundary $\partial D^{2}$ traversed in the positive sense to $\bar{C}_{i}$, that is, to $C_{i}$ traversed in the negative sense. In this way, the gluing maps are positively oriented with respect to the orientation on $\hat{\Sigma}$.

We let $\hat{g}$ be any extension of $g$ to $\Sigma$ as a Riemannian metric. On the bundle $\left.E^{a}\right|_{D_{i}}$ we choose a nowhere vanishing section $M^{i}$ defined as the reflection along the direction of $\partial_{x} \in T D^{2}$. Let $w_{\bar{C}_{i}}\left(H^{a}, M^{i}\right)$ be the winding number of $H^{a}$ with respect to $M^{i}$ along $C_{i}$ traversed in the negative direction. Then

$$
\begin{aligned}
w_{\bar{C}_{i}}\left(H^{a}, M^{i}\right) & =w_{\bar{C}_{i}}\left(H^{a}, R^{i}\right)+w_{\bar{C}_{i}}\left(R^{i}, M^{i}\right) \\
& =-w_{C_{i}}\left(H^{a}, R^{i}\right)+w_{\partial D^{2}}\left(R^{i}, M^{i}\right)=-w_{i}(g, h)+2,
\end{aligned}
$$

where we have used that $\bar{C}_{i}$ is identified with $\partial D^{2}$ and that the unoriented line tangent to $\partial D^{2}$ rotates twice with respect to the horizontal unoriented line. By the Extension Theorem in [10, p. 145], it is possible to construct an extension $\hat{h}$ of $h$ to $\hat{\Sigma}$ such that $\mathcal{C}(\hat{g}, \hat{h})=\mathcal{C}(g, h) \cup\left\{z_{1}, \ldots, z_{n}\right\}$, where $z_{1}, \ldots, z_{n}$ are the centers of the discs $D_{1}, \ldots, D_{n}$ and

$$
\begin{equation*}
\operatorname{ind}_{(\hat{g}, \hat{h})}\left(z_{i}\right)=w_{\bar{C}_{i}}\left(H^{a}, M^{i}\right)=2-w_{i}(g, h) . \tag{2.1}
\end{equation*}
$$

Therefore,

$$
\begin{aligned}
{[\mathcal{C}(g, h)] } & =[\mathcal{C}(\hat{g}, \hat{h})]-\sum_{i=1}^{n} \operatorname{ind}_{(\hat{g}, \hat{h})}\left(z_{i}\right) \\
& =2 \chi(\hat{\Sigma})-2 n+\sum_{i=1}^{n} w_{i}(g, h)=2 \chi(\Sigma)+\sum_{i=1}^{n} w_{i}(g, h),
\end{aligned}
$$

where we used that $\chi(\Sigma)+n=\chi(\hat{\Sigma})$ as follows from the formula $\chi(A \cup B)=$ $\chi(A)+\chi(B)-\chi(A \cap B)$. We have thus completed the proof of Theorem 1.2(1).

Let us first prove Theorem 1.2(2) when $\Sigma$ has no boundary. Let us consider an embedded closed disc $D$ containing $\mathcal{C}$ in its interior. There is a section $H^{\text {out }}$ of $E^{a}$ which is nowhere vanishing on $\Sigma \backslash D$ and there is a section $H^{\text {in }}$ which is nowhere vanishing over $D$. The winding number of $H^{\text {out }}$ with respect to $H^{\text {in }}$ along $\partial D$ is $w\left(H^{\text {out }}, H^{\text {in }}\right)=2 \chi(\Sigma)$. For each $z \in \mathcal{C}$ consider an embedded closed disc $D^{z}$ centered at $z$ and contained in $D$. After shrinking the discs $D^{z}$ we may assume that they are pairwise disjoint. Let $H^{z}$ be a section of $\left.E^{a}\right|_{D^{z}}$ which has just one zero at $z$
with index $\operatorname{ind}(z)=l(z)$. Thus the winding number of $H^{z}$ with respect to $H^{\text {in }}$ along $\partial D^{z}$ is $w\left(H^{z}, H^{\mathrm{in}}\right)=l(z)$. Since $2 \chi(\Sigma)=\sum_{z \in \mathcal{C}} \iota(z)$ by assumption, we get

$$
w\left(H^{\text {out }}, H^{\text {in }}\right)=\sum_{z \in \mathcal{C}} w\left(H^{z}, H^{\text {in }}\right)
$$

Consider the surface

$$
\tilde{\Sigma}:=D \backslash \bigsqcup_{z \in \mathcal{C}} D^{z} .
$$

It satisfies $\partial \tilde{\Sigma}=\partial D \sqcup\left(\sqcup_{z \in \mathcal{C}} \overline{\partial D^{z}}\right)$. Since $w\left(H^{\text {out }}, H^{\text {in }}\right)-\sum_{z \in \mathcal{C}} w\left(H^{z}, H^{\text {in }}\right)=0$, the Extension Theorem in [10, p. 145] implies that there is a nowhere vanishing section $\tilde{H}$ of $\left.E^{a}\right|_{\tilde{\Sigma}}$ coinciding with $H^{\text {out }}$ on $\partial D$ and with $H^{z}$ on every $\partial D^{z}$. Thus, $H^{\text {out }}, \tilde{H}$, and all $H^{z}$ glue together to yield a section $H$ of $E^{a} \rightarrow \Sigma$ having the desired properties.

When $\Sigma$ has boundary, we construct the closed surface $\hat{\Sigma}$ as in the proof of Theorem 1.2(1). We define $\hat{\mathcal{C}}:=\mathcal{C} \cup\left\{z_{1}, \ldots, z_{n}\right\}$ and $\hat{\imath}: \hat{\mathcal{C}} \rightarrow \mathbb{Z}$ as the extension of $\iota$ such that $l\left(z_{i}\right)=2-w_{i}$ for all $i=1, \ldots, n$. Applying Theorem 1.2(2) for closed surfaces to $\hat{\Sigma}$ and $\hat{\imath}$ and using (2.1) yields Theorem 1.2(2) for the case of surfaces with boundary, as well.

## 3 Proof of Theorem 1.7

Let $C_{i}$ be a component of $\partial \Sigma$ for some $i \in\{1, \ldots, n\}$. There is $j \in\{1, \ldots, n\}$ such that $F\left(C_{i}\right)=C_{j}$. Recall that $\left.\mathrm{d} F\right|_{C_{i}}$ is expressed by the matrix

$$
N_{i}=c_{i}\left(\begin{array}{ll}
a_{i} & 0 \\
b_{i} & 1
\end{array}\right)
$$

with respect to the positive orthonormal bases $\nu_{i}, \tau_{i}$ and $\nu_{j}, \tau_{j}$.
The metric $\left.F^{*} g\right|_{C_{i}}$ is represented by the endomorphism $\mathrm{d} F^{T} \cdot \mathrm{~d} F$ via (1.3). A computation shows that the matrix representing $\mathrm{d} F^{T} \cdot \mathrm{~d} F$ with respect to the basis $\nu_{i}, \tau_{i}$ is

$$
N_{i}^{T} N_{i}=c_{i}^{2} Q_{i}, \quad \text { with } Q_{i}=\left(\begin{array}{cc}
a_{i}^{2}+b_{i}^{2} & b_{i} \\
b_{i} & 1
\end{array}\right) .
$$

We point out that the condition that $\left(a_{i}, b_{i}\right)$ is never equal to $(1,0)$ is equivalent to $Q_{i}$ having distinct eigenvalues, since $Q_{i}$ is symmetric. Let $q_{i}: C_{i} \rightarrow \mathbb{R} P^{1} \cong \mathbb{R} / \pi \mathbb{Z}$ be the eigendirection of $Q_{i}$ with larger eigenvalue. By (1.6), $w_{i}\left(g, F^{*} g\right)$ is the degree of the map $q_{i}: C_{i} \rightarrow \mathbb{R} / \pi \mathbb{Z}$. Therefore, our goal is to show that the degree of $q_{i}$ is equal to the winding number of $\left(a_{i}-1, b_{i}\right): C_{i} \rightarrow \mathbb{R}^{2}$ around the origin. To this purpose, let us parametrize $C_{i}$ in the positive direction by $\theta_{i} \in \mathbb{R} / 2 \pi \mathbb{Z}$ and, to ease notation, let us drop all the subscripts $i$ in what follows.

We may assume without loss of generality that the curve $(a-1, b): \mathbb{R} / 2 \pi \mathbb{Z} \rightarrow \mathbb{R}^{2}$ intersects the positive real axis transversely. In this case $w(a-1, b)$ counts the number of points $\theta_{0} \in \mathbb{R} / 2 \pi \mathbb{Z}$ such that $\left(a\left(\theta_{0}\right)-1, b\left(\theta_{0}\right)\right)$ lies on the positive real axis, namely $a\left(\theta_{0}\right)>1$ and $b\left(\theta_{0}\right)=0$, with sign: the intersection is counted positively if $b^{\prime}\left(\theta_{0}\right)>0$ and negatively if $b^{\prime}\left(\theta_{0}\right)<0$.

On the other hand, the degree of $q$ is computed using a regular value $\xi \in \mathbb{R} / \pi \mathbb{Z}$ of $q$. Being regular means that $q^{\prime}\left(\theta_{0}\right) \neq 0$ for all $\theta_{0} \in q^{-1}(\xi)$. In this case, the degree of $q$ counts number of points $\theta_{0} \in q^{-1}(\xi)$ with sign: the point $\theta_{0}$ is counted positively if $q^{\prime}\left(\theta_{0}\right)>0$ and negatively if $q^{\prime}\left(\theta_{0}\right)<0$.

Choosing $\xi=0$, we see that $\theta_{0} \in q^{-1}(0)$ if and only if $(1,0) \in \mathbb{R}^{2}$ is an eigenvector of $Q$ with eigenvalue larger than 1 . This happens exactly when $b\left(\theta_{0}\right)=0$ and $a\left(\theta_{0}\right)>$ 1 , that is when $(a-1, b)$ intersects the positive real axis. Therefore, we prove that 0 is a regular value of $q$ and that $w(a-1, b)$ is the degree of $q$ if we can show that for every such $\theta_{0}$ the numbers $b^{\prime}\left(\theta_{0}\right)$ and $q^{\prime}\left(\theta_{0}\right)$ have the same sign.

For this purpose, let $v(\theta)=(x(\theta), y(\theta)) \in \mathbb{R}^{2}$ be a generator of the line $q(\theta)$ such that $v\left(\theta_{0}\right)=(1,0)$ and write $\lambda(\theta)$ for the corresponding eigenvalue of $Q(\theta)$, so that $\lambda\left(\theta_{0}\right)=a\left(\theta_{0}\right)$. Then $q^{\prime}\left(\theta_{0}\right)=y^{\prime}\left(\theta_{0}\right)$. To compute $y^{\prime}\left(\theta_{0}\right)$ we differentiate the vector equation $(Q(\theta)-\lambda(\theta) I) v(\theta)=0$ at $\theta_{0}$ :

$$
\left(Q\left(\theta_{0}\right)-\lambda\left(\theta_{0}\right) I\right) v^{\prime}\left(\theta_{0}\right)+\left(Q^{\prime}\left(\theta_{0}\right)-\lambda^{\prime}\left(\theta_{0}\right) I\right) v\left(\theta_{0}\right)=0
$$

Therefore, substituting the values for $Q\left(\theta_{0}\right), \lambda\left(\theta_{0}\right)$ and $Q^{\prime}\left(\theta_{0}\right)$ and taking the $y$ component of the vector equation, we get

$$
\left(1-a\left(\theta_{0}\right)\right) y^{\prime}\left(\theta_{0}\right)+b^{\prime}\left(\theta_{0}\right)=0
$$

Thus,

$$
q^{\prime}\left(\theta_{0}\right)=y^{\prime}\left(\theta_{0}\right)=\frac{b^{\prime}\left(\theta_{0}\right)}{a\left(\theta_{0}\right)-1}
$$

from which we see that $q^{\prime}\left(\theta_{0}\right)$ and $b^{\prime}\left(\theta_{0}\right)$ have the same sign since $a\left(\theta_{0}\right)>1$. This completes the proof.

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