**Research Article** 

# One-step synthesis and electromagnetic absorption properties of high entropy rare earth hexaborides (HE REB<sub>6</sub>) and high entropy rare earth hexaborides/borates (HE REB<sub>6</sub>/HE REBO<sub>3</sub>) composite powders

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Abstract: Considering the emergence of severe electromagnetic interference problems, it is vital to develop electromagnetic (EM) wave absorbing materials with high dielectric, magnetic loss and optimized impedance matching. However, realizing the synergistic dielectric and magnetic losses in a single phase material is still a challenge. Herein, high entropy (HE) rare earth hexaborides ( $REB_6$ ) powders with coupling of dielectric and magnetic losses were designed and successfully synthesized through a facial one-step boron carbide reduction method, and the effects of high entropy borates intermedia phases on the EM wave absorption properties were investigated. Five HE REB<sub>6</sub> ceramics including  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0,2}Y_{0,2}Eu_{0,2}Er_{0,2}Yb_{0,2})B_{6,2}$  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$ , and  $(Nd_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$  possess CsCl-type cubic crystal structure, and their theoretical densities range from 4.84 to 5.25 g/cm<sup>3</sup>. (Ce<sub>0.2</sub>Y<sub>0.2</sub>Sm<sub>0.2</sub>Er<sub>0.2</sub>Yb<sub>0.2</sub>)B<sub>6</sub> powders with the average particle size of 1.86 µm were found to possess the best EM wave absorption properties among these hexaborides. The RL<sub>min</sub> value of (Ce<sub>0.2</sub>Y<sub>0.2</sub>Sm<sub>0.2</sub>Er<sub>0.2</sub>Yb<sub>0.2</sub>)B<sub>6</sub> reaches -33.4 dB at 11.5 GHz at thickness of 2 mm; meanwhile, the optimized effective absorption bandwidth  $(E_{AB})$  is 3.9 GHz from 13.6 to 17.5 GHz with a thickness of 1.5 mm. The introduction of HE REBO<sub>3</sub> (RE = Ce, Y, Sm, Eu, Er, Yb) as intermediate phase will give rise to the mismatching impedance, which will further lead to the reduction of reflection loss. Intriguingly, the HEREB<sub>6</sub>/HEREBO<sub>3</sub> still possess wide effective absorption bandwidth of 4.1 GHz with the relative low thickness of 1.7 mm. Considering the better stability, low density, and good EM wave absorption properties, HE REB<sub>6</sub> ceramics are promising as a new type of EM wave absorbing materials.

**Keywords:** high entropy rare earth hexaborides (HE REB<sub>6</sub>); one-step synthesis; electromagnetic wave absorbing properties; synergistic dielectric and magnetic losses; wide effective absorption bandwidth

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# 1 Introduction

With the electromagnetic (EM) wave as an important carrier of information exchange penetrating into all aspects in human life, it is critical to develop high performance EM wave absorbing materials to eliminate the adverse effects of EM wave, such as the electromagnetic interference and microwave pollution [1–3]. Under increasingly complex serving environment, EM wave absorbing materials are required to possess the following properties: (1) strong absorption, broad efficient absorption bandwidth, and optimized impedance matching; (2) lightweight, good mechanical property, and favorable high-temperature stability; (3) resistance to corrosion and oxidation; (4) easy fabrication and cost-efficient [3-5]. However, achieving high EM wave absorption performance with foregoing features in a single phase material is still a challenge [6]. For examples, magnetite with superior dual electromagnetic properties (ferrimagnetic and dielectric properties) has shortcomings of high density and inferior stability [7,8]; magnetic metals with high permeability have disadvantages of high density and poor resistance to environment corrosion [9,10]; lightweight carbon-based nanomaterials and conducting polymers both take the advantage of high conductivity but also suffer from dissatisfactory electromagnetic impedance matching [11–13]. Through element doping [14], structure design [15], surface treatment [10], and composite construction [16], weight reduction and impedance matching adjustment can be effectively realized. However, realizing synergistic dielectric and magnetic losses in a single phase material is still difficult. Thus, searching for materials with good high-temperature stability, corrosion and oxidation resistance is appealing. Consequently, it is of great significance to develop lightweight high temperature EM wave absorbing ceramics with exceptional corrosion and oxidation resistance.

EM wave absorbing materials are capable of converting absorbed EM wave into thermal energy in terms of dielectric loss and/or magnetic loss [17]. Considering the conductivity and magnetism of rare-earth (RE) hexaborides (REB<sub>6</sub>), it is expected that lightweight REB<sub>6</sub> ceramics possess promising EM wave absorption properties. REB<sub>6</sub> exhibits a CsCl-type crystal structure, with an RE atom surrounded by eight boron octahedra [18]. The boron framework is deficient of 2 electrons that can be donated by RE atoms, making divalent REB<sub>6</sub> (EuB<sub>6</sub>, YbB<sub>6</sub>) semiconductors or semimetals

while other trivalent REB<sub>6</sub> conductors [19–22]. Besides, the rigidity of octahedral boron cage endows REB<sub>6</sub> with high melting point, high hardness in lattice a direction, and superior chemical stability [23,24]. The substitution of RE atoms leads to the adjustment of the distances of intra-octahedron bond (B-Bintra) and inter-octahedra bond (B-Binter), which will affect the electrical properties and then give rise to the transition of conductivity from semiconductors to conductors [25,26]. It is also found that with the exception of paramagnetic LaB<sub>6</sub>, trivalent REB<sub>6</sub> are antiferromagnetic at low temperature and EuB<sub>6</sub> exhibits unusual ferromagnetism below 15.3 K [27–31]. Surprisingly,  $(Ca_{1-x}La_x)B_6$  has a ferromagnetic Curie temperature of 600 K, which equals to that of transition-metal ferromagnets such as Fe [32]. Combining excellent high-temperature property, good electrical conductivity and magnetic property, REB<sub>6</sub> are expected to be a new type of promising EM wave absorbing materials.

The similar lattice parameters of different REB<sub>6</sub> is beneficial to tuning the properties of REB<sub>6</sub> through element doping [26,33]. Besides, boron carbide reduction method is found effective to synthesize REB<sub>6</sub> without post-synthesis treatment such as acid washing [34,35]. Consequently, it is feasible to design and synthesize high-entropy rare-earth hexaborides (HE REB<sub>6</sub>) with multi-principal rare-earth elements. It is expected that the introduction of configuration entropy can yield a batch of unexpected properties including better resistance to environment attack, better oxidation and corrosion resistance, good high-temperature stability, and lower thermal conductivity [36-42]. Most importantly, the "cocktail" effect on electrical and magnetic properties is anticipated which will open a new door to tailor EM wave absorbing properties including absorbing capability, effective absorbing bandwidth, and impedance matching. Previous work in the authors' lab has found that high-entropy rare-earth silicide carbide/rare-earth oxide composites show a wide efficient absorption bandwidth of 4.5 GHz with a maximum reflection loss of -50.9 dB at thickness of 2.0 mm [43]. However, by far EM wave absorption properties of HE REB<sub>6</sub> have not been explored.

The aims of this work are to design and synthesize HE REB<sub>6</sub> and to explore the EM wave absorption properties of HE REB<sub>6</sub>. To achieve these goals, five high-entropy rare-earth hexaborides  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$ , and  $(Nd_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2})E_{0.2}Eu_{0.2}$ 

 $Yb_{0,2})B_6$  are designed and successfully synthesized. Hereafter they are referred to as HE REB<sub>6</sub>-1, HE REB<sub>6</sub>-2, HE REB<sub>6</sub>-3, HE REB<sub>6</sub>-4, and HE REB<sub>6</sub>-5. The choice of rare-earth elements is due to the following considerations. Firstly, hexaborides containing these RE elements exhibit the same crystal structure. Secondly, as shown in Table 1 [44], the difference of selected ion radius is smaller than 15%, which makes the formation of solid solution with a single phase easy. Thirdly, the selected compounds including NdB<sub>6</sub>, SmB<sub>6</sub>, YB<sub>6</sub>, ErB<sub>6</sub>, and CeB<sub>6</sub> are conductors with antiferromagnetism, while  $EuB_6$  is a semiconductor with two ferromagnetic transition temperature of 12.5 and 15.3 K. Fourthly, although some REB<sub>6</sub> such as YbB<sub>6</sub> is a semiconductor, the formation of high-entropy hexaborides could lead to the band gap narrowing which can tune the electrical conductivity [45]. The electrical conductivity and magnetic ordering in hexaborides endow these new types of HE REB<sub>6</sub> promising EM wave absorption properties. Finally, the selected compounds also take advantages of lightweight, exceptional high-temperature stability, and favorable resistant to corrosion and oxidation. The formation of HE REB<sub>6</sub> could lead to further improvement of the comprehensive properties including mechanical property, thermal property, and EM wave absorption property.

# 2 Experimental

#### 2.1 Synthesis of HE REB<sub>6</sub> powders

Five HE REB<sub>6</sub>, i.e.,  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Y_{0.2}Eu_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$ , and  $(Nd_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$  powders were synthesized by a one-step boron

Table 1Ion radius and radius difference of the selectedRE ions for the design of HE REB6 (data were obtainedfrom the revised list of effective ionic radii [44])

carbide reduction process using commercially available
Vo Co Co F O F O VI O
$Y_2O_3$ , $CeO_2$ , $Sm_2O_3$ , $Eu_2O_3$ , $Er_2O_3$ , $Yo_2O_3$ powders
(99.9% purity, 1 µm, HWRK Chem Co., Ltd., Beijing,
China) and B <sub>4</sub> C powders (98% purity, -200 mesh,
Mudanjiang Qianjin Reagent Co. Ltd., Heilongjiang,
China) as starting materials. Five rare earth oxide
powders in the stoichiometric ratio of required HE
REB <sub>6</sub> were mixed together with B <sub>4</sub> C powders to
prepare HE REB <sub>6</sub> powders. The initial stage of the
boron carbide reduction process leads to the formation
of REBO3 as an intermediate phase as described by
Eqs. (1)-(4), where RE refers to Y, Sm, Eu, Er, or Yb.
In particular, the molar ratio of reactants is listed in
Table 2.

$$4RE_2O_3 + 7B_4C = 4REBO_3 + REB_6 + 7C$$
 (1)

$$4\text{REBO}_{3} + 7\text{C} + 5\text{B}_{4}\text{C} = 4\text{REB}_{6} + 12\text{CO}$$
(2)

$$6CeO_2 + 4B_4C = 4CeBO_3 + 2CeB_6 + 4C$$
 (3)

$$4CeBO_3 + 4C + 5B_4C = 4CeB_6 + 9CO + \frac{3}{2}O_2$$
 (4)

Each mixture was ball-milled with ethanol and agate media in a polytetrafluoroethylene jar for 4 h. After ball milling, the homogeneously mixed slurry was fully dried in an oven at 60 °C. After drying and being sieved through a 120-mesh screen, the mixed powders were uniaxially cold-pressed into cylindrical pellets of 30 mm diameter in a stainless-steel die. Eventually, the green bodies were heated to 1900  $\,^{\circ}\!C$  and held for 1 h in vacuum to synthesize pure HE REB<sub>6</sub> powders. In order to study the influence of intermediate phases, i.e., high-entropy rare-earth borates (HE REBO<sub>3</sub>) on the EM wave absorption performance of HE REB<sub>6</sub>, the pellets were heated to 1600 °C and kept for 1 h in Ar atmosphere to obtain HE REB6/HE REBO3 composites. The as-prepared samples were ground into powders in an agate mortar for subsequent characterizations.

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Table 2 Constituent of raw materials of the KED		bbl	<u> </u>		Ongt	+++++++++++++++++++++++++++++++++++++++	• • •	MONT	mot	01010	a ot		
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RE	Ion radius (Å)	Radius difference (%)	HE compound	Molar ratio of raw materials	
Nd	0.98	11	$(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$	$\begin{array}{c} 2 \ CeO_2: 1 \ Y_2O_3: 1 \ Sm_2O_3: \\ 1 \ Er_2O_3: 1 \ Yb_2O_3: 15 \ B_4C \end{array}$	
Sm	0.96	9	$(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_{6}$	$2\ CeO_2:1\ Eu_2O_3:1\ Sm_2O_3:$	
Eu	0.95	8	(0.2 -0.2 - 0.2 - 0.2) 0	$1 \text{ Er}_2\text{O}_3 : 1 \text{ Yb}_2\text{O}_3 : 15 \text{ B}_4\text{C}$	
Y	0.90	3	$(Ce_{0.2}Y_{0.2}Eu_{0.2}Er_{0.2}Yb_{0.2})B_6$	$\begin{array}{c} 2 \ CeO_2: 1 \ Y_2O_3: 1 \ Eu_2O_3: \\ 1 \ Er_2O_3: 1 \ Yb_2O_3: 15 \ B_4C \end{array}$	
Er	0.89	2	$(C_{2}, \mathbf{V}, \mathbf{Sm}, \mathbf{Eu}, \mathbf{Vh})$ P	$2 \text{ CeO}_2$ : $1 \text{ Y}_2\text{O}_3$ : $1 \text{ Sm}_2\text{O}_3$ :	
Ce	0.87	0	$(Ce_{0.2} I_{0.2}SII_{0.2}Eu_{0.2} I_{0.2})B_6$	$1 Eu_2O_3$ : $1 Yb_2O_3$ : $15 B_4C$	
Yb	0.87	0	$(Nd_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$	$\begin{array}{l} 1 \ Nd_2O_3: 1 \ Y_2O_3: 1 \ Sm_2O_3: \\ 1 \ Eu_2O_3: 1 \ Yb_2O_3: 15 \ B_4C \end{array}$	

# 2.2 Phase composition, microstructure, and EM absorbing properties

Phase composition was analyzed by an X-ray diffractometer (XRD, D8 Advanced, Bruker, Germany) utilizing Cu K $\alpha$  radiation ( $\lambda = 1.5406$  Å) with a step size of 0.02° at a scanning rate of 2 (°)/min. The weight fraction of intermediate phase is quantitatively analyzed using the total pattern solution software (TOPAS, Bruker Corp., Karlsruhe, Germany). The microstructure of as-prepared HE REB<sub>6</sub> powders was investigated by a scanning electron microscope (SEM, Apollo 300, CamScan, UK) equipped with an energy dispersive spectroscopy system (EDS Inca X-Max 80 T, Oxford, UK).

The relative complex permeability ( $\mu_r = \mu' + j\mu''$ ) and permittivity ( $\varepsilon_r = \varepsilon' + j\varepsilon''$ ) were determined by a vector network analyzer (Agilent N5244A, USA) using the coaxial method in the frequency range of 1.0–18.0 GHz. High-entropy REB<sub>6</sub> powders and paraffin wax were mixed at a mass ratio of 7:3 and compacted into a toroidal shape ( $\Phi_{out} = 7.00 \text{ mm}, \Phi_{in} = 3.04 \text{ mm}$ ). The reflection loss (*RL*), which is one of the evaluation criteria for the EM wave absorption properties, was determined from the relative complex permeability and permittivity for a given frequency and absorber thickness based on the transmission-line theory and metal back-panel model [46], according to Eqs. (5) and (6):

$$RL(dB) = 20\log|(Z_{in} - Z_0)/(Z_{in} + Z_0)| \qquad (5)$$

$$Z_{\rm in} = Z_0 \sqrt{\frac{\mu_{\rm r}}{\varepsilon_{\rm r}}} \tanh\left(j\frac{2\pi f d\sqrt{\mu_{\rm r}\varepsilon_{\rm r}}}{c}\right) \tag{6}$$

where  $Z_0$  is the impedance of free space,  $Z_{in}$  is the normalized input characteristic impedance at the interface between atmosphere and absorber,  $\mu_r$  is the relative complex permeability,  $\varepsilon_r$  is the relative complex permittivity, f is the frequency, d is the thickness of samples, c is the velocity of light (3×10<sup>8</sup> m/s). The effective absorption bandwidth ( $E_{AB}$ ) is defined by the range of frequencies over which RL is smaller than -10 dB, which represents that 90% of the microwave energy is absorbed while only 10% of that is reflected [47].

### 3 Results and discussion

# 3.1 Synthesis and microstructure of pure HE REB<sub>6</sub> powders

As shown in Fig. 1, rare earth hexaborides exhibit

CsCl-type cubic structure with a space group of  $Pm\overline{3}m$ (No. 221). The rare-earth and boron atoms are located at 1a(0, 0, 0) and 6f(0.5, 0.5, z) positions, respectively, where the position parameter z of B atom is related to the lattice parameter (a) and B-B bond lengths [48]. On the one hand, the structure of REB<sub>6</sub> is constructed by the strong covalent  $\sigma$  bonds that link boron octahedrons, which limits the change of the lattice parameters but only allows the adjustments of the lengths of inter-octahedra  $\sigma$  bond (B-B<sub>inter</sub>) and intra-octahedron  $\tau$  bond (B–B<sub>intra</sub>) [19,23,24,33,49]. The changes in B-Binter and B-Bintra distances have great impact on the electrical properties, leading to the electric conduction transition. Each boron octahedron in REB<sub>6</sub> is deficient of two electrons and thus the stability is ensured by the transformation of an electron from rare earth metals to boron octahedron, which give rise to good electrical conductivity in trivalent hexaborides [25,26,50]. As for the intriguing magnetism found in REB<sub>6</sub>, the magnetic ordering might be attributed to an indirect exchange of the RKKY type via conduction electrons donated from trivalent rare-earth ions or impurities [51,52].

Figure 2 shows the XRD pattern of  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$  (HE REB<sub>6</sub>-1) synthesized at 1900 °C for 1 h in vacuum. Standard ICDD/JCPDS cards of CeB<sub>6</sub>, YB<sub>6</sub>, SmB<sub>6</sub>, ErB<sub>6</sub>, and YbB<sub>6</sub> are also presented to confirm the formation of HE REB<sub>6</sub>, which is structural ordered but compositional disordered with multi-principle RE elements.

To characterize the particle size and morphology, the HE REB<sub>6</sub>-1 powders after grinding in an agate mortar are observed by SEM as shown in Fig. 3(a). The powders are equiaxial and the average particle size is about 1.86  $\mu$ m as shown in Fig. 3(b).



Fig. 1 Crystal structure of  $REB_6$ . Rare earth (RE) atoms (brown) and  $B_6$  octahedron (cyan) are arranged in a CsCl-like packing.



Fig. 2 XRD pattern of  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$  (HE REB<sub>6</sub>-1) together with those of CeB<sub>6</sub>, YB<sub>6</sub>, SmB<sub>6</sub>, ErB<sub>6</sub>, and YbB<sub>6</sub> obtained from ICDD/JCPDS cards.



Fig. 3 (a) SEM image and (b) particle size distribution of  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$  powders.

Figure 4 compares the XRD patterns of five HE REB<sub>6</sub> synthesized at 1900 °C for 1 h in vacuum. The diffraction peaks of five samples show consistency in spite of different constituents, which indicate that single-phase high-entropy solid-solutions are formed among these rare-earth hexaborides. The lattice parameters of these HE REB<sub>6</sub> are calculated according to Bragg's law as shown in Eq. (7):

$$2d_{hkl} \cdot \sin \theta = \lambda \tag{7}$$

J Adv Ceram 2021, 10(1): 62-77

where  $\lambda = 1.5406$  Å,  $\theta$  is obtained from peak position of {100} by fitted Gaussian peak functions. The lattice parameter of a cubic lattice is equal to the interplanar spacing of  $\{100\}$  plane. Then the theoretical density is calculated by the lattice parameter. The lattice parameter (a) and theoretical density ( $\rho_{th}$ ) of five HE  $REB_6$  are listed in Table 3. The variation of lattice parameter is derived from the substitutions of RE elements, which leads to different degrees of lattice distortions. By replacing rare-earth ions (Y<sup>3+</sup>, Sm<sup>3+</sup>, Er<sup>3+</sup> with ionic radius of 0.9 Å, 0.96 Å, 0.89 Å respectively) in HE REB<sub>6</sub>-1 with  $Eu^{2+}$  (r = 1.17 Å),  $(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ (HE REB<sub>6</sub>-2),  $(Ce_{0.2}Y_{0.2}Eu_{0.2}Er_{0.2}Yb_{0.2})B_6$ (HE REB<sub>6</sub>-3), and (Ce<sub>0.2</sub>Y<sub>0.2</sub>Sm<sub>0.2</sub>Eu<sub>0.2</sub>Yb<sub>0.2</sub>)B<sub>6</sub> (HE REB<sub>6</sub>-4) with larger lattice parameters are obtained, which can be attributed to the substitution of  $\mathrm{Eu}^{2+}$  that has larger ionic radius. Likewise, replacing  $Ce^{3+}$  (r = 0.87 Å) in HE REB<sub>6</sub>-4 with  $Nd^{3+}$  (r = 0.98 Å) will lead to ( $Nd_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}$ )  $Yb_{0,2}B_6$  (HE REB<sub>6</sub>-5) with a larger lattice parameter.



Fig. 4 XRD patterns of HE REB<sub>6</sub>-1  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Y_{0.2})B_6$ , HE REB<sub>6</sub>-2  $(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ , HE REB<sub>6</sub>-3  $(Ce_{0.2}Y_{0.2}Eu_{0.2}Er_{0.2}Yb_{0.2})B_6$ , HE REB<sub>6</sub>-4  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$ , and HE REB<sub>6</sub>-5  $(Nd_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$ .

Table 3Lattice parameter (a) and theoretical density $(\rho_{th})$  of HE REB6

Hexaborides	Composition	a (Å)	$ ho_{ m th}({ m g/cm^3})$
HE REB <sub>6</sub> -1	$(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$	4.1092±0.0005	5.00
HE REB <sub>6</sub> -2	$(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$	4.1215±0.0005	5.25
HE REB <sub>6</sub> -3	$(Ce_{0.2}Y_{0.2}Eu_{0.2}Er_{0.2}Yb_{0.2})B_6$	4.1279±0.0005	4.94
HE REB <sub>6</sub> -4	$(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$	4.1321±0.0005	4.84
HE REB <sub>6</sub> -5	$(Nd_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6$	4.1355±0.0005	4.85

# 3. 2 Synthesis and microstructure of HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite powders

From Fig. 5, standard ICDD/JCPDS cards of borates show that cerium borate exhibits orthorhombic aragonitetype structure which is different from the rest borates with hexagonal vaterite-type structure. After synthesized at 1600 °C for 1 h in Ar atmosphere, the main phases of the product are  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$  and  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})BO_3$ , with the existence of a small amount of cerium borate and boron carbide, which is attributed to the uncompleted reduction.

Figure 6 presents the backscattered electron image (BSEI) and EDS patterns of two selected points in  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6/(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})$  BO<sub>3</sub> composites. Through semi-quantitative analysis, the atomic ratio of O/B of two selected points in the golden phase and the blue phase are 1.602 and 0.018, respectively. Combining with XRD analysis, it can be concluded that the golden phase is HE REBO<sub>3</sub> and the blue phase is HE REB6. Therefore, it comes to the conclusion that HE REBO<sub>3</sub> phase is distributed on the surface of HE REB<sub>6</sub> in HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite powders, as shown in Fig. 7.

Figure 8 shows the XRD patterns of the composite powders synthesized at 1600  $^{\circ}$ C for 1 h in Ar atmosphere. The diffraction peaks of these five composites are similar to each other and all can be indexed as vaterite-type structure although they possess different compositions. Therefore, it can be concluded



Fig. 5 XRD pattern of  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$  and  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})BO_3$  composite powders together with those of CeBO<sub>3</sub>, YBO<sub>3</sub>, SmBO<sub>3</sub>, ErBO<sub>3</sub>, and YbBO<sub>3</sub> obtained from ICDD/JCPDS cards.



Fig. 6 (a) BESI of  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6/(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})BO_3$  composite powders, (b) EDS spectrum of selected point in the golden phase, and (c) EDS spectrum of the selected point in the blue phase.



Fig. 7 BESI of  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6/(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})BO_3$  composite powers indicating the distribution of  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})BO_3$  phase (golden color).

that composites with HE REB<sub>6</sub> and HE REBO<sub>3</sub> phases are formed. Hereafter, the five composites synthesized at 1600 °C in Ar atmosphere are referred to as HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1, HE REB<sub>6</sub>/HE REBO<sub>3</sub>-2, HE REB<sub>6</sub>/HE REBO<sub>3</sub>-3, HE REB<sub>6</sub>/HE REBO<sub>3</sub>-4, and HE REB<sub>6</sub>/HE REBO<sub>3</sub>-5, respectively. To estimate the relative content of different phases in these composites, quantitative analysis was performed using the total pattern solution software (TOPAS). The weight fraction of HE REBO<sub>3</sub> and reliability factors,  $R_p$  and  $R_{wp}$ , are listed in Table 4.

### 3.3 EM wave absorption properties

### 3.3.1 EM wave absorption properties of HE REB<sub>6</sub>

Figure 9 reveals the frequency dependences of the electromagnetic parameters for HE REB<sub>6</sub>/paraffin composites. As illustrated in Fig. 9(a) and Fig. 9(b), across the testing frequency range, the values of  $\varepsilon'$  and  $\varepsilon''$  show a decreasing tendency from HE REB<sub>6</sub>-1 to HE REB<sub>6</sub>-4, which is attributed to the changes of electrical conductivity of HE REB<sub>6</sub> as a result of the changes of lattice constant a due to the substitution of  $Eu^{2+}$ . The decrease of lattice constant a will lead to the increasement of the bond length difference between B-Binter and B-B<sub>intra</sub>, which will further give rise to the increasement of shear anisotropy. On this basis, REB<sub>6</sub> with smaller lattice constant would possess better electrical conductivity [49]. Therefore, the conductivity of HE REB<sub>6</sub> decreases from HE REB<sub>6</sub>-1 to HE REB<sub>6</sub>-4 as the lattice constant increases. In particularly, in exception of HE REB<sub>6</sub>-3,  $\varepsilon''$  values of HE REB<sub>6</sub> decrease with the increasing of frequency at low frequency, which indicates the dispersion phenomenon; while the value of  $\varepsilon''$  exhibits growing tendency as the frequency increases. Generally,



**Fig. 8** XRD patterns of HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1, HE REB<sub>6</sub>/HE REBO<sub>3</sub>-2, HE REB<sub>6</sub>/HE REBO<sub>3</sub>-3, HE REB<sub>6</sub>/HE REBO<sub>3</sub>-4, and HE REB<sub>6</sub>/HE REBO<sub>3</sub>-5 composite powders.

Table 4	W	eight	fract	ion	of H	E RE	BO <sub>3</sub> and re	eliability
factors,	<b>R</b> <sub>p</sub>	and	<b>R</b> <sub>wp</sub>	in	five	HE	REB <sub>6</sub> /HE	REBO <sub>3</sub>
composi	tes							

Composites	Weight fraction of HE REBO <sub>3</sub> (%)	$R_{\rm p}$ (%)	$R_{wp}$ (%)
HE REB <sub>6</sub> /HE REBO <sub>3</sub> -1	58.34	8.49	14.69
HE REB <sub>6</sub> /HE REBO <sub>3</sub> -2	49.22	7.25	11.93
HE REB <sub>6</sub> /HE REBO <sub>3</sub> -3	64.78	8.32	14.69
HE REB <sub>6</sub> /HE REBO <sub>3</sub> -4	38.67	7.95	13.26
HE REB <sub>6</sub> /HE REBO <sub>3</sub> -5	41.95	6.90	10.85

the magnitude of dielectric loss can be estimated by the dielectric loss tangent (  $\tan \delta_{\varepsilon} = \frac{\varepsilon''}{\varepsilon'}$ ). From Fig. 10(a),

the variation of dielectric loss tangent shows a similar tendency to that of  $\varepsilon''$  value. It is expected that HE REB<sub>6</sub>-1 and HE REB<sub>6</sub>-2 have considerable dielectric loss ability, which is derived from dipolar polarization, conduction loss, and polarization loss [53,54]. In specific, numerous dipoles in HE REB<sub>6</sub> can provide multiple ways for the transmission of the electromagnetic wave [55]. Moreover, abundant polarization charges and a sufficient number of interfaces between the HE REB<sub>6</sub> ceramics and paraffin wax are beneficial to spur interfacial polarization (called the Maxwell-Wagner effect) and related relaxation, which contribute to the dielectric loss [56,57]. The electron migration because of natural electrical conductivity will also lead to the reduction of electromagnetic energy. The polarization mechanism can be described in the assist of Debye theory [58] as shown in Eq. (8):

$$\left(\varepsilon' - \frac{\varepsilon_{\rm s} + \varepsilon_{\infty}}{2}\right)^2 + (\varepsilon'')^2 = \left(\frac{\varepsilon_{\rm s} - \varepsilon_{\infty}}{2}\right)^2 \tag{8}$$

where  $\varepsilon_s$ ,  $\varepsilon_{\infty}$  are the static permittivity, and relative dielectric permittivity at high-frequency limit, respectively. Therefore, the relationship between  $\varepsilon'$  and  $\varepsilon''$  is represented by a curve in the shape of semicircle (called the Cole–Cole semicircle), and each semicircle corresponds to one Debye relaxation process [59]. Figure 10(b) presents the  $\varepsilon' - \varepsilon''$  plots of HE REB<sub>6</sub>, where clear semicircles are pictured to reflect the polarization relaxations of HE REB<sub>6</sub> powders.

The relative complex permeability and magnetic loss tangent of HE REB<sub>6</sub> are shown in Figs. 9(c), 9(d), and 10(c). From Fig. 9(c), the  $\mu'$  values of HE REB<sub>6</sub>-2, HE REB<sub>6</sub>-4, and HE REB<sub>6</sub>-5 decrease in the range of 2–8 GHz and then remain steady because of the relaxation of magnetic moments procession, while those



**Fig. 9** Electromagnetic parameters of HE REB<sub>6</sub>: (a) real permittivity ( $\varepsilon'$ ), (b) imaginary permittivity ( $\varepsilon''$ ), (c) real permeability ( $\mu''$ ), (d) imaginary permeability ( $\mu''$ ).



Fig. 10 Frequency dependences of (a) the dielectric loss tangent ( $\varepsilon''/\varepsilon'$ ) and (c) the magnetic loss tangent ( $\mu''/\mu'$ ); (b) Cole–Cole semicircles and (d)  $C_0$ –*f* curves of HE REB<sub>6</sub>.

of HE REB<sub>6</sub>-1 and HE REB<sub>6</sub>-3 exhibit a small amount of increasement. Figure 9(d) shows frequency dependent of  $\mu''$  of HE REB<sub>6</sub>. Resonance peaks are observed, which can be described by natural resonance [60]. The corresponding values of magnetic loss tangent of HE REB<sub>6</sub> exhibit similar fluctuation with  $\mu''$  values. Generally, the magnetic loss is derived from hysteresis loss, domain wall displacement loss, natural resonance loss, and eddy current loss [61]. Since the characterization is performed in a weak external

magnetic field at the frequency range of gigahertz, hysteresis loss and domain wall displacement loss can be ruled out from the mechanism of magnetic loss in HE REB<sub>6</sub>. The contribution of Eddy current loss to the reflected loss is related to the electrical conductivity of materials ( $\sigma$ ) and the matching thickness (*d*), as represented by Eq. (9)[62]:

$$\mu'' \approx \frac{2}{3}\pi\mu_0(\mu')d^2f\sigma \tag{9}$$

where  $\mu_0$  is the vacuum permeability. If the eddy current loss is the dominate factor for the magnetic loss, the value of  $C_0$  ( $C_0 = \mu''(\mu')^{-2} f^{-1}$ ) should be a constant [63]; otherwise, the natural resonance loss is preponderance. Figure 10(d) presents the values of  $C_0$ of HE REB<sub>6</sub> as a function of frequency. In the frequency range of 1–12 GHz, the values of  $C_0$  decrease with the increase of frequency. With raising frequency, the values of  $C_0$  remain as a constant. Therefore, it can be concluded that both the eddy current loss and the natural resonance loss contribute to the magnetic loss. It is worth noting that the  $\mu''$  values of HE REB<sub>6</sub> in Fig. 9(d) are close to that of Fe<sub>3</sub>O<sub>4</sub>, which is a typical magnetic EM wave absorber with exceptional magnetic features [64,65]. Drawn from the analysis of dielectric and magnetic loss tangents, it is concluded that the absorption abilities of HE REB<sub>6</sub> are generated by the coupling of dielectric and magnetic losses.

Based on above discussions, high electric and magnetic loss tangents will lead to high reflection loss (RL), which is one of the vital evaluation indexes of the EM wave absorption abilities. With relatively higher dielectric and magnetic loss tangents, HE REB<sub>6</sub>-2 might possess the highest RL value. However, from Fig. 11(a), it can be found that HE REB<sub>6</sub>-1 rather than HE REB<sub>6</sub>-2 exhibits the highest RL value, which is due to the better impedance match. It has come to light that the closer the impedance match values to 1, the better the impedance match [66]. From this angle, HE REB<sub>6</sub>-1 possesses the best impedance matching, and HE  $REB_6$ -2 ranks the second, as shown in Fig. 11(b). As shown in Figs. 11(a) and 11(c), the  $RL_{min}$  value of HE REB<sub>6</sub>-1 reaches -33.4 dB at 11.5 GHz with a thickness of 2 mm, and the optimized effective absorption bandwidth ( $E_{AB}$ ) is 3.9 GHz from 13.6 to 17.5 GHz with a thickness of 1.5 mm. As shown in Fig. 11(d), HE REB<sub>6</sub>-2 reaches the largest  $E_{AB}$  of 4.3 GHz from 12.6 to 16.9 GHz at thickness of 1.7 mm. Therefore, the superior EM absorption properties require not only strong absorption, but also good impedance matching. In addition, the variation of impedance match values of HE  $REB_6$  indicates the impedance matching can be tuned via compositional design. The addition of multi-principle elements with different atomic sizes in rare-earth hexaborides brings about severe lattice



Fig. 11 Comparison of (a) reflection loss (*RL*) values and (b) impedance match ( $Z_{in}/Z_0$ ) for HE REB<sub>6</sub>. (c) and (d) Frequency dependence of *RL* values at various thickness for HE REB<sub>6</sub>-1 and HE REB<sub>6</sub>-2, respectively.

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distortion effect, which has an impact on the changes of B–B bonding lengths and electronic structure [25,26]. Thus, different electronic conductivity and magnetic performance in HE REB<sub>6</sub> can be achieved through composition design, which will lead to the difference of EM wave absorption properties in HE REB<sub>6</sub> with different constituent.

# 3.3.2 EM wave absorption properties of HE REB<sub>6</sub>/HE REBO<sub>3</sub>

Combining the microstructure shown in Fig. 7 and the XRD patterns of HE REB<sub>6</sub>/HE REBO<sub>3</sub> powders shown in Fig. 8, it comes to the conclusion that composites of HE REB<sub>6</sub> and HE REBO<sub>3</sub> are formed. From Fig. 12(a), except for HE REB<sub>6</sub>/HE REBO<sub>3</sub>-4,  $\varepsilon'$  values of HE composites are smaller than those of HE REB<sub>6</sub> in the whole frequency range, indicating that the conductivity is reduced because of the instinct insulating feature of HE REBO<sub>3</sub>, which will further decrease  $\varepsilon''$  value according to free electron theory  $\varepsilon'' \approx \sigma / (2\pi f \varepsilon_0)$ , where  $\sigma$  is the conductivity and  $\varepsilon_0$  is the permittivity of vacuum [67]. However, as shown in Fig. 12(b), the considerable improvement of  $\varepsilon''$  value of HE REB<sub>6</sub>/HE REBO<sub>3</sub> compared with those of HE REB<sub>6</sub> are observed, which indicates that there is another dominant mechanism contributing to electric loss. In specific, the coverage of HE REBO3 intermediate phase will spur more interfacial

polarizations and related relaxations among HE REB<sub>6</sub>, HE REBO<sub>3</sub> and paraffin wax phases, which will lead to more severe interface polarization under the action of electromagnetic waves. Therefore, compared to HE REB<sub>6</sub> samples, the values of electric dissipation factor  $(\varepsilon''/\varepsilon')$  of HE REB<sub>6</sub>/HE REBO<sub>3</sub> samples are higher (Fig. 13(a)). As for the permeability shown in Figs. 12(c)and 12(d), the  $\mu'$  values of HE REB<sub>6</sub>/HE REBO<sub>3</sub> are close to those of HE REB<sub>6</sub>, while the  $\mu''$  values are lower than those of HE REB<sub>6</sub>. According to Eq. (9), the decrease of electrical conductivity due to the formation of instinct insulating HE REBO<sub>3</sub> phase will give rise to reduction of  $\mu''$  value. Then, from Fig. 13(c), the magnetic loss tangent  $(\mu''/\mu')$  of HE REB<sub>6</sub>/HE REBO<sub>3</sub> samples is reduced, which indicates relatively lower magnetic loss. These foregoing results indicate that both dielectric loss and magnetic loss of HE REB<sub>6</sub>/HE REBO<sub>3</sub> contribute to dissipating the electromagnetic wave energy, but the dominant effect is due to the dielectric loss.

The dielectric loss is caused by Debye dipolar relaxation [59], which is proven by the presence of semicircle curves as shown in Fig. 13(b). The  $C_0$  values in Fig. 13(d) decrease dramatically as frequency increase in the frequency range of 1–6 GHz; then, they keep unchanged as the frequency increases. Therefore, the magnetic loss is derived from eddy current loss and



**Fig. 12** Electromagnetic parameters of HE REB<sub>6</sub>/HE REBO<sub>3</sub>: (a) real permittivity ( $\varepsilon'$ ), (b) imaginary permittivity ( $\varepsilon''$ ), (c) real permeability ( $\mu'$ ), (d) imaginary permeability ( $\mu''$ ).

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Fig. 13 Frequency dependences of (a) the dielectric loss tangent ( $\varepsilon''/\varepsilon'$ ) and (c) the magnetic loss tangent ( $\mu''/\mu'$ ); (b) Cole–Cole semicircles and (d)  $C_0$ –*f* curves of HE REB<sub>6</sub>/HE REBO<sub>3</sub>.

natural resonance loss [60,63]. As has been discussed in the previous section, the EM wave absorption properties are strongly affected by the impedance match. Otherwise, electromagnetic wave will be reflected on the surface of the EM wave absorbing materials, and the reflect loss ability is in vain when it comes to the fact that the entering of EM wave into the absorber is difficult to realize [68]. Therefore, the reduction of reflection losses of HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite powders shown in Fig. 14(a) is due to the unfitting impedance match compared with those of HE REB<sub>6</sub> samples. For example, from Fig. 14(a), the value of impedance match  $(Z_{in}/Z_0)$  of HE REB<sub>6</sub>-1 and HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1 increases from 1.06 to 1.19, which gives rise to reduction of the  $RL_{min}$  value from -33.4 to -24.1 dB with the same thickness of 2 mm.

Although the reflection loss abilities of HE REB<sub>6</sub>/HE REBO<sub>3</sub> samples become weaker, HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite powders still exhibit considerable effective absorption bandwidth with low thickness. As shown in Figs. 14(c) and 14(d), both HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1 and HE REB<sub>6</sub>/HE REBO<sub>3</sub>-2 exhibit wide  $E_{AB}$  of 4.1 GHz with the same width of 1.7 mm. Thus, the introduction of HE REBO<sub>3</sub> as an intermediate phase into HE REB<sub>6</sub> will lead to greater dielectric loss, relatively poorer magnetic loss, and worsen impedance matching. For example, in the frequency range of

10–14 GHz, the value of dielectric tangent loss of HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1 reaches 0.35 which is 1.75 times of that of HE REB<sub>6</sub>-1, while the value of magnetic tangent loss decreases from 0.12 to 0.05, and the value of impedance match ( $Z_{in}/Z_0$ ) gets further away from 1. As a result, the value of  $RL_{min}$  of HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1 decreases, which indicates that the EM wave absorption abilities become weaker. The above results infer that although the coverage of insulator HE REBO<sub>3</sub> phase is beneficial to spur interfacial polarization and related relaxation and give rise to greater dielectric loss, worsen of the magnetic loss and impedance matching, which will finally lead to the weaken of overall reflection loss ability.

Figure 15 summarizes the optimized  $E_{AB}$  value and the sample thickness of some reported EM absorbing materials of various kinds (SiC based [69–73], carbons based [74–78], ferrite metal-organic framework (Fe-MOF) [47], phosphides based [79], sulfides based [68,80], polymers based [81–83], metals/alloys based [62,63,67, 84,85], and oxides based materials [64,66,86–89]). As we have mentioned in previous sections, seeking for materials with thin thickness and wide effective bandwidth is the trend for the development of EM wave absorbing materials. From this aspect, searching for materials in region II are the purposes of many researches. As shown in Fig. 15, although materials in



**Fig. 14** Comparison of (a) reflection loss (*RL*) values and (b) impedance match ( $Z_{in}/Z_0$ ) for HE REB<sub>6</sub>/HE REBO<sub>3</sub>. (c, d) Frequency dependence of *RL* values at various thickness for HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1 and HE REB<sub>6</sub>/HE REBO<sub>3</sub>-2, respectively.



**Fig. 15** Effective bandwidth and corresponding sample thickness of various kinds of reported EM absorbing materials.

region I possess wide  $E_{AB}$ , they are still restricted by thicker thickness. In contract, while materials in region III exhibit small thickness, their narrow effective bandwidth is the hindrances. As is self-evident, materials in region IV have disadvantages of both heavy thickness and narrow  $E_{AB}$ . It is worthy of noting that HE REB<sub>6</sub>-1, HE REB<sub>6</sub>-2, HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1, and HE REB<sub>6</sub>/HE REBO<sub>3</sub>-2 are located in region II, which exhibit wide  $E_{AB}$  with considerable thin thickness, indicating their superior EM wave absorption properties among these materials.

### 4 Conclusions

In conclusion, five HE REB<sub>6</sub> powders including  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$ ,  $(Ce_{0.2}Y_{0.2}Eu_{0.2}Er_{0.2}Yb_{0.2})B_{6}$ ,  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_{6}$ , and (Nd<sub>0.2</sub>Y<sub>0.2</sub>Sm<sub>0.2</sub>Eu<sub>0.2</sub>Yb<sub>0.2</sub>)B<sub>6</sub> as well as five HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite powders including  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6/(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})BO_3$  $(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_{6}/(Ce_{0.2}Eu_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})BO_{3}$ (Ce<sub>0.2</sub>Y<sub>0.2</sub>Eu<sub>0.2</sub>Er<sub>0.2</sub>Yb<sub>0.2</sub>)B<sub>6</sub>/(Ce<sub>0.2</sub>Y<sub>0.2</sub>Eu<sub>0.2</sub>Er<sub>0.2</sub>Yb<sub>0.2</sub>)BO<sub>3</sub>,  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})B_6/(Ce_{0.2}Y_{0.2}Sm_{0.2}Eu_{0.2}Yb_{0.2})BO_3$ and (Nd<sub>0.2</sub>Y<sub>0.2</sub>Sm<sub>0.2</sub>Eu<sub>0.2</sub>Yb<sub>0.2</sub>)B<sub>6</sub>/(Nd<sub>0.2</sub>Y<sub>0.2</sub>Sm<sub>0.2</sub>Eu<sub>0.2</sub>Yb<sub>0.2</sub>) BO3 were designed and successfully synthesized through a facile one-step boron carbide reduction method. XRD analysis reveals that the five as-synthesized HE REB<sub>6</sub> are phase-pure solid solutions with CsCl-type cubic structure, and the five as-synthesized HE REB<sub>6</sub>/HE REBO<sub>3</sub> contain two main phases of HE REB<sub>6</sub> and HE REBO<sub>3</sub>. The theoretical density of HE REB<sub>6</sub> ranges from 4.84 to 5.25 g/cm<sup>3</sup>. The weight fraction of HE REBO<sub>3</sub> in HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite powders was estimated through quantitative phase analysis to be ranging from 38.67% to 64.78%. SEM analysis indicates that the  $(Ce_{0.2}Y_{0.2}Sm_{0.2}Er_{0.2}Yb_{0.2})B_6$  powders possess equiaxial morphology with the mean particle size of 1.86 µm and HE REBO3 is distributed on the surface of HE REB<sub>6</sub> phases in HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite

powders.

The synergetic effects of dielectric and magnetic losses and tunable impedance matching are attributed to favorable EM wave absorption properties. The  $RL_{min}$ value of HE REB<sub>6</sub>-1 reaches -33.4 dB at 11.5 GHz at thickness of 2 mm, and the optimized effective absorption bandwidth ( $E_{AB}$ ) is 3.9 GHz from 13.6–17.5 GHz at thickness of 1.5 mm. The introduction of HE REBO3 as an intermediate phase into HE REB<sub>6</sub> leads to greater dielectric loss, relatively poor magnetic loss, and worse impedance matching, which further leads to the reduction of EM wave reflection loss ability in HE REB<sub>6</sub>/HE REBO<sub>3</sub> composite powders. Although the  $RL_{min}$  value of HE REB<sub>6</sub>/HE REBO<sub>3</sub>-1 decreases to -24.1 dB, it still reaches wide  $E_{AB}$  of 4.1 GHz from 13.4–17.5 GHz at thin thickness of 1.7 mm. Single-phase HE REB<sub>6</sub> ceramics with coupling of dielectric and magnetic losses possess thin thickness, wide effective absorption bandwidth, and considerable EM absorption capability. Combining good high-temperature stability, oxidation and corrosion resistance, lightweight HE REB<sub>6</sub> ceramics are promising as a new type of EM absorbing materials that can meet the requirements of high temperature or corrosive serving environments.

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