Role of *in-situ* formed free carbon on electromagnetic absorption properties of polymer-derived SiC ceramics

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Abstract: In order to enhance dielectric properties of polymer-derived SiC ceramics, a novel single-source-precursor was synthesized by the reaction of an allylhydrido polycarbosilane (AHPCS) and divinyl benzene (DVB) to form carbon-rich SiC. As expected, the free carbon contents of resultant SiC ceramics annealed at 1600 °C are significantly enhanced from 6.62 wt% to 44.67 wt%. After annealing at 900–1600 °C, the obtained carbon-rich SiC ceramics undergo phase separation from amorphous to crystalline feature where superfine SiC nanocrystals and turbostratic carbon networks are dispersed in an amorphous SiC(O) matrix. The dielectric properties and electromagnetic (EM) absorption performance of as-synthesized carbon-rich SiC ceramics are significantly improved by increasing the structural order and content of free carbon. For the 1600 °C ceramics mixed with paraffin wax, the minimum reflection coefficient ($RC_{\min}$) reaches -56.8 dB at 15.2 GHz with the thickness of 1.51 mm and a relatively broad effective bandwidth (the bandwidth of RC values lower than -10 dB) of 4.43 GHz, indicating the excellent EM absorption performance. The carbon-rich SiC ceramics have to be considered as harsh environmental EM absorbers with excellent chemical stability, high temperature, and oxidation and corrosion resistance.

Keywords: polymer-derived ceramics (PDCs); microstructural evolution; dielectric properties; electromagnetic properties

1 Introduction

Centimeter-wave frequency electromagnetic (EM) waves (2–18 GHz) are widely used in various environments, which require high efficient EM absorbing (EMA) materials due to the massive increasing use of high

operating frequencies and bandwidths in electronic systems [1–7]. Recently, polymer-derived ceramics (PDCs) attract great attention as EM absorbers since they possess tunable dielectric properties and conductivities, low density, excellent high-temperature mechanical properties, oxidation, and corrosion resistance as well as thermophysical and chemical stabilities [8,9]. However, due to their intrinsic low permittivity, high-temperature resistant ceramics including most PDCs in general cannot meet the basic requirements for the application as EMA materials [8].

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Pristine PDCs without additional conductive fillers mostly show low dielectric loss [9–17]. Due to the excellent electrical conductivity of carbon nanophase-based families, such as carbon nanotubes, graphene, and reduced graphene oxide, great efforts have therefore been made to introduce them into the PDCs to enhance the EMA performance [18–26]. However, as an EM absorbent, the effective utilization of carbon nanophase in composites depends strongly on their dispersed state since the agglomeration can lead to an inferior polarization relaxation and electrical conductivity [21]. Therefore, complicated chemistry is always involved to homogeneously distribute the carbon nanophase in the matrix [22,23,26]. It is worth pointing out that a phenomenon associated with the PDC, namely, the *in-situ* formation of free carbon within the generated silicon-based ceramic matrix has been widely reported [27–30]. Among the phases occurring in the PDCs, most of them are insulators (e.g., SiO₂ and Si₃N₄) or semiconductors (SiC, SiO_xC_y, SiC_xN_y, and amorphous carbon) with the room-temperature DC electrical conductivity in the ranges of 10^{−14}–10^{−12} S/cm for insulators and 10^{−4}–10² S/cm for semiconductors [30]. Only the segregated free carbon which exists in the form of turbostratic carbon or graphitic carbon within the PDCs has a significantly higher electrical conductivity in the range of 1–10⁵ S/cm [31]. Compared with the incorporated carbon nanophase, the free carbon is *in-situ* formed and homogeneously distributed with the PDC-based ceramic matrix to avoid the agglomeration, which facilitates the resultant dielectric and EM absorbing properties. Therefore, the *in-situ* generated carbon within the PDC matrix has been considered as conductive phase and plays a critical role to improve the final EM absorbing properties [32–41].

Among the above-mentioned studies, metal-based compounds are often involved in the complicated composites and influence the final EM properties, which makes it difficult to evaluate the contribution of *in-situ* formed free carbon [33,38,40,41]. Therefore, the effect of *in-situ* formed carbon including the concentration and/or structural evolution on the dielectric properties and EM absorbing performance of the PDCs has never been systematically discussed up to now. To keep this consideration in mind, we tried to synthesize metal-free and carbon-rich SiC with tunable free carbon contents. The amount and morphology of free carbon within the PDCs strongly depend on the molecular structure of preceramic polymers and the nature of organic

substituents [42]. Divinyl benzene (DVB) is rich in unsaturated groups (C=C and phenyl), which is often used as an additional carbon source to enhance the carbon content in final PDCs [43, 44]. Therefore, in the present study, DVB was introduced into an allylhydrido polycarbosilane (AHPCS) to form a carbon-rich single-source-precursor. The free carbon content of resulting ceramics can be easily tuned by the amount of DVB in the feed. The dielectric properties and EM absorption performance of as-synthesized carbon-rich SiC ceramics mixed with a paraffin wax matrix were systematically investigated. A relationship between the content/microstructure of *in-situ* formed free carbon and dielectric/EM properties of the resultant carbon-rich SiC ceramics will be acquired.

2 Experimental

2.1 Materials

The AHPCS with a formal composition [SiH_{1.26}(CH₃)_{0.60}(CH₂CH=CH₂)_{0.14}CH₂]_n was synthesized as previously reported [45]. DVB was purchased from Aldrich. Paraffin wax (Ar) was bought from Shanghai Hualing Health Machinery Co., China. The Ar atmosphere was used for pyrolysis and annealing process with a purity of 99.99%. Other commercially available reagents were used as received.

2.2 Synthesis of single-source-precursors

One typical synthesis of single-source-precursor was achieved as follows: Firstly, 3.000 g AHPCS and 3.000 g DVB were added to the Schlenk flask equipped with a gas inlet, reflux condenser, and a gas outlet. The obtained yellow solution was then heated up to 160 °C in an oil bath for 6 h in the Ar atmosphere with magnetic stirring. Finally, a pale yellow solid can be obtained and denoted as PAD-2. A series of single-source-precursor PADs with different DVB contents in the feed were prepared (Table 1). To make a comparison, a control test was performed under the same conditions but no

Table 1 Content of AHPCS and DVB in the feed

Precursor	AHPCS (wt%)	DVB (wt%)
AHPCS-blank	100	0
PAD-1	67	33
PAD-2	50	50
PAD-3	40	60

DVB was added, and the resultant product is nominated as AHPCS-blank.

2.3 Polymer-to-ceramic transformation

The single-source-precursors were pyrolyzed in the Ar atmosphere as follows: heating up to the target temperature (300, 600, or 900 °C) hold for 2 h and followed by cooling down to room temperature. The typical heating rate is as follows: 20–500 °C, 2 °C/min; 500–900 °C, 5 °C/min. Finally, the amorphous ceramic (pre-pyrolyzed at 900 °C) was annealed at different temperatures in the range of 1200–1600 °C. The 900 °C ceramic was transferred to a graphite crucible in a tube furnace and then heated rapidly in Ar to a target temperature at a rate of 5 °C/min and kept at this temperature for 2 h. After annealing, the resulting ceramic was furnace-cooled to room temperature. The synthesis procedure from precursors to ceramic powders is illustrated in Fig. 1.

2.4 Characterization

Fourier transform infrared spectroscopy (FT-IR) was recorded on a Nicolet Avator 360 apparatus (Nicolet, Madison, WI) in the range of 4000–500 cm^{-1} with KBr plates for liquid samples and KBr disks for solid samples. Elemental analysis of the ceramic powders was carried out by Horiba Carbon/Sulfur Analyzer EMIA-320V (Horiba, Kyoto, Japan) for carbon element, Horiba Oxygen/Nitrogen Analyzer EMGA-620W (Horiba) for oxygen element, and Elemental Analyzer EA/MA1110 (Carlo Erba, Italy) for hydrogen element. The Si fraction was calculated as the difference between 100% and the values of the other elements. Raman spectra were recorded from 100 to 4000 cm^{-1} employing a micro-Raman spectrometer (XploRA, HORIBA Scientific) using a laser wavelength of 532 nm. Transmission electron microscope (TEM, JEM-2100, JEOL, Tokyo, Japan) was used to observe the microstructure of the ceramics annealed at different temperatures, and the TEM specimens were prepared by mechanical grinding.

The electromagnetic parameters of the ceramics have been determined by the coaxial transmission/reflection method in 2–18 GHz. The ceramic powders were homogeneously mixed with molten paraffin (weight ratio ceramic/paraffin wax of 4:6) and pressed into hollow cylinder-shaped samples (inner diameter: 3.04 mm; external diameter: 7.00 mm; height: 3.04 mm). The complex permittivity of the samples was calculated from the scattering parameters (S-parameters) corresponding to transmission (S_{21} and S_{12}) and reflection (S_{11} and S_{22}) measured on an N5222A PNA-X Vector Network Analyzer (10 MHz to 26.5 GHz, Agilent Technologies Inc., Santa Clara, CA). The hollow cylinder-shaped samples were tightly fixed into the 50 Ω Airline (85051-60007, Agilent) and the measurements were calibrated using an electronic calibration module (N4696-60004, 300 kHz–18 GHz, Agilent).

3 Results and discussion

3.1 Characterization of single-source-precursors

To trace the molecular structure evolution during the single-source-precursor synthesis, the FT-IR spectra of the original AHPCS, DVB, unreacted AHPCS/DVB mixture, and obtained PAD-2 (taking the PAD-2 as an example) were measured, and the results are shown in Fig. 2.

The major absorption bands of DVB were observed at 3088 cm^{-1} due to C–H stretch in $-\text{CH}=\text{CH}_2$ and at 1630 cm^{-1} due to C=C stretch in $-\text{CH}=\text{CH}_2$. The broadband centered at 3025 cm^{-1} is due to phenyl C–H stretching. Other major bonds in the DVB include phenyl C–C (1594, 1574 cm^{-1}), phenyl C–H (1510 cm^{-1}), and phenyl C–C out-of-plane bending (900–650 cm^{-1}) can be assigned [46,47]. According to our previous study [45], the typical AHPCS characteristic peaks are marked in Fig. 2(b). It is worth mentioning that the C–H stretch in C=C–H groups (3077 cm^{-1}) and C=C stretch (1630 cm^{-1}) peaks significantly decrease after the reaction, evidencing the structural differences between

Fig. 1 Photographic illustration of the synthesis procedure.

Fig. 2 FT-IR spectra of original AHPCS, DVB, unreacted AHPCS/DVB mixture, and single-source-precursor PAD-2.

the original AHPCS and obtained PAD. To make a comparison between the unreacted AHPCS/DVB mixture and resultant PDA-2, the absorption bands at 3088 (C–H stretch in $-\text{CH}=\text{CH}_2$ of DVB) and 3077 cm^{-1} (C–H stretch in $-\text{CH}=\text{CH}_2$ of AHPCS) nearly vanish, indicating that the C=C groups from both the DVB and the AHPCS took part in chemical reactions. In our previous studies [48,49], it was found that two major reactions can consume the C=C bonds as follows: hydrosilylation (Si-H/C=C) and vinyl polymerization (C=C/C=C). The hydrosilylation reaction is fast even at low temperatures (it starts at 100–120 $^{\circ}\text{C}$) [28], whereas vinyl polymerizations start at higher temperatures (*ca.* 400 $^{\circ}\text{C}$) [48,49]. With such a low reaction temperature of 160 $^{\circ}\text{C}$, it is assumed that the hydrosilylation (Si-H/C=C) contributes to the consumption of C=C groups of DVB. Besides, the intensity of the Si–H bands (centered at 2100 cm^{-1}) from the PAD-2 decreases compared with that of unreacted AHPCS/DVB mixture, which again confirms the occurrence of hydrosilylation (Si-H/C=C). The FT-IR results clearly show that the DVB was chemically bonded to the AHPCS chains to form the carbon-rich single-source-precursors.

Based on the FT-IR results, the reaction pathways and polymer chain morphologies during the synthesis of the single-source-precursor are demonstrated in Fig. 3. The hydrosilylation between the C=C groups of DVB and the Si–H groups of AHPCS, and the self-hydrosilations of AHPCS are involved to form the single-source-precursors.

3.2 Composition and microstructure of ceramics

Figure 4(a) shows the Raman spectra of the PAD-2

derived ceramics annealed at different temperatures. All spectra indicate the presence of both disorder-induced D and graphite-like G bands. The D-band at 1352 cm^{-1} corresponds to the presence of amorphous carbonaceous products, disordered carbon or defects in graphite, and the G-band at 1580 cm^{-1} is related to the signal from graphite [50]. Additionally, two bands at around 2690 and 2920 cm^{-1} are attributed to overtones or combination modes of the D- and G-bands [41]. With increasing annealing temperature, the D- and G- bands show smaller line width and the intensities of the G- and 2D-bands (2700 cm^{-1}) grow gradually, indicating progressive structural ordering of the carbon. The Raman spectra of ceramics derived from different precursors annealed at the same temperature of 1600 $^{\circ}\text{C}$ were also measured and the results are shown in Fig. 4(b). The graphitization grade (G_G) of carbon-containing material can be calculated as follows: $I_G / (I_G + I_D) \times 100\%$, where I_G and I_D are the intensity of the G- and D-bands, respectively [50]. The G_G of 1600 $^{\circ}\text{C}$ ceramics derived from AHPCS-blank, PAD-1, PAD-2, and PAD-3 is 0.30, 0.36, 0.39, and 0.40, respectively, indicating that the ordering of carbon (degree of graphitization) is higher with the increase of free carbon of final ceramics.

Moreover, with the DVB content increasing, the carbon content of obtained ceramics significantly increases from 33.44 wt% to 59.66 wt% (Table 2). Correspondingly, the free carbon content increases from 6.62 wt% to 44.67 wt%. The results indicate that as expected, the free carbon content of ceramics is significantly increased by the introduction of DVB.

The PAD-2 derived ceramics annealed at different temperatures were carefully studied by transmission

Fig. 3 (a) Reaction pathways and (b) polymer chain morphologies during the synthesis of the single-source-precursor.

Fig. 4 Raman spectra of (a) PAD-2 derived ceramics annealed at different temperatures and (b) 1600 °C ceramics derived from different precursors.

electron microscopy (TEM) with respect to the evolution of its phase composition and microstructure (Fig. 5). The 900 °C sample shows no crystalline phases or segregated phase separation. The 1200 and 1400 °C

samples show a gradually clear Debye ring rather than diffuse halo. High-resolution TEM (HRTEM) imaging shows the characteristic nanostructure of the ceramic where β -SiC crystallites and poorly organized turbostratic

Table 2 Chemical composition of 1600 °C ceramics derived from different precursors

Sample	Chemical composition (wt%)			Empirical formula	Mole fraction	Free carbon (wt%)
	Si ^d	C ^a	O ^b			
AHPCS-blank-1600	64.44	33.44	2.12	SiC _{1.21} O _{0.06}	0.97SiC·0.03SiO ₂ ·0.24C	6.62
PAD-1-1600	47.58	50.33	2.09	SiC _{2.47} O _{0.08}	0.96SiC·0.04SiO ₂ ·1.51C	30.72
PAD-2-1600	41.14	56.25	2.61	SiC _{3.19} O _{0.11}	0.94SiC·0.06SiO ₂ ·2.25C	39.60
PAD-3-1600	37.48	59.66	2.86	SiC _{3.71} O _{0.13}	0.93SiC·0.07SiO ₂ ·2.78C	44.67

^aMeasured by Carbon/Sulfur Analyzer; ^bMeasured by Oxygen/Nitrogen Analyzer; ^dDetermined by the difference.

Fig. 5 (a) Bright field lattice images (insets: selected area electron diffraction) and (b) HRTEM images of PAD-2 derived ceramics annealed at different temperatures (insets: diffraction patterns obtained by Fourier filtered transformation).

carbon are dispersed in an amorphous SiC(O) matrix. With increasing the annealing temperature, the obtained ceramic reveals an enhanced crystallization as shown in the selected area electron diffraction (SAED) insets. Even annealed at the highest temperature of 1600 °C, the resultant ceramic is composed of superfine SiC nanocrystals with an average size of smaller than 20 nm and turbostratic carbon networks. On the basis of Raman spectroscopy and TEM results, the resultant carbon-rich SiC ceramics are characterized as a SiC/C nanocomposite.

Based on the TEM observation, a scheme shown in Fig. 6 is used to demonstrate the microstructural evolution of resultant carbon-rich SiC ceramics. With the annealing temperature increasing, the carbon-rich SiC ceramics undergo phase separation, from amorphous to crystalline feature. Meanwhile, free carbon is segregated at 1200 °C and is better organized to form conductive networks with the temperature increasing. For PDCs, the amorphous phase acts as A phase (electrically insulating matrix), the turbostratic carbon and nanocrystalline SiC can be considered as the so-called B phase (electrically lossy phase) [9].

Fig. 6 Schematic illustration of the microstructural evolution of carbon-rich SiC ceramics.

Therefore, the resultant carbon-rich SiC ceramics are expected to have superior EM absorption properties.

3.3 Dielectric properties

To determine the complex permittivity ($\epsilon_r = \epsilon' - j\epsilon''$) and dielectric loss ($\tan\delta = \epsilon''/\epsilon'$) in 2–18 GHz frequency range, a vector network analyzer was employed and the results are shown in Figs. 7 and 8. All samples were prepared by mixing the as-synthesized ceramic powders with paraffin wax. As shown in Fig. 7, the average ϵ' of PAD-2 derived ceramics annealed at 900, 1200, 1400, and 1600 °C amounts 3.93, 9.01, 9.21, and 12.45,

Fig. 7 (a) The real permittivity, (b) the imaginary permittivity, and (c) the dielectric loss of PAD-2 derived ceramics annealed at different temperatures.

Fig. 8 (a) The real permittivity, (b) the imaginary permittivity, and (c) the dielectric loss of AHPCS-blank, PAD-1, PAD-2, and PAD-3 derived ceramics annealed at 1600 °C.

respectively. The results indicate that the ϵ' increases with the annealing temperature increasing. According to the Debye theory, the ϵ' is related to the dipolar polarization and interfacial polarization of the electrons. As shown in Figs. 5 and 6, phase segregation occurs at 1200 °C to form β -SiC nanocrystallites and turbostratic carbon, which significantly increases polarization loss at the homo-interfaces and hetero-interfaces. As a result, the ϵ' dramatically increases from 900 to 1200 °C during the phase separation from amorphous to crystallization features. In general, the imaginary part ϵ'' is related to dielectric loss and a higher loss tangent

$\tan\delta_e$ implies higher attenuation ability under the condition of impedance match. With respect to the average ε'' , it reaches 0.92, 1.13, 2.38, and 3.65 corresponds to the ceramics annealed at 900, 1200, 1400, and 1600 °C, respectively. For the average $\tan\delta_e$, it amounts 0.23, 0.13, 0.28, and 0.29, respectively. The results indicate that with the annealing temperature increasing, the ε'' gradually increases because the better organized free carbon (Figs. 7(a) and 8) can enhance the electrical conductivity of resultant ceramics.

The influence of free carbon content of 1600 °C ceramics on dielectric properties was also investigated and the results are displayed in Fig. 8. The average ε' of AHPCS-blank derived ceramic is 5.83, whereas that of PAD-1, PAD-2, and PAD-3 derived ceramics is 7.76, 12.45, and 13.76, respectively. With the DVB content increasing, the enhanced free carbon content of resulting ceramics (Table 2) can increase the polarization at interfaces including C/C and SiC/C interfaces, leading to the increasing of ε' . On the other hand, the increased ε' should be attributed to electrical loss, because increasing free carbon content can enhance the conductivity. Besides, the average ε'' ($\tan\delta_e$) of AHPCS-blank, PAD-1, PAD-2, and PAD-3 derived ceramics is 0.78 (0.14), 0.91 (0.12), 3.65 (0.29), and 8.46 (0.62), respectively. As expected, the enhanced free carbon in final ceramics can form electrical conductive networks, leading to an increase of both ε'' and $\tan\delta_e$. Especially, the average ε'' and $\tan\delta_e$ of PAD-2 derived ceramic significantly increase compared with those of AHPCS and PAD-1, probably due to the percolation threshold.

3.4 Electromagnetic absorption properties

For EM wave absorption materials, reflection coefficient (RC) is used to characterize the absorption properties. Based on the transmission-line theory and the metal back-panel model, the RC can be calculated from the relative complex permeability and permittivity [9]:

$$RC(\text{dB}) = 20 \cdot \log_{10} \left| \frac{(Z_{\text{in}} - 1)}{(Z_{\text{in}} + 1)} \right| \quad (1)$$

$$Z_{\text{in}} = \sqrt{\frac{\mu}{\varepsilon}} \tan h \left(\frac{j2\pi\sqrt{\mu\varepsilon}fd}{c} \right) \quad (2)$$

where Z_{in} , d , and μ are the normalized input impedance, thickness, and permeability of the material, respectively; $\tanh$ presents hyperbolic tangent function; c is the light velocity in a vacuum; and f is the microwave frequency.

Figures 9 and 10 show the theoretical RC values of obtained ceramics. The frequency range corresponding to an RC value < -10 dB is defined as the effective absorption bandwidth (EAB), which means that more than 90% of the EM waves can be attenuated. As shown in Fig. 9(a), the PAD-2 derived ceramics annealed at 900–1400 °C possess weak EM absorption capability with the RC higher than -5 dB. The 1600 °C ceramic exhibits the best absorption properties with a $RC_{\text{min}} = -35.3$ dB at 15.20 GHz (more than 99.9% EM waves can be attenuated), $d = 1.50$ mm, and $EAB = 4.48$ GHz, due to the better ordering of free carbon with the annealing temperature increasing. The EM absorption performance of ceramics with different free carbon contents annealed at 1600 °C was calculated and the results are shown in Fig. 9(b). The ceramics derived from the AHPCS-blank and PAD-1 show weak EM absorption properties with the RC values higher than -3 dB. For the PAD-2 and PAD-3 derived ceramics, they possess much lower RC_{min} values with -35.3 and -24.8 dB, respectively, corresponding to much better EM absorption performance. The results show that

Fig. 9 Calculated theoretical RC values of (a) PAD-2 derived ceramics annealed at different temperatures with a thickness of 1.50 mm and (b) AHPCS-blank, PAD-1, PAD-2, and PAD-3 derived ceramics annealed at 1600 °C with a thickness of 1.50 mm.

Fig. 10 (a) Calculated theoretical RC values of composites of PAD-2 derived carbon-rich SiC mixed with wax versus frequency at different sample thickness variation and (b) corresponding 3D images.

with the increase of free carbon content, the EM absorption performance is significantly enhanced. Among all the samples, the PAD-2 derived ceramics show the best EM absorption properties with the lowest $RC_{\min}$.

The RC of PAD-2 derived ceramic annealed at $1600\text{ }^{\circ}\text{C}$ at different thicknesses were also calculated, as shown in Fig. 10.

It is worth mentioning that the $RC_{\min}$ reaches -56.8 dB at 15.22 GHz , $d = 1.51\text{ mm}$, and $EAB = 4.43\text{ GHz}$, indicating that more than 99.999% EM waves can be attenuated by the obtained carbon-rich SiC/paraffin wax composites. The results show that the present carbon-rich SiC ceramics have to be considered as promising EM wave absorbers which can be applied in harsh environments.

Table 3 summarizes EMA properties for various PDCs in references. In comparison to the previous studies, either the broadest EAB of 4.48 GHz or the ultra-small thickness in this study is excellent. The results indicate that the present carbon-rich SiC ceramics exhibit promising EM absorption performance among all the reported PDCs containing *in-situ* generated free-carbon.

The associated EM wave (EMW) absorption mechanisms based on the carbon-rich SiC/paraffin wax composites are proposed, as illustrated in Fig. 11. After

Fig. 11 Schematic illustration of EMW absorption for carbon-rich SiC/paraffin wax composites.

Table 3 Summary of EMW absorption properties for various PDCs containing *in-situ* formed free carbon in references

Sample	Composition	$RC_{\min}$ (dB)	Absorption bandwidth		Matrix	Per. (wt%)	Ref.
			EAB (GHz)	Thickness (mm)			
SiC/C	SiC, amorphous SiC, free carbon	-56.8	4.48	1.51	Paraffin	40	This study
SiCN	SiC, Si_3N_4 , amorphous SiCN, CNWs, graphene-like free carbon	-38.0	4.20	4.50+3.00 (multilayer)	—	100	[32]
SiOC	SiC, Fe_3Si , amorphous SiOC, free carbon	-46.0	3.50	2.93	—	100	[33]
SiCN	Amorphous SiCN, free carbon	-12.5	10.0	8.00	Paraffin	33	[34]
SiCN NWs	SiC, Si_3N_4 , amorphous SiO_xC_y and $\text{SiO}_x\text{C}_y\text{N}_z$, free carbon	-53.1	5.60	1.95	Paraffin	35	[35]
SiC/C	SiC, amorphous SiC, free carbon	-35.0	1.34	1.50	Paraffin	75	[36]
SiCN	Amorphous SiCN, free carbon	-59.6	4.00	2.70	—	100	[37]
Si-C-Co(O)	SiC, CoSi, Co_2Si , amorphous SiC, free carbon	-42.4	3.94	2.65	Paraffin	33	[38]
SiBCN	Amorphous SiBCN, free carbon	-71.8	3.65	3.08	—	100	[39]
SiC/HfC _x N _{1-x} /C	SiC, HfC _x N _{1-x} , free carbon, amorphous SiHfCN	-47.0	3.10	2.40	Paraffin	40	[41]

the phase separation, the hierarchical structures of SiC nanocrystals and turbostratic carbon induce lots of interfaces in the nanocomposites, which are of great benefit to the EMW absorption.

4 Conclusions

In summary, the carbon-rich SiC ceramics were obtained by a facile PDC approach, starting from the AHPCS and carbon-rich DVB. Using FT-IR, it is confirmed that the hydrosilation between Si–H groups of AHPCS and C=C bonds of DVB occurs to form a carbon-rich single-source precursor. Compared to the AHPCS-derived ceramic, the free carbon contents of single-source-precursor derived SiC ceramics annealed at 1600 °C are tailorable and significantly enhanced from 6.62 wt% to 44.67 wt%. After annealing at the temperature range of 1200–1600 °C, the obtained carbon-rich SiC ceramics are composed of superfine SiC nanocrystals and turbostratic carbon networks dispersed in an amorphous SiC(O) matrix. The dielectric properties and EM absorption performance of resultant carbon-rich SiC ceramics are significantly improved by the better ordering and higher free carbon contents. For the 1600 °C ceramics derived from the PAD-2, the $RC_{\min}$ reaches –56.8 dB at 15.2 GHz with a thickness of 1.51 mm and a relatively broad EAB of 4.43 GHz, indicating the excellent EM absorption performance. The present carbon-rich SiC ceramics open a facile strategy to fabricate excellent EMW absorbers applied in harsh environments.

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