




Hadley Circulation in the Present and Future Climate Simulations of the K-ACE Model

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Abstract

Hadley circulation (HC) is a planetary-scale overturning circulation in the tropics that transports momentum, heat, and moisture poleward. In this study, we evaluate the strength and extent of the HC in the historical and future climate simulations of the Korean Meteorological Administration (KMA) Advanced Community Earth system model (K-ACE), which was recently developed by the National Institute of Meteorological Sciences of Korea. Compared with a reanalysis product, the overall structure of the HC is reasonably reproduced by the K-ACE. At the same time, it is also found that the Northern Hemisphere HC in the K-ACE is shifted southward by a few degrees, while the strength of the Southern Hemisphere (SH) HC is under-represented by approximately 20%. These biases in the strength and extent of the HC can be explained by biases in the eddy momentum flux and precipitation in the tropics. In the future climate simulations under the Shared Socioeconomic Pathway 5-Representative Concentration Pathway 8.5 scenario, the HCs in the K-ACE show a weakening and widening trend in both hemispheres, which is consistent with the projections of many Coupled Model Intercomparison Project Phase 6 models. A notable feature of the K-ACE is the widening of the SH HC, which takes place at a rate that is about double the multi-model mean. Climate models that share the component models with the K-ACE, such as UKESM, HadGEM3-GC31-LL, and ACCESS-CM2/ESM1, also show enhanced poleward expansion of the HC in the SH. This strong expansion is shown to be dominated by the expansion of the regional HC over the Pacific.

Keywords K-ACE model · Hadley circulation · CMIP · Historical simulation · Scenario simulation

1 Introduction

Hadley circulation (HC) refers to a thermally direct, meridional overturning circulation in the tropics. The ascending motion of the HC at the equator is associated with the convection of warm and moist air, which leads to the formation of clouds and precipitation. The descending motion of the HC in the subtropics is accompanied by the subsidence of relatively cool and dry air. As a result, the

HC has a great influence on the global hydrological cycle, energy balance, and the circulation of the atmosphere. Furthermore, regional climates, such as monsoon climates, and sea surface temperatures are also known to be strongly connected to the strength and meridional extent of the HC (Bindoff et al. 2013; Zhou et al. 2020). For example, when the HC is widened, precipitation over the southeast Australia decreases, and the dry season of the southern Amazon gets enhanced (Cai and Cowan 2013; Fu et al. 2013).

For the width of the HC, many climate models project the widening in response to anthropogenic forcing, including increases in greenhouse gas concentration (Chemke and Polvani 2019; Hu et al. 2018). When the carbon dioxide concentration is quadrupled (i.e., the $4\times\text{CO}_2$ forcing simulations), many models in the Coupled Model Intercomparison Project Phase 6 (CMIP6) show expansions of the HC by a few degrees over decades, while the degree of expansion varies between models, hemispheres, and seasons (Grise and Davis 2020). The same conclusion has previously been drawn from CMIP Phase 5 (CMIP5) model

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simulations (Tao et al. 2016), despite some differences in the extent of widening being greater for the CMIP6 models in the Northern Hemisphere (NH), possibly because of the differences in external forcings, such as solar/volcanic forcing, anthropogenic aerosols, and ozone concentrations (Allen and Ajoku 2016; Grise and Davis 2020). Interestingly, the widening of the HC of about 0.2° per decade in both hemispheres has been observed since the late 1970s (Hu and Fu 2007; Nguyen et al. 2013; Rollings and Merlis 2021; Staten et al. 2019). This trend is due to enhanced static stability in the tropics (e.g., Lu et al. 2007) and stratospheric ozone depletion in the Southern Hemisphere (SH; Son et al. 2010; Kim et al. 2017), as well as the role of natural variability (Grise et al. 2019).

For the strength of the HC, weakening is projected to be another important feature of the HC in the future climate (Chemke and Polvani 2019; Hu et al. 2018; Liu et al. 2012). According to Chemke and Polvani (2019), CMIP5 models, forced by climate forcings (e.g., greenhouse gas concentrations and aerosols) of the Representative Concentration Pathway 8.5 scenario, show a statistical reduction in HC strength of less than 0.1% per year, i.e., approximately $-6.8 \times 10^7 \text{ kg s}^{-1} \text{ yr}^{-1}$ during 1979–2017 (Fig. 1 in Chemke and Polvani 2019). A similar trend has been obtained in their study from large ensemble simulations of the Community Earth System Model, while observations in recent decades show a trend with the opposite sign, that is, a strengthening of the HC. This discrepancy in HC between the climate models and observations may lead to inaccurate projections of future climates because of the immense role of the HC in global energy, momentum, and hydrological cycles. Therefore, it is important to understand how the HCs are represented in climate models.

The National Institute of Meteorological Sciences/Korea Meteorological Administration (NIMS/KMA) has recently developed the KMA Advanced Community Earth system model (K-ACE) by collaboration with the UK Met Office. The K-ACE is a coupled atmosphere–ocean–land–sea ice model that is designed to provide multi-decadal climate simulations as for other CMIP6 models. A detailed description of the K-ACE is presented in Lee et al. (2020), such as the model components and coupling techniques. A brief introduction to the K-ACE is given in the following. The K-ACE consists of the Unified Model Global Atmosphere 7.1 (Walters et al. 2019), the Modular Ocean Model (MOM; Griffies et al. 2004), the Joint UK Land Environment Simulator (Best et al. 2011), and the Los Alamos Sea Ice Model (Hunke et al. 2015) through the coupler OASIS3-MCT. The key difference between the K-ACE and the UK Earth System Model is the ocean component model and the representation of aerosols, among many others. Sung et al. (2021) documented the performance of the K-ACE for climate mean states (e.g., surface temperature, precipitation, and sea ice extent) and climate variability (e.g., Northern and Southern

Annular Modes) in the present and future climate. However, it has not been reported whether the HC is reasonably represented in the model or how it is projected to the changes under the future scenarios.

Therefore, the objective of the present study is to evaluate the HC in historical and future climate scenario simulations using the K-ACE. For the historical simulation, the annual mean and seasonal means of the HC are compared against those from the reanalysis. Possible causes of the model bias are examined in terms of the diagnostic equation of the zonal mean meridional streamfunction. For the future projection, the K-ACE simulation forced by the Shared Socioeconomic Pathway 5-Representative Concentration Pathway 8.5 (SSP5-8.5) scenario is compared with corresponding CMIP6 model simulations in terms of the extent and strength of the HC.

The rest of this paper is organized as follows. Section 2 describes the observational and model simulation data and analysis methods. Sections 3 and 4 present evaluations of the historical and future scenario simulations of the K-ACE, respectively. Discussion and Conclusions and discussion are provided in Section 5.

2 Data and Methods

2.1 Data

We employ three ensemble members of the historical (years 1850–2014) and the SSP5-8.5 scenario (years 2015–2100) simulations of the K-ACE. Both simulations are conducted following the experimental design of the CMIP6 (O'Neill et al. 2016), which is described in detail by Sung et al. (2021). To diagnose the HC, monthly fields of meridional wind, surface pressure, and precipitation rate are analyzed. Also, daily fields of horizontal wind and temperature are used to measure the eddy fluxes. The horizontal resolution of the data is 1.875° longitude \times 1.25° latitude, however, for comparison with the reanalysis, the resolution of the data is reduced to 2.5° longitude \times 2.5° latitude through bilinear interpolation. Vertically, the monthly and daily data have 19 and 8 levels, respectively, up to the 50-hPa pressure level.

Furthermore, to provide a comparison with other climate models participating in the CMIP6, we additionally analyze the nine climate models listed in Table 1. The models were chosen simply because they had data available from the early phase of the present study. For the intercomparison between the models and the reanalysis, bilinear horizontal interpolation is applied to conform the grid to a resolution of 2.5° longitude \times 2.5° latitude resolution. For the computation of the HC extent, the model data were used at their original resolution for better accuracy. Regional Hadley cells of the

Table 1 The global climate models and the spatial resolution used in this study. One ensemble member, which is arbitrarily selected, is used for each model

CMIP6 Model	Horizontal Resolution (° lon × ° lat)
ACCESS-CM2	1.875° × 1.25°
ACCESS-ESM1	1.875° × 1.25°
CanESM5	2.8125° × 2.7906°
FGOALS-g3	2.25° × 2.0°
HadGEM3-GC31-LL	1.875° × 1.25°
IPSL-CM6A-LR	2.5° × 1.2676°
MIROC6	1.40625° × 1.40625°
MRI-ESM2-0	1.125° × 1.125°
UKESM1-0-LL	1.875° × 1.25°

CMIP6 are presented at 1.875° longitude × 1.25° latitude to match that of the K-ACE.

We use the European Centre for Medium-range Weather Forecasts (ECMWF) interim reanalysis data (ERA-Interim; Dee et al. 2011). As for the model simulations, monthly variables are taken for 1979–2014. It should be noted that the period of the ERA-interim is much shorter than that of the historical run, and therefore, we match the time length of the data when the data are compared. The horizontal resolution of the ERA-interim is also set to 2.5° longitude × 2.5° latitude, and 19 and 8 vertical levels out of 37 vertical levels are taken to match that of the monthly and daily model simulations, respectively.

2.2 Definition of the HC

The HC is defined as the zonal mean meridional streamfunction (Ψ):

$$\Psi(p, \phi) = \frac{2\pi a \cos(\phi)}{g} \int_p^{p_s} \bar{v}(p, \phi) dp, \tag{1}$$

where ϕ is the latitude, p is the pressure, p_s is the surface pressure, a is the radius of the earth, g is the gravitational acceleration, and v is the meridional wind. Here, the overbar represents the zonal mean. The strength and extent of the HC are quantified by the cross-section of the Ψ at 500 hPa (Ψ_{500}); that is, the strengths of the NH (SH) HCs are defined by the maximum (minimum) value of Ψ_{500} over the latitudes between 15°S–30°N (15°N–30°S). The extent of the NH (SH) HC is defined by the latitude of sign reversal in Ψ_{500} over the latitudes 20°N–50°N (20°S–50°S).

For the regional definition of the HC, we employ the meridional streamfunction of the divergent horizontal wind (Ψ_d ; Schwendike et al. 2014):

$$\Psi_d(p, \phi, \lambda) = \frac{2\pi a \cos(\phi)}{g} \int_p^{p_s} v_d(p, \phi, \lambda) dp, \tag{2}$$

where λ is the longitude. The underlying assumption for this regional HC definition is that the vertical motion in the tropical atmosphere can be decomposed into a zonal Walker circulation component and a meridional HC component.

2.3 Diagnostic Equation for the HC

To examine the cause of model bias, we examine the terms in the diagnostic equation of the zonal mean meridional streamfunction, which can be obtained through the quasi-geostrophy assumption (Chapter 10 of Holton and Hakim 2013):

$$\frac{\partial^2 \bar{\Psi}}{\partial y^2} + \frac{f_0^2}{N^2} \rho_0 \frac{\partial}{\partial z} \left(\frac{1}{\rho_0} \frac{\partial \bar{\Psi}}{\partial z} \right) = \frac{\rho_0}{N^2} \left[\frac{\partial}{\partial y} \left(\frac{\kappa \bar{J}}{H} - \frac{R}{H} \frac{\partial \overline{v'T'}}{\partial y} \right) - f_0 \left(\frac{\partial^2}{\partial z \partial y} \overline{u'v'} - \frac{\partial \bar{X}}{\partial z} \right) \right], \tag{3}$$

where the prime denotes the deviation from the zonal mean. This equation accounts for the fact that the zonal mean meridional circulation ($\bar{\Psi}$) is determined by the zonally averaged diabatic heating (\bar{J}), the eddy heat flux ($\overline{v'T'}$), and the eddy momentum flux ($\overline{u'v'}$), while the zonal drag force (\bar{X}) is neglected. Because of the elliptic operator to the streamfunction in the left-hand side, one may expect that $\bar{\Psi} > 0$ over the regions where the diabatic heating decreases poleward, the eddy heat flux is minimum, and the eddy vorticity flux decreases with height. As a representation of the diabatic heating, the precipitation rate will be examined.

3 The Hadley Circulation in the Historical Simulation

3.1 Climatology

Starting with the historical simulation, we first examine whether the time mean HCs are well represented in the K-ACE. Note that for comparisons with the ERA-Interim, we use the ensemble mean of the simulations for the period between 1979–2014 in this section. When averaged for all seasons (ANN; Fig. 1a), the upward branch of the HCs is located near 5°N, while downward branches are located at approximately 30°S and 30°N. The maxima of the HCs are reached near $\pm 10^{11}$ kg s⁻¹ at pressure levels between 500 and 700 hPa. For December–February (DJF), the streamfunction exhibits a strengthened NH winter HC between 10°S–30°N,

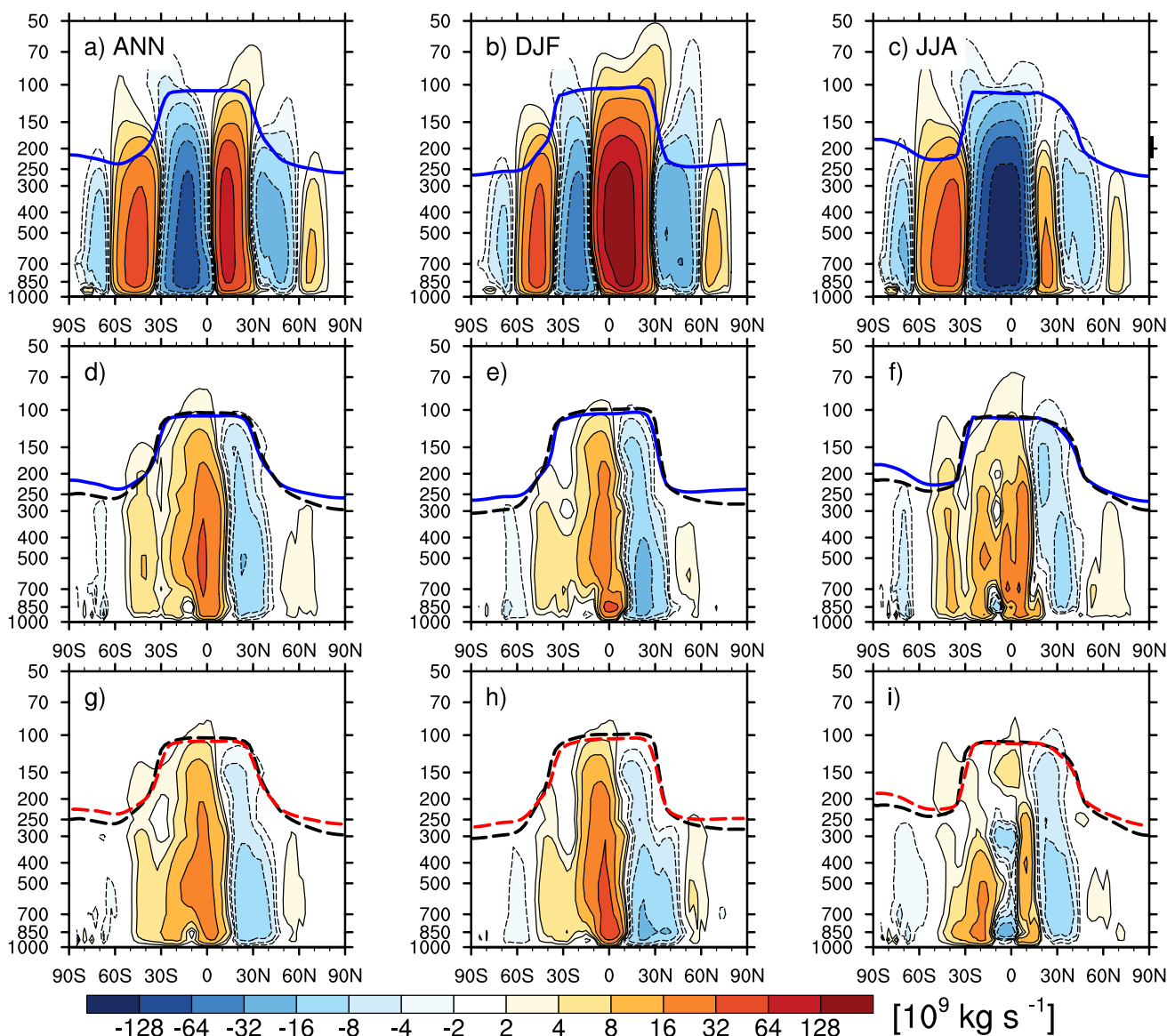


Fig. 1 The zonal mean meridional streamfunction $\bar{\Psi}$ of the historical simulation of the K-ACE are time-averaged during years 1979–2014 for (a) ANN, (b) DJF, and (c) JJA. The values obtained by subtracting the corresponding means of the ERA-interim from those of the K-ACE are shown in the middle row (d–f) and subtracting the means

of ERA-interim from those of the CMIP6 are in the bottom (g–i). The unit for the streamfunction is 10^9 kg s^{-1} . Superposed are the tropopause heights according to the WMO definition for the K-ACE (solid blue curve), for the ERA-interim (dashed black curve) and the CMIP6 (red dashed curve) of the corresponding seasons

with a maximum value exceeding $2.1 \times 10^{11} \text{ kg s}^{-1}$ (Fig. 1b). For June–August (JJA), the SH HC extends up to 15°N , and its strength is about $-2.2 \times 10^{11} \text{ kg s}^{-1}$ (Fig. 1c). Vertically, the HCs are roughly bounded by the tropopause height (solid blue curves in Fig. 1), which is determined by World Meteorological Organization (WMO) definition as the lowest level at which the lapse rate decreases to 2 K km^{-1} or less. As such, the location, extent, and strength of the HC match reasonably well with the reanalysis, with clear hemispheric symmetry and asymmetry for the annual and seasonal means, respectively (Dima and Wallace 2003).

To elaborate the bias of the HCs in the K-ACE, we take the time mean of the K-ACE minus that from the ERA-Interim. The difference for all seasons exhibits in general positive and negative anomalies in the tropics of the SH and NH, respectively. This indicates that the strength of both HCs in the NH and SH are underestimated, although it is more obvious for the SH HC (Fig. 1d). For the extent, the positive anomaly, centered at the equator, suggests that the upward branch of the HCs is shifted southward. The negative anomaly, near 25°N , indicates that the northward extent of the NH HC is reduced. These biases in the strength and extent can be clearly seen

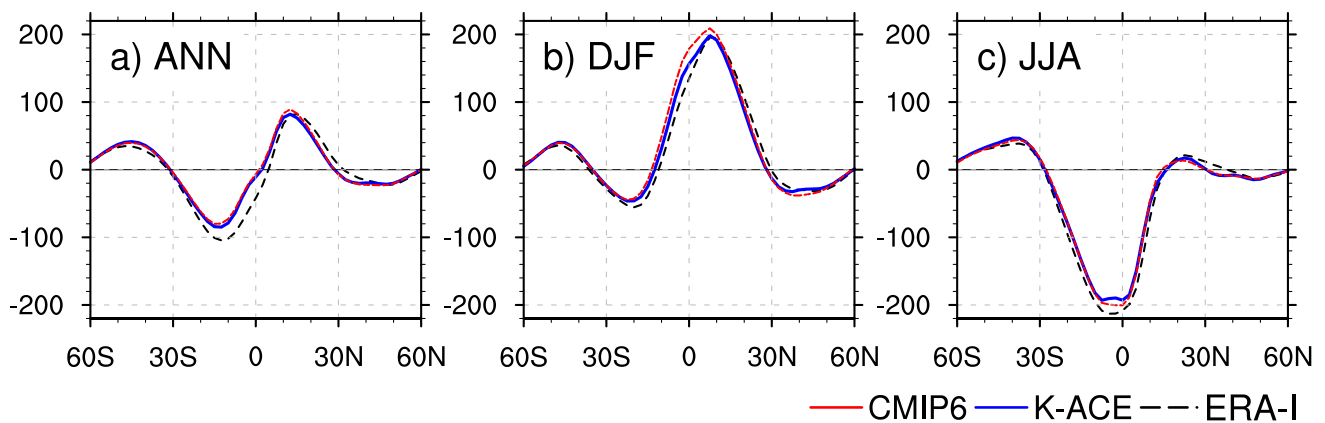


Fig. 2 The zonal mean meridional streamfunction $\bar{\Psi}$ time-averaged as for Fig. 1, except that the cross-section at 500 hPa is shown. The time-averaging is done for (a) ANN, (b) DJF, and (c) JJA. The

unit for the streamfunction is 10^9 kg s^{-1} . The result for the historical K-ACE is shown by the solid blue line, while that for the ERA-interim and CMIP6 are shown by the dashed black and red line, respectively

from the cross-section at 500 hPa pressure level (Fig. 2a). The SH HC of the K-ACE (solid blue curve in Fig. 2a), which has strength of about $-8.49 \times 10^{10} \text{ kg s}^{-1}$, is underestimated by up to approximately 20% compared to that of the ERA-interim (dashed black curve in Fig. 2a), of which strength is about $-10.45 \times 10^{10} \text{ kg s}^{-1}$. We also find that both the southern and northern edges of the NH HC are shifted southward by 1.4 and 1.6 degrees, respectively.

The biases in the annual mean HCs can be accounted for by the biases in the seasonal mean HCs (Figs. 1e, f and 2b, c). The underestimation of the strength of the SH HC takes place primarily during JJA, particularly between 10°S – 0° (Figs. 1f and 2c), although a weakened cell can also be found during DJF near 20°S (Figs. 1e and 2b). The JJA zonal mean streamfunctions of the K-ACE and ERA-Interim cross the zero line at 28.2°S and 28.8°S , respectively (Fig. 2c), indicating 0.6° equatorward bias. During DJF, however, the extent of the K-ACE is 28.1°N , which is substantially located south than that of the ERA-Interim (29.7°N), showing 1.6° bias (Fig. 2b). The underestimation of the strength of the SH HC takes place primarily during JJA, particularly between 10°S – 0° (Figs. 1f and 2c), although a weakened cell can also be found during DJF near 20°S (Figs. 1e and 2b). We note that the extent and strength bias can be similarly seen from the nine CMIP6 models (Figs. 1g–h and the red lines in Fig. 2).

To examine whether the HC biases are accompanied by biases in the zonal mean temperature, the zonal mean temperature field and its meridional gradient are examined in Fig. 3. Previous studies have found a linear correlation between the poleward extent of the HC and the meridional gradient of zonal mean temperature over the hemisphere (Choi et al. 2019; Wilcox et al. 2012). However, we find warm and cold biases in the tropical and polar near tropopause, respectively, indicating enhanced meridional

temperature gradient in the K-ACE. During DJF, for instance, the lower stratosphere in the Arctic shows a strong cold bias, implying steepened meridional temperature gradient. To be consistent with the previous studies, this should have contributed to widened NH HC, which turned out to be the opposite (Fig. 2b). Thus, in the next section, we use the diagnostic equation and examine the role of eddy fluxes and diabatic heating in the HC formation.

3.2 Diagnostic Equation

Having shown the HC climatology, we now investigate possible causes of the biases through the use of diagnostic Eq. (3). That is, the roles of the time mean eddy momentum flux ($\overline{u'v'}$), eddy heat flux ($\overline{v'T'}$), and diabatic heating (\overline{J}) in HC strength and extent biases are examined.

The time mean eddy momentum and eddy heat fluxes of the K-ACE simulation are presented in Fig. 4 (contours). The difference in the fluxes between the K-ACE and the ERA-interim is shown by shading. For the eddy momentum flux (Fig. 4a–c), the spatial structure and seasonal variation are reasonably captured by the K-ACE with some location and strength biases. The momentum flux is the maximum at the tropopause height with flux convergence between 45° – 50° region in both hemispheres (Fig. 4a,b). On the other hand, the eddy heat flux in the SH does not show a realistic structure, although in the NH it is reasonably centered at 40° – 50°N in the lower and the upper troposphere, except for JJA, (Fig. 4d–e). However, the heat flux bias is in a distance from the tropics, which may not have large influence on the HC.

Unlike the eddy heat flux, the momentum flux bias occurs in the tropics and subtropics. Compared to the ERA-interim, the eddy momentum flux is underestimated in the subtropics by approximately 11–16% in the annual and DJF

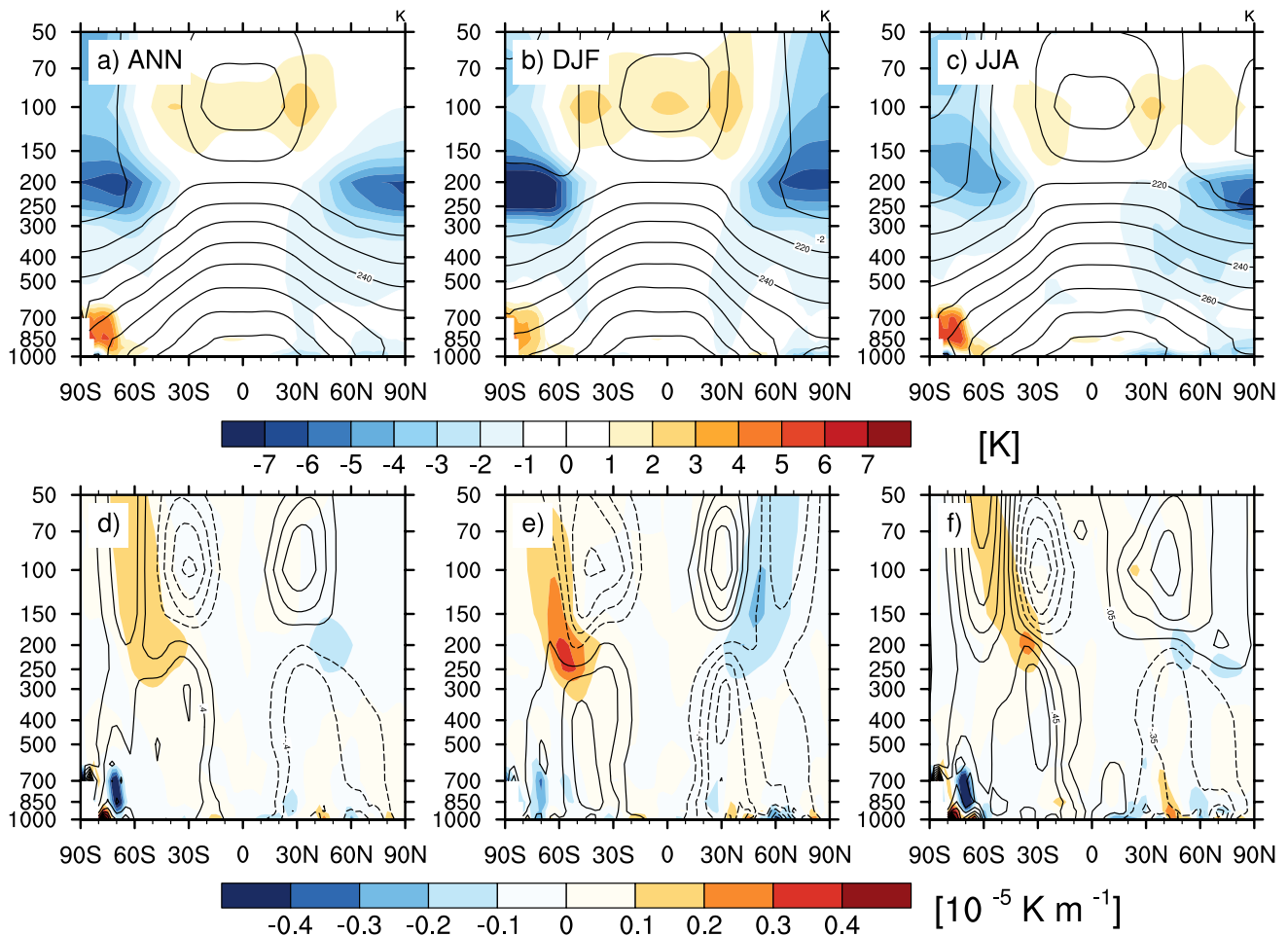


Fig. 3 The time-averaged zonal mean temperature (top) and its meridional gradient (bottom) of the historical K-ACE between 1979–2014 are shown (the contours). The contour of the zonal mean

temperature and its gradient are 10 K and 0.2 K m^{-1} , respectively. The difference between the K-ACE and the ERA-interim is shown by shading

means (Fig. 4a,b). The weakened eddy momentum flux during DJF, and hence anomalous convergence near 30°N (Fig. 4b), contribute to the weakening of the HC in the K-ACE. During JJA, an underestimation of approximately 10% can be observed in the tropical momentum convergence (shaded area in Fig. 4c). For the HC extent bias, the difference fields in the NH for all seasons and DJF (shaded area in Fig. 4a,b) are located to the north of the eddy momentum flux in the ERA-Interim (contours in Fig. 4a, b), implying that their contribution may not dominate the southward shift bias in the HC. Similarly, for JJA, the contribution of the momentum flux to the extent bias of the HC is weak because the difference field is jointly located with the climatology (Fig. 4c).

To estimate the role of diabatic heating (i.e., \bar{J} in Eq. (3)) in the HC bias, we examine the time mean zonal mean precipitation rate (Fig. 5). The K-ACE captures the precipitation maxima in the intertropical convergence zone (ITCZ) near 5°N and in the South Pacific convergence zone (SPCZ) near 10°S . As has previously been observed for climate

models (e.g., Tian and Dong 2020), the amount of precipitation at the SPCZ (solid blue curve in Fig. 5) is exaggerated in the K-ACE compared to that from the ERA-interim (dashed black curve), while the opposite is true for the ITCZ. For DJF, the mean difference in precipitation rate over the latitudes 15°S – 15°N between the model simulation and the reanalysis is approximately $-0.20 \text{ mm day}^{-1}$, which may not result in a net change in the HC strength. However, for JJA, the deficit at the ITCZ is greater than the exaggeration at the SPCZ, which leads to the weakened HC in the K-ACE. Furthermore, the peaks at ITCZ and SPCZ are shifted equatorward, contributing to the HC extent bias in the model.

4 The HC in the SSP5-8.5 Simulation

To evaluate the future projection of the HC strength and extent using the K-ACE, we first compare the time-averaged zonal mean streamfunction of the SSP5-8.5 simulation (years

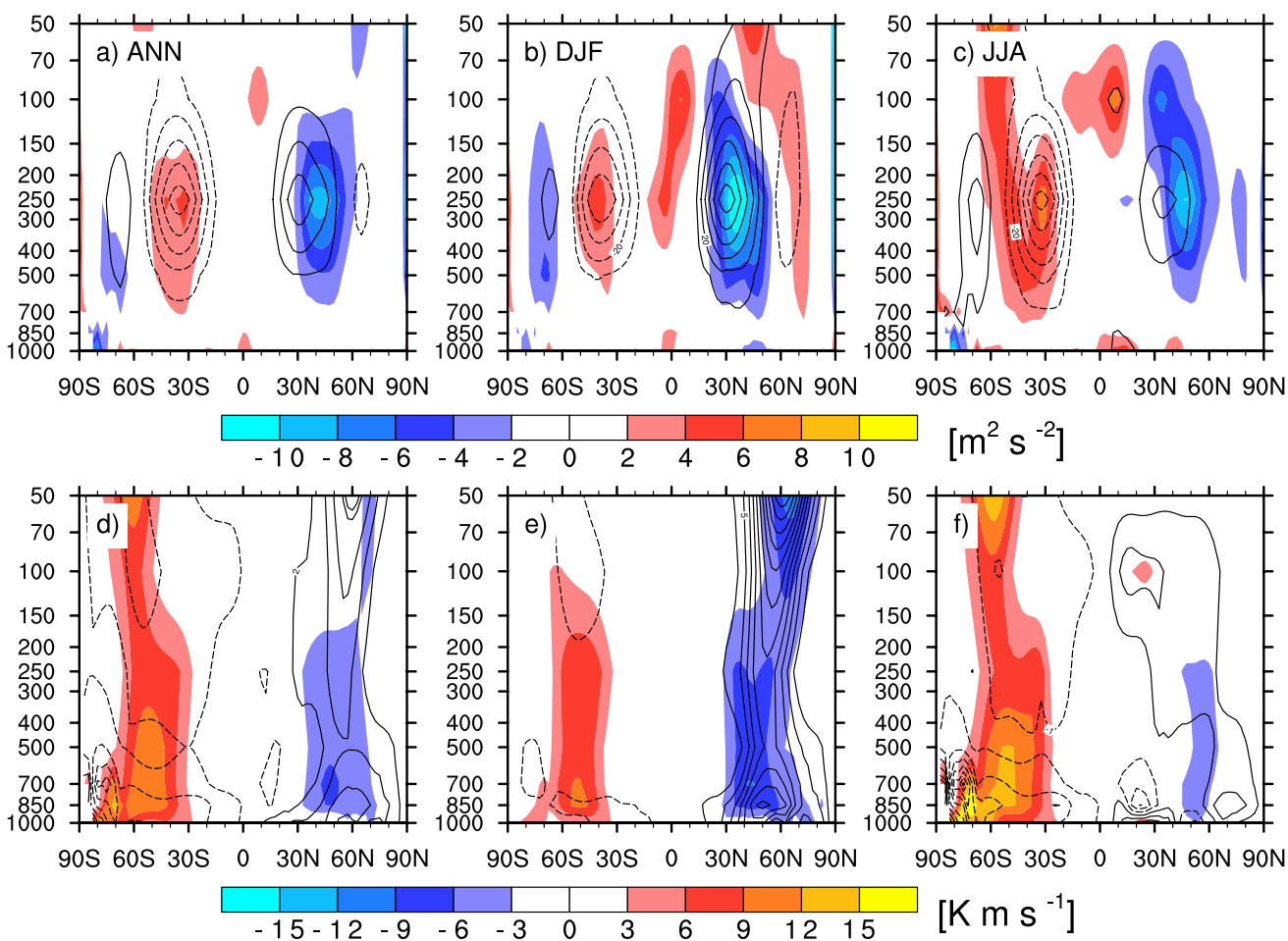


Fig. 4 The same as Fig. 1, except that the time-averaged (top) eddy momentum and (bottom) eddy heat fluxes of the historical K-ACE between 1979–2014 are shown (the contours). The contour intervals

for the momentum and heat fluxes are $10 \text{ m}^2 \text{ s}^{-2}$ and 2.5 K m s^{-1} , respectively. The difference between the K-ACE and the ERA-interim is shown by shading

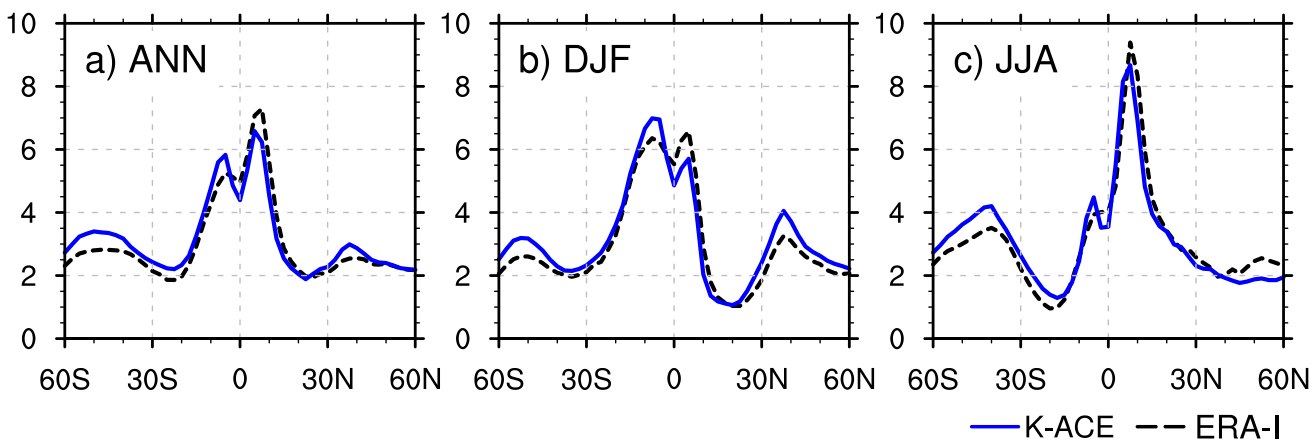


Fig. 5 The same as for Fig. 2, except that the time-averaged zonal mean precipitation (mm day^{-1}) is shown for (a) ANN, (b) DJF, and (c) JJA. The values for the historical K-ACE are shown by the solid blue line, while those for the ERA-interim are shown by the dashed black line

2015–2100) with that of the historical simulation (years 1979–2014) (Fig. 6). As was reported in previous studies (e.g., Chemke and Polvani 2019), weakening of the HC by approximately 10–15% can be identified; for the all-season mean, the difference between the SSP5-8.5 and the historical simulations shows negative and positive streamfunction anomalies in the tropics of the NH and SH, respectively (Fig. 6a). For both DJF and JJA, the weakening of the HC in the tropics is evident, as indicated by the negative and positive anomalies, respectively (Fig. 6b, c). Interestingly, the anomalies change the sign across the tropopause near 250 hPa, suggesting the HCs vertically stretched in the future, which is consistent with the rise of tropopause height (i.e., annual mean tropopause height between 20°S–20°N being approximately 100.5 hPa for the SSP5-8.5, in contrast to that of the historical being 106.8 hPa). Such a baroclinic structure can also be found from the CMIP6 models (not shown). For the extent, the widening of the HC can be observed from the streamfunction anomalies in the subtropics. For example, the positive anomaly at 30°N during DJF suggests a poleward expansion of the NH HC (Fig. 6b); similarly, the negative anomaly at 30°S during JJA displays the poleward expansion of the SH HC (Fig. 6c).

The temporal evolution of the HC weakening under the SSP5-8.5 scenario can be examined through the HC strength measured by the 500-hPa zonal mean streamfunction maximum. In Fig. 7, the HC strengths for three ensemble members of the K-ACE (dashed blue lines) and its ensemble mean (solid blue lines) are compared with those of nine CMIP6 models (black lines) and their multi-model mean (red lines). Overall, weakening can be observed in the K-ACE for both January (top panels in Fig. 7) and July (bottom panels in Fig. 7), and the weakening trend resembles that of the CMIP6 models. For example, during January, the boreal winter HC in the NH from

the K-ACE simulations shows a weakening with the slope of $-0.75 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$ for the entire period, which is comparable to the slope of the CMIP6 ($-0.90 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$) (Fig. 7a). These changes become steeper by the late 2000s. Winter NH HC in the K-ACE has negative trends of $-0.9 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$ for 50 years of 1950–2000. During 2050–2100, the trend is drastically decreased to $-2.94 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$. Such a steep trend between 2050–2100 is also seen in CMIP6 (-0.23 to $-2.96 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$). For the austral winter HC in the SH, the weakening in the K-ACE occurs at a rate of $0.47 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$, while it occurs at $0.50 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$ in CMIP6 (Fig. 7d). For the NH summer and SH summer HCs, a weakening trend can be observed in both the K-ACE and the CMIP6 (Fig. 7b, c). Like the boreal winter NH HC, the trend in the SH winter HC has steepened from $0.23 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$ for 1950–2000 to $1.81 \times 10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$ for 2050–2100.

The widening of the HC in the SSP5-8.5 scenario shows large inter-model variations. Figure 8 exhibits the temporal evolution of the poleward edge of the HC for the NH (panel a) and the SH (panel b). Note that the 24-year running mean is applied to the poleward latitude of the streamfunction sign reversal, which is computed from the annual mean zonal mean streamfunction at 500 hPa. The NH HC edge in K-ACE (Fig. 8a) is projected to expand from approximately 28.5°N to 29.4°N in the second half of the twenty-first century (blue line in Fig. 8). This is similar to the multi-model mean of the CMIP6, which shows an expansion from approximately 28.6°N to 29.6°N (red). An expansion of the HC can also be seen in the SH (Fig. 8b). However, for the K-ACE, the expansion ranges from near 30.5°S to 33.4°S (blue line), substantially exceeding the projected range obtained using the CMIP6 multi-model mean, which ranges from near 30.4°S to 32.3°S (red line). The widening trend increases from 0.04°

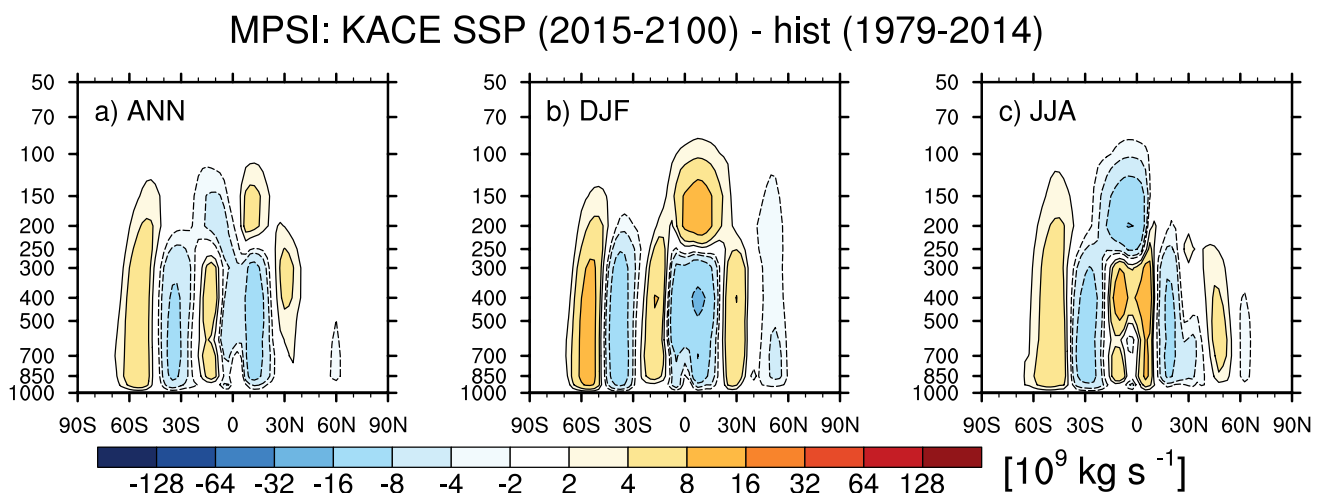


Fig. 6 The difference in zonal mean meridional streamfunction (10^9 kg s^{-1}) obtained by subtracting the values of the SSP5-8.5 run (2015–2100) from the current climate (1979–2014) of K-ACE and time-averaged for (a) ANN, (b) DJF, and (c) JJA

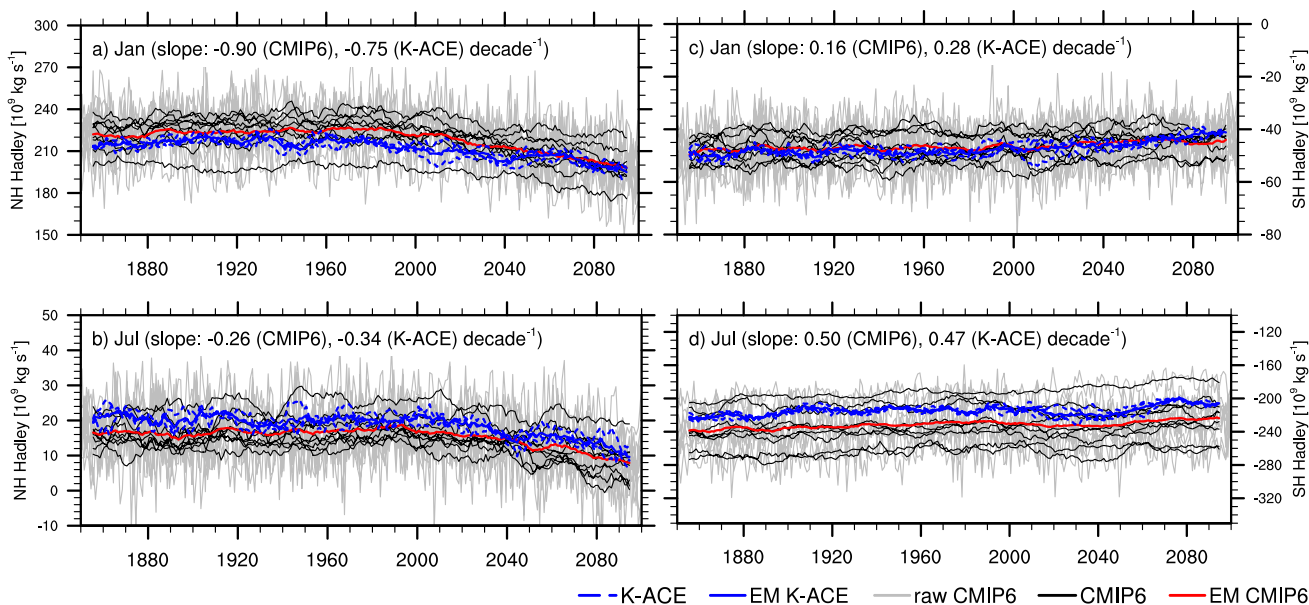


Fig. 7 The NH HC strength shown by the maximum values of the zonal mean meridional streamfunction Ψ at 500 hPa (a) between 5°S–30°N for January and (b) between 30°S–5°N for July. Similarly, the SH HC strength is shown by the minimum values (c) between 30°S–5°N for January and (d) between 5°S–30°N for July. The values for the three ensemble members of the K-ACE are shown by the dashed blue line, while the ensemble mean (EM K-ACE) is shown by the solid blue line. The results for the nine CMIP6 models are

shown by the black line, while their multi-model mean (EM CMIP6) is shown by the red line. To investigate long-term changes in the strength time series, the 12-year running average is applied to all time series, except for the original CMIP6 results, by the gray line. The unit for the streamfunction is shown in 10^9 kg s^{-1} . The linear trend for the K-ACE ensemble mean and CMIP6 multi-model mean is denoted in each panel with units of $10^9 \text{ kg s}^{-1} \text{ decade}^{-1}$

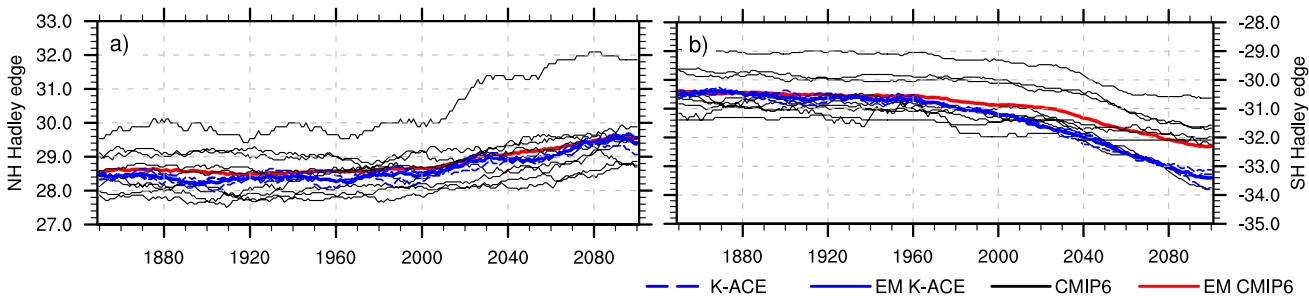


Fig. 8 The Hadley cell edge measured by the latitude of the sign reversal in the annual mean zonal mean streamfunction shown for (a) the NH HC and (b) the SH HC. To remove short-time scale noise, the 24-year running mean is applied. Blue dashed lines indicate the val-

ues of the three ensemble members of the K-ACE, and their ensemble mean is illustrated by the blue solid line. The results for the CMIP6 models and their multi-model mean are shown by black and red lines, respectively

($0.12^\circ \text{ decade}^{-1}$ of years 1950–2000 to 0.14° ($0.24^\circ \text{ decade}^{-1}$) of years 2050–2100 in NH (SH) HC.

Interestingly, the models that share similar component models with the K-ACE all show extensive widening of the SH HC in the SSP5-8.5 scenario. Those models are the UKESM, the HadGEM3-GC31-LL, and the ACCESS-CM2/ESM1, and they are shown by the thin black lines in Fig. 8b. The difference between the K-ACE and the HadGEM3-GC31-LL is the ocean model (i.e., MOM4 vs. NEMO), while it is the land model for the ACCESS-CM2/ESM1. The UKESM differs by the ocean model, in

addition to comprehensive treatment in carbon cycles and atmospheric chemistry. Although beyond the scope of this study, it would be an interesting question to study why these family models project substantial widening in the SSP5-8.5 scenario. If including all these models, we find that the HC expansion range obtained using the K-ACE still resides within the multi-model spread of the CMIP6 (black lines).

To identify regions where the amplified widening of the K-ACE SH HC takes place, we analyze the regional HC that is defined for each longitude. Figure 9 presents the meridional streamfunction for the historical simulations (years 1979–2014;

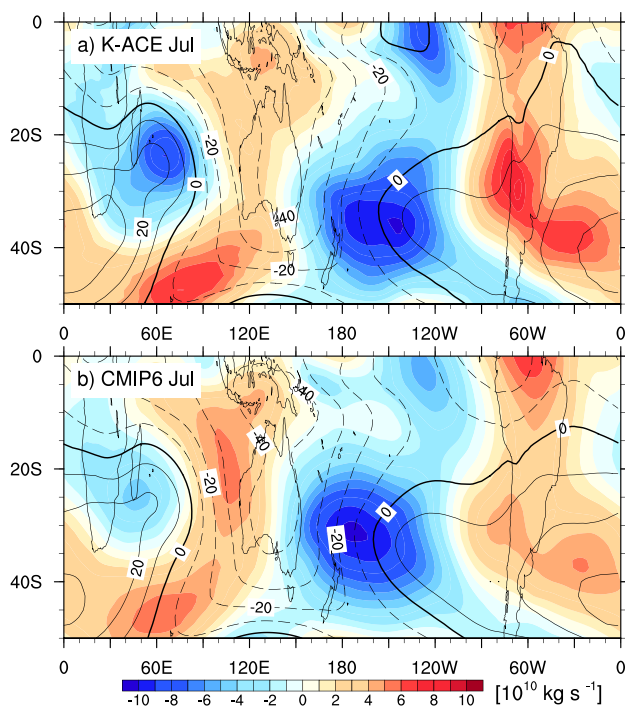


Fig. 9 The time mean meridional streamfunction at 500 hPa, defined by the divergent wind of the historical simulation (July in 1979–2014) using Eq. (2), is shown by the contour for the (a) K-ACE and (b) CMIP6. The contour interval is 10 kg s^{-1} . Positive (negative) values are indicated by solid (dashed) contours, and the thick black contours show zero values. The difference between the historical and SSP5-8.5 scenario simulations (July in 2015–2100) are shown by shading

contours) and its changes in the SSP5-8.5 scenario simulations (years 2015–2100; shaded). Note that the regional HC extents are defined by the latitude at which the meridional streamfunction based on the divergent horizontal wind (i.e., Eq. (2)) changes its sign (thick black contour in Fig. 9). The strong SH HC is located over 80°E – 150°W , and we can discern that its edge is placed south of 50°S in several regions, which also notes that it is hard to define the edge for each longitude. However, the strong negative anomaly is shown in the Pacific Ocean near 35°S , 150°W for both K-ACE and CMIP6 (Fig. 9a, b, respectively), indicating that the expansion for the July SH HC primarily occurs in the Pacific. However, the negative anomaly of the K-ACE is located further southward compared to that of the CMIP6 multi-model mean. In addition, the negative anomaly over the Indian Ocean near 20°S , 60°E is much stronger in the K-ACE than in the CMIP6 models, which contributes to the amplified widening in the K-ACE.

5 Discussion and Conclusions

In this study, we evaluated the HC in the historical and SSP5-8.5 scenario climate simulations using the K-ACE climate model. Overall, the K-ACE presented a reasonable

performance in terms of the HC strength and extent, as well as HC seasonality and future projection. However, several features were found in the K-ACE compared to reanalysis data and simulations by nine CMIP6 models:

- In the historical simulations, the K-ACE generates an NH HC that is shifted slightly southward, by 1 – 2° , for both its northern and southern extent, despite its reasonable strength.
- In the historical simulations, the SH HC strength is underestimated by approximately 20%, most pronouncedly during JJA.
- The biases in the HC extent and strength are consistent with the biases in the eddy momentum flux, which shows anomalous convergence near 5°N , and in the precipitation rate, which exhibits a weaker intensity in the ITCZ compared to observations.
- In the SSP5-8.5 scenario simulations, the weakening and widening of the HC are projected by the K-ACE. Compared to the CMIP6 models, the widening for the SH HC is higher by 1 – 2° in the second half of the twenty-first century.
- The overestimated widening in the wintertime SH HC of the K-ACE occurs primarily over the Indian and Pacific oceans.

However, our analysis using the diagnostic equation has a limitation, in that the contribution from each term to the HC formation is not computed quantitatively, although the connection between the HC bias and the eddy flux and precipitation bias is identified. The relative contribution of each term can be examined by solving the Laplacian equation, which requires performing extensive sensitivity tests on the boundary conditions when constructing the inverse Laplacian. We plan to investigate this problem in our follow-up study.

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Declarations

Conflict of Interest The authors declare that they have no conflict of interest.

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