

# Inorganic Nanoribbons with Unpassivated Zigzag Edges: Half Metallicity and Edge Reconstruction

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## ABSTRACT

We have investigated the electronic and structural properties of inorganic nanoribbons (BN, AlN, GaN, SiC, and ZnO) with unpassivated zigzag edges using density functional theory calculations. We find that, in general, the unpassivated zigzag edges can lead to spin-splitting of energy bands. More interestingly, the inorganic nanoribbons AlN and SiC with either one or two edges unpassivated are predicted to be half metallic. Possible structural reconstruction at the unpassivated edges and its effect on the electronic properties are investigated. The unpassivated N edge in the BN nanoribbon and P edge in the AlP nanoribbon are energetically less stable than the corresponding reconstructed edge. Hence, edge reconstruction at the two edges may occur at high temperatures. Other unpassivated edges of the inorganic nanoribbons considered in this study are all robust against edge reconstruction.

## KEYWORDS

AlN and SiC nanoribbons, half metallicity, unpassivated zigzag edge, edge reconstruction, density functional theory

## 1. Introduction

Graphene, a two-dimensional (2D) material, has attracted considerable research attention due to its unique physical properties such as massless Dirac fermion behavior [1–3], room-temperature quantum Hall effect, [4, 5] high carrier mobility and thermal conductivity, and high electronic coherence [6]. Moreover, when a quasi-one-dimensional (Q1D) nanoribbon is cut out from the 2D graphene sheet, the energy band gap of the graphene nanoribbon (GNR) is no longer zero but dependent on the width and crystallographic orientation of the GNR [7–12].

In particular, zigzag-edged GNRs (zGNRs) show distinct electronic properties due to their special localized edge states which are antiferromagnetically coupled [11]. Recent theoretical studies suggest that half metallicity can be also achieved in zGNRs [13–17]. If confirmed, the half metallicity could be exploited for future spintronic applications because the electric current can be fully spin polarized when going through a half metallic system [18–20].

In addition to GNRs, some other Q1D inorganic structures such as BN, AlN, GaN, SiC, and ZnO nanotubes and nanoribbons have been investigated by theoretical means [21–26]. In Ref. [21], Zheng et al.

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showed from first-principles calculations that the zigzag BN nanoribbon (zBNNR) can be half metallic if its N edge is unpassivated. It would be interesting to know whether the same electronic behavior can be seen in other inorganic nanoribbons with unpassivated edges. Nevertheless, unpassivated edges may be very unstable and, if so, edge reconstruction may take place, akin to structural reconstruction of the unpassivated edges of GNRs [27]. In this article, we present calculated electronic structures of a number of inorganic zigzag nanoribbons, including the main group III nitrides, the group III phosphides, SiC, and ZnO, based on density functional theory (DFT) calculations. We find that some of these nanoribbons with either one or two unpassivated edges can be half metallic. Possible edge reconstruction on all unpassivated edges is also investigated in detail. Some of the unpassivated edges are quite unstable, and edge reconstructions can take place above room temperature.

## 2. Computational methods

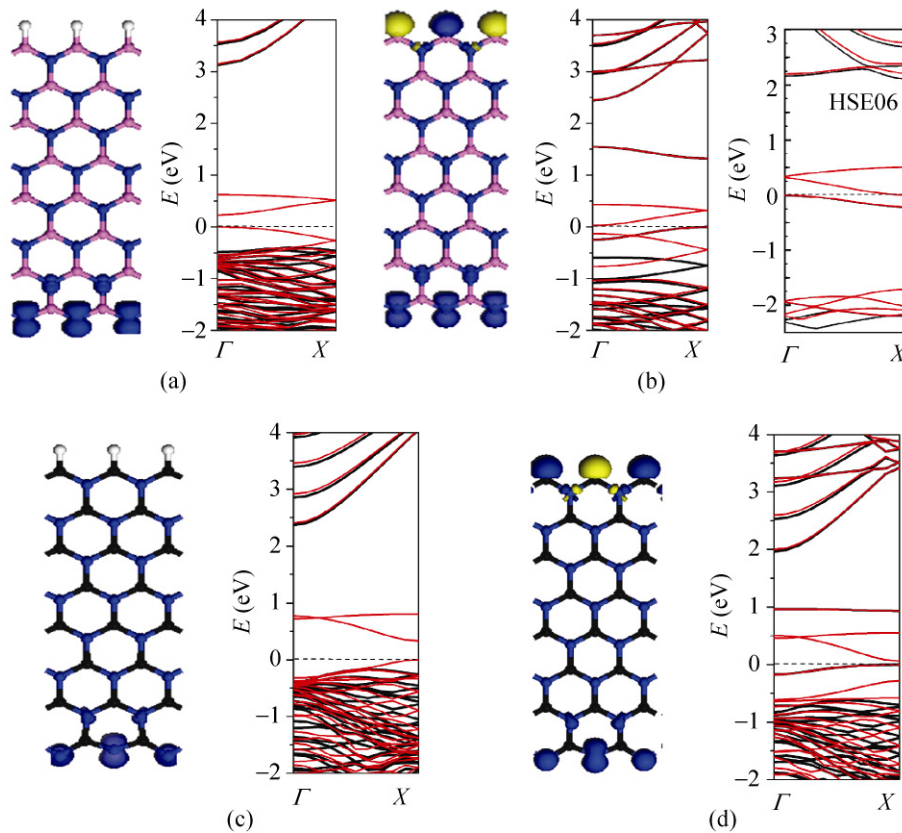
The DFT calculations were carried out using the DMol<sup>3</sup> package [28–30]. The generalized gradient approximation (GGA) in the Perdew–Burke–Ernzerhof (PBE) form, together with the all-electron double numerical (DNP) basis set with polarized functions were chosen for the spin-unrestricted DFT computations [31]. The real-space global cutoff radius was set to be 3.7 Å. To simulate the inorganic nanoribbons, a rectangular supercell with dimensions  $36 \times L \times 16 \text{ \AA}^3$  was built, where  $L$  was chosen to be twice the width of the hexagonal ring and hence depended on the nanoribbon concerned. The nearest distance between the nanoribbon and its nearest image (under periodic boundary conditions) was greater than 16 Å. For geometric optimization, the Brillouin zone was sampled with  $1 \times 20 \times 1$   $k$  points using the Monkhorst–Pack scheme [32]. After geometric optimization, the forces on all atoms were less than 0.0002 Ha/Å. To further confirm the predicted half metallicity for certain inorganic nanotubes, we re-examined their band structures using the screened exchange hybrid Heyd–Scuseria–Ernzerhof (HSE) [33, 34] functional and the 6-31G basis set (implemented in the Gaussian 09 package [35]).

## 3. Results and discussion

### 3.1 Half metallicity

We use the number of parallel zigzag chains  $n$  in a zigzag-edged nanoribbon to describe the width of the nanoribbon. Here,  $n = 8$  was chosen for all the nanoribbons considered. First, we focus on two group III nitrides, AlN and GaN. In the zigzag-edged AlN nanoribbon (zAlNNR), all the outermost atoms at the Al edge are Al atoms while all the outermost atoms at the opposing edge (i.e., the N edge) are N atoms. As shown in Fig. 1(a), when the Al edge of a zAlNNR is passivated by hydrogen atoms (designated as zAlNNR-HAl), the calculated band structure suggests that the zAlNNR-HAl is a spin-polarized semiconductor with a band gap of ~0.2 eV. When the H atoms on the Al edge are removed, however, the calculated band structure suggests that the pristine zAlNNR is a half metal, which is also confirmed by DFT calculations using the HSE functional. As can be seen in Fig. 1(b), the N edge is ferromagnetic while the Al edge is antiferromagnetic, similar to the pristine zBNNR [19]. For zGaNNR-HGa, the N atoms at the N-edge no longer lie on a single straight line along the ribbon direction, but become slightly modulated in a periodic fashion as shown in Fig. 1(c). The calculated band structure suggests that it is a spin-polarized semiconductor. When the H atoms on the Ga edge are removed, it is still a semiconductor, although the band gap is reduced to 0.05 eV as shown in Fig. 1(d).

For group III phosphides, we consider BP and AlP nanoribbons. As shown in Fig. 2(a), the band structure calculations suggest zBPNR-HB is metallic if the P edge is unpassivated. For zAlPNR-HAl, the P edge is antiferromagnetic and the system is a spin-unpolarized semiconductor as shown in Fig. 2(b). Besides the group III/V compounds, we also consider SiC and ZnO nanoribbons. As shown in Fig. 2(c), when the Si edge of zSiCNR-HC is unpassivated, the spin distribution at the two edges in the ground state is ferromagnetically coupled. The nanoribbon is predicted to be a half metal based on the calculated band structure using both the PBE and HSE functionals. For zZnONR-HZn, shown in Fig. 2(d), the spin distribution at the two edges is antiferromagnetically coupled, and the



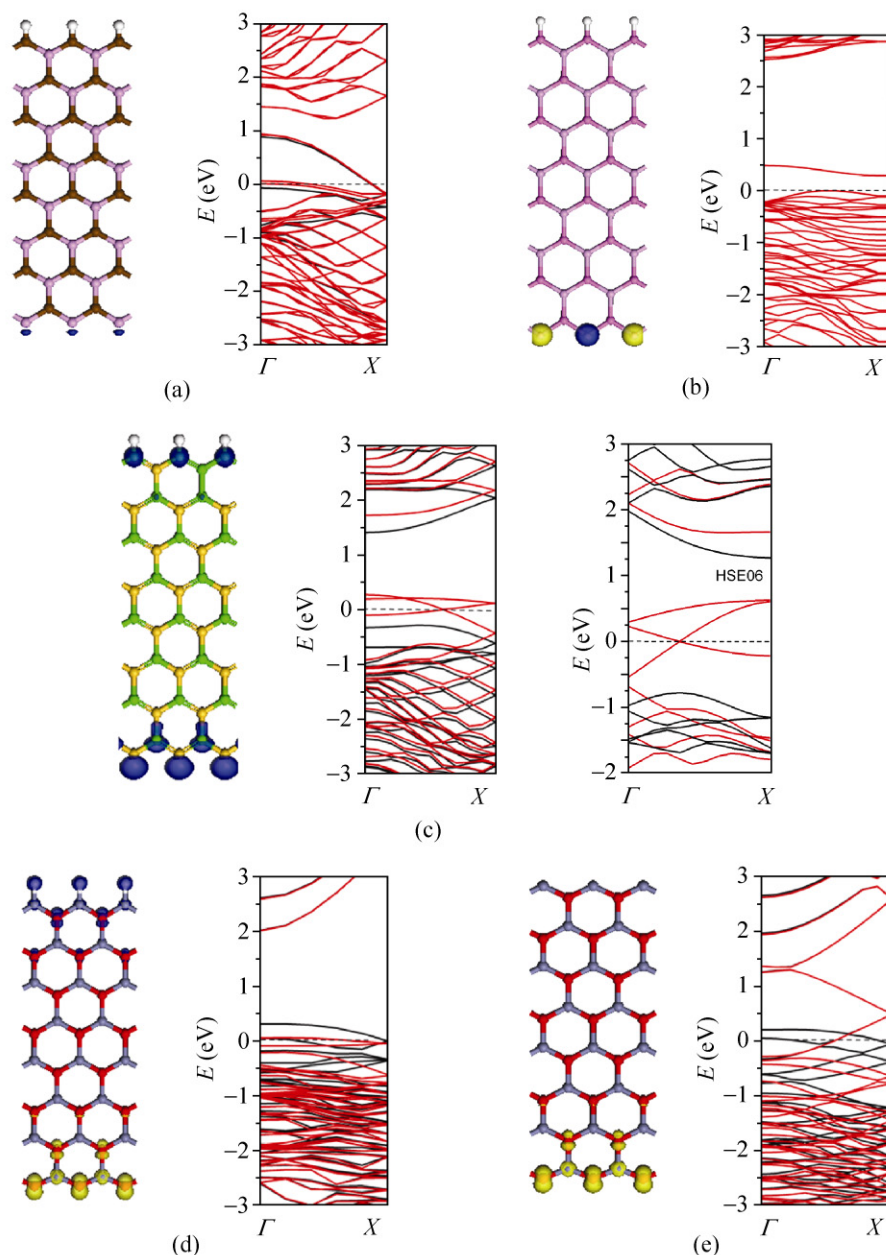
**Figure 1** The spin density (left panel) and spin-polarized band structures (right panel) of (a) zAlNNR-HAl, (b) pristine zAlNNR, (c) zGaNNR-HGa, and (d) pristine zGaNNR. Small balls in white, pink, blue, and black denote hydrogen, aluminum, nitrogen, and gallium atoms, respectively. For spin density, blue and yellow lobes represent spin-up and spin-down regions, respectively; and the range of the iso-values of spin density is  $[-0.03, 0.03] \text{ e}/\text{\AA}^3$ . In the band structure graphs, black represents the spin-up channel and red represents the spin-down channel. The Fermi level is denoted by the horizontal dashed line

nanoribbon is metallic. If all the H atoms at the Zn edge are removed, the nanoribbon is still metallic, but also becomes ferromagnetic as shown in Fig. 2(e). Such ferromagnetism has also been recently predicted to appear in ZnO nanowires [36].

### 3.2 Edge reconstruction

According to Ref. [27], in zGNRs, the reconstruction of an unpassivated zigzag edge may spontaneously take place at room temperature, where two neighboring hexagons can transform into a pentagon and heptagon. Here we investigate a similar edge reconstruction that may take place at the unpassivated edges of any of the inorganic nanoribbons discussed above. First, we consider the half metallic zBNNR-HB reported in Ref. [21]. As shown in Figs. 3(a)–3(c), the N edge has two ways of reconstruction: (1) it can reconstruct into a symmetric structure as shown in Fig. 3(b), where the

pentagon is formed by two N atoms and three B atoms. This way of edge reconstruction is denoted as the A style; (2) The N edge may also reconstruct into an unsymmetrical structure, as shown in Fig. 3(c), where each pentagon is formed by three N atoms and two B atoms. This way of edge reconstruction is denoted as the B style. After the edge reconstruction in the A style, the system is 1.06 eV/supercell higher in energy, while in the B style, the energy is lowered by 0.22 eV/supercell. Hence, the B style is energetically more favorable. Our calculations, using the nudged elastic band method [37, 38] implemented in DMol<sup>3</sup>, also showed that the activation barrier for the edge reconstruction in the B style is about 1.0 eV. Quantum molecular dynamics simulations (for 5 ps) of this system at 800 K were performed and no sign of spontaneous reconstruction of the N edge was observed, suggesting the edge reconstruction would occur at a very high temperature.



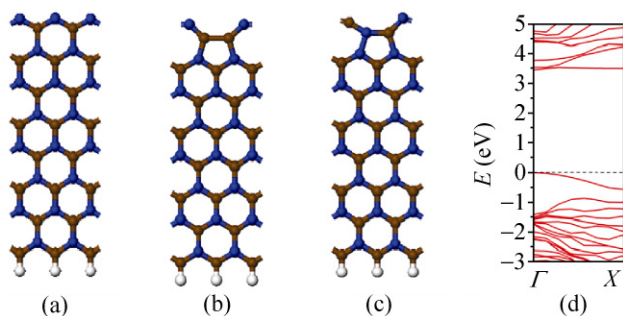
**Figure 2** The spin density and spin-polarized band structures of (a) zBPNR-HB, (b) zAIPNR-HAI, (c) zSiCNR-HC, (d) zZnONR-HZn, and (e) pristine zZnONR. Small balls in green, white, pink, blue, black, brown, purple, yellow, grey, and red denote carbon, hydrogen, aluminum, nitrogen, gallium, boron, phosphorus, silicon, zinc, and oxygen atoms, respectively. For spin density, blue, and yellow lobes represent spin-up and spin-down regions, respectively; the range of the iso-values of spin density is  $[-0.03, 0.03] \text{ e}/\text{\AA}^3$ . In the band structure graphs, black represents the spin-up channel and red represents the spin-down channel. The Fermi level is denoted by the horizontal dashed line

The calculated band structure for the reconstructed nanoribbon (in the B style) is shown in Fig. 3(d). Interestingly, after edge reconstruction, the nanoribbon is no longer a half metal but becomes a wide-gap semiconductor.

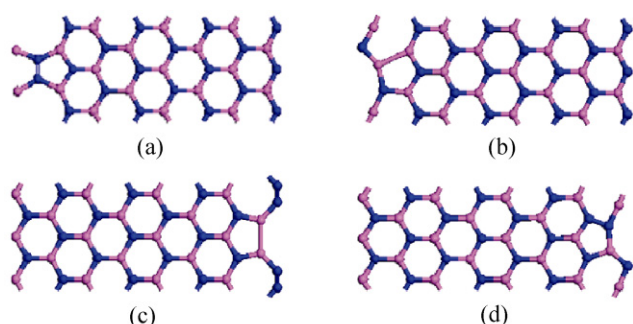
For the pristine zAlNNR, we first consider reconstruction of the Al edge. The edge reconstructions in

the A and B styles are shown in Figs. 4(a) and 4(b), respectively. The calculations show that after edge reconstruction in the A or B style the energy will be raised by 3.4 eV or 1.4 eV, respectively. Hence, the unpassivated Al edge appears to be robust against edge reconstruction. Figures 4(c) and 4(d) display the edge reconstruction at the N edge in the A and B style,





**Figure 3** The optimized structures of zBNNR-HB (a) before, and after the edge reconstruction in (b) the A style and (c) B style at the N edge. White, blue, and brown balls denote hydrogen, nitrogen, and boron atoms, respectively. The band structure of nanoribbon (c) is shown in (d)

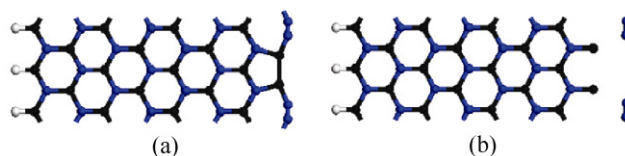


**Figure 4** The optimized structures of zAlNNR after edge reconstruction in (a) the A style and (b) the B style at the Al edge, and in (c) the A style and (d) the B style after reconstruction at the N edge. Blue and pink balls denote nitrogen and aluminum atoms, respectively

respectively. Our calculations show that upon edge reconstruction in the A style, the energy will be lowered by 0.26 eV while in the B style it will be increased by 0.64 eV. However, even though the edge reconstruction in the A style is energetically favorable, the activation barrier for the edge reconstruction is typically very high (~3.5 eV). Hence, it is unlikely for an N atom at the edge to swap position with its nearby Al atom. As a consequence, the unpassivated N edge can still be considered as being very stable against edge reconstruction.

In Table 1, we list the energy changes for edge reconstruction in the A and B styles at the unpassivated edge in zGaNNR-HGa, zBPNR-HB, zAlPNR-HAl, zSiCNR-HC, and zZnONR-ZnH. For zGaNNR-HGa, the pentagon and heptagon are not stable after edge reconstruction in the style A at the N edge because the N atoms will be converted into free N<sub>2</sub> molecules, as shown in Fig. 5. The energy change for the edge reconstruction of zZnONR-HZn in the B style is not given in Table 1, since the unpassivated O edge recovers its original zigzag structure after the geometric optimization, when the edge reconstruction is initially imposed on it.

As shown in Table 1, only the energy of zAlPNR-HAl is lowered on edge reconstruction in the B style. Moreover, our calculations show that the activation barrier for edge reconstruction in this case is about 0.65 eV, implying that spontaneous edge reconstruction is feasible at room temperature. The normal bond lengths of Al–Al, Al–P, and P–P are 2.22 Å, 2.21 Å, and 2.20 Å, respectively, which are nearly the same. This observation is similar to the case of zBNNR-HB where the normal bond lengths of B–B, B–N, and N–N are nearly the same (1.42 Å, 1.45 Å, and 1.48 Å, respectively). Considering the edge reconstruction of GNR reported in Ref. [27], it seems that edge reconstruction in the B style tends to occur at the unpassivated edge with two types of constituent atoms giving rise to nearly the same bond lengths between them.



**Figure 5** For the edge reconstruction of (a) zGaNNR-HGa in the A style, (b) the edge N atoms form free N<sub>2</sub> molecules after geometric optimization. White, blue, and black balls denote the hydrogen, nitrogen, and gallium atoms, respectively

**Table 1** Energy changes associated with edge reconstruction in the A and B styles at the unpassivated edge in zGaNNR-HGa zBPNR-HB, zAlPNR-HAl, zSiCNR-HC, and zZnONR-HZn

Reconstruction	zGaNNR-HGa	zBPNR-HB	zAlPNR-HAl	zSiCNR-HC	zZnONR-HZn
A style (eV)	-2.92	1.03	0.18	1.19	2.13
B style (eV)	0.11	1.00	-0.015	1.35	N/A



## 4. Conclusions

We have shown that the unpassivated zigzag edges of the inorganic nanoribbons (BN, AlN, GaN, SiC, and ZnO) can lead to spin-splitting energy bands in general. More interestingly, the inorganic nanoribbons AlN and SiC with either one or two edges unpassivated are predicted to be half metals. Possible structural reconstruction at the unpassivated edges and its effect on the electronic properties have been investigated. The unpassivated N edge in the BN nanoribbon and the P edge in the AlP nanoribbon are energetically less stable compared to the reconstructed edge. Hence, edge reconstruction is possible at high temperatures. Other unpassivated edges in the inorganic nanoribbons considered here are robust against edge reconstruction.

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