CORRECTION



Correction to: Uniformly Compressing Mean Curvature Flow

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In the sketch of the proof of Theorem 4 we claimed that a weak solution (ξ, κ, σ) can be obtained by following the same lines as in the proof of [1, Proposition 3.4 and Theorem 3]. However, an appropriate argument should be slightly different due to different boundary conditions. The aim of this addendum is to fill that gap.

More precisely, in [1] at the free end one has $\sigma^{\epsilon}(t, 0) = \kappa^{\epsilon}(t, 0) = 0$ for all $t \in [0, T]$. Therefore, the uniform (w.r.t. to ϵ) L^2 bound for $(\sigma^{\epsilon}, \kappa^{\epsilon})$ follows directly from the Poincaré inequality and the uniform L^2 bound for $(\partial_s \sigma^{\epsilon}, \partial_s \kappa^{\epsilon})$. In our case, the vanishing boundary conditions for σ^{ϵ} and κ^{ϵ} are no longer satisfied. We thus need some additional estimates to obtain the uniform L^2 bound.

We first estimate the spatial average $\overline{\sigma^{\epsilon}}(t) := \int_{\mathbb{S}^1} \sigma^{\epsilon}(s, t) \, ds$. Indeed, an integration by parts and Cauchy–Schwarz yield

$$\left|\overline{\sigma^{\epsilon}}(t)\right| = \left|\int_{\mathbb{S}^{1}} \kappa^{\epsilon} \cdot \partial_{s} \xi^{\epsilon} \, \mathrm{d}s\right| = \left|-\int_{\mathbb{S}^{1}} \partial_{s} \kappa^{\epsilon} \cdot \xi^{\epsilon} \, \mathrm{d}s\right| \le \|\partial_{s} \kappa^{\epsilon}(t, \cdot)\|_{L^{2}(\mathbb{S}^{1})} \|\xi^{\epsilon}\|_{L^{2}(\mathbb{S}^{1})}, \forall t \in [0, T].$$

Thus

$$\|\overline{\sigma^{\epsilon}}\|_{L^{2}([0,T])} \leq \sup_{t \in [0,T]} \|\xi^{\epsilon}(t,\cdot)\|_{L^{2}(\mathbb{S}^{1})} \|\partial_{s}\kappa^{\epsilon}\|_{L^{2}(\mathcal{Q}_{T})}$$

where the right-hand side is uniformly bounded (cf. [1, Proposition 3.1]). Thus from Poincaré inequality $\|\sigma^{\epsilon} - \overline{\sigma^{\epsilon}}\|_{L^{2}(Q_{T})} \leq C(T) \|\partial_{s}\sigma^{\epsilon}\|_{L^{2}(Q_{T})}$ we obtain $\|\sigma^{\epsilon}\|_{L^{2}(Q_{T})} \leq C$ with *C* independent of ϵ .

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It remains to show that κ^{ϵ} is uniformly bounded in $L^2(Q_T)$. To this end, we note that from the definition of κ^{ϵ} one has $\partial_s \xi^{\epsilon} = \epsilon \kappa^{\epsilon} + \frac{\kappa^{\epsilon}}{\sqrt{\epsilon + |\kappa^{\epsilon}|^2}}$, whence

$$\sigma^{\epsilon} = \kappa^{\epsilon} \cdot \partial_{s} \xi^{\epsilon} = |\kappa^{\epsilon}| \left(\epsilon |\kappa^{\epsilon}| + \frac{|\kappa^{\epsilon}|}{\sqrt{\epsilon + |\kappa^{\epsilon}|^{2}}} \right) \ge |\kappa^{\epsilon}| |\partial_{s} \xi^{\epsilon}|.$$

Observing that

$$|\partial_s \xi^{\epsilon}| = \epsilon |\kappa^{\epsilon}| + \frac{|\kappa^{\epsilon}|}{\sqrt{\epsilon + |\kappa^{\epsilon}|^2}} \ge \epsilon + \frac{1}{\sqrt{1 + \epsilon}} > 1$$

provided $|\kappa^{\epsilon}| \ge 1$, we infer that $|\kappa^{\epsilon}| \le \sigma^{\epsilon}$ when $|\kappa^{\epsilon}| \ge 1$. Thus,

$$\begin{split} \int_{Q_T} |\kappa^{\epsilon}|^2 \, \mathrm{d}s \, \mathrm{d}t &= \int_{\{(s,t) \in Q_T : |\kappa^{\epsilon}| < 1\}} |\kappa^{\epsilon}|^2 \, \mathrm{d}s \, \mathrm{d}t + \int_{\{(s,t) \in Q_T : |\kappa^{\epsilon}| \ge 1\}} |\kappa^{\epsilon}|^2 \, \mathrm{d}s \, \mathrm{d}t \\ &\leq |Q_T| + \int_{Q_T} |\sigma^{\epsilon}|^2 \, \mathrm{d}s \, \mathrm{d}t \le C, \end{split}$$

where C does not depend on ϵ .

Reference

 Shi, W., Vorotnikov, D.: The gradient flow of the potential energy on the space of arcs. Calc. Var. 58, 59 (2019)

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