



Mixed Martingale Hardy Spaces

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Abstract

In this paper, we consider the martingale Hardy spaces defined with the help of the mixed $L_{\vec{p}}$ -norm. Five mixed martingale Hardy spaces will be investigated: $H_{\vec{p}}^S$, $H_{\vec{p}}^{\mathcal{S}}$, $H_{\vec{p}}^M$, $\mathcal{P}_{\vec{p}}$, and $\mathcal{Q}_{\vec{p}}$. Several results are proved for these spaces, like atomic decompositions, Doob's inequality, boundedness, martingale inequalities, and the generalization of the well-known Burkholder–Davis–Gundy inequality.

Keywords Mixed Lebesgue spaces · Mixed martingale Hardy spaces · Atomic decompositions · Martingale inequalities · Doob's inequality · Weighted maximal inequality · Burkholder–Davis–Gundy inequality

Mathematics Subject Classification Primary 42B30 · Secondary 60G42 · 42B35 · 42B25

1 Introduction

The mixed Lebesgue spaces were introduced in 1961 by Benedek and Panzone [2]. They considered the Descartes product $(\Omega, \mathcal{F}, \mathbb{P})$ of the probability spaces $(\Omega_i, \mathcal{F}^i, \mathbb{P}_i)$, where $\Omega = \prod_{i=1}^d \Omega_i$, \mathcal{F} is generated by $\prod_{i=1}^d \mathcal{F}^i$ and \mathbb{P} is generated by $\prod_{i=1}^d \mathbb{P}_i$. The mixed $L_{\vec{p}}$ -norm of the measurable function f is defined as a number obtained after taking successively the L_{p_1} -norm of f in the variable x_1 , the L_{p_2} -norm of f in the variable x_2, \dots , the L_{p_d} -norm of f in the variable x_d . Some basic proper-

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ties of the spaces $L_{\vec{p}}$ were proved in [2], such as the well-known Hölder's inequality or the duality theorem for $L_{\vec{p}}$ -norm (see Lemma 2). The boundedness of operators on mixed-norm spaces has been studied for instance by Fernandez [29] and Stefanov and Torres [36]. Using the mixed Lebesgue spaces, Lizorkin [32] considered Fourier integrals and estimations for convolutions. Torres and Ward [38] gave the wavelet characterization of the space $L_{\vec{p}}(\mathbb{R}^n)$. For more about mixed-norm spaces, the papers [1, 8, 9, 18, 20, 21, 24, 25, 35, 36] are referred.

Since 1970, the theory of Hardy spaces has been developed very quickly (see, e.g., Fefferman and Stein [11], Stein [37], Grafakos [17]). Parallel, a similar theory was evolved for martingale Hardy spaces (see, e.g., Garsia [12], Long [33] and Weisz [40]). Recently several papers were published about the generalization of Hardy spaces. For example, (anisotropic) Hardy spaces with variable exponents were considered in Nakai and Sawano [34], Yan et al. [43], Jiao et al. [28], Liu et al. [30, 31]. Moreover, Musielak–Orlicz–Hardy spaces were studied in Yang et al. [44]. These results were also investigated for martingale Hardy spaces in Jiao et al. [26, 27] and Xie et al. [42]. The mixed-norm classical Hardy spaces were intensively studied by Huang et al. [22, 23] and Huang and Yang [24]. In this paper, we will develop a similar theory for mixed-norm martingale Hardy spaces.

The classical martingale Hardy spaces $(H_p^s, H_p^S, H_p^M, \mathcal{P}_p, \mathcal{Q}_p)$ have a long history and the results of this topic can be well applied in the Fourier analysis. In the celebrated work of Burkholder and Gundy [6], it was proved that the L_p norms of the maximal function and the quadratic variation, that is, the spaces H_p^M and H_p^S , are equivalent for $1 < p < \infty$. In the same year, Davis [10] extended this result for $p = 1$. In the classical case, Herz [19] and Weisz [40] gave one of the most powerful techniques in the theory of martingale Hardy spaces, the so-called atomic decomposition. Some boundedness results, duality theorems, martingale inequalities, and interpolation results can be proved with the help of atomic decomposition. Details for the martingale Hardy spaces can be found in Burkholder [4, 5], Burkholder and Gundy [6], Garsia [12], Long [33], or Weisz [39, 40]. For the application of martingale Hardy spaces in Fourier analysis, see Gát [13, 14], Goginava [15, 16] or Weisz [40, 41].

In this paper, we will introduce five mixed martingale Hardy spaces: $H_{\vec{p}}^s, H_{\vec{p}}^S, H_{\vec{p}}^M, \mathcal{P}_{\vec{p}}$ and $\mathcal{Q}_{\vec{p}}$. In Sect. 3, Doob's inequality will be proved, that is, we will show that

$$\left\| \sup_{n \in \mathbb{N}} |\mathbb{E}_n f| \right\|_{\vec{p}} \leq C \|f\|_{\vec{p}}$$

for all $f \in L_{\vec{p}}$, where $1 < \vec{p} < \infty$. In Sect. 4, we give the atomic decomposition for the five mixed martingale Hardy spaces. Using the atomic decompositions and Doob's inequality, the boundedness of general σ -subadditive operators from $H_{\vec{p}}^s$ to $L_{\vec{p}}$, from $\mathcal{P}_{\vec{p}}$ to $L_{\vec{p}}$ and from $\mathcal{Q}_{\vec{p}}$ to $L_{\vec{p}}$ can be proved (see Theorems 7 and 8). With the help of these general boundedness theorems, several martingale inequalities will be proved in Sect. 5 (see Corollary 21). We will show, that if the stochastic basis (\mathcal{F}_n) is regular, then the five martingale Hardy spaces are equivalent. As a consequence of Doob's inequality, the well-known Burkholder–Davis–Gundy inequality

can be shown. Moreover, if the stochastic basis is regular, then the so-called martingale transform is bounded on $L_{\vec{p}}$.

We denote by C a positive constant, which can vary from line to line, and denote by C_p a constant depending only on p . The symbol $A \sim B$ means that there exist constants $\alpha, \beta > 0$ such that $\alpha A \leq B \leq \beta A$.

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2 Backgrounds

2.1 Mixed Lebesgue Spaces

We will start with the definition of the mixed Lebesgue spaces. To this end, let $1 \leq d \in \mathbb{N}$ and $(\Omega_i, \mathcal{F}^i, \mathbb{P}_i)$ be probability spaces for $i = 1, \dots, d$, and $\vec{p} := (p_1, \dots, p_d)$ with $0 < p_i \leq \infty$. Consider the product space $(\Omega, \mathcal{F}, \mathbb{P})$, where $\Omega = \prod_{i=1}^d \Omega_i$, \mathcal{F} is generated by $\prod_{i=1}^d \mathcal{F}^i$ and \mathbb{P} is generated by $\prod_{i=1}^d \mathbb{P}_i$. A measurable function $f : \Omega \rightarrow \mathbb{R}$ belongs to the mixed $L_{\vec{p}}$ space if

$$\begin{aligned} \|f\|_{\vec{p}} &:= \|f\|_{(p_1, \dots, p_d)} := \left\| \dots \|f\|_{L_{p_1}(dx_1)} \dots \right\|_{L_{p_d}(dx_d)} \\ &= \left(\int_{\Omega_d} \dots \left(\int_{\Omega_1} |f(x_1, \dots, x_d)|^{p_1} dx_1 \right)^{p_2/p_1} \dots dx_d \right)^{1/p_d} < \infty \end{aligned}$$

with the usual modification if $p_j = \infty$ for some $j \in \{1, \dots, d\}$. If for some $0 < p \leq \infty$, $\vec{p} = (p, \dots, p)$, then we get back the classical Lebesgue space, that is, in this case, $L_{\vec{p}} = L_p$. Throughout the paper, $0 < \vec{p} \leq \infty$ will mean that the coordinates of \vec{p} satisfy the previous condition, e.g., for all $i = 1, \dots, d$, $0 < p_i \leq \infty$. The conjugate exponent vector of \vec{p} will be denoted by $(\vec{p})'$, that is, if $(\vec{p})' = (p'_1, \dots, p'_d)$, then $1/p_i + 1/p'_i = 1$ ($i = 1, \dots, d$). For $\alpha > 0$, $\vec{p}/\alpha := (p_1/\alpha, \dots, p_d/\alpha)$. Benedek and Panzone [2] proved some basic properties for the mixed Lebesgue space.

Lemma 1 *If $1 \leq \vec{p} \leq \infty$, then for all $f \in L_{\vec{p}}$ and $g \in L_{(\vec{p})'}$, $fg \in L_1$ and*

$$\|fg\|_1 = \int_{\Omega_d} \dots \int_{\Omega_1} |f(x)g(x)| d\mathbb{P}(x) \leq \|f\|_{\vec{p}} \|g\|_{(\vec{p})'}$$

Lemma 2 *If $1 \leq \vec{p} \leq \infty$ and $f \in L_{\vec{p}}$, then*

$$\|f\|_{\vec{p}} = \sup_{\|g\|_{(\vec{p})'} \leq 1} \left| \int_{\Omega} fg d\mathbb{P} \right|$$

2.2 Martingale Hardy Spaces

Suppose that the σ -algebra $\mathcal{F}_n^i \subset \mathcal{F}^i$ ($n \in \mathbb{N}$, $i = 1, \dots, d$), $(\mathcal{F}_n^i)_{n \in \mathbb{N}}$ is increasing and $\mathcal{F}^i = \sigma(\cup_{n \in \mathbb{N}} \mathcal{F}_n^i)$. Let $\mathcal{F}_n = \sigma(\prod_{i=1}^d \mathcal{F}_n^i)$. The expectation and conditional expectation operators relative to \mathcal{F}_n are denoted by \mathbb{E} and \mathbb{E}_n , respectively. An integrable sequence $f = (f_n)_{n \in \mathbb{N}}$ is said to be a *martingale* if

- (i) $(f_n)_{n \in \mathbb{N}}$ is *adapted*, that is for all $n \in \mathbb{N}$, f_n is \mathcal{F}_n -measurable;
- (ii) $\mathbb{E}_n f_m = f_n$ in case $n \leq m$.

For $n \in \mathbb{N}$, the *martingale difference* is defined by $d_n f := f_n - f_{n-1}$, where $f = (f_n)_{n \in \mathbb{N}}$ is a martingale and $f_0 := f_{-1} := 0$. If for all $n \in \mathbb{N}$, $f_n \in L_{\vec{p}}$, then f is called an $L_{\vec{p}}$ -martingale. Moreover, if

$$\|f\|_{\vec{p}} := \sup_{n \in \mathbb{N}} \|f_n\|_{\vec{p}} < \infty,$$

then f is called an $L_{\vec{p}}$ -bounded martingale and it will be denoted by $f \in L_{\vec{p}}$. The map $\nu : \Omega \rightarrow \mathbb{N} \cup \{\infty\}$ is called a *stopping time relative to* (\mathcal{F}_n) if for all $n \in \mathbb{N}$, $\{\nu = n\} \in \mathcal{F}_n$.

For a martingale $f = (f_n)$ and a stopping time ν , the *stopped martingale* is defined by

$$f_n^\nu = \sum_{m=0}^n d_m f \chi_{\{\nu \geq m\}}.$$

Let us define the *maximal function*, the *quadratic variation* and the *conditional quadratic variation* of the martingale f relative to $(\Omega, \mathcal{F}, \mathbb{P}, (\mathcal{F}_n)_{n \in \mathbb{N}})$ by

$$\begin{aligned} M_m(f) &:= \sup_{n \leq m} |f_n|, & M(f) &:= \sup_{n \in \mathbb{N}} |f_n| \\ S_m(f) &:= \left(\sum_{n=0}^m |d_n f|^2 \right)^{1/2}, & S(f) &:= \left(\sum_{n=0}^{\infty} |d_n f|^2 \right)^{1/2} \\ s_m(f) &:= \left(\sum_{n=0}^m \mathbb{E}_{n-1} |d_n f|^2 \right)^{1/2}, & s(f) &:= \left(\sum_{n=0}^{\infty} \mathbb{E}_{n-1} |d_n f|^2 \right)^{1/2}. \end{aligned}$$

The set of the sequences $(\lambda_n)_{n \in \mathbb{N}}$ of non-decreasing, non-negative, and adapted functions with $\lambda_\infty := \lim_{n \rightarrow \infty} \lambda_n$ is denoted by Λ . With the help of the previous operators, the mixed martingale Hardy spaces can be defined as follows:

$$\begin{aligned} H_{\vec{p}}^M &:= \left\{ f = (f_n)_{n \in \mathbb{N}} : \|f\|_{H_{\vec{p}}^M} := \|M(f)\|_{\vec{p}} < \infty \right\}; \\ H_{\vec{p}}^S &:= \left\{ f = (f_n)_{n \in \mathbb{N}} : \|f\|_{H_{\vec{p}}^S} := \|S(f)\|_{\vec{p}} < \infty \right\}; \end{aligned}$$

$$\begin{aligned}
 H_{\vec{p}}^S &:= \left\{ f = (f_n)_{n \in \mathbb{N}} : \|f\|_{H_{\vec{p}}^S} := \|s(f)\|_{\vec{p}} < \infty \right\}; \\
 \mathcal{Q}_{\vec{p}} &:= \left\{ f = (f_n)_{n \in \mathbb{N}} : \exists (\lambda_n)_{n \in \mathbb{N}} \in \Lambda, \text{ such that } S_n(f) \leq \lambda_{n-1}, \lambda_\infty \in L_{\vec{p}} \right\}, \\
 \|f\|_{\mathcal{Q}_{\vec{p}}} &:= \inf_{(\lambda_n) \in \Lambda} \|\lambda_\infty\|_{\vec{p}}; \\
 \mathcal{P}_{\vec{p}} &:= \left\{ f = (f_n)_{n \in \mathbb{N}} : \exists (\lambda_n)_{n \in \mathbb{N}} \in \Lambda, \text{ such that } |f_n| \leq \lambda_{n-1}, \lambda_\infty \in L_{\vec{p}} \right\}, \\
 \|f\|_{\mathcal{P}_{\vec{p}}} &:= \inf_{(\lambda_n) \in \Lambda} \|\lambda_\infty\|_{\vec{p}}.
 \end{aligned}$$

3 Doob's Inequality

In this section, we will prove that the maximal operator M is bounded on the space $L_{\vec{p}}$ for $1 < \vec{p} < \infty$. If $1 < \vec{p} < \infty$ and the martingale $(f_n)_{n \in \mathbb{N}} \in L_{\vec{p}}$, then there is a function $g \in L_{\vec{p}}$ such that for all $n \in \mathbb{N}$, $f_n = \mathbb{E}_n g$. For $k \in \{1, \dots, d\}$ and $m \in \mathbb{N}$, let us denote

$$\mathcal{F}_{\infty, \dots, \infty, m, \infty, \dots, \infty} := \sigma \left(\mathcal{F}^1 \times \dots \times \mathcal{F}^{k-1} \times \mathcal{F}_m^k \times \mathcal{F}^{k+1} \times \dots \times \mathcal{F}^d \right),$$

where m stands in the k th position. The conditional expectation operator relative to $\mathcal{F}_{\infty, \dots, \infty, m, \infty, \dots, \infty}$ is denoted by $\mathbb{E}_{\infty, \dots, \infty, m, \infty, \dots, \infty}$. We need the following maximal operators: for an integrable function f , let

$$M_k f := \sup_{m \in \mathbb{N}} \left| \mathbb{E}_{\infty, \dots, \infty, m, \infty, \dots, \infty} f \right|, \quad \tilde{M} f := M_d \circ M_{d-1} \circ \dots \circ M_1 f.$$

It is clear that

$$M f \leq \tilde{M} f.$$

For $0 < p < \infty$ and $w > 0$, the weighted space $L_p(w)$ consists of all functions f , for which

$$\|f\|_{L_p(w)} := \left(\int_{\Omega} |f|^p w d\mathbb{P} \right)^{1/p} < \infty.$$

We need the following lemma.

Lemma 3 *Let φ be a positive function. Then for all $r > 1$, M is bounded from $L_r(M\varphi)$ to $L_r(\varphi)$, that is for all $f \in L_r(M\varphi)$, we have*

$$\int_{\Omega} |Mf|^r \varphi d\mathbb{P} \leq C_r \int_{\Omega} |f|^r M\varphi d\mathbb{P}. \tag{1}$$

Proof It is easy to see that the operator M is bounded from $L_\infty(M\varphi)$ to $L_\infty(\varphi)$. We will prove that M is bounded from $L_1(M\varphi)$ to $L_{1,\infty}(\varphi)$ as well, where $L_{1,\infty}(\varphi)$

denotes the weak- $L_1(\varphi)$ space. From this, it follows by interpolation (see, e.g., [3]) that for all $r > 1$, the operator M is bounded from $L_r(M\varphi)$ to $L_r(\varphi)$, in other words, (1) holds. Let $\varrho > 0$ arbitrary and let

$$v_\varrho := \inf \{n \in \mathbb{N} : |f_n| > \varrho\}.$$

Since $\{v_\varrho < \infty\} = \{Mf > \varrho\}$, we get that

$$\begin{aligned} \|Mf\|_{L_{1,\infty}(\varphi)} &:= \varrho \int_{\{Mf > \varrho\}} \varphi \, d\mathbb{P} = \varrho \sum_{k \in \mathbb{N}} \int_{\{v_\varrho = k\}} \varphi \, d\mathbb{P} \leq \sum_{k \in \mathbb{N}} \int_{\{v_\varrho = k\}} |fk| \varphi \, d\mathbb{P} \\ &= \sum_{k \in \mathbb{N}} \int_{\{v_\varrho = k\}} \mathbb{E}_k |f| \varphi \, d\mathbb{P} = \sum_{k \in \mathbb{N}} \int_{\{v_\varrho = k\}} |f| \mathbb{E}_k \varphi \, d\mathbb{P} \\ &\leq \sum_{k \in \mathbb{N}} \int_{\{v_\varrho = k\}} |f| M\varphi \, d\mathbb{P} \leq \int_\Omega |f| M\varphi \, d\mathbb{P} = \|f\|_{L_1(M\varphi)}, \end{aligned}$$

which finishes the proof. □

Now we prove that M_d is bounded on $L_{\vec{p}}$.

Theorem 1 *For all $1 < \vec{p} < \infty$, M_d is bounded on $L_{\vec{p}}$, that is for all $f \in L_{\vec{p}}$,*

$$\|M_d f\|_{\vec{p}} \leq C \|f\|_{\vec{p}}.$$

Proof We will prove the theorem by induction in d . If $d = 1$, then the function f has only 1 variable and the theorem holds (see, e.g., Weisz [40]). Suppose that the theorem is true for some fixed $d \in \mathbb{N}$ and for all $1 < \vec{p} = (p_1, \dots, p_d) < \infty$ and $f \in L_{\vec{p}}$. For a function f with d variables and for the vector (p_1, \dots, p_k) , let us denote

$$\begin{aligned} T_{(p_1, \dots, p_k)} f(x_{k+1}, \dots, x_d) &:= \|f(\cdot, \dots, \cdot, x_{k+1}, \dots, x_d)\|_{(p_1, \dots, p_k)} \\ &= \left(\int_{\Omega_k} \dots \left(\int_{\Omega_1} |f(x_1, \dots, x_d)|^{p_1} dx_1 \right)^{p_2/p_1} \dots dx_k \right)^{1/p_k}. \end{aligned}$$

Using this notation, the condition of the induction can be written in the form

$$\int_{\Omega_d} [T_{(p_1, \dots, p_{d-1})}(M_d f)]^{p_d}(x_d) dx_d \leq C \int_{\Omega_d} [T_{(p_1, \dots, p_{d-1})}(f)]^{p_d}(x_d) dx_d.$$

We will show that if $1 < p_{d+1} < \infty$, then for all $f \in L_{(p_1, \dots, p_{d+1})}$,

$$\begin{aligned} &\int_{\Omega_{d+1}} [T_{(p_1, \dots, p_d)}(M_{d+1} f)]^{p_{d+1}}(x_{d+1}) dx_{d+1} \\ &\leq C \int_{\Omega_{d+1}} [T_{(p_1, \dots, p_d)}(f)]^{p_{d+1}}(x_{d+1}) dx_{d+1}, \end{aligned}$$

where f has $d + 1$ variable and the maximal operator M_{d+1} is taken in the variable x_{d+1} . If $p_1 = \infty$, then

$$|f(x_1, \dots, x_d, x_{d+1})| \leq \sup_{x_1 \in \Omega_1} |f(x_1, \dots, x_d, x_{d+1})| = T_\infty f(x_2, \dots, x_d, x_{d+1}).$$

Hence $M_{d+1}f(x_1, \dots, x_d, x_{d+1}) \leq M_{d+1}(T_\infty f)(x_2, \dots, x_d, x_{d+1})$, and therefore

$$\begin{aligned} T_\infty(M_{d+1}f)(x_2, \dots, x_d, x_{d+1}) &= \sup_{x_1 \in \Omega_1} M_{d+1}f(x_1, \dots, x_d, x_{d+1}) \\ &\leq M_{d+1}(T_\infty f)(x_2, \dots, x_d, x_{d+1}). \end{aligned}$$

So we get that

$$\begin{aligned} &\int_{\Omega_{d+1}} [T_{(p_2, \dots, p_d)}(T_\infty(M_{d+1}f))]^{p_{d+1}}(x_{d+1}) dx_{d+1} \\ &\leq \int_{\Omega_{d+1}} [T_{(p_2, \dots, p_d)}(M_{d+1}(T_\infty f))]^{p_{d+1}}(x_{d+1}) dx_{d+1}. \end{aligned} \tag{2}$$

Here the function $M_{d+1}(T_\infty f)$ has d variables: x_2, \dots, x_{d+1} and the maximal operator is taken over the d th variable, that is over x_{d+1} . Therefore by induction, we have that (2) can be estimated by

$$C \int_{\Omega_{d+1}} [T_{(p_2, \dots, p_d)}(T_\infty f)]^{p_{d+1}}(x_{d+1}) dx_{d+1},$$

which means that

$$\|M_{d+1}f\|_{(\infty, p_2, \dots, p_d, p_{d+1})} \leq C \|f\|_{(\infty, p_2, \dots, p_d, p_{d+1})}, \tag{3}$$

so the theorem holds for $p_1 = \infty$. Now let choose a number r for which $1 < r < \min\{p_2, \dots, p_d, p_{d+1}\}$. It will be shown that

$$\|M_{d+1}f\|_{(r, p_2, \dots, p_d, p_{d+1})} \leq C \|f\|_{(r, p_2, \dots, p_d, p_{d+1})}.$$

It is easy to see that

$$\|M_{d+1}f\|_{(r, p_2, \dots, p_{d+1})} = \|[T_r(M_{d+1}f)]^r\|_{\left(\frac{p_2}{r}, \dots, \frac{p_{d+1}}{r}\right)}^{1/r}.$$

Let $\vec{q} := (\frac{p_2}{r}, \dots, \frac{p_{d+1}}{r})$. Then, the vector \vec{q} has d coordinates and $1 < \vec{q} < \infty$. Using Lemma 2,

$$\|[T_r(M_{d+1}f)]^r\|_{\vec{q}}$$

$$= \sup_{\substack{\varphi \in L_{(\vec{q})'} \\ \|\varphi\|_{(\vec{q})'} \leq 1}} \left| \int_{\Omega_{d+1}} \cdots \int_{\Omega_2} [T_r(M_{d+1}f)]^r(x_2, \dots, x_{d+1}) \varphi(x_2, \dots, x_{d+1}) dx_2 \cdots dx_{d+1} \right|.$$

We can suppose that $\varphi > 0$. Then

$$\begin{aligned} & \int_{\Omega_{d+1}} \cdots \int_{\Omega_2} [T_r(M_{d+1}f)]^r(x_2, \dots, x_{d+1}) \varphi(x_2, \dots, x_{d+1}) dx_2 \cdots dx_{d+1} \\ &= \int_{\Omega_1} \left(\int_{\Omega_d} \cdots \int_{\Omega_2} \left(\int_{\Omega_{d+1}} |M_{d+1}f|^r(x_1, \dots, x_{d+1}) \varphi(x_2, \dots, x_{d+1}) dx_{d+1} \right) dx_2 \cdots dx_d \right) dx_1. \end{aligned} \tag{4}$$

Since $1 < r < \infty$, applying Lemma 3 for the variable x_{d+1} , we have that for all fixed x_1, \dots, x_d ,

$$\begin{aligned} & \int_{\Omega_{d+1}} |M_{d+1}f|^r(x_1, \dots, x_{d+1}) \varphi(x_2, \dots, x_{d+1}) dx_{d+1} \\ & \leq C \int_{\Omega_{d+1}} |f(x_1, \dots, x_{d+1})|^r M_{d+1}\varphi(x_2, \dots, x_{d+1}) dx_{d+1} \end{aligned}$$

Hence (4) can be estimated by

$$\begin{aligned} & C \int_{\Omega_{d+1}} \cdots \int_{\Omega_2} \left(\int_{\Omega_1} |f(x_1, \dots, x_{d+1})|^r dx_1 \right) M_{d+1}\varphi(x_2, \dots, x_{d+1}) dx_2 \cdots dx_{d+1} \\ &= C \int_{\Omega_{d+1}} \cdots \int_{\Omega_2} [T_r(f)]^r(x_2, \dots, x_{d+1}) M_{d+1}\varphi(x_2, \dots, x_{d+1}) dx_2 \cdots dx_{d+1} \\ & \leq C \| [T_r(f)]^r \|_{\vec{q}} \| M_{d+1}\varphi \|_{(\vec{q})'}. \end{aligned} \tag{5}$$

Here $M_{d+1}\varphi$ is a function with d variables, the vector $(\vec{q})'$ has d coordinates such that $1 < (\vec{q})' < \infty$ and the maximal operator is taken in the d th coordinate, that is over the variable x_{d+1} . So, by induction, we get that

$$\| M_{d+1}\varphi \|_{(\vec{q})'} \leq C \| \varphi \|_{(\vec{q})'} \leq C.$$

Therefore, we can estimate (5) by

$$\begin{aligned} & C \| [T_r(f)]^r \|_{\vec{q}} \\ &= C \left(\int_{\Omega_{d+1}} \cdots \left(\int_{\Omega_2} \left(\int_{\Omega_1} |f(x_1, x_2, \dots, x_{d+1})|^r dx_1 \right)^{p_2/r} dx_2 \right)^{p_3/r-r/p_2} \cdots dx_{d+1} \right)^{r/p_{d+1}} \\ &= \| f \|_{(r, p_2, \dots, p_{d+1})}^r. \end{aligned}$$

Consequently,

$$\|M_{d+1} f\|_{(r, p_2, \dots, p_{d+1})} \leq C \|f\|_{(r, p_1, \dots, p_{d+1})}. \tag{6}$$

Combining the results (3) and (6), we get by interpolation that for all $1 < p_1 < \infty$

$$\|M_{d+1} f\|_{(p_1, p_2, \dots, p_{d+1})} \leq C \|f\|_{(p_1, p_1, \dots, p_{d+1})}.$$

Using induction, the proof is complete. □

Remark 1 Using the proof of the previous theorem, Theorem 1 can be generalized for \vec{p} -s, for which its first k coordinates are ∞ , but the others are strongly between 1 and ∞ , that is, for \vec{p} -s, with

$$\vec{p} = (\infty, \infty, \dots, \infty, p_{k+1}, \dots, p_d), \quad 1 < p_{k+1}, \dots, p_d < \infty \tag{7}$$

for some $k \in \{1, \dots, d\}$.

Now, we can generalize the well-known Doob’s inequality. Using the previous theorem, we get that the maximal operator M is bounded on $L_{\vec{p}}$ in case $1 < \vec{p} < \infty$.

Theorem 2 *If $1 < \vec{p} < \infty$ or \vec{p} satisfies (7), then the maximal operator M is bounded on $L_{\vec{p}}$, that is, for all $f \in L_{\vec{p}}$,*

$$\|Mf\|_{\vec{p}} \leq C \|f\|_{\vec{p}}.$$

Proof It is clear that $Mf \leq \tilde{M}f = M_d \circ \dots \circ M_1 f$, therefore by Theorem 1 and Remark 1,

$$\begin{aligned} \|Mf\|_{\vec{p}} &\leq \|M_d \circ M_{d-1} \circ \dots \circ M_1 f\|_{\vec{p}} \leq C \|M_{d-1} \circ \dots \circ M_1 f\|_{\vec{p}} \\ &\leq C \|M_{d-2} \circ \dots \circ M_1 f\|_{\vec{p}} \leq \dots \leq C \|f\|_{\vec{p}} \end{aligned}$$

and the proof is complete. □

A weighted version of Doob’s inequality can be found in Chen et al. [7]. The following corollary is well known for classical Hardy spaces with $\vec{p} = (p, \dots, p)$.

Corollary 1 *If $1 < \vec{p} < \infty$, or \vec{p} satisfies (7), then $H_{\vec{p}}^M$ is equivalent to $L_{\vec{p}}$.*

Theorem 3 *Theorem 2 is not true for all $1 < \vec{p} \leq \infty$.*

Proof We prove the theorem for two dimensions and for the exponent $\vec{p} := (p, \infty)$, where $1 < p < \infty$. The proof is similar for higher dimensions. Let us define the following sequence of functions

$$f_n(x, y) := \sum_{k=1}^n 2^{k/p} \chi_{[2^{-k}, 2^{-k+1})^2}(x, y) \quad ((x, y) \in [0, 1) \times [0, 1)).$$

Then for an arbitrary fixed $y \in [2^{-k}, 2^{-k+1})$ ($k = 1, \dots, n$),

$$\int_{[0,1)} |f_n(x, y)|^p dx = 2^k \frac{1}{2^k} = 1$$

and for all fixed $y \notin [2^{-n}, 1)$, the previous integral is 0. From this follows that for all $n \in \mathbb{N}$,

$$\|f_n\|_{(p,\infty)} = \sup_{y \in [0,1)} \left(\int_{[0,1)} |f(x, y)|^p dx \right)^{1/p} = 1.$$

At the same time, for $x \in [2^{-k}, 2^{-k+1})$ ($k = 1, \dots, n$) and $y \in [0, 2^{-n})$,

$$Mf_n(x, y) \geq \frac{1}{|[0, 2^{-k+1})^2|} \int_{[0, 2^{-k+1})^2} |f_n(u, v)| dudv \geq \frac{1}{2^{-2k+2}} 2^{k/p} 2^{-2k} = \frac{2^{k/p}}{4}.$$

Hence we get that for all $y \in [0, 2^{-n})$,

$$\begin{aligned} \int_{[0,1)} |Mf_n(x, y)|^p dx &\geq \sum_{k=1}^n \int_{[2^{-k}, 2^{-k+1})} |Mf_n(x, y)|^p dx \\ &\geq \sum_{k=1}^n \frac{2^k}{4} 2^{-k} = \frac{n}{4} \rightarrow \infty \quad (n \rightarrow \infty) \end{aligned}$$

and therefore

$$\|Mf_n\|_{(p,\infty)} \rightarrow \infty \quad (n \rightarrow \infty),$$

which means, that M is not bounded on $L_{(p,\infty)}$. □

Remark 2 This counterexample proves also that M_2 is not bounded on $L_{(p,\infty)}$. Moreover, the counterexample shows also that the classical Hardy–Littlewood maximal operator considered in Huang et al. [22] is not bounded on $L_{(p,\infty)}$ (cf. Lemma 3.5 in [22] and Lemma 4.8 in [35]).

4 Atomic Decomposition

First of all, we need the definition of the *atoms*. For \vec{p} , a measurable function a is called an (s, \vec{p}) -atom (or (S, \vec{p}) -atom or (M, \vec{p}) -atom) if there exists a stopping time τ such that

- (i) $\mathbb{E}_n a = 0$ for all $n \leq \tau$,
- (ii) $\|s(a)\chi_{\{\tau < \infty\}}\|_\infty \leq \|\chi_{\{\tau < \infty\}}\|_{\vec{p}}^{-1}$
 (or $\|S(a)\chi_{\{\tau < \infty\}}\|_\infty \leq \|\chi_{\{\tau < \infty\}}\|_{\vec{p}}^{-1}$, or $\|M(a)\chi_{\{\tau < \infty\}}\|_\infty \leq \|\chi_{\{\tau < \infty\}}\|_{\vec{p}}^{-1}$, respectively).

Now we can give the atomic decomposition of the space $H_{\vec{p}}^s$.

Theorem 4 *Let $0 < \vec{p} < \infty$. A martingale $f = (f_n)_{n \in \mathbb{N}} \in H_{\vec{p}}^s$ if and only if there exists a sequence $(a^k)_{k \in \mathbb{Z}}$ of (s, \vec{p}) -atoms and a sequence $(\mu_k)_{k \in \mathbb{Z}}$ of real numbers such that*

$$f_n = \sum_{k \in \mathbb{Z}} \mu_k \mathbb{E}_n a^k \quad \text{a. e. } (n \in \mathbb{N}) \tag{8}$$

and

$$\|f\|_{H_{\vec{p}}^s} \sim \inf \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}}, \tag{9}$$

where $0 < t \leq 1$ and the infimum is taken over all decompositions of the form (8).

Proof Let $f \in H_{\vec{p}}^s$ and let us define the following stopping times:

$$\tau_k := \inf \left\{ n \in \mathbb{N} : s_{n+1}(f) > 2^k \right\}.$$

Obviously f_n can be written in the form

$$f_n = \sum_{k \in \mathbb{Z}} (f_n^{\tau_{k+1}} - f_n^{\tau_k}).$$

Let

$$\mu_k := 3 \cdot 2^k \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}, \quad \text{and} \quad a_n^k := \frac{f_n^{\tau_{k+1}} - f_n^{\tau_k}}{\mu_k}. \tag{10}$$

If $\mu_k = 0$, then let $a_n^k = 0$. If $n \leq \tau_k$, then $a_n^k = 0$ and naturally

$$f_n = \sum_{k \in \mathbb{Z}} \mu_k a_n^k.$$

Moreover, (a_n^k) is L_2 -bounded (see [40]), therefore there exists $a^k \in L_2$ such that $\mathbb{E}_n a^k = a_n^k$. Because of $s(f^{\tau_k}) = s_{\tau_k}(f) \leq 2^k$, we have that

$$s(a^k) \leq \frac{s(f^{\tau_{k+1}}) + s(f^{\tau_k})}{\mu_k} \leq \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}^{-1},$$

thus a^k is an (s, \vec{p}) -atom.

Since $s^2(f - f^{\tau_k}) = s^2(f) - s^2(f^{\tau_k})$, thus $s(f - f^{\tau_k}) \leq s(f)$ and $s(f^{\tau_k}) \leq s(f)$. Using that $\lim_{k \rightarrow \infty} s(f - f^{\tau_k}) = \lim_{k \rightarrow \infty} s(f^{\tau_k}) = 0$ almost everywhere, by the dominated convergence theorem (see, e.g., [2]) we get that

$$\left\| f - \sum_{k=-l}^m \mu_k a^k \right\|_{H_{\vec{p}}^s} \leq \|f - f^{\tau_{m+1}}\|_{H_{\vec{p}}^s} + \|f^{\tau_{-l}}\|_{H_{\vec{p}}^s} \rightarrow 0 \quad (l, m \rightarrow \infty),$$

this means that $f = \sum_{k \in \mathbb{Z}} \mu_k a^k$ in the $H^s_{\vec{p}}$ -norm.

Denote $\mathcal{O}_k := \{\tau_k < \infty\} = \{s(f) > 2^k\}$. Then for all $k \in \mathbb{Z}$, $\mathcal{O}_{k+1} \subset \mathcal{O}_k$. Moreover, for all $x \in \Omega$ and for all $0 < t \leq 1$,

$$\sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \chi_{\mathcal{O}_k}(x)\right)^t \leq C \left(\sum_{k \in \mathbb{Z}} 3 \cdot 2^k \chi_{\mathcal{O}_k \setminus \mathcal{O}_{k+1}}(x)\right)^t.$$

Since the sets $\mathcal{O}_k \setminus \mathcal{O}_{k+1}$ are disjoint, we have

$$\begin{aligned} \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}} &= \left\| \left(\sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \chi_{\{\tau_k < \infty\}} \right)^t \right)^{1/t} \right\|_{\vec{p}} \\ &\leq C \left\| \sum_{k \in \mathbb{Z}} 3 \cdot 2^k \chi_{\mathcal{O}_k \setminus \mathcal{O}_{k+1}} \right\|_{\vec{p}} \\ &\leq C \left\| \sum_{k \in \mathbb{Z}} s(f) \chi_{\mathcal{O}_k \setminus \mathcal{O}_{k+1}} \right\|_{\vec{p}} \\ &= C \|s(f)\|_{\vec{p}}. \end{aligned}$$

Conversely, if f has a decomposition of the form (8), then

$$s(f) \leq \sum_{k \in \mathbb{Z}} \mu_k s(a^k) \leq \sum_{k \in \mathbb{Z}} \mu_k \frac{\chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}},$$

and so for all $0 < t \leq 1$,

$$\|f\|_{H^s_{\vec{p}}} \leq \left\| \sum_{k \in \mathbb{Z}} \mu_k \frac{\chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right\|_{\vec{p}} \leq \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}},$$

which proves the theorem. □

For the classical martingale Hardy space H^s_p , this result is due to the second author (see [40]). For the spaces $\mathcal{Q}_{\vec{p}}$ and $\mathcal{P}_{\vec{p}}$, we can give similar decompositions.

Theorem 5 *Let $0 < \vec{p} < \infty$. A martingale $f = (f_n)_{n \in \mathbb{N}} \in \mathcal{P}_{\vec{p}}$ (or $\in \mathcal{Q}_{\vec{p}}$) if and only if there exists a sequence $(a^k)_{k \in \mathbb{Z}}$ of (M, \vec{p}) -atoms (or (S, \vec{p}) -atoms) and a sequence $(\mu_k)_{k \in \mathbb{Z}}$ of real numbers such that (8) holds and*

$$\|f\|_{\mathcal{P}_{\vec{p}}} \text{ (or } \|f\|_{\mathcal{Q}_{\vec{p}}}) \sim \inf \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}},$$

where $0 < t \leq 1$ and the infimum is taken over all decompositions of the form (8).

Proof Let $f \in \mathcal{P}_{\vec{p}}$ and $(\lambda_n)_{n \in \mathbb{N}}$ be a sequence such that $|f_n| \leq \lambda_{n-1}$ and $\lambda_\infty = \sup_n \lambda_n \in L_{\vec{p}}$. Let the stopping time τ_k be defined by

$$\tau_k := \inf \left\{ n \in \mathbb{N} : \lambda_n > 2^k \right\}$$

and μ_k and a_n^k be given by (10). Then again $f_n = \sum_{k \in \mathbb{Z}} \mu_k a_n^k$ and we can prove as before that

$$\left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}} \leq C \|f\|_{\mathcal{P}_{\vec{p}}}.$$

Conversely, assume that for some μ_k and a_n^k , the martingale (f_n) can be written in the form (8). For $n \in \mathbb{N}$, let us define

$$\lambda_n := \sum_{k \in \mathbb{Z}} \mu_k \chi_{\{\tau_k \leq n\}} \|M(a^k)\|_\infty.$$

It is clear, that (λ_n) is a non-negative adapted sequence and for all $n \in \mathbb{N}$, $|f_n| \leq \lambda_{n-1}$. Therefore, for all $0 < t \leq 1$,

$$\|f\|_{\mathcal{P}_{\vec{p}}} \leq \|\lambda_\infty\|_{\vec{p}} \leq \left\| \sum_{k \in \mathbb{Z}} \frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right\|_{\vec{p}} \leq \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}}.$$

The case of $\mathcal{Q}_{\vec{p}}$ is similar. □

The stochastic basis (\mathcal{F}_n) is said to be *regular*, if there exists $R > 0$ such that for all non-negative martingales (f_n) ,

$$f_n \leq R f_{n-1}. \tag{11}$$

If the stochastic basis is regular, then atomic decomposition can also be proved for the remainder two martingale Hardy spaces, $H_{\vec{p}}^M$ and $H_{\vec{p}}^S$.

Theorem 6 *Let $0 < \vec{p} < \infty$ and the stochastic basis (\mathcal{F}_n) be regular. A martingale $f = (f_n)_{n \in \mathbb{N}} \in H_{\vec{p}}^M$ (or $\in H_{\vec{p}}^S$) if and only if there exists a sequence $(a^k)_{k \in \mathbb{Z}}$ of (M, \vec{p}) -atoms (or (S, \vec{p}) -atoms) and a sequence $(\mu_k)_{k \in \mathbb{Z}}$ of real numbers such that (8) holds and*

$$\|f\|_{H_{\vec{p}}^M} \left(\text{or } \|f\|_{H_{\vec{p}}^S} \right) \sim \inf \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}},$$

where $0 < t < \min\{p_1, \dots, p_d, 1\}$ and the infimum is taken over all decompositions of the form (8).

Proof We will prove the theorem only for $H^M_{\vec{p}}$. The case of $H^S_{\vec{p}}$ is similar. Suppose that $f \in H^M_{\vec{p}}$ and for $k \in \mathbb{Z}$, let us define the stopping times

$$\varrho_k := \inf\{n \in \mathbb{N} : |f_n| > 2^k\}.$$

Moreover, let

$$I_{k,j} := \{\varrho_k = j\} \in \mathcal{F}_j \quad (j \in \mathbb{N})$$

and

$$\bar{I}_{k,j} := \left\{ \mathbb{E}_{j-1} \chi_{I_{k,j}} \geq \frac{1}{R} \right\} \in \mathcal{F}_{j-1} \quad (j \in \mathbb{N}).$$

Then (11) implies that $I_{k,j} \subset \bar{I}_{k,j}$. Moreover, for $I \in \mathcal{F}_{j-1}$, we get that

$$\mathbb{P}(I \cap \bar{I}_{k,j}) = \int_{I \cap \bar{I}_{k,j}} 1 \, d\mathbb{P} \leq R \int_I \mathbb{E}_{j-1} \chi_{I_{k,j}} \, d\mathbb{P} = R \mathbb{P}(I \cap I_{k,j}).$$

In other words

$$\int_{\bar{I}_{k,j}} \chi_I \, d\mathbb{P} \leq R \int_{I_{k,j}} \chi_I \, d\mathbb{P}$$

for all $I \in \mathcal{F}_{j-1}$. By a usual density argument, we obtain

$$\int_{\bar{I}_{k,j}} h \, d\mathbb{P} \leq R \int_{I_{k,j}} h \, d\mathbb{P} \tag{12}$$

for all non-negative \mathcal{F}_{j-1} -measurable function h .

Let us define a new family of stopping times by

$$\tau_k(x) := \inf\{n \in \mathbb{N} : x \in \bar{I}_{k,n+1}\} \quad (x \in \Omega, k \in \mathbb{Z}).$$

Then τ_k is non-decreasing and using Lemma 4,

$$\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}} \leq C \|\chi_{\{\varrho_k < \infty\}}\|_{\vec{p}} = \|\chi_{Mf > 2^k}\|_{\vec{p}} \leq 2^{-k} \|Mf\|_{\vec{p}} \rightarrow 0 \quad (k \rightarrow \infty),$$

which means that $\lim_{k \rightarrow \infty} \mathbb{P}(\{\tau_k < \infty\}) = 0$. Hence $\lim_{k \rightarrow \infty} \tau_k = \infty$ almost everywhere and therefore

$$\lim_{k \rightarrow \infty} f_n^{\tau_k} = f_n \quad a.e. \quad (n \in \mathbb{N}).$$

Let μ_k and a_n^k be defined again as in (10). Then $a^k = (a_n^k)$ is an (M, \vec{p}) -atom. Since $\{\tau_k < \infty\} = \cup_{j \in \mathbb{N}} \bar{I}_{k,j}$, we have

$$\begin{aligned} \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}} &\leq \left\| \left[\sum_{k \in \mathbb{Z}} \left(\sum_{j \in \mathbb{N}} \frac{\mu_k \chi_{\bar{I}_{k,j}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right]^{1/t} \right\|_{\vec{p}} \\ &\leq \left\| \sum_{k \in \mathbb{Z}} \sum_{j \in \mathbb{N}} (3 \cdot 2^k)^t \chi_{\bar{I}_{k,j}} \right\|_{\vec{p}/t}^{1/t}, \end{aligned}$$

where $0 < t < \min\{p_1, \dots, p_d, 1\}$. By Lemma 2 there exists a non-negative $g \in L_{(\vec{p}/t)'}$ with $\|g\|_{(\vec{p}/t)'} \leq 1$, such that

$$\begin{aligned} \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}}^t &= \int_{\Omega} \sum_{k \in \mathbb{Z}} \sum_{j \in \mathbb{N}} (3 \cdot 2^k)^t \chi_{\bar{I}_{k,j}} g \, d\mathbb{P} \\ &= \sum_{k \in \mathbb{Z}} \sum_{j \in \mathbb{N}} (3 \cdot 2^k)^t \int_{\bar{I}_{k,j}} \mathbb{E}_{j-1} g \, d\mathbb{P}. \end{aligned}$$

By (12),

$$\begin{aligned} \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}}^t &\leq R \sum_{k \in \mathbb{Z}} \sum_{j \in \mathbb{N}} (3 \cdot 2^k)^t \int_{I_{k,j}} \mathbb{E}_{j-1} g \, d\mathbb{P} \\ &\leq R \int_{\Omega} \sum_{k \in \mathbb{Z}} \sum_{j \in \mathbb{N}} (3 \cdot 2^k)^t \chi_{I_{k,j}} M(g) \, d\mathbb{P}. \end{aligned}$$

Since $(\vec{p}/t)' > 1$ and M is bounded on $L_{(\vec{p}/t)'}$, we conclude

$$\begin{aligned} \left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}}^t &\leq R \left\| \sum_{k \in \mathbb{Z}} \sum_{j \in \mathbb{N}} (3 \cdot 2^k)^t \chi_{I_{k,j}} \right\|_{\vec{p}/t} \|M(g)\|_{(\vec{p}/t)'} \\ &\leq CR \left\| \sum_{k \in \mathbb{Z}} (3 \cdot 2^k \chi_{\{\ell_k < \infty\}}) \right\|_{\vec{p}/t} \|g\|_{(\vec{p}/t)'} \\ &= CR \left\| \left[\sum_{k \in \mathbb{Z}} (3 \cdot 2^k \chi_{\{Mf > 2^k\}}) \right]^{1/t} \right\|_{\vec{p}}^t, \end{aligned}$$

where we have used that $\{\varrho_k < \infty\} = \{Mf > 2^k\}$. So we have that

$$\left\| \left(\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right)^{1/t} \right\|_{\vec{p}} \leq CR \left\| \left[\sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \chi_{\{Mf > 2^k\}} \right)^t \right]^{1/t} \right\|_{\vec{p}},$$

where the right-hand side can be estimated by $\|f\|_{H^M_{\vec{p}}}$, similarly as in the proof of Theorem 4.

Conversely, if f has a decomposition of the form (8), then

$$\|f\|_{H^M_{\vec{p}}} \leq \left\| \sum_{k \in \mathbb{Z}} \mu_k \frac{\chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right\|_{\vec{p}} \leq \left\| \left[\sum_{k \in \mathbb{Z}} \left(\mu_k \frac{\chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right]^{1/t} \right\|_{\vec{p}}$$

can be proved similarly as in Theorem 4. □

Lemma 4 *Let $0 < \vec{p} < \infty$ and the stochastic basis (\mathcal{F}_n) be regular. If ϱ_k and τ_k are the stopping times defined in the proof of Theorem 6, then*

$$\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}} \leq C \|\chi_{\{\varrho_k < \infty\}}\|_{\vec{p}}.$$

Proof It is enough to prove that for some $0 < \varepsilon < \min\{1, p_1, \dots, p_d\}$, the inequality

$$\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}/\varepsilon} \leq C \|\chi_{\{\varrho_k < \infty\}}\|_{\vec{p}/\varepsilon}$$

holds. Notice that

$$\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}/\varepsilon} \leq \left\| \sum_{j \in \mathbb{N}} \chi_{\bar{I}_{k,j}} \right\|_{\vec{p}/\varepsilon} \quad \text{and} \quad \|\chi_{\{\varrho_k < \infty\}}\|_{\vec{p}/\varepsilon} = \left\| \sum_{j \in \mathbb{N}} \chi_{I_{k,j}} \right\|_{\vec{p}/\varepsilon}.$$

By Lemma 2, there exists a non-negative function g with $\|g\|_{(\vec{p}/\varepsilon)'} \leq 1$, such that

$$\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}/\varepsilon} \leq \int_{\Omega} \sum_{j \in \mathbb{N}} \chi_{\bar{I}_{k,j}} g \, d\mathbb{P}.$$

Using (12), we obtain

$$\begin{aligned} \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}/\varepsilon} &= \sum_{j \in \mathbb{N}} \int_{\bar{I}_{k,j}} \mathbb{E}_{j-1} g \, d\mathbb{P} \\ &\leq R \sum_{j \in \mathbb{N}} \int_{I_{k,j}} \mathbb{E}_{j-1} g \, d\mathbb{P} \\ &\leq R \int_{\Omega} \sum_{j \in \mathbb{N}} \chi_{I_{k,j}} M(g) \, d\mathbb{P}. \end{aligned}$$

Lemma 1 implies

$$\begin{aligned} \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}/\varepsilon} &\leq R \left\| \sum_{j \in \mathbb{N}} \chi_{I_{k,j}} \right\|_{\vec{p}/\varepsilon} \|M(g)\|_{\vec{p}/\varepsilon} \\ &\leq CR \|\chi_{\{\varrho_k < \infty\}}\|_{\vec{p}/\varepsilon}, \end{aligned}$$

where we have used that $(\vec{p}/\varepsilon)' > 1$ and therefore M is bounded on $L_{(\vec{p}/\varepsilon)'}$. The proof is complete. \square

Corollary 2 *If the stochastic basis (\mathcal{F}_n) is regular, then*

$$H_{\vec{p}}^S = \mathcal{Q}_{\vec{p}} \quad \text{and} \quad H_{\vec{p}}^M = \mathcal{P}_{\vec{p}} \quad (0 < \vec{p} < \infty)$$

with equivalent quasi-norms.

5 Martingale Inequalities

We will prove the analogous version of the classical martingale inequalities (see, e.g., Weisz [40]) for the five mixed martingale Hardy spaces. To this end, we need the following boundedness results.

Let X be a martingale space, Y be a measurable function space. Then, the operator $U : X \rightarrow Y$ is called σ -sublinear operator, if for any $\alpha \in \mathbb{C}$,

$$\left| U \left(\sum_{k=1}^{\infty} f_k \right) \right| \leq \sum_{k=1}^{\infty} |U(f_k)| \quad \text{and} \quad |U(\alpha f)| = |\alpha| |U(f)|.$$

The σ -algebra generated by the stopping time τ is denoted by

$$\mathcal{F}_{\tau} = \{F \in \mathcal{F} : F \cap \{\tau \leq n\} \in \mathcal{F}_n, \quad n \geq 1\}.$$

\mathcal{F}_{τ} is a sub- σ -algebra of \mathcal{F} . Then, the conditional expectation with respect to \mathcal{F}_{τ} is denoted by \mathbb{E}_{τ} .

Theorem 7 *Let $0 < \vec{p} < \infty$ and suppose that the σ -sublinear operator $T : H_r^S \rightarrow L_r$ is bounded, where $\vec{p} = (p_1, \dots, p_d)$ and $r > p_i$ ($i = 1, \dots, d$). If for all (s, \vec{p}) -atom a*

$$(Ta)\chi_A = T(a\chi_A) \quad (A \in \mathcal{F}_{\tau}), \tag{13}$$

where τ is the stopping time associated with the (s, \vec{p}) -atom a , then for all $f \in H_{\vec{p}}^S$,

$$\|Tf\|_{\vec{p}} \leq C \|f\|_{H_{\vec{p}}^S}.$$

Proof By the σ -sublinearity of T and the atomic decomposition of $H_{\vec{p}}^s$ given in Theorem 4, we have

$$|Tf| \leq \sum_{k \in \mathbb{Z}} 3 \cdot 2^k \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}} \left| T(a^k) \right|,$$

If we choose $0 < t < \min \{p_1, \dots, p_d, 1\} \leq 1$, then

$$\begin{aligned} \|Tf\|_{\vec{p}} &\leq \left\| \left[\sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}} \left| T(a^k) \right| \right)^t \right]^{1/t} \right\|_{\vec{p}} \\ &= \left\| \sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \right)^t \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}^t \left| T(a^k) \right|^t \right\|_{\vec{p}/t}^{1/t}. \end{aligned}$$

By Lemma 2, there exists a function $g \in L_{(\vec{p}/t)'}$ with $\|g\|_{(\vec{p}/t)'} \leq 1$, such that

$$\|Tf\|_{\vec{p}}^t \leq \int_{\Omega} \sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \right)^t \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}^t \left| T(a^k) \right|^t g \, d\mathbb{P}.$$

Since $\{\tau_k < \infty\} \in \mathcal{F}_{\tau_k}$, using the fact that $a^k = a^k \chi_{\{\tau_k < \infty\}}$ and Eq. (13), we have $T(a^k \chi_{\{\tau_k < \infty\}}) = T(a^k) \chi_{\{\tau_k < \infty\}}$. Since $t < r$, the previous expression can be estimated by Hölder’s inequality

$$\begin{aligned} \|Tf\|_{\vec{p}}^t &\leq \sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \right)^t \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}^t \int_{\Omega} \chi_{\{\tau_k < \infty\}} \mathbb{E}_{\tau_k} \left(\left| T(a^k) \right|^t g \right) \, d\mathbb{P} \\ &\leq C \sum_{k \in \mathbb{Z}} \left(3 \cdot 2^k \right)^t \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}^t \int_{\Omega} \chi_{\{\tau_k < \infty\}} \\ &\quad \left[\mathbb{E}_{\tau_k} \left(\left| T(a^k) \right|^r \right) \right]^{t/r} \left[\mathbb{E}_{\tau_k} \left(|g|^{(r/t)'} \right) \right]^{1/(r/t)'} \, d\mathbb{P}. \end{aligned} \tag{14}$$

Here, by the boundedness of T and by the fact that a^k is an (s, \vec{p}) -atom, we get

$$\begin{aligned} \int_A |T(a^k)|^r \, d\mathbb{P} &= \int_{\Omega} |T(a^k \chi_A)|^r \, d\mathbb{P} \leq \int_{\Omega} |s(a^k \chi_A)|^r \, d\mathbb{P} \\ &\leq \int_A |s(a^k)|^r \, d\mathbb{P} \leq \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}^{-r} \mathbb{P}(A), \end{aligned}$$

where $A \in \mathcal{F}_{\tau_k}$. This implies that

$$\left[\mathbb{E}_{\tau_k} \left(\left| T(a^k) \right|^r \right) \right]^{t/r} \leq \|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}^{-t}.$$

Hence, (14) can be estimated by

$$\begin{aligned} \|Tf\|_{\vec{p}}^t &\leq C \int \sum_{k \in \mathbb{Z}} (3 \cdot 2^k)^t \chi_{\{\tau_k < \infty\}} \left[\mathbb{E}_{\tau_k} \left(|g|^{(r/t)'} \right) \right]^{1/(r/t)'} d\mathbb{P} \\ &\leq C \int \sum_{k \in \mathbb{Z}} (3 \cdot 2^k)^t \chi_{\{\tau_k < \infty\}} \left[M \left(|g|^{(r/t)'} \right) \right]^{1/(r/t)'} d\mathbb{P} \\ &\leq C \left\| \sum_{k \in \mathbb{Z}} (3 \cdot 2^k)^t \chi_{\{\tau_k < \infty\}} \right\|_{\vec{p}/t} \left\| \left[M \left(|g|^{(r/t)'} \right) \right]^{1/(r/t)'} \right\|_{(\vec{p}/t)'}. \end{aligned}$$

Since $\vec{p} < r$, therefore $(\vec{p}/t)'/(r/t)' > 1$, so by the boundedness of M (see Theorem 2), we obtain that

$$\begin{aligned} \left\| \left[M \left(|g|^{(r/t)'} \right) \right]^{1/(r/t)'} \right\|_{(\vec{p}/t)'} &= \left\| M \left(|g|^{(r/t)'} \right) \right\|_{(\vec{p}/t)'/(r/t)'}^{1/(r/t)'} \leq C \left\| |g|^{(r/t)'} \right\|_{(\vec{p}/t)'/(r/t)'}^{1/(r/t)'} \\ &= C \|g\|_{(\vec{p}/t)'} \leq C. \end{aligned}$$

By (9), we have

$$\|Tf\|_{\vec{p}} \leq C \left\| \sum_{k \in \mathbb{Z}} (3 \cdot 2)^t \chi_{\{\tau_k < \infty\}} \right\|_{\vec{p}/t}^{1/t} = C \left\| \left[\sum_{k \in \mathbb{Z}} \left(\frac{\mu_k \chi_{\{\tau_k < \infty\}}}{\|\chi_{\{\tau_k < \infty\}}\|_{\vec{p}}} \right)^t \right]^{1/t} \right\|_{\vec{p}} \leq C \|f\|_{H_{\vec{p}}^s}$$

and we get that T is bounded from $H_{\vec{p}}^s$ to $L_{\vec{p}}$. □

The following theorem can be proved similarly.

Theorem 8 *Let $0 < \vec{p} < \infty$ and suppose that the σ -sublinear operator $T : \mathcal{Q}_r \rightarrow L_r$ (resp. $T : \mathcal{P}_r \rightarrow L_r$) is bounded, where $\vec{p} = (p_1, \dots, p_d)$ and $r > p_i$ ($i = 1, \dots, d$). If all (S, \vec{p}) -atoms (resp. (M, \vec{p}) -atoms) satisfy (13), then for all $f \in \mathcal{Q}_{\vec{p}}$ (resp. $f \in \mathcal{P}_{\vec{p}}$),*

$$\|Tf\|_{\vec{p}} \leq C \|f\|_{\mathcal{Q}_{\vec{p}}} \quad \left(\text{resp. } \|f\|_{\mathcal{P}_{\vec{p}}}\right).$$

It is easy to see that for all (s, \vec{p}) -atoms a , (S, \vec{p}) -atoms a or (M, \vec{p}) -atoms a and $A \in \mathcal{F}_\tau$, $s(a\chi_A) = s(a)\chi_A$, $S(a\chi_A) = S(a)\chi_A$ and $M(a\chi_A) = M(a)\chi_A$. This means that the operators s , S and M satisfy condition (13).

Let $f \in H_{\vec{p}}^s$. The σ -sublinear operator M is bounded from H_2^s to L_2 (see, e.g., Weisz [40]), that is $\|Mf\|_2 \leq C \|f\|_{H_2^s}$. So we can apply Theorem 7 with the choice $r = 2$ and $\vec{p} := (p_1, \dots, p_d)$, where $p_i < 2$ and we get that

$$\|f\|_{H_{\vec{p}}^M} = \|M(f)\|_{\vec{p}} \leq C \|f\|_{H_{\vec{p}}^s} \quad (0 < \vec{p} < 2). \tag{15}$$

The operator S is also bounded from H_2^S to L_2 (see [40]), hence using Theorem 7, we obtain

$$\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{H_{\vec{p}}^s} \quad (0 < \vec{p} < 2). \tag{16}$$

From the definition of the Hardy spaces, it follows immediately that

$$\|f\|_{H_{\vec{p}}^M} \leq \|f\|_{\mathcal{P}_{\vec{p}}}, \quad \|f\|_{H_{\vec{p}}^S} \leq \|f\|_{\mathcal{Q}_{\vec{p}}} \quad (0 < \vec{p} < \infty). \tag{17}$$

By the Burkholder–Gundy and Doob’s inequality, for all $1 < r < \infty$, $\|S(f)\|_r \approx \|M(f)\|_r \approx \|f\|_r$ (see Weisz [40]). Using this, inequality (17) and Theorem 8, we have

$$\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{\mathcal{P}_{\vec{p}}} \quad \text{and} \quad \|f\|_{H_{\vec{p}}^M} \leq C \|f\|_{\mathcal{Q}_{\vec{p}}} \quad (0 < \vec{p} < \infty). \tag{18}$$

For $f = (f_n)_{n \in \mathbb{N}} \in \mathcal{Q}_{\vec{p}}$, there exists a sequence $(\lambda_n)_{n \in \mathbb{N}}$ for which $S_n(f) \leq \lambda_{n-1}$ and $\lambda_\infty \in L_{\vec{p}}$. Using the inequality $|f_n| \leq M_{n-1}(f) + \lambda_{n-1}$ and the second inequality in (18), we get that

$$\|f\|_{\mathcal{P}_{\vec{p}}} \leq \|M(f)\|_{\vec{p}} + \|\lambda_\infty\|_{\vec{p}} \leq \|f\|_{H_{\vec{p}}^M} + C \|f\|_{\mathcal{Q}_{\vec{p}}} \leq C \|f\|_{\mathcal{Q}_{\vec{p}}}. \tag{19}$$

Similarly, if $f = (f_n)_{n \in \mathbb{N}} \in \mathcal{P}_{\vec{p}}$, then $|f_n| \leq \lambda_{n-1}$ with a suitable sequence $(\lambda_n)_{n \in \mathbb{N}}$ for which $\lambda_\infty \in L_{\vec{p}}$. Since

$$S_n(f) = \left(\sum_{k=0}^n |d_k f|^2 \right)^{1/2} \leq S_{n-1}(f) + |d_n f| \leq S_{n-1}(f) + 2\lambda_{n-1},$$

using the first inequality in (18), we have that

$$\|f\|_{\mathcal{Q}_{\vec{p}}} \leq \|S(f)\|_{\vec{p}} + 2\|\lambda_\infty\|_{\vec{p}} = \|f\|_{H_{\vec{p}}^S} + 2\|f\|_{\mathcal{P}_{\vec{p}}} \leq C \|f\|_{\mathcal{P}_{\vec{p}}} \tag{20}$$

for all $0 < \vec{p} < \infty$.

From [40] Proposition 2.11 (ii), we get that the operator s is bounded from H_r^M to L_r and from H_r^S to L_r if $2 \leq r < \infty$. Again, using Theorem 8, we obtain

$$\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{\mathcal{P}_{\vec{p}}} \quad \text{and} \quad \|f\|_{H_{\vec{p}}^M} \leq C \|f\|_{\mathcal{Q}_{\vec{p}}} \quad (0 < \vec{p} < \infty). \tag{21}$$

The inequalities (15), (16), (17), (18), (19), (20), and (21) are collected in the following corollary.

Corollary 3 *We have the following martingale inequalities:*

(i) $\|f\|_{H_{\vec{p}}^M} \leq C \|f\|_{H_{\vec{p}}^s}, \quad \|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{H_{\vec{p}}^s} \quad (0 < \vec{p} < 2).$

(ii) $\|f\|_{H_{\vec{p}}^M} \leq \|f\|_{\mathcal{P}_{\vec{p}}}, \quad \|f\|_{H_{\vec{p}}^S} \leq \|f\|_{\mathcal{Q}_{\vec{p}}} \quad (0 < \vec{p} < \infty).$

- (iii) $\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{\mathcal{P}_{\vec{p}}}, \quad \|f\|_{H_{\vec{p}}^M} \leq C \|f\|_{\mathcal{Q}_{\vec{p}}} \quad (0 < \vec{p} < \infty).$
- (iv) $\|f\|_{\mathcal{P}_{\vec{p}}} \leq C \|f\|_{\mathcal{Q}_{\vec{p}}}, \quad \|f\|_{\mathcal{Q}_{\vec{p}}} \leq C \|f\|_{\mathcal{P}_{\vec{p}}} \quad (0 < \vec{p} < \infty).$
- (v) $\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{\mathcal{P}_{\vec{p}}} \quad \text{and} \quad \|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{\mathcal{Q}_{\vec{p}}} \quad (0 < \vec{p} < \infty).$ (22)

Theorem 9 *If the stochastic basis (\mathcal{F}_n) is regular, then the five Hardy spaces are equivalent, that is*

$$H_{\vec{p}}^S = \mathcal{Q}_{\vec{p}} = \mathcal{P}_{\vec{p}} = H_{\vec{p}}^M = H_{\vec{p}}^s \quad (0 < \vec{p} < \infty)$$

with equivalent quasi-norms.

Proof We know (see, e.g., Weisz [40]) that $S_n(f) \leq R^{1/2} s_n(f)$ and from this, it follows that $\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{H_{\vec{p}}^s}$. Using the definition of $\mathcal{Q}_{\vec{p}}$ and the fact that $s_n(f) \in \mathcal{F}_{n-1}$, we get

$$\|f\|_{\mathcal{Q}_{\vec{p}}} \leq C \|s(f)\|_{\vec{p}} = C \|f\|_{H_{\vec{p}}^s}.$$

By inequalities (22), we obtain that $\mathcal{Q}_{\vec{p}} = H_{\vec{p}}^s$. Since the stochastic basis (\mathcal{F}_n) is regular, the theorem follows from Corollary 2 and from (20). □

Theorem 10 *Suppose that $1 < \vec{p} < \infty$, or*

$$\vec{p} = (1, \dots, 1, p_{k+1}, \dots, p_d), \quad 1 < p_{k+1}, \dots, p_d < \infty \tag{23}$$

for some $k \in \{1, \dots, d\}$. Then for all non-negative, measurable function sequence $(f_n)_{n \in \mathbb{N}}$,

$$\left\| \sum_{n \in \mathbb{N}} \mathbb{E}_n(f_n) \right\|_{\vec{p}} \leq C \left\| \sum_{n \in \mathbb{N}} f_n \right\|_{\vec{p}}.$$

Proof From Lemma 2, we know that there exists a function $g \in L_{(\vec{p})'}$ with $\|g\|_{(\vec{p})'} \leq 1$ such that

$$\left\| \sum_{n \in \mathbb{N}} \mathbb{E}_n(f_n) \right\|_{\vec{p}} = \int_{\Omega} \sum_{n \in \mathbb{N}} \mathbb{E}_n(f_n) g \, d\mathbb{P}.$$

Since $\mathbb{E}_n(f_n)$ is \mathcal{F}_n -measurable, we obtain

$$\int_{\Omega} \sum_{n \in \mathbb{N}} \mathbb{E}_n(f_n) g \, d\mathbb{P} \leq \sum_{n \in \mathbb{N}} \int_{\Omega} f_n M(g) \, d\mathbb{P} = \int_{\Omega} \sum_{n \in \mathbb{N}} f_n M(g) \, d\mathbb{P}.$$

Using Hölder’s inequality and Theorem 2, we have

$$\left\| \sum_{n \in \mathbb{N}} \mathbb{E}_n(f_n) \right\|_{\vec{p}} \leq C \left\| \sum_{n \in \mathbb{N}} f_n \right\|_{\vec{p}} \|Mg\|_{(\vec{p})'} \leq C \left\| \sum_{n \in \mathbb{N}} f_n \right\|_{\vec{p}} \|g\|_{(\vec{p})'} \leq C \left\| \sum_{n \in \mathbb{N}} f_n \right\|_{\vec{p}}$$

and the proof is complete. □

As an application of the previous theorem, we get the following martingale inequality.

Corollary 4 *If $2 < \vec{p} < \infty$, or $\vec{p}/2$ satisfies (23), then*

$$\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{H_{\vec{p}}^S}.$$

Proof Indeed, using Theorem 10 with the choice $f_n := |d_n f|^2$, we have

$$\begin{aligned} \left\| \left(\sum_{n \in \mathbb{N}} \mathbb{E}_{n-1} |d_n f|^2 \right)^{1/2} \right\|_{\vec{p}} &= \left\| \sum_{n \in \mathbb{N}} \mathbb{E}_{n-1} |d_n f|^2 \right\|_{\vec{p}/2}^{1/2} \leq C \left\| \sum_{n \in \mathbb{N}} |d_n f|^2 \right\|_{\vec{p}/2}^{1/2} \\ &= C \|S(f)\|_{\vec{p}}, \end{aligned}$$

which finishes the proof. □

To prove the Burkholder–Davis–Gundy inequality, we introduce the norm

$$\|f\|_{\mathcal{G}_{\vec{p}}} := \left\| \sum_{n \in \mathbb{N}} |d_n f| \right\|_{\vec{p}}.$$

Lemma 5 *Suppose that $1 < \vec{p} < \infty$ or \vec{p} satisfies (23). If $f \in H_{\vec{p}}^S$, then there exists $h \in \mathcal{G}_{\vec{p}}$ and $g \in \mathcal{Q}_{\vec{p}}$ such that $f = h + g$ and*

$$\|h\|_{\mathcal{G}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^S} \quad \text{and} \quad \|g\|_{\mathcal{Q}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^S}.$$

Proof Let $f \in H_{\vec{p}}^S$ and let (λ_n) be an adapted, non-decreasing sequence such that $\lambda_0 = 0$, $S_n f \leq \lambda_n$ and $\lambda_\infty \in L_{\vec{p}}$. Let us define the functions

$$\begin{aligned} h_n &:= \sum_{k=1}^n (d_k f \chi_{\{\lambda_k > 2\lambda_{k-1}\}} - \mathbb{E}_{k-1} (d_k f \chi_{\{\lambda_k > 2\lambda_{k-1}\}})), \\ g_n &:= \sum_{k=1}^n (d_k f \chi_{\{\lambda_k \leq 2\lambda_{k-1}\}} - \mathbb{E}_{k-1} (d_k f \chi_{\{\lambda_k \leq 2\lambda_{k-1}\}})). \end{aligned}$$

Then $f_n = h_n + g_n$ ($n \in \mathbb{N}$). On the set $\{\lambda_k > 2\lambda_{k-1}\}$, we have $\lambda_k < 2(\lambda_k - \lambda_{k-1})$, henceforth

$$|d_k f| \chi_{\{\lambda_k > 2\lambda_{k-1}\}} \leq \lambda_k \chi_{\{\lambda_k > 2\lambda_{k-1}\}} \leq 2(\lambda_k - \lambda_{k-1})$$

and so

$$\sum_{k=1}^n |d_k h| \leq 2\lambda_n + 2 \sum_{k=1}^n \mathbb{E}_{k-1} (\lambda_k - \lambda_{k-1}).$$

Using Theorem 10, we have that

$$\|h\|_{\mathcal{G}_{\vec{p}}} \leq 2\|\lambda_\infty\|_{\vec{p}} + C \left\| \sum_{n \in \mathbb{N}} (\lambda_{n+1} - \lambda_n) \right\|_{\vec{p}} \leq C \|\lambda_\infty\|_{\vec{p}}.$$

At the same time,

$$|d_k f| \chi_{\{\lambda_k \leq 2\lambda_{k-1}\}} \leq \lambda_k \chi_{\{\lambda_k \leq 2\lambda_{k-1}\}} \leq 2\lambda_{k-1},$$

which implies that $|d_k g| \leq 4\lambda_{k-1}$. Then

$$\begin{aligned} S_n(g) &\leq S_{n-1}(g) + |d_n g| \leq S_{n-1}(f) + S_{n-1}(h) + 4\lambda_{n-1} \\ &\leq \lambda_{n-1} + 2\lambda_{n-1} + 2 \sum_{k=1}^{n-1} E_{k-1}(\lambda_k - \lambda_{k-1}) + 4\lambda_{n-1} \end{aligned}$$

and therefore

$$\|g\|_{\mathcal{Q}_{\vec{p}}} \leq C \|\lambda_\infty\|_{\vec{p}}.$$

Choosing $\lambda_n := S_n(f)$, we get that

$$\|h\|_{\mathcal{G}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^S} \quad \text{and} \quad \|g\|_{\mathcal{Q}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^S},$$

which proves the theorem. □

A similar lemma can be proved for $H_{\vec{p}}^M$ in the same way.

Lemma 6 *Suppose that $1 < \vec{p} < \infty$ or \vec{p} satisfies (23). If $f \in H_{\vec{p}}^M$, then there exists $h \in \mathcal{G}_{\vec{p}}$ and $g \in \mathcal{P}_{\vec{p}}$ such that $f = h + g$ and*

$$\|h\|_{\mathcal{G}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^M} \quad \text{and} \quad \|g\|_{\mathcal{P}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^M}.$$

Now we are ready to generalize the well-known Burkholder–Davis–Gundy inequality.

Theorem 11 *If $1 < \vec{p} < \infty$ or \vec{p} satisfies (23), then the spaces $H_{\vec{p}}^S$ and $H_{\vec{p}}^M$ are equivalent, that is*

$$H_{\vec{p}}^S = H_{\vec{p}}^M$$

with equivalent norms.

Proof Let $f \in H_{\vec{p}}^S$. By Lemma 5, there exists $h \in \mathcal{G}_{\vec{p}}$ and $g \in \mathcal{Q}_{\vec{p}}$ such that $f = h + g$ and

$$\|h\|_{\mathcal{G}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^S} \quad \text{and} \quad \|g\|_{\mathcal{Q}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^S}.$$

Then

$$\|f\|_{H_{\vec{p}}^M} \leq \|h\|_{H_{\vec{p}}^M} + \|g\|_{H_{\vec{p}}^M} \leq \|h\|_{\mathcal{G}_{\vec{p}}} + C \|g\|_{\mathcal{Q}_{\vec{p}}} \leq C \|f\|_{H_{\vec{p}}^S}.$$

The reverse inequality

$$\|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{H_{\vec{p}}^M}$$

can be proved similarly. □

For a martingale f , the *martingale transform* is defined by

$$(\mathcal{T}f)_n := \sum_{k=1}^{\infty} b_{k-1} d_k f,$$

where b_k are \mathcal{F}_k -measurable and $|b_k| \leq 1$. The martingale transform is bounded on $L_{\vec{p}}$, if $1 < \vec{p} < \infty$.

Theorem 12 *If $1 < \vec{p} < \infty$, then for all $f \in L_{\vec{p}}$,*

$$\|\mathcal{T}f\|_{\vec{p}} \leq C \|f\|_{\vec{p}}.$$

Proof Because of $|b_k| \leq 1$, it is clear that $S(\mathcal{T}f) \leq S(f)$. By Theorem 11, the spaces $H_{\vec{p}}^M$ and $H_{\vec{p}}^S$ are equivalent. Therefore using Theorem 2,

$$\|\mathcal{T}f\|_{\vec{p}} \leq \|\mathcal{T}f\|_{H_{\vec{p}}^M} \leq C \|\mathcal{T}f\|_{H_{\vec{p}}^S} \leq C \|f\|_{H_{\vec{p}}^S} \leq C \|f\|_{H_{\vec{p}}^M} \leq C \|f\|_{\vec{p}},$$

which proves the theorem. □

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