# Non-local to local transition for ground states of fractional Schrödinger equations on $\mathbb{R}^{N}$ 

Bartosz Bieganowski© and Simone Secchi©


#### Abstract

We consider the nonlinear fractional problem $$
(-\Delta)^{s} u+V(x) u=f(x, u) \quad \text { in } \mathbb{R}^{N}
$$

We show that ground state solutions converge (along a subsequence) in $L_{\text {loc }}^{2}\left(\mathbb{R}^{N}\right)$, under suitable conditions on $f$ and $V$, to a ground state solution of the local problem as $s \rightarrow 1^{-}$. Mathematics Subject Classification. Primary 35Q55; Secondary 35A15, 35R11. Keywords. Variational methods, fractional Schrödinger equation, nonlocal to local transition, ground state, Nehari manifold.


## 1. Introduction

The aim of this paper is to analyse the asymptotic behavior of least-energy solutions to the fractional Schrödinger problem:

$$
\left\{\begin{array}{l}
(-\Delta)^{s} u+V(x) u=f(x, u) \quad \text { in } \mathbb{R}^{N}  \tag{1.1}\\
u \in H^{s}\left(\mathbb{R}^{N}\right),
\end{array}\right.
$$

under suitable assumptions on the scalar potential $V: \mathbb{R}^{N} \rightarrow \mathbb{R}$ and on the nonlinearity $f: \mathbb{R}^{N} \times \mathbb{R} \rightarrow \mathbb{R}$. We recall that the fractional laplacian is defined as the principal value of a singular integral via the formula:

$$
(-\Delta)^{s} u(x)=C(N, s) \lim _{\varepsilon \rightarrow 0} \int_{\mathbb{R}^{N} \backslash B_{\varepsilon}(x)} \frac{u(x)-u(y)}{|x-y|^{N+2 s}} \mathrm{~d} y
$$

with

$$
\frac{1}{C(N, s)}=\int_{\mathbb{R}^{N}} \frac{1-\cos \zeta_{1}}{|\zeta|^{N+2 s}} \mathrm{~d} \zeta_{1} \cdots \mathrm{~d} \zeta_{N}
$$

This formal definition needs of course a function space in which problem (1.1) becomes meaningful: we will come to this issue in Sect. 2.

Several models have appeared in recent years that involve the use of the fractional laplacian. We only mention elasticity, turbulence, porous media
flow, image processing, wave propagation in heterogeneous high contrast media, and stochastic models: see $[1,11,13,19]$.

Instead of fixing the value of the parameter $s \in(0,1)$, we will start from the well-known identity (see [10, Proposition 4.4]):

$$
\begin{equation*}
\lim _{s \rightarrow 1^{-}}(-\Delta)^{s} u=-\Delta u \tag{1.2}
\end{equation*}
$$

valid for any $u \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$, and investigate the convergence properties of solutions to (1.1) as $s \rightarrow 1^{-}$.

In view of (1.2), it is somehow natural to conjecture that solutions to (1.1) converge to solutions of the problem:

$$
\left\{\begin{array}{l}
-\Delta u+V(x) u=f(x, u) \quad \text { in } \mathbb{R}^{N}  \tag{1.3}\\
u \in H^{1}\left(\mathbb{R}^{N}\right)
\end{array}\right.
$$

We do not know if this conjecture is indeed correct with this degree of generality.

In this paper, we will always assume that both $V$ and $f$ are $\mathbb{Z}^{N}$-periodic in the space variables. Hence equations (1.1) and (1.3) are invariant under $\mathbb{Z}^{N}$-translations, and their solutions are not unique. We will prove that-up to $\mathbb{Z}^{N}$-translations and along a subsequence-least-energy solutions of (1.1) converge to a ground state solution to the local problem (1.3). Our result is a continuation of the previous paper [5], in which we consider the equation on a bounded domain and extend the very recent analysis of Biccari et al. (see [2]) in the linear case for the Poisson problem to the semilinear case. See also [6].

We collect our assumptions.
(N) $N \geq 3,1 / 2<s<1$;
(V) $V \in L^{\infty}\left(\mathbb{R}^{N}\right)$ is $\mathbb{Z}^{N}$-periodic and $\inf _{\mathbb{R}^{N}} V>0$;
(F1) $f: \mathbb{R}^{N} \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function, namely $f(\cdot, u)$ is measurable for any $u \in \mathbb{R}$ and $f(x, \cdot)$ is continuous for a.e. $x \in \mathbb{R}^{N}$. Moreover $f$ is $\mathbb{Z}^{N}$-periodic in $x \in \mathbb{R}^{N}$ and there are numbers $C>0$ and $p \in\left(2, \frac{2 N}{N-1}\right)$ such that

$$
|f(x, u)| \leq C\left(1+|u|^{p-1}\right)
$$

for $u \in \mathbb{R}$ and a.e. $x \in \mathbb{R}^{N}$.
(F2) $f(x, u)=o(u)$ as $u \rightarrow 0$, uniformly with respect to $x \in \mathbb{R}^{N}$.
(F3) $\lim _{|u| \rightarrow+\infty} \frac{F(x, u)}{u^{2}}=+\infty$ uniformly with respect to $x \in \mathbb{R}^{N}$, where $F(x, u)=\int_{0}^{u} f(x, s) d s$.
(F4) The function $\mathbb{R} \backslash\{0\} \ni u \mapsto f(x, u) / u$ is strictly increasing on $(-\infty, 0)$ and on $(0, \infty)$, for a.e. $x \in \mathbb{R}^{N}$.

Remark 1.1. It follows from (F1) and (F2) that for every $\varepsilon>0$ there is $C_{\varepsilon}>0$ such that

$$
|f(x, u)| \leq \varepsilon|u|+C_{\varepsilon}|u|^{p-1}
$$

for every $u \in \mathbb{R}$ and a.e $x \in \mathbb{R}^{N}$. Furthermore, assumption (F4) implies the validity of the inequality:

$$
0 \leq 2 F(u) \leq f(x, u) u
$$

for every $u \in \mathbb{R}$ and a.e. $x \in \mathbb{R}^{N}$.
We can now state our main result.
Theorem 1.2. Suppose that assumptions $(N),(V),(F 1)-(F 4)$ hold. Let $u_{s} \in$ $H^{s}\left(\mathbb{R}^{N}\right)$ be a ground state solution of problem (1.1). Then, there exists a sequence $\left\{s_{n}\right\}_{n} \subset(1 / 2,1)$, such that $s_{n} \rightarrow 1$ as $n \rightarrow+\infty$ and there exists a sequence of translations $\left\{z_{n}\right\}_{n}$, such that $u_{s_{n}}\left(\cdot-z_{n}\right)$ converges in $L_{\mathrm{loc}}^{2}\left(\mathbb{R}^{N}\right)$ to a ground state solution $u_{0} \in H^{1}\left(\mathbb{R}^{N}\right)$ of the problem (1.3).

## 2. The variational setting

In this section we collect the basic tools from the theory of fractional Sobolev spaces we will need to prove our results. For a thorough discussion, we refer to $[10,14]$ and to the references therein.

For $0<s<1$, we define a Sobolev space on $\mathbb{R}^{N}$ as

$$
H^{s}\left(\mathbb{R}^{N}\right)=\left\{u \in L^{2}\left(\mathbb{R}^{N}\right) \left\lvert\, \int_{\mathbb{R}^{N} \times \mathbb{R}^{N}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 s}} \mathrm{~d} x \mathrm{~d} y<+\infty\right.\right\}
$$

endowed with the norm:

$$
\|u\|_{H^{s}\left(\mathbb{R}^{N}\right)}^{2}=\|u\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}+\int_{\mathbb{R}^{N} \times \mathbb{R}^{N}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 s}} \mathrm{~d} x \mathrm{~d} y
$$

One can show that $C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ is dense in $H^{s}\left(\mathbb{R}^{N}\right)$. For $u \in H^{s}\left(\mathbb{R}^{N}\right)$, an equivalent norm of $u$ is (see [14, Proposition 1.18])

$$
u \mapsto\left(\|u\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}+\left\|(-\Delta)^{\frac{s}{2}} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}\right)^{1 / 2}
$$

More explicitly, for every $u \in H^{s}\left(\mathbb{R}^{N}\right)$

$$
\int_{\mathbb{R}^{N} \times \mathbb{R}^{N}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 s}} \mathrm{~d} x \mathrm{~d} y=\frac{2}{C(N, s)}\left\|(-\Delta)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2},
$$

where

$$
\begin{aligned}
C(N, s) & =\frac{s(1-s)}{A(N, s) B(s)} \\
A(N, s) & =\int_{\mathbb{R}^{N-1}} \frac{\mathrm{~d} \eta}{\left(1+|\eta|^{2}\right)^{(N+2 s) / 2}} \\
B(s) & =s(1-s) \int_{\mathbb{R}} \frac{1-\cos t}{|t|^{1+2 s}} \mathrm{~d} t .
\end{aligned}
$$

Lemma 2.1. For every $u \in H^{1}\left(\mathbb{R}^{N}\right)$, there results

$$
\lim _{s \rightarrow 1^{-}}\left\|(-\Delta)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}=\|\nabla u\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2} .
$$

Proof. From [10, Proposition 3.6], we know that

$$
\left\|(-\Delta)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}=\frac{C(N, s)}{2} \int_{\mathbb{R}^{N} \times \mathbb{R}^{N}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 s}} \mathrm{~d} x \mathrm{~d} y .
$$

From [10, Remark 4.3], we know that

$$
\lim _{s \rightarrow 1^{-}}(1-s) \int_{\mathbb{R}^{N} \times \mathbb{R}^{N}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 s}} \mathrm{~d} x \mathrm{~d} y=\frac{\omega_{N-1}}{2 N}\|\nabla u\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}
$$

Therefore, recalling [10, Corollary 4.2],

$$
\begin{aligned}
\lim _{s \rightarrow 1^{-}}\left\|(-\Delta)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2} & =\lim _{s \rightarrow 1^{-}} \frac{C(N, s)}{2(1-s)}\left((1-s) \int_{\mathbb{R}^{N} \times \mathbb{R}^{N}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 s}} \mathrm{~d} x \mathrm{~d} y\right) \\
& =\frac{1}{2} \frac{4 N}{\omega_{N-1}} \frac{\omega_{N-1}}{2 N}\|\nabla u\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}=\|\nabla u\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2} .
\end{aligned}
$$

On $H^{s}\left(\mathbb{R}^{N}\right)$ we introduce a new norm

$$
\begin{equation*}
\|u\|_{s}^{2}:=\left\|(-\Delta)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}+\int_{\mathbb{R}^{N}} V(x) u^{2} \mathrm{~d} x, \quad u \in H^{s}\left(\mathbb{R}^{N}\right) \tag{2.1}
\end{equation*}
$$

which is, under $(\mathrm{V})$, equivalent to $\|\cdot\|_{H^{s}\left(\mathbb{R}^{N}\right)}$. Similarly we introduce the norm on $H^{1}\left(\mathbb{R}^{N}\right)$ by putting

$$
\begin{equation*}
\|u\|^{2}:=\int_{\mathbb{R}^{N}}|\nabla u|^{2}+V(x) u^{2} \mathrm{~d} x, \quad u \in H^{1}\left(\mathbb{R}^{N}\right) . \tag{2.2}
\end{equation*}
$$

Corollary 2.2. For every $u \in H^{1}\left(\mathbb{R}^{N}\right)$, we have

$$
\lim _{s \rightarrow 1^{-}}\|u\|_{s}=\|u\| .
$$

The following convergence result will be used in the sequel.
Lemma 2.3. For every $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$, there results

$$
\lim _{s \rightarrow 1^{-}}\left\|(-\Delta)^{s} \varphi-(-\Delta) \varphi\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}=0
$$

Proof. We notice that

$$
\begin{aligned}
& \left\|(-\Delta)^{s} \varphi-(-\Delta) \varphi\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}=\left\|\mathcal{F}_{\xi}^{-1}\left(\left(|\xi|^{2 s}-|\xi|^{2}\right) \hat{\varphi}(\xi)\right)\right\|_{L^{2}\left(\mathbb{R}^{N}\right)} \\
& \quad \leq C\left\|\left(|\cdot|^{2 s}-|\cdot|^{2}\right) \hat{\varphi}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}
\end{aligned}
$$

where $C>0$ is a constant, independent of $s$, that depends on the definition of the Fourier transform $\mathcal{F}$. It is now easy to conclude, since the Fourier transform of a test function is a rapidly decreasing function.

We will need some precise information on the embedding constant for fractional Sobolev spaces.

Theorem 2.4. [9] Let $N>2 s$ and $2_{s}^{*}=2 N /(N-2 s)$. Then

$$
\|u\|_{L^{2 *}\left(\mathbb{R}^{N}\right)}^{2} \leq \frac{\Gamma\left(\frac{N-2 s}{2}\right)}{\Gamma\left(\frac{N+2 s}{2}\right)}|\mathbb{S}|^{-\frac{2 s}{N}}\left\|(-\Delta)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}
$$

for every $u \in H^{s}\left(\mathbb{R}^{N}\right)$, where $\mathbb{S}$ denotes the $N$-dimensional unit sphere and $|\mathbb{S}|$ its surface area.

The following inequality in an easy consequence of Theorem 2.4, see also [5, Lemma 2.7].

Lemma 2.5. Let $N \geq 3$ and $q \in[2,2 N /(N-1)]$. Then there exists a constant $C=C(N, q)>0$ such that, for every $s \in[1 / 2,1]$ and every $u \in H^{s}\left(\mathbb{R}^{N}\right)$, we have

$$
\|u\|_{L^{q}\left(\mathbb{R}^{N}\right)} \leq C(N, q)\|u\|_{s}
$$

Definition 2.6. A weak solution to problem (1.1) is a function $u \in H^{s}\left(\mathbb{R}^{N}\right)$, such that

$$
\left\langle(-\Delta)^{s / 2} u \mid(-\Delta)^{s / 2} \varphi\right\rangle_{L^{2}\left(\mathbb{R}^{N}\right)}+\int_{\mathbb{R}^{N}} V(x) u \varphi \mathrm{~d} x=\int_{\mathbb{R}^{N}} f(x, u) \varphi \mathrm{d} x
$$

for every $\varphi \in H^{s}\left(\mathbb{R}^{N}\right)$.
Weak solutions are therefore critical points of the associated energy functional $\mathcal{J}_{s}: H^{s}\left(\mathbb{R}^{N}\right) \rightarrow \mathbb{R}$ defined by

$$
\mathcal{J}_{s}(u)=\frac{1}{2}\left\|(-\Delta)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}+\frac{1}{2} \int_{\mathbb{R}^{N}} V(x) u^{2} \mathrm{~d} x-\int_{\mathbb{R}^{N}} F(x, u) \mathrm{d} x
$$

We recall also the definition of a weak solution in the local case.
Definition 2.7. A weak solution to problem (1.3) is a function $u \in H^{1}\left(\mathbb{R}^{N}\right)$ such that

$$
\int_{\mathbb{R}^{N}} \nabla u \cdot \nabla \varphi \mathrm{~d} x+\int_{\mathbb{R}^{N}} V(x) u \varphi \mathrm{~d} x=\int_{\mathbb{R}^{N}} f(x, u) \varphi \mathrm{d} x
$$

for every $\varphi \in H^{1}\left(\mathbb{R}^{N}\right)$.
For the local problem (1.3), we put $\mathcal{J}: H^{1}\left(\mathbb{R}^{N}\right) \rightarrow \mathbb{R}$

$$
\begin{equation*}
\mathcal{J}(u)=\frac{1}{2} \int_{\mathbb{R}^{N}}|\nabla u|^{2}+V(x) u^{2} \mathrm{~d} x-\int_{\mathbb{R}^{N}} F(x, u) \mathrm{d} x \tag{2.3}
\end{equation*}
$$

Recalling the notation (2.1) and (2.2), we can rewrite our functionals in the form:

$$
\begin{aligned}
\mathcal{J}_{s}(u) & =\frac{1}{2}\|u\|_{s}^{2}-\int_{\mathbb{R}^{N}} F(x, u) \mathrm{d} x, \\
\mathcal{J}(u) & =\frac{1}{2}\|u\|^{2}-\int_{\mathbb{R}^{N}} F(x, u) \mathrm{d} x,
\end{aligned} \quad u \in H^{N}\left(\mathbb{R}^{N}\right),
$$

## 3. Uniform Lions' concentration-compactness principle

Since the summability exponent of our space is not fixed, we need a "uniform" version of a celebrated result by Lions.
Theorem 3.1. Let $r>0,2 \leq q<\frac{2 N}{N-1}$ and $N \geq 3$. Suppose moreover that $\left\{s_{n}\right\}_{n} \subset(1 / 2,1), u_{n} \in H^{s_{n}}\left(\mathbb{R}^{N}\right)$ and

$$
\left\|u_{n}\right\|_{s_{n}} \leq M
$$

where $M>0$ does not depend on $s_{n}$. If

$$
\lim _{n \rightarrow+\infty} \sup _{y \in \mathbb{R}^{N}} \int_{B(y, r)}\left|u_{n}\right|^{q} \mathrm{~d} x=0
$$

then $u_{n} \rightarrow 0$ in $L^{p}\left(\mathbb{R}^{N}\right)$ for all $p \in\left(2, \frac{2 N}{N-1}\right)$.
Proof. Let $t \in\left(q, \frac{2 N}{N-1}\right)$. Then

$$
\begin{aligned}
\left\|u_{n}\right\|_{L^{t}(B(y, r))} & \leq\left\|u_{n}\right\|_{L^{q}(B(y, r))}^{1-\lambda}\left\|u_{n}\right\|_{L^{\frac{2 N}{N-1}}(B(y, r))}^{\lambda} \\
& \leq C\left\|u_{n}\right\|_{L^{q}(B(y, r))}^{1-\lambda}\left\|u_{n}\right\|_{s_{n}}^{\lambda},
\end{aligned}
$$

where $C>0$ is independent of $s_{n}$ and $\lambda=\frac{t-q}{\frac{2 N}{N-1}-q} \frac{2 N}{(N-1) t}$. Choose $t$ such that $\lambda=\frac{2}{t}$. Then

$$
\int_{B(y, r)}\left|u_{n}\right|^{t} \mathrm{~d} x \leq C^{t}\left\|u_{n}\right\|_{L^{q}(B(y, r))}^{(1-\lambda) t}\left\|u_{n}\right\|_{s_{n}}^{2}
$$

Covering space $\mathbb{R}^{N}$ by balls of radius $r$, in a way that each point is contained in at most $N+1$ balls, we get

$$
\begin{aligned}
\int_{\mathbb{R}^{N}}\left|u_{n}\right|^{t} \mathrm{~d} x & \leq(N+1) C^{t} \sup _{y \in \mathbb{R}^{N}}\left(\int_{B(y, r)}\left|u_{n}\right|^{q} \mathrm{~d} x\right)^{\frac{(1-\lambda) t}{q}}\left\|u_{n}\right\|_{s_{n}}^{2} \\
& \leq(N+1) M^{2} C^{t} \sup _{y \in \mathbb{R}^{N}}\left(\int_{B(y, r)}\left|u_{n}\right|^{q} \mathrm{~d} x\right)^{\frac{(1-\lambda) t}{q}} \rightarrow 0
\end{aligned}
$$

Hence $u_{n} \rightarrow 0$ in $L^{t}\left(\mathbb{R}^{N}\right)$. Note that

$$
\left\|u_{n}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2} \leq D\left\|u_{n}\right\|_{s_{n}}^{2} \leq D M^{2}
$$

where $D$ does not depend on $s_{n}$ and $n$. Similarly, from Lemma 2.5 , there follows that $\left\{u_{n}\right\}_{n}$ is bounded in $L^{\frac{2 N}{N-1}}\left(\mathbb{R}^{N}\right)$. From the interpolation inequality, since $\left\{u_{n}\right\}_{n}$ is bounded in $L^{2}\left(\mathbb{R}^{N}\right)$ and in $L^{\frac{2 N}{N-1}}\left(\mathbb{R}^{N}\right)$, we obtain $u_{n} \rightarrow 0$ in $L^{p}\left(\mathbb{R}^{N}\right)$ for all $p \in\left(2, \frac{2 N}{N-1}\right)$.

Finally, we extend the locally compact embedding into Lebesgue spaces in a uniform way.

Theorem 3.2. Let $\left\{s_{n}\right\}_{n}$ be a sequence such that $1 / 2<s_{n}<1$ and $s_{n} \rightarrow 1$, and let $\left\{v_{s_{n}}\right\}_{n} \subset H^{s_{n}}\left(\mathbb{R}^{N}\right)$ be such that

$$
M=\sup _{n}\left\|v_{s_{n}}\right\|_{s_{n}}<\infty
$$

Then, the sequence $\left\{v_{s_{n}}\right\}_{n}$ converges, up to a subsequence, to some $v \in$ $H^{1}\left(\mathbb{R}^{N}\right)$ in $L_{\mathrm{loc}}^{q}\left(\mathbb{R}^{N}\right)$ for every $q \in[2,2 N /(N-1)$ ), and pointwise almost everywhere.

Proof. Note that $H^{s_{n}}\left(\mathbb{R}^{N}\right) \subset H^{1 / 2}\left(\mathbb{R}^{N}\right)$ and

$$
\|\cdot\|_{1 / 2} \leq C\|\cdot\|_{s_{n}}
$$

where $C>0$ does not depend on $s_{n}$ (and therefore also on $n$ ): see for instance [14, Proposition 1.1]. In particular, for every $n \in \mathbb{N}$, we have

$$
\begin{equation*}
\left\|v_{s_{n}}\right\|_{1 / 2} \leq C\left\|v_{s_{n}}\right\|_{s_{n}} \leq C M \tag{3.1}
\end{equation*}
$$

Thus, $\left\{v_{s_{n}}\right\}_{n}$ is bounded in $H^{1 / 2}\left(\mathbb{R}^{N}\right)$. Hence, passing to a subsequence, there exists a function $v$, such that $v_{s_{n}} \rightharpoonup v$ in $H^{1 / 2}\left(\mathbb{R}^{N}\right), v_{s_{n}} \rightarrow v$ pointwise almost everywhere, and $v_{s_{n}} \rightarrow v$ in $L_{\mathrm{loc}}^{q}\left(\mathbb{R}^{N}\right)$ for every $q \in[2,2 N /(N-1))$. From [7, Corollary 7], it follows that $v \in H_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N}\right)$. To complete the proof, we need to show that $v \in H^{1}\left(\mathbb{R}^{N}\right)$.

Let $\widehat{v_{s_{n}}}$ denote the Fourier transform of $v_{s_{n}}$, similarly for $\widehat{v}$. We may assume, without loss of generality, that $\widehat{v_{s_{n}}} \rightharpoonup \widehat{v}$ in $L^{2}\left(\mathbb{R}^{N}\right)$. Note that (3.1) implies that

$$
\sup _{n} \int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)^{s_{n}}\left|\widehat{v_{s_{n}}}\right|^{2} \mathrm{~d} \xi \leq K
$$

for some constant $K>0$. For $1 / 2<t \leq 1$, we define

$$
B_{t}:=\left\{\left.w \in L^{2}\left(\mathbb{R}^{N}\right)\left|\int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)^{t}\right| w(\xi)\right|^{2} \mathrm{~d} \xi \leq K\right\}
$$

First of all, we observe that

$$
\begin{equation*}
\bigcap_{1 / 2<t<1} B_{t}=B_{1} \tag{3.2}
\end{equation*}
$$

Indeed, for any $1 / 2<t<1$ we have $\left(1+|\xi|^{2}\right)^{t} \leq 1+|\xi|^{2}$. Take $w \in B_{1}$ and note that

$$
\int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)^{t}|w(\xi)|^{2} \mathrm{~d} \xi \leq \int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)|w(\xi)|^{2} \mathrm{~d} \xi \leq K
$$

Hence $w \in B_{t}$ for any $t<1$. Thus

$$
\bigcap_{1 / 2<t<1} B_{t} \supset B_{1}
$$

On the other hand, fix $w \in \bigcap_{1 / 2<t<1} B_{t}$. Take any sequence $t_{n} \rightarrow 1^{-}$with $t_{n}>1 / 2$. Then, obviously

$$
\liminf _{n \rightarrow+\infty}\left(1+|\xi|^{2}\right)^{t_{n}}|w(\xi)|^{2}=\left(1+|\xi|^{2}\right)|w(\xi)|^{2}
$$

and Fatou's lemma yields

$$
\int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)|w(\xi)|^{2} \mathrm{~d} \xi \leq \liminf _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)^{t_{n}}|w(\xi)|^{2} \mathrm{~d} \xi \leq K
$$

Hence $w \in B_{1}$, or

$$
\bigcap_{1 / 2<t<1} B_{t} \subset B_{1}
$$

and (3.2) is proved. Fix now any $t \in(1 / 2,1)$ and choose $n_{0}$ such that $s_{n}>t$ for all $n \geq n_{0}$. Then

$$
\left(1+|\xi|^{2}\right)^{t} \leq\left(1+|\xi|^{2}\right)^{s_{n}} \quad \text { for every } \xi \in \mathbb{R}^{N}
$$

and

$$
\int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)^{t}\left|\widehat{v_{s_{n}}}\right|^{2} \mathrm{~d} \xi \leq \int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)^{s_{n}}\left|\widehat{v_{s_{n}}}\right|^{2} \mathrm{~d} \xi \leq K \quad \text { for } n \geq n_{0}
$$

Hence, $\widehat{v_{s_{n}}} \in B_{t}$ for $n \geq n_{0}$. Each $B_{t}$ is a closed and convex subset in $L^{2}\left(\mathbb{R}^{N}\right)$, and from $\left[8\right.$, Theorem 3.7] it is also weakly closed. Hence, $\widehat{v} \in B_{t}$. Therefore, recalling (3.2),

$$
\widehat{v} \in \bigcap_{\frac{1}{2}<t<1} B_{t}=\left\{\left.w \in L^{2}\left(\mathbb{R}^{N}\right)\left|\int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)\right| w(\xi)\right|^{2} \mathrm{~d} \xi \leq K\right\}
$$

This implies that

$$
\int_{\mathbb{R}^{N}}\left(1+|\xi|^{2}\right)|\hat{v}(\xi)|^{2} \mathrm{~d} \xi \leq K<+\infty
$$

and $v \in H^{1}\left(\mathbb{R}^{N}\right)$.

## 4. Existence of ground states

It is easy to check that the energy functional $\mathcal{J}$ has the mountain-pass geometry. In particular, there is radius $r>0$, such that

$$
\inf _{\|u\|=r} \mathcal{J}(u)>0
$$

The following existence result is well-known in the literature, and has been shown in various ways, see e.g. $[4,12,17,18]$.

Theorem 4.1. Suppose that assumptions ( $N$ ), ( $V$ ), ( $F 1$ )-(F4) hold. Then there exists a ground state solution $u_{0} \in H^{1}\left(\mathbb{R}^{N}\right)$ to (1.3), i.e., a critical point of the functional $\mathcal{J}$ given by (2.3), such that

$$
\mathcal{J}\left(u_{0}\right)=\inf _{\mathcal{N}} \mathcal{J}=\inf _{u \in H^{1}\left(\mathbb{R}^{N}\right) \backslash\{0\}} \sup _{t \geq 0} \mathcal{J}(t u)=\inf _{\gamma \in \Gamma} \sup _{t \in[0,1]} \mathcal{J}(\gamma(t))
$$

where $\mathcal{N}$ is the so-called Nehari manifold

$$
\mathcal{N}:=\left\{u \in H^{1}\left(\mathbb{R}^{N}\right) \backslash\{0\} \mid \mathcal{J}^{\prime}(u)(u)=0\right\}
$$

and

$$
\Gamma:=\left\{\gamma \in \mathcal{C}\left([0,1], H^{1}\left(\mathbb{R}^{N}\right)\right) \mid \gamma(0)=0,\|\gamma(1)\|>r, \mathcal{J}(\gamma(1))<0\right\}
$$

The same methods can be applied also in the nonlocal case, and the following existence result can be shown, see e.g. $[3,15,16]$. In what follows, $r_{s}>0$ is the radius chosen so that

$$
\inf _{\|u\|_{s}=r_{s}} \mathcal{J}_{s}(u)>0
$$

Theorem 4.2. Suppose that assumptions (N), (V), (F1)-(F4) hold and $1 / 2<$ $s<1$. Then there exists a ground state solution $u_{s} \in H^{s}\left(\mathbb{R}^{N}\right)$ to (1.1), i.e. a critical point of the functional $\mathcal{J}_{\text {s }}$ given by (2.3), such that

$$
\begin{equation*}
\mathcal{J}_{s}\left(u_{s}\right)=\inf _{\mathcal{N}_{s}} \mathcal{J}_{s}=\inf _{u \in H^{s}\left(\mathbb{R}^{N}\right) \backslash\{0\}} \sup _{t \geq 0} \mathcal{J}_{s}(t u)=\inf _{\gamma \in \Gamma_{s}} \sup _{t \in[0,1]} \mathcal{J}_{s}(\gamma(t)) \tag{4.1}
\end{equation*}
$$

where $\mathcal{N}_{s}$ is the corresponding Nehari manifold

$$
\mathcal{N}_{s}:=\left\{u \in H^{s}\left(\mathbb{R}^{N}\right) \backslash\{0\} \mid \mathcal{J}_{s}^{\prime}(u)(u)=0\right\}
$$

and

$$
\Gamma_{s}:=\left\{\gamma \in \mathcal{C}\left([0,1], H^{s}\left(\mathbb{R}^{N}\right)\right) \mid \gamma(0)=0,\|\gamma(1)\|>r_{s}, \mathcal{J}_{s}(\gamma(1))<0\right\} .
$$

## 5. Non-local to local transition

For any $s \in(1 / 2,1)$ we define

$$
c_{s}:=\inf _{\mathcal{N}_{s}} \mathcal{J}_{s}>0
$$

Similarly, we put also

$$
c:=\inf _{\mathcal{N}} \mathcal{J}>0
$$

For any $v \in H^{s}\left(\mathbb{R}^{N}\right) \backslash\{0\}$ let $t_{s}(v)>0$ be the unique positive real number such that $t_{s}(v) \in \mathcal{N}_{s}$. Then, we put $m_{s}(v):=t_{s}(v) v$.

Lemma 5.1. There results

$$
\limsup _{s \rightarrow 1^{-}} c_{s} \leq c
$$

Proof. Take $u \in H^{1}\left(\mathbb{R}^{N}\right) \subset H^{s}\left(\mathbb{R}^{N}\right)$ as a ground state solution of (1.3), in particular $u \in \mathcal{N}$ and $\mathcal{J}(u)=c$, where $\mathcal{J}$ is given by (2.3). Consider the function $m_{s}(u) \in \mathcal{N}_{s}$. Obviously

$$
c_{s} \leq \mathcal{J}_{s}\left(m_{s}(u)\right)
$$

Hence

$$
\begin{aligned}
& \limsup _{s \rightarrow 1^{-}} c_{s} \leq \limsup _{s \rightarrow 1^{-}} \mathcal{J}_{s}\left(m_{s}(u)\right) \\
& \quad=\limsup _{s \rightarrow 1^{-}}\left\{\mathcal{J}_{s}\left(m_{s}(u)\right)-\frac{1}{2} \mathcal{J}_{s}^{\prime}\left(m_{s}(u)\right)\left(m_{s}(u)\right)\right\} \\
& \quad=\limsup _{s \rightarrow 1^{-}}\left\{\frac{1}{2} \int_{\mathbb{R}^{N}} f\left(x, m_{s}(u)\right) m_{s}(u)-2 F\left(x, m_{s}(u)\right) \mathrm{d} x\right\} .
\end{aligned}
$$

Recall that $m_{s}(u)=t_{s} u$ for some real numbers $t_{s}>0$. Suppose by contradiction that $t_{s} \rightarrow+\infty$ as $s \rightarrow 1^{-}$. Then, in view of the Nehari identity

$$
\|u\|_{s}^{2}=\int_{\mathbb{R}^{N}} \frac{f\left(x, t_{s} u\right)}{t_{s}^{2}} t_{s} u \mathrm{~d} x \geq 2 \int_{\mathbb{R}^{N}} \frac{F\left(x, t_{s} u\right)}{t_{s}^{2} u^{2}} u^{2} \mathrm{~d} x \rightarrow+\infty
$$

but the left-hand side stays bounded (see Corollary 2.2). Hence $\left(t_{s}\right)_{s}$ is bounded. Take any convergent subsequence $\left(t_{s_{n}}\right)$ of $\left(t_{s}\right)$, i.e. $t_{s_{n}} \rightarrow t_{0}$ as
$n \rightarrow+\infty$. Obviously $t_{0} \geq 0$. We will show that $t_{0} \neq 0$. Indeed, suppose that $t_{0}=0$, i.e. $t_{s_{n}} \rightarrow 0$. Then, in view of the Nehari identity

$$
\|u\|_{s_{n}}^{2}=\int_{\mathbb{R}^{N}} \frac{f\left(x, t_{s_{n}} u\right)}{t_{s_{n}} u} u^{2} \mathrm{~d} x
$$

By Corollary 2.2, $\|u\|_{s_{n}}^{2} \rightarrow\|u\|^{2}>0$. Hence, in view of (F2),

$$
\|u\|^{2}+o(1)=\int_{\mathbb{R}^{N}} \frac{f\left(x, t_{s_{n}} u\right)}{t_{s_{n}} u} u^{2} \mathrm{~d} x \rightarrow 0
$$

a contradiction. Hence $t_{0}>0$. Again, by Corollary 2.2,

$$
t_{s_{n}}^{2}\|u\|_{s_{n}}^{2} \rightarrow t_{0}^{2}\|u\|^{2} \quad \text { as } n \rightarrow+\infty
$$

Moreover, in view of Remark 1.1,

$$
\left|f\left(x, t_{s_{n}} u\right) t_{s_{n}} u\right| \leq \varepsilon t_{s_{n}}^{2}|u|^{2}+C_{\varepsilon} t_{s_{n}}^{p}|u|^{p} \leq C\left(|u|^{2}+|u|^{p}\right)
$$

for some constant $C>0$, independent of $n$. In view of the Lebesgue's convergence theorem,

$$
\int_{\mathbb{R}^{N}} f\left(x, t_{s_{n}} u\right) t_{s_{n}} u \mathrm{~d} x \rightarrow \int_{\mathbb{R}^{N}} f\left(x, t_{0} u\right) t_{0} u \mathrm{~d} x
$$

Thus, the limit $t_{0}$ satisfies

$$
t_{0}^{2}\|u\|^{2}=\int_{\mathbb{R}^{N}} f\left(x, t_{0} u\right) t_{0} u \mathrm{~d} x
$$

Taking the Nehari identity into account we see that $t_{0}=1$. Hence $t_{s} \rightarrow 1$ as $s \rightarrow 1^{-}$. Repeating the same argument we see that

$$
\begin{aligned}
& \limsup _{s \rightarrow 1^{-}}\left\{\frac{1}{2} \int_{\mathbb{R}^{N}} f\left(x, m_{s}(u)\right) m_{s}(u)-2 F\left(x, m_{s}(u)\right) \mathrm{d} x\right\} \\
& \quad=\frac{1}{2} \int_{\mathbb{R}^{N}} f(x, u) u-2 F(x, u) \mathrm{d} x=\mathcal{J}(u)=c
\end{aligned}
$$

and the proof is completed.
Lemma 5.2. There exists a constant $M>0$, such that

$$
\left\|u_{s}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}+\left\|u_{s}\right\|_{s}+\left\|u_{s}\right\|_{L^{\frac{2 N}{N-1}}\left(\mathbb{R}^{N}\right)} \leq M
$$

for every $s \in(1 / 2,1)$.
Proof. Note that $\left\|u_{s}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}+\left\|u_{s}\right\|_{L^{\frac{2 N}{N-1}}\left(\mathbb{R}^{N}\right)} \leq C\left\|u_{s}\right\|_{s}$, for some $C>0$ independent of $s$. So it is enough to show that $\left\|u_{s}\right\|_{s} \leq M$. Suppose by contradiction that

$$
\left\|u_{s}\right\|_{s} \rightarrow+\infty \quad \text { as } s \rightarrow 1^{-}
$$

Put $v_{s}:=\frac{u_{s}}{\left\|u_{s}\right\|_{s}}$. Then $\left\|v_{s}\right\|_{s}=1$. In particular, $\left\{v_{s}\right\}$ is bounded in $L^{2}\left(\mathbb{R}^{N}\right)$. Suppose that

$$
\begin{equation*}
\sup _{y \in \mathbb{R}^{N}} \int_{B(y, 1)}\left|v_{s}\right|^{2} \mathrm{~d} x \rightarrow 0 \tag{5.1}
\end{equation*}
$$

Then $v_{s} \rightarrow 0$ in $L^{p}\left(\mathbb{R}^{N}\right)$. Fix any $t>0$. By (4.1) we obtain

$$
\mathcal{J}_{s_{n}}\left(u_{s_{n}}\right) \geq \mathcal{J}_{s_{n}}\left(\frac{t}{\left\|u_{s_{n}}\right\|_{s_{n}}} u_{s_{n}}\right)=\mathcal{J}_{s_{n}}\left(t v_{s_{n}}\right)=\frac{t^{2}}{2}-\int_{\mathbb{R}^{N}} F\left(x, t v_{s_{n}}\right) \mathrm{d} x
$$

From Remark 1.1 we see that

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} F\left(x, t v_{s_{n}}\right) \mathrm{d} x \leq \varepsilon t^{2}\left\|v_{s_{n}}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}+C_{\varepsilon} t^{p}\left\|v_{s_{n}}\right\|_{L^{p}\left(\mathbb{R}^{N}\right)}^{p} \\
& \quad \rightarrow \varepsilon t^{2} \limsup _{n \rightarrow \infty}\left\|v_{s_{n}}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}
\end{aligned}
$$

for every $\varepsilon>0$. Thus $\int_{\mathbb{R}^{N}} F\left(x, t v_{s_{n}}\right) \mathrm{d} x \rightarrow 0$ and for any $t>0$

$$
\mathcal{J}_{s_{n}}\left(u_{s_{n}}\right) \geq \frac{t^{2}}{2}+o(1)
$$

which is a contradiction with the boundedness of $\left\{\mathcal{J}_{s_{n}}\left(u_{s_{n}}\right)\right\}_{n}$. Hence (5.1) does not hold, i.e. there is a sequence $\left\{z_{n}\right\} \subset \mathbb{Z}^{N}$, such that

$$
\liminf _{n \rightarrow \infty} \int_{B\left(z_{n}, 1+\sqrt{N}\right)}\left|v_{n}\right|^{2} \mathrm{~d} x>0
$$

or, equivalently

$$
\liminf _{n \rightarrow \infty} \int_{B(0,1+\sqrt{N})}\left|v_{n}\left(x-z_{n}\right)\right|^{2} \mathrm{~d} x>0
$$

From Theorem 3.2, $v_{n}\left(\cdot-z_{n}\right) \rightarrow v_{0}$ in $L_{\text {loc }}^{2}\left(\mathbb{R}^{N}\right)$ and pointwise a.e., moreover $v_{0} \neq 0$. See that, for a.e. $x \in \operatorname{supp} v_{0}$ we have

$$
\left|u_{s_{n}}\left(x-z_{n}\right)\right|=\left\|u_{s_{n}}\right\|_{s_{n}}\left|v_{s_{n}}\left(x-z_{n}\right)\right| \rightarrow+\infty
$$

Thus

$$
\begin{aligned}
o(1) & =\frac{\mathcal{J}_{s_{n}}\left(u_{s_{n}}\right)}{\left\|u_{s_{n}}\right\|_{s_{n}}^{2}}=\frac{1}{2}-\int_{\mathbb{R}^{N}} \frac{F\left(x, u_{s_{n}}\right)}{u_{s_{n}}^{2}} v_{s_{n}}^{2} \mathrm{~d} x \\
& =\frac{1}{2}-\int_{\mathbb{R}^{N}} \frac{F\left(x, u_{s_{n}}\left(x-z_{n}\right)\right)}{u_{s_{n}}\left(x-z_{n}\right)^{2}} v_{s_{n}}\left(x-z_{n}\right)^{2} \mathrm{~d} x \\
& \leq \frac{1}{2}-\int_{\operatorname{supp} v_{0}} \frac{F\left(x, u_{s_{n}}\left(x-z_{n}\right)\right)}{u_{s_{n}}\left(x-z_{n}\right)^{2}} v_{s_{n}}\left(x-z_{n}\right)^{2} \mathrm{~d} x \rightarrow-\infty
\end{aligned}
$$

a contradiction.
Lemma 5.3. Since $u_{s} \in \mathcal{N}_{s}$ there is (independent of s) constant $\rho$, such that

$$
\left\|u_{s}\right\|_{s} \geq \rho>0
$$

Proof. Since $u_{s} \in \mathcal{N}_{s}$, we can write by Remark 1.1

$$
\begin{aligned}
\left\|u_{s}\right\|_{s}^{2}=\int_{\mathbb{R}^{N}} f\left(x, u_{s}\right) u_{s} \mathrm{~d} x & \leq \varepsilon\left\|u_{s}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}^{2}+C_{\varepsilon}\left\|u_{s}\right\|_{L^{p}\left(\mathbb{R}^{N}\right)}^{p} \\
& \leq C\left(\varepsilon\left\|u_{s}\right\|_{s}^{2}+C_{\varepsilon}\left\|u_{s}\right\|_{s}^{p}\right)
\end{aligned}
$$

for a constant $C>0$ independent of $s$. Choosing $\varepsilon>0$ small enough, we conclude that

$$
\left\|u_{s}\right\|_{s}^{p-2} \geq \frac{1-C \varepsilon}{C \cdot C_{\varepsilon}}=\rho>0
$$

Corollary 5.4. There exist $u_{0} \in H^{1}\left(\mathbb{R}^{N}\right)$, a sequence $\left\{z_{n}\right\}_{n} \subset \mathbb{Z}^{N}$ and a sequence $\left\{s_{n}\right\}_{n}$ such that $s_{n} \rightarrow 1^{-}$and

$$
u_{s_{n}}\left(\cdot-z_{n}\right) \rightarrow u_{0} \neq 0 \quad \text { in } L_{\mathrm{loc}}^{\nu}\left(\mathbb{R}^{N}\right) \quad \text { as } n \rightarrow+\infty
$$

for all $\nu \in[2,2 N /(N-1))$.
Proof. From Lemma 5.2 and Theorem 3.2 we note that

$$
u_{s_{n}} \rightarrow u_{0} \quad \text { in } L_{\mathrm{loc}}^{\nu}\left(\mathbb{R}^{N}\right) \quad \text { as } n \rightarrow+\infty
$$

for all $\nu \in[2,2 N /(N-1))$. If $u_{0} \neq 0$, we can take $z_{n}=0$ and the proof is completed. Otherwise $u_{s_{n}} \rightarrow 0$ in $L_{\mathrm{loc}}^{2}\left(\mathbb{R}^{N}\right)$ and therefore, $u_{s_{n}}(x) \rightarrow 0$ for a.e. $x \in \mathbb{R}^{N}$. Assume that

$$
\sup _{y \in \mathbb{R}^{N}} \int_{B(y, 1)}\left|u_{s_{n}}\right|^{2} \mathrm{~d} x \rightarrow 0
$$

Then from Theorem 3.1 we know that $u_{s_{n}} \rightarrow 0$ in $L^{\nu}\left(\mathbb{R}^{N}\right)$ for all $\nu \in$ $[2,2 N /(N-1))$. Then

$$
\int_{\mathbb{R}^{N}} f\left(x, u_{s_{n}}\right) u_{s_{n}} \mathrm{~d} x \rightarrow 0
$$

and $\left\|u_{s_{n}}\right\|_{s_{n}}^{2}=\int_{\mathbb{R}^{N}} f\left(x, u_{s_{n}}\right) u_{s_{n}} \mathrm{~d} x \rightarrow 0$, which is a contradiction with Lemma 5.3. Hence there is a sequence $\left\{z_{n}\right\} \subset \mathbb{Z}^{N}$ such that

$$
\begin{equation*}
\liminf _{n \rightarrow+\infty} \int_{B(0,1+\sqrt{N})}\left|u_{s_{n}}\left(\cdot-z_{n}\right)\right|^{2} \mathrm{~d} x>0 \tag{5.2}
\end{equation*}
$$

Moreover $\left\|u_{s_{n}}\left(\cdot-z_{n}\right)\right\|_{s_{n}}=\left\|u_{s_{n}}\right\|_{s_{n}}$, so that $\left\|u_{s_{n}}\left(\cdot-z_{n}\right)\right\|_{s_{n}}$ is bounded (see Lemma 5.2). Hence, in view of Theorem 3.2

$$
u_{s_{n}\left(--z_{n}\right)} \rightarrow \tilde{u}_{0} \quad \text { in } L_{\mathrm{loc}}^{\nu}\left(\mathbb{R}^{N}\right) \quad \text { as } n \rightarrow+\infty
$$

for some $\tilde{u}_{0}$. Moreover, in view of (5.2), $\tilde{u}_{0} \neq 0$.
Lemma 5.5. The limit $u_{0} \in H^{1}\left(\mathbb{R}^{N}\right) \backslash\{0\}$ is a weak solution for (1.3).
Proof. Take any test function $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ and note that by [20, Section 6] we have

$$
\int_{\mathbb{R}^{N}}(-\Delta)^{s / 2} u_{s_{n}}(-\Delta)^{s_{n} / 2} \varphi \mathrm{~d} x=\int_{\mathbb{R}^{N}} u_{s_{n}}(-\Delta)^{s_{n}} \varphi \mathrm{~d} x
$$

Moreover

$$
\begin{aligned}
& \left|\int_{\mathbb{R}^{N}} u_{s_{n}}(-\Delta)^{s_{n}} \varphi \mathrm{~d} x-\int_{\mathbb{R}^{N}} u_{0}(-\Delta \varphi) \mathrm{d} x\right| \\
& \quad=\left|\int_{\mathbb{R}^{N}} u_{s_{n}}\left((-\Delta)^{s_{n}} \varphi-(-\Delta \varphi)\right) \mathrm{d} x+\int_{\operatorname{supp} \varphi}\left(u_{s_{n}}-u_{0}\right)(-\Delta \varphi) \mathrm{d} x\right| \\
& \quad \leq\left\|u_{s_{n}}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}\left\|(-\Delta)^{s_{n}} \varphi-(-\Delta \varphi)\right\|_{L^{2}\left(\mathbb{R}^{N}\right)} \\
& \quad+\|(-\Delta \varphi)\|_{L^{2}\left(\mathbb{R}^{N}\right)}\left\|u_{s_{n}}-u_{0}\right\|_{L^{2}(\operatorname{supp} \varphi)} \rightarrow 0 .
\end{aligned}
$$

Hence

$$
\begin{aligned}
\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}(-\Delta)^{s_{n} / 2} u_{s_{n}}(-\Delta)^{s_{n} / 2} \varphi \mathrm{~d} x & =\int_{\mathbb{R}^{N}} u_{0}(-\Delta \varphi) \mathrm{d} x \\
& =\int_{\mathbb{R}^{N}} \nabla u_{0} \cdot \nabla \varphi \mathrm{~d} x
\end{aligned}
$$

Obviously

$$
\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}} V(x) u_{s_{n}} \varphi \mathrm{~d} x=\lim _{n \rightarrow+\infty} \int_{\operatorname{supp} \varphi} V(x) u_{s_{n}} \varphi \mathrm{~d} x=\int_{\mathbb{R}^{N}} V(x) u_{0} \varphi \mathrm{~d} x
$$

Take any measurable set $E \subset \operatorname{supp} \varphi$ and note that, taking into account Remark 1.1,

$$
\begin{aligned}
& \int_{E}\left|f\left(x, u_{s_{n}}\right) \varphi\right| \mathrm{d} x \leq \varepsilon\left\|u_{s_{n}}\right\|_{L^{2}\left(\mathbb{R}^{N}\right)}\left\|\varphi \chi_{E}\right\|_{L^{2}(\operatorname{supp} \varphi)} \\
& \quad+C_{\varepsilon}\left\|u_{s_{n}}\right\|_{L^{p}\left(\mathbb{R}^{N}\right)}^{p-1}\left\|\varphi \chi_{E}\right\|_{L^{p}(\operatorname{supp} \varphi)} .
\end{aligned}
$$

Hence, the family $\left\{f\left(\cdot, u_{s_{n}}\right) \varphi\right\}_{n}$ is uniformly integrable on $\operatorname{supp} \varphi$ and in view of the Vitali convergence theorem

$$
\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}} f\left(x, u_{s_{n}}\right) \varphi \mathrm{d} x=\int_{\mathbb{R}^{N}} f\left(x, u_{0}\right) \varphi \mathrm{d} x
$$

Therefore $u_{0}$ satisfies

$$
\int_{\mathbb{R}^{N}} \nabla u_{0} \cdot \nabla \varphi \mathrm{~d} x+\int_{\mathbb{R}^{N}} V(x) u_{0} \varphi \mathrm{~d} x=\int_{\mathbb{R}^{N}} f\left(x, u_{0}\right) \varphi \mathrm{d} x
$$

i.e. $u_{0}$ is a weak solution to (1.3).

Proof of Theorem 1.2. Recalling Corollary 5.4 and Lemma 5.5, it is sufficient to check that $u_{0}$ is a ground state solution, i.e. $\mathcal{J}\left(u_{0}\right)=c$. From Lemma 5.5 it follows that $u_{0} \in H^{1}\left(\mathbb{R}^{N}\right) \backslash\{0\}$ is a weak solution, so that $u_{0} \in \mathcal{N}$. Note that, from Corollary 5.4 and Fatou's lemma,

$$
\begin{aligned}
\liminf _{n \rightarrow+\infty} c_{s_{n}} & =\liminf _{n \rightarrow+\infty} \mathcal{J}_{s_{n}}\left(u_{s_{n}}\right)=\liminf _{n \rightarrow+\infty}\left\{\mathcal{J}_{s_{n}}\left(u_{s_{n}}\right)-\frac{1}{2} \mathcal{J}_{s_{n}}^{\prime}\left(u_{s_{n}}\right)\left(u_{s_{n}}\right)\right\} \\
& =\liminf _{n \rightarrow+\infty}\left\{\frac{1}{2} \int_{\mathbb{R}^{N}} f\left(x, u_{s_{n}}\right) u_{s_{n}}-2 F\left(x, u_{s_{n}}\right) \mathrm{d} x\right\} \\
& =\liminf _{n \rightarrow+\infty}\left\{\frac{1}{2} \int_{\mathbb{R}^{N}} f\left(x, u_{s_{n}}\left(\cdot-z_{n}\right)\right) u_{s_{n}}\left(\cdot-z_{n}\right)-2 F\left(x, u_{s_{n}}\left(\cdot-z_{n}\right)\right) \mathrm{d} x\right\} \\
& \geq \frac{1}{2} \int_{\mathbb{R}^{N}} f\left(x, u_{0}\right) u_{0}-2 F\left(x, u_{0}\right) \mathrm{d} x=\mathcal{J}\left(u_{0}\right) \geq c .
\end{aligned}
$$

Taking into account Lemma 5.1 we see that

$$
c \leq \mathcal{J}\left(u_{0}\right) \leq \liminf _{n \rightarrow+\infty} c_{s_{n}} \leq \limsup _{n \rightarrow+\infty} c_{s_{n}} \leq c
$$

Hence $\lim _{n \rightarrow+\infty} c_{s_{n}}$ exists and $\lim _{n \rightarrow+\infty} c_{s_{n}}=c=\mathcal{J}\left(u_{0}\right)$.

## Acknowledgements

The authors would like to thank an anonymous referee for several valuable comments helping to improve the original version of the manuscript. Bartosz Bieganowski was partially supported by the National Science Centre,

Poland (Grant No. 2017/25/N/ST1/00531). Simone Secchi is member of the Gruppo Nazionale per l'Analisi Matematica, la Probabilità e le loro Applicazioni (GNAMPA) of the Istituto Nazionale di Alta Matematica (INdAM).

Funding Open access funding provided by Università degli Studi di Milano Bicocca within the CRUI-CARE Agreement.

Open Access. This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http:// creativecommons.org/licenses/by/4.0/.

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

## References

[1] Bakunin, O.G.: Turbulence and Diffusion: Scaling Versus Equations. Springer, Berlin (2008)
[2] Biccari, U., Hernández-Santamaría, V.: The Poisson equation from non-local to local. Electron. J. Differ. Equ. 2018(145), 1-13 (2018)
[3] Bieganowski, B.: Solutions of the fractional Schrödinger equation with a signchanging nonlinearity. J. Math. Anal. Appl. 450(1), 461-479 (2017)
[4] Bieganowski, B., Mederski, J.: Nonlinear Schrödinger equations with sum of periodic and vanishing potentials and sign-changing nonlinearities. Commun. Pure Appl. Anal. 17(1), 143-161 (2018)
[5] Bieganowski, B., Secchi, S.: Non-local to local transition for ground states of fractional Schrödinger equations on bounded domains. Topol. Methods Nonlinear Anal. (to appear)
[6] Borthagaray, J.P., Ciarlet Jr., P.: On the convergence in the $H^{1}$-norm for the fractional Laplacian. arXiv:1810.07645
[7] Bourgain, J., Brezis, H., Mironescu, P.: Another look at Sobolev spaces, Optimal Control and Partial Differential Equations, pp. 439-455. IOS, Amsterdam (2001)
[8] Brezis, H.: Functional Analysis. Sobolev Spaces and Partial Differential Equations. Springer, Berlin (2011)
[9] Cotsiolis, A., Tavoularis, N.: Best constants for Sobolev inequalities for higher order fractional derivatives. J. Math. Anal. Appl. 295, 225-236 (2004)
[10] Di Nezza, E., Palatucci, G., Valdinoci, E.: Hitchhiker's guide to the fractional Sobolev spaces. Bull. Sci. Math. 136, 521-573 (2012)
[11] Dipierro, S., Palatucci, G., Valdinoci, E.: Dislocation dynamics in crystals: a macroscopic theory in a fractional Laplace setting. Commun. Math. Phys. 333(2), 1061-1105 (2015)
[12] de Paiva, F.O., Kryszewski, W., Szulkin, A.: Generalized Nehari manifold and semilinear Schrödinger equation with weak monotonicity condition on the nonlinear term. Proc. Am. Math. Soc. 145(11), 4783-4794 (2017)
[13] Gilboa, G., Osher, S.: Nonlocal operators with applications to image processing. Multiscale Model. Simul. 7(3), 1005-1028 (2008)
[14] Molica Bisci, G., Radulescu, V., Servadei, R.: Variational Methods for Nonlocal Fractional Problems. Encyclopedia of Mathematics and Its Applications, vol. 162. Cambridge University Press, Cambridge (2016)
[15] Secchi, S.: Ground state solutions for nonlinear fractional Schrödinger equations in $\mathbb{R}^{N}$. J. Math. Phys. 54, 031501 (2013)
[16] Secchi, S.: On fractional Schrödinger equations in $\mathbb{R}^{N}$ without the AmbrosettiRabinowitz condition. Topol. Methods Nonlinear Anal. 47, 19-41 (2016)
[17] Szulkin, A., Weth, T.: Ground state solutions for some indefinite variational problems. J. Funct. Anal. 257(12), 3802-3822 (2009)
[18] Szulkin, A., Weth, T.: The method of Nehari manifold. In: Gao, D.Y., Motreanu, D. (eds.) Handbook of Nonconvex Analysis and Applications. International Press, Boston, pp. 597-632 (2010)
[19] Vázquez, J.L.: Nonlinear diffusion with fractional Laplacian operators. In: Nonlinear Partial Differential Equations (Oslo 2010), Abel Symp., vol. 7. Springer, Heidelberg, pp. 271-298 (2012)
[20] Warma, M.: The fractional relative capacity and the fractional Laplacian with Neumann and Robin boundary conditions on open sets. Potential Anal. 42(2), 499-547 (2015)

Bartosz Bieganowski
Faculty of Mathematics and Computer Science
Nicolaus Copernicus University
ul. Chopina 12/18
87-100 Toruń
Poland
e-mail: bartoszb@mat.umk.pl
Simone Secchi
Dipartimento di Matematica e Applicazioni
Università degli Studi di Milano-Bicocca
via Roberto Cozzi 55
20125 Milan
Italy
e-mail: simone.secchi@unimib.it

