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Tunable Color Emissions upon UV Irradiation from Tb³⁺:Y₂SiO₅ Phosphor

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Abstract

A series of terbium-activated yttrium oxyorthosilicate $(Tb^{3+}:Y_2SiO_5)$ phosphor was prepared through the sol-gel technique at a high calcination temperature of 1100°C. The phase purity was examined by the x-ray diffraction (XRD) analysis and found to be X1 type for all series of samples. The existence of functional groups in the prepared samples was confirmed by Fourier-transform infrared (FT-IR) spectroscopy. The sample exhibited absorption bands of Si-O-Si and Y-O groups at wavenumbers between 717 cm⁻¹ and 550 cm⁻¹, and the NO₃ groups appeared between 800 cm⁻¹ and 1020 cm⁻¹ in the FT-IR spectra. Under 244-nm excitation, the photoluminescence spectra were recorded in the 350 to 650 nm wavelength range. A possible cross-relaxation mechanism between the ⁵D₃ and the ⁵D₄ levels for the concentration quenching of two distinct emission transitions has been described in detail. The dipole–dipole interaction is most responsible for concentration quenching, as evidenced by the Dexter theory. The CIE chromaticity analysis of the Y₂SiO₅:Tb³⁺ phosphor was illustrated by a CIE 1931 color calculator. These findings demonstrate the suitability of the Tb³⁺-activated Y₂SiO₅ matrix as a UV-excitable, color-tunable emitter with Tb³⁺ ion variations and an efficient green light-emitting phosphor in the fields of solid-state lighting devices.

Keywords $Y_2SiO_5 \cdot photoluminescence \cdot Tb^{3+} ions \cdot sol-gel$

Introduction

Luminescent phosphors have been extensively studied for their use as color-emitting components in the lighting industry, such as displays, plasma display panels, solid-state lighting, electroluminescent panels, light emitting diodes (LEDs), and scintillator panels for x-ray radiography.^{1–4} To meet technological requirements, the materials should be characterized by long periods of work with high transparency and brightness. Therefore, within this framework, finding novel functional materials is a challenging task for

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researchers who have turned their attention to inorganic luminescent materials. Several crystal hosts doped with rare-earth elements, such as aluminate garnet ($Y_3Al_5O_{12}$), alkaline orthophosphates, and oxyorthosilicates (such as Re_2SiO_5) have shown good luminescence properties.^{5–7} In this paper, the authors focus on Tb^{3+} -doped yettrium oxyorthosilicate (Y_2SiO_5) host matrix. Generally, the Y_2SiO_2 host is known to be an excellent cathodoluminescent phosphor due to its unique physicochemical characteristics, such as efficient chemical and thermal stability with the large stopping potential of the absorption coefficient.^{8,9} The persistence of crystalline Y_2SiO_5 monoclinic phases of X1 or X2 type is visible in the preparation of materials depending on the temperature, and, usually, at a higher temperature (\geq 1300° C), the X2 crystalline phase can be observed.^{10,11}

The development of Tb^{3+} -doped phosphor stands as a cornerstone in the realm of luminescent materials, heralding advances in various technological domains. These phosphors, comprised of terbium ions integrated into host matrices, exhibit unique photoluminescent properties, making them indispensable in diverse applications.³ The journey of these phosphors traces back to their inception, when researchers endeavored to harness the distinctive emission characteristics of Tb³⁺ ions to create efficient light-emitting materials.⁹ The inception of Tb³⁺ -doped phosphors stemmed from the quest to harness the unique luminescence properties of terbium ions within a host matrix. Researchers aimed to exploit the characteristic emission spectrum of Tb³⁺ ions, which emits intense and well-defined green light upon excitation by external energy sources.^{3,9,11} This distinctive luminescence property, coupled with the ability to manipulate its emission wavelengths and intensity, sparked immense interest in the scientific community. The early stages of development witnessed pioneering efforts in synthesizing and characterizing these phosphors, laying the groundwork for subsequent breakthroughs. Fundamental research into crystal structures, doping concentrations, and fabrication methodologies paved the way for enhanced understanding and control over the luminescence behavior of Tb^{3+} -doped phosphors. As technological demands evolved, so did the pursuit of optimizing these materials for a myriad of applications. The continual refinement of synthesis techniques, from traditional powder-based formulations to cutting-edge nanostructured designs, has been instrumental in achieving superior efficiency, stability, and tailored functionalities. The ongoing development and advancement of Tb³⁺ -doped phosphors underscore the relentless pursuit of efficiency, reliability, and versatility in luminescent materials.¹² This journey of innovation not only reshapes the landscape of technological applications but also continues to inspire further exploration, promising new frontiers in the realm of luminescence and materials science.

Nevertheless, trivalent terbium (Tb³⁺) ions have emissions in the green region with numerous host materials as described in the local crystal environment.^{12,13} Properly excited, they typically exhibit emissions at ~ 489 nm, ~ 540 nm, ~ 580 nm, and ~ 620 nm, which can be attributed to ${}^{5}D_{4} \rightarrow {}^{7}F_{6}$, ${}^{7}F_{5}$, ${}^{7}F_{4}$, and ${}^{7}F_{3}$ transitions, respectively.¹⁴ While the narrow emission band was attributed to the spin-forbidden 4f-4f transitions, the intense emission of the ${}^{5}D_{4} \rightarrow {}^{7}F_{5}$ electronic transition is the prime contributor to the phosphor's green emission. The high emission band that leads to spin-allowed transitions is unsuitable for LEDs, hindering the application of most Tb³⁺-based host phosphors for w-LEDs development despite its optimal green luminescence.^{15–17} This has motivated thorough research into the development of Tb³⁺ions as luminescent phosphors for green light emission for application in various fields, including mercury-excited lamps and display devices. Orthosilicate-based luminescent phosphor materials are generally environmentally friendly as they do not contain toxic materials like sulfur.¹⁸ Various inorganic luminescent materials composed of rare-earth-activated oxyorthosilicates have been widely described, including Y₂SiO₅:Tb³⁺, Eu³⁺,¹⁹ $Y_2SiO_5:Tb^{3+}$, ²⁰ X_2 - $Y_2SiO_5:Tb^{3+}$, ²¹ Tb-doped Y_2SiO_5 , ²² and Y_2SiO_5 :Ce.²³ Oxide phosphors have been synthesized using various methods, such as solid-state reaction,²⁴ spray pyrolysis,²⁵ combustion synthesis,²⁶ solid–liquid reaction,²⁷ and co-precipitation,²¹ respectively.

It has recently been noticed that Tb³⁺ -activated phosphors exhibit tunable blue-to-green emissions due to luminescence quenching, which are associated with the crossrelaxation mechanism between the two excited levels, ⁵D₃ and ⁵D₄ of Tb³⁺ ions, occurring in the host crystal lattices and activator concentration.^{28–30} Kang et al.,³¹ Li et al.,³² and Zang et al.³³ investigated and then concluded that the crystal structure of X2 presents better photoluminescence characteristics than the X1-type in Tb³⁺:Y₂SiO₅ phosphor. Afandi et al.³⁴ prepared an X2-type crystal structure of Tb³⁺:Y₂SiO₅ phosphors at high temperatures through solid-state reaction. They observed tunable emissions with a variation of Tb³⁺ ions, typically blue in color with 0.1% of Tb³⁺ ions and green in color with 1% of Tb³⁺ ions, and concluded that phosphor particles with an X2-type crystal structure in Tb³⁺:Y₂SiO₅ phosphors are useful for the development of display technology.

In this paper, the structural and luminescence properties of Tb³⁺ -activated Y₂SiO₅ phosphor samples were studied in order to observe tunable emissions with the doping of Tb³⁺ ions concentrations over a wide range from 1 mol% to 11 mol%. Structural and functional group characteristics were investigated by x-ray diffraction (XRD) and Fouriertransform infrared (FT-IR) spectroscopy. The photoluminescence properties and CIE (Commission International de L'Eclairage) coordinates of the prepared Tb³⁺:Y₂SiO₅ phosphor have been examined in detail. The novelty of the obtained results lies in the ability to achieve tunable emissions of blue-green by altering the Tb³⁺ ion concentration in the X1-type crystallite structure of Tb³⁺:Y₂SiO₅ phosphors which is now limited to X2-type only. To summarize our investigation, the Tb³⁺:Y₂SiO₅ exhibited outstanding photoluminescence characteristics that support its promising potential as a blue-green color-emitting phosphor material for displays.

Experimental

A series of $Y_2SiO_5:xTb^{3+}$ (x = 0.01 mol, 0.03 mol, 0.05 mol, 0.07 mol, 0.09 mol, 0.11 mol) luminescent samples were prepared using the sol-gel method. Stoichiometric amounts of the starting materials $Y(NO_3)_3 \cdot 6H_2O$, SiO_2 , $Tb(NO_3)_3 \cdot xH_2O$, and citric acid (the molar ratio of metal ion to citric acid was 1:2) were used for the synthesis process. The compounds were accurately weighed and then added to a 0.05-L glass beaker containing 0.012 L of deionized water. The resulting solution was thoroughly mixed with a magnetic stirrer for over 1 h at room

temperature. The homogeneous mixture was initially inserted into an oven to obtain the dry gel. The dried gels were heated at 400°C for 20 min, then crushed, and finally sintered at 1100°C for 3 h. The entire sintering procedure of the prepared sample was performed in the air. The block diagram of the synthesis process is depicted in Fig. 1. The samples' phase identification was conducted by analyzing their XRD patterns using a RIGAKU (Miniflex-II) diffractometer in the 2θ range of 12–60 degrees. Diffraction patterns were examined at a scan rate of 0.08/s with a scanning angular range of 10°–60°. A Thermo Fisher Nicolet (6700) FT-IR spectrometer was used to detect the various structural groups present in the samples. Photoluminescence



Fig. 1 Flow chart of the synthesis method.

characterization was performed with RF-5301PC spectro-fluorophotometer (Shimadzu, Japan).

Results and Discussion

According to previous works on Tb³⁺:Y₂SiO₅ phosphor matrices, the crystal lattice distortion or phase transition can occur at the synthesis temperature, although Y^{3+} and Tb^{3+} ionic radii are almost equal. It is known that phase purity is determined by evaluating XRD data. Figure 2 shows the XRD pattern of the prepared Tb³⁺:Y₂SiO₅ phosphor samples. All the diffraction peaks were observed to be indexed with the low-temperature monoclinic phase Y₂SiO₅ (41-0004, JCPDS file) with a space group of P21/c and the high-temperature monoclinic phase Y₂SiO₅ (36-1476, JCPDS file) with a space symmetry of I^*/a .^{35,36} It is interesting to note that increasing the concentration of Tb³⁺ in the Y₂SiO₅ host matrix does not induce any distortions in the space group and that the major planes correspond to the lowtemperature phase of Y₂SiO₅. Therefore, we achieved a stable and homogeneous crystal plane and a uniform increase in the concentration of Tb^{3+} ions (1–11 mol%) during the synthesis of the samples at 1100 °C by the sol-gel method. These results suggest that the observed crystal phase is of the X1 type for all the series of samples, and thus the sol-gel synthesis offers excellent intercalation of elements in the starting materials and high cation homogeneity, which lowers the diffusion resistance. The low-temperature phase of Y_2SiO_5 plays a crucial role in influencing the structural and optical properties of the synthesized Tb^{3+} :Y₂SiO₅ phosphor. The choice of a lower synthesis temperature, in comparison to higher-temperature phases, is intentional and designed



Fig. 2 Powder XRD patterns of $Y_2SiO_5:xTb^{3+}(0.01 \le x \le 0.11)$ phosphors.

to favor the formation of a specific crystal structure. Lowtemperature phases often exhibit distinct crystallographic arrangements, and, in our case, it contributes to achieving the desired structural characteristics in Y_2SiO_5 .

Figure 3 depicts the crystal structure of the Y_2SiO_5 host material by using the VESTA software. The CIF file number -1001838 was obtained directly from open crystallographic data.^{37,38} The crystal structure parameters of the yttrium oxyorthosilicate (Y₂SiO₅) phosphor are (α , $\beta = 90$, $\gamma = 122.25$), (a = 14.58320 Å, b = 10.5120 Å, c = 6.8152 Å) and the cell volume, V = 884.2910 Å³. The resulting X1-type crystal phase structure of the Tb³⁺:Y₂SiO₅ phosphor was monoclinic with P21/c and the local symmetry consisted of two distinct Y^{3+} sites: (1) 9-coordinated with 8 Si-O bonded atoms and 1 without silicon-bonded atom (free oxygen (O) atoms); and (2) 7-coordinated with 4 Si-O bonded atoms and 3 free O atoms. Due to the equal ionic radii of Tb³⁺ and Y³⁺, Tb³⁺ ions can be expected to be situated at local Y^{3+} sites. The percentage difference in radius between the active ion (Tb^{3+}) and the host cation (Y^{3+}) can be determined by³⁹:

$$\delta_{\rm r} = \frac{R_{Tb}(\rm CN) - R_Y(\rm CN)}{R_{Tb}(\rm CN)} \times 100 \tag{1}$$

where δ_r is the percentage difference in the radii, CN is the coordination number, and R_{Tb} and R_Y are the radii of the active ion and the host cation, respectively. R_{Tb} and R_Y with coordination number 6 are 0.92 Å and 0.9 Å. The estimated δ_r is ~ 2.2, which is an excellent acceptable range for effective ionic substitution ($\leq 3\%$).⁴⁰

Figure 4 shows the FI-IR spectra of Tb³⁺:Y₂SiO₅ phosphors. The band between 4000 cm⁻¹ to 3500 cm⁻¹ is related to the stretching mode of O-H molecules. The band at 1528 cm⁻¹ is attributed to the C=C bond stretching those results from the alkoxy groups. The Y(NO₃)₃ absorption bands exhibited vibration peaks of the NO₃ group at $\sim 1029 \text{ cm}^{-1}$, $\sim 933 \text{ cm}^{-1}$, and $\sim 885 \text{ cm}^{-1}$, which are associated with the $m(NO_2)$, m(NO) and $m(NO_3)$ bands, respectively.^{41,42} This confirms the existence of these ions in the gels. The peak at 698 cm⁻¹ corresponds to the stretching vibrations of the Si-O-Si bonds within the Y₂SiO₅ crystal lattice. This frequency is typically associated with the symmetric stretching mode of the Si-O-Si unit in the structure.⁴³ The peak at approximately 2354 cm⁻¹ could correspond to a variety of possibilities depending on the specific environment of the material. It could potentially indicate the presence of C-H stretching vibrations if there are any organic contaminants or residual carbon-based compounds present. The sample exhibited a Y-O bond absorption band appearing at ~ 555 cm⁻¹. The FT-IR spectrum of Y_2SiO_5 shows a weaker peak at 492 cm⁻¹, which corresponds to the Si-O bending modes.



Fig. 3 Crystal structure of Y₂SiO₅ (a) standard orientation, (b, c, and d) three different orientations, 100, 010, and 001, respectively.



Fig. 4 FT-IR spectrum of $Y_2SiO_5:0.11Tb^{3+}$ phosphor.

Figure 5 shows the photoluminescence excitation and emission spectra of Tb^{3+} :Y₂SiO₅ phosphors. Photoluminescence excitation spectra are recorded in the 220–400 nm wavelength range by monitoring the emission at 543 nm (Tb³⁺: ⁵D₄ \rightarrow ⁷F₅), shown in Fig. 5a. Two significant wide bands were discovered in the ultraviolet spectral region. The

bands at 244 nm and 302 nm would be attributed to the mixing of a charge-transfer band due to Tb-O interaction and the 4f-5d transition bands in the Tb³⁺ center. The several other small peaks in the range of 300-400 nm represent the f-f transition of Tb^{3+.44} It can be seen that the excitation band with an intensity of 302 nm changes with the concentration's variation, whereas a clear increase in the intensity of the 244-nm peak was observed. In this context, we assume that the Tb^{3+} ions are in an octahedral environment because they excited Tb³⁺ levels of the 4f⁷5d¹ configuration in the Tb³⁺ centers.⁴⁵ The yttrium sites in the Y₂SiO₅ host lattice are 7 and 9 coordinated with oxides (P21/c symmetry group), which restricts the crystal field splitting of the current system to octahedral crystal field-splitting approximations.^{46,47} However, the crystal field splitting helps to understand the excitation of Tb³⁺:Y₂SiO₅. Given the spin selection rules, a weak intense peak has a weak oscillator strength originating from a spin-forbidden, and a strong intense peak has a strong oscillator strength coming from a spin-allowed. The excitation peak at 244 nm was most appropriate for the excellent emissions from the prepared phosphor.

The photoluminescence spectra of the Tb^{3+} : Y_2SiO_5 phosphors were measured in the wavelength range of 350–650 nm upon a 244-nm exciter and are depicted in Fig. 5b. The observed emission peaks are associated with the transitions originating from the ⁵D₃ and ⁵D₄ levels to the lower ⁷F_J levels. The electronic transitions ⁵D₃ \rightarrow ⁷F₆,



Fig. 5 (a) Excitation ($\lambda_{em} = 543 \text{ nm}$) and (b) photoluminescence spectra ($\lambda_{exc} = 244 \text{ nm}$) of Y₂SiO₅:*x*Tb³⁺ (0.01 $\leq x \leq 0.11$) phosphors.

 ${}^{7}F_{5}$, ${}^{7}F_{4}$ and ${}^{7}F_{3}$ correspond to the emission peak centers at ~ 383 nm, ~ 417 nm, ~ 438 nm and ~ 460 nm, respectively, and ${}^{5}D_{4} \rightarrow {}^{7}F_{6}$, ${}^{7}F_{5}$, ${}^{7}F_{4}$, and ${}^{7}F_{3}$ transitions are corresponding to the peak centers at ~ 490 nm, ~ 543 nm, ~ 589 nm and ~ 624 nm, respectively. These two distinct ${}^{5}D_{I}$ levels originating transitions are created by the splitting effect of the crystal field. Moreover, the ${}^{5}D_{4} \rightarrow {}^{7}F_{6}$ transition in blue light emission was attributed to the electric dipole transition described by the local environment of the crystal field. The band at ~ 543 nm is the symmetry of the crystal fielddependent one, whereas the ${}^{5}D_{4} \rightarrow {}^{7}F_{5}$ located in the green light emission region corresponds to the magnetic dipole transition and is free from the crystal field strength.⁴⁴ The decrease in the emission intensity of transitions occurring from higher energy states in Tb-doped Y₂SiO₅ can be influenced by several factors. As the energy levels increase, there is a higher probability of non-radiative decay pathways becoming more dominant. Non-radiative processes, such as phonon-assisted relaxation or energy transfer to defects or impurities within the crystal lattice, can reduce the efficiency of radiative emissions from higher energy states. These nonradiative pathways compete with the radiative transitions, leading to a decrease in the observed emission intensity. In complex materials like Y₂SiO₅, the crystal field around the Tb ions can lead to Stark splitting, causing the energy levels to separate. Transitions originating from higher Stark levels might have lower probabilities of radiative emission due to their altered energy landscapes, contributing to reduced emission intensity. Figure 6 shows an energy-level schematic of the excitation and emission mechanism of Tb³⁺:Y₂SiO₅ phosphors (cross-relaxation mechanisms, non-radiative transition).

Figure 7 shows the concentration-quenching behavior of the strong emission transitions associated with both the ${}^{5}D_{3}$ and ${}^{5}D_{4}$ levels to the ${}^{7}F_{J(J=4,5\&6)}$ levels as the Tb³⁺ concentration increases. It can be seen that the maximum emission intensity of the ${}^{5}D_{3} \rightarrow {}^{7}F_{5}$, ${}^{7}F_{4}$, and ${}^{5}D_{4} \rightarrow {}^{7}F_{6}$, ${}^{7}F_{5}$ transitions was obtained for 1 mol% and 9 mol% of Tb³⁺ ions in Tb³⁺:Y₂SiO₅ phosphors. The observed concentration quenching is caused by a non-radiative transition from ${}^{5}D_{3}$ to $^{5}D_{4}$ levels through a cross-relaxation mechanism. The concentration quenching observed ranged from approximately 1 mol% to 9 mol% of Tb³⁺ ions for the ${}^{5}D_{3}$ and ${}^{5}D_{4}$ excited levels due to the energy transfer process, multi-ion interaction, and localized crystal field effect. In the ${}^{5}D_{3}$ to ${}^{5}D_{4}$ cross-relaxation mechanism, in which the energy transfer is possible from Tb³⁺-Tb³⁺ ions in single-doped systems, the critical distance between donor and acceptor ions can be estimated using Bassle's equation⁴⁸:

$$R_c = 2 \left[\frac{3V}{4\pi N X_c} \right]^{1/3} \tag{2}$$

where *V* is the cell volume (884.2910 Å³), *N* is the availability of four Y³⁺ in the unit cell, and X_c is the critical concentration of Tb³⁺ ions. The R_c for 1 mol% and 9 mol% of Tb³⁺ was found to be about 38.54 Å and 16.39 Å, respectively, from the concentration-quenching data from the ⁵D₃ \rightarrow ⁷F₅ and ⁵D₄ \rightarrow ⁷F₅ emission peaks. The R_c value from the ⁵D₄ \rightarrow ⁷F₅ was more than the ⁵D₃ \rightarrow ⁷F₅ for the ⁵D₃ \rightarrow ⁵D₄ cross-relaxation, and the R_c values are much higher than 5 Å, so there is an exchange interaction responsible for the cross-relaxation between ⁵D₃ and ⁵D₄ levels.

The multipolar interaction mechanism⁴⁹ was adopted to provide a more reasonable interpretation of the concentration



Fig. 6 Partial energy level schematic showing the emission mechanism of Tb^{3+} ions under UV light excitations and possible cross-relaxation channels (i and ii) for concentration quenching in $Y_2SiO_5:Tb^{3+}$ phosphors.



Fig. 7 Variation of emission intensity as a function of Tb^{3+} concentration for the ${}^{5}D_{3} \rightarrow {}^{7}F_{5}$, ${}^{7}F_{4}$, transitions (a) and the ${}^{5}D_{4} \rightarrow {}^{7}F_{6}$, ${}^{7}F_{5}$ transitions (b) in Y₂SiO₅:xTb³⁺ phosphors.

quenching of Tb^{3+} ions in the Y₂SiO₅ phosphor. According to Dexter's theory, the exact interaction type between dipole–dipole ($d \leftrightarrow d$), dipole–quadrupole ($d \leftrightarrow q$), and

quadrupole–quadrupole ($q \leftrightarrow q$) interactions for the occurrence of the energy transfer between Tb³⁺ ions is obtained by^{50,51}:

$$\frac{I}{x} = K \left\{ 1 + \beta(x)^{\frac{Q}{3}} \right\}^{-1}$$
(3)

where I is the intensity of the emission peak, x is the Tb concentration (1 and 9 mol%), and K and β are constants. Q is also a constant for the $d \leftrightarrow d(6), d \leftrightarrow q(8)$, and $q \leftrightarrow q(10)$ interactions. The estimated Q values in this work are 5.64 and 4.69 for the ${}^{5}D_{3} \rightarrow {}^{7}F_{5}$ and ${}^{5}D_{4} \rightarrow {}^{7}F_{5}$ emission bands, respectively, from the slope value of the linear fit to Log(I/x)versus Log(x) (see Fig. 8). The Q value suggests that the specific energy transfer occurs due to dipole-dipole interactions. According to the concentration-quenching luminescence spectral results, two cross-relaxation schemes may be responsible: (1) cross-relaxation from the ${}^{5}D_{3}$ level to the ${}^{5}D_{4}$ level, and (2) cross-relaxation from the ${}^{5}D_{3}$ and/or ${}^{5}D_{4}$ to the ⁷D level (see Fig. 6). Normally, possible quenching may be an energy transfer from the Tb-Tb ion at the ${}^{5}D_{3}$ level to the ${}^{5}D_{4}$ levels. One of the Tb ions is excited to the ${}^{5}D_{3}$ or ${}^{5}D_{4}$ level with selective excitation and decay to the ${}^{7}F_{6}$ ground level, then adjacent Tb ions at ⁵D₃ or ⁵D₄ level are raised to the ⁷D excited state of the 4f⁷5d¹ configuration. Therefore, the luminescence quenching of ${}^{5}D_{3} \rightarrow {}^{7}F_{5}$, ${}^{7}F_{4}$, and ${}^{5}\text{D}_{4} \rightarrow {}^{7}\text{F}_{6}, {}^{7}\text{F}_{5}$ transitions are caused by cross-relaxations (1) and (2), respectively.⁴⁹ Among the numerous radiative and non-radiative processes that occur, the leading increase in intensity is noticed for the green emission of the ${}^{5}D_{4} \rightarrow {}^{7}F_{5}$ transition of Tb^{3+} ions in Y_2SiO_5 phosphors.

The CIE chromaticity plot of $Y_2SiO_5:xTb^{3+}$ phosphors excited at 244 nm is shown in Fig. 9. The CIE coordinates for the Tb³⁺ -activated Y_2SiO_5 phosphors have been determined since 1931 CIE and are listed in Table I. The variation of the Tb³⁺ dopant leads to a shift in the color coordinates of the tunable emission from bluish (x = 0.2223, y = 0.3072) to green (x = 0.2720, y = 0.5470). Therefore, the Tb³⁺:Y₂SiO₅ phosphor is a potential color-adjustable phosphor candidate for UV-based wLEDs. The color purity of the Y₂SiO₅:xTb³⁺ phosphor was evaluated.⁵²



Fig. 9 CIE chromaticity diagram for $Y_2SiO_5:xTb^{3+}$ phosphors.



Fig. 8 Plots of Log(x) versus Log(I/x) for (a) ${}^{5}\text{D}_{3} \rightarrow {}^{7}\text{F}_{5}$ and (b) ${}^{5}\text{D}_{4} \rightarrow {}^{7}\text{F}_{5}$ emission transitions in $\text{Y}_{2}\text{SiO}_{5}$: Tb³⁺ phosphors.

Color purity (%) =
$$\frac{\sqrt{(x-x_i)^2 + (y-y_i)^2}}{\sqrt{(x_d-x_i)^2 + (y_d-y_i)^2}} \times 100$$

The CIE color purity for the 0.09 and 0.11 mol concentrations of Tb^{3+} in Y_2SiO_5 phosphor indicates the suitability of the prepared phosphor for applications in UV excitable solid-state lighting. The color purity of the Y_2SiO_5 phosphor doped with a higher concentration of Tb^{3+} is approximately 60.96%. The observed lower color purity in the emitted light can be influenced by several factors. The presence of impurities, defects, or variations in the host lattice could lead to deviations in the emission spectrum, affecting the color purity. Also, energy transfer processes among luminescent centers can contribute to a broadening of the emission spectrum, leading to a lower color purity. Understanding and controlling these energy transfer mechanisms are crucial for

Table I CIE chromaticity for $Y_2SiO_5:xTb^{3+}$ phosphors $(0.01 \le x \le 0.11)$ using PL emission ($\lambda_{exc} = 244$ nm)

$\overline{Y_2SiO_5}$: xTb^{3+}	$\lambda_{\rm exc} = 244 \text{ nm}$		Color purity%: (sign)
	x	у	
x = 0.01	0.2223	0.3072	64.35 (●)
x = 0.03	0.2493	0.4503	55.81 (●)
x = 0.05	0.2628	0.5076	56.70 (●)
x = 0.07	0.2653	0.5224	58.79 (●)
x = 0.09	0.2706	0.5443	60.69 ()
x = 0.11	0.2720	0.5470	60.96 (•)

optimizing color quality. In summary, these results confirm the development potential of new luminescent devices that are appropriate as color-tunable emitting and efficient single components of green phosphors in the lighting industry. Figure 10 shows photographic images of (1) $Y_2SiO_5:0.01Tb^{3+}$ and (2) $Y_2SiO_5:0.11Tb^{3+}$ green-emitting sol–gel-synthesized luminescent phosphors. These phosphor images show a whitish color in room light and show a prominent greenish color under UV light (excitation wavelength of 254 nm), demonstrating its suitability for as-sold lighting applications.

Conclusions

A Y₂SiO₅:Tb³⁺ phosphor set was efficiently prepared by the sol-gel route. The formation of Y₂SiO₅ was examined by XRD. The functional group was confirmed by FT-IR study. Under 244 nm excitation, the phosphor exhibited emission spectrum lines in the 350-650 nm wavelength range. The emission peaks are responsible for the transitions originating from ${}^{5}D_{3}$ and ${}^{5}D_{4}$ levels to the lower ${}^{7}F_{1}$ levels of Tb³⁺. The quenching concentration for 1 mol% and 9 mol% of Tb^{3+} ions in Tb^{3+} : Y₂SiO₅ is due to the non-radiative transition from the ${}^{5}D_{3}$ to ${}^{5}D_{4}$ levels via a cross-relaxation mechanism. The detailed cross-relaxation mechanisms by energy transfer are explained using an energy level diagram of Tb^{3+} ions. The chromaticity coordinates of the synthesized Tb³⁺ ions in the Y₂SiO₅ phosphor were calculated and found to range from a bluish color to a green color in the CIE plot with increasing Tb³⁺ ions. The color purity of 0.11 mol of Tb³⁺ -doped Y₂SiO₅ phosphors is 60.9%. Photoluminescence



Fig. 10 Typical photographs of the (1) $Y_2SiO_5:0.01Tb^{3+}$ and (2) $Y_2SiO_5:0.11Tb^{3+}$ phosphors: (a) at room light (white powder) and (b) at 254-nm UV light (green powder) (Color figure online).

results suggest that the synthesized Tb^{3+} -doped Y_2SiO_5 phosphor produces tunable color emissions upon UV excitation and an efficient, promising single-green color-emitting phosphor for display device applications.

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