CALCULATION AND COMPARATIVE ANALYSIS OF IR SPECTRA OF A NUMBER OF BRASSINOLIDES WITH DIFFERENT SIDE-CHAIN STRUCTURES

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Experimental IR spectra and theoretical spectral absorption curves of the practically important biologically active steroid phytohormones (22S,23S)-24-epibrassinolide, (22S,23S)-28-homobrassinolide, and homobrassinolide, which have similar chemical structures, were compared and analyzed to establish spectral and structural correlations of them. This enabled an evaluation of the influence of side-chain structural features on the formation of IR spectra and the establishment of connections between their structural differences and the observed changes in the IR spectra in the range $1500-950 \text{ cm}^{-1}$.

Keywords: IR spectroscopy, normal vibration analysis, absolute *IR* intensities, characteristic frequencies, interpretation of *IR* spectra, brassinosteroids, brassinolides.

Introduction. Brassinolides and castasterones are biologically active hormones of the phytohormonal steroid (PS) class. They help to improve the quality of plant products, are toxic to rapidly spreading tumor cells, and are considered potential anticancer pharmacological agents [1–4]. This group of steroids numbers greater than 70 compounds [5–7]. The biological activity of these compounds is known to depend substantially on the stereochemical configuration of the C22 and C23 atoms and the structure and configuration of the side-chain C24 alkyl substituent [8]. These could be important to the toxicity of brassinolide and castasterone derivatives for a whole series of cancer cells [9, 10].

IR spectroscopy offers huge possibilities for obtaining information about molecular structures. However, recorded IR spectra of series of PSs have been analyzed only in a few studies [11–13]. Theoretical investigations are needed to establish structure–spectra correlations in practically important biologically active PSs. These include modeling of the molecular structure, calculation of vibrational spectra, and comparative analysis of experimental and calculated IR spectra of structurally similar molecules of this compound class. Previously, vibrational spectra were fully calculated and optical-density spectral curves were modeled by us for the biologically active PSs (22S,23S)-24-epicastasterone [14], 24-epicastasterone [15], homobrassinolide, and (22S,23S)-24-epibrassinolide [16].

The aim of the present work was to compare and analyze IR spectra of natural homobrassinolide and synthetic (22S,23S)-24-epibrassinolide and (22S,23S)28-homobrassinolide, which have similar chemical structures, to evaluate the influence of structural differences in the side chains of these molecules on the formation of their IR spectra using total calculation of the frequencies and intensities of normal modes and to model the optical-density spectral curves of them in the crystalline state.

Experimental. The starting geometric parameters of (22S,23S)-28-homobrassinolide and (22S,23S)-24epibrassinolide were taken from x-ray structure analyses [17, 18]; of homobrassinolide, from molecular modeling of its crystal structure. An original combined approach to the analysis of IR spectra of complicated organic compounds that was developed and programmed for studying spectra of carbohydrates was used to calculate the frequencies and intensities of normal modes. It combined a classical analysis of normal modes with quantum-chemical evaluation of absolute intensities corresponding to integrated intensities of IR absorption bands. The algorithm for calculating intensities has been described in detail [19].

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Frequencies, shapes, potential-energy distributions (PEDs) of normal modes, and atomic shifts in Cartesian coordinates in each normal mode for calculating absolute intensities of IR absorption bands were calculated using a molecular mechanics (MM) method. The force field developed by us for calculating the frequencies and shapes of normal modes by the MM method for brassinosteroids was published [14]. Optical-density spectral curves corresponding to crystalline synthetic (22S,23S)-24-epibrassinolide, (22S,23S)-28-homobrassinolide, and natural homobrassinolide were modeled based on the results.

Results and Discussion. (22S,23S)-24-Epibrassinolide (I) consists of 82 atoms; (22S,23S)-28-homobrassinolide (II) and homobrassinolide (III), 85 atoms that form two six-membered (A and C), one seven-membered (B), and one fivemembered ring (D) and a bulky side chain. The structural differences of these molecules are localized within the side chains, which have different configurations of the 22- and 23-diol groups (22S,23S in I and II and 22R,23R in III) and different substituents in (methyl in I and ethyl in II and III) and configurations of the 24-position (24R in I and 24S in II and III). For convenience of recognizing the differences in the configurations of neighboring side-chain C atoms in the 22,-, 23-, and 24-positions, we will use the designations SSR (I), SSS (II), and RRS (III).



Homobrassinolide

IR spectra of I and III were previously analyzed by comparing experimental and theoretical absorption spectra [16]. Structurally similar II was also included in the analysis in the present work to establish more reliable structure–spectra correlations. The analyzed IR spectra were recorded on a NEXUS FTIR spectrometer. Samples of the compounds were ground with dried finely disperse KBr. The resulting powders were compressed under vacuum by the standard method. Theoretical spectra were calculated using quasi-isolated molecules. The effect of intermolecular interactions, in particular intermolecular H-bonds, on the vibrational spectra were not considered. Most attention was paid to the range $1500-950 \text{ cm}^{-1}$, which is often used for analytical purposes.

Many of the numerous normal modes are weak in the IR spectra. However, Fig. 1 shows that the density of vibrational states in this range is rather large. Theoretical spectra are presented as an envelope of Gaussian bands (Fig. 1a–c) and as vertical lines (Fig. 1d–f) showing the positions and intensities of the calculated normal modes. Figure 2 shows the calculated and experimental absorption spectra. Table 1 lists the assignments of the absorption bands of the studied molecules in the analyzed spectral range. A comparison of the theoretical and experimental spectra showed that



Fig. 1. Theoretical IR absorption spectra as Gaussian band envelopes (a–c) and as vertical lines (d–f) of (22S,23S)-24-epibrassinolide (a, d), (22S,23S)-28-homobrassinolide (b, e), and homobrassinolide (c, f) in the range 1500–1200 cm⁻¹; numbers denote normal modes.



Fig. 2. Experimental (a–c) and theoretical IR absorption spectra (d–f) of (22S,23S)-24epibrassinolide (a, d), (22S,23S)-28-homobrassinolide (b, e), and homobrassinolide (c, f) in the range $1500-1200 \text{ cm}^{-1}$.

TABLE 1. Experimental (IR) (v_{exp}) and Calculated (v_{calc}) Frequencies (cm⁻¹), Absolute Intensities *A* (10¹⁶ cm²·mol⁻¹·s⁻¹), and PED of Normal Modes of Crystalline (22S,23S)-24-Epibrassinolide (I) [16], (22S,23S)-28-homobrassinolide (II), and Homobrassinolide (III) [16] in the Range 1500–950 cm⁻¹

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TABLE 1. (

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$v_{exp}(v_p)$		1225	(1239)	1225	(1238)	1225	(1225)							1182	(1159)	1182	(1152)	1185	(1180)					1141	(1131)	1141	(1132)	1143	(1149)		

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	Ι	C 7 (s.c.) CC26H 6 CC25H 5		א אנטט פאוענ	216H 7 CC12H 6 CC15H 5 (с7н) 7 сс16н 7 осн 7 (s		C7H 5 OCH 6 (s.c.) CC16H	IC7H 6 CC2H 6 02C3H 5 0		C15H 7 04C7H 6		227H 8 CCC 7 (s.c.) CC15H	C16H 5	CC 8 04C7H 7 CC19H 5	CC 8 (s.c.) CC18H 6		5H 10 CC18H 8 CC 6 CC16I		IH 8 CC19H 8 04C7H 6 CC		15H 8 CC 6 CC17H 5	C26H 12 CC28H 12 CC21H		ости и сола « со з со о		C27H 13 CC26H 11 CCC 7	
		25H 5 CC20H 5 CC4			5H 5 CC3H 5 CC2	H 7 (s.c.) CC15H (040	CC1	C16H 5 (040	3H 5 CC16H 5 CC1	CC1	0C2	5 + 6	C15H 5 CC16H 5 CC1	CC1	5 CC1	C25(CC16H 6 CC8H 5 CC2		6 CC12H 6 CCC		CCC	CC21H & CCC 5 CCC		CC27H 6 CC1		CC 7 (s.c.)	
PED, %	Π	tH 13 CC3H 5		10H 10 CC4H 6 CC22H 5 OSC22H 5	1911 / CC411 9 CC311 3 CC1311 3 28H 14 CCC 13 (s.e.) CC29H 8 CC25H 6 CC24H 5	C7H + CC7H) 12 CC25H 7 CC26H 5 CC16H 5	12H 5	C7H + CC7H) 10 CC25H 9 CC26H 7 CC16H 6	15H 7 CC2H 6 OC7H 6 CCC 6 CC 5 O2C3H 5	16H 5	22H 7 OC23H 6 OC7H 6 CC21H 5 CC 5 + CC0	CCC 6 (s.c.)	16H 15 CC15H 8	ISH 22 CC16H 7	15H 12 CC27H 6 CC26H 6 04C7H 5 CC12H 5	C26H 8 C25C27H 8 CCC 6		21H 10 CCC 10 (s.c.) CC18H 5 CC27H 5 CC26H 5		C 15 CC15H 10 CC18H 8 CC16H 6 CC 6 CC8H 5		C 10 CC11H 8 CC19H 7 04C7H 5 CC12H 5	C 17 CC15H 10 CC 7 CC17H 5		1911 23 C24C2011 10 CC2/11 10 CC2011	1		
	III	СС12H 10 СС2H 9 СС3H 7 ССО 7 СС4H 6 СС 6	02C3H 5	ССІБН 13 04С/Н 11 С8С/Н 9 ССІ2Н / ССІЗН 6 ССЛН 17 ССІКН 7 ССС 7	CC12H 15 CC4H 12 CC16H 5	CCH 14 (s.c.) CC16H 12 CCH 8 (skel.) CCC 5		CCH 20 (s.c.) CC15H 8 CCH 6 (skel.) (CC0+CCC)	6 (s.c.)	CCH 16 (skelet.) CCC 11 02C3H 6 CC2H 6 CC4H	6 CCO 5 CC 5	(CCH 18 CCC 13) (s.c.) CC15H 7 CC 5 CC21H 5	05CH 5 CC27H 5	CC16H 25 CC15H 15 CCH 7 (skelet.)	CC15H 16 CCH 9 (skelet.) CCH 7(s.c.) CC16H 5	CC15H 16 CCH 11 (skelet.) CCC 9 CC14H 5	CC12H	CC26H 15 CC27H 12 (CCH 13 CCC 7 CC 6) (s.c.)	C24C28H 5 –	CCC 16 CCH 10 (skelet.) 04C7H 8 CC15H 7	CC16H 5 CC8H 5 CC18H 5	CCH 17 (skelet.) CCH1 11 (s.c.) CCC 5	CC29H 15 CC21H 8 CC27H 8 CC26H 8 CCC 8	(s.c.) CC28H 6	CCC 12 CC1311 IV CC11 7 (SKC161.) UC 1011 3 CC 2711	C28C29H 16 CC21H 13 C24C28H 7 C29C28H 6	I	

TABLE 1. (Continued)

	Ш	19H 13 CCC 11 CC3H 7 CC 5 CC4H 5		29H 14 CC21H 12 CCC 9 CC18H 7	18H 20 CC21H 14 CC15H 6 CC16H 5 CC29H 5		19H 22 C8C7H 5 CC11H 5 CC21H 5	H 12 (skelet.) CC19H 11 CC18H 7 CC15H 6	tH 6 CC11H 5 CCC 5	15H 10 CC18H 9 CC21H 8 CC19H 7 CC16H 6
PED, %	П	CC18H 12 CC21H 10 CC29H 7 CC15H 6 CC16H 5 CC	CCC 5	CC18H 20 CC21H 19 CC	CC18H 11 CC15H 9 CC21H 7 CC16H 5 CC27H 5 CC	CC19H 5	CC19H 18 CC4H 6 CC27H 5 CC	- CC	CC	- CC
	I	CC19H 21 CCC 9 CC7H 8 HC19H 6 CC18H 5		CC18H 14 CC15H 12 CC16H 9 CC21H 9	CC19H 17 CC11H 8 CC21H 6 HC19H 5		I	I		1
1	III I	3 3		8	7 1		4	1		ŝ
4	I I	7 43) 1(1 37		- 62	 		
	III	92		73 5	158		151 -	- 120		945 -
'calc	П	67 9		58 9	45 9		41 9	6		6
1	Ι	354 5		946 5	35 5		- 2			
$v_{exp}(v_p)$		964	(954)	5 62	(967)	985	(974)			

Note. s.c., side chain; frequencies v_p of absorption band maxima of theoretical spectra of I, II, and III are given in parentheses.

the theoretical spectra were in general satisfactorily described by the main and strongest characteristic absorption bands of the examined region. For example, the maximum frequency shift of corresponding v_{calc} and v_{exp} bands was from 14 cm⁻¹ (1500–1200 cm⁻¹ range) to 23 cm⁻¹ (1200–950 cm⁻¹) for I, from 13 to 30 cm⁻¹ for II, and from 10 to 23 cm⁻¹ for III.

Table 1 lists the calculated frequencies, PEDs, and absolute intensities of normal modes located in the studied region that contributed most to the intensity of the summed band and the frequencies of band maxima in the theoretical spectra. It can be seen that each band of the experimental spectra had a complicated origin and consisted of several bands of various integrated intensities (the sum of Gaussian bands over the calculated absolute intensities and half-widths set from 8 to 16 cm^{-1}). For example, the strong band with two maxima at 1467 (I and II), 1469 (III), and 1445 cm⁻¹ in the experimental spectra (1500–1425 cm⁻¹) consisted of 15 (I and III) and 18 (II) constituent components.

The high-frequency part of this band with maxima at 1469 (III) and 1467 cm⁻¹ (I and II) in the studied compounds was formed mostly by bending vibrations of methyl and methylene groups. The low-frequency part with a maximum at 1445 cm⁻¹ originated in bending vibrations involving methyl and methylene groups and hydroxyls. An analysis of the data in Table 1 also suggested that slight structural differences in the molecules shifted the frequencies of several normal modes of similar shapes. In particular, conformational differences in the terminal part of their side chains affected the frequencies and intensities of normal vibrations with the main contribution to the PED of methyl and methylenes located in just this part of the side chain. For example, the normal mode with the main contribution to the PED from bending of C27H₃ methyl bonds had a frequency of 1424 cm⁻¹ in I, 1457 cm⁻¹ in II, and 1445 cm⁻¹ in III.

Replacing the C28H₃ methyl (I) in the side-chain C24-position by a C28H₂C29H₃ ethyl (II and III) led to the appearance (instead of the normal mode with the main contribution to the PED from bending of C28H₃ methyl bonds with a frequency of 1451 cm⁻¹) of a normal mode with the main contribution to the PED from bending of groups in the C28H₂C29H₃ ethyl in II at 1437 cm⁻¹ and in III at 1472 cm⁻¹. Figure 1 shows lines corresponding to these normal modes that are designated 1 and 2. The high-frequency shift [1398 (II), 1402 (III), 1437 cm⁻¹ (I)] (Fig. 1, line 3) of the normal mode for the O5H hydroxyl and C22H and C23H methyls of the side-chain diol group could be explained by the configuration change of the C24–C28 bond [24S (II and III) and 24R (I)].

A change of the intensity ratio of the spectral maxima was characteristic of the second strong band in the range $1500-1200 \text{ cm}^{-1}$ with two maxima at 1404, 1385 (III), and 1386 cm⁻¹ (I and II) in the experimental spectra (1425–1350 cm⁻¹) of the studied molecules. Also, the ratio of the peak intensity changed as compared to peak intensities of close-lying bands in the experimental spectra. This band strengthened with respect to bands with maxima at 1467 (I and II) and 1319 (I and II) and 1318 cm⁻¹ (III) on going from I and II to III. The ratio of peak intensities in these molecules followed an analogous trend in the theoretical spectra (Fig. 2).

An analysis of Table 1 showed that bending vibrations of methyls C26H₃, C27H₃, and C29H₃ situated in the terminal part of the side chain in III (1422 and 1420 cm⁻¹, Fig. 1, lines 4 and 5) contributed most to the intensity of the high-frequency portion of the band with a maximum at 1404 cm⁻¹; in II, bending vibrations of methyl C19H₃ located at the junction of rings A and B of the molecular skeleton (1412 and 1410 cm^{-1} , lines 4' and 5'). Highly active bending vibrations of methyl C27H₃ (1424 and 1404 cm⁻¹, lines 1 and 5) and bending vibrations of methylenes C16H₂, C11H₂, and C1H₂ of the molecular skeleton (1415 and 1411 cm⁻¹, lines 4' and 5') in I contributed noticeably to the high-frequency portion of this band. Normal modes with vibrational energy localized on both atoms incorporated into the molecular skeleton (ring A, 1386 cm⁻¹, line 6) and in the terminal part of the side chain (1381 cm⁻¹, line 7) contributed most to the calculated integrated intensity of the low-frequency portion of this band with a maximum at 1385 cm⁻¹ in III. Conversely, normal modes with the main contribution to the PED of HCH and CCH bending vibrations of ethyl C28H₂C29H₃ that was immediately adjacent to the diol group (1395 and 1384 cm⁻¹, lines 6' and 7') had the greatest absolute intensity in the calculated spectrum of II. Normal modes with the main contribution to the PED of HCH and CCH bending vibrations in the calculated spectrum of I both in the molecular skeleton (1391 and 1387 cm⁻¹, lines 6 and 7) and in the diol group and the atomic group immediately adjacent to it (1371 and 1355 cm⁻¹, lines 6' and 7') had the greatest absolute intensities. Because the diol located on C22 and C23 had different configurations in I and II (22S,23S) and III (22R,23R), it obviously had different effects on the frequencies and intensities of the normal modes in this range. Thus, the main changes in the ratios of both the maximum intensities at 1404 and 1385 (III) and 1386 cm⁻¹ (I and II) and its intensity and close-lying bands in experimental spectra of the studied molecules were related to the different configurations of the diol and conformations of the side-chain terminal part as this band formed in them.

A third strong band of the experimental spectra in the range $1500-1200 \text{ cm}^{-1}$ was located in the range $1350-1300 \text{ cm}^{-1}$; had maxima at 1319 (I and II) and 1318 cm⁻¹ (III) [at 1309 (I), 1307 (II), and 1308 cm⁻¹ (III) in the

calculated spectra] and shoulders at ~1331 (I and II) and 1327 cm⁻¹ (III) [1337 (I), 1331 (II), and 1325 cm⁻¹ (III)] (Fig. 2); and consisted according to the calculations of 12 (I) and 13 (II and III) components. Delocalized CCH bending vibrations (and OCH to a lesser extent) of all methine (CH) and methylene (CH₂) groups of the molecular skeleton made the main contribution to the formation of this band. These vibrations [1339 and 1309 cm⁻¹ (I), 1330 and 1308 cm⁻¹ (II), 1343 and 1307 cm⁻¹ (III)] (Fig. 1, lines 8 and 9) in rings A and C of the molecular skeleton contributed most to the intensity (Table 1). Two weaker bands in this range with maxima at 1275 (I and II) and 1276 cm⁻¹ (III) [1271 (I), 1273 (II), and 1273 and 1260 cm^{-1} (III) in the calculated spectra] and 1225 cm^{-1} [1239 (I), 1238 (II), and 1225 cm⁻¹ (III)] (ranges 1300–1250 and 1250–1200 cm⁻¹) (Fig. 2) consisted according to the calculation of six (I and II) and seven (III) and eight (I and II) and seven band components (III), respectively. Like for the band discussed above, delocalized CCH bending vibrations of methines and methylenes made the main contribution to their formation (Table 1). A change in the intensity ratio of their peak intensities in experimental spectra of the studied molecules, i.e., a decrease in the intensity of the first band on going from III to I and II, was a characteristic feature of these two bands. An analogous trend was observed in the calculated spectra of these molecules (Fig. 2). Table 1 also shows that normal modes in which the vibrational energy was localized on diol groups and neighboring side-chain groups (1259 cm⁻¹, Fig. 1, line 10) contributed most to the integrated intensity of the first band in the calculated spectrum of III; vibrations involving the molecular skeleton [1296 cm⁻¹ (I), line 10 and 1290 cm^{-1} (II), line 10'] and the diol and neighboring side-chain groups [1268 cm⁻¹ (I), line 10'], in I and II.

Delocalized CCH bending vibrations of methine (CH) and methylene groups (CH₂) of the molecular skeleton, among which normal modes with frequencies 1248 (III), 1239 and 1233 (II) (Fig. 1, lines 11, 11', and 12'), and 1241 and 1232 cm⁻¹ (I) (lines 11 and 12) contributed most to the integrated intensity despite the structural differences of these molecules being localized in the side chains, made the main contribution to the formation of the band at 1250–1200 cm⁻¹ in these molecules. It could be assumed that the change in the intensity ratio of the absorption bands in this range was related to the different configurations of the side-chain diol.

Experimental spectra of the studied compounds in the range 1200–900 cm⁻¹ consisted of a strong band with maxima at 1182 (I and II) and 1185 cm⁻¹ (III); a doublet with maxima at 1141 (I and II) and 1143 (III) and 1119 (I and II) and 1126 cm⁻¹ (III); and a broad absorption band with two clearly resolved maxima at 1064 (I), 1061 (II), and 1063 (III) and 1025 (I), 1023 (II) and 1027 cm⁻¹ (III) (Fig. 2). The strong band with maxima at 1185 (III) and 1182 cm⁻¹ (I and II) in the experimental spectrum (1200–1150 cm⁻¹) consisted according to the calculations for the studied compounds of 11 components, the main contribution to the formation of which came from bending vibrations of CCH and OCH methines and methylenes, CCC bending vibrations in the skeleton and side chain, and C–C stretching vibrations. This band strengthened relative to the absorption band with a maximum at ~1225 cm⁻¹ on going from III to I and II.

Distinguishing features of this band as compared to those discussed above were a shape change of the vibrations in most normal modes forming it and a more noticeable role of the side chain, in particular, ethyl C28H₂C29H₃. The absolute intensity of the normal mode at 1198 cm⁻¹ in II increased significantly as compared to the analogous vibration at 1202 cm⁻¹ in III. Bending vibrations of atomic moieties located in rings *A* and *B* in III and in rings *B*, *C*, and *D* in I made the main contribution to its formation. These were normal modes with frequencies 1181 and 1174 (III) and 1162 and 1156 cm⁻¹ (I). It could be assumed that these changes in the formation of this band were associated with the different diol configurations of side-chain C22 and C23 of I, II, and III and with the different configurations of the C24–C28 bond [R (I) and S (II and III)].

Also, ethyl C28H₂C29H₃ in II played an important role in the calculations in forming a doublet with maxima at ~1141 (I and II) and 1143 (III) and 1119 (I and II) and 1126 cm⁻¹ (III) in the experimental IR spectra (1150–1100 cm⁻¹). Bands with maxima at 1131 and 1112 (I), 1132 and 1104 (II), and 1149 and 1122 cm⁻¹ (III) in the theoretical spectra corresponded to them (Fig. 2). Bending vibrations of ethyl CC28H contributed greatly in addition to CCH and OCH bending vibrations of methines and methylenes of the molecular skeleton to the total intensity of this doublet in II. This affected the intensity ratio of both maxima in this doublet absorption band (Fig. 2).

Experimental spectra of the molecules in the range $1100-950 \text{ cm}^{-1}$ showed a broad absorption band with two maxima that were strongest in the range $1500-950 \text{ cm}^{-1}$ at 1064 (I), 1061 (II), and 1063 (III) and 1025 (I), 1023 (II), and 1027 cm⁻¹ (III) and a weaker band at ~985 (III) and a shoulder at ~964 (I) and 967 cm⁻¹ (II) (Fig. 2). According to the calculations, these bands consisted of 20 component bands. An analysis of Table 1 showed that CCH bending vibrations of methyls, CCH and OCH bending vibrations of methylenes, and bending (CCC) and stretching vibrations (C-C) of the skeleton and side chain made the main contribution to the formation of these strong maxima [at 1074 (I), 1075 (II), and 1086 (III) and 1003 (I), 1001 (II), and 1025 cm⁻¹ (III) in the theoretical spectra].

The activities of the C15H₂ and C16H₂ methylenes in five-membered ring *D* and the C7H₂ methylene in sevenmembered ring *B* were noteworthy. The CCH and OCH bending vibrations involving these groups contributed noticeably to the integrated intensity of experimental bands with maxima near 1064 (I), 1061 (II), and 1063 (III) and 1025 (I), 1023 (II), and 1027 cm⁻¹ (III). Furthermore, the ethyl (methyl in I) on side-chain C24 (and the C24–C28 configuration) had a significant effect on the formation of bands at 1025 (I), 1023 (II), and 1027 cm⁻¹ (III) in experimental spectra of the studied molecules. For example, the normal mode in II at 998 cm⁻¹ that made the main contribution to the PED of the bending vibration of the CCH localized in this group had the maximum absolute intensity as compared to normal modes of this frequency range and contributed most to the integrated intensity of the band maximum at 1001 cm⁻¹ in the theoretical spectrum, which corresponded to the band maximum at 1023 cm⁻¹ in the experimental spectrum. Normal modes close to it in III (1010 and 994 cm⁻¹) with noticeable contributions to the PED of the ethyl CCH bending vibrations contributed less to the integrated intensity of the analogous band (1027 cm⁻¹). As a result, the intensities of band maxima at ~1027 (III) and 1023 cm⁻¹ (II) changed on going from III to II in the theoretical and experimental spectra of the studied molecules.

Bending vibrations of CCH groups localized in this group and those adjacent to it in I made the main contribution to the PED of normal modes with frequencies 1009, 992, and 975 cm⁻¹ and contributed noticeably to the integrated intensity of the band maximum at 1003 cm⁻¹ in the theoretical spectrum, which corresponded to the band with a maximum at 1025 cm⁻¹ in the experimental spectrum. A weaker band at ~985 cm⁻¹ in the low-frequency region of this spectral range of III transformed into a shoulder at ~967 cm⁻¹ in the spectrum of II and at ~964 cm⁻¹ in that of I (Fig. 2).

According to Table 1, the normal mode (973 cm⁻¹) with a significant contribution to the PED of the bending vibration of the CC29H methyl in the side-chain ethyl contributed noticeably to the integrated intensity of this band in the experimental spectrum of III. However, normal modes of atomic moieties in the molecular skeleton contributed noticeably to the integrated intensity in this range for I and II. Therefore, it could be proposed that the intensity change (transformation into a shoulder) of this absorption band in spectra of I and II was related to the different configurations of the side-chain diol.

Conclusions. A comparative analysis of calculated vibrational spectra of structurally similar PSs allowed a relationship to be found between their structural differences and changes of the IR spectra. It was found that small differences in the molecular structures led to either shape changes of several normal modes with similar frequencies or frequency shifts of several normal modes with similar shapes. As a result, the integrated intensities, shapes, and half-widths of the corresponding absorption bands changed. It was also found that differences in the C22 and C23 diol configuration (22S,23S in I and II and 22R,23R in III) and the conformation of the terminal part of the side chain in the studied molecules had large effects on the integrated intensities, shapes, and half-widths of absorption bands with maxima near 1404 and 1385 (III), 1386 (I and II), 1276 (III), 1275 cm⁻¹ (I and II) and 1225, 1143 (III), 1141 (I and II) and 1126 (III), 1119 (I and II) and 1027 (III), 1025 (I), 1023 (II), and 985 cm⁻¹ in the experimental IR spectra.

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