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Investigation of native cellulose under high pressure using microfocused synchrotron radiation

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Abstract The mechanical properties of native cellulose are critical for understanding the properties of natural biomaterials. To investigate the elastic moduli of the cellulose crystalline fraction an isotropic mechanical load using hydrostatic pressure ranging from 0.01 to 0.5 GPa was applied to flax fibers, pine wood and tension wood samples. The response of the crystalline part was monitored by using microfocused synchrotron radiation. The compressibility of the crystalline fraction of native cellulose was anisotropic and dependent on the crystal size and possibly on the composition of the sample. The compressibilities along the [001] direction, where covalent bonds hold the cellulose chains together, varied between 2.1 and 2.9 TPa⁻¹ for the different samples and confirm

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M. Burghammer · E. di Cola European Synchrotron Radiation Facility, 71, avenue des Martyrs, CS40220, 38043 Grenoble Cedex 9, France similar values found for cotton fibres. However, the compressibilities along the [100]-direction range from 56.2 TPa⁻¹ to 63.5 TPa⁻¹, slightly exceeding the previously determined value of 50 TPa⁻¹, which can possibly be attributed to differences between individual samples.

Keywords Cellulose microfibrils · Mechanical properties · High pressure · Synchrotron radiation · Microbeam

Introduction

Native cellulose often appears in alliance with an amorphous matrix of lignin, hemicellulose and other components, building a semi-crystalline composite material with outstanding mechanical properties. Despite its low material density, it allows plants to withstand continuous high mechanical forces such as wind. Cellulose possesses a hierarchical structure from the molecular over the crystalline level to the so called microfibrils forming a helix with a certain angle to the cell axis referred to as microfibril angle (Salmén and Burgert 2009). Being synthesized in sixfold rosette complexes also referred to as terminal complexes, parallel orientated cellulose chains form crystals via hydrogen bonds and van-der-Waals force (Kölln 2004; Taylor 2008). The microfibrils have a small, specific diameter ranging from 2 to 25 nanometers and are over several micrometers in length (Nishiyama 2009). They twist around their long axis (Willhammar et al. 2021) and are considered to contain unordered regions (Fengel and Wegener 1983). The unit cell parameters for the crystal structure of the dominant allomorph cellulose I β provide a monoclinic unit cell with the parameters of a=7.78 Å, b=8.20 Å, c=10.38 Å and γ =96.5° (Nishiyama et al. 2002).

Determining the elastic modulus of cellulose is critical to understand the material properties of naturally occurring biomaterials. The elastic tensor of the monoclinic cellulose allomorph IB consists of 13 independent elements and is highly anisotropic (Nye 1985; Song et al. 2021). To understand the interplay of the different levels and the elasticity of native cellulose, in situ tensile tests have been performed on wood and flax fibers using synchrotron radiation (Reiterer et al. 1999; Kölln et al. 2005a; Clair et al. 2006). The effect on the microfibrils is visible through the shift of the Bragg reflections 200 and 004 in the diffraction pattern, roughly corresponding to the deformation in [100]- and [001] direction. In other approaches, nanoindentation and in situ bending techniques were applied to determine two diagonal elements of the elastic tensor using materials containing native cellulose (Konnerth and Gindl 2006; Orso et al. 2006). Diddens at al. performed inelastic X-ray scattering and calculated the two diagonal elastic tensor elements of native cellulose in flax via the measured sound velocity to 220 GPa in [001]-direction and 15 GPa in [100]-direction. These were also calculated theoretically with molecular dynamics simulations and vary between 98 and 202 GPa in in [001]-direction (Kroon-Batenburg and Kroon 1997; Eichhorn and Davies 2006; Dri et al. 2013; Song et al. 2021) and 14.5-19 GPa (Dri et al. 2013; Song et al. 2021) in [100]-direction.

However, it is important to take into account that often the samples used in many experiments have been chemically treated. Therefore, the results can only be transferred to a limited degree to native cellulose, where various parameters such as the interplay of the different cell wall layers, the deformation of the amorphous matrix under uniaxial load (Keckes et al. 2003; Peura et al. 2007) as well as changes in the alignment of the microfibrils towards the fiber axis (Montero et al. 2011) cannot be neglected. When interpreting the elastic modulus of native cellulose, it is important to take into account both the rotation of bond angles (Krywka et al. 2014) and hydrogen bonds (Mariano et al. 2014) at the atomic level.

In the present study native cellulose was subjected to high hydrostatic pressure ranging from 0.01 to 0.5 GPa. By exposing matrix and crystals to hydrostatic pressure, the mechanical load is considered to be homogeneous and non-directional. It is expected that the microfibrils and matrix will not be structurally affected, but rather the cellulose crystals will be compressed according to their anisotropic compressibility due to the heterogenous structure of the crystals.

The hydrostatic compression allows the direct investigation of the a- and b-axis compressibility undisturbed by effects usually present in tensile tests (e.g. the Poisson's ratio). However, the hydrostatic method implies that the results obtained here is not directly comparable to elastic constants measured in non-hydrostatic conditions.

Materials and methods

Experimental set-up and data collection

For this experiment, a bundle of industrially bleached flax fibers (Diddens et al. 2008), two untreated pieces of pine wood and two tension wood samples from poplar (Müller et al. 2006) with a thickness of 0.5 mm and a length of 20 mm each were chosen. A hydrostatic high pressure sample cell was used to perform the measurements (Krywka et al. 2008). The cell body is machined of a high tensile strength alloy, allows for pressures up to 7 kbar and contains the sample inside the pressure chamber (see Fig. 1a). Two diamond windows each with a thickness of 1 mm are stable against the pressure and allow the radiation to pass. The sample is mounted to the reusable windowless sample carrier, allowing the sample to be directly exposed to the water pressure (compare Fig. 1b).

The measurements with microfocused synchrotron radiation were carried out at the microfocus beamline ID13 of the European Synchrotron Radiation Facility (ESRF, Grenoble, France). The experiment was performed at an X-ray energy of 22.977 keV. After the monochromator, the beam was focussed through compound refractive lenses and conditioned by a slit system, resulting in a size of $30 \times 30 \ \mu\text{m}^2$ with a flux of 10^8 – 10^9 photons/s. After passing through an ionization chamber, the beam was directed towards



Fig. 1 a Assembly of the hydrostatic pressure cell with diamond windows including the sample holder with a fibre bundle; **b** schematic cross-section of the hydrostatic high pressure cell showing how the sample is subjected to hydrostatic pressure



Fig. 2 Diffraction image of a pine sample submerged in water at the initial pressure of 10 MPa applied. The different reflections and the referring crystal planes are shown for the 200 and the 004 reflection (adapted from (Diddens et al.

the sample through the diamond windows of the high pressure sample cell. The first diffraction image

2008)). The linear artefacts are Kossel lines originating from the diamond windows. The fiber axis in the experiment was oriented horizontally

(compare Fig. 2) was taken at a pressure of about 10 MPa, in order to ensure that no bubbles remained

in the pressure chamber. The pressure was then increased up to 0.5 GPa with a step size of 100 MPa. Afterwards, the pressure was reduced again with the same step size. The procedure with the increasing pressure was repeated without the sample to obtain the background data.

Data analysis

After normalization for exposure time and ionisation chamber values the data were subjected to background subtraction. The transmission was estimated taking into account the pressure-dependent attenuation by water and further used to normalize the scattered intensity.

To understand the effect of pressure on the reflection position, azimuthal integration was conducted on the respective areas of the 200 reflection for the a-axis, the 110 and 1 1 0 reflection and the 004 reflection for the c-axis at each pressure step (compare Fig. 3). When performing a complete azimuthal integration (see Fig. 4), it can be seen that depending on the sample the 110 and 1 1 0 reflections are resolved differently. The variation in crystallinity, crystal sizes, and orientations in the different samples accounts for this difference. The lower limit of the material's

$$L_{hkl} = \frac{2\sqrt{(3\ln 2)\lambda}}{\pi FWHM \cos\theta} \tag{1}$$

not taking into account other effects broadening the peak. Table 1 shows the calculated lower limits of the mean cross section diagonal (L_{200}) and the length of the crystals calculated from the 004 reflection (L_{004}).

The positions of the integrated reflections in reciprocal space were determined by fitting them with Gaussian functions. The corresponding d-spacing in real space were calculated using $d(p) = \frac{2\pi}{q(p)}$. The reflection positions were plotted as a function of

	Pine wood	Flax	Tension wood	
L ₂₀₀ /Å	35.5	55.2	71.0	
L ₀₀₄ /Å	356.0	293.2	356.0	



Fig. 3 left: Tension wood raw data at 0.5 GPa; right: tension wood data at 0.5 GPa after background subtraction. The area marked in orange was integrated to determine the 200 reflection

pressure and fitted with a linear function. Extrapolating this function to zero pressure allows for the calculation of strain $\epsilon = \left| \frac{l - l_0}{l_0} \right|$ which, since pressure is always present, can be considered as stress (σ) on the cross-section perpendicular to the crystal axis. With this, Young's modulus (in specific crystallographic directions) can be calculated from Hooke's Law: $\varepsilon = \frac{\sigma}{r}$. Compressibilities can directly be determined for the crystallographic [100] and [001] directions of native cellulose crystals via $\kappa = \frac{3(1-2)}{E}$ with $\upsilon = 0.377$ as determined before (Nakamura et al. 2004). Comparison with data collected while decreasing the pressure showed some differences, which indicate that viscoelastic effects and potentially also humidity effects could play a role here. Therefore, only data recorded while the pressure was increased were considered. The final data point, collected after depressurization, was used for comparison with the first data point to verify that no radiation damage occurred during the data collection.

To determine the unit cell parameter along the a-axis, the formula $a = 2 \frac{d_{200}}{\sin \gamma}$ was employed., γ was assumed to be constant with a value of 96.5° according to the literature. The unit cell parameter of the c-axis was calculated by $c = 4 d_{004}$. Because the unit cell parameter for the b-axis cannot be determined directly from the available 110 and 110 reflections, which are additionally quite noisy, the d₁₁₀ and d₁₁₀ values are given directly (compare Table 2).



Fig. 4 Radially integrated data after the subtraction of the background image showing the 110, the $1\overline{10}$, the 200 and the 004 reflection for flax, pine and tension wood

Results and discussion

As shown in Table 2, the values of the different samples are all slightly larger than the theoretically determined value of 7.78 Å for cellulose I β , as reported before for native cellulose. The expected compression in the [100] direction is clearly visible for all samples (see Fig. 5a). The compression of pine wood and tension wood samples is very similar when compared to each other. The weak van-der-Waals-bonds allow a compression between 2.83% in the case of tension wood and 3.23% in the case of pine wood when exposed to 0.5 GPa. Relating the observed compression to the crystal sizes, it becomes clear that the compression is relatively stronger for smaller crystals

Table 2 For the a- and
the c-axis, the unit cell
lengths at 10 MPa as well
as the compression and
the strain at 0.5 GPa were
calculated. Additionally,
the compression and strain
were also computed for the
d-spacings obtained from
the 110 and 110 reflections

	Pine wood 1	Pine wood 2	Tension wood 1	Tension wood 2	flax
a/Å	7.93	7.94	7.84	7,84	7.83
a-axis к/ TPa ⁻¹	62.7	63.5	56.2	58.2	56.9
ε/%	3.17	3.23	2.8	2.95	2.88
$d_{110}/\text{\AA}$	6.06	6.05	5.91	5.92	5.99
ε/%	0.30	1.40	0.70	0.90	1.90
d ₁₁₀ /Å	5.28	5.29	5.41	5.41	5.34
ε/%	1.3	1.2	0.6	1.0	1.6
c/Å	10.34	10.33	10.37	10.37	10.36
c-axis κ/TPa ⁻¹	2.1	2.2	2.9	2.1	2.3
ε/%	0.11	0.11	0.15	0.11	0.12





Fig. 5 Linear compression of for flax, pine and tension wood of the \mathbf{a} [100] direction as derived from the 200 reflection for the unit cell parameter a and \mathbf{b} of the [001] direction derived

of pine in flax compared to the bigger crystals of tension wood.

The compressibility along the [100] direction for the bundle of flax fibers (s. Table 2) is 58.7 TPa^{-1} . If the off-diagonal element of the elastic tensor is neglected, this corresponds to a modulus of 17.2 GPa, which is slightly higher than the value of 15 GPa obtained with inelastic X-ray scattering from a bundle of fibers (Diddens et al. 2008) and significantly lower than the 24.7 GPa obtained from a single flax fiber along the fiber direction. (Kölln et al. 2005b). Furthermore, a change of the humidity caused the surrounding water may have an influence at high pressure as Young's modulus for flax seems to decrease with humidity (Davies and Bruce 1998).

Along the [100] direction, the compressibilities for pine wood are with 62.6 TPa^{-1} and 63.5 TPa^{-1} higher than for tension wood with 56.2 TPa^{-1} and 58.2 TPa^{-1} , respectively. Results from other studies (Quesada Cabrera et al. 2011; Song et al. 2021) show slightly lower values of 50 TPa^{-1} . More systematic studies are required to investigate whether this is sample-dependent or whether there are other explanations for these values.

The length of the unit cell along the c-axis is slightly smaller, ranging from 10.33 to 10.37 Å, compared to the value of 10.38 Å previously determined by Nishiyama. The degree of compression in the [001]-direction is relatively smaller than that in the

from the change of the position of the 004 reflection for the unit cell parameter c, both fitted with a straight line.

a-direction, resulting in compressions of 0.11-0.15%. The compressibilities range from 2.1 TPa⁻¹ for pine wood sample 1 up to 2.9 TPa⁻¹ for tension wood sample 1 and are therefore higher than the value of 1.71 TPa⁻¹ determined by Song et al. for chemically treated tunicate, but in the range of the value of 2.8 TPa⁻¹ determined by Quesada Cabrera et al. for cotton cellulose fibres. In the case of tension wood, the differences between samples of the same type are also significant. However, in contrast to the other samples with an error of less than 2%, the error for the compressibility of the tension wood samples amounts for 3.8% and 10.7% respectively.

Given that there are only covalent bonds in [001]-direction, one would actually expect only little compression, based on previous experiments indicating that covalent bonds in proteins are pressure insensitive up to at least 10–15 kbar (Mozhaev et al. 1996). The observed compression can most likely be explained by other processes.

High pressure studies indicate that the most likely explanation is a change of the crystal packing. The largest contributions are expected to arise from hydration changes that accompany non-covalent interactions (Mozhaev et al. 1996). Similar compressibilities found by Quesada Cabrera at al were attributed to hydrogen bonds which are likely to strengthen upon compression. However, this study was performed without a pressure-transmitting media. As water is used in the study presented here, protein-water interactions can form, promoted by pressure producing stronger and shorter bonds that decrease the overall volume (Boonyaratanakornkit et al. 2002).

The intensity of the 110 and the 110 reflections are low compared to the background, the fits of the reflection positions come with significantly higher error bars compared to the other reflections. In Fig. 3, one can clearly see that the two reflections are best resolved for flax. The values for the $d_{1\bar{1}0}$ distances ranging from 5.91 Å for tension wood up to 6.06 Å for pine wood as well as the d_{110} distances with values from 5.28 Å for pine wood and up to 5.41 Å for tension wood are in good agreement with values reported before (Quesada Cabrera et al. 2011).

Conclusion and outlook

We have shown here that the unit cell of three different types of native cellulose shows a highly anisotropic behaviour under hydrostatic pressure, which can be attributed to the difference between covalent bonds and van-der-Waals interactions. The compressibilities along the [001] direction of the cellulose crystals are in line previous reports. For the [100] direction, the compressibilities are higher than previously reported. More systematic studies are required to investigate whether this is sampledependent or whether this be attribute to the method applied. In order to exclude a potential influence of the pressurizing water onto the cellulose structure, similar experiments could be performed with a water-free pressurizing medium (e.g. alcohols, silicone oil).

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Declarations

Conflict of interest The authors declare no conflict of interest.

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