Defining an Affine Partition Algebra



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Abstract

We define an affine partition algebra by generators and relations and prove a variety of basic results regarding this new algebra analogous to those of other affine diagram algebras. In particular we show that it extends the Schur-Weyl duality between the symmetric group and the partition algebra. We also relate it to the affine partition category recently defined by J. Brundan and M. Vargas. Moreover, we show that this affine partition category is a *full* monoidal subcategory of the Heisenberg category.

Keywords Affine partition algebra · Schur-Weyl duality · Heisenberg category

Mathematics Subject Classification (2010) $16 \cdot 20$

1 Introduction

Classical Schur-Weyl duality relates the representations of the symmetric group and the general linear group via their commuting actions on tensor space. The Brauer algebra was introduced in [1] to play the role of the symmetric group in a corresponding duality for the symplectic and orthogonal groups. The partition algebra was originally defined by P. Martin in [15] in the context of Statistical Mechanics. V. Jones showed in [10] that it appears in another version of Schur-Weyl duality. More precisely, if one replaces the general linear group by the finite subgroup of all permutation matrices then the centraliser algebra of its action on tensor space is precisely the partition algebra. The aim of this paper is to define an affine version of the partition algebra.

There are different 'affinization' processes for such algebras. One such process amounts to making the Jucys-Murphy elements of the ordinary algebra into variables, retaining some of the relations between these variables and the standard generators of the ordinary algebra. Starting with the symmetric group algebra, this 'affinization' process gives rise to the much-studied degenerate affine Hecke algebra (see for example [13]). In the case of the Brauer

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¹ Department of Mathematics, City University of London, Northampton square, London, EC1V0HB, UK algebra, M. Nazarov used this process in [16] to define the affine Wenzl algebra (also called the Nazarov-Wenzl algebra or the degenerate affine BMW algebra in the literature). This process was also employed independently in both [19] and [20] to define a degenerate affine walled Brauer algebra.

R. Orellana and A. Ram introduced a different 'affinization' process in [17] focussed on extending the Schur-Weyl dualities, via the affine braid group. They have applied it to the symmetric group, Brauer algebra and their quantum analogues. This process naturally leads to cyclotomic quotients of the affine algebras.

In this paper, we follow the first approach to define an affine partition algebra by turning the Jucys-Murphy elements for the partition algebra into variables and ask them to retain certain relations with the generators. But we also show that this affine partition algebra naturally extends the commuting action on tensor space with the symmetric group. We started this work by trying to use the presentation for the partition algebra given by T. Halverson and A. Ram in [8] but were unable to define an algebra with the expected properties in this way. So we instead used the more recent presentation given by J. Enyang in [6] (which uses a new set of generators) to define the affine partition algebra \mathcal{A}_{2k}^{aff} . We prove that it satisfies many properties analogous to those for other affine diagram algebras.

While writing this paper, J. Brundan and M. Vargas produced a preprint [2] defining an affine partition category APar as a monoidal subcategory of the Heisenberg category generated by some objects and morphisms. Taking an endomorphism algebra in their category gives an alternative definition of an affine partition algebra, which they denote by AP_k . They prove many properties for this category and use it to give a new approach to the representation theory of the partition category. However, as they note in [2, Remark 4.12] they have not attempted to give a basis for the morphism spaces in their category, or to give a presentation for it. Inspired by their work, we have explored the connection between our affine partition algebra and the Heisenberg category. We have added a section at the end of our paper where we construct a surjective homomorphism from $\mathcal{A}_{2k}^{\mathrm{aff}}$ to an endomorphism algebra in the Heisenberg category. Our argument generalises to show that the affine partition category APar of Brundan and Vargas is in fact the full monoidal subcategory of the Heisenberg category generated by one object. Using work of Khovanov [12], this gives a basis for all morphism spaces in APar and hence also for AP_k . We also obtain as a corollary that AP_k is a quotient of $\mathcal{A}_{2k}^{\text{aff}}$. We do not know whether these two algebras are in fact isomorphic. If they were, then our definition of $\mathcal{A}_{2k}^{\text{aff}}$ would also give a presentation for AP_k .

The paper is structured as follows. Section 2 deals with the ordinary partition algebra. In Section 2.1 we recall the diagram basis and the original presentation of the partition algebra given by T. Halverson and A. Ram. In Section 2.2, we recall the definition of the Jucys-Murphy elements and the more recent presentation of the algebra given by J. Enyang. Section 2.3 introduces a new normalisation of the Jucys-Murphy elements and of Enyang's generators which has the advantage of simplifying many of the relations in our definition. We also collect many of the relations which will be needed in defining an affine partition. Finally, in Section 2.4, we recall explicitly the Schur-Weyl duality between the partition algebra and the symmetric group.

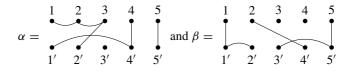
Section 3 gives the definition of the affine partition algebra $\mathcal{A}_{2k}^{\text{aff}}$ in terms of generators and relations and proves some properties. In particular, we show in Section 3.1 that the ordinary partition algebra appears both as a subalgebra and as a quotient of $\mathcal{A}_{2k}^{\text{aff}}$. In Section 3.2, we describe a family of central elements in $\mathcal{A}_{2k}^{\text{aff}}$ and formulate a conjecture about its centre. Finally, in Section 3.3 we show that $\mathcal{A}_{2k}^{\text{aff}}$ extends the action of the partition algebra on tensor space as desired. Section 4 deals with the connections with the Heisenberg category and the work of J. Brundan and M. Vargas on their affine partition category. In Section 4.1, we recall the definition of the Heisenberg category including the basis of the morphism spaces given by M. Khovanov in [12]. In Section 4.2 we define a homomorphism from \mathcal{A}_{2k}^{aff} to the endomorphism space of a particular object in the Heisenberg category and prove that it is surjective. In Section 4.3, we generalise the arguments from Section 4.2 to show that APar is the full monoidal subcategory of the Heisenberg category generated by one object and deduce that AP_k is a quotient of \mathcal{A}_{2k}^{aff} .

2 Partition Algebra

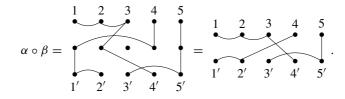
2.1 Diagrammatics and Presentation

For this section we give the definition of the partition algebra $\mathcal{A}_{2k}(z)$ and its presentation established in [8] (and independently in [5]). For $k \in \mathbb{N}$, we let $[k] := \{1, 2, ..., k\}$, and $[k'] := \{1', 2', ..., k'\}$. We view $[k] \cup [k']$ as a formal set on 2k elements, and let Π_{2k} denote the set of all set partitions of $[k] \cup [k']$. Given any $\alpha \in \Pi_{2k}$, we say a partition diagram of α is any graph with vertex set $[k] \cup [k']$ whose connected components partition the vertices according to the blocks of α . We do not distinguish between α and any partition diagram of α , in particular we will only care about the connected components of such graphs, not the particular edges which form the components. When drawing such a diagram, we will arrange the vertices in two rows with the top row going from 1 to k, and the bottom row from 1' to k'. For example, in Π_{10} we have the identification

We define a product \circ on Π_{2k} as follows: Given $\alpha, \beta \in \Pi_{2k}$, we let $\alpha \circ \beta \in \Pi_{2k}$ be the set partition obtained by first stacking the diagram of α on top of that of β , identifying the bottom row of α with the top row of β , removing any connected components lying entirely within the middle row, and then reading off the connected components formed between the top row of α and the bottom row of β . For example consider



in Π_{10} . Then we have



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Clearly this product is associative and independent of the choice of graphs used to represent the set partitions. The element $1 = \{\{i, i'\} \mid i \in [k]\} \in \Pi_{2k}$ is an identity element, and thus (Π_{2k}, \circ) is in fact a monoid. Given $\alpha, \beta \in \Pi_{2k}$, we let $m(\alpha, \beta)$ denote the number of middle components removed in evaluating $\alpha \circ \beta$. In the above example, we have $m(\alpha, \beta) = 1$. Now let z be a formal variable and $\mathbb{C}[z]$ the polynomial ring. The partition algebra $\mathcal{A}_{2k}(z)$ is the $\mathbb{C}[z]$ -algebra whose basis as a free $\mathbb{C}[z]$ -module is given by the set Π_{2k} , and whose product is given by

$$\alpha\beta := z^{m(\alpha,\beta)}\alpha \circ \beta$$

for all $\alpha, \beta \in \Pi_{2k}$, extended linearly over $\mathbb{C}[z]$.

For $0 \le l \le k$, we identify Π_{2l} as a submonoid of Π_{2k} given diagrammatically by

for any $\alpha \in \Pi_{2l}$. Define Π_{2k-1} to be the submonoid of Π_{2k} consisting of all set partitions of $[k] \cup [k']$ where k and k' belong to the same block. We have a chain of monoids $\emptyset = \Pi_0 \subset \Pi_1 \subset \Pi_2 \subset \ldots$. For any $0 \le r \le 2k$, we let $\mathcal{A}_r(z)$ denote the subalgebra of $\mathcal{A}_{2k}(z)$ generated by Π_r . We obtain an analogous chain

$$\mathbb{C}[z] = \mathcal{A}_0(z) \subseteq \mathcal{A}_1(z) \subset \mathcal{A}_2(z) \subset \dots$$

of $\mathbb{C}[z]$ -algebras. The rank of $\mathcal{A}_r(z)$ over $\mathbb{C}[z]$ is $|\Pi_r| = B_r$, where B_r is the r^{th} Bell number. We can view $\mathcal{A}_r(z)$ as an infinite dimensional algebra over \mathbb{C} with basis $\{z^n \alpha \mid n \in \mathbb{Z}_{\geq 0}, \alpha \in \Pi_r\}$. When we do so, we use the notation \mathcal{A}_r instead. For any $\delta \in \mathbb{C}$, let $(z - \delta)$ denote the ideal of \mathcal{A}_r generated by $z - \delta$. Then we let $\mathcal{A}_r(\delta) := \mathcal{A}_r/(z - \delta)$, which is a finite dimensional \mathbb{C} -algebra with dim $(\mathcal{A}_r(\delta)) = B_r$.

For $i \in [k-1]$ and $j \in [k]$, we define the following elements of Π_{2k} :

$$s_{i} = \underbrace{1}_{1'} \cdots \underbrace{1}_{i' \quad (i+1)'} \underbrace{1}_{k'} \cdots \underbrace{1}_{k'}, e_{2j-1} = \underbrace{1}_{1'} \cdots \underbrace{1}_{j'} \underbrace{1}_{k'} \cdots \underbrace{1}_{k'}_{k'}$$
$$e_{2i} = \underbrace{1}_{1'} \cdots \underbrace{1}_{i' \quad (i+1)'} \underbrace{1}_{i' \quad (i+1)'} \underbrace{1}_{k'} \cdots \underbrace{1}_{k'}^{k}.$$

These elements generate the monoid (Π_{2k}, \circ) , and in turn the algebra $\mathcal{A}_{2k}(z)$. Moreover, a presentation in terms of these generators, which we display below, was given in [8, Theorem 1.11], see also *Theorem 36* and *Section 6.3* of [5].

Theorem 2.1.1 The partition algebra $A_{2k}(z)$ has a presentation with generating set

$$\{s_i, e_j \mid i \in [k-1], j \in [2k-1]\}$$

and relations

(HR1) (Coxeter relations)

(i) $s_i^2 = 1$, for $i \in [k-1]$.

- (ii) $s_i s_j = s_j s_i$, for $j \neq i + 1$.
- (iii) $s_i s_{i+1} s_i = s_{i+1} s_i s_{i+1}$, for $i \in [k-2]$.
- (HR2) (Idempotent relations)
 - (i) $e_{2i-1}^2 = ze_{2i-1}$, for $i \in [k]$. (ii) $e_{2i}^2 = e_{2i}$, for $i \in [k-1]$.

 - (iii) $s_i e_{2i} = e_{2i} s_i = e_{2i}$, for $i \in [k-1]$.
 - (iv) $s_i e_{2i-1} e_{2i+1} = e_{2i-1} e_{2i+1} s_i = e_{2i-1} e_{2i+1}$, for $i \in [k-1]$.

(HR3) (Commutation relations)

- (i) $e_{2i-1}e_{2i-1} = e_{2i-1}e_{2i-1}$, for $i, j \in [k]$.
- (ii) $e_{2i}e_{2j} = e_{2j}e_{2i}$, for $i, j \in [k-1]$.
- (iii) $e_{2i-1}e_{2i} = e_{2i}e_{2i-1}$, for $j \neq i-1, i$.
- (iv) $s_i e_{2i-1} = e_{2i-1}s_i$, for $j \neq i, i+1$.
- (v) $s_i e_{2i} = e_{2i} s_i$, for $j \neq i 1, i + 1$.
- (vi) $s_i e_{2i-1} s_i = e_{2i+1}$, for $i \in [k-1]$.
- (vii) $s_i e_{2i-2} s_i = s_{i-1} e_{2i} s_{i-1}$, for $i \in [k-1]$.

(HR4) (Contraction relations)

- (i) $e_i e_{i+1} e_i = e_i$ for $i \in [2k-2]$.
- (ii) $e_{i+1}e_ie_{i+1} = e_{i+1}$, for $i \in [2k-2]$.

The presentation above extends to one for the \mathbb{C} -algebra \mathcal{A}_{2k} by simply adding z as a central generator. The \mathbb{C} -algebra $\mathcal{A}_{2k}(\delta)$ has a presentation identical to above with the exception of replacing z with δ in relation (*HR2*)(i). From the symmetry of the above presentation, one can deduce that we have an anti-involution $*: \mathcal{A}_{2k}(z) \to \mathcal{A}_{2k}(z)$ given by flipping a partition diagram up-side-down, and extending linearly over $\mathbb{C}[z]$. We denote the image of an element $a \in A_{2k}(z)$ under this anti-involution by a^* .

2.2 Jucys-Murphy Elements and Enyang's Presentation

In this section we give the definition of the Jucys-Murphy elements of the partition algebra. These elements were originally defined diagrammatically by Halverson and Ram in [8]. They were later given a recursive definition by Enyang in [6]. For this recursive definition, Envang introduced new elements σ_i which resemble the Coxeter generators s_i . We recall this recursive definition, and a new presentation of the partition algebra given in [6] in terms of the generators e_i and σ_i . The following definition is the one given in Section 2.3 of [7].

Definition 2.2.1 Let $L_1 = 0$, $L_2 = e_1$, $\sigma_2 = 1$, and $\sigma_3 = s_1$. Then for i = 1, 2, ..., define

$$L_{2i+2} = s_i L_{2i} s_i - s_i L_{2i} e_{2i} - e_{2i} L_{2i} s_i + e_{2i} L_{2i} e_{2i+1} e_{2i} + \sigma_{2i+1},$$

where, for $i = 2, 3, \ldots$, we have

$$\sigma_{2i+1} = s_{i-1}s_i\sigma_{2i-1}s_is_{i-1} + s_ie_{2i-2}L_{2i-2}s_ie_{2i-2}s_i + e_{2i-2}L_{2i-2}s_ie_{2i-2} -s_ie_{2i-2}L_{2i-2}s_{i-1}e_{2i}e_{2i-1}e_{2i-2} - e_{2i-2}e_{2i-1}e_{2i}s_{i-1}L_{2i-2}e_{2i-2}s_i.$$

Also for $i = 1, 2, \ldots$, define

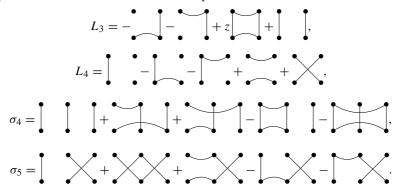
$$L_{2i+1} = s_i L_{2i-1} s_i - L_{2i} e_{2i} - e_{2i} L_{2i} + (z - L_{2i-1}) e_{2i} + \sigma_{2i},$$

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where, for $i = 2, 3, \ldots$, we have

$$\sigma_{2i} = s_{i-1}s_i\sigma_{2i-2}s_is_{i-1} + e_{2i-2}L_{2i-2}s_ie_{2i-2}s_i + s_ie_{2i-2}L_{2i-2}s_ie_{2i-2}$$
$$-e_{2i-2}L_{2i-2}s_{i-1}e_{2i}e_{2i-1}e_{2i-2} - s_ie_{2i-2}e_{2i-1}e_{2i}s_{i-1}L_{2i-2}e_{2i-2}s_i.$$

Example 2.2.2 The first few non-trivial examples are



We will refer to the elements L_i as the JM-elements, and the elements σ_i as Enyang's generators. A simple proof by induction tells us the following:

Lemma 2.2.3 For each $i \in \mathbb{N}$ we have that $L_i \in \mathcal{A}_i(z)$ and $\sigma_i \in \mathcal{A}_{i+1}(z)$.

It was shown in [6] that these elements are invariant under the anti-automorphism *. They also commute with smaller partition algebras with respect to the chain described in the previous section: For any $i \le r$, let $Z_i(\mathcal{A}_r(z)) := \langle a \in \mathcal{A}_r(z) | ab = ba, \forall b \in \mathcal{A}_i(z) \rangle$, then it was shown in [6, Theorem 3.8] that

$$L_i, \sigma_{i+1} \in Z_{i-1}(\mathcal{A}_i(z)). \tag{1}$$

In particular this shows that the JM-elements pairwise commute. We now give the new presentation of $A_{2k}(z)$ established in [6]. This presentation is given in terms of the generators e_i and Enyang's generator's σ_i . Remarkably, although the definition and diagrammatic description of the σ_i is rather complicated, the defining relations in the following presentation are very simple. This is less surprising when one considers how these elements act on tensor space. This will be discussed in Section 2.4.

Theorem 2.2.4 [6, Theorem 4.1] *The partition algebra* $A_{2k}(z)$ *has a presentation with generating set*

$$\{\sigma_i, e_j \mid 3 \le i \le 2k - 1, \ j \in [2k - 1]\}$$

and relations:

(i)
$$\sigma_{2i+2}^2 = 1$$
 for $i \in [k-2]$.

(ii)
$$\sigma_{2i+1}^2 = 1$$
 for $i \in [k-1]$.

(E2) (Braid relations)

(i)
$$\sigma_{2i+1}\sigma_{2i} = \sigma_{2i}\sigma_{2i+1}$$
 for $j \neq i+1$.

- (ii) $\sigma_{2i+1}\sigma_{2i+1} = \sigma_{2i+1}\sigma_{2i+1}$ for $j \neq i \pm 1$.
- (iii) $\sigma_{2i}\sigma_{2j} = \sigma_{2j}\sigma_{2i}$ for $j \neq i \pm 1$.

(iv)
$$s_i s_{i+1} s_i = s_{i+1} s_i s_{i+1}$$
, for $i \in [k-2]$, where $s_j = \begin{cases} \sigma_3, & j = 1 \\ \sigma_{2j} \sigma_{2j+1}, & j > 1 \end{cases}$

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(E3) (Idempotent relations)

- (i) $e_{2i-1}^2 = ze_{2i-1}$ for $i \in [k]$.
- (ii) $e_{2i}^2 = e_{2i}$ for $2 \ge i \le k 1$.
- (iii) $\sigma_{2i+1}e_{2i} = e_{2i}\sigma_{2i+1} = e_{2i}$ for $i \in [k-1]$.
- (iv) $\sigma_{2i}e_{2i} = e_{2i}\sigma_{2i} = e_{2i}$ for $2 \le i \le k 1$.
- (v) $\sigma_{2i}e_{2i-1}e_{2i+1} = \sigma_{2i+1}e_{2i-1}e_{2i+1}$ for $2 \le i \le k-1$.
- (vi) $e_{2i+1}e_{2i-1}\sigma_{2i} = e_{2i+1}e_{2i-1}\sigma_{2i+1}$ for $2 \le i \le k-1$.

(E4) (Commutation relations)

- (i) $e_i e_j = e_j e_i$, if $|i j| \ge 2$.
- (ii) $\sigma_{2i-1}e_{2j-1} = e_{2j-1}\sigma_{2i-1}$, if $j \neq i-1, i$.
- (iii) $\sigma_{2i-1}e_{2i} = e_{2i}\sigma_{2i-1}$, if $j \neq i$.
- (iv) $\sigma_{2i}e_{2i-1} = e_{2i-1}\sigma_{2i}$, if $j \neq i, i+1$.
- (v) $\sigma_{2i}e_{2j} = e_{2j}\sigma_{2i}$, if $j \neq i 1$.

(E5) (Contractions)

- (i) $e_i e_{i+1} e_i = e_i$ and $e_{i+1} e_i e_{i+1} = e_{i+1}$, for $i \in [2k-2]$.
- (ii) $\sigma_{2i}e_{2i-1}\sigma_{2i} = \sigma_{2i+1}e_{2i+1}\sigma_{2i+1}$, for $2 \le i \le k-1$.
- (iii) $\sigma_{2i}e_{2i-2}\sigma_{2i} = \sigma_{2i-1}e_{2i}\sigma_{2i-1}$, for $2 \le i \le k-1$.

Note we only worked with the elements σ_i for $i \ge 3$, since $\sigma_2 = 1$. The elements s_j in the above presentation are precisely the Coxeter generators. From the involution relations we have that $s_i \sigma_{2i} = \sigma_{2i} s_i = \sigma_{2i+1}$. From (1) one can deduce that L_i and σ_j commute whenever $j \ne i - 1, i, i + 1$. We end this section by giving relations which tell us how the JM-elements interact with Enyang's generators when they do not commute. We use results established in [6], although we have adopted the notation of [7].

Remark 2.2.5 The change of notation between [6] and [7] is given respectively by $p_i \sim e_{2i-1}$, $p_{i+\frac{1}{2}} \sim e_{2i}$, $\sigma_i \sim \sigma_{2i-1}$, $\sigma_{i+\frac{1}{2}} \sim \sigma_{2i}$, $L_i \sim L_{2i}$, and $L_{i+\frac{1}{2}} \sim L_{2i+1}$.

Proposition 2.2.6 The following relations hold:

(i) $L_{2i+1} = \sigma_{2i}L_{2i-1}\sigma_{2i} - e_{2i}e_{2i-1}\sigma_{2i} - \sigma_{2i}e_{2i-1}e_{2i} + e_{2i}e_{2i+1}\sigma_{2i}e_{2i+1}e_{2i} + \sigma_{2i}$

(ii) $L_{2i+2} = \sigma_{2i+1}L_{2i}\sigma_{2i+1} - e_{2i}e_{2i+1} - e_{2i+1}e_{2i} + e_{2i}e_{2i+1}\sigma_{2i+1}e_{2i+1}e_{2i} + \sigma_{2i+1}$.

(iii) $L_{2i} = \sigma_{2i}L_{2i}\sigma_{2i} + e_{2i}e_{2i-1}\sigma_{2i} + \sigma_{2i}e_{2i-1}e_{2i} - e_{2i}e_{2i+1} - e_{2i+1}e_{2i}$.

(iv)
$$L_{2i+1} = \sigma_{2i+1}L_{2i+1}\sigma_{2i+1} - e_{2i}e_{2i+1}\sigma_{2i+1} - \sigma_{2i+1}e_{2i+1}e_{2i} + e_{2i}e_{2i+1} + e_{2i+1}e_{2i}$$

Proof (*i*): By definition,

$$L_{2i+1} = s_i L_{2i-1} s_i - L_{2i} e_{2i} - e_{2i} L_{2i} + (z - L_{2i-1}) e_{2i} + \sigma_{2i}.$$
 (2)

We examine the right hand side term by term. For the first term we have

$$s_i L_{2i-1} s_i = \sigma_{2i} \sigma_{2i+1} L_{2i-1} \sigma_{2i+1} \sigma_{2i} = \sigma_{2i} \sigma_{2i+1}^2 L_{2i-1} \sigma_{2i} = \sigma_{2i} L_{2i-1} \sigma_{2i}.$$

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For the second term, multiplying [6, Proposition 3.2 (3)] on the left by s_i we get $\sigma_{2i}e_{2i-1}e_{2i} = L_{2i}e_{2i}$. Acting by the anti-automorphism * yields $e_{2i}e_{2i-1}\sigma_{2i} = e_{2i}L_{2i}$ for the third term. Lastly for the forth term

$$(z - L_{2i-1})e_{2i} = (z - L_{2i-1})e_{2i}e_{2i+1}e_{2i}$$

= $e_{2i}(z - L_{2i-1})e_{2i+1}e_{2i}$
= $e_{2i}e_{2i+1}\sigma_{2i}e_{2i+1}e_{2i}$

where the last equality follows by [6, Proposition 4.3 (2)]. Substituting these terms back into (2) yields (i).

(ii): By definition,

$$L_{2i+2} = s_i L_{2i} s_i - s_i L_{2i} e_{2i} - e_{2i} L_{2i} s_i + e_{2i} L_{2i} e_{2i+1} e_{2i} + \sigma_{2i+1}.$$
 (3)

Multiplying this equation on the left and right by σ_{2i} gives

$$L_{2i+2} = \sigma_{2i+1}L_{2i}\sigma_{2i+1} - \sigma_{2i+1}L_{2i}e_{2i} - e_{2i}L_{2i}\sigma_{2i+1} + e_{2i}L_{2i}e_{2i+1}e_{2i} + \sigma_{2i+1}$$

= $\sigma_{2i+1}L_{2i}\sigma_{2i+1} - \sigma_{2i+1}^2e_{2i+1}e_{2i} - e_{2i}e_{2i+1}\sigma_{2i+1}^2 + e_{2i}e_{2i+1}\sigma_{2i+1}e_{2i} + e_{2i}e_{2i+1}e_{2i} + \sigma_{2i+1}e_{2i} + \sigma_$

which gives (*ii*), where the second equality follows by relation $\sigma_{2i+1}e_{2i+1}e_{2i} = L_{2i}e_{2i}$ and its dual $e_{2i}e_{2i+1}\sigma_{2i+1} = e_{2i}L_{2i}$, which follows from [6, Proposition 3.2 (3)].

(*iii*): It was shown in [6, Proposition 3.10] that the element $L_1 + L_2 + \cdots + L_r$ belongs to the center of $\mathcal{A}_r(z)$. From this, and the fact that L_i and σ_j commute whenever $j \neq i-1, i, i+1$, one may deduce that

$$\sigma_{2i}(L_{2i-1} + L_{2i} + L_{2i+1})\sigma_{2i} = L_{2i-1} + L_{2i} + L_{2i+1}$$

Rearranging gives

$$L_{2i} = \sigma_{2i}L_{2i}\sigma_{2i} + (\sigma_{2i}L_{2i-1}\sigma_{2i} - L_{2i+1}) + (\sigma_{2i}L_{2i+1}\sigma_{2i} - L_{2i-1}).$$
(4)

We examine the bracketed terms in (4). Rearranging (i) gives the first bracketed term as

$$\sigma_{2i}L_{2i-1}\sigma_{2i} - L_{2i+1} = e_{2i}e_{2i-1}\sigma_{2i} + \sigma_{2i}e_{2i-1}e_{2i} - e_{2i}e_{2i+1}\sigma_{2i}e_{2i+1}e_{2i} - \sigma_{2i}.$$

Multiplying this on the left and right by σ_{2i} , and then rearranging gives the second bracketed term

$$\sigma_{2i}L_{2i+1}\sigma_{2i} - L_{2i-1} = -e_{2i}e_{2i-1} - e_{2i-1}e_{2i} + e_{2i}e_{2i+1}\sigma_{2i}e_{2i+1}e_{2i} + \sigma_{2i}.$$

Substituting these back into (4) yields (iii).

(iv): Analogously to the previous case, one can deduce that

$$\sigma_{2i+1}(L_{2i} + L_{2i+1} + L_{2i+2})\sigma_{2i+1} = L_{2i} + L_{2i+1} + L_{2i+2}.$$

Rearranging gives

$$L_{2i+1} = \sigma_{2i+1}L_{2i+1}\sigma_{2i+1} + (\sigma_{2i+1}L_{2i}\sigma_{2i+1} - L_{2i+2}) + (\sigma_{2i+1}L_{2i+2}\sigma_{2i+1} - L_{2i}).$$
(5)

We examine the bracketed terms in (5). Rearranging (2)(ii) gives the first bracketed term as

$$\sigma_{2i+1}L_{2i}\sigma_{2i+1} - L_{2i+2} = e_{2i}e_{2i+1} + e_{2i+1}e_{2i} - e_{2i}e_{2i+1}\sigma_{2i+1}e_{2i+1}e_{2i} - \sigma_{2i+1}.$$

Multiplying this on the left and right by σ_{2i+1} , and then rearranging gives the second bracketed term

$$\sigma_{2i+1}L_{2i+2}\sigma_{2i+1} - L_{2i} = -e_{2i}e_{2i+1}\sigma_{2i+1} - \sigma_{2i+1}e_{2i+1}e_{2i} + e_{2i}e_{2i+1}\sigma_{2i+1}e_{2i+1}e_{2i} + \sigma_{2i+1}.$$

Substituting these back into (5) yields (*iv*).

2.3 Normalisation

As mentioned in the introduction, we seek to 'affinize' the partition algebra by replacing the Jucys-Murphy elements with commuting variables, and asking them to retain various relations with the generators. In preparation for this construction, this section collects all the relations we seek to retain in one place. However, instead of working with the JM-elements and Enyang's generators, it turns out to be easier to work with the following elements: For each $i \in \mathbb{N}$ we set

$$t_{2i} := \sigma_{2i} - e_{2i}, t_{2i+1} := \sigma_{2i+1} - e_{2i}.$$

For each $i \in \mathbb{N}$ we set

$$X_i := \begin{cases} z - 1 - L_i, & \text{if } i \text{ odd} \\ L_i - 1, & \text{if } i \text{ even} \end{cases}$$

We also call the elements X_i the JM-elements and the elements t_i Enyang's generators. By definitions we have that $t_i \in A_{i+1}(z)$, $X_i \in A_i(z)$, and that $t_i^* = t_i$ and $X_i^* = X_i$. One can also deduce that $s_i t_{2i} = t_{2i} s_i = t_{2i+1}$. We briefly collect some simple relations to ease the proof of the proceeding proposition.

Lemma 2.3.1 The following relations hold:

(i) $e_{2i+1}t_{2i}e_{2i+1} = X_{2i-1}e_{2i+1}$

(ii) $t_{2i}e_{2i-1}e_{2i} = X_{2i}e_{2i}$, and $e_{2i}e_{2i-1}t_{2i} = e_{2i}X_{2i}$

(iii) $t_{2i+1}e_{2i+1}e_{2i} = X_{2i}e_{2i}$, and $e_{2i}e_{2i+1}t_{2i+1} = e_{2i}X_{2i}$

Proof (*i*): We have that

$$e_{2i+1}t_{2i}e_{2i+1} = e_{2i+1}(\sigma_{2i} - e_{2i})e_{2i+1}$$

= $e_{2i+1}\sigma_{2i}e_{2i+1} - e_{2i+1}$ by (E5)
= $(z - L_{2i-1})e_{2i+1} - e_{2i+1}$ by [6 Proposition 4.3 (2)]
= $(X_{2i-1} + 1)e_{2i+1} - e_{2i+1}$
= $7X_{2i-1}e_{2i+1}$

(ii): We have that

$$t_{2i}e_{2i-1}e_{2i} = (\sigma_{2i} - e_{2i})e_{2i-1}e_{2i}$$

= $\sigma_{2i}e_{2i-1}e_{2i} - e_{2i}$ by (E5)
= $L_{2i}e_{2i} - e_{2i}$ by [6 Proposition 3.2 (3)]
= $(X_{2i} + 1)e_{2i} - e_{2i}$
= $X_{2i}e_{2i}$

The relation $e_{2i}e_{2i-1}t_{2i} = e_{2i}X_{2i}$ is obtained by acting by *.

(*iii*): We have $t_{2i+1}e_{2i+1}e_{2i} = t_{2i}s_ie_{2i+1}e_{2i} = t_{2i}e_{2i-1}e_{2i} = X_{2i}e_{2i}$. Again the relation $e_{2i}e_{2i+1}t_{2i+1} = e_{2i}X_{2i}$ is obtained by acting by *.

The following proposition contains all the relations we seek to retain for our construction of the affine partition algebra, as such some are identical to relations we have already stated. It provides a presentation of the partition algebra $A_{2k}(z)$ which is simply Enyang's presentation 2.2.4 except working with the generators t_i instead of σ_i . For those relations we have adopted the same naming conventions given in 2.2.4. **Proposition 2.3.2** The partition algebra $A_{2k}(z)$ has a presentation with generating set

$$\{t_i, e_j \mid 3 \le i \le 2k - 1, j \in [2k - 1]\}$$

and relations:

- (1) (Involutions)
 - (i) $t_{2i+2}^2 = 1 e_{2i}$, for $i \in [k-2]$.

(ii)
$$t_{2i+1}^2 = 1 - e_{2i}$$
, for $i \in [k-1]$.

(2) (Braid relations)

- (i) $t_{2i+1}t_{2j} = t_{2j}t_{2i+1}$ for $j \neq i+1$.
- (ii) $t_{2i+1}t_{2j+1} = t_{2j+1}t_{2i+1}$ for $j \neq i \pm 1$.
- (iii) $t_{2i}t_{2j} = t_{2j}t_{2i}$ for $j \neq i \pm 1$.

(iv)
$$s_i s_{i+1} s_i = s_{i+1} s_i s_{i+1}$$
, for $i \in [k-2]$, where $s_j = \begin{cases} t_3 + e_2, & j = 1 \\ t_{2j} t_{2j+1} + e_{2j}, & j > 1 \end{cases}$

(3) (Idempotent relations)

- (i) $e_{2i-1}^2 = ze_{2i-1}$ for $i \in [k]$.
- (ii) $e_{2i}^2 = e_{2i}$ for $i \in [k-1]$.
- (iii) $t_{2i+1}e_{2i} = e_{2i}t_{2i+1} = 0$ for $i \in [k-1]$.
- (iv) $t_{2i}e_{2i} = e_{2i}t_{2i} = 0$ for $2 \le i \le k 1$.
- (v) $t_{2i}e_{2i-1}e_{2i+1} = t_{2i+1}e_{2i-1}e_{2i+1}$ for $2 \le i \le k-1$.
- (vi) $e_{2i+1}e_{2i-1}t_{2i} = e_{2i+1}e_{2i-1}t_{2i+1}$ for $2 \le i \le k-1$.

(4) (Commutation relations)

- (i) $e_i e_j = e_j e_i$, if $|i j| \ge 2$.
- (ii) $t_{2i-1}e_{2j-1} = e_{2j-1}t_{2i-1}$, if $j \neq i-1, i$.
- (iii) $t_{2i-1}e_{2j} = e_{2j}t_{2i-1}$, if $j \neq i$.
- (iv) $t_{2i}e_{2j-1} = e_{2j-1}t_{2i}$, if $j \neq i, i+1$.
- (v) $t_{2i}e_{2j} = e_{2j}t_{2i}$, if $j \neq i 1$.

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(5) (Contractions)
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- (i) $e_i e_{i+1} e_i = e_i$ and $e_{i+1} e_i e_{i+1} = e_{i+1}$, for $i \in [2k-2]$.
- (ii) $t_{2i}e_{2i-1}t_{2i} = t_{2i+1}e_{2i+1}t_{2i+1}$, for $i \in [k-1]$.
- (iii) $t_{2i}e_{2i-2}t_{2i} = t_{2i-1}e_{2i}t_{2i-1}$, for $2 \le i \le k-1$.

Furthermore, the following relations are satisfied in $A_{2k}(z)$:

(6) (JM Commutation Relations)

- (i) $X_i X_j = X_j X_i$ for all $i, j \in [2k]$
- (ii) $t_i X_i = X_i t_i$ for $j \neq i 1, i, i + 1$
- (iii) $e_i X_j = X_j e_i$ for $j \neq i, i+1$
- (7) (Braid-like Relations)
 - (i) $t_{2i-2}t_{2i}t_{2i-2} = t_{2i}t_{2i-2}t_{2i}(1-e_{2i-2})$
 - (ii) $t_{2i+1}t_{2i-1}t_{2i+1} = t_{2i-1}t_{2i+1}t_{2i-1}(1-e_{2i})$
 - (iii) $t_{2i-1}t_{2i}t_{2i-1} = t_{2i} e_{2i-2}t_{2i} t_{2i}e_{2i-2}$
 - (iv) $t_{2i}t_{2i-1}t_{2i} = t_{2i-1} e_{2i}t_{2i-1} t_{2i-1}e_{2i}$

(8) (Skein-like Relations)

- (i) $X_{2i+1} = t_{2i}X_{2i-1}t_{2i} + e_{2i}e_{2i-1}t_{2i} + t_{2i}e_{2i-1}e_{2i} t_{2i}$.
- (ii) $X_{2i+2} = t_{2i+1}X_{2i}t_{2i+1} + e_{2i}e_{2i+1}t_{2i+1}e_{2i+1}e_{2i} + t_{2i+1}$.
- (iii) $X_{2i} = t_{2i}X_{2i}t_{2i} + e_{2i}e_{2i-1}t_{2i} + t_{2i}e_{2i-1}e_{2i}$.
- (iv) $X_{2i+1} = t_{2i+1}X_{2i+1}t_{2i+1} + e_{2i}e_{2i+1}t_{2i+1} + t_{2i+1}e_{2i+1}e_{2i}$.

(9) (Anti-symmetry Relations)

- (i) $e_i(X_i X_{i+1}) = 0$ for $i \in [2k 1]$.
- (ii) $(X_i X_{i+1})e_i = 0$ for $i \in [2k 1]$.

(10) (Bubble Relations)

(i) $e_1 X_1^l e_1 = z(z-1)^l e_1$, for all $l \in \mathbb{Z}_{\geq 0}$.

Proof Although lengthy, it is simple to check that relations (1) to (5) give an alternative presentation for $A_{2k}(z)$ since we merely exchanged the elements σ_i with t_i from Enyang's presentation given in Theorem 2.2.4.

(6): Follows from (1) and Lemma 2.2.3.

(7): These relations will be proven in the next section in Lemma 2.4.5.

(9): Follows from [6, Proposition 3.9] (1) and (2).

(10): We have that $X_1 = z - 1 - L_1 = z - 1$. Thus for any $l \in \mathbb{N}$,

$$e_1 X_1^l e_1 = (z-1)^l e_1^2 = z(z-1)^l e_1.$$

(8)(i): From Proposition 2.2.6 (i) we have

$$L_{2i+1} = \sigma_{2i}L_{2i-1}\sigma_{2i} - e_{2i}e_{2i-1}\sigma_{2i} - \sigma_{2i}e_{2i-1}e_{2i} + e_{2i}e_{2i+1}\sigma_{2i}e_{2i+1}e_{2i} + \sigma_{2i}.$$
 (6)

Examining the right hand side term by term: For the first term,

$$\sigma_{2i}L_{2i-1}\sigma_{2i} = (t_{2i} + e_{2i})(-X_{2i-1})(t_{2i} + e_{2i}) + (z-1)$$

= $-t_{2i}X_{2i-1}t_{2i} - t_{2i}X_{2i-1}e_{2i} - e_{2i}X_{2i-1}t_{2i} - e_{2i}X_{2i-1}e_{2i} + (z-1)$
= $-t_{2i}X_{2i-1}t_{2i} - X_{2i-1}e_{2i} + (z-1)$

where the last equality follows since X_{2i-1} commutes with e_{2i} and $t_{2i}e_{2i} = e_{2i}t_{2i} = 0$. For the second and third term of (6), we have

$$-e_{2i}e_{2i-1}\sigma_{2i} = -e_{2i} - e_{2i}e_{2i-1}t_{2i}$$
, and $-\sigma_{2i}e_{2i-1}e_{2i} = -t_{2i}e_{2i-1}e_{2i} - e_{2i}$.

For the forth term of (6),

$$e_{2i}e_{2i+1}\sigma_{2i}e_{2i+1}e_{2i} = e_{2i}e_{2i+1}t_{2i}e_{2i+1}e_{2i} + e_{2i}e_{2i+1}e_{2i}e_{2i+1}e_{2i}$$

$$= e_{2i}e_{2i+1}t_{2i}e_{2i+1}e_{2i} + e_{2i}$$

$$= e_{2i}e_{2i-1}t_{2i+1}e_{2i+1}e_{2i} + e_{2i} \text{ by } t_{2i} = s_it_{2i+1}$$

$$= e_{2i}e_{2i-1}X_{2i+1}e_{2i} + e_{2i} \text{ by Lemma 2.3.1 (iii)}$$

$$= e_{2i}e_{2i-1}e_{2i}X_{2i-1} + e_{2i} \text{ by } (9)(i), (ii)$$

$$= e_{2i}X_{2i-1} + e_{2i}.$$

Substituting all these back into (6) yields

$$z - 1 - X_{2i+1} = -t_{2i}X_{2i-1}t_{2i} - X_{2i-1}e_{2i} + (z - 1)$$

$$-e_{2i} - e_{2i}e_{2i-1}t_{2i} - t_{2i}e_{2i-1}e_{2i} - e_{2i}$$

$$+e_{2i}X_{2i-1} + e_{2i} + t_{2i} + e_{2i}$$

$$\iff X_{2i+1} = t_{2i}X_{2i-1}t_{2i} + e_{2i}e_{2i-1}t_{2i} + t_{2i}e_{2i-1}e_{2i} - t_{2i}$$

giving (8)(i).

(8)(ii): From Proposition 2.2.6 (ii) we have

$$L_{2i+2} = \sigma_{2i+1}L_{2i}\sigma_{2i+1} - e_{2i}e_{2i+1} - e_{2i+1}e_{2i} + e_{2i}e_{2i+1}\sigma_{2i+1}e_{2i+1}e_{2i} + \sigma_{2i+1}.$$
 (7)

We examine two terms on the right hand side: The first term gives

$$\sigma_{2i+1}L_{2i}\sigma_{2i+1} = (t_{2i+1} + e_{2i})(X_{2i} + 1)(t_{2i+1} + e_{2i})$$

= $t_{2i+1}X_{2i}t_{2i+1} + t_{2i+1}X_{2i}e_{2i} + e_{2i}X_{2i}t_{2i+1} + e_{2i}X_{2i}e_{2i} + 1$
= $t_{2i+1}X_{2i}t_{2i+1} + t_{2i+1}^2e_{2i+1}e_{2i} + e_{2i}e_{2i+1}t_{2i+1}^2 + 1$
= $t_{2i+1}X_{2i}t_{2i+1} + e_{2i+1}e_{2i} + e_{2i}e_{2i+1} - 2e_{2i} + 1$

where the second equality follows since $(t_{2i+1} + e_{2i})^2 = 1$, and the third from Lemma 2.3.1 (iii) and since $e_{2i}X_{2i}e_{2i} = e_{2i}e_{2i-1}t_{2i}e_{2i} = 0$. The forth term in (7) gives

$$e_{2i}e_{2i+1}\sigma_{2i+1}e_{2i+1}e_{2i} = e_{2i}e_{2i+1}t_{2i+1}e_{2i+1}e_{2i} + e_{2i}e_{2i+1}e_{2i}e_{2i+1}e_{2i}$$

= $e_{2i}e_{2i+1}t_{2i+1}e_{2i+1}e_{2i} + e_{2i}$.

Substituting these back into (7) yields

$$X_{2i+2} + 1 = t_{2i+1}X_{2i}t_{2i+1} + e_{2i+1}e_{2i} + e_{2i}e_{2i+1} - 2e_{2i} + 1$$

$$-e_{2i}e_{2i+1} - e_{2i+1}e_{2i} + e_{2i}e_{2i+1}t_{2i+1}e_{2i+1}e_{2i} + e_{2i} + t_{2i+1} + e_{2i}$$

$$\iff X_{2i+2} = t_{2i+1}X_{2i}t_{2i+1} + e_{2i}e_{2i+1}t_{2i+1}e_{2i+1}e_{2i} + t_{2i+1}$$

giving (8)(ii).

(8)(iii): From Proposition 2.2.6 (iii) we have

$$L_{2i} = \sigma_{2i}L_{2i}\sigma_{2i} + e_{2i}e_{2i-1}\sigma_{2i} + \sigma_{2i}e_{2i-1}e_{2i} - e_{2i}e_{2i+1} - e_{2i+1}e_{2i}.$$
 (8)

We have that

$$\sigma_{2i} L_{2i} \sigma_{2i} = (t_{2i} + e_{2i})(X_{2i} + 1)(t_{2i} + e_{2i})$$

= $t_{2i} X_{2i} t_{2i} + t_{2i} X_{2i} e_{2i} + e_{2i} X_{2i} t_{2i} + e_{2i} X_{2i} e_{2i} + 1$
= $t_{2i} X_{2i} t_{2i} + t_{2i}^2 e_{2i-1} e_{2i} + e_{2i} e_{2i-1} t_{2i}^2 + 1$
= $t_{2i} X_{2i} t_{2i} + e_{2i-1} e_{2i} + e_{2i} e_{2i-1} - 2e_{2i} + 1$

where the second equality follows since $(t_{2i} + e_{2i})^2 = 1$, and the third equality from Lemma 2.3.1 (ii) and the since $t_{2i}e_{2i} = e_{2i}t_{2i} = 0$. Substituting this, and relations

$$e_{2i}e_{2i-1}\sigma_{2i} = e_{2i}e_{2i-1}t_{2i} + e_{2i}$$
 and $\sigma_{2i}e_{2i-1}e_{2i} = t_{2i}e_{2i-1}e_{2i} + e_{2i}$,

back into (8) yields

$$X_{2i} + 1 = t_{2i}X_{2i}t_{2i} + e_{2i-1}e_{2i} + e_{2i}e_{2i-1} - 2e_{2i} + 1 + e_{2i}e_{2i-1}t_{2i} + e_{2i}$$

+ $t_{2i}e_{2i-1}e_{2i} + e_{2i} - e_{2i}e_{2i+1} - e_{2i+1}e_{2i}$
 $\iff X_{2i} = t_{2i}X_{2i}t_{2i} + e_{2i}e_{2i-1}t_{2i} + t_{2i}e_{2i-1}e_{2i} + 1$

giving (8)(iii).

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(8)(iv): From Proposition 2.2.6 (iv) we have

$$L_{2i+1} = \sigma_{2i+1}L_{2i+1}\sigma_{2i+1} - e_{2i}e_{2i+1}\sigma_{2i+1} - \sigma_{2i+1}e_{2i+1}e_{2i} + e_{2i}e_{2i+1} + e_{2i+1}e_{2i}.$$
 (9)

We have that

$$\sigma_{2i+1}L_{2i+1}\sigma_{2i+1} = (t_{2i+1} + e_{2i})(-X_{2i+1})(t_{2i+1} + e_{2i}) + (z-1)$$

$$= -t_{2i+1}X_{2i+1}t_{2i+1} - t_{2i+1}X_{2i+1}e_{2i}$$

$$-e_{2i}X_{2i+1}t_{2i+1} - e_{2i}X_{2i+1}e_{2i} + (z-1)$$

$$= -t_{2i+1}X_{2i+1}t_{2i+1} - t_{2i+1}^2e_{2i+1}e_{2i} - e_{2i}e_{2i+1}t_{2i+1}^2 + (z-1)$$

$$= -t_{2i+1}X_{2i+1}t_{2i+1} - e_{2i+1}e_{2i} - e_{2i}e_{2i+1} + 2e_{2i} + (z-1)$$

where the third equality follows from Lemma 2.3.1 (iii), and noting that $e_{2i}X_{2i+1}e_{2i} = e_{2i}X_{2i}e_{2i} = e_{2i}e_{2i-1}t_{2i}e_{2i} = 0$. Substituting this, and the equations

$$-e_{2i}e_{2i+1}\sigma_{2i+1} = -e_{2i}e_{2i+1}t_{2i+1} - e_{2i}$$
 and $\sigma_{2i}e_{2i-1}e_{2i} = -t_{2i+1}e_{2i+1}e_{2i} - e_{2i}$,

back into (9) yields

$$\begin{aligned} (z-1) - X_{2i+1} &= -t_{2i+1}X_{2i+1}t_{2i+1} - e_{2i+1}e_{2i} - e_{2i}e_{2i+1} \\ &+ 2e_{2i} + (z-1) - e_{2i}e_{2i+1}t_{2i+1} \\ &- e_{2i} - t_{2i+1}e_{2i+1}e_{2i} - e_{2i} + e_{2i}e_{2i+1} + e_{2i+1}e_{2i} \\ &\iff X_{2i+1} &= t_{2i+1}X_{2i+1}t_{2i+1} + t_{2i+1}e_{2i+1}e_{2i} + e_{2i}e_{2i+1}t_{2i+1} \end{aligned}$$

giving (8)(iv).

2.4 Schur-Weyl Duality

In this section we recall the Schur-Weyl duality between the partition algebra $A_{2k}(n)$ and the group algebra of the symmetric group $\mathbb{C}S(n)$ via their actions on tensor space. We will also highlight how X_i and t_i act on this tensor space, and complete the proof of Proposition 2.3.2. Consider the permutation module $V = \text{Span}_{\mathbb{C}}\{v_1, \ldots, v_n\}$ of $\mathbb{C}S(n)$ with action given by $\pi v_a = v_{\pi(a)}$ for all $\pi \in S(n)$ and $a \in [n]$, extended \mathbb{C} -linearly to $\mathbb{C}S(n)$. For $k \ge 0$, the tensor space

 $V^{\otimes k} := V \otimes \cdots \otimes V(k \text{ tensor components})$

is an $\mathbb{C}S(n)$ -module via the diagonal action. For any k-tuple $\mathbf{a} = (a_1, \ldots, a_k) \in [n]^k$ we let $v_{\mathbf{a}} := v_{a_1} \otimes \cdots \otimes v_{a_k} \in V^{\otimes k}$, and for any $\pi \in S(n)$ let $\pi(\mathbf{a}) := (\pi(a_1), \ldots, \pi(a_k))$. Then $V^{\otimes k} = \operatorname{Span}_{\mathbb{C}}\{v_{\mathbf{a}} \mid \mathbf{a} \in [n]^k\}$, and the diagonal action is given by the \mathbb{C} -linear extension of $\pi v_{\mathbf{a}} = v_{\pi(\mathbf{a})}$ for all $\pi \in S(n)$ and $\mathbf{a} \in [n]^k$. We let $\operatorname{End}(V^{\otimes k})$ be the algebra of all vector space endomorphims $V^{\otimes k} \to V^{\otimes k}$. We identify any $g \in \mathbb{C}S(n)$ with the corresponding endomorphism in $\operatorname{End}(V^{\otimes k})$ given by the diagonal action. Then consider the subalgebra

$$\operatorname{End}_{S(n)}(V^{\otimes k}) := \{ f \in \operatorname{End}(V^{\otimes k}) \mid f\pi = \pi f, \text{ for all } \pi \in S(n) \}$$

of all S(n) commuting endomorphisms. For the following result see [8, Section 3].

Theorem 2.4.1 For any $n, k \ge 0$, we have a surjective \mathbb{C} -algebra homomorphism

$$\psi_{n,k}: \mathcal{A}_{2k} \to \operatorname{End}_{S(n)}(V^{\otimes k})$$

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defined on the generators z, s_i , and e_j , by letting $z \mapsto n$ and

$$\psi_{n,k}(s_i)(v_a) = v_{a_1} \otimes \cdots \otimes v_{a_{i-1}} \otimes v_{a_{i+1}} \otimes v_{a_i} \otimes v_{a_{i+2}} \otimes \cdots \otimes v_{a_k}$$

$$\psi_{n,k}(e_{2j-1})(v_a) = \sum_{b=1}^n v_{a_1} \otimes \cdots \otimes v_{a_{j-1}} \otimes v_b \otimes v_{a_{j+1}} \otimes \cdots \otimes v_{a_k}$$

$$\psi_{n,k}(e_{2j})(v_a) = \delta_{a_j,a_{j+1}} v_a$$

for all $\mathbf{a} = (a_1, \ldots, a_k) \in [n]^k$, $i \in [k-1]$, $j \in [2k-1]$, where $\delta_{a,b}$ is the Kronecker delta. Moreover, we have that $\text{Ker}(\psi_{n,k}) = (z-n)$ if and only if $n \ge 2k$, in which case $\mathcal{A}_{2k}(n) \cong \text{End}_{S(n)}(V^{\otimes k})$.

For any $a, b \in [n]$, we let $(a, b) \in S(n)$ denote the transposition exchanging a and b, and let $\varepsilon_{a,b} := 1 - \delta_{a,b}$. We now recall how the elements X_i and t_i act under $\psi_{n,k}$, which was proven in [6].

Proposition 2.4.2 [6, Proposition 5.2] For any $a \in [n]^k$, we have

$$\psi_{n,k}(t_{2i})(v_{\boldsymbol{a}}) = \varepsilon_{a_i,a_{i+1}}(a_i, a_{i+1})(v_{a_1} \otimes \cdots \otimes v_{a_{i-1}}) \otimes v_{a_i} \otimes \cdots \otimes v_{a_k}$$

$$\psi_{n,k}(t_{2i+1})(v_{\boldsymbol{a}}) = \varepsilon_{a_i,a_{i+1}}(a_i, a_{i+1})(v_{a_1} \otimes \cdots \otimes v_{a_{i+1}}) \otimes v_{a_{i+2}} \otimes \cdots \otimes v_{a_k}$$

for all $\boldsymbol{a} = (a_1, \ldots, a_k) \in [n]^k$, and $i \in [k-1]$.

Proposition 2.4.3 [6, Proposition 5.3] For any $a \in [n]^k$, we have

$$\psi_{n,k}(X_{2i-1})(v_{\boldsymbol{a}}) = \sum_{\substack{b=1\\b\neq a_i}}^n (a_i, b)(v_{a_1} \otimes \cdots \otimes v_{a_{i-1}}) \otimes v_{a_i} \otimes \cdots \otimes v_{a_k}$$
$$\psi_{n,k}(X_{2i})(v_{\boldsymbol{a}}) = \sum_{\substack{b=1\\b\neq a_i}}^n (a_i, b)(v_{a_1} \otimes \cdots \otimes v_{a_i}) \otimes v_{a_{i+1}} \otimes \cdots \otimes v_{a_k}$$

for all $a = (a_1, ..., a_k) \in [n]^k$, and $i \in [k]$.

The following result tells us that a relation holds in A_{2k} if and only if it holds under $\psi_{n,k}$ for all *n*. We will use this result to complete Proposition 2.3.2.

Lemma 2.4.4 Let $R_1, R_2 \in A_{2k}$. If $\psi_{n,k}(R_1) = \psi_{n,k}(R_2)$ for all $n \ge 1$, then $R_1 = R_2$.

Proof Follows since

$$R_1 - R_2 \in \bigcap_{n=1}^{\infty} \operatorname{Ker}(\psi_{n,k}) \subset \bigcap_{n \ge 2k} (z - n) = 0.$$

Lemma 2.4.5 The relations

(7) (Braid-like relations)

- (i) $t_{2i-2}t_{2i}t_{2i-2} = t_{2i}t_{2i-2}t_{2i}(1-e_{2i-2})$
- (ii) $t_{2i+1}t_{2i-1}t_{2i+1} = t_{2i-1}t_{2i+1}t_{2i-1}(1-e_{2i})$

- (iii) $t_{2i-1}t_{2i}t_{2i-1} = t_{2i} e_{2i-2}t_{2i} t_{2i}e_{2i-2}$
- (iv) $t_{2i}t_{2i-1}t_{2i} = t_{2i-1} e_{2i}t_{2i-1} t_{2i-1}e_{2i}$

hold in A_{2k} , thus completing the proof of Proposition 2.3.2.

Proof To ease notation, for any tuple $\mathbf{a} = (a_1, \ldots, a_k) \in [n]^k$, we represent a simple tensor in $V^{\otimes k}$ by a word in the entries of \mathbf{a} , that is $a_1 \cdots a_k := v_{a_1} \otimes \cdots \otimes v_{a_k}$. We will prove these relations by showing that they hold under $\psi_{n,k}$ for all $n \ge 1$, and then employ Lemma 2.4.4. For each relation we will have to consider different cases based on the relative values of the entries a_{i-1}, a_i , and a_{i+1} , although most cases are trivial. Also note that $\psi_{n,k}(1 - e_{2i})(\mathbf{a}) = \varepsilon_{a_i,a_{i+1}}\mathbf{a}$.

(7)(*i*): If $a_{i-1} = a_i$ or $a_i = a_{i+1}$, then it is easy to check that both $t_{2i-2}t_{2i}t_{2i-2}$ and $t_{2i}t_{2i-2}t_{2i}(1 - e_{2i-2})$ will act on **a** by 0. Assume that $a_i \neq a_{i-1} = a_{i+1}$, then

$$\psi_{n,k}(t_{2i-2}t_{2i}t_{2i-2})(\boldsymbol{a}) = \psi_{n,k}(t_{2i-2}t_{2i})\Big((a_{i-1}, a_i)(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k\Big)$$
$$= \psi_{n,k}(t_{2i-2})\Big((a_i, a_{i+1})(a_{i-1}, a_i)(a_1 \cdots a_{i-2})a_ia_i \cdots a_k\Big) = 0.$$

Similarly one can show that $\psi_{n,k}(t_{2i}t_{2i-2}t_{2i}(1-e_{2i-2}))(a) = 0$ when $a_i \neq a_{i-1} = a_{i+1}$. Lastly assume that a_{i-1}, a_i , and a_{i+1} are pairwise distinct, in particular $\varepsilon_{a,b} = 1$ for any $a, b \in \{a_{i-1}, a_i, a_{i+1}\}$. Then

$$\begin{split} \psi_{n,k}(t_{2i}t_{2i-2}t_{2i}(1-e_{2i-2}))(a) &= \psi_{n,k}(t_{2i}t_{2i-2}t_{2i})(a) \\ &= \psi_{n,k}(t_{2i}t_{2i-2})\Big((a_i, a_{i+1})(a_1 \cdots a_{i-1})a_i \cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i}t_{2i-2})\Big((a_i, a_{i+1})(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i})\Big((a_{i-1}, a_i)(a_i, a_{i+1})(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k\Big) \\ &= \Big((a_i, a_{i+1})(a_{i-1}, a_i)(a_i, a_{i+1})(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k\Big) \\ &= \Big((a_{i-1}, a_i)(a_i, a_{i+1})(a_{i-1}, a_i)(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i-2})\Big((a_i, a_{i+1})(a_{i-1}, a_i)(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i-2}t_{2i})\Big((a_{i-1}, a_i)(a_{i-1} \cdots a_{i-2})a_{i-1} \cdots a_k\Big)$$

(7)(*ii*): If $a_i = a_{i+1}$ then its clear that both $t_{2i+1}t_{2i-1}t_{2i+1}$ and $t_{2i-1}t_{2i+1}t_{2i-1}(1-e_{2i})$ act on **a** by 0. Assume that $a_i \neq a_{i+1}$ and $a_{i-1} \in \{a_i, a_{i+1}\}$, then

$$\begin{aligned} \psi_{n,k}(t_{2i+1}t_{2i-1}t_{2i+1})(\boldsymbol{a}) &= \psi_{n,k}(t_{2i+1}t_{2i-1})\Big((a_i, a_{i+1})(a_1 \cdots a_{i-1})a_{i+1}a_i a_{i+2} \cdots a_k\Big) \\ &= \varepsilon_{b,a_{i+1}}\psi_{n,k}(t_{2i+1})\Big((b, a_{i+1})(a_i, a_{i+1})(a_1 \cdots a_{i-1})ba_i a_{i+2} \cdots a_k\Big) \\ &= \varepsilon_{b,a_i}\varepsilon_{b,a_{i+1}}\Big((b, a_i)(b, a_{i+1})(a_i, a_{i+1})(a_1 \cdots a_{i-1})a_i ba_{i+2} \cdots a_k\Big)\end{aligned}$$

where $b = (a_i, a_{i+1})(a_{i-1})$. Since $a_{i-1} \in \{a_i, a_{i+1}\}$, we have that $\varepsilon_{b,a_i}\varepsilon_{b,a_{i+1}} = 0$, and so $t_{2i+1}t_{2i-1}t_{2i+1}$ acts on a by 0. Similarly one can check that $t_{2i-1}t_{2i+1}t_{2i-1}(1-e_{2i})$ also acts on a by 0. Lastly assume that a_{i-1}, a_i , and a_{i+1} are pairwise distinct. Then

$$\begin{split} \psi_{n,k}(t_{2i-1}t_{2i+1}t_{2i-1}(1-e_{2i}))(\boldsymbol{a}) &= \psi_{n,k}(t_{2i-1}t_{2i+1}t_{2i-1})(\boldsymbol{a}) \\ &= \psi_{n,k}(t_{2i-1}t_{2i+1})\Big((a_{i-1},a_i)(a_1\cdots a_{i-2})a_ia_{i-1}a_{i+1}\cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i-1})\Big((a_{i-1},a_{i+1})(a_{i-1},a_i)(a_1\cdots a_{i-2})a_ia_{i+1}a_{i-1}a_{i+2}\cdots a_k\Big) \\ &= (a_i,a_{i+1})(a_{i-1},a_{i+1})(a_{i-1},a_i)(a_1\cdots a_{i-2})a_{i+1}a_ia_{i-1}a_{i+2}\cdots a_k \\ &= (a_{i-1},a_i)(a_{i-1},a_{i+1})(a_i,a_{i+1})(a_1\cdots a_{i-2})a_{i+1}a_ia_{i-1}a_{i+2}\cdots a_k \\ &= \psi_{n,k}(t_{2i+1})\Big((a_{i-1},a_{i+1})(a_i,a_{i+1})(a_1\cdots a_{i-2})a_{i+1}a_{i-1}a_ia_{i+2}\cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i+1}t_{2i-1})\Big((a_i,a_{i+1})(a_1\cdots a_{i-2})a_{i-1}a_{i+1}a_ia_{i+2}\cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i+1}t_{2i-1})\Big((a_i,a_{i+1})(a_1\cdots a_{i-2})a_{i-1}a_{i+1}a_ia_{i+2}\cdots a_k\Big) \\ &= \psi_{n,k}(t_{2i+1}t_{2i-1}t_{2i-1})(a) \end{split}$$

(7)(*iii*): Assume $a_i = a_{i+1}$, then it is easy to check that $t_{2i} - e_{2i-2}t_{2i} - t_{2i}e_{2i-2}$ acts on a by 0. Similarly

$$\begin{aligned} \psi_{n,k}(t_{2i-1}t_{2i}t_{2i-1})(\boldsymbol{a}) &= \varepsilon_{a_{i-1},a_i}\psi_{n,k}(t_{2i-1}t_{2i})\Big((a_{i-1},a_i)(a_1\cdots a_{i-2})a_ia_{i-1}a_{i+1}\cdots a_k\Big) \\ &= \varepsilon_{a_{i-1},a_{i+1}}\varepsilon_{a_{i-1},a_i}\psi_{n,k}(t_{2i-1})\Big((a_{i-1},a_{i+1})(a_{i-1},a_i)(a_1\cdots a_{i-2})a_{i-1}a_{i-1}a_{i+1}\cdots a_k\Big) \\ &= 0.\end{aligned}$$

Now assume $a_i \neq a_{i+1}$ and $a_{i-1} \in \{a_i, a_{i+1}\}$. Then

 $\psi_{n,k}(t_{2i} - e_{2i-2}t_{2i} - t_{2i}e_{2i-2})(\mathbf{a}) = (1 - \delta_{(a_i,a_{i+1})(a_{i-1}),a_i})(a_i, a_{i+1})(a_1 \cdots a_{i-1})a_i \cdots a_k.$ In either case for $a_{i-1} = a_i$ or $a_{i-1} = a_{i+1}$, we have $\psi_{n,k}(t_{2i} - e_{2i-2}t_{2i} - t_{2i}e_{2i-2})(\mathbf{a}) = 0.$ Also, from above we see that $\psi_{n,k}(t_{2i-1}t_{2i}t_{2i-1})(\mathbf{a}) = 0$ since the factor $\varepsilon_{a_{i-1},a_{i+1}}\varepsilon_{a_{i-1},a_i}$ comes into play. Lastly, assume that a_{i-1}, a_i , and a_{i+1} are pairwise distinct, then it is easy to check that $\psi_{n,k}(e_{2i-2}t_{2i})(\mathbf{a}) = \psi_{n,k}(t_{2i}e_{2i-2})(\mathbf{a}) = 0.$ Also,

$$\psi_{n,k}(t_{2i-1}t_{2i}t_{2i-1})(\boldsymbol{a}) = \psi_{n,k}(t_{2i-1}t_{2i})\Big((a_{i-1}, a_i)(a_1 \cdots a_{i-2})a_ia_{i-1}a_{i+1} \cdots a_k\Big)$$

$$= \psi_{n,k}(t_{2i-1})\Big((a_{i-1}, a_{i+1})(a_{i-1}, a_i)(a_1 \cdots a_{i-2})a_ia_{i-1}a_{i+1} \cdots a_k\Big)$$

$$= (a_{i-1}, a_i)(a_{i-1}, a_{i+1})(a_{i-1}, a_i)(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k$$

$$= (a_i, a_{i+1})(a_1 \cdots a_{i-2})a_{i-1} \cdots a_k$$

$$= \psi_{n,k}(t_{2i})(\boldsymbol{a}) = \psi_{n,k}(t_{2i} - e_{2i-2}t_{2i} - t_{2i}e_{2i-2})(\boldsymbol{a}).$$

(7)(iv): This relation can be proved by analogous computations to (7)(iii) above.

3 Affine Partition Algebra

3.1 Definition of $\mathcal{A}_{2k}^{\text{aff}}(z)$ and basic results

In this section we give the definition of the affine partition algebra $\mathcal{A}_{2k}^{\text{aff}}$ by generators and relations. We prove some basic properties about this algebra including the fact that the

partition algebra \mathcal{A}_{2k} is both a quotient and subalgebra of \mathcal{A}_{2k}^{aff} . We also show that the polynomial algebra $\mathbb{C}[x_1, \ldots, x_{2k}]$ is a subalgebra and that $\mathcal{H}_k \otimes \mathcal{H}_k$ is a quotient, where \mathcal{H}_k is the degenerate affine Hecke algebra. We will prove a variety of relations in $\mathcal{A}_{2k}^{\text{aff}}$ including counterparts to the recursive definition of both the Jucys-Murphy elements and Enyang's generators.

Definition 3.1.1 We define the *affine partition algebra* \mathcal{A}_{2k}^{aff} to be the associative unitial C-algebra with set of generators

$$\{\tau_i, e_j, x_r, z_l \mid 2 \le i \le 2k - 1, 1 \le j \le 2k - 1, r \in [2k], l \in \mathbb{Z}_{\ge 0}\}$$

and defining relations

- (1) (Involutions)
 - (i) $\tau_{2i}^2 = 1 e_{2i}$, for $i \in [k 1]$. (ii) $\tau_{2i+1}^2 = 1 e_{2i}$, for $i \in [k 1]$.
- (2) (Braid relations)
 - (i) $\tau_{2i+1}\tau_{2i} = \tau_{2i}\tau_{2i+1}$ for $j \neq i+1$.
 - (ii) $\tau_{2i+1}\tau_{2i+1} = \tau_{2i+1}\tau_{2i+1}$ for $j \neq i \pm 1$.
 - (iii) $\tau_{2i}\tau_{2i} = \tau_{2i}\tau_{2i}$ for $j \neq i \pm 1$.
 - (iv) $s_i s_{i+1} s_i = s_{i+1} s_i s_{i+1}$, for $i \in [k-2]$, where $s_i := \tau_{2i} \tau_{2i+1} + e_{2i}$.
- (3) (Idempotent relations)
 - (i) $e_{2i-1}^2 = z_0 e_{2i-1}$ for $i \in [k]$. (ii) $e_{2i}^2 = e_{2i}$ for $i \in [k-1]$.

 - (iii) $\tau_{2i+1}e_{2i} = e_{2i}\tau_{2i+1} = 0$ for $i \in [k-1]$.
 - (iv) $\tau_{2i}e_{2i} = e_{2i}\tau_{2i} = 0$ for $i \in [k-1]$.
 - (v) $\tau_{2i}e_{2i-1}e_{2i+1} = \tau_{2i+1}e_{2i-1}e_{2i+1}$ for $i \in [k-1]$.
 - (vi) $e_{2i+1}e_{2i-1}\tau_{2i} = e_{2i+1}e_{2i-1}\tau_{2i+1}$ for $i \in [k-1]$.
- (4) (Commutation relations)
 - (i) $e_i e_j = e_j e_i$, if $|i j| \ge 2$.
 - (ii) $\tau_{2i-1}e_{2i-1} = e_{2i-1}\tau_{2i-1}$, if $j \neq i-1, i$.
 - (iii) $\tau_{2i-1}e_{2j} = e_{2j}\tau_{2i-1}$, if $j \neq i$.
 - (iv) $\tau_{2i}e_{2i-1} = e_{2i-1}\tau_{2i}$, if $j \neq i, i+1$.
 - (v) $\tau_{2i}e_{2j} = e_{2j}\tau_{2i}$, if $j \neq i 1$.
- (5) (Contractions)
 - (i) $e_i e_{i+1} e_i = e_i$ and $e_{i+1} e_i e_{i+1} = e_{i+1}$, for $i \in [2n-2]$.
 - (ii) $\tau_{2i}e_{2i-1}\tau_{2i} = \tau_{2i+1}e_{2i+1}\tau_{2i+1}$, for $i \in [k-1]$.
 - (iii) $\tau_{2i}e_{2i-2}\tau_{2i} = \tau_{2i-1}e_{2i}\tau_{2i-1}$, for $2 \le i \le k-1$.
- (6) (Affine Commuting Relations)
 - (i) $x_i x_j = x_j x_i$ for all $i, j \in [2k]$
 - (ii) $\tau_i x_j = x_j \tau_i$ for $j \neq i 1, i, i + 1$
 - (iii) $e_i x_j = x_j e_i$ for $j \neq i, i + 1$
- (7) (Braid-like relations)

- (i) $\tau_{2i-2}\tau_{2i}\tau_{2i-2} = \tau_{2i}\tau_{2i-2}\tau_{2i}(1-e_{2i-2}).$
- (ii) $\tau_{2i+1}\tau_{2i-1}\tau_{2i+1} = \tau_{2i-1}\tau_{2i+1}\tau_{2i-1}(1-e_{2i}).$
- (iii) $\tau_{2i-1}\tau_{2i}\tau_{2i-1} = \tau_{2i} e_{2i-2}\tau_{2i} \tau_{2i}e_{2i-2}$.
- (iv) $\tau_{2i}\tau_{2i-1}\tau_{2i} = \tau_{2i-1} e_{2i}\tau_{2i-1} \tau_{2i-1}e_{2i}$.

(8) (Skein-like Relations)

- (i) $x_{2i+1} = \tau_{2i} x_{2i-1} \tau_{2i} + e_{2i} e_{2i-1} \tau_{2i} + \tau_{2i} e_{2i-1} e_{2i} \tau_{2i}$
- (ii) $x_{2i+2} = \tau_{2i+1}x_{2i}\tau_{2i+1} + e_{2i}e_{2i+1}\tau_{2i+1}e_{2i+1}e_{2i} + \tau_{2i+1}$.
- (iii) $x_{2i} = \tau_{2i} x_{2i} \tau_{2i} + e_{2i} e_{2i-1} \tau_{2i} + \tau_{2i} e_{2i-1} e_{2i}$.
- (iv) $x_{2i+1} = \tau_{2i+1} x_{2i+1} \tau_{2i+1} + e_{2i} e_{2i+1} \tau_{2i+1} + \tau_{2i+1} e_{2i+1} e_{2i}$
- (9) (Anti-symmetry Relations)
 - (i) $e_i(x_i x_{i+1}) = 0$ for $i \in [2k 1]$.
 - (ii) $(x_i x_{i+1})e_i = 0$ for $i \in [2k 1]$.

(10) (Bubble Relations)

- (i) $e_1 x_1^l e_1 = z_l e_1$, for all $l \in \mathbb{N}$. (ii) z_l is central for all $l \in \mathbb{Z}_{\geq 0}$.

Note we have overloaded the symbols e_i and s_j as elements in \mathcal{A}_{2k} and $\mathcal{A}_{2k}^{\text{aff}}$, however we will show shortly that the mapping $\mathcal{A}_{2k} \to \mathcal{A}_{2k}^{\text{aff}}$ via $z \mapsto z_0, e_i \mapsto e_i$, and $s_j \mapsto s_j$ realises the subalgebra $\langle e_i, s_j, z_0 \rangle$ of $\mathcal{A}_{2k}^{\text{aff}}$ as an isomorphic copy of the partition algebra \mathcal{A}_{2k} . The defining relations above are those present in Proposition 2.3.2, except where the Jucys-Murphy elements X_i have been replaced with the affine generators x_i , Enyang's generators t_i have been replaced by new generators τ_i , and the polynomials $z(z-1)^l$ have been replaced by central generators z_l . It is worth mentioning that the map $\mathcal{A}_{2k} \to \mathcal{A}_{2k}^{aff}$ given by $z \mapsto z_0, e_i \mapsto e_i$, and $t_j \mapsto \tau_j$ does not realise an algebra homomorphism. This is since τ_2 is a non-trivial generator in $\mathcal{A}_{2k}^{\text{aff}}$, while t_2 is abscent in the presentation of Theorem 2.2.4 since it equals $1 - e_2$, hence the braid relation (E2)(iv) is not respected under such a map. The subalgebra $\langle e_i, \tau_j, z_0 \rangle$ of $\mathcal{A}_{2k}^{\text{aff}}$ is not isomorphic to the partition algebra, and in fact one can show that this subalgebra is infinite dimensional as an $\mathbb{C}[z_0]$ -module (see Corollary 3.3.3 below).

Replacing the Jucys-Murphy elements with commuting variables, and introducing new central generators is very much analogous to the 'affinization' process employed on other diagram algebras. In particular relations (6) to (10) (except (7)) are comparable to the relations in [16, Section 4] which were chosen as the defining relations for the affine Wenzl algebra. The Skein-like relations (8) tell us how the affine generators x_i interact with the generators τ_j when they do not commute. These relations are to $\mathcal{A}_{2k}^{\text{aff}}$ what the defining relation $y_{i+1} = s_i y_i s_i + s_i$ is to the degenerate affine Hecke algebra \mathcal{H}_k . In the next section we provide a projection of $\mathcal{A}_{2k}^{\text{aff}}$ onto a diagram algebra living within the Heisenberg category. Under this projection the Skein-like relations will correspond to moving a decoration over crossings.

We have also chosen to replace the generators t_j with new generators τ_j , which appears to be a departure from the 'affinization' process. However, we will show that these elements are not needed to generate the algebra, that is $\mathcal{A}_{2k}^{\text{aff}} = \langle e_i, s_i, x_i, z_l \rangle$. Hence to go from \mathcal{A}_{2k} to $\mathcal{A}_{2k}^{\text{aff}}$ we have indeed just adjoined new affine and central genrators. The reason for letting the elements τ_j play the role of generators is to allow us to give a cleaner presentation which is more comparable to its counterparts within the literature. We have chosen to include

the Braid-like relations (7) as they tell us how none commuting τ_j generators interact in a manner which resembles the braid relations of the Coxeter generators s_i . These relations will allow us to give counterparts to the recursive definition of Enyang's generators (see Lemma 3.1.11 below).

We begin by showing that the partition algebra is a quotient of the affine partition algebra. This follows naturally from its construction.

Lemma 3.1.2 We have a surjective \mathbb{C} -algebra homomorphism pr : $\mathcal{A}_{2k}^{aff} \to \mathcal{A}_{2k}$, given on the generators by

 $pr(\tau_i) = t_i$, $pr(e_i) = e_i$, $pr(x_i) = X_i$, $pr(z_l) = z(z-1)^l$.

Proof This follows by comparing the defining relations with those of the same numbering in Proposition 2.3.2, and surjectivity follows since $\langle t_i, e_j, z \rangle = A_{2k}$.

Similar to the partition algebra, the affine partition algebra has a corresponding antiautomorphism which fixes the generators.

Lemma 3.1.3 The mapping $* : \mathcal{A}_{2k}^{aff} \to \mathcal{A}_{2k}^{aff}$ which fixes the generators, extended \mathbb{C} -linearly, gives an anti-automorphism.

Proof All defining relations of Definition 3.1.1 are symmetric in the generators except relations (7)(i) and (7)(ii). Thus it is clear that the result holds if we can show that e_{2i-2} and $\tau_{2i}\tau_{2i-2}\tau_{2i}$ commute, and that e_{2i} and $\tau_{2i-1}\tau_{2i+1}\tau_{2i-1}$ commute. For the former we have

$$\tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2} = \tau_{2i}\tau_{2i-2}\tau_{2i-1}e_{2i}\tau_{2i-1}\tau_{2i}$$

= $\tau_{2i}\tau_{2i-1}e_{2i}\tau_{2i-1}\tau_{2i-2}\tau_{2i}$
= $e_{2i-2}\tau_{2i}\tau_{2i-2}\tau_{2i}$

where the first equality can be deduced from relation (5)(iii) of Definition 3.1.1, the second relations follows since τ_{2i-2} commutes with τ_{2i-1} and e_{2i} , then the last equality again is deducable from relation (5)(iii) of Definition 3.1.1. Showing that e_{2i} and $\tau_{2i-1}\tau_{2i+1}\tau_{2i-1}$ commute follows in a similar manner.

We now seek to show that A_{2k} is the subalgebra $\langle s_i, e_j, z_0 \rangle$ of A_{2k}^{aff} . We first prove a few helpful relations.

Lemma 3.1.4 *The following relations hold:*

(i) $e_{2i}x_{2i} = e_{2i}e_{2i-1}\tau_{2i}$, and $x_{2i}e_{2i} = \tau_{2i}e_{2i-1}e_{2i}$

(ii) $e_{2i}x_{2i+1} = e_{2i}e_{2i+1}\tau_{2i+1}$, and $x_{2i+1}e_{2i} = \tau_{2i+1}e_{2i+1}e_{2i}$

(ii) $e_{2i}e_{2i-1}\tau_{2i} = e_{2i}e_{2i+1}\tau_{2i+1}$, and $\tau_{2i}e_{2i-1}e_{2i} = \tau_{2i+1}e_{2i+1}e_{2i}$

Proof (*i*): Multiplying (8)(*iii*) of Definition 3.1.1 on the left by e_{2i} gives

 $e_{2i}x_{2i} = e_{2i}\tau_{2i}x_{2i}\tau_{2i} + e_{2i}e_{2i-1}\tau_{2i} + e_{2i}\tau_{2i}e_{2i-1}e_{2i} = e_{2i}e_{2i-1}\tau_{2i}$

since $e_{2i}\tau_{2i} = 0$ and $e_{2i}e_{2i} = e_{2i}$. The relation $x_{2i}e_{2i} = \tau_{2i}e_{2i-1}e_{2i}$ follows by *.

(*ii*): Multiplying (8)(*iv*) of Definition 3.1.1 on the left by e_{2i} gives

$$e_{2i}x_{2i+1} = e_{2i}\tau_{2i+1}x_{2i+1}\tau_{2i+1} + e_{2i}e_{2i}e_{2i+1}\tau_{2i+1} + e_{2i}\tau_{2i+1}e_{2i+1}e_{2i} = e_{2i}e_{2i+1}\tau_{2i+1}$$

since $e_{2i}\tau_{2i+1} = 0$ and $e_{2i}e_{2i} = e_{2i}$. The relation $x_{2i+1}e_{2i} = \tau_{2i+1}e_{2i+1}e_{2i}$ follows by *. (*iii*): By (9)(*i*), (*ii*) of Definition 3.1.1, $e_{2i}x_{2i} = e_{2i}x_{2i+1}$ and $x_{2i}e_{2i} = x_{2i+1}e_{2i}$. So (*i*) and (*ii*) imply (*iii*).

Proposition 3.1.5 We have an injective \mathbb{C} -algebra homomorphism $\iota : \mathcal{A}_{2k} \to \mathcal{A}_{2k}^{aff}$ given on the generators by $\iota(z) = z_0$, $\iota(s_i) = \tau_{2i}\tau_{2i+1} + e_{2i}$, and $\iota(e_i) = e_i$.

Proof We first prove that ι is a homomorphism. To do this we show that each of the defining relations of A_{2k} given in Theorem 2.1.1 is respected under ι . We only check the relations involving s_i since the others are accounted for in the definition of A_{2k}^{aff} .

(HR1)(i):

$$\iota(s_i^2) = (\tau_{2i}\tau_{2i+1} + e_{2i})(\tau_{2i}\tau_{2i+1} + e_{2i}) = \tau_{2i}^2\tau_{2i+1}^2 + e_{2i} = (1 - e_{2i})(1 - e_{2i}) + e_{2i} = 1 - 2e_{2i} + 2e_{2i} = 1$$

where we used (1), (2)(i), (3)(ii), (3)(iii), and (3)(iv).

(*HR1*)(*ii*): This holds by relations (2)(*i*), (2)(*ii*), (2)(*iii*) and (4). (*HR1*)(*iii*): This is precisely (2)(*iv*). (*HR2*)(*iii*):

$$\iota(e_{2i}s_i) = e_{2i}(\tau_{2i}\tau_{2i+1} + e_{2i}) = e_{2i} = \iota(e_{2i})$$

where we used (3)(*iii*) and (3)(*ii*). Similarly we have $\iota(s_i e_{2i}) = \iota(e_{2i})$. (*HR2*)(*iv*):

$$\iota(s_i e_{2i-1} e_{2i+1}) = (\tau_{2i} \tau_{2i+1} + e_{2i}) e_{2i-1} e_{2i+1}$$

$$= \tau_{2i} \tau_{2i+1} e_{2i-1} e_{2i+1} + e_{2i} e_{2i-1} e_{2i+1}$$

$$= \tau_{2i}^2 e_{2i-1} e_{2i+1} + e_{2i} e_{2i-1} e_{2i+1}$$

$$= e_{2i-1} e_{2i+1} - e_{2i} e_{2i-1} e_{2i+1} + e_{2i} e_{2i-1} e_{2i+1}$$

$$= e_{2i-1} e_{2i+1} = \iota(e_{2i-1} e_{2i+1})$$

where the third equality follows from (3)(v) and the forth from (1)(i). Similarly we have $\iota(e_{2i-1}e_{2i+1}s_i) = \iota(e_{2i-1}e_{2i+1})$.

(HR3)(iv): Follows from commuting relations (4)(i), (4)(ii), and (4)(iv). (HR3)(v): Follows from commuting relations (4)(i), (4)(iii), and (4)(v). (HR3)(vi):

$$\begin{split} \iota(s_i e_{2i-1} s_i) &= (\tau_{2i} \tau_{2i+1} + e_{2i}) e_{2i-1} (\tau_{2i} \tau_{2i+1} + e_{2i}) \\ &= \tau_{2i+1} \tau_{2i} e_{2i-1} \tau_{2i} \tau_{2i+1} + \tau_{2i+1} \tau_{2i} e_{2i-1} e_{2i} + e_{2i} e_{2i-1} \tau_{2i} \tau_{2+1} + e_{2i} \\ &= \tau_{2i+1}^2 e_{2i+1} \tau_{2i+1}^2 + \tau_{2i+1}^2 e_{2i+1} e_{2i} + e_{2i} e_{2i+1} \tau_{2i+1}^2 + e_{2i} \\ &= (1 - e_{2i}) e_{2i+1} (1 - e_{2i}) + e_{2i+1} e_{2i} - e_{2i} + e_{2i} e_{2i+1} - e_{2i} + e_{2i} \\ &= e_{2i+1} - e_{2i} e_{2i+1} - e_{2i+1} e_{2i} + e_{2i} + e_{2i+1} e_{2i} - e_{2i} + e_{2i} e_{2i+1} \\ &= e_{2i+1} = \iota(e_{2i+1}) \end{split}$$

where the third equality follows by Lemma 3.1.4 (iii) and (5)(*ii*), and the forth from $\tau_{2i+1}^2 = 1 - e_{2i}$.

$$(HR3)(vii):$$

$$\iota(s_ie_{2i-2}s_i) = (\tau_{2i+1}\tau_{2i} + e_{2i})e_{2i-2}(\tau_{2i}\tau_{2i+1} + e_{2i})$$

$$= \tau_{2i}\tau_{2i+1}e_{2i-2}\tau_{2i+1}\tau_{2i} + \tau_{2i+1}\tau_{2i}e_{2i-2}e_{2i} + e_{2i}e_{2i-2}\tau_{2i}\tau_{2i+1} + e_{2i}e_{2i-2}e_{2i}$$

$$= \tau_{2i}\tau_{2i+1}^2e_{2i-2}\tau_{2i} + e_{2i}e_{2i-2}$$

$$= \tau_{2i}e_{2i-2}\tau_{2i} + e_{2i}e_{2i-2}$$

$$= \tau_{2i-1}e_{2i}\tau_{2i-1} + e_{2i}e_{2i-2}$$

where the third equality follows since τ_{2i+1} and e_{2i} commute with e_{2i-2} , $e_{2i}^2 = e_{2i}$, and $e_{2i}\tau_{2i} = \tau_{2i}e_{2i} = 0$. We also have

$$\begin{split} \iota(s_{i-1}e_{2i}s_{i-1}) &= (\tau_{2i-2}\tau_{2i-1} + e_{2i-2})e_{2i}(\tau_{2i-2}\tau_{2i-1} + e_{2i-2}) \\ &= \tau_{2i-1}\tau_{2i-2}e_{2i}\tau_{2i-2}\tau_{2i-1} + \tau_{2i-1}\tau_{2i-2}e_{2i}e_{2i-2} + e_{2i-2}e_{2i}\tau_{2i-2}\tau_{2i-1} + e_{2i-2}e_{2i}e_{2i-2} \\ &= \tau_{2i-1}\tau_{2i-2}^{2}e_{2i}\tau_{2i-1} + e_{2i}e_{2i-2} \\ &= \tau_{2i-1}e_{2i}\tau_{2i-1} - \tau_{2i-1}e_{2i-2}e_{2i}\tau_{2i-1} + e_{2i}e_{2i-2} \\ &= \tau_{2i-1}e_{2i}\tau_{2i-1} + e_{2i}e_{2i-2} \end{split}$$

where the third equality follows since τ_{2i-2} and e_{2i-2} commute with e_{2i} , $e_{2i-2}^2 = e_{2i-2}$, and $e_{2i-2}\tau_{2i-2} = \tau_{2i-2}e_{2i-2} = 0$. The forth equality follows since $\tau_{2i-1}e_{2i-2} = 0$. Comparing to above, we see that $\iota(s_ie_{2i-2}s_i) = \iota(s_{i-1}e_{2i}s_{i-1})$.

Hence we have shown that ι is indeed an algebra homomorphism. For injectivity, note that pr $\circ \iota = id$ where id : $\mathcal{A}_{2k} \to \mathcal{A}_{2k}$ is the identity morphism. Thus ι has a left inverse, and so is injective.

Therefore the partition algebra A_{2k} is both a subalgebra and quotient of the affine partition algebra A_{2k}^{aff} . Also note that restricting * down to the partition algebra coincides with the anti-automorphism of flipping a diagram. We now seek to give affine counterparts to the recursive definition of the Jucys-Murphy elements given in Definition 2.2.1.

Lemma 3.1.6 The following relations hold in \mathcal{A}_{2k}^{aff} :

- (i) $x_{2i+1} = s_i x_{2i-1} s_i + x_{2i} e_{2i} + e_{2i} x_{2i} x_{2i-1} e_{2i} \tau_{2i}$
- (ii) $x_{2i+2} = s_i x_{2i} s_i s_i x_{2i} e_{2i} e_{2i} x_{2i} s_i + e_{2i} x_{2i} e_{2i+1} e_{2i} + \tau_{2i+1}$

Proof (*i*): Multiplying on the left and right of equation (8)(*i*) in Definition 3.1.1 by τ_{2i+1} gives

$$\tau_{2i+1}x_{2i+1}\tau_{2i+1} = \tau_{2i+1}\tau_{2i}x_{2i-1}\tau_{2i}\tau_{2i+1} - \tau_{2i+1}\tau_{2i}\tau_{2i+1}$$

= $(s_i - e_{2i})x_{2i-1}(s_i - e_{2i}) - (s_i - e_{2i})\tau_{2i+1}$
= $s_ix_{2i-1}s_i - e_{2i}x_{2i-1}s_i - s_ix_{2i-1}e_{2i} + x_{2i+1} - \tau_{2i}$
= $s_ix_{2i-1}s_i - x_{2i-1}e_{2i} - \tau_{2i}$

where, in the first equality we used the fact that $\tau_{2i+1}e_{2i} = e_{2i}\tau_{2i+1} = 0$, the second equality we used the substitution $\tau_{2i}\tau_{2i+1} = \tau_{2i+1}\tau_{2i} = s_i - e_{2i}$, and the last equality we used the fact that e_{2i} and x_{2i-1} commute. Now applying (8)(*iv*) from Definition 3.1.1 to the left hand side of above, we obtain

$$x_{2i+1} - e_{2i}e_{2i+1}\tau_{2i+1} - \tau_{2i+1}e_{2i+1}e_{2i} = s_ix_{2i-1}s_i - x_{2i-1}e_{2i} - \tau_{2i}.$$

By applying Lemma 3.1.4 (ii), and rearranging, we arrive at (*i*). Item (*ii*) is proved in an analogous manner were we instead employ relations (8)(ii) and (8)(iii) from Definition 3.1.1.

By rearranging the relations in the above lemma in terms of the generators τ_{2i} and τ_{2i+1} , we immediately obtain the following:

Corollary 3.1.7 We have that $\mathcal{A}_{2k}^{aff} = \langle e_i, s_j, x_k, z_l \rangle_{i,j,k,l}$.

Recall that the degenerate affine Hecke algebra \mathcal{H}_k is the \mathbb{C} -algebra given as a vector space by the tensor product $\mathbb{C}[y_1, \ldots, y_k] \otimes \mathbb{C}S(k)$, where $\mathbb{C}[y_1, \ldots, y_k]$ is the polynomial algebra in commuting variables y_1, \ldots, y_k . The defining relations of \mathcal{H}_k are such that $\mathbb{C}[y_1, \ldots, y_k]$ and $\mathbb{C}S(k)$ are subalgebras, and

$s_i y_j = y_j s_i,$	for all $j \neq i, i + 1$,
$y_{i+1} = s_i y_i s_i + s_i,$	for each $i \in [k-1]$.

It turns out that $\mathcal{H}_k \otimes \mathcal{H}_k$ is a quotient of $\mathcal{A}_{2k}^{\text{aff}}$.

Proposition 3.1.8 Let $\lambda = (\lambda_l)_{l=0}^{\infty}$ be any sequence of constants in \mathbb{C} . Then we have a surjective \mathbb{C} -algebra homomorphism $f_{\lambda} : \mathcal{A}_{2k}^{aff} \to \mathcal{H}_k \otimes \mathcal{H}_k$ given on the generators by

$f_{\lambda}(\tau_{2i+1}) = s_i \otimes 1,$	$f_{\lambda}(x_{2i-1}) = -1 \otimes y_i,$
$f_{\lambda}(\tau_{2i}) = 1 \otimes s_i,$	$f_{\lambda}(x_{2i}) = y_i \otimes 1,$
$f_{\lambda}(e_i) = 0,$	$f_{\lambda}(z_l) = \lambda_l.$

Proof We show that each of the defining relations of A_{2k}^{aff} are upheld under f_{λ} . Since $f_{\lambda}(e_i) = 0$, one may observe that most of the defining relations involving generators e_i are trivially upheld.

(1)(*i*): $f_{\lambda}(\tau_{2i}^2) = (1 \otimes s_i)(1 \otimes s_i) = 1 \otimes s_i^2 = 1 = f_{\lambda}(1 - e_{2i}).$ (1)(*ii*): Similar to (1)(*i*) above.

(2)(i): For any $j \neq i + 1$, $f_{\lambda}(\tau_{2i+1}\tau_{2j}) = (s_i \otimes 1)(1 \otimes s_j) = (1 \otimes s_j)(s_i \otimes 1) = f_{\lambda}(\tau_{2j}\tau_{2i+1}).$

(2)(*ii*): For any $j \neq i \pm 1$,

 $f(\tau_{2i+1}\tau_{2j+1}) = (s_i \otimes 1)(s_j \otimes 1) = s_i s_j \otimes 1 = s_j s_i \otimes 1 = (s_j \otimes 1)(s_i \otimes 1) = f(\tau_{2j+1}\tau_{2i+1}).$

(2)(*iii*): Similar to (2)(*ii*) above. (2)(*iv*): Noting that $f_{\lambda}(s_i) = f_{\lambda}(\tau_{2i}\tau_{2i+1} + e_{2i}) = f_{\lambda}(\tau_{2i})f_{\lambda}(\tau_{2i+1}) = s_i \otimes s_i$, then

$$f_{\lambda}(s_i s_{i+1} s_i) = s_i s_{i+1} s_i \otimes s_i s_{i+1} s_i = s_{i+1} s_i s_{i+1} \otimes s_{i+1} s_i s_{i+1} = f_{\lambda}(s_{i+1} s_i s_{i+1}).$$

(6)(*i*): Follows since y_1, \ldots, y_k pariwise commute. (6)(*ii*): Follows since $s_i y_j = y_j s_i$ whenever $j \neq i, i + 1$. (7)(*i*):

$$f_{\lambda}(\tau_{2i-2}\tau_{2i}\tau_{2i-2}) = 1 \otimes s_{i-1}s_i s_{i-1} = 1 \otimes s_i s_{i-1}s_i = f_{\lambda}(\tau_{2i}\tau_{2i-2}\tau_{2i}) = f_{\lambda}(\tau_{2i}\tau_{2i-2}\tau_{2i}(1-e_{2i-2}))$$

(7)(*ii*): Similar to (7)(*i*). (7)(*iii*): $f_{\lambda}(\tau_{2i-1}\tau_{2i}\tau_{2i-1}) = s_{i-1}^2 \otimes s_i = 1 \otimes s_i = f_{\lambda}(\tau_{2i}) = f_{\lambda}(\tau_{2i} - e_{2i-2}\tau_{2i} - \tau_{2i}e_{2i-2}).$ (7)(*iv*): Similar to (7)(*iii*).

(8)(i): $f_{\lambda}(\tau_{2i}x_{2i-1}\tau_{2i} + e_{2i}e_{2i-1}\tau_{2i} + \tau_{2i}e_{2i-1}e_{2i} - \tau_{2i}) = f_{\lambda}(\tau_{2i}x_{2i-1}\tau_{2i}) - f_{\lambda}(\tau_{2i})$ $= (1 \otimes s_i)(-1 \otimes y_i)(1 \otimes s_i) - 1 \otimes s_i$ $= -1 \otimes s_i v_i s_i - 1 \otimes s_i$ $= -1 \otimes (v_{i+1} - s_i) - 1 \otimes s_i$ $= -1 \otimes y_{i+1}$ $= f_{\lambda}(x_{2i+1})$ where the forth equality follows since $s_i y_i s_i = y_{i+1} - s_i$ in \mathcal{H}_k . (8)(*ii*): $f_{\lambda}(\tau_{2i+1}x_{2i}\tau_{2i+1} + e_{2i}e_{2i+1}\tau_{2i+1}e_{2i+1}e_{2i} + \tau_{2i+1}) = f_{\lambda}(\tau_{2i+1}x_{2i}\tau_{2i+1}) + f_{\lambda}(\tau_{2i+1})$ $= (s_i \otimes 1)(v_i \otimes 1)(s_i \otimes 1) + s_i \otimes 1$ $= (s_i y_i s_i + s_i) \otimes 1$ $= y_{i+1} \otimes 1$ $= f_{\lambda}(x_{2i+2})$ where the forth equality follows since $y_{i+1} = s_i y_i s_i + s_i$ in \mathcal{H}_k . (8)(*iii*): $f_{\lambda}(\tau_{2i}x_{2i}\tau_{2i} + e_{2i}e_{2i-1}\tau_{2i-1} + \tau_{2i-1}e_{2i-1}e_{2i}) = f_{\lambda}(\tau_{2i}x_{2i}\tau_{2i}),$ $= (1 \otimes s_i)(y_i \otimes 1)(1 \otimes s_i),$ $= v_i \otimes 1$. $= f_{\lambda}(x_{2i}).$ (8)(iv): $f_{\lambda}(\tau_{2i+1}x_{2i+1}\tau_{2i+1} + e_{2i}e_{2i+1}\tau_{2i+1} + \tau_{2i+1}e_{2i+1}e_{2i}) = f_{\lambda}(\tau_{2i+1}x_{2i+1}\tau_{2i+1}),$ $= (s_i \otimes 1)(-1 \otimes y_i)(s_i \otimes 1),$

(10)(*i*) and (10)(*ii*): Immediate.

Thus f_{λ} is a homomorphism. Surjectivity follows as $\langle f_{\lambda}(\tau_i), f_{\lambda}(x_i) \rangle_{i,i} = \mathcal{H}_k \otimes \mathcal{H}_k$. \Box

 $= -1 \otimes y_i,$ $= f_{\lambda}(x_{2i+1}).$

Corollary 3.1.9 The polynomial algebra $\mathbb{C}[x_1, \ldots, x_{2k}]$ is a subalgebra of \mathcal{A}_{2k}^{aff} .

Proof This is the same as asking that all monomials in the generators of the subalgebra $\langle x_1, \ldots, x_{2k} \rangle$ of $\mathcal{A}_{2k}^{\text{aff}}$ are linearly independent, which follows since their images under f_{λ} are.

To end this section we establish a counterpart to the recursive relations of Enyang's generators. To do so, we collect the more technical relations needed into the following lemma:

Lemma 3.1.10 The following relations hold in \mathcal{A}_{2k}^{aff} :

(i) $e_{2i}x_{2i}e_{2i} = 0$ (ii) $e_{2i}\tau_{2i-1}e_{2i} = 0$ (iii) $e_{2i-2}\tau_{2i}e_{2i-2} = 0$

- (iv) $e_{2i-2}\tau_{2i} = e_{2i-2}x_{2i-2}s_ie_{2i-2}s_i$
- (v) $\tau_{2i}e_{2i-2} = s_ie_{2i-2}s_ix_{2i-2}e_{2i-2}$
- (vi) $\tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2} = e_{2i-2}x_{2i-2}s_{i-1}e_{2i}e_{2i-1}e_{2i-2}$
- (vii) $\tau_{2i-1}e_{2i}s_{i-1} = s_ie_{2i-2}e_{2i-1}e_{2i}s_{i-1}x_{2i-2}e_{2i-2}s_i$
- (viii) $\tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2} = e_{2i-2}\tau_{2i}\tau_{2i-2}\tau_{2i}$

Proof (*i*): We have $e_{2i}x_{2i}e_{2i} = e_{2i}e_{2i-1}\tau_{2i}e_{2i} = 0$, by employing Lemma 3.1.4 (i) and Definition 3.1.1 (3)(iv).

(*ii*): By rearranging (7)(*iv*) of Definition 3.1.1 in terms of τ_{2i-1} , we have that

 $e_{2i}\tau_{2i-1}e_{2i} = e_{2i}(\tau_{2i}\tau_{2i-1}\tau_{2i} + e_{2i}\tau_{2i-1} + \tau_{2i-1}e_{2i})e_{2i} = e_{2i}\tau_{2i-1}e_{2i} + e_{2i}\tau_{2i-1}e_{2i},$

where we used relation $e_{2i}\tau_{2i} = 0$. Rearranging gives $e_{2i}\tau_{2i-1}e_{2i} = 0$. Item (*iii*) follows in a similar manner.

(*iv*): We have

$$e_{2i-2}x_{2i-2}s_ie_{2i-2}s_i = e_{2i-2}x_{2i-1}s_ie_{2i-2}s_i$$

= $e_{2i-2}(s_ix_{2i+1} - s_ix_{2i}e_{2i} - e_{2i}x_{2i} + x_{2i-1}e_{2i} + \tau_{2i+1})e_{2i-2}s_i$
= $e_{2i-2}s_ix_{2i+1}e_{2i-2}s_i - e_{2i-2}s_ix_{2i}e_{2i-2}s_i - e_{2i-2}e_{2i}x_{2i}e_{2i-2}s_i$
+ $e_{2i-2}x_{2i-1}e_{2i}e_{2i-2}s_i + e_{2i-2}\tau_{2i+1}e_{2i-2}s_i$

where the first equality follows from (9)(i) of Definition 3.1.1, and the second from Lemma 3.1.6 (i). We examine the five terms above:

- (1) $e_{2i-2}s_ix_{2i+1}e_{2i-2}s_i = e_{2i-2}s_ie_{2i-2}x_{2i+1}s_i = e_{2i-2}e_{2i}x_{2i+1}s_i$,
- (2) $-e_{2i-2}s_ix_{2i}e_{2i-2}s_i = -e_{2i-2}s_ie_{2i-2}x_{2i}e_{2i} = -e_{2i-2}e_{2i}x_{2i}e_{2i} = 0,$

 $(3) \quad -e_{2i-2}e_{2i}x_{2i}e_{2i-2}s_i = -e_{2i-2}e_{2i}x_{2i}s_i = -e_{2i-2}e_{2i}x_{2i+1}s_i,$

- (4) $e_{2i-2}x_{2i-1}e_{2i}e_{2i-2}s_i = e_{2i-2}x_{2i-1}e_{2i-2}e_{2i}s_i = 0,$
- (5) $e_{2i-2}\tau_{2i+1}e_{2i-2}s_i = e_{2i-2}\tau_{2i+1}s_i = e_{2i-2}\tau_{2i}$.

Substituting back into the above equation gives $e_{2i-2}x_{2i-2}s_ie_{2i-2}s_i = e_{2i-2}\tau_{2i}$ as desired.

$$\begin{aligned} \tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2} &= \tau_{2i}\tau_{2i-2}(s_ie_{2i-2}s_ix_{2i-2}e_{2i-2}), \\ &= \tau_{2i}\tau_{2i-2}s_{i-1}e_{2i}s_{i-1}x_{2i-2}e_{2i-2}, \\ &= \tau_{2i}\tau_{2i-1}e_{2i}s_{i-1}x_{2i-2}e_{2i-2}, \\ &= \tau_{2i}(\tau_{2i}e_{2i-2}\tau_{2i}\tau_{2i-1})s_{i-1}x_{2i-2}e_{2i-2}, \\ &= (1-e_{2i})e_{2i-2}\tau_{2i}\tau_{2i-2}x_{2i-2}e_{2i-2}, \\ &= e_{2i-2}\tau_{2i}\tau_{2i-2}x_{2i-2}e_{2i-2}, \\ &= e_{2i-2}\tau_{2i}(x_{2i-2}\tau_{2i-2}+e_{2i-3}e_{2i-2}-e_{2i-2}e_{2i-3})e_{2i-2}, \\ &= e_{2i-2}\tau_{2i}e_{2i-3}e_{2i-2}, \\ &= e_{2i-2}\tau_{2i}e_{2i-3}e_{2i-2}, \\ &= e_{2i-2}\tau_{2i}e_{2i-3}e_{2i-2}, \\ &= e_{2i-2}\tau_{2i}e_{2i-3}e_{2i-2}, \\ &= (e_{2i-2}x_{2i-2}s_ie_{2i-2}s_i)e_{2i-3}e_{2i-2}, \\ &= e_{2i-2}x_{2i-2}s_ie_{2i-2}s_i)e_{2i-3}e_{2i-2}, \\ &= e_{2i-2}x_{2i-2}s_ie_{2i-2}s_ie_{2i-2}e_{2i-2}, \end{aligned}$$

The first equality follows by (v), the forth from (5)(iii) of Lemma 3.1.4, the sixth since $e_{2i}\tau_{2i} = 0$, the seventh from (8)(iii) of Definition 3.1.1, the ninth from $\tau_{2i-2}e_{2i-2} = 0$ and (iii), and the tenth from (iv).

(vii):

$$s_i e_{2i-2} e_{2i-1} e_{2i} s_{i-1} x_{2i-2} e_{2i-2} s_i = s_i e_{2i-2} e_{2i-1} s_{i-1} s_i e_{2i-2} s_i x_{2i-2} e_{2i-2} s_i$$

$$= s_i e_{2i-2} s_i e_{2i-3} e_{2i-2} x_{2i-2} s_i e_{2i-2} s_i$$

$$= s_{i-1} e_{2i} e_{2i-1} e_{2i-2} x_{2i-2} s_{i-1} e_{2i} s_{i-1}$$

$$= s_{i-1} e_{2i} e_{2i-1} e_{2i-2} e_{2i-1} \tau_{2i-1} s_{i-1} e_{2i} s_{i-1}$$

$$= s_{i-1} e_{2i} e_{2i-1} \tau_{2i-2} e_{2i} s_{i-1}$$

$$= s_{i-1} e_{2i} \tau_{2i-1}$$

$$= s_{i-1} e_{2i} \tau_{2i-2} s_{i-1}$$

$$= s_{i-1} e_{2i} \tau_{2i-2} s_{i-1}$$

$$= s_{i-1} \tau_{2i-2} e_{2i} s_{i-1}$$

$$= \tau_{2i-1} e_{2i} s_{i-1}$$

where the forth equality follows from Lemma 3.1.4 (i).

(viii):

$$\tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2} = \tau_{2i}\tau_{2i-2}(\tau_{2i-1}e_{2i}\tau_{2i-1}\tau_{2i})$$

= $\tau_{2i}(s_{i-1} - e_{2i-2})e_{2i}\tau_{2i-1}\tau_{2i}$
= $\tau_{2i}s_{i-1}e_{2i}\tau_{2i-1}\tau_{2i}$
= $\tau_{2i}s_ie_{2i-2}s_is_{i-1}\tau_{2i-1}\tau_{2i}$
= $\tau_{2i+1}e_{2i-2}s_i\tau_{2i-2}\tau_{2i}$
= $e_{2i-2}\tau_{2i}\tau_{2i-2}\tau_{2i}$

where the first equality follows from (5)(*iii*) of Definition 3.1.1, the second since $s_{i-1} = \tau_{2i-1}\tau_{2i-2} + e_{2i-2}$, the third since $e_{2i-2}\tau_{2i-1} = 0$, and the sixth since τ_{2i+1} and e_{2i-2} commute.

Lemma 3.1.11 The following relations hold in \mathcal{A}_{2k}^{aff} :

$$\tau_{2i} = s_{i-1}s_i\tau_{2i-2}s_is_{i-1} + e_{2i-2}x_{2i-2}s_ie_{2i-2}s_i + s_ie_{2i-2}x_{2i-2}s_ie_{2i-2}$$

$$-e_{2i-2}x_{2i-2}s_{i-1}e_{2i}e_{2i-1}e_{2i-2} - s_ie_{2i-2}e_{2i-1}e_{2i}s_{i-1}x_{2i-2}e_{2i-2}s_i.$$

and

$$\tau_{2i+1} = s_{i-1}s_i\tau_{2i-1}s_is_{i-1} + s_ie_{2i-2}x_{2i-2}s_ie_{2i-2}s_i + e_{2i-2}x_{2i-2}s_ie_{2i-2}$$

-s_ie_{2i-2}x_{2i-2}s_{i-1}e_{2i}e_{2i-1}e_{2i-2} - e_{2i-2}e_{2i-1}e_{2i}s_{i-1}x_{2i-2}e_{2i-2}s_i

Proof We prove the first relation, the second follows from by multiplying on the left by s_i . We have that

$$s_i \tau_{2i-2} s_i = (\tau_{2i} \tau_{2i+1} + e_{2i}) \tau_{2i-2} (\tau_{2i+1} \tau_{2i} + e_{2i})$$

= $\tau_{2i} \tau_{2i+1}^2 \tau_{2i-2} \tau_{2i} + \tau_{2i-2} e_{2i}$
= $\tau_{2i} \tau_{2i-2} \tau_{2i} + \tau_{2i-2} e_{2i}$

where the second equality follows since τ_{2i-2} commutes with τ_{2i+1} and $e_{2i}\tau_{2i} = \tau_{2i}e_{2i} = 0$. Substituting the above we get

$$s_{i-1}s_i\tau_{2i-2}s_is_{i-1} = s_{i-1}\tau_{2i}\tau_{2i-2}\tau_{2i}s_{i-1} + \tau_{2i-1}e_{2i}s_{i-1}.$$
(10)

For the first term in equation (10) we have

$$s_{i-1}\tau_{2i}\tau_{2i-2}\tau_{2i}s_{i-1} = s_{i-1}(\tau_{2i-2}\tau_{2i}\tau_{2i-2} + \tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2})s_{i-1}$$

= $\tau_{2i-1}\tau_{2i}\tau_{2i-1} + s_{i-1}\tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2}$
= $\tau_{2i-1}\tau_{2i}\tau_{2i-1} + \tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2}$
= $\tau_{2i} - e_{2i-2}\tau_{2i} - \tau_{2i}e_{2i-2} + \tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2}$

where the first equality follows by (7)(*i*) of Definition 3.1.1, the second from $s_{i-1}\tau_{2i-2} = \tau_{2i-2}s_{i-1} = \tau_{2i-1}$, the third from Lemma 3.1.10 (viii), and the forth from (7)(*iii*) of Definition 3.1.1. Substituting this back into equation (10), and rearranging yields

 $\tau_{2i} = s_{i-1}s_i\tau_{2i-2}s_is_{i-1} + e_{2i-2}\tau_{2i} + \tau_{2i}e_{2i-2} - \tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2} - \tau_{2i-1}e_{2i}s_{i-1}.$

The desired relation is obtained by applying relations (*iv*) to (*vii*) of Lemma 3.1.10. \Box

3.2 Central Elements in $\mathcal{A}_{2k}^{\text{aff}}$

In this section we describe a central subalgebra of $\mathcal{A}_{2k}^{\text{aff}}$ consisting of certain polynomials in the affine generators. We end the section with a conjecture describing the center of $\mathcal{A}_{2k}^{\text{aff}}$.

Lemma 3.2.1 The following relations hold:

(i) $\tau_{2i}x_{2i+1} = x_{2i-1}\tau_{2i} + e_{2i-1}e_{2i} - 1.$

- (ii) $\tau_{2i+1}x_{2i+2} = x_{2i}\tau_{2i+1} e_{2i}e_{2i+1} + 1.$
- (iii) $\tau_{2i}x_{2i} = x_{2i}\tau_{2i} + e_{2i-1}e_{2i} e_{2i}e_{2i-1}$.

(iv) $\tau_{2i+1}x_{2i+1} = x_{2i+1}\tau_{2i+1} - e_{2i}e_{2i+1} + e_{2i+1}e_{2i}$.

Proof (*i*): Multiplying (8)(*i*) of Definition 3.1.1 on the left by τ_{2i} gives

$$\tau_{2i}x_{2i+1} = \tau_{2i}^2 x_{2i-1}\tau_{2i} + \tau_{2i}e_{2i}e_{2i-1}\tau_{2i} + \tau_{2i}^2 e_{2i-1}e_{2i} - \tau_{2i}^2$$

= $(1 - e_{2i})x_{2i-1}\tau_{2i} + (1 - e_{2i})e_{2i-1}e_{2i} - (1 - e_{2i})$
= $x_{2i-1}\tau_{2i} + x_{2i-1}e_{2i}\tau_{2i} + e_{2i-1}e_{2i} - e_{2i} - 1 + e_{2i}$
= $x_{2i-1}\tau_{2i} + e_{2i-1}e_{2i} - 1$

where the second equality follows as $\tau_{2i}^2 = 1 - e_{2i}$ and $t_{2i}e_{2i} = 0$, and the third since x_{2i-1} and e_{2i} commute.

(*ii*): Multiplying (8)(*ii*) of Definition 3.1.1 on the left by τ_{2i+1} gives

$$\begin{aligned} \tau_{2i+1} x_{2i+2} &= \tau_{2i+1}^2 x_{2i} \tau_{2i+1} + \tau_{2i+1} e_{2i} e_{2i+1} \tau_{2i+1} e_{2i+1} e_{2i} + \tau_{2i+1}^2 \\ &= (1 - e_{2i}) x_{2i} \tau_{2i+1} + 1 - e_{2i} \\ &= x_{2i} \tau_{2i+1} - e_{2i} x_{2i} \tau_{2i+1} + 1 - e_{2i} \\ &= x_{2i} \tau_{2i+1} - e_{2i} e_{2i-1} \tau_{2i+1}^2 + 1 - e_{2i} \\ &= x_{2i} \tau_{2i+1} - e_{2i} e_{2i-1} \tau_{2i+1}^2 + 1 - e_{2i} \\ &= x_{2i} \tau_{2i+1} - e_{2i} e_{2i-1} + e_{2i} + 1 - e_{2i} \\ &= x_{2i} \tau_{2i+1} - e_{2i} e_{2i-1} + e_{2i} + 1 - e_{2i} \end{aligned}$$

where the second equality follows since $\tau_{2i+1}e_{2i} = 0$ and $\tau_{2i+1}^2 = 1 - e_{2i}$, the forth equality follows since $e_{2i}x_{2i} = e_{2i}x_{2i+1}$, and the fifth equality follows since $e_{2i}x_{2i} = e_{2i}e_{2i-1}\tau_{2i+1}$ (by Lemma 3.1.4 (ii) and (*iii*)).

(*iii*): Multiplying (8)(*iii*) of Definition 3.1.1 on the left by τ_{2i} gives

$$\tau_{2i}x_{2i} = \tau_{2i}^2 x_{2i}\tau_{2i} + \tau_{2i}e_{2i}e_{2i-1}\tau_{2i} + \tau_{2i}^2 e_{2i-1}e_{2i}$$

= $(1 - e_{2i})x_{2i}\tau_{2i} + (1 - e_{2i})e_{2i-1}e_{2i}$
= $x_{2i}\tau_{2i} - e_{2i}x_{2i}\tau_{2i} + e_{2i-1}e_{2i} - e_{2i}$
= $x_{2i}\tau_{2i} - e_{2i}e_{2i-1}\tau_{2i}^2 + e_{2i-1}e_{2i} - e_{2i}$
= $x_{2i}\tau_{2i} - e_{2i}e_{2i-1} + e_{2i} + e_{2i-1}e_{2i} - e_{2i}$
= $x_{2i}\tau_{2i} - e_{2i}e_{2i-1} + e_{2i} + e_{2i-1}e_{2i} - e_{2i}$

where the second equality follows since $\tau_{2i}e_{2i} = 0$ and $\tau_{2i+1}^2 = 1 - e_{2i}$, and the forth equality follows since $e_{2i}x_{2i} = e_{2i}e_{2i-1}\tau_{2i}$ (by Lemma 3.1.4 (i)).

(*iv*): Multiplying (8)(*iv*) of Definition 3.1.1 on the left by τ_{2i+1} gives

$$\tau_{2i+1}x_{2i+1} = \tau_{2i+1}^2 x_{2i+1}\tau_{2i+1} + \tau_{2i+1}e_{2i}e_{2i+1}\tau_{2i+1} + \tau_{2i+1}^2 e_{2i+1}e_{2i}$$

$$= (1 - e_{2i})x_{2i+1}\tau_{2i+1} + (1 - e_{2i})e_{2i+1}e_{2i}$$

$$= x_{2i+1}\tau_{2i+1} - e_{2i}x_{2i+1}\tau_{2i+1} + e_{2i+1}e_{2i} - e_{2i}$$

$$= x_{2i+1}\tau_{2i+1} - e_{2i}e_{2i+1}\tau_{2i+1}^2 + e_{2i+1}e_{2i} - e_{2i}$$

$$= x_{2i}\tau_{2i} - e_{2i}e_{2i+1} + e_{2i} + e_{2i+1}e_{2i} - e_{2i}$$

$$= x_{2i}\tau_{2i} - e_{2i}e_{2i+1} + e_{2i+1}e_{2i}$$

where the second equality follows since $\tau_{2i+1}e_{2i} = 0$ and $\tau_{2i+1}^2 = 1 - e_{2i}$, and the forth equality follows since $e_{2i}x_{2i+1} = e_{2i}e_{2i+1}\tau_{2i+1}$ (by Lemma 3.1.4 (ii)).

Lemma 3.2.2 For any $n \ge 1$, the following relations hold:

(i)
$$\tau_{2i}x_{2i+1}^n = x_{2i-1}^n\tau_{2i} + \sum_{\substack{a+b=n-1\\a,b>0}} x_{2i-1}^a (e_{2i-1}e_{2i} - 1)x_{2i+1}^b.$$

(ii)
$$\tau_{2i} x_{2i}^n = x_{2i}^n \tau_{2i} + \sum_{\substack{a+b=n-1\\a,b>0}} \overline{x_{2i}^a} (e_{2i-1}e_{2i} - e_{2i}e_{2i-1}) x_{2i}^b.$$

(iii)
$$\tau_{2i+1} x_{2i+2}^n = x_{2i}^n \tau_{2i+1} + \sum_{\substack{a+b=n-1\\a,b>0}} x_{2i}^a (-e_{2i}e_{2i+1} + 1) x_{2i+2}^b.$$

(iv)
$$\tau_{2i+1}x_{2i+1}^n = x_{2i+1}^n \tau_{2i+1} + \sum_{\substack{a+b=n-1\\a,b\geq 0}} x_{2i+1}^a (-e_{2i}e_{2i+1} + e_{2i+1}e_{2i})x_{2i+1}^b.$$

Proof This follows from Lemma 3.2.1 by induction on *n*.

Let y_1, \ldots, y_{2k} be commuting variables. We let $SSym[y_1, \ldots, y_{2k}]$ denote the subalgebra of the polynomial algebra $\mathbb{C}[y_1, \ldots, y_{2k}]$ generated by the supersymmetric power-sum polynomials

$$p_n(y_1, \dots, y_{2k}) := y_1^n + y_3^n + \dots + y_{2k-1}^n - (y_2^n + y_4^n + \dots + y_{2k}^n)$$

for all $n \ge 1$. We have an injective algebra homomorphism $SSym[y_1, \ldots, y_{2k}] \to \mathcal{A}_{2k}^{aff}$ via $y_i \mapsto x_i$. We denote the image by $SSym[x_1, \ldots, x_{2k}]$, and let p_n denote $p_n(x_1, \ldots, x_{2k})$. Let $Z(\mathcal{A}_{2k}^{aff})$ denote the center of \mathcal{A}_{2k}^{aff} .

Proposition 4.2.1 We have that $SSym[x_1, \ldots, x_{2k}] \subset Z(\mathcal{A}_{2k}^{aff})$.

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Proof We simply show that each generator of $\mathcal{A}_{2k}^{\text{aff}}$ commutes with each polynomial p_n . It is immediate that the generators z_l and x_i commute with p_n for any $n \ge 1$ by (10)(ii) and (6)(i) of Definition 3.1.1. Let [-, -] denote the commutator bracket.

For the generators e_{2i} we have

$$[p_n, e_{2i}] = (-x_{2i}^n + x_{2i+1}^n)e_{2i} - e_{2i}(-x_{2i}^n + x_{2i+1}^n) = (-x_{2i}^n + x_{2i}^n)e_{2i} - e_{2i}(-x_{2i}^n + x_{2i}^n) = 0,$$

where the first equality follows from the commuting relation (6)(*iii*) of Definition 3.1.1, and the second equality follows since $x_{2i+1}e_{2i} = x_{2i}e_{2i}$ and $e_{2i}x_{2i+1} = e_{2i}x_{2i}$ by (9)(*ii*) and (9)(*i*) of Definition 3.1.1. Similarly we have $[p_n, e_{2i-1}] = 0$.

For the generator τ_{2i} , the commuting relation (6)(ii) of Definition 3.1.1 tells us that

$$[\tau_{2i}, p_n] = \tau_{2i}(x_{2i-1}^n - x_{2i}^n + x_{2i+1}^n) - (x_{2i-1}^n - x_{2i}^n + x_{2i+1}^n)\tau_{2i}.$$

By acting on relation (*i*) of Lemma 3.2.2 by the anti-automorphism *, and rearranging, we obtain

$$\tau_{2i}x_{2i-1}^n = x_{2i+1}^n \tau_{2i} - \sum_{\substack{a+b=n-1\\a,b\ge 0}} x_{2i+1}^a (e_{2i}e_{2i-1} - 1)x_{2i-1}^b.$$

Employing this and relations (i) and (ii) of Lemma 3.2.2, we have

$$\tau_{2i}(x_{2i-1}^n - x_{2i}^n + x_{2i+1}^n) = (x_{2i-1}^n - x_{2i}^n + x_{2i+1}^n)\tau_{2i} + \sum_{\substack{a+b=n-1\\a,b\geq 0}} x_{2i-1}^a(e_{2i-1}e_{2i} - 1)x_{2i+1}^b$$
$$- \sum_{\substack{a+b=n-1\\a,b\geq 0}} x_{2i}^a(e_{2i-1}e_{2i} - e_{2i}e_{2i-1})x_{2i}^b - \sum_{\substack{a+b=n-1\\a,b\geq 0}} x_{2i+1}^a(e_{2i}e_{2i-1} - 1)x_{2i-1}^b$$

Hence showing that $[\tau_{2i}, p_n] = 0$ is equivalent to showing that the three summations above sum to zero. This follows by changing the second summation accordingly:

$$-\sum_{\substack{a+b=n-1\\a,b\geq 0}} x_{2i}^{a} (e_{2i-1}e_{2i} - e_{2i}e_{2i-1}) x_{2i}^{b} = -\sum_{\substack{a+b=n-1\\a,b\geq 0}} x_{2i}^{a} e_{2i-1}e_{2i} x_{2i}^{b} - x_{2i}^{a}e_{2i}e_{2i-1} x_{2i}^{b}$$
$$= -\sum_{\substack{a+b=n-1\\a,b\geq 0}} x_{2i-1}^{a}e_{2i-1}e_{2i} x_{2i+1}^{b} - x_{2i+1}^{a}e_{2i}e_{2i-1} x_{2i-1}^{b}$$

by repeat application of relations (9)(*i*) and (9)(*ii*) of Definition 3.1.1. One shows $[\tau_{2i+1}, p_n] = 0$ analogously.

Under the projection pr : $\mathcal{A}_{2k}^{\text{aff}} \to \mathcal{A}_{2k}$ the subalgebra SSym $[x_1, \ldots, x_{2k}]$ gets sent to SSym $[X_1, \ldots, X_{2k}]$, showing that such a subalgebra is central in \mathcal{A}_{2k} . It was shown in [3, Theorem 4.2.6] that this is in fact the whole centre of \mathcal{A}_{2k} . Note that in [3], the centre is given as SSym $[N_1, \ldots, N_{2k}]$ where

$$N_i := \begin{cases} \frac{z}{2} - 1 - X_i, & \text{if } i \text{ odd,} \\ X_i - \frac{z}{2} + 1, & \text{if } i \text{ even.} \end{cases}$$

But one can easily see that $SSym[X_1, ..., X_{2k}] = SSym[N_1, ..., N_{2k}]$. Based on this and comparing with the centers of other affine diagram algebras, see [16, Corollary 4.10] and [4, Theorem 4.2], leads to a natural conjecture for the center of \mathcal{A}_{2k}^{aff} .

Conjecture 3.2.4 $Z(\mathcal{A}_{2k}^{aff}) = \langle z_l, p_n | l, n \in \mathbb{Z}_{\geq 0} \rangle.$

3.3 Extending the Action on Tensor Spaces

We now seek to extend the action of \mathcal{A}_{2k} on $V^{\otimes k}$ to one of $\mathcal{A}_{2k}^{\text{aff}}$ on $M \otimes V^{\otimes k}$, where M is any $\mathbb{C}S(n)$ -module. The tensor space $M \otimes V^{\otimes k}$ is also viewed as an $\mathbb{C}S(n)$ -module by the diagonal action. Before extending the action, we briefly define some central elements in $\mathbb{C}S(n)$. For each $b \in [n]$ and $l \in \mathbb{N}$, we let

$$T_{n,b} := \sum_{a \in [n] \setminus \{b\}} (a, b), \text{ and } Z_{n,l} := \sum_{b \in [n]} T_{n,b}^l.$$

So $T_{n,b}$ is the sum of all transposition containing b, and $Z_{n,l}$ is the *l*-power sum in $T_{n,b}$ as b runs from 1 to n.

Lemma 3.3.1 For each $l \in \mathbb{N}$, we have that $Z_{n,l}$ belongs to the center of $\mathbb{C}S(n)$.

Proof This follows since $\pi T_{n,b} = T_{n,\pi(b)}\pi$ for any $\pi \in S(n)$.

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Theorem 3.3.2 Given any $\mathbb{C}S(n)$ -module $M = \text{Span}_{\mathbb{C}}\{m_1, \ldots, m_d\}$, we have a \mathbb{C} -algebra homomorphism

$$\mathcal{V}_{n,k}^{(M)}: \mathcal{A}_{2k}^{a\!f\!f} \to \mathsf{End}_{S(n)}(M \otimes V^{\otimes k})$$

defined on the generators by

$$\begin{split} \psi_{n,2k}^{(M)}(e_{2i-1})(m_{a_{0}}\otimes v_{a}) &= \sum_{b=1}^{n} m_{a_{0}}\otimes v_{a_{1}}\otimes \cdots \otimes v_{a_{i-1}}\otimes v_{b}\otimes v_{a_{i+1}}\otimes \cdots \otimes v_{a_{k}}, \\ \psi_{n,2k}^{(M)}(e_{2i})(m_{a_{0}}\otimes v_{a}) &= \delta_{a_{i},a_{i+1}}m_{a_{0}}\otimes v_{a}, \\ \psi_{n,2k}^{(M)}(\tau_{2i})(m_{a_{0}}\otimes v_{a}) &= \varepsilon_{a_{i},a_{i+1}}(a_{i},a_{i+1})(m_{a_{0}}\otimes v_{a_{1}}\otimes \cdots \otimes v_{a_{i-1}})\otimes v_{a_{i}}\otimes \cdots \otimes v_{a_{k}}, \\ \psi_{n,2k}^{(M)}(\tau_{2i+1})(m_{a_{0}}\otimes v_{a}) &= \varepsilon_{a_{i},a_{i+1}}(a_{i},a_{i+1})(m_{a_{0}}\otimes v_{a_{1}}\otimes \cdots \otimes v_{a_{i+1}})\otimes v_{a_{i+2}}\otimes \cdots \otimes v_{a_{k}}, \\ \psi_{n,2k}^{(M)}(x_{2i-1})(m_{a_{0}}\otimes v_{a}) &= \sum_{\substack{b=1\\b\neq a_{i}}}^{n}(b,a_{i})(m_{a_{0}}\otimes v_{a_{1}}\otimes \cdots \otimes v_{a_{i-1}})\otimes v_{a_{i}}\otimes \cdots \otimes v_{a_{k}}, \\ \psi_{n,2k}^{(M)}(x_{2i})(m_{a_{0}}\otimes v_{a}) &= \sum_{\substack{b=1\\b\neq a_{i}}}^{n}(b,a_{i})(m_{a_{0}}\otimes v_{a_{1}}\otimes \cdots \otimes v_{a_{i}})\otimes v_{a_{i+1}}\otimes \cdots \otimes v_{a_{k}}, \\ \psi_{n,2k}^{(M)}(z_{2i})(m_{a_{0}}\otimes v_{a}) &= \sum_{\substack{b=1\\b\neq a_{i}}}^{n}(b,a_{i})(m_{a_{0}}\otimes v_{a_{1}}\otimes \cdots \otimes v_{a_{i}})\otimes v_{a_{i+1}}\otimes \cdots \otimes v_{a_{k}}, \end{split}$$

for all $(a_0, \mathbf{a}) \in [d] \times [n]^k$, extended \mathbb{C} -linearly across $M \otimes V^{\otimes k}$.

Proof This can been shown by direct computations, much of which are fairly simple but lengthy. To ease notation, for any tuple $a = (a_0, a_1, ..., a_k) \in [d] \times [n]^k$, we represent a simple tensor in $M \otimes V^{\otimes k}$ by a word in the entries of a, that is $a_0a_1...a_k := m_{a_0} \otimes v_{a_1} \otimes \cdots \otimes v_{a_k}$. We begin by showing that $\psi_{n,k}^{(M)}$ is well-defined, that is to confirm that these endomorphisms do indeed commute with the diagonal action of S(n). We do this by showing for any $\pi \in S(n)$, that $\pi \psi_{n,k}^{(M)}(g)\pi^{-1} = \psi_{n,k}^{(M)}(g)$ for each generator g of $\mathcal{A}_{2k}^{\text{aff}}$.

One can deduce that $\pi \psi_{n,k}^{(M)}(e_i)\pi^{-1} = \psi_{n,k}^{(M)}(e_i)$ since the action of the generators e_i ignores the *M* component, and hence this follows from Theorem 2.4.1.

For the generators τ_{2i} ,

$$\begin{aligned} \pi \psi_{n,k}^{(M)}(\tau_{2i}) \pi^{-1}(\boldsymbol{a}) &= \pi \psi_{n,k}^{(M)}(\tau_{2i}) \Big(\pi^{-1}(a_0 a_1 \dots a_k) \Big) \\ &= \varepsilon_{a_i,a_{i+1}} \pi \Big((\pi^{-1}(a_i), \pi^{-1}(a_{i+1})) \pi^{-1}(a_0 a_1 \dots a_{i-1}) \pi^{-1}(a_i \dots a_k) \Big) \\ &= \varepsilon_{a_i,a_{i+1}} \pi (\pi^{-1}(a_i), \pi^{-1}(a_{i+1})) \pi^{-1}(a_0 a_1 \dots a_{i-1}) a_i \dots a_k \\ &= \varepsilon_{a_i,a_{i+1}}(a_i, a_{i+1})(a_0 a_1 \dots a_{i-1}) a_i \dots a_k \\ &= \psi_{n,k}^{(M)}(\tau_{2i})(\boldsymbol{a}) \end{aligned}$$

noting $\varepsilon_{\pi^{-1}(a_i),\pi^{-1}(a_{i+1})} = \varepsilon_{a_i,a_{i+1}}$. One can show $\pi \psi_{n,k}^{(M)}(\tau_{2i+1})\pi^{-1} = \psi_{n,k}(\tau_{2i+1})$ in a similar manner.

For the generators x_{2i-1} ,

$$\pi \psi_{n,k}^{(M)}(x_{2i-1})\pi^{-1}(\boldsymbol{a}) = \pi \psi_{n,k}^{(M)}(x_{2i-1}) \left(\pi^{-1}(a_0a_1\dots a_k)\right)$$

$$= \pi \left(\sum_{\substack{b \in [n] \\ b \neq \pi^{-1}(a_i)}} (b, \pi^{-1}(a_i))\pi^{-1}(a_0a_1\dots a_{i-1})\pi^{-1}(a_i\dots a_k)\right)$$

$$= \sum_{\substack{b \in [n] \\ b \neq \pi^{-1}(a_i)}} \pi(b, \pi^{-1}(a_i))\pi^{-1}(a_0a_1\dots a_{i-1})a_i\dots a_k$$

$$= \sum_{\substack{b \in [n] \\ b \neq \pi^{-1}(a_i)}} (\pi(b), a_i)(a_0a_1\dots a_{i-1})a_i\dots a_k$$

$$= \sum_{\substack{b' \in [n] \\ b' \neq a_i}} (b', a_i)(a_0a_1\dots a_{i-1})a_i\dots a_k$$

$$= \psi_{n,2k}^{(M)}(x_{2i-1})(\boldsymbol{a})$$

by the substitution $b' = \pi(b)$. One can show $\pi \psi_{n,k}^{(M)}(x_{2i})\pi^{-1} = \psi_{n,k}(x_{2i})$ in a similar manner. Lastly $\pi \psi_{n,k}^{(M)}(z_l)\pi^{-1} = \psi_{n,k}(z_l)$ can be seen since $Z_{n,l}$ are central in $\mathbb{C}S(n)$.

One now needs to confirm that the defining relations of $\mathcal{A}_{2k}^{\text{aff}}$ in Definition 3.1.1 are upheld under $\psi_{n,k}^{(M)}$. As mentioned, these can be shown by direct, but lengthy computations. With this in mind, we will only give details of some of the more difficult relations, namely relations (8) through (10). Note that the Braid-like relations (7) follow in a analogous manner to the proof of Lemma 2.4.5.

(8)(i): We seek to show that

$$\psi_{n,k}^{(M)}(x_{2i+1}) = \psi_{n,k}^{(M)}(\tau_{2i}x_{2i-1}\tau_{2i} + e_{2i}e_{2i-1}\tau_{2i} + \tau_{2i}e_{2i-1}e_{2i} - \tau_{2i}).$$

To show this we examine how each term on the hand right side acts on the simple tensor a, and show that the sum recovers the action of x_{2i+1} . It proves easier to do this by tackling two cases, when $a_i \neq a_{i+1}$ and when $a_i = a_{i+1}$.

(Case 1): Assume $a_i \neq a_{i+1}$, then for the first term we have

$$\begin{split} \psi_{n,k}(\tau_{2i}x_{2i-1}\tau_{2i})(a) &= \psi_{n,k}(\tau_{2i}x_{2i-1})\Big((a_i,a_{i+1})(a_0a_1\dots a_{i-1})a_i\dots a_k\Big) \\ &= \psi_{n,k}(\tau_{2i})\left(\sum_{\substack{b\in[n]\\b\neq a_i}} (b,a_i)(a_i,a_{i+1})(a_0a_1\dots a_{i-1})a_i\dots a_k\right) \\ &= \sum_{\substack{b\in[n]\\b\neq a_i}} (a_i,a_{i+1})(b,a_i)(a_i,a_{i+1})(a_0a_1\dots a_{i-1})a_i\dots a_k \\ &= \sum_{\substack{b\in[n]\\b\neq a_i}} ((a_i,a_{i+1})(b),a_{i+1})(a_0a_1\dots a_{i-1})a_i\dots a_k \\ &= \sum_{\substack{c\in[n]\\c\neq a_{i+1}}} (c,a_{i+1})(a_0a_1\dots a_i)a_{i+1}\dots a_k + (a_i,a_{i+1})(a_0a_1\dots a_{i-1})a_i\dots a_k \\ &= \sum_{\substack{c\in[n]\\c\neq a_{i+1}}} (c,a_{i+1})(a_0a_1\dots a_i)a_{i+1}\dots a_k + (a_i,a_{i+1})(a_0a_1\dots a_{i-1})a_i\dots a_k \\ &= \psi_{n,k}^{(M)}(x_{2i+1})(a) + \psi_{n,k}^{(M)}(\tau_{2i})(a) - (a_i,a_{i+1})(a_0a_1\dots a_i)a_{i+1}\dots a_k \end{split}$$

where we employed the substitution $c = (a_i, a_{i+1})(b)$. For the second term,

$$\psi_{n,k}^{(M)}(e_{2i}e_{2i-1}\tau_{2i})(\boldsymbol{a}) = \psi_{n,k}^{(M)}(e_{2i}e_{2i-1})\Big((a_i, a_{i+1})(a_0a_1 \dots a_{i-1})a_i \dots a_k\Big)$$
$$= \psi_{n,k}^{(M)}(e_{2i})\left(\sum_{b=1}^n (a_i, a_{i+1})(a_0a_1 \dots a_{i-1})ba_{i+1} \dots a_k\right)$$
$$= (a_i, a_{i+1})(a_0a_1 \dots a_{i-1})a_{i+1}a_{i+1} \dots a_k$$

 $= (a_i, a_{i+1})(a_0a_1 \dots a_i)a_{i+1} \dots a_k.$

For the third term $\psi_{n,k}^{(M)}(\tau_{2i}e_{2i-1}e_{2i})(a) = 0$ since $a_i \neq a_{i+1}$. Thus collectively,

$$\begin{split} \psi_{n,k}^{(M)} & (\tau_{2i}x_{2i-1}\tau_{2i} + e_{2i}e_{2i-1}\tau_{2i} + \tau_{2i}e_{2i-1}e_{2i} - \tau_{2i})^{a} \\ &= \psi_{n,k}^{(M)}(\tau_{2i}x_{2i-1}\tau_{2i})^{a} + \psi_{n,k}^{(M)}(e_{2i}e_{2i-1}\tau_{2i})^{a} + \psi_{n,k}^{(M)}(\tau_{2i}e_{2i-1}e_{2i})^{a} - \psi_{n,k}^{(M)}(\tau_{2i})^{a} \\ &= \psi_{n,k}^{(M)}(x_{2i+1})(a) + \psi_{n,k}^{(M)}(\tau_{2i})(a) - (a_{i}, a_{i+1})(a_{0}a_{1} \dots a_{i})a_{i+1} \dots a_{k} \\ &+ (a_{i}, a_{i+1})(a_{0}a_{1} \dots a_{i})a_{i+1} \dots a_{k} - \psi_{n,k}^{(M)}(\tau_{2i})^{a} \\ &= \psi_{n,k}^{(M)}(x_{2i+1})(a). \end{split}$$

(Case 2): Assume $a_i = a_{i+1}$. Then $\psi_{n,k}^{(M)}(\tau_{2i})(\boldsymbol{a}) = 0$, and so

$$\psi_{n,k}^{(M)}(\tau_{2i}x_{2i-1}\tau_{2i}+e_{2i}e_{2i-1}\tau_{2i}+\tau_{2i}e_{2i-1}e_{2i}-\tau_{2i})^{a}=\psi_{n,k}^{(M)}(\tau_{2i}e_{2i-1}e_{2i})^{a}.$$

Hence we just need to confirm that $\psi_{n,k}^{(M)}(x_{2i+1})^a = \psi_{n,k}^{(M)}(\tau_{2i}e_{2i-1}e_{2i})^a$. Well,

$$\psi_{n,k}^{(M)}(\tau_{2i}e_{2i-1}e_{2i})(a) = \psi_{n,k}^{(M)}(\tau_{2i})\left(\sum_{b=1}^{n} a_0a_1\dots a_{i-1}ba_{i+1}\dots a_k\right) = \sum_{b=1}^{n} (b, a_{i+1})(a_0a_1\dots a_{i-1})ba_{i+1}\dots a_k$$
$$= \sum_{b=1}^{n} (b, a_{i+1})(a_0a_1\dots a_i)a_{i+1}\dots a_k = \psi_{n,k}^{(M)}(x_{2i+1})^a.$$

The remaining Skein-like relations follow by employing similar arguments.

(9)(i): We seek to show $\psi_{n,k}^{(M)}(e_i x_i) = \psi_{n,k}^{(M)}(e_i x_{i+1})$. We show this first when working with e_{2i} , then with e_{2i-1} . Assume $a_i \neq a_{i+1}$, then

$$\psi_{n,k}^{(M)}(e_{2i}x_{2i})(\boldsymbol{a}) = \psi_{n,k}^{(M)}(e_{2i})\left(\sum_{\substack{b=1\\b\neq a_i}} (b, a_i)(a_0a_1\dots a_i)a_{i+1}\dots a_k\right) = (a_i, a_{i+1})(a_0a_1\dots a_i)a_{i+1}\dots a_k,$$

$$\psi_{n,k}^{(M)}(e_{2i}x_{2i+1})(a) = \psi_{n,k}^{(M)}(e_{2i})\left(\sum_{\substack{b=1\\b\neq a_{i+1}}} (b, a_{i+1})(a_0a_1 \dots a_i)a_{i+1} \dots a_k\right) = (a_i, a_{i+1})(a_0a_1 \dots a_i)a_{i+1} \dots a_k.$$

When $a_i = a_{i+1}$ one can check that $\psi_{n,k}^{(M)}(e_{2i}x_{2i})^a = \psi_{n,k}^{(M)}(e_{2i}x_{2i+1})^a = 0$, thus $\psi_{n,k}^{(M)}(e_{2i}x_{2i})^a = \psi_{n,k}^{(M)}(e_{2i}x_{2i+1})^a$. For odd indices we have

$$\psi_{n,k}^{(M)}(e_{2i-1}x_{2i-1})(\boldsymbol{a}) = \psi_{n,k}^{(M)}(e_{2i-1}) \left(\sum_{\substack{b=1\\b\neq a_i}} (b, a_i)(a_0a_1\dots a_{i-1})a_i\dots a_k \right)$$
$$= \sum_{\substack{c=1\\b\neq a_i}}^n \sum_{\substack{b=1\\b\neq a_i}} (b, a_i)(a_0a_1\dots a_{i-1})ca_{i+1}\dots a_k,$$

$$\psi_{n,k}^{(M)}(e_{2i-1}x_{2i})(\boldsymbol{a}) = \psi_{n,k}^{(M)}(e_{2i-1}) \left(\sum_{\substack{b=1\\b\neq a_i}} (b, a_i)(a_0a_1\dots a_i)a_{i+1}\dots a_k \right)$$
$$= \sum_{\substack{c=1\\b\neq a_i}}^n \sum_{\substack{b=1\\b\neq a_i}} (b, a_i)(a_0a_1\dots a_{i-1})ca_{i+1}\dots a_k.$$

Thus $\psi_{n,k}^{(M)}(e_i x_i)^a = \psi_{n,k}^{(M)}(e_i x_{i+1})^a$. Relation (9)(*ii*) may be shown in a similar manner.

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(10)(i):

$$\begin{split} \psi_{n,k}^{(M)}(e_1 x_1^l e_1)(\boldsymbol{a}) &= \psi_{n,k}^{(M)}(e_1 x_1^l) \left(\sum_{b=1}^n a_0 b a_2 \dots a_k \right) = \psi_{n,k}^{(M)}(e_1) \left(\sum_{b=1}^n (T_{n,b}^l a_0) b a_2 \dots a_k \right) \\ &= \sum_{c=1}^n \left(\sum_{b=1}^n T_{n,b}^l a_0 \right) c a_2 \dots a_k = \sum_{c=1}^n \left(Z_{n,l} a_0 \right) c a_2 \dots a_k \\ &= \psi_{n,k}^{(M)}(z_l) \left(\sum_{c=1}^n a_0 c a_2 \dots a_k \right) = \psi_{n,k}^{(M)}(z_l) \left(\psi_{n,k}^{(M)}(e_1)(a_0 a_1 a_2 \dots a_k) \right) \\ &= \psi_{n,k}^{(M)}(z_l e_1)(\boldsymbol{a}). \end{split}$$

Lastly relation (10)(ii) is simple to check since $Z_{n,l}$ belongs to the center of $\mathbb{C}S(n)$.

Corollary 3.3.3 The subalgebra $\langle z_0, \tau_i, e_j \rangle_{i,j}$ of \mathcal{A}_{2k}^{aff} is infinite dimensional over $\mathbb{C}[z_0]$.

Proof Set $d_m := x_1^m e_2 e_1$ for all $m \in \mathbb{N}$. We first show that $d_m \in \langle \tau_i, e_j \rangle$ by induction on m. By Lemma 3.1.4 (i) we see that $e_2 x_2 \in \langle \tau_i, e_j \rangle$. Then multiplying on the right by e_1 yeilds $e_2 x_2 e_1 = e_2 x_1 e_1 = x_1 e_2 e_1 = d_1$, where the first equality follows from (9)(ii) of Definition 3.1.1, and the second from (6)(iii). Thus we have the base case $d_1 \in \langle \tau_i, e_j \rangle$. Assume $d_{m'} \in \langle \tau_i, e_j \rangle$ for all m' < m with $m \ge 2$, we seek to show that $d_m \in \langle \tau_i, e_j \rangle$. Well

$$d_{m-1}\tau_2 e_1 = x_1^{m-1} e_2 e_1 \tau_2 e_1 = x_1^{m-1} e_2 x_2 e_1 = x_1^{m-1} e_2 x_1 e_1 = x_1^m e_2 e_1 = d_n,$$

where the second equality follows from Lemma 3.1.4 (i), and the remaining equalities follow in the same manner as the base case. Hence $d_m \in \langle \tau_i, e_j \rangle$ completing induction. We now seek to show that the set $\{d_m \mid m \in \mathbb{N}\}$ is $\mathbb{C}[z_0]$ -linearly independent in $\mathcal{A}_{2k}^{\text{aff}}$, which will complete the proof. Let $I \subset \mathbb{N}$ be finite and assume

$$\sum_{m\in I}h_m(z_0)d_m=0,$$

where $h_m(z_0)$ are polynomials in $\mathbb{C}[z_0]$. We seek to show that $h_m(z_0) = 0$ for each $m \in I$. Let $M \in I$ be the maximal element, and let R be the set of roots for each $h_m(z_0)$. Pick an $n \in \mathbb{N}$ such that n > M + 1 and $n \notin R$. Let F be any free $\mathbb{C}S(n)$ -module. For any $f \in F$ and $(a_1, \ldots, a_k) \in [n]^k$, we have

$$\psi_{n,k}^{(F)}(d_m)(f \otimes v_{a_1} \otimes v_{a_2} \otimes \cdots \otimes v_{a_k}) = (T_{n,a_2}^m f) \otimes v_{a_2} \otimes v_{a_2} \otimes v_{a_3} \otimes \cdots \otimes v_{a_k}.$$

Since *F* is free, it will follow that the set $\{\psi_{n,k}^{(F)}(d_m) \mid m \in I\}$ is linear independent in $\operatorname{End}_{S(n)}(F \otimes V^{\otimes k})$ if the set $\{T_{n,a_2}^m \mid m \in I\}$ is linearly independent in $\mathbb{C}S(n)$. This follows since n > M + 1, and hence T_{n,a_2}^m contains a permutation consisting of a single cycle of size m + 1, while all permutations in T_{n,a_2}^m must have smaller support whenever m' < m. Now consider the equation

$$\psi_{n,k}^{(F)}\left(\sum_{m\in I}h_m(z_0)d_m\right) = \sum_{m\in I}h_m(n)\psi_{n,k}^{(F)}(d_m) = 0.$$

Since *n* is not a root of any $h_m(z_0)$, and the set $\{\psi_{n,k}^{(F)}(d_m) \mid m \in I\}$ is linear independent, we must have that $h_m(z_0) = 0$ for each $m \in I$.

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4 Connections with the Heisenberg Category

J. Brundan and M. Vargas recently defined in [2] an affine partition category APar as a monoidal subcategory of the Heisenberg category introduced by Khovanov in [12] generated by certain objects and morphisms. This was based on the observation made by S. Likeng and A. Savage in [14] that the partition category can be realised inside the Heisenberg category. This affine partition category naturally gives rise to another definition of an affine partition algebra, which they denote by AP_k by taking the endomorphism algebra $End_{APar}((\uparrow \downarrow)^k)$ for the object $(\uparrow \downarrow)^k$ in APar (see Section 4.1. below).

Inspired by the work of Brundan and Vargas, we construct a surjective homomorphism φ from $\mathcal{A}_{2k}^{\text{aff}}$ to $\text{End}_{\text{Heis}}((\uparrow\downarrow)^k)$. In fact, our argument generalises to show that Brundan and Vargas' affine partition category APar is the *full* monoidal subcategory in Heis generated by the object $\uparrow\downarrow$. As a corollary we obtain that AP_k is a quotient of $\mathcal{A}_{2k}^{\text{aff}}$.

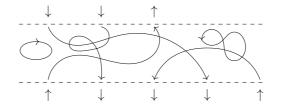
We start by recalling the definition of the Heisenberg category.

4.1 Heisenberg Category

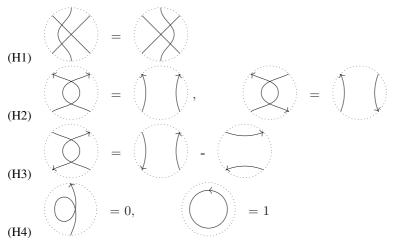
The Heisenberg Category Heis is a \mathbb{C} -linear monoidal category originally defined by M. Khovanov in [12]. The objects of Heis are generated, as a monoidal category, by the two objects \uparrow and \downarrow . We use juxtaposition to denote the tensor product of objects, and the monoidal identity object is the empty word \emptyset . Hence we view the free monoid $\langle\uparrow,\downarrow\rangle$ as the set of objects in Heis. Consider two objects $a = a_1 \cdots a_n$ and $b = b_1 \cdots b_m$ for $a_i, b_i \in \{\uparrow,\downarrow\}$. The space of morphisms $\text{Hom}_{\text{Heis}}(a, b)$ is the \mathbb{C} -vector space generated by certain diagrams modulo local relations. We call such diagrams (a, b)-diagrams and define them as follows: Firstly, we work in the strip $\mathbb{R} \times [0, 1]$ with boundary $\mathbb{B} := \mathbb{R} \times \{1\} \cup \mathbb{R} \times \{0\}$. We call an orientated immersion of the interval [0, 1] and circle S^1 a *string* and *loop* respectively. We denote orientations by drawing an arrow on the curve. Now consider the set of points $E = [n] \times \{1\} \cup [m] \times \{0\}$, and colour $(i, 1) \in [n] \times \{1\}$ and $(j, 0) \in [m] \times \{0\}$ with the symbols a_i and b_j respectively. We say that a set partition of E into pairs is an (a, b)matching if pairs of points in the same row are coloured by opposite arrows, while pairs of points in different rows are coloured by the same arrow. Then an (a, b)-diagram is a finite collection of strings and loops, modulo rel boundary isotopies, such that:

- (D1) The endpoints of the strings induce an (a, b)-matching on E
- (D2) There are only finitely many points of intersection, and no triple or tangential intersections occur
- (D3) The boundary \mathbb{B} doesn't intersect any loops, and only intersects strings at the endpoints *E*

For example let $a = \downarrow \downarrow \uparrow$ and $b = \uparrow \downarrow \downarrow \downarrow \uparrow$, then



is a (a, b)-diagram. Isotopic deformation of the interior of $\mathbb{R} \times [0, 1]$ is allowed, and will preserve the relative structure of the points of intersection. If a loop contains no intersections we call it a *bubble*. Bubbles can have clockwise or anticlockwise orientation. If the endpoints of a string occur in different rows we call it a *vertical* string, and it has either a *down* or *up* orientation. If the endpoints belong to the same row then we call it an *arc*. Non self-intersecting arcs have either a *clockwise* or *anti-clockwise* orientation. In the above example there are two loops, one of which is a bubble, and four strings, three of which are vertical and one an arc. We call an endpoint of a string a *source* if the arrow of orientation points away from it, and a *target* otherwise. We consider (a, b)-diagrams modulo the following local relations:



Relation (H1) holds regardless of orientations. To apply such a local relation to an (a, b)diagram one locates a disk which is isotopic to one of the disks above, then replace such a disk according to the corresponding equation. Note that any of the local relations may be rotated in any way to give an equivalent relation. Relation (H1) tells us that any curve may past over a crossing, and relations (H2) and (H3) tells us how to pull part orientated curves, where (H3) shows that this can not always be done for free. Relation (H4) tells us that left curls kill (a, b)-diagrams, and that any anti-clockwise bubble may be removed for free.

The composition of morphisms is given by vertical concatenation of diagrams, and rescaling (and extending \mathbb{C} -linearly). We denote composition by juxtaposition of symbols. When a = b we write *a*-diagram instead of (a, a)-diagram. The morphism space End_{Heis}(a) is a \mathbb{C} -algebra with identity given by the diagram of non-intersecting vertical strings. Now for later use, we collect some relations regarding arbitrary (a, b)-diagrams. The following local relation follows from (H2) and (H3), see also [14, (3.5)]:

Lemma 4.1.1 Clockwise bubbles satisfy the commuting relation

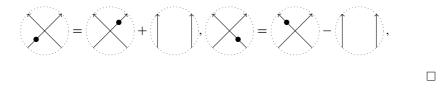
$$\uparrow \quad \downarrow \bigcirc = \bigcirc \uparrow \quad \downarrow$$

Although left curls annihilate diagrams, right curls do not, and they play an important role. We will represent right curls by a decoration, and label such decorations with weights to denote multiplicity:



The following result is a simple application of the local relations.

Lemma 4.1.2 *The following two local relations hold:*

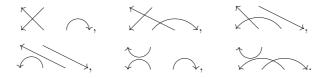


We now recall a basis for the morphism spaces $\text{Hom}_{\text{Heis}}(a, b)$ presented in [12]. We first introduce a few definitions to help us describe this basis in a manner which will lend itself better for later results.

Definition 4.1.3 For $a, b \in \langle \uparrow, \downarrow \rangle$, we say an (a, b)-diagram is *simple* if it contains no loops, no self-intersections, and two strings intersect at most once. Let Sim(a, b) denote the set of simple (a, b)-diagrams, and write Sim(a) for Sim(a, a).

Given words $a = a_1 \cdots a_n$, $b = b_1 \cdots b_m \in \langle \uparrow, \downarrow \rangle$ with $a_i, b_j \in \{\uparrow, \downarrow\}$, let b^* denote the word obtained from b by replacing up arrows with down arrows, and down arrows with up arrows. Let u equal the number of up arrows appearing in a and b^* , and d the number of a down arrows. Then by (D1), one can deduce that $\text{Hom}_{\text{Heis}}(a, b)$ is non-empty if and only if u = d. In this situation we have that |Sim(a, b)| = u!, since there is one simple (a, b)diagram for every (a, b)-matching. Such a correspondence is given by reading the pairings of endpoints formed from the strings of a simple diagram.

Example 4.1.4 Consider the words $a = \uparrow \downarrow$ and $b = \downarrow \uparrow \uparrow \downarrow$. Then the 6 = 3! simple (a, b)-diagrams are



These diagrams are in a one-to-one correspondence with the (a, b)-matchings of the set of endpoints $E = \{(i, 1), (j, 0) \mid i \in [2], j \in [4]\}$. As an example, for the first diagram above we have

$$\left\{ \{(1,1),(2,0)\},\{(2,1),(1,0)\},\{(3,0),(4,0)\} \right\}.$$

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The basis for $Hom_{Heis}(a, b)$ we describe below is obtained by adding decorations (right curls) and decorated clockwise loops to all the simple (a, b)-diagrams in a particular manner. We describe this by introducing some basic diagrams and using the composition of diagrams.

Definition 4.1.5 Let $a = a_1 \cdots a_n \in \langle \uparrow, \downarrow \rangle$ for $a_i \in \{\uparrow, \downarrow\}$. For $i \in [n], w \in \mathbb{Z}_{\geq 0}$, define the *a*-diagrams

$$r_i = \left| \begin{array}{ccc} a_1 & a_{i-1} & a_i & a_{i+1} & a_n \\ \vdots & \vdots & \vdots & \vdots \\ \mathbf{r}_i & \mathbf{r}_i & \mathbf{r}_i \end{array} \right| \left| \begin{array}{ccc} a_1 & a_n \\ \vdots & \vdots & \vdots \\ \mathbf{r}_i & \mathbf{r}_i \end{array} \right| \left| \begin{array}{ccc} a_1 & a_n \\ \vdots & \vdots & \vdots \\ \mathbf{r}_i & \mathbf{r}_i \end{array} \right|$$

The orientation of strings is taken to match a. Although both r_i and c_w depend on a, we surpress this fact as it should be clear from context.

Definition 4.1.6 Given any $a = a_1 \cdots a_n$, $b = b_1 \cdots b_m \in \langle \uparrow, \downarrow \rangle$ for $a_i, b_j \in \{\uparrow, \downarrow\}$, let B(a, b) be the set of (a, b)-diagrams of the form

$$c_w^{k_w}\cdots c_1^{k_1}c_0^{k_0}r_1^{s_1}\cdots r_n^{s_n}\alpha r_1^{t_1}\cdots r_m^{t_m}$$

where $\alpha \in Sim(a, b)$, $w, k_l, s_i, t_j \in \mathbb{Z}_{\geq 0}$, and $s_i = t_j = 0$ whenever (i, 1) and (j, 0) are sources respectively. We write B(a) for B(a, a).

Example 4.1.7 Given $a = \uparrow \downarrow$ and $b = \downarrow \uparrow \uparrow \downarrow$, an example of an element of B(a, b) is

$$c_5 c_0^2 r_1^3 \alpha r_4^2 = \underbrace{\bigcirc}_{5} \underbrace{\bigcirc}_{5} \underbrace{\bigcirc}_{5} \underbrace{\bigcirc}_{6} \underbrace{\bigcirc}_{7} \underbrace{$$

where α is the third simple (a, b)-diagram in the list given in 4.1.4.

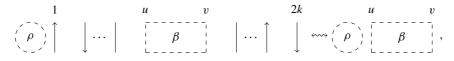
The following result is *Proposition 5* of [12].

Theorem 4.1.8 The set B(a, b) is a basis for $Hom_{Heis}(a, b)$.

Remark 4.1.9 The description of this basis is analogous to the basis given by regular monomials presented in [16, Theorem 4.6]. Note that no decorated anti-clockwise bubbles appear. This is due to the fact that any decorated anti-clockwise bubble may be expressed as a linear combination of decorated clockwise bubbles, see for example [12, Proposition 2].

4.2 A surjective homomorphism $\varphi : \mathcal{A}_{2k}^{\text{aff}} \to \text{End}_{\text{Heis}}((\uparrow\downarrow)^k)$

When drawing a $(\uparrow\downarrow)^k$ -diagram, instead of labelling the endpoints with arrows, we instead will label the points (i, 1) with *i* for each $1 \le i \le 2k$, since the parity of the label recovers the orientation of the arrow. Also to ease notation we employ the following diagrammatic shorthand for elements of B($(\uparrow\downarrow)^k$):

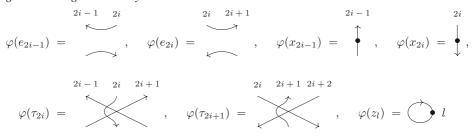


where β is loopless, ρ is a collection of (possibly decorated) clockwise bubbles, and $u, v \in [2k]$. Hence we drop the trivial vertical strings but retain the labels u through v, allowing one to recover the original diagram.

Proposition 4.2.1 We have a C-algebra homomorphism

$$\varphi: \mathcal{A}_{2k}^{a\!f\!f} \to \mathsf{End}_{\mathsf{Heis}}((\uparrow \!\!\downarrow)^k)$$

given on the generators by



Proof This will be shown by directly checking that each of the defining relations in Definition 3.1.1 is satisfied under the map φ . Most of these are simple to check but lengthy, hence for such relations we do not give full details.

(1)(i):

$$\varphi(\tau_{2i}^2) = \begin{array}{c} 2i-1 & 2i & 2i+1 \\ \hline & & & \\ 2i-1 & 2i & 2i+1 \\ = \\ 1 & & \\ 2i-1 & 2i & 2i+1 \\ 2i-1 & 2i & 2i+1 \\ = \\ 1 & & \\ 2i-1 & 2i & 2i+1 \\ = \\ 1 & & \\ 1 &$$

which equals $\varphi(1 - e_{2i})$. One can show that relation (1)(ii) is upheld in a similar manner.

(2): Relation (2)(*i*) is $\tau_{2i+1}\tau_{2j} = \tau_{2j}\tau_{2i+1}$ for all $j \neq i + 1$. When $j \neq i$, it is clear to see diagrammatically that this relation is upheld under φ . For case j = i, one applies (H1) and then (H2) to see that

Both relations (2)(*ii*) and (2)(*iii*) can be seen to hold under φ diagrammatically. For relation (2)(*iv*), we have that

$$\varphi(s_i) = \varphi(\tau_{2i+1}\tau_{2i} + e_{2i}) = \begin{array}{cccc} 2i - 1 & 2i & 2i - 1 & 2i & 2i + 12i + 2 \\ & & & & & & \\ \varphi(s_i) = \varphi(\tau_{2i+1}\tau_{2i} + e_{2i}) = \begin{array}{cccc} & & & & & \\ & & & & \\ & & &$$

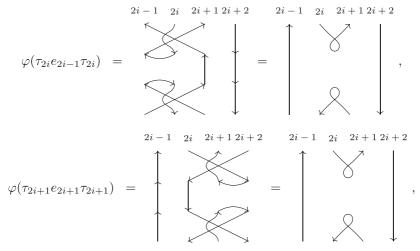
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Such elements satisfy the braid relation $s_i s_{i+1} s_i = s_{i+1} s_i s_{i+1}$ by [14, Theorem 4.1].

(3): Relation (3)(*i*) is upheld under φ by applying Lemma 4.1.1, and (3)(*ii*) is upheld by (H4). Relations (3)(*iii*) and (3)(*iv*) are upheld by the fact that left curls are annihilated. For relation (3)(*v*), it is clear that applying (H1) allows one to go from the diagram $\varphi(\tau_{2i}e_{2i-1}e_{2i+1})$ to $\varphi(\tau_{2i+1}e_{2i-1}e_{2i+1})$, and similarly for relation (3)(*vi*).

(4): All of these relations follow diagrammatically and from (3)(iii) and (3)(iv).

(5): Relation (5)(i) is simple to check since the diagrams contain no points of intersection. For (5)(ii), applying (H1) and (H2) we see that

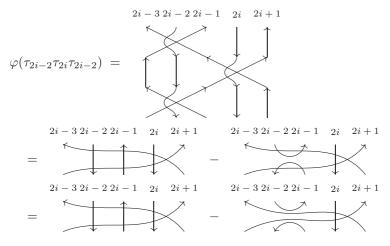


thus $\varphi(\tau_{2i}e_{2i-1}\tau_{2i}) = \varphi(\tau_{2i+1}e_{2i+1}\tau_{2i+1})$. In a similar manner, for relation (5)(*iii*) one can show,

$$\varphi(\tau_{2i}e_{2i-2}\tau_{2i}) = \varphi(\tau_{2i-1}e_{2i}\tau_{2i-1}).$$

(6): These relations are immediately seen to be upheld diagrammatically.

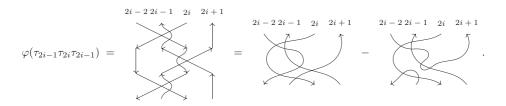
(7)(*i*): We seek to show $\varphi(\tau_{2i-2}\tau_{2i}\tau_{2i-2}) = \varphi(\tau_{2i}\tau_{2i-2}\tau_{2i}(1-e_{2i-2}))$. The left hand side gives



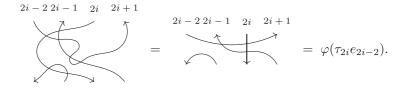
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where the second equality follows by applying (H3), and the third equality follows from (H2). By applying (H1) and (H2), one can check that the first term above is $\varphi(\tau_{2i}\tau_{2i-2}\tau_{2i})$ and the second term above is $\varphi(\tau_{2i}\tau_{2i-2}\tau_{2i}e_{2i-2})$, hence (7)(*i*) holds. Relation (7)(*ii*) can be shown in an analogous manner.

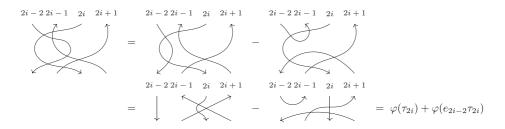
(7)(*iii*): We seek to show that $\varphi(\tau_{2i-1}\tau_{2i}\tau_{2i-1}) = \varphi(\tau_{2i} - e_{2i-2}\tau_{2i} - \tau_{2i}e_{2i-2})$. The left hand side gives



By applying (H2) twice and (H1), the second term above straightens out to



For the first term we get



where the first equality follows by applying (H3), and the second equality by (H1) and (H2). Therefore collectively we have show (7)(iii). Relation (7)(iv) follows in an analogous manner.

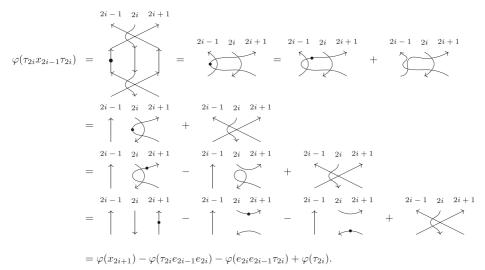
(8)(i): We seek to show that

$$\varphi(x_{2i+1}) = \varphi(\tau_{2i}x_{2i-1}\tau_{2i}) + \varphi(e_{2i}e_{i-1}\tau_{2i}) + \varphi(\tau_{2i}e_{2i-1}e_{2i}) - \varphi(\tau_{2i}).$$
(11)

One can check that

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By (H1), (H2), (H3), and applying Lemma 4.1.2 (and a 90° clockwise rotation of Lemma 4.1.2), we have



Rearranging yields (11). The remaining Skein-like relations (8)(ii), (8)(iii), and (8)(iv), following in a similar manner where we employ Lemma 4.1.2 to pull the decoration over various oriented crossings.

(9) and (10): These relations are immediately seen to be upheld diagrammatically. \Box

The remainder of this section seeks to show that the algebra homomorphism φ in the above proposition is surjective. Firstly, from Theorem 4.1.8 we know that $\operatorname{End}_{\operatorname{Heis}}((\uparrow\downarrow)^k)$ has a basis given by

$$c_w^{k_w} \dots c_1^{k_1} c_0^{k_0} r_1^{s_1} r_3^{s_3} \dots r_{2k-1}^{s_{2k-1}} \alpha r_2^{t_2} r_4^{t_4} \dots r_{2k}^{t_{2k}}$$

where $\alpha \in \text{Sim}((\uparrow\downarrow)^k)$. Since $\varphi(z_l) = c_l$ and $\varphi(x_i) = r_i$, to prove that φ is surjective it is enough to show that $\text{Sim}((\uparrow\downarrow)^k) \subset \text{Im}(\varphi)$. We will prove that $\text{Sim}((\uparrow\downarrow)^k) \subset \langle \varphi(e_i), \varphi(\tau_j) \rangle_{i,j} \subset \text{Im}(\varphi)$. We say that a simple diagram is *planar* if no intersections occur among its strings, for example the diagrams $\varphi(e_i)$ are all planar for each $i \in [2k - 1]$. The total number of planar diagrams in $\text{Sim}((\uparrow\downarrow)^k)$ is C_{2k} , the 2k-th Catalan number. These diagrams are precisely oriented versions of the Temperley-Lieb diagrams. The Jones normal form gives a way of writing the Temperley-Lieb diagrams as a product of generators (see [9], and also [11, Theorem 4.3 and Figure 16]) which does not involve bubbles, and so may be applied here for the elements $\varphi(e_i)$ to show that any planar diagram belongs to $\langle \varphi(e_i) \rangle_i$ and hence to $\text{Im}(\varphi)$.

Definition 4.2.2 Let $\pi \in S(k)$. Then we define the following simple $(\uparrow \downarrow)^k$ -diagrams:

- (i) π^{\uparrow} by pairings of endpoints {(2*i* 1, 0), (2 π (*i*) 1, 1)} and {(2*i*, 0), (2*i*, 1)} for each $1 \le i \le k$.
- (ii) π^{\downarrow} by pairings of endpoints {(2*i* 1, 0), (2*i* 1, 1)} and {(2 $\pi(i)$, 0), (2*i*, 1)} for each $1 \le i \le k$.

Example 4.2.3 For k = 3 and $\pi = (1, 2, 3) \in S(3)$, we have

$$\pi^{\uparrow} = \begin{array}{c} & & \\ & &$$

For any $\pi \in S(k)$, it is shown in [21] that we have a reduced expression of the form

$$\pi = (s_{m_1}s_{m_1+1}\cdots s_{n_1})(s_{m_2}s_{m_2+1}\cdots s_{n_2})\cdots (s_{m_l}s_{m_l+1}\cdots s_{n_l}),$$

where $k > n_1 > n_2 > \cdots > n_l$ and $n_i \ge m_i$. Noting that $s_i^{\downarrow} = \varphi(\tau_{2i+1})$, consider

$$\alpha^{\downarrow}(w) := (s_{m_1}^{\downarrow}s_{m_1+1}^{\downarrow}\cdots s_{n_1}^{\downarrow})(s_{m_2}^{\downarrow}s_{m_2+1}^{\downarrow}\cdots s_{n_2}^{\downarrow})\cdots (s_{m_l}^{\downarrow}s_{m_l+1}^{\downarrow}\cdots s_{n_l}^{\downarrow}) \in \mathsf{Im}(\varphi).$$

Strings in $\alpha^{\downarrow}(w)$ may intersect one another more than once, but we can resolve such double crossings by pulling strings apart via the local relations. The descending condition on the indices in this reduced expression means we will never need to employ (H3) to pull strings apart, and thus we must have that $\alpha^{\downarrow}(w) = \pi^{\downarrow}$. Hence $\pi^{\downarrow} \in \text{Im}(\varphi)$. Rotating π^{\downarrow} by 180° yields $(\rho \pi \rho^{-1})^{\uparrow}$ where ρ is the product of transposition (i, k - i + 1) for each $i \in [k]$. Thus we also have that $\pi^{\uparrow} \in \text{Im}(\varphi)$ for all $\pi \in S(k)$.

To aid upcoming proofs we define a collection of diagrams which loosen the conditions on simple diagrams.

Definition 4.2.4 We call an (a, b)-diagram *semisimple* if the following hold:

- (1) It contains no loops or self intersections.
- (2) No top arc intersects a bottom arc.

Let SSim(a, b) denote the set of semisimple (a, b)-diagrams, and write SSim(a) for SSim(a, a).

From definitions we have that $Sim(a, b) \subset SSim(a, b)$. Any diagram $\alpha \in SSim((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ contains precisely k + l strings, and the endpoints of these strings induce an $((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ -matching of the endpoints *E*. We let $\overline{\alpha}$ denote the unique simple diagram corresponding to such a matching (recalling the discussion after Definition 4.1.3).

Lemma 4.2.5 Given any simple diagram $\alpha \in \text{Sim}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$, there exists $\pi \in S(k)$, $\sigma \in S(l)$, and a planar diagram $\beta \in \text{Sim}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ such that $\pi^{\uparrow}\beta\sigma^{\downarrow}$ is semisimple and

$$\alpha = \pi^{\uparrow}\beta\sigma^{\downarrow}.$$

Proof Given any simple diagram $\gamma \in \text{Sim}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ let $(2i, 0), (2j, 0) \in E$ (respectively (2i - 1, 1), (2i - j, 1)) be two \downarrow (respectively \uparrow) endpoints in the bottom row (respectively top row) of γ . Let γ' be the simple diagram obtained from γ by permuting these two endpoints around. It can be seen that $\gamma(i, j)^{\downarrow}$ (respectively $(i, j)^{\uparrow}\gamma$) is semisim-

ple as long as the permutation doesn't swap the orianetation of an arc around, since that is the only way a self intersection can occur. In this situation, one can see that

$$\gamma' = \overline{\gamma(i, j)^{\downarrow}} \Big(\text{respectively} = \overline{(i, j)^{\uparrow} \gamma} \Big).$$

Hence to prove this lemma it is enough to show that we can reach a planar diagram β from α by repeatably permuting the endpoints in the bottom row coloured by \downarrow , and top row coloured by \uparrow , in such a way that the orientations of arcs are preserved. We focus on the bottom row, where the top row will follow in the same manner by a 180° rotation of the diagrammatics. Starting with α we remove intersections one at a time by employing a suitable permutation of endpoints. There are a few cases to consider, and in each such case the endpoints of the strings in the following diagrams will be arbitrary:

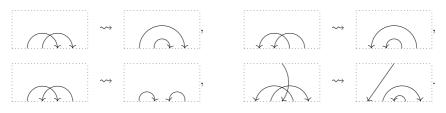
(Case 1): Crossing of two down strings:



(Case 2): Crossing of a down string with an clockwise/anti-clockwise arc:

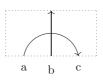


Note in either situation the orientation of the arc is preserved by the permutation. (Case 3): Crossing of two arcs: There are four cases based on the orientations of the two arcs given by

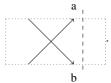


noting that in the last case such a down string must exist. Again, the orientations of the arcs are preserved under the permutation of endpoints in all four of the above situations.

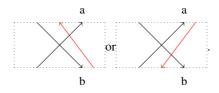
In all three of the cases above, it can be seen that the new simple diagram we obtain after the permutation of endpoints has strictly less number of intersections. We claim that applying the moves above on the bottom row, and their 180° counterparts on the top row, until all such intersections are removed will yield a planar diagram. For contradiction, suppose this is not the case. Thus even after removing all such types of intersections, the diagram still contains some other type of intersection. The other such intersections are either between an up string and arc on the bottom row, a down string and arc on the top row, or an up string and a down string. The former two are 180° counterparts to one another, hence we only need to consider one such type. Firstly, if an up string intersects a clockwise arc on the bottom we have



Note that the parity of the number of endpoints on the bottom row strictly between a and b must be different to the partity of endpoints strictly between b and c. Thus on can deduce that such an endpoint must be a target to a string which intersects the arc, and such an intersection would be accounted for by Case 2 or 3, hence a contradiction. The same argument can be used to show that the case of an up string intersecting an anti-clockwise arc on the bottom is also impossible. Note all intersections involving arcs have now been accounted for. Lastly assume an up string intersects a down string. We have two cases, one of which is



The dashed vertical line is simply an aid for arguments to come, and has been draw so that the endpoints a and b are the closest endpoints to its left. The other case is given by rotating the above by 180° and will follow analogously. The parity of the number of endpoints to the right of a is odd, while the parity of the number of endpoints to the right of b is even. This implies that there exists a string s such that one of its endpoints belongs to the right of the dashed line, while the other belongs to the left. Moreover, since the right-most endpoint on the top and bottom row are coloured by \downarrow , we can say that the endpoint of s which is to the right of the dashed line is a source while the endpoint to the left is a target. So s must intersect one of the above strings, and must be a vertical string since all intersections with arcs are accounted for. Hence, colouring the string s in red we have

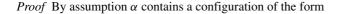


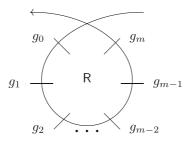
Note that *s* may intersect both of the other strings and not just one, but it is always forced to intersect the string depicted. As such each situation exhibits an intersection accounted for in Cases 1 (or its 180° counterpart), giving the desired contradiction. Thus removing all intersections of the types presented in Cases 1 to 3 (and their 180° rotated counterparts) will result in a planar diagram, completing the proof.

Let R be an open subspace of $\mathbb{R} \times [0, 1]$ and let α be an (a, b)-diagram. Examining α locally in R will give a configuration of curve segments, and we refer to such as a *region* of α . Within a given region we treat distinct curve segments as different curves, even if in α itself the two segments belong to the same curve. In particular, if in R two distinct curve segments intersect one another, and in α these two segments belong to the same curve, we will not call such an intersection a self-intersection in R, but it is a self-intersection in α .

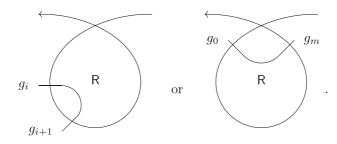
Recall that the local relation (H4) tells us that if a left curl appears in a diagram we may annihilate such a diagram. This relation asks that the region enclosed in the curl is absent of any other strings. The following result shows that even if such a region is non-empty, as long as its contains no loops or self-intersections, we can annihilate the diagram.

Lemma 4.2.6 Let α be an (a, b)-diagram containing a left curl where the region bounded by the curl contains no loops or self-intersecting curve segments, then $\alpha = 0$.



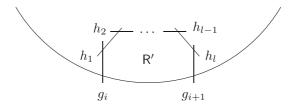


where we let *R* denote the interior region bounded by the curl, which contains no loops or self-intersecting curve segments, and g_0, \ldots, g_m account for all the intersections which occur on the curl. Note we have only drawn the segments of the g_i 's which realise the intersection on the curl. We prove the result by induction on the number of intersections occurring in R. Assume that no intersections occur in R, hence R gives a planar configuration of strings. One can deduce that there exists neighbours $g_i \mod(m+1)$ and $g_{(i+1) \mod(m+1)}$ such that either

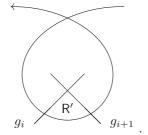


In the former situation, since we are dealing with a left curl, one can check that regardless of the orientation of the depicted string in R, it may be pulled outside the curl by (H2). For the latter situation we may employ (H1) to pull the string out of the curl over the crossing at the top. Continually pulling out such strings one at a time will result in making R empty, and then applying (H4) gives $\alpha = 0$.

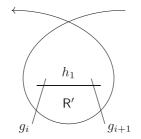
Now suppose that the result holds whenever R contains n or less intersections for some $n \ge 0$, and assume that R contains n + 1 intersections. It is clear that there must exist an empty region R' in R bounded by the curl and various segments. Diagrammatically we have



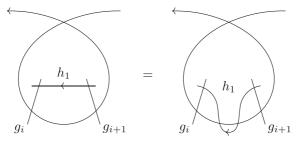
where g_i , g_{i+1} , and the (possibly empty) set of curve segments $H = \{h_1, \ldots, h_l\}$ make up the remainder of the boundary of R'. Note such curve segments may not be pairwise distinct in R. In the case when H is empty, we simply have the situation



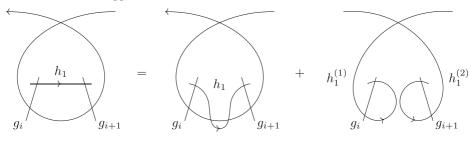
Since R' is empty we may pull this crossing out of the curl by (H1), which will decrease the number of intersection in R and thus by induction $\alpha = 0$. Hence we may assume that H is non-empty. The general case $H = \{h_1, \ldots, h_l\}$ is solved by focusing on h_1 , and in fact solving the case $H = \{h_1\}$ is sufficient to understand the general case, hence we only prove this case. So we are working with the situation



There are two cases to consider based on the orientation of h_1 . For the first case we have



by (H2). Then we may pull the crossing between either g_i and h_1 , or g_{i+1} and h_1 out of the curl by (H1), which will decrease the number of intersections in R by one and so $\alpha = 0$ by induction. With the opposite orientation on h_1 we have



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by (H3). Here denote the first diagram on the right of the above equation by α_1 and the second by α_2 . For α_1 , as was done in the previous case we may pull one of the crossings outside of the curl, and thus decrease the number of intersections in R by one, and hence $\alpha_1 = 0$ by induction. For α_2 the curve containing h_1 and the original left curl have been turned into the two new curves $h_1^{(1)}$ and $h_1^{(2)}$. Note the original left curl is no longer present, but regardless of how the original curve containing h_1 intersected the curl, at least one of the new curves $h_1^{(1)}$ and $h_1^{(2)}$ must form a new, smaller, left curl. The region bounded by this new curl is a subregion of R containing strictly less number of intersections. Hence by induction $\alpha_2 = 0$, and so collectively $\alpha = \alpha_1 + \alpha_2 = 0$ completing the proof by induction. Note the general case for $H = \{h_1, \ldots, h_l\}$ is tackled in the exact same manner by pulling h_1 out of the curl, the diagrammatics are just more cluttered, but the remaining segments h_2, \ldots, h_l do not interfer with the above argumenets.

Let $\boldsymbol{a} = a_1 \cdots a_k$ and $\boldsymbol{b} = b_1 \cdots b_l$ for $a_i, b_i \in \{\uparrow, \downarrow\}$, and consider the map deg : SSim $(\boldsymbol{a}, \boldsymbol{b}) \to \mathbb{Z}_{\geq 0} \times \mathbb{Z}_{\geq 0}$ given by deg $(\alpha) = (A(\alpha), C(\alpha))$ where $A(\alpha)$ is the number of arcs in α , and $C(\alpha)$ is the number of clockwise arcs in α . We order the degrees by using the lexicographical ordering < on $\mathbb{Z}_{\geq 0} \times \mathbb{Z}_{\geq 0}$. Note that for any $\alpha \in$ SSim $(\boldsymbol{a}, \boldsymbol{b})$ we have deg $(\overline{\alpha}) = \text{deg}(\alpha)$.

Proposition 4.2.7 *Let* $\alpha \in SSim(a, b)$ *. Then*

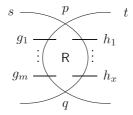
$$\alpha = \overline{\alpha} + \sum_{\substack{\beta \in \operatorname{Sim}(a, b) \\ \operatorname{deg}(\beta) > \operatorname{deg}(\alpha)}} c_{\beta}\beta$$

where $c_{\beta} \in \mathbb{Z}$.

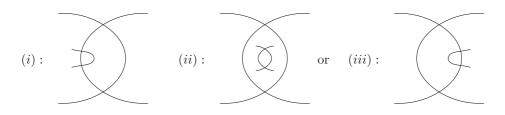
Proof Given two distinct strings *s* and *t* in α , let *n* be the number of intersections occurring between the two strings. If *n* is even set $\mu(\{s, t\}) = n$, while if *n* is odd set $\mu(\{s, t\}) = n-1$. Note $\mu(\{s, t\})$ is always even. We let

$$\eta(\alpha) = \sum_{s,t} \mu(\{s,t\}),$$

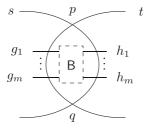
where the sum runs over all unordered pairs of distinct strings of α . Informally, $\eta(\alpha)$ is the number of intersections of α which prevent it from being simple, in particular $\eta(\alpha) = 0$ if and only if $\alpha \in \text{Sim}(a, b)$. We will prove this proposition by induction on $\eta(\alpha)$, where the base case of $\eta(\alpha) = 0$ follows immediately since $\alpha = \overline{\alpha}$. Assume the result holds for all $\alpha \in \text{SSim}(a, b)$ such that $\eta(\alpha) < n$ for some n > 0. Now let α be such that $\eta(\alpha) = n$. Pick two strings *s* and *t* in α such that $\mu(\{s, t\}) \ge 2$. Order the points of intersections between *s* and *t* according to when they appear as one travels from the source of *s* to its target. Under this ordering pick two neighbouring points of intersection *p* and *q*. Then diagrammatically we have a configuration of strings of the form



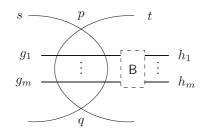
where R is the interior region bounded by the curve segments of *s* and *t* between the points of intersection *p* and *q*, and the two (possible empty) sets of string segments $G = \{g_1, \ldots, g_m\}$ and $H = \{h_1, \ldots, h_x\}$ account for all the intersections of the boundary of R through *t* and *s* respectively. We may assume that we are not in one of the following three situations:



since otherwise we may pick the more nested pair of intersections to work with instead. Since situations (*i*) and (*iii*) are not present, any string segment g_i must connect to a h_j (rather than another segment in *G*). Hence m = x and R realises a pairing of the string segments *G* with *H*. Diagrammatically we have



where B is some permutation connecting segments in G with those in H. Moreover, since situation (ii) is not present, this means that no string segments in B can intersect more that once. In other words B is built out of crossings, and so we may pull all of B outside of the region R one crossing at a time by (H1), and thus obtain



Lastly we may pull these horizontal strings out of R through the top or bottom crossing by (H1). Hence we have emptied R by employing only local relation (H1), and so the value $\eta(\alpha)$ has remained the same. Now there are four different cases depending on the orientations of the strings *s* and *t*. In three of these cases, since R is empty, we may pull the strings *s* and *t* apart by applying (H2), and thus remove the two intersections *p* and *q*. This decreases

 $\eta(\alpha)$ by two, and so the result follows by induction. The last case is given with orientations as follows



where we have applied (H3). Let the two diagrams on the right hand side of the above equation be denoted by α_1 and α_2 respectively, hence $\alpha = \alpha_1 - \alpha_2$. It is clear that $\overline{\alpha_1} = \overline{\alpha}$ and deg(α_1) = deg(α). Moreover we have that $\eta(\alpha_1) = \eta(\alpha) - 2$, and so by induction

$$\alpha_1 = \overline{\alpha} + \sum_{\substack{\beta \in \operatorname{Sim}(a,b) \\ \deg(\beta) > \deg(\alpha)}} d_\beta \beta$$
(12)

where $d_{\beta} \in \mathbb{Z}$. As for α_2 , the original strings *s* and *t* have been replaced by *u* and *v*. Although the points of intersection *p* and *q* have been removed, in general we cannot apply the inductive step for α_2 as it may not be semisimple, since the new strings *u* and *v* may contain self-intersections. This occurs precisely when there are more intersections between the strings *s* and *t* than just *p* and *q*. So we break this situation into two cases:

(Case 1) Assume that p and q are the only intersections between the strings s and t in α , and so α_2 is semisimple. Thus by induction we have

$$\alpha_{2} = \overline{\alpha_{2}} + \sum_{\substack{\beta \in \operatorname{Sim}(\boldsymbol{a}, \boldsymbol{b}) \\ \deg(\beta) > \deg(\alpha_{2})}} f_{\beta}\beta$$
(13)

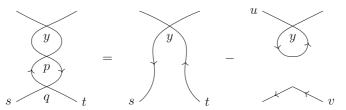
where $f_{\beta} \in \mathbb{Z}$. We seek to show that $\deg(\alpha_2) > \deg(\alpha)$, and then subtracking (13) away from (12) will prove this case. One can show this by comparing the string types of the sets $\{s, t\}$ and $\{u, v\}$. We have the following to consider:

- (1) The set $\{s, t\}$ contains a down and up string.
- (2) The set $\{s, t\}$ contains a vertical string and clockwise arc.
- (3) The set $\{s, t\}$ contains two arcs on the same row, but not both anti-clockwise.

Note $\{s, t\}$ cannot contain a top and bottom arc since α is semisimple. The remaining cases which have been left out are due to the fact they can never realise the orientated double crossing of the strings *s* and *t* which we are considering. For (1) it is easy to see that $\{u, v\}$ consists of two arcs. For (2) one can deduce that $\{u, v\}$ contains a vertical string and anti-clockwise arc. For (3), when $\{s, t\}$ consists of two clockwise arcs one can check that $\{u, v\}$ consists of a clockwise arc and an anti-clockwise arc. When $\{s, t\}$ contains a clockwise and anti-clockwise arc. For all these case we have deg(α_2) > deg(α), completing Case 1.

(Case 2) Assume now that there is at least one more point of intersection between the strings s and t beside p or q. In the ordering of intersections discussed previously, pick a

neighbouring point which either preceeds p or proceeds q, say y. Without loss of generality assume y preceeds p. Then diagrammatically the equation $\alpha = \alpha_1 - \alpha_2$ is given by



by (H3). In α_2 the interior region bounded by the left curl cannot contain loops or string segments with self-intersections since α is semisimple. Hence by Lemma 4.2.6 $\alpha_2 = 0$, and so $\alpha = \alpha_1$ and thus the result follows by (12).

Theorem 4.2.8 The algebra homomorphism $\varphi : \mathcal{A}_{2k}^{aff} \to \text{End}_{\text{Heis}}((\uparrow \downarrow)^k)$ given in Proposition 4.2.1 is surjective.

Proof As discussed previously, this will follow by showing that $\alpha \in \langle \varphi(e_i), \varphi(\tau_j) \rangle_{i,j}$ for all $\alpha \in \text{Sim}((\uparrow \downarrow)^k)$. We prove this by downwards induction on deg(α). It's easy to see that the maximum degree is deg(α) = (2k, 2k). By considering what endpoints can be targets and sources of clockwise arcs, one can deduce the only element $\alpha \in \text{Sim}((\uparrow \downarrow)^k)$ satisfying deg(α) = (2k, 2k) is given by

$$\overbrace{\longrightarrow}^{\longleftarrow} \qquad \overbrace{\longrightarrow}^{\longleftarrow} \qquad \overbrace{\longrightarrow}^{\longleftarrow} = \varphi \left(\prod_{i \in [k]} e_{2i-1} \right)$$

This completes the base case. Now, pick α such that $\deg(\alpha) = (x, y) < (2k, 2k)$ and assume that $\gamma \in \langle \varphi(e_i), \varphi(\tau_j) \rangle_{i,j}$ for all $\gamma \in \text{Sim}((\uparrow \downarrow)^k)$ such that $\deg(\gamma) > (x, y)$. By Lemma 4.2.5 there exists $\pi, \sigma \in S(k)$ and a planar diagram $\beta \in \text{Sim}((\uparrow \downarrow)^k)$ such that $\pi^{\uparrow}\beta\sigma^{\downarrow}$ is semisimple and $\alpha = \pi^{\uparrow}\beta\sigma^{\downarrow}$, in particular $\deg(\alpha) = \deg(\pi^{\uparrow}\beta\sigma^{\downarrow})$. Hence by Proposition 4.2.7 we have that

$$\pi^{\uparrow}\beta\sigma^{\downarrow} = \alpha + \sum_{\substack{\gamma \in \operatorname{Sim}((\uparrow\downarrow)^{k}) \\ \operatorname{deg}(\gamma) > \operatorname{deg}(\alpha)}} c_{\gamma}\gamma, \tag{14}$$

where $c_{\gamma} \in \mathbb{Z}$. By induction all the simple terms in the above summation belong to $\operatorname{Im}(\varphi)$. Also from previous discussions we know that $\pi^{\uparrow}\beta\sigma^{\downarrow} \in \operatorname{Im}(\varphi)$, thus rearranging the above equation shows that $\alpha \in \operatorname{Im}(\varphi)$, completing the proof by induction.

Remark 4.2.9 (14) is the key to Theorem 4.2.8, and follows from Proposition 4.2.7. This proposition applies to all semisimple diagrams which are much more general than those appearing here. Ideally, one would like to prove that (14) holds for $\pi^{\uparrow}\beta\sigma^{\downarrow}$ by some inductive argument without needing to show it for all semisimple diagrams. However it is a very delicate task to check which properties are preserved by an inductive process. So we ended up using this more general approach instead, even though many of the cases considered in proving Proposition 4.2.7 probably won't occur in this case.

4.3 The Affine Partition Category of Brundan and Vargas

In this last section we relate our affine partition algebra to the work of J. Brundan and M. Vargas in [2] and prove a new result on their category. We start by recalling the definition of their affine partition category APar as a subcategory of Heis generated by certain objects and morphisms, and of their affine partition algebra AP_k , which is an endomorphism algebra within APar.

Definition 4.3.1 [2, Definition 4.6 and Equation 4.47] The *affine partition category* APar is the monoidal subcategory of Heis generated by the object $\uparrow\downarrow$ and the following morphisms:

 $+ 1 \qquad (15)$

(16)

$$\begin{array}{c} \swarrow \\ \swarrow \\ \end{array} + \uparrow \\ \\ (19)$$

The *affine partition algebra* is defined to be $AP_k := \text{End}_{\text{APar}}((\uparrow\downarrow)^k)$.

We can generalise the arguments in the proof of Theorem 4.2.8 to show the following result.

Theorem 4.3.2 *The affine partition category* APar *is the full monoidal subcategory of* Heis *generated by the object* $\uparrow \downarrow$.

Proof We need to show that

$$\operatorname{Hom}_{\mathsf{APar}}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l) = \operatorname{Hom}_{\mathsf{Heis}}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l).$$

Using Theorem 4.1.8, we need to show that any element of the form

$$c_w^{k_w} \dots c_1^{k_1} c_0^{k_0} r_1^{s_1} r_3^{s_3} \dots r_{2k-1}^{s_{2k-1}} \alpha r_2^{t_2} r_4^{t_4} \dots r_{2l}^{t_{2l}}$$

where $\alpha \in \text{Sim}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ can be written in terms of the generating morphisms in APar. The morphisms r_i can be obtained by tensoring the generators (18) with the appropriate identity morphisms on the left and right (and subtracting the identity). Moreover, the morphisms c_i can be obtained by concatenating r_1^i with the generators (17) on top and bottom. Thus, it remains to show that any diagram $\alpha \in \text{Sim}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ can be written in terms of the generating morphism in APar. A generalisation of Jones' normal form shows that any planar $\alpha \in \text{Sim}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ can be written in terms of the generators (16) and (17) (see for example [18, Proof of Lemma 2.1] for an explicit construction). Now Lemma 4.2.5 allows us the write any $\alpha \in \text{Sim}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$ as $\alpha = \pi^{\uparrow}\beta\sigma^{\downarrow}$ where $\pi \in S(k), \sigma \in S(l)$ and β is planar. Note that s_i^{\uparrow} and s_i^{\downarrow} can be written using the generators (19) and the composition of the generators (16) (and tensoring with the appropriate identity morphism on the left and right). So using the discussion following Example 4.2.3 we know that π^{\uparrow} and σ^{\downarrow} belong to $\text{Hom}_{\text{APar}}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$. Now we can follow exactly the same proof as for Theorem 4.2.8 noting that in this case the maximum degree is (k + l, k + l) and the only simple diagram with that degree is the one containing *k* consecutive arcs at the top and *l* consecutive arcs at the bottom, which is planar. The rest of the proof can be followed verbatum simply replacing $\text{Im}\varphi$ by $\text{Hom}_{\text{APar}}((\uparrow\downarrow)^k, (\uparrow\downarrow)^l)$.

We immediately obtain the following consequences.

Corollary 4.3.3 The map φ gives a surjective homomorphism for \mathcal{A}_{2k}^{aff} to AP_k .

Corollary 4.3.4 The set $B((\uparrow\downarrow)^k)$ gives a basis for AP_k .

We do not know whether the map φ is an isomorphism. If it were, then we would also have a presentation for AP_k .

Data Availability All data generated or analysed during this study are included in this published article.

Declarations

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References

- Brauer, J.: On algebras which are connected with the semisimple continuous groups. Annals Math. 38, 857–872 (1937). https://doi.org/10.2307/1968843
- 2. Brundan, J., Vargas, M.: A new approach to the representation theory of the partition category. arXiv:2107.05099 (2022)
- 3. Creedon, S.: The center of the partition algebra. J. Algebra **570**, 215–266 (2021). https://doi.org/10.1016/j.jalgebra.2020.10.041
- Daugherty, Z., Ram, A., Virk, R.: Affine and degenerate affine BMW algebras: the center Osaka. J. Math., vol. 51. https://projecteuclid.org/euclid.ojm/1396966236 (2011)
- James East: Generators and relations for partition monoids and algebras. J. Algebra 339(1), 1–26 (2011). https://doi.org/10.1016/j.jalgebra.2011.04.008
- Enyang, J.: Jucys-murphy elements and a presentation for the partition algebra. J. Algebraic Combin. 37(3), 401–454 (2013). MR 3035512, 2012. https://doi.org/10.1007/s10801-012-0370-4
- Enyang, J.: A seminormal form for partition algebras. J. Combinatorial Theory, Series A 120, 1737–1785 (2013). 2013. https://doi.org/10.1016/j.jcta.2013.06.001
- Halverson, T., Ram, A.: Partition algebras. European J. Combinatorics 26, 869–921 (2005). https://doi.org/10.1016/j.ejc.2004.06.005
- 9. Jones, V.: Index for subfactors. Invent Math. 72, 1-25 (1983)
- Jones, V.F.R.: The potts model and the symmetric group, Subfactors (Kyuzeso, 1993), World Sci. Publ., River Edge, NJ. pp. 259–267 (1994)
- Kauffman, L.H.: An invariant of regular isotopy. Trans. Amer. Math. Soc. 318(2), 417–471 (1990). https://doi.org/10.1090/S0002-9947-1990-0958895-7

- Khovanov, M.: Heisenberg algebra and a graphical calculus. Fundam. Math. 225, 169–210 (2014). arXiv:1009.3295
- Kleshchev, A.: Linear and projective representations of the symmetric group. Cambridge tract in mathematics, series number 163, CUP (2005)
- 14. Likeng, S., Savage, A.: Embedding Deligne's category $\underline{\text{Rep}}(S_t)$ in the Heisenberg category. Quantum Topol. **12**, 211–242 (2021)
- Martin, P.P.: Potts Models and Related Problems in Statistical Mechanics Series on Advances in Statistical Mechanics, vol. 5. World Scientific Publishing Co. Inc., NJ (1991)
- Nazarov, M.: Young's orthogonal form for Brauer's centralizer algebra. J. Algebra 182, 664–693 (1996). https://doi.org/10.1006/jabr.1996.0195
- Orellana, R., Ram, A.: Affine braids, Markov traces and the category O, Algebraic groups and homogeneous spaces. Tata Inst. Fund. Res. Stud. Math., Tata Inst, Fund. Res., Mumbai, pp. 423–473 (2007)
- Ridout, D., Saint-Aubin, Y.: Standard modules, induction and the structure of the Temperley-Lieb algebra. Adv. Theor. Math. Phys. 18, 957–1041 (2014)
- Rui, H., Su, Y.: Affine walled Brauer algebras and super Schur–Weyl duality. Adv. Math. 285, 28–71 (2015). https://doi.org/10.1016/j.aim.2015.07.018
- 20. Sartori, A.: The degenerate affine walled Brauer algebra, arXiv:1305.2347 (2014)
- Stembridge, J.R.: Some combinatorial aspects of reduced words in finite coxeter groups. Trans. Amer. Math. Soc 349, 1285–1332 (1997). https://doi.org/10.1090/S0002-9947-97-01805-9

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