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# A priori and universal estimates for global solutions of superlinear degenerate parabolic equations

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**Abstract**. We prove an a priori estimate and a universal bound for any global solution of the nonlinear degenerate reaction-diffusion equation  $u_t = \Delta u^m + u^p$  in a bounded domain with zero Dirichlet boundary conditions.

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**Key words.** nonlinear degenerate parabolic equations – global solutions – a priori estimates – universal bounds

# 1. Introduction and main results

Consider the following problem:

(1.1) 
$$\begin{cases} u_t = \Delta u^m + u^p, & 0 < t < T, \quad x \in \Omega, \\ u(t, x) = 0, & 0 < t < T, \quad x \in \partial\Omega, \\ u(0, x) = u_0(x), & x \in \partial\Omega. \end{cases}$$

Throughout the paper,  $\Omega$  is a  $C^3$ -smooth bounded domain of  $\mathbb{R}^N$ ,  $p > \max(1, m)$ and m > 0. We consider solutions which may change sign and define  $u^k$  as  $|u|^k \operatorname{sign}(u)$  for all real k > 0. We assume that the initial datum  $u_0 \in L^{\infty}(\Omega)$ satisfies  $u_0^m \in H_0^1(\Omega)$ .

It is well known that solutions of (1.1) blow up in finite time if  $u_0$  is suitably large, while they exist globally and decay as  $t \to \infty$  if  $u_0$  is small. In view of a classification of all solutions of (1.1), it is then a natural question to ask whether unbounded global solutions may exist or not.

The question of the boundedness of global solutions of (1.1) was initiated in [15] and further investigated in [4]. Denoting  $p_S = (N+2)/(N-2)$  ( $\infty$  if  $N \le 2$ ), the result of [4] says that if  $p/m < p_S$ , then any global solution of (1.1) is uniformly bounded for  $t \ge 0$ , that is,

(1.2) 
$$\sup_{t\geq 0} \|u(t)\|_{\infty} < \infty.$$

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On the other hand, it is known from [15] that some unbounded global solutions do exist if  $N \ge 3$  and  $p/m \ge p_S$  (see also [9]).

In the semilinear case m = 1, the conclusion of [4] had been obtained earlier in [3]. A more precise result was given in [10] for m = 1 and  $u_0 \ge 0$ , where an *a priori* estimate of the form

(1.3) 
$$\sup_{t\geq 0} \|u(t)\|_{\infty} \leq C(\|u_0\|_{\infty}, \Omega, p)$$

was established for  $p < p_S$ . In [16], (1.3) was then proved for m = 1 without the restriction  $u_0 \ge 0$ . It is to be noted that boundedness of global solutions does not imply the a priori estimate (1.3) in general (see [9, 6] for some counter-examples).

Recently, the question of whether the bound in (1.3) might be independent of  $u_0$  for  $t \ge \tau > 0$  was raised in [7], and it was proved there that

(1.4) 
$$\sup_{t \ge \tau} \|u(t)\|_{\infty} \le C(\tau, \Omega, p), \quad \text{for all } \tau > 0.$$

holds for all global solutions of (1.1), provided m = 1,  $u_0 \ge 0$  and p < (N + 1)/(N - 1) or N = 1. We call (1.4) a *universal* bound. In other words, (1.4) says that after any positive time  $\tau$ , every global non-negative trajectory of (1.1) enters into an absorbing bounded set  $A_{\tau}$ . Shortly thereafter, it was shown in [17] that (1.4) still holds for  $N \le 2$  or N = 3 and  $p < 5 = p_S$  (with m = 1 and  $u_0 \ge 0$ ).

For more results on boundedness of global solutions and on a priori estimates for other classes of evolution equations, we refer to the surveys [18, 5], to [13, 20] and to the references in [7]. Let us also mention the paper [19] which contains some results on semilinear parabolic systems.

The aim of the present paper is to establish both a priori estimates and universal bounds for global solutions of (1.1) in the degenerate case m > 1. By a solution of (1.1), we mean a weak solution (see Section 2 below for a precise definition). Our main results are the following:

# **Theorem 1** (a priori estimate). Assume that

(1.5) 
$$1 < m < p < p_1(m, N) := \begin{cases} \infty, & \text{if } N = 1\\ m + \frac{10m+2}{3N-4}, & \text{if } N \ge 2. \end{cases}$$

Then any global solution of (1.1) satisfies

(1.6) 
$$\sup_{t \ge 0} \|u(t)\|_{\infty} \le C(\|u_0\|_{\infty}, \Omega, p, m).$$

where the constant  $C(||u_0||_{\infty}, \Omega, p, m) > 0$  remains bounded for  $||u_0||_{\infty}$  bounded.

### **Theorem 2** (universal bound). Assume that

(1.7) 
$$1 < m < p < p_2(m, N) := \begin{cases} \infty, & \text{if } N = 1 \\ \frac{N+2}{N}m, & \text{if } N \ge 2. \end{cases}$$

Then, for all  $\tau > 0$ , there exists a constant  $C(\tau, \Omega, p, m) > 0$  such that any global non-negative solution of (1.1) satisfies

(1.8) 
$$\sup_{t \ge \tau} \|u(t)\|_{\infty} \le C(\tau, \Omega, p, m).$$

*Remark 1.1.* It is easy to show that (1.8) implies (1.6). On the other hand, (1.8) cannot be true for  $\tau = 0$ , since there exist global solutions starting from unbounded initial data in  $L^q$  (see, e.g., [1]). Furthermore, Theorem 2 cannot be true for solutions of a mixed sign. Indeed, if  $1 < p/m < p_S$ , then there exist sign-changing stationary solutions of arbitrary large sup norm.

*Remark 1.2.* The paper [4] on boundedness of global solutions also gives an a priori estimate similar to (1.6) under the condition m (see [4, Remark 1.10 p. 236], where <math>m + 1 is substituted by 2 due to a misprint). Our condition (1.5) in Theorem 1 is weaker. However, we restrict to the "slow-diffusion" case m > 1, while the results of [4] also work for the "fast-diffusion" case m < 1, with m > (N - 2)/(N + 2).

The proof of our a priori estimate relies on suitable modifications of ideas in [3], based on energy and interpolation arguments. For the proof of universal bounds, besides the previously established a priori estimate, we use suitable test functions and certain smoothing properties of solutions of (1.1).

The rest of the paper is organized as follows. Section 2 contains some useful preliminary material, including local existence and smoothing properties of solutions. Sections 3 and 4 are then devoted to the proofs of Theorems 1 and 2, respectively.

# 2. Preliminaries

For  $u_0 \in L^{\infty}(\Omega)$ ,  $u_0^m \in H_0^1(\Omega)$ , by a (weak) solution of (1.1) on [0, T], we mean a function *u* such that

(2.1) 
$$\begin{cases} u \in C([0,T]; L^{2}(\Omega)) \cap L^{\infty}(0,T; L^{\infty}(\Omega), \quad u^{m} \in L^{\infty}(0,T; H_{0}^{1}(\Omega)), \\ \int_{0}^{t} \int_{\Omega} \left( u\varphi_{t} - \nabla u^{m} \cdot \nabla \varphi + u^{p}\varphi \right) = \int_{\Omega} u(t)\varphi(t) - \int_{\Omega} u_{0}\varphi(0) \\ \text{for all } t \in (0,T] \text{ and } \varphi \in L^{2}(0,T; H_{0}^{1}(\Omega)) \cap H^{1}(0,T; L^{2}(\Omega)). \end{cases}$$

It is known (see, e.g., [12]) that there exists  $T^* = T^*(u_0) \in (0, \infty]$  such that for each  $T \in (0, T^*)$ , (P) admits a unique solution on [0, T]. If  $T^* < \infty$ , then  $\|u(t)\|_{\infty} \to \infty$  as  $t \to T^*$ . Moreover, u satisfies the energy inequality

(2.2) 
$$E(t_1) + m \int_{t_0}^{t_1} \int_{\Omega} |u|^{m-1} u_t^2 \le E(t_0), \quad 0 \le t_0 \le t_1 < T^*,$$

where the energy is defined by

(2.3) 
$$E(t) := \int_{\Omega} \left( \frac{1}{2} |\nabla u^m(t)|^2 - \frac{m}{p+m} |u(t)|^{p+m} \right).$$

In particular, the energy is non-increasing in time. Also, we have the useful identity

(2.4) 
$$\frac{d}{dt} \int_{\Omega} \frac{|u|^{m+1}}{m+1} = -\int_{\Omega} |\nabla u^m(t)|^2 + \int_{\Omega} |u(t)|^{p+m}, \quad \text{a.e. } t \in (0, T^*).$$

(This follows from (2.1) with the test function  $\varphi = u^m$  – note that  $\varphi_t \in L^2(0, T; L^2(\Omega))$  is due to (2.2) and  $u \in L^{\infty}(0, T; L^{\infty}(\Omega))$ .)

Finally, in what follows, C(...) will denote positive constants which may vary from line to line and depend only on the indicated arguments.

In the proof of Theorems 1 and 2, we will need the following two lemmas which give useful regularizing properties for local solutions of (1.1). In particular Lemma 2.2 (i), which is essentially a consequence of results in [1], provides smoothing properties from  $L^q$  into  $L^{\infty}$ .

**Lemma 2.1.** Assume p > m > 1. For all  $t \in (0, T^*)$ , it holds

(2.5) 
$$\|\nabla u^m(t)\|_2^2 \le C(p,m) \cdot \left(t^{-1} \|u_0\|_{m+1}^{m+1} + \sup_{s \in (0,t)} \|u(s)\|_{m+p}^{m+p}\right).$$

Moreover, we have

(2.6) 
$$T^* > T_1 := C(p) \|u_0\|_{\infty}^{1-p}, \qquad \|u(t)\|_{\infty} \le 2\|u_0\|_{\infty}, \quad 0 < t \le T_1$$

and

(2.7) 
$$\|\nabla u^m(T_1)\|_2^2 \le C(p,m)|\Omega| \|u_0\|_{\infty}^{p+m}.$$

**Lemma 2.2.** Assume p > m > 1 and  $q \ge 1$ ,  $q > \frac{N}{2}(p-m)$ . (i) There exist positive constants  $L_1, L_2, \alpha, \beta, \gamma$  depending only on N, p, m, q, such that  $T^* > T_0 := L_2(1 + ||u_0||_q)^{-\gamma}$  and

(2.8) 
$$||u(t)||_{\infty} \leq L_1 ||u_0||_q^{\beta} t^{-\alpha}, \quad 0 < t \leq T_0.$$

(ii) For all M > 0, there exists C(M, N, p, m, q) > 0 such that if

$$0 < T \le T^*$$
,  $||u_0||_{\infty} \le M$  and  $\sup_{t \in (0,T)} ||u(t)||_q \le M$ ,

then

$$||u(t)||_{\infty} \le C(M, N, p, m, q), \quad 0 \le t < T.$$

*Proof of Lemma 2.1.* Integrating (2.4) over (0, t), we have

$$\int_{0}^{t} \int_{\Omega} |\nabla u^{m}(s)|^{2} \leq \int_{0}^{t} \int_{\Omega} |u(s)|^{m+p} + \int_{\Omega} \frac{|u_{0}|^{m+1}}{m+1}.$$

Due to the non-increasing property of the energy, it follows that

$$tE(t) \leq \int_0^t E(s) \, ds \leq \frac{1}{2} \int_0^t \int_{\Omega} |\nabla u^m(s)|^2.$$

By combining these two inequalities, we deduce that

$$\int_{\Omega} |\nabla u^{m}(t)|^{2} = 2E(t) + \frac{2m}{m+p} \int_{\Omega} |u(t)|^{m+p}$$

$$\leq \frac{1}{t} \left( \int_{0}^{t} \int_{\Omega} |u(s)|^{m+p} + \int_{\Omega} \frac{|u_{0}|^{m+1}}{m+1} \right) + \frac{2m}{m+p} \int_{\Omega} |u(t)|^{m+p},$$
high implies (2.5)

which implies (2.5).

On the other hand, by comparing with the solution of the ODE  $y' = y^p$  for  $y(0) = ||u_0||_{\infty}$ , we get

$$\|u(t)\|_{\infty} \le y(t) = \left(\|u_0\|_{\infty}^{1-p} - (p-1)t\right)^{-1/(p-1)}, 0 < t < \min\left(T^*, (p-1)^{-1}\|u_0\|_{\infty}^{1-p}\right),$$

which immediately yields (2.6).

Finally, (2.7) follows from (2.6) and (2.5) with  $t = T_1$ .

*Proof of Lemma 2.2.* (i) For each  $u_0 \in L^q(\Omega)$  with  $u_0 \ge 0$ , it follows from [1, Theorem 3.1 and Remark 3.3] that there exists (at least) a solution v of (1.1) which exists on  $[0, T_0]$  with  $T_0 := L_2(1 + ||u_0||_q)^{-\gamma}$  and satisfies (2.8). However, the notion of solution in [1] is weaker than (2.1), so that we cannot immediately identify v with our solution u.

To do so, we note that the solution v in [1] is constructed as a limit, pointwise on  $(0, T_0) \times \Omega$ , of solutions of the approximating problems

(2.9)<sub>n</sub> 
$$\begin{cases} v_{n,t} = \Delta v_n^m + \min(v_n^p, n), & t > 0, \quad x \in \Omega, \\ v_n(t, x) = 0, & t > 0, \quad x \in \partial \Omega, \\ v_n(0, x) = \min(u_0(x), n), & x \in \partial \Omega. \end{cases}$$

For each  $n \ge 1$ ,  $(2.9)_n$  admits a unique (global) solution  $v_n$  in the sense of (2.1). For each  $0 < T' < T^*(u_0)$ , since  $u_0 \in L^{\infty}$  and  $u \in L^{\infty}((0, T') \times \Omega)$ , it follows by uniqueness for  $(2.9)_n$  that  $v_n = u$  on  $(0, T') \times \Omega$  for all  $n \ge n_0(T')$  large enough. Passing to the limit, we deduce that v = u on  $(0, \min(T', T_0)) \times \Omega$ , hence on  $(0, \min(T^*(u_0), T_0)) \times \Omega$ . Since v satisfies (2.8), we conclude that u satisfies (2.8).

Finally, in the case when  $u_0$  changes sign, it suffices to compare u with the solutions  $u^{\pm}$  corresponding to the initial data  $\pm |u_0|$ .

(ii) Let  $T_2 = \min(T_0, T_1)$ , where  $T_0$  is defined in Lemma 2.2 (i) and  $T_1$  is defined in (2.6). By Lemma 2.1 and Lemma 2.2 (i), we have

$$\|u(t)\|_{\infty} \leq \begin{cases} 2\|u_0\|_{\infty}, & \text{if } 0 \leq t \leq T_2\\ L_1\|u(t-T_2)\|_q^{\beta} T_2^{-\alpha}, & \text{if } T_2 \leq t < T, \end{cases}$$

and the result follows.

*Remark 2.1.* The result of Lemma 2.2 (ii) can also be found in [8] under the assumption  $q > \max(1, N/2)(p - m)$  (stronger if N = 1). Under this assumption, the result of Lemma 2.2 (i) can be proved alternatively by combining the arguments in [4, Lemma 1.6] and [14, pp. 46–48] with [11, Théorème 1].

## 3. Proof of Theorem 1

The proof of Theorem 1 relies on a suitable adaptation of the arguments in [3], based on energy estimates and interpolation.

*Proof of Theorem 1.* Let *u* be a global solution of (1.1). From (2.4) and (2.3) we have, for a.e.  $t \ge 0$ ,

(3.1) 
$$\frac{d}{dt} \int_{\Omega} \frac{|u|^{m+1}}{m+1} = -2E(t) + \frac{p-m}{p+m} \int_{\Omega} |u(t)|^{p+m}$$

By Hölder's inequality and the non-increasing property of the energy, it follows that for each  $t_0 \ge 0$ ,

$$\frac{d}{dt}\int_{\Omega}\frac{|u|^{m+1}}{m+1}\geq C(p,m,\Omega)\Big(\int_{\Omega}|u(t)|^{m+1}\Big)^{\alpha}-2E(t_0),\quad \text{ a.e. }t\geq t_0,$$

where  $\alpha = \frac{m+p}{m+1} > 1$ . Since *u* exists globally, this implies

$$(3.2) E(t_0) \ge 0, \text{for all } t_0 \ge 0$$

and

(3.3) 
$$\int_{\Omega} |u(t)|^{m+1} \le C(p, m, \Omega) E^{\frac{m+1}{m+p}}(0), \quad \text{for all } t \ge 0.$$

On the other hand, by (2.2) and (3.2), we have

(3.4) 
$$\int_0^\infty \int_\Omega |u|^{m-1} u_t^2 \le \frac{E(0)}{m}.$$

Now using (3.1) and  $E(t) \leq E(0)$ , we get

$$\frac{p-m}{p+m}\int_{\Omega}|u(t)|^{p+m} \leq 2E(0) + \int_{\Omega}|u(t)|^{(m+1)/2}|u(t)|^{(m-1)/2}u_t;$$

hence, by Cauchy-Schwarz's inequality and (3.3),

(3.5) 
$$\left( \int_{\Omega} |u(t)|^{p+m} \right)^2 \leq 2 \left( \frac{p+m}{p-m} \right)^2 \left( 4E^2(0) + \int_{\Omega} |u(t)|^{m+1} \int_{\Omega} |u(t)|^{m-1} u_t^2 \right)$$
$$\leq C(p,m,\Omega) \left( E^2(0) + E^{\frac{m+1}{m+p}}(0) \int_{\Omega} |u(t)|^{m-1} u_t^2 \right).$$

By integrating over (t, t + 1), setting  $a = \frac{p+2m+1}{p+m}$ , it follows from (3.4) that

(3.6) 
$$\int_{t}^{t+1} \left( \int_{\Omega} |u|^{p+m} \right)^{2} \leq C(p,m,\Omega) \left( E^{2}(0) + E^{a}(0) \right), \quad t \geq 0.$$

Let  $v = u^{(m+1)/2}$  and  $r = \frac{2(m+p)}{m+1} > 2$ . In the rest of the proof, k will denote various positive constants depending only on p and m. Rewrite the inequalities (3.4) and (3.6) as

$$\int_0^\infty \int_\Omega v_t^2 \le C(m) E(0)$$

and

$$\int_t^{t+1} \left( \int_{\Omega} |v|^r \right)^2 \le C(p, m, \Omega) \left( 1 + E^k(0) \right), \quad t \ge 0.$$

hence, in particular,

 $\|v\|_{L^{2r}(t,t+1;L^r(\Omega))} + \|v\|_{H^1(t,t+1;L^2(\Omega))} \le C(p,m,\Omega)(1+E^k(0)), \quad t\ge 0.$ 

By interpolation (see [3, Appendice] or also [2] for more general results), it follows that

$$\|v\|_{L^{\infty}(t,t+1;L^{a}(\Omega))} \le C(p,m,\Omega,a)(1+E^{k}(0)), \quad t \ge 0,$$

for all  $a \in [1, a_0)$ , where  $a_0 = r - \frac{r-2}{3} = \frac{2(r+1)}{3} > 2$ , so that

(3.7) 
$$\|u\|_{L^{\infty}(0,\infty;L^{q}(\Omega))} \leq C(p,m,\Omega,q)(1+E^{k}(0)),$$

for all  $q \in [1, q_0)$ , where  $q_0 = \frac{(m+1)(r+1)}{3} = \frac{3m+2p+1}{3} > 1$ . By combining (2.6) and (2.7) from Lemma 2.2 with (3.7) (shifting the origin

By combining (2.6) and (2.7) from Lemma 2.2 with (3.7) (shifting the origin of time from 0 to  $T_1$  in (3.7)), we obtain

$$||u||_{L^{\infty}(0,\infty;L^{q}(\Omega))} \le C(p,m,\Omega,q) (1+||u_{0}||_{\infty}^{k}), \text{ for all } q \in [1,q_{0}).$$

Finally, since the assumption (1.5) implies that  $q_0 > \frac{N}{2}(p-m)$ , the conclusion follows from Lemma 2.2 (ii).

*Remark 3.1.* By using the estimates (3.5), (3.4), Lemma 2.2 (i) and arguing similarly as in the proof of [3, Proposition 6], we obtain a new proof of the boundedness of global solutions of (1.1) under the assumption  $1 < p/m < p_S$ , m > 1. This proof is completely different from that in [4] (note, however, that the proof in [4] works also for (N - 2)/(N + 2) < m < 1).

# 4. Proof of Theorem 2

In this section, we denote, respectively, by  $\lambda_1 > 0$  and  $\varphi_1(x) > 0$  the first eigenvalue and first eigenfunction of  $-\Delta$  in  $\Omega$  with Dirichlet boundary conditions, satisfying

$$\begin{cases} -\Delta \varphi_1 = \lambda_1 \varphi_1, & x \in \Omega, \\ \varphi_1 \in H_0^1(\Omega), \\ \int_{\Omega} \varphi_1 = 1. \end{cases}$$

Let u be a non-negative global solution of (1.1). We first claim that

(4.1) 
$$\int_{\Omega} u(t)\varphi_1 \le C(p,m,\lambda_1), \quad t \ge 0$$

and

(4.2) 
$$\int_0^{\tau} \int u^p \varphi_1 \le C(p, m, \lambda_1, \tau), \quad \tau > 0.$$

Indeed, taking  $\varphi = \varphi_1$  in (2.1), we get for a.e.  $t \ge 0$ ,

$$\frac{d}{dt}\int_{\Omega}u(t)\varphi_{1}=\int_{\Omega}u^{p}(t)\varphi_{1}-\lambda_{1}\int_{\Omega}u^{m}(t)\varphi_{1};$$

hence, by Young's and Jensen's inequalities, and since p > m,

(4.3) 
$$\frac{d}{dt} \int_{\Omega} u(t)\varphi_1 \ge \frac{1}{2} \int_{\Omega} u^p(t)\varphi_1 - C(p, m, \lambda_1)$$
$$\ge \frac{1}{2} \left( \int_{\Omega} u(t)\varphi_1 \right)^p - C(p, m, \lambda_1).$$

This implies (4.1) since otherwise  $\int_{\Omega} u(t)\varphi_1$  would blow up in finite time, contradicting  $T^* = \infty$ . Integrating (4.3) over (0,  $\tau$ ) and using (4.1) we deduce

$$\frac{1}{2}\int_0^\tau\int_\Omega u^p(t)\varphi_1\leq C(p,m,\lambda_1)\tau+\int_\Omega u(\tau)\varphi_1\leq C(p,m,\lambda_1,\tau),$$

which is (4.2).

We next claim that there exists  $t_0 \in (0, \tau/2)$  such that for all  $l \in (0, p/2)$ ,

(4.4) 
$$\int_{\Omega} u^m(t_0) + \int_{\Omega} u^l(t_0) \le C(p, m, l, \Omega, \tau).$$

To this end, we introduce the solution  $\chi$  of the problem

$$\begin{cases} -\Delta \chi = 1, & x \in \Omega, \\ \chi \in H_0^1(\Omega). \end{cases}$$

Using (2.1) with test function  $\chi$ , we get

$$\int_0^\tau \int_\Omega u^m = \int_0^\tau \int_\Omega u^p \chi + \int_\Omega u_0 \chi - \int_\Omega u(\tau) \chi.$$

Since

$$c\chi(x) \le \operatorname{dist}(x, \partial\Omega) \le c^{-1}\varphi_1(x), \quad x \in \Omega,$$

for some  $c = c(\Omega) > 0$ , we deduce from (4.1) and (4.2) that

$$\int_0^\iota \left( \int_\Omega u^m + \int_\Omega u^p \varphi_1 \right) \le C(p, m, \Omega, \tau).$$

In particular, there exists  $t_0 \in (0, \tau/2)$  such that  $\int_{\Omega} u^m(t_0) + \int_{\Omega} u^p(t_0)\varphi_1 \le C(p, m, \Omega, \tau)$ . Now, by Hölder's inequality,

$$\int_{\Omega} u^{l}(t_{0}) = \int_{\Omega} u^{l}(t_{0})\varphi_{1}^{l/p}.\varphi_{1}^{-l/p} \le \left(\int_{\Omega} u^{p}(t_{0})\varphi_{1}\right)^{l/p} \left(\int_{\Omega} \varphi_{1}^{-l/(p-l)}\right)^{(p-l)/p}$$

Since l/(p-l) < 1 and  $\varphi_1(x) \ge c$  dist $(x, \partial \Omega)$ , we have  $\int_{\Omega} \varphi_1^{-l/(p-l)} < \infty$  and (4.4) follows.

Now, let us first consider the case  $N \ge 2$ , for which  $p < p_2(N, m)$  implies that  $m > \frac{N}{2}(p - m)$ . We may thus apply Lemma 2.1 (i) with q = m. From  $||u(t_0)||_m \le C(p, m, \Omega, \tau)$  (see (4.4)), we then deduce that

(4.5) 
$$||u(t_1)||_{\infty} \leq C(p, m, \Omega, \tau) \text{ for some } t_1 \in (t_0, \tau).$$

In the case N = 1, we have  $p/2 > \frac{N}{2}(p-m)$  and since m > 1, we can choose  $\tilde{m} \ge 1$  such that  $\frac{N}{2}(p-m) < \tilde{m} < \max(m, p/2)$ . We may thus apply Lemma 2.1 (i) with  $q = \tilde{m}$ . Since  $||u(t_0)||_{\tilde{m}} \le C(p, m, \Omega, \tau)$  due to (4.4), we then deduce that (4.5) is still true.

Finally, since  $p_1(N, m) > p_2(N, m)$ , we conclude from (4.5) and Theorem 1 that

$$\sup_{t \ge \tau} \|u(t)\|_{\infty} \le C(p, m, \Omega, \tau).$$

*Remark 4.1.* The control on  $\int_{\Omega} u^l(t_0)$  for all l < p/2 (cf. (4.4)) does not enable one to improve the value of  $p_2(N, m)$  if  $N \ge 2$ .

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