



Expansion of eigenvalues of the perturbed discrete bilaplacian

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Abstract

We consider the family

$$\widehat{H}_\mu := \widehat{\Delta}\widehat{\Delta} - \mu\widehat{V}, \quad \mu \in \mathbb{R},$$

of discrete Schrödinger-type operators in d -dimensional lattice \mathbb{Z}^d , where $\widehat{\Delta}$ is the discrete Laplacian and \widehat{V} is of rank-one. We prove that there exist coupling constant thresholds $\mu_o, \mu^o \geq 0$ such that for any $\mu \in [-\mu^o, \mu_o]$ the discrete spectrum of \widehat{H}_μ is empty and for any $\mu \in \mathbb{R} \setminus [-\mu^o, \mu_o]$ the discrete spectrum of \widehat{H}_μ is a singleton $\{e(\mu)\}$, and $e(\mu) < 0$ for $\mu > \mu_o$ and $e(\mu) > 4d^2$ for $\mu < -\mu^o$. Moreover, we study the asymptotics of $e(\mu)$ as $\mu \searrow \mu_o$ and $\mu \nearrow -\mu^o$ as well as $\mu \rightarrow \pm\infty$. The asymptotics highly depends on d and \widehat{V} .

Keywords Discrete bilaplacian · Essential spectrum · Discrete spectrum · Eigenvalues · Asymptotics · Expansion

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1 Introduction

In this paper we investigate the spectral properties of the perturbed discrete biharmonic operator

$$\widehat{H}_\mu := \widehat{\Delta}\widehat{\Delta} - \mu\widehat{V}, \quad \mu \in \mathbb{R}, \quad (1.1)$$

in the d -dimensional cubical lattice \mathbb{Z}^d , where $\widehat{\Delta}$ is the discrete Laplacian and \widehat{V} is a rank-one potential with a generating potential \widehat{v} . This model is associated to a one-particle system in \mathbb{Z}^d with a potential field \widehat{v} , in which the particle freely “jumps” from a node X of the lattice not only to one of its nearest neighbors Y (similar to the discrete Laplacian case), but also to the nearest neighbors of the node Y . From the mathematical point of view, the discrete bilaplacian represents a discrete Schrödinger operator with a degenerate bottom, i.e., $\widehat{\Delta}\widehat{\Delta}$ is unitarily equivalent to a multiplication operator by a function ϵ which behaves as $o(|p - p_0|^2)$ close to its minimum point p_0 .

The spectral properties of discrete Schrödinger operators with non-degenerate bottom (i.e., ϵ behaves as $O(|p - p_0|^2)$ close to its minimum point p_0), in particular with discrete Laplacian, have been extensively studied in recent years (see e.g. [1, 2, 7, 8, 10, 11, 20, 21, 23, 26, 28] and references therein) because of their applications in the theory of ultracold atoms in optical lattices [16, 24, 35, 36]. In particular, it is well-known that the existence of the discrete spectrum is strongly connected to the threshold phenomenon [18, 20–22], which plays a role in the existence of the Efimov effect in three-body systems [31, 32, 34]: *if any two-body subsystem in a three-body system has no bound state below its essential spectrum and at least two two-body subsystems have a zero-energy resonance, then the corresponding three-body system has infinitely many bound states whose energies accumulate at the lower edge of the three-body essential spectrum.*

Recall that the Efimov effect may appear only for certain attractive systems of particles [29]. However, recent experimental results in the theory of ultracold atoms in an optical lattice have shown that two-particle systems can have repulsive bound states and resonances (see e.g. [36]), thus, one expects the Efimov effect to hold also for some repulsive three-particle systems in \mathbb{Z}^3 .

The strict mathematical justification of the Efimov effect including the asymptotics for the number of negative eigenvalues of the three-body Hamiltonian has been successfully established in 3-space dimensions (for both \mathbb{R}^3 and \mathbb{Z}^3) see e.g., [1, 4, 13, 19, 29, 31, 32, 34] and the references therein. In particular, the non-degeneracy of the bottom of the (reduced) one-particle Schrödinger operator played an important role in the study of resonance states of the associated two-body system [1, 31]. Another keypoint in the proof of the Efimov effect in \mathbb{Z}^3 was the asymptotics of the (unique) smallest eigenvalue of the (reduced) one-particle discrete Schrödinger operator which creates a singularity in the kernel of a Birman-Schwinger-type operator which used to obtain an asymptotics to the number of three-body bound states.

To the best of our knowledge, there are no published results related to the Efimov effect in lattice three-body systems in which associated (reduced) one-body Schrödinger operator has degenerate bottom.

We also recall that fourth order elliptic operators in \mathbb{R}^d in particular, the biharmonic operator, play also a central role in a wide class of physical models such as linear elasticity theory, rigidity problems (for instance, construction of suspension bridges) and in streamfunction formulation of Stokes flows (see e.g. [9, 25, 27] and references therein). Moreover, recent investigations have shown that the Laplace and biharmonic operators have high potential in image compression with the optimized and sufficiently sparse stored data [15]. The need for corresponding numerical simulations has led to a vast literature devoted to a variety of discrete approximations to the solutions of fourth order equations [5, 12, 33]. The question of stability of such models is basically related to their spectral properties and therefore, numerous studies have been dedicated to the numerical evaluation of the eigenvalues [3, 6, 30].

The aim of the present paper is the study of the existence and asymptotics of eigenvalues as well as threshold resonance and bound states of \widehat{H}_μ defined in (1.1), which corresponds to the one-body Schrödinger operator with degenerate bottom. Namely, we study the discrete spectrum of \widehat{H}_μ depending on μ and on \widehat{v} . For simplicity we assume the generator \widehat{v} of \widehat{V} to decay exponentially at infinity, however, we urge that our methods can also be adjusted to less regular cases (see Remark 2.6). Since the spectrum of $\widehat{\Delta}$ consists of $[0, 2d]$ (see e.g., [1]), by the compactness of \widehat{V} and Weyl's Theorem, the essential spectrum of \widehat{H}_μ fills the segment $[0, 4d^2]$ independently of μ . Moreover, the essential spectrum does not give birth to a new eigenvalue while μ runs in some real interval $[-\mu^o, \mu_o]$, and it turns out as soon as μ leaves this interval through μ_o resp. through $-\mu^o$, a unique negative resp. a unique positive eigenvalue $e(\mu)$ releases from the essential spectrum (Theorem 2.2).

Now we are interested in the absorption rate of $e(\mu)$ as $\mu \rightarrow \mu_o$ and $\mu \rightarrow -\mu^o$. The associated asymptotics are highly dependent not only on the dimension d of the lattice (as in the discrete Laplacian case [20, 21]), but also values on the multiplicity $2n_o$ and $2n^o$ of $0 \in \{v = 0\}$ (if $v(0) = 0$) and $\vec{\pi} \in \{v = 0\}$ (if $v(\vec{\pi}) = 0$), respectively. More precisely, depending on d and n_o , $e(\mu)$ has a convergent expansion

- in $(\mu - \mu_o)^{1/3}$ for $2n_o + d = 1, 7$;
- in $\mu - \mu_o$ for $2n_o + d = 3, 5$;
- in $(\mu - \mu_o)^{1/4}$ for $2n_o + d \geq 9$ with d odd;
- in $\mu - \mu_o$ and $-(\mu - \mu_o) \ln(\mu - \mu_o)$ for $2n_o + d = 2, 6$;
- in $\mu - \mu_o$ and $e^{-1/(\mu - \mu_o)}$ for $2n_o + d = 4$;
- in $(\mu - \mu_o)^{1/2}$, $-(\mu - \mu_o) \ln(\mu - \mu_o)$, $(-\frac{1}{\ln(\mu - \mu_o)})^{1/2}$ and $-\frac{\ln \ln(\mu - \mu_o)^{-1}}{\ln(\mu - \mu_o)}$ for $2n_o + d = 8$;
- in $(\mu - \mu_o)^{1/2}$ and $-(\mu - \mu_o)^{1/2} \ln(\mu - \mu_o)$ for $2n^o + d \geq 10$ with d even

(see Theorem 2.4). Moreover, resonance states of 0-energy, i.e. non-zero solutions f of $\widehat{H}_{\mu_o} f = 0$ not belonging to $\ell^2(\mathbb{Z}^d)$ appear if and only if $2n_o + d \in \{5, 6, 7, 8\}$. Recall that the emergence of 0-energy resonances in more lattice dimensions could allow the Efimov effect to be observed in other dimensions than $d = 3$.

Furthermore, observing that the top $e(\vec{\pi}) = 4d^2$ of the essential spectrum is non-degenerate, one expects the asymptotics of $e(\mu)$ as $\mu \rightarrow -\mu^o$ to be similar as in the discrete Laplacian case [20, 21]; more precisely, depending on d and n^o , $e(\mu)$ has a convergent expansion

- in $\mu + \mu^o$ for $2n^o + d = 1, 3$;

- in $(\mu + \mu^o)^{1/2}$ for $2n^o + d \geq 5$ with d odd;
- in $\mu + \mu^o$ and $e^{-1/(\mu+\mu^o)}$ for $2n^o + d = 2$;
- in $\mu + \mu^o$, $-\frac{1}{\ln(\mu+\mu^o)}$ and $-\frac{\ln \ln(\mu+\mu^o)^{-1}}{\ln(\mu+\mu^o)}$ for $2n^o + d = 4$;
- in $\mu + \mu^o$ and $-(\mu + \mu^o) \ln(\mu + \mu^o)$ for $2n^o + d \geq 6$ with d even

(see Theorem 2.5). Moreover, the resonance states of energy $4d^2$, i.e. non-zero solutions f of $\widehat{H}_{-\mu^o} f = 4d^2 f$ not belonging to $\ell^2(\mathbb{Z}^d)$ appear if and only if $2n^o + d = 3, 4$.

The threshold analysis for more general class of nonlocal discrete Schrödinger operators with δ -potential of type

$$\widehat{H}_\mu = \Psi(-\widehat{\Delta}) + \mu\delta_{x0},$$

can be found in [14], where Ψ is some strictly increasing C^1 -function and δ_{x0} is the Dirac's delta-function supported at 0. Besides the existence of eigenvalues, authors of [14] classify (embedded) threshold resonances and threshold eigenvalues depending on the behaviour of Ψ at the edges of the essential spectrum of $-\widehat{\Delta}$ and on the lattice dimension d . The eigenvalue expansions for the discrete bilaplacian with δ -perturbation have been established in [17] for $d = 1$ using the complex analytic methods.

The paper is organized as follows. In Sect. 2 after introducing some preliminaries we state the main results of the paper. In Theorem 2.2 we establish necessary and sufficient conditions for non-emptiness of the discrete spectrum of \widehat{H}_μ , and in case of existence, we study the location and the uniqueness, analyticity, monotonicity and convexity properties of eigenvalues $e(\mu)$ as a function of μ . In particular, we study the asymptotics of $e(\mu)$ as $\mu \rightarrow \mu^o$ and $\mu \rightarrow -\mu^o$ as well as $\mu \rightarrow \pm\infty$. As discussed above in Theorems 2.4 and 2.5 we obtain expansions of $e(\mu)$ for small and positive $\mu - \mu^o$ and $\mu + \mu^o$. In Sect. 3 we prove the main results. The main idea of the proof is to obtain a nonlinear equation $\Delta(\mu; z) = 0$ with respect to the eigenvalue $z = e(\mu)$ of \widehat{H}_μ and then study properties of $\Delta(\mu; z)$. Finally, in appendix Section A we obtain the asymptotics of certain integrals related to $\Delta(\mu; z)$ which will be used in the proofs of main results.

Data availability statement

We confirm that the current manuscript has no associated data.

2 Preliminary and main results

Let \mathbb{Z}^d be the d -dimensional lattice and $\ell^2(\mathbb{Z}^d)$ be the Hilbert space of square-summable functions on \mathbb{Z}^d . Consider the family

$$\widehat{H}_\mu := \widehat{H}_0 - \mu\widehat{V}, \quad \mu \geq 0,$$

of self-adjoint bounded discrete Schrödinger operators in $\ell^2(\mathbb{Z}^d)$. Here $\widehat{H}_0 := \widehat{\Delta}\widehat{\Delta}$ is discrete bilaplacian, where

$$\widehat{\Delta}f(x) = \frac{1}{2} \sum_{|s|=1} [f(x) - f(x + s)], \quad f \in \ell^2(\mathbb{Z}^d),$$

is the discrete Laplacian, and \widehat{V} is a rank-one operator

$$\widehat{V}\widehat{f}(x) = \widehat{v}(x) \sum_{y \in \mathbb{Z}^d} \widehat{v}(y)\widehat{f}(y),$$

where $\widehat{v} \in \ell^2(\mathbb{Z}^d) \setminus \{0\}$ is a given real-valued function.

Let \mathbb{T}^d be the d -dimensional torus equipped with the Haar measure and $L^2(\mathbb{T}^d)$ be the Hilbert space of square-integrable functions on \mathbb{T}^d . By \mathcal{F} we denote the the standard Fourier transform

$$\mathcal{F} : \ell^2(\mathbb{Z}^d) \rightarrow L^2(\mathbb{T}^d), \quad \mathcal{F}\widehat{f}(p) = \frac{1}{(2\pi)^{d/2}} \sum_{x \in \mathbb{Z}^d} \widehat{f}(x)e^{ix \cdot p}.$$

Further we always assume that \widehat{v} and its Fourier image

$$v(p) := \mathcal{F}\widehat{v}(p) = \frac{1}{(2\pi)^{d/2}} \sum_{x \in \mathbb{Z}^d} \widehat{v}(x)e^{ix \cdot p}$$

satisfy the following assumptions:

There exist reals $C, a > 0$ and nonnegative integers $n_0, n^o \geq 0$ such that

$$|\widehat{v}(x)| \leq Ce^{-a|x|} \quad \text{for all } x \in \mathbb{Z}^d, \tag{H1}$$

$$|v(0)|^2 = D^2|v(0)|^2 = \dots = D^{2n_0-2}|v(0)|^2 = 0, \quad D^{2n_0}|v(0)|^2 \neq 0, \tag{H2}$$

$$|v(\vec{\pi})|^2 = D^2|v(\vec{\pi})|^2 = \dots = D^{2n^o-2}|v(\vec{\pi})|^2 = 0, \quad D^{2n^o}|v(\vec{\pi})|^2 \neq 0, \tag{H3}$$

here $D^j f(p)$ is the j -th order differential of f at p , i.e. the j -th order symmetric tensor

$$D^j f(p) \underbrace{[w, \dots, w]}_{j\text{-times}} = \sum_{i_1 + \dots + i_d = j, i_k \geq 0} \frac{\partial^j f(p)}{\partial^{i_1} p_1 \dots \partial^{i_d} p_d} w_1^{i_1} \dots w_d^{i_d},$$

$$w = (w_1, \dots, w_d) \in \mathbb{R}^d,$$

and $\vec{\pi} = (\pi, \dots, \pi) \in \mathbb{T}^d$. Notice that under assumption (H1), v is analytic on \mathbb{T}^d .

Recall that $\sigma(\widehat{\Delta}) = \sigma_{\text{ess}}(\widehat{\Delta}) = [0, 2d]$ (see e.g. [1]). Hence, $\sigma(\widehat{H}_0) = \sigma_{\text{ess}}(\widehat{H}_0) = [0, 4d^2]$, and by the compactness of \widehat{V} and Weyl's Theorem,

$$\sigma_{\text{ess}}(\widehat{H}_\mu) = \sigma_{\text{ess}}(\widehat{H}_0) = [0, 4d^2]$$

for any $\mu \in \mathbb{R}$.

Before stating the main results let us introduce the constants

$$\mu_o := \left(\int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{\epsilon(q)} \right)^{-1}, \quad \mu^o := \left(\int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{4d^2 - \epsilon(q)} \right)^{-1}, \quad (2.2)$$

$$\widehat{c}_v := \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{\epsilon(q)^2}, \quad \widehat{C}_v := \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{(4d^2 - \epsilon(q))^2}, \quad (2.3)$$

and

$$c_v := \frac{2^{2n_o+d}}{(2n_o)!} \int_{\mathbb{S}^{d-1}} D^{2n_o} |v(0)|^2 [w, \dots, w] d\mathcal{H}^{d-1}(w), \quad (2.4)$$

$$C_v := \frac{2^{2n_o+d-1}}{(8d)^{n_o+d/2} (2n_o)!} \int_{\mathbb{S}^{d-1}} D^{2n_o} |v(\vec{\pi})|^2 [w, \dots, w] d\mathcal{H}^{d-1}(w), \quad (2.5)$$

where \mathbb{S}^{d-1} is the unit sphere in \mathbb{R}^d and

$$\epsilon(q) := \left(\sum_{i=1}^d (1 - \cos q_i) \right)^2.$$

Remark 2.1 Under assumptions (H1)–(H3), $\mu_o, \mu^o \geq 0$, $c_v, C_v > 0$, and $\widehat{c}_v, \widehat{C}_v \in (0, +\infty]$. Moreover, by Propositions A.1 and A.2:

- $\mu_o = 0$ (resp. $\mu^o = 0$) if and only if $2n_o + d \leq 4$ (resp. $2n_o + d \leq 2$);
- $\widehat{c}_v < \infty$ (resp. $\widehat{C}_v < \infty$) if $2n_o + d \geq 9$ (resp. $2n_o + d \geq 5$).

2.1 Main results

First we concern with the existence of the discrete spectrum of \widehat{H}_μ .

Theorem 2.2 *Let $\mu_o, \mu^o \geq 0$ be given by (2.2). Then $\sigma_{\text{disc}}(\widehat{H}_\mu) = \emptyset$ for any $\mu \in [-\mu^o, \mu_o]$ and $\sigma_{\text{disc}}(\widehat{H}_\mu)$ is a singleton $\{e(\mu)\}$ for any $\mu \in \mathbb{R} \setminus [-\mu^o, \mu_o]$. Moreover, the associated eigenfunction \widehat{f}_μ to $e(\mu)$ is given by $\widehat{f}_\mu := \mathcal{F}^* f_\mu$, where*

$$f_\mu(p) = \frac{v(p)}{\epsilon(p) - e(\mu)}.$$

Furthermore, if $\mu < -\mu^o$ (resp. $\mu > \mu_o$), then $e(\mu) > 4d^2$ (resp. $e(\mu) < 0$). Moreover, the function $\mu \in \mathbb{R} \setminus [-\mu^o, \mu_o] \mapsto e(\mu)$ is real-analytic strictly decreasing, convex in $(-\infty, -\mu^o)$ and concave in $(\mu_o, +\infty)$, and satisfies

$$\lim_{\mu \searrow -\mu^o} e(\mu) = 0 \quad \text{and} \quad \lim_{\mu \nearrow -\mu^o} e(\mu) = 4d^2 \quad (2.6)$$

and

$$\lim_{\mu \rightarrow \pm\infty} \frac{e(\mu)}{\mu} = - \int_{\mathbb{T}^d} |v(q)|^2 dq. \tag{2.7}$$

Next we study the threshold resonances of \widehat{H}_μ .

Theorem 2.3 *Let $n_o, n^o \geq 0$ be given by (H2)–(H3).*

(a) *Let $2n_o + d \geq 5$. Then $\widehat{f} := \mathcal{F}^* f \in c_0(\mathbb{Z}^d)$, i.e., $\widehat{f}(x) \rightarrow 0$ as $|x| \rightarrow +\infty$, where*

$$f(p) = \frac{v(p)}{\epsilon(p)} \in L^1(\mathbb{T}^d).$$

Moreover, $\widehat{f} \in c_0(\mathbb{Z}^d) \setminus \ell^2(\mathbb{Z}^d)$ for $2n_o + d \in \{5, 6, 7, 8\}$, $\widehat{f} \in \ell^2(\mathbb{Z}^d)$ for $2n_o + d \geq 9$, and \widehat{f} solves the equation $\widehat{H}_{\mu_o} f = 0$.

(b) *Let $2n^o + d \geq 3$. Then $\widehat{g} := \mathcal{F}^* g \in \ell^0(\mathbb{Z}^d)$, where*

$$g(p) = \frac{v(p)}{4d^2 - \epsilon(p)}.$$

Moreover, $\widehat{g} \in \ell^0(\mathbb{Z}^d) \setminus \ell^2(\mathbb{Z}^d)$ for $2n^o + d \in \{3, 4\}$, $\widehat{g} \in \ell^2(\mathbb{Z}^d)$ for $2n^o + d \geq 5$, and \widehat{g} solves the equation $\widehat{H}_{-\mu_o} f = 4d^2 f$.

We recall that in the literature the non-zero solutions of equations $\widehat{H}_\mu \widehat{f} = 0$ and $\widehat{H}_\mu \widehat{g} = 4d^2 \widehat{g}$ not belonging to $\ell^2(\mathbb{Z}^d)$ are called the *resonance states* [1, 2].

Now we study the rate of the convergences in (2.6).

Theorem 2.4 (Expansions of $e(\mu)$ at $\mu = \mu_o$) *For $\mu > \mu_o$ let $e(\mu) < 0$ be the eigenvalue of \widehat{H}_μ .*

(a) *Suppose that d is odd:*

(a1) *if $2n_o + d = 1, 3$, then $\mu_o = 0$ and for sufficiently small and positive μ ,*

$$(-e(\mu))^{1/4} = \begin{cases} \left(\frac{\pi c_v}{4}\right)^{1/3} \mu^{1/3} + \sum_{n \geq 1} c_{1,n} \mu^{\frac{n+1}{3}}, & 2n_o + d = 1, \\ \frac{\pi c_v}{8} \mu + \sum_{n \geq 1} c_{3,n} \mu^{n+1}, & 2n_o + d = 3, \end{cases}$$

where $\{c_{1,n}\}$ and $\{c_{3,n}\}$ are some real coefficients;

(a2) *if $2n_o + d = 5, 7$, then $\mu_o > 0$ and for sufficiently small and positive $\mu - \mu_o$,*

$$(-e(\mu))^{1/4} = \begin{cases} \frac{8}{\pi c_v \mu_o^2} (\mu - \mu_o) + \sum_{n \geq 1} c_{5,n} (\mu - \mu_o)^{n+1}, & 2n_o + d = 5, \\ \left(\frac{8}{\pi c_v \mu_o^2}\right)^{1/3} (\mu - \mu_o)^{1/3} + \sum_{n \geq 1} c_{7,n} (\mu - \mu_o)^{\frac{n+1}{3}}, & 2n_o + d = 7, \end{cases}$$

where $\{c_{5,n}\}$ and $\{c_{7,n}\}$ are some real coefficients;

(a3) if $2n_o + d \geq 9$, then $\mu_o > 0$ and for sufficiently small and positive $\mu - \mu_o$,

$$(-e(\mu))^{1/4} = (\mu_o^2 \widehat{c}_v)^{-1/4} (\mu - \mu_o)^{1/4} + \sum_{n \geq 1} c_{9,n} (\mu - \mu_o)^{n/4},$$

where $\{c_{9,n}\}$ are some real coefficients.

(b) Suppose that d is even:

(b1) if $2n_o + d = 2, 4$, then $\mu_o = 0$ and for sufficiently small and positive μ ,

$$(-e(\mu))^{1/2} = \begin{cases} \frac{\pi c_v}{8} \mu + \sum_{n+m \geq 1, n, m \geq 0} c_{2,nm} \mu^{n+1} (-\mu \ln \mu)^m, & 2n_o + d = 2, \\ c e^{-\frac{8}{c_v \mu}} + \sum_{n+m \geq 1, n, m \geq 0} c_{4,nm} \mu^{n+1} \left(\frac{1}{\mu} e^{-\frac{8}{c_v \mu}}\right)^{m+1}, & 2n_o + d = 4, \end{cases}$$

where $\{c_{2,nm}\}$ and $\{c_{4,nm}\}$ are some real coefficients and $c > 0$;

(b2) if $2n_o + d = 6, 8$, then $\mu_o > 0$ and for sufficiently small and positive $\mu - \mu_o$,

$$(-e(\mu))^{1/2} = \begin{cases} \frac{8}{\pi c_v \mu_o^2} \tau^2 + \sum_{n+m \geq 1, n, m \geq 0} c_{6,nm} \tau^{2n+2} \theta^m, & 2n_o + d = 6, \\ \left(\frac{8}{c_v \mu_o^2}\right)^{1/2} \tau \sigma + \sum_{n+m+k \geq 1, n, m, k \geq 0} c_{8,nmk} \tau^{n+1} \sigma^{m+1} \eta^k, & 2n_o + d = 8, \end{cases}$$

where $\{c_{4,nm}\}$ and $\{c_{8,nmk}\}$ are some real coefficients and

$$\tau := (\mu - \mu_o)^{1/2}, \quad \theta := -\tau^2 \ln \tau, \quad \sigma := \left(-\frac{1}{\ln \tau}\right)^{1/2}, \quad \eta := -\frac{\ln \ln \tau^{-1}}{\ln \tau}, \tag{2.8}$$

(b3) if $2n_o + d \geq 10$, then $\mu_o > 0$ and for sufficiently small and positive $\mu - \mu_o$,

$$(-e(\mu))^{1/2} = (\mu_o^2 \widehat{c}_v)^{-1/2} \tau + \sum_{n+m \geq 1, n, m \geq 0} c_{10,nm} \tau^{n+1} \theta^m,$$

where $\{c_{10,nm}\}$ are some real coefficients.

Here $c_v > 0$ and $\widehat{c}_v > 0$ are given by (2.4) and (2.3), respectively.

Theorem 2.5 (Expansions of $e(\mu)$ at $\mu = -\mu^o$) For let $\mu < -\mu^o$ let $e(\mu) > 4d^2$ be the eigenvalue of \widehat{H}_μ .

(a) Suppose that d is odd:

(a1) if $2n^o + d = 1$, then $\mu^o = 0$ and for sufficiently small and negative μ ,

$$(e(\mu) - 4d^2)^{1/2} = -\pi C_v \mu + \sum_{n \geq 1} C_{1,n} \mu^{n+1},$$

where $\{C_{1,n}\}$ are some real coefficients;

(a2) if $2n^o + d = 3$, then $\mu^o > 0$ and for sufficiently small and positive $\mu + \mu^o$,

$$(e(\mu) - 4d^2)^{1/2} = (\pi C_v \mu^{o2})^{-1} (\mu + \mu^o) + \sum_{n \geq 1} C_{3,n} (\mu + \mu^o)^{n+1},$$

where $\{C_{3,n}\}$ and $\{C_{7,n}\}$ are some real coefficients;

(a3) if $2n^o + d \geq 5$, then $\mu^o > 0$ and for sufficiently small and positive $\mu + \mu^o$,

$$(e(\mu) - 4d^2)^{1/2} = (\widehat{C}_v \mu^{o2})^{-1/2} (\mu + \mu^o)^{1/2} + \sum_{n \geq 1} C_{5,n} (\mu + \mu^o)^{(n+1)/2},$$

where $\{C_{5,n}\}$ are some real coefficients.

(b) Suppose that d is even:

(b1) if $2n_o + d = 2$, then $\mu_o = 0$ and for sufficiently small and negative μ ,

$$e(\mu) - 4d^2 = C e^{\frac{1}{C_v \mu}} + \sum_{n+m \geq 1, n, m \geq 0} C_{2,nm} \mu^{n+1} \left(-\frac{1}{\mu} e^{\frac{1}{C_v \mu}}\right)^{m+1},$$

where $\{C_{2,nm}\}$ are some real coefficients and $C > 0$;

(b2) if $2n_o + d = 4$, then $\mu^o > 0$ and for sufficiently small and positive $\mu + \mu^o$,

$$e(\mu) - 4d^2 = (C_v \mu^{o2})^{-1} \mu \sigma + \sum_{n+m+k \geq 1, n, m, k \geq 0} C_{4,nmk} \tau^{n+1} \sigma^{m+1} \eta^k,$$

where $\{C_{4,nm}\}$ are some real coefficients and

$$\tau := \mu + \mu^o, \quad \sigma := -\frac{1}{\ln \tau}, \quad \eta := -\frac{\ln \ln \tau^{-1}}{\ln \tau};$$

(b3) if $2n_o + d \geq 6$, then $\mu^o > 0$ and for sufficiently small and positive $\mu + \mu^o$,

$$e(\mu) - 4d^2 = (\widehat{C}_v \mu^{o2})^{-1} (\mu + \mu^o) + \sum_{n+m \geq 1, n, m \geq 0} C_{6,nm} (\mu + \mu^o)^{n+1} [-(\mu + \mu^o) \ln(\mu + \mu^o)]^m,$$

where $\{C_{6,nm}\}$ are some real coefficients.

Here C_v and \widehat{C}_v are given by (2.5) and (2.3), respectively.

Remark 2.6 Few comments on the main results are in order.

1. The assertions of Theorem 2.2 hold in fact for any $\widehat{v} \in \ell^2(\mathbb{Z}^d)$ (see Remark 3.2);
2. Similar expansions of $e(\mu)$ in Theorems 2.4 and 2.5 at $\mu = \mu_o$ and $\mu = -\mu^o$, respectively, still hold for any exponentially decaying $\widehat{v} : \mathbb{Z}^d \rightarrow \mathbb{C}$ (see Remark 3.3);
3. If \widehat{v} decays at most polynomially at infinity, i.e. $\widehat{v}(x) = O(|x|^{-\alpha})$ for some $\alpha > 0$, then instead of the expansions in Theorem 2.4 and 2.5 we obtain only asymptotics of $e(\mu)$ (see Remark 3.4).

3 Proof of main results

In this section we prove the main results. By the Birman-Schwinger principle and the Fredholm Theorem we have

Lemma 3.1 *A complex number $z \in \mathbb{C} \setminus [0, 4d^2]$ is an eigenvalue of \widehat{H}_μ if and only if*

$$\Delta(\mu; z) := 1 - \mu \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{\epsilon(q) - z} = 0.$$

Proof of Theorem 2.2 By the definition of μ_o , for any $\mu < -\mu^o$:

$$\lim_{z \nearrow -\mu^o} \Delta(\mu; z) = 1 + \frac{\mu}{\mu^o} < 0, \quad \lim_{z \rightarrow +\infty} \Delta(\mu; z) = 1.$$

Since $\Delta(\mu; z) > 1$ for $z < 0$ and $\mu > -\mu^o$, in view of the strict monotonicity $\Delta(\mu; \cdot)$ in $(4d^2, \infty)$, for any $\mu < -\mu^o$ there exists a unique $e(\mu) \in (4d^2, +\infty)$ such that $\Delta(\mu; e(\mu)) = 0$. Analogously, for any $\mu > \mu_o$ there exists a unique $e(\mu) \in (-\infty, 0)$ such that $\Delta(\mu; e(\mu)) = 0$. By the Implicit Function Theorem the function $\mu \in \mathbb{R} \setminus [-\mu^o, \mu_o] \mapsto e(\mu)$ is real-analytic. Moreover, computing the derivatives of the implicit function $e(\mu)$ we find:

$$e'(\mu) = -\frac{1}{\mu} \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{\epsilon(q) - e(\mu)} \left(\int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{(\epsilon(q) - e(\mu))^2} \right)^{-1}, \quad \mu \neq 0, \quad (3.1)$$

thus, using $\mu(\epsilon(q) - e(\mu)) > 0$ we get $e'(\mu) < 0$, i.e. $e(\cdot)$ is strictly decreasing in $\mathbb{R} \setminus \{0\}$. Differentiating (3.1) one more time we get

$$e''(\mu) = \frac{2e'(\mu)}{\mu} \left(1 - \mu e'(\mu) \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{(\epsilon(q) - e(\mu))^3} \left(\int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{(\epsilon(q) - e(\mu))^2} \right)^{-1} \right).$$

Therefore, $e''(\mu) > 0$ (i.e. $e(\cdot)$ is strictly convex) for $\mu < 0$ and $e''(\mu) < 0$ (i.e. $e(\cdot)$ is strictly concave) for $\mu > 0$.

To prove (2.7), first we let $\mu \rightarrow \pm\infty$ in

$$1 = \mu \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{\epsilon(q) - e(\mu)} \quad (3.2)$$

and find $\lim_{\mu \rightarrow \pm\infty} e(\mu) = \mp\infty$. In particular, if $|\mu|$ is sufficiently large, $|\frac{e(q)}{e(\mu)}| < \frac{1}{2}$ and hence, by (3.2) and the Dominated Convergence Theorem,

$$\lim_{\mu \rightarrow \pm\infty} \frac{e(\mu)}{\mu} = - \lim_{\mu \rightarrow \pm\infty} \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{1 - \frac{e(q)}{e(\mu)}} = - \int_{\mathbb{T}^d} |v(q)|^2 dq.$$

To prove that \widehat{f}_μ solves $\widehat{H}_\mu \widehat{f}_\mu = e(\mu) \widehat{f}_\mu$ we consider the equivalent equality $\mathcal{F}\widehat{H}_\mu \mathcal{F}^* f_\mu = e(\mu) f_\mu$, which is easily reduced to the equality $\Delta(\mu; e(\mu)) = 0$. □

Remark 3.2 In view of Lemma 3.1 and the proof of Theorem 2.2, their assertions still hold for any $v \in \ell^2(\mathbb{Z}^d)$.

Proof of Theorem 2.3 We prove only (a), the proof of (b) being similar. Repeating the proof of the continuity (resp. differentiability) of ι_f at $z = 0$ in Proposition A.1 one can show that $f \in L^1(\mathbb{T}^d) \setminus L^2(\mathbb{T}^d)$ for $2n_o + d \in \{5, 6, 7, 8\}$ and $f \in L^2(\mathbb{T}^d)$ for $2n_o + d \geq 9$. Thus, by the Riemann-Lebesgue Lemma, $\widehat{f} \in \ell^0(\mathbb{Z}^d)$. To show that $\widehat{H}_{\mu_o} \widehat{f} = 0$ it suffices to observe that $\mathcal{F}\widehat{H}_{\mu_o} \mathcal{F}^* f = 0$. □

Proof of Theorem 2.4 Since

$$|v(p)|^2 = (2\pi)^{-d} \left(\sum_{x \in \mathbb{Z}^d} \widehat{v}(x) \cos p \cdot x \right)^2 + (2\pi)^{-d} \left(\sum_{x \in \mathbb{Z}^d} \widehat{v}(x) \sin p \cdot x \right)^2, \tag{3.3}$$

the function $p \in \mathbb{T}^d \mapsto |v(p)|^2$ is nonnegative even real-analytic function. Notice also that if $n_o \geq 1$, then by the nonnegativity of $|v|^2$, $p = 0$ is a global minimum for $|v|^2$. Therefore, the tensor $D^{2n_o} |v(0)|^2$ is positively definite and

$$c_v := \frac{2^{2n_o+d}}{(2n_o)!} \int_{\mathbb{S}^{d-1}} D^{2n_o} |v(0)|^2 [w, \dots, w] d\mathcal{H}^{d-1} > 0.$$

Note that

$$\widehat{c}_v = \iota'_{|v|^2}(0) = \int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{\epsilon(q)^2},$$

where ι_f is defined in (A.1). By Proposition A.1, $f(p) = \frac{v(p)}{\epsilon(p)} \in L^2(\mathbb{T}^d)$ if and only if $2n_o + d \geq 9$. Moreover, by definition, $\mu_o > 0$ and $\Delta(\mu_o; 0) = 0$ for $2n_o + d \geq 5$, and hence, as in the proof of Lemma 3.1 for such d one can show that $H_{\mu_o} f = 0$.

In view of the strict monotonicity and (2.6) there exists a unique $\mu_1 > 0$ such that $e(\mu) \in (-\frac{1}{128}, 0)$ for any $\mu \in (0, \mu_1)$. Since

$$\mu = (\iota_{|v|^2}(e(\mu)))^{-1}, \tag{3.4}$$

we can use Proposition A.1 with $f = |v|^2$ and $e := e(\mu)$, to find the expansions of the inverse function $\mu := \mu(e)$. Then applying the appropriate versions of the Implicit Function Theorem in analytical case we get the expansions of $e = e(\mu)$. Notice that from (A.3) and (A.4) as well as (3.5) it follows that $\mu_o = 0$ for $2n_o + d \leq 4$ and $\mu_o = \left(\int_{\mathbb{T}^d} \frac{|v(q)|^2 dq}{\epsilon(q)} \right)^{-1} > 0$ for $2n_o + d \geq 5$.

(a) Suppose that d is odd. In view of the expansions (A.3) of l_f , in this case, (3.4) is reduced to the inverting the equation

$$\mu = g(\alpha), \tag{3.5}$$

where $\alpha := (-e)^{1/4}$ and g is an analytic function around $\alpha = 0$.

Case $2n_o + d = 1$. In this case by (A.3),

$$g(\alpha) := \frac{\alpha^3}{c_1^3 + \sum_{n \geq 1} a_n \alpha^n},$$

where $\{a_n\} \subset \mathbb{R}$ and $c_1 := (\pi c_v/4)^{1/3}$ and (3.5) is equivalently represented as

$$\alpha = \mu \left(c_1^3 + \sum_{n \geq 1} a_n \alpha^n \right)^{1/3}, \tag{3.6}$$

where $\mu = \mu^{1/3}$. Now setting

$$\alpha = \mu(c_1 + u), \tag{3.7}$$

and using the Taylor series of $(c_1^3 + x)^{1/3}$, for μ and u sufficiently small we rewrite (3.6) as

$$F(u, \mu) := u - \sum_{n \geq 1} \tilde{a}_n \mu^n (c_1 + u)^n = 0, \tag{3.8}$$

where $F(\cdot, \cdot)$ is analytic at $(u, \mu) = (0, 0)$, $F(0, 0) = 0$ and $F_u(0, 0) = 1$. Hence, by the Implicit Function Theorem, there exists $\gamma_1 > 0$ such that for $|\mu| < \gamma_1$, (3.8) has a unique real-analytic solution $u = u(\mu)$ which can be represented as an absolutely convergent series $u = \sum_{n \geq 1} b_n \mu^n$. Putting this in (3.7) and recalling the definitions of α and μ we get the expansion of $(-e(\mu))^{1/4}$ for $\mu > 0$ small.

Case $2n_o + d = 3$. By (A.3),

$$g(\alpha) = \alpha \left(c_3 + \sum_{n \geq 1} a_n \alpha^n \right)^{-1}, \tag{3.9}$$

where $\{a_n\} \subset \mathbb{R}$ and $c_3 := \pi c_v/8$, and hence, (3.5) is represented as

$$\alpha = \mu \left(c_3 + \sum_{n \geq 1} a_n \alpha^n \right).$$

Then setting $\alpha = \mu(c_3 + u)$ we rewrite (3.9) in the form (3.8), and as in the case of $2n_o + d = 1$, we get the expansion of $(-e(\mu))^{1/4}$.

Case $2n_o + d = 5$. In this case by (A.3)

$$g(\alpha) = \left(\frac{1}{\mu_o} - \frac{\pi c_v \alpha}{8} \left(1 + \sum_{n \geq 1} a_n \alpha^n \right) \right)^{-1},$$

where $\{a_n\} \subset \mathbb{R}$, and hence, by (3.5),

$$\frac{\mu - \mu_o}{\mu \mu_o} = \frac{\pi c_v \alpha}{8} \left(1 + \sum_{n \geq 1} a_n \alpha^n \right). \tag{3.10}$$

Note that if $|\mu - \mu_o| < \mu_o$, then

$$\frac{\mu - \mu_o}{\mu \mu_o} = \frac{\mu - \mu_o}{\mu_o^2 + \mu_o(\mu - \mu_o)} = \frac{\mu - \mu_o}{\mu_o^2} \sum_{n \geq 0} \left(\frac{\mu - \mu_o}{\mu_o} \right)^n, \tag{3.11}$$

thus from (3.10) we get

$$\alpha = (\mu - \mu_o) \left(c_5 + c_5 \sum_{n \geq 1} \mu_o^{-n} (\mu - \mu_o)^n \right) \left(1 + \sum_{n \geq 1} a_n \alpha^n \right)^{-1}.$$

and $c_5 := 8/(\pi c_v \mu_o^2)$. Now setting $\alpha = (\mu - \mu_o)(c_5 + u)$ for sufficiently small and positive $\mu - \mu_o$ we get

$$u = \sum_{n,m \geq 1} \tilde{c}_{n,m} (\mu - \mu_o)^n (c_5 + u)^m,$$

where $\tilde{c}_{n,m} \subset \mathbb{R}$. By the Implicit Function Theorem, for sufficiently small $\mu - \mu_o$ there exists a unique real-analytic function $u = u(\mu)$ given by the absolutely convergent series $u(\mu) = \sum_{n \geq 1} b_n (\mu - \mu_o)^n$. By the definition of α , this implies the expansion of $(-e(\mu))^{1/4}$.

Case $2n_o + d = 7$. As the previous case, by (A.3) and (3.11), the equation (3.5) is represented as

$$(\mu - \mu_o) \left(c_7^3 + c_7^3 \sum_{n \geq 1} \mu_o^{-n} (\mu - \mu_o)^n \right) = \alpha^3 \left(1 + \sum_{n \geq 1} a_n \alpha^n \right), \tag{3.12}$$

where $\{a_n\} \subset \mathbb{R}$ and $c_7 := [8/(\pi c_v \mu_o^2)]^{1/3}$. When $\mu - \mu_o > 0$ is small enough, by the Taylor series of $(1 + x)^{\pm 1/3}$ at $x = 0$, (3.12) is equivalently rewritten as

$$\alpha = (\mu - \mu_o)^{1/3} \left(c_7 + \sum_{n \geq 1} \tilde{c}_n (\mu - \mu_o)^n \right) \left(1 + \sum_{n \geq 1} \tilde{a}_n \alpha^n \right), \tag{3.13}$$

Thus, for $\rho = (\mu - \mu_o)^{1/3}$, setting $\alpha = \rho (c_7 + u)$ in (3.13), for sufficiently small and positive ρ we get

$$u = \sum_{n,m \geq 1} \tilde{c}_{n,m} \rho^n (c_7 + u)^m.$$

By the Implicit Function Theorem, this equation has a unique real-analytic solution $u = u(\rho)$ given by the absolutely convergent series $u = \sum_{n \geq 1} b_n \rho^n$. This, definitions

of α and ρ imply the expansion of $(-e(\mu))^{1/4}$.

Case $2n_o + d = 9$. In this case by (A.3) and (3.11)

$$(\mu - \mu_o) \left(c_9^4 + c_9^4 \sum_{n \geq 1} \mu_o^{-n} (\mu - \mu_o)^n \right) = \alpha^4 \left(1 + \sum_{n \geq 1} a_n \alpha^n \right), \tag{3.14}$$

where $\{a_n\} \subset \mathbb{R}$ and $c_9 := (\mu_o^2 \hat{c}_v)^{-1/4}$. Thus, for sufficiently small and positive $\mu - \mu_o$ using the Taylor series of $(1 + x)^{\pm 1/4}$ at $x = 0$, this equation can also be represented as

$$\alpha = \rho \left(c_9 + \sum_{n \geq 1} \tilde{b}_n \rho^{4n} \right) \left(1 + \sum_{n \geq 1} \tilde{a}_n \alpha^n \right),$$

where $\rho := (\mu - \mu_o)^{1/4}$. Now setting $\alpha = \rho (c_9 + u)$ in (3.14) we get

$$u = \sum_{n,m \geq 1} \tilde{c}_{n,m} \rho^n (c_9 + u)^m,$$

and the expansion of $(-e(\mu))^{1/4}$ follows as in the case of $2n_o + d = 7$.

(b) Suppose that d is even. In view of the expansion (A.3) of l_f , in this case, (3.4) is reduced to the inverting the equation

$$\mu = \frac{\alpha^l}{g(\alpha) + h(\alpha) \ln \alpha}, \tag{3.15}$$

where $\alpha := (-e)^{1/2}$, $l \in \mathbb{N}_0$, and g and h are analytic around $\alpha = 0$. Presence of $\ln \alpha$ implies that unlike the case of odd dimensions, α is not necessarily analytic with respect to μ^s . Therefore, we need to introduce new variables dependent on $\ln \mu$ to reduce the problem to the Implicit Function Theorem.

Case $2n_o + d = 2$. By (A.4), in this case for $c_2 := \pi c_v/8$

$$l = 1, \quad g(\alpha) = c_2 + \sum_{n \geq 1} a_n \alpha^n, \quad h(\alpha) = \sum_{n \geq 1} b_n \alpha^{2n}.$$

Hence, setting

$$\alpha = \mu(c_2 + u) \tag{3.16}$$

and $\tau = -\mu \ln \mu$ we represent (3.15) as

$$\begin{aligned} F(u, \mu, \tau) := & u - \sum_{n \geq 1} a^n \mu^n (c_2 + u)^n + \ln(c_2 + u) \sum_{n \geq 1} b^n \mu^n (c_2 + u)^n \\ & - \tau \sum_{n \geq 1} b^n \mu^{n-1} (c_2 + u)^n = 0, \end{aligned}$$

where F is analytic around $(0, 0, 0)$, $F(0, 0, 0) = 0$, $F_u(0, 0, 0) = 1$. Hence, by the Implicit Function Theorem, there exists a unique real-analytic function $u = u(\mu, \tau)$ given by the convergent series $u(\mu, \tau) = \sum_{n+m \geq 1, n, m \geq 0} \tilde{c}_{n,m} \mu^n \tau^m$ for sufficiently small $|\mu|$ and $|\tau|$, which satisfies $F(u(\mu, \tau), \mu, \tau) \equiv 0$. Inserting u in (3.16) we get the expansion of $\alpha = (-e)^{1/2}$.

Case $2n_o + d = 4$. In this case, by (A.4) for $c_4 := 8/c_v$

$$l = 0, \quad g(\alpha) = \sum_{n \geq 0} a_n \alpha^n, \quad h(\alpha) = -c_4 + \sum_{n \geq 1} b_n \alpha^{2n}.$$

Letting $\alpha = e^{-\frac{1}{c_4 \mu}}(c + u)$, where $c = e^{a_0/c_4} > 0$, we represent (3.15) as

$$\begin{aligned} \ln(c + u) - b_0 = & \frac{1}{\mu} e^{-\frac{1}{c_4 \mu}} \sum_{n \geq 1} a^n e^{-\frac{n-1}{c_4 \mu}} (c + u)^n \\ & + \ln(c + u) \sum_{n \geq 1} b^n e^{-\frac{n}{c_4 \mu}} (c + u)^n - \sum_{n \geq 1} a^n e^{-\frac{n}{c_4 \mu}} (c + u)^n = 0. \end{aligned} \tag{3.17}$$

Writing $\tau := \frac{1}{\mu} e^{-\frac{1}{c_4\mu}}$ so that $e^{-\frac{1}{c_4\mu}} = \mu\tau$, (3.17) is represented as

$$F(u, \mu, \tau) := \ln(c + u) - b_0 - \mu \sum_{n \geq 1} a^n \mu^{n-1} \tau^{n-1} (c + u)^n - \ln(c + u) \sum_{n \geq 1} b^n \mu^n \tau^n (c + u)^n + \sum_{n \geq 1} a^n \mu^n \tau^n (c + u)^n = 0,$$

where F is analytic around $(0, 0, 0)$, $F(0, 0, 0) = 0$, and $F_u(0, 0, 0) = \frac{1}{c} > 0$. Thus, by the Implicit Function Theorem, for $|\mu|$, $|\tau|$ and $|u|$ small there exists a unique real analytic function $u = u(\mu, \tau)$ given by the convergent series $u = \sum_{n+m \geq 1, n, m \geq 0} \tilde{c}_{n,m} \mu^n \tau^m$ such that $F(u(\mu, \tau), \mu, \tau) \equiv 0$. Since $\tau = \frac{1}{\mu} e^{-\frac{1}{c_4\mu}}$, this implies

$$\alpha = e^{-\frac{1}{c_4\mu}} (c + u) = c e^{-\frac{1}{c_4\mu}} + \sum_{n+m \geq 1, n, m \geq 0} \tilde{c}_{n,m} \mu^{n+1} \left(\frac{1}{\mu} e^{-\frac{1}{c_4\mu}} \right)^{m+1}.$$

Case $2n_o + d = 6$. In this case, by (A.4), for $c_6 := 8/(\pi c_v \mu_o^2)$

$$l = 0, \quad g(\alpha) = \frac{1}{\mu_o} - \frac{1}{c_6 \mu_o^2} \left(\alpha + \sum_{n \geq 2} a_n \alpha^n \right), \quad h(\alpha) = \frac{1}{c_6 \mu_o^2} \sum_{n \geq 1} b_n \alpha^{2n},$$

and hence, (3.15) is represented as

$$\frac{1}{\mu} - \frac{1}{\mu_o} = \frac{1}{c_6 \mu_o^2} \left(\alpha + \sum_{n \geq 2} a_n \alpha^n + \ln \alpha \sum_{n \geq 1} b_n \alpha^{2n} \right),$$

or equivalently, by (3.11),

$$\alpha = c_6 (\mu - \mu_o) \sum_{n \geq 0} \left(\frac{\mu - \mu_o}{\mu_o} \right)^n - \sum_{n \geq 2} a_n \alpha^n - \ln \alpha \sum_{n \geq 1} b_n \alpha^{2n}. \tag{3.18}$$

Recalling the definitions of τ and θ in (2.8), setting $\alpha = \tau^2 (c_6 + u)$, we represent (3.18) as

$$F(u, \tau, \theta) := u - c_6 \sum_{n \geq 1} \frac{\tau^{2n}}{\mu_o^n} - \sum_{n \geq 2} a_n \tau^{2n-2} (c_6 + u)^n - \ln(c_6 + u) \sum_{n \geq 1} b_n \tau^{4n} (c_6 + u)^{2n} - \theta \sum_{n \geq 1} b_n \tau^{4n-4} (c_6 + u)^{2n} = 0,$$

where F is real-analytic around $(0, 0, 0)$, $F(0, 0, 0) = 0$ and $F_u(0, 0, 0) = 1$, and F is even in τ . Thus, by the Implicit Function Theorem, for $|u|$, $|\tau|$ and $|\theta|$ small there exists a unique real analytic function $u = u(\tau, \theta)$, even in τ , given by the convergent series $u = \sum_{n+m \geq 1, n, m \geq 0} \tilde{c}_{n,m} \tau^{2n} \theta^m$ such that $F(u(\tau, \theta), \tau, \theta) \equiv 0$. Thus,

$$\alpha = \tau^2 (c_6 + u) = c_6 \sigma + \sum_{n+m \geq 1, n, m \geq 0} \tilde{c}_{n,m} \tau^{2n+2} \theta^m.$$

Case $2n_o + d = 8$. By (A.4), for $c_8 := [8/c_v \mu_o^2]^{-1/2}$,

$$l = 0, \quad g(\alpha) = \frac{1}{\mu_o^2 c_8^2} \sum_{n \geq 2} a_n \alpha^n, \quad h(\alpha) = \frac{1}{\mu_o^2 c_8^2} \left(\alpha^2 + \sum_{n \geq 2} b_n \alpha^{2n} \right),$$

thus, as in the case of $2n_o + d = 6$, (3.15) is represented as

$$c_8^2 (\mu - \mu_o) \sum_{n \geq 0} \left(\frac{\mu - \mu_o}{\mu_o} \right)^n = \alpha^2 \ln \alpha + \ln \alpha \sum_{n \geq 2} b_n \alpha^{2n} + \sum_{n \geq 2} a_n \alpha^n. \quad (3.19)$$

For τ , σ and η given in (2.8) set $\alpha = \tau \sigma (c_8 + u)$ and represent (3.19) as

$$\begin{aligned} 2c_8 u + u^2 = & c_8^2 \sum_{n \geq 1} \frac{\tau^{2n}}{\mu_o^n} + \sum_{n \geq 2} a_n \tau^{n-1} \sigma^{n+1} (c_8 + u)^{n+2} \\ & - \sum_{n \geq 2} b_n (\tau \sigma)^{2n-2} (c_8 + u)^{2n+2} \\ & + \left(\sigma^2 \ln(c_8 + u) - \frac{\eta}{2} \right) \left((c_8 + u)^2 + \sum_{n \geq 2} b_n (\tau \sigma)^{2n-2} (c_8 + u)^{2n+2} \right). \end{aligned}$$

This equation is represented as $F(u, \tau, \sigma, \eta) = 0$, where F is real-analytic in a neighborhood of $(0, 0, 0, 0)$, $F(0, 0, 0, 0) = 0$ and $F_u(0, 0, 0, 0) = 2c_8 > 0$. Hence, for $|u|$, $|\tau|$, $|\sigma|$ and $|\eta|$ small, by the Implicit Function Theorem, there exists a unique real-analytic function $u = u(\tau, \sigma, \eta)$ given by the convergent series $u = \sum_{n+m+k \geq 1, n, m, k \geq 0} \tilde{c}_{n,m,k} \tau^n \sigma^m \mu^k$ such that $F(u(\tau, \sigma, \eta), \tau, \sigma, \eta) \equiv 0$. Thus,

$$\alpha = \tau \sigma (c_8 + u) = c_8 \tau \sigma + \sum_{n+m+k \geq 1, n, m, k \geq 0} \tilde{c}_{n,m,k} \tau^{n+1} \sigma^{m+1} \eta^k.$$

Case $2n_o + d \geq 10$. By (A.4) for $c_{10} := (\mu_o^2 \hat{c}_v)^{-1/2}$,

$$l = 0, \quad g(\alpha) = \frac{1}{\mu_o} + \hat{c}_v \alpha^2 + \sum_{n \geq 2} a_n \alpha^{n+2}, \quad h(\alpha) = \sum_{n \geq 2} b_n \alpha^{2n},$$

and as in the case of $2n_o + d = 6$, (3.15) is represented as

$$\frac{\mu - \mu_o}{\mu_o^2} \sum_{n \geq 0} \left(\frac{\mu - \mu_o}{\mu_o} \right)^n = \widehat{c}_v \alpha^2 + \sum_{n \geq 2} a_n \alpha^{n+2} + \ln \alpha \sum_{n \geq 2} b_n \alpha^{2n}. \tag{3.20}$$

Recalling the definitions of τ and θ in (2.8), we set $\alpha = \tau(c_{10} + u)$. Then (3.20) is represented as

$$\begin{aligned} F(u, \tau, \theta) := & 2c_{10}u + u^2 - c_{10}^2 \sum_{n \geq 1} \frac{\tau^{2n}}{\mu_o^n} + \sum_{n \geq 2} a_n \tau^n (c_{10} + u)^{n+2} \\ & - \theta \sum_{n \geq 2} b_n \tau^{2n-4} (c_8 + u)^{2n} + \ln(c_{10} + u) \sum_{n \geq 2} b_n \tau^{2n-2} (c_8 + u)^{2n} = 0, \end{aligned}$$

where F is analytic at $(0, 0, 0)$, $F(0, 0, 0) = 0$ and $F_u(0, 0, 0) = 2c_{10} > 0$. Thus, by the Implicit Function Theorem, for $|u|$, $|\tau|$ and $|\theta|$ small there exists a unique real-analytic function $u = u(\tau, \theta)$ given by the convergent series $u = \sum_{n+m \geq 1, n, m \geq 0} \tilde{c}_{n,m} \tau^n \theta^m$ such that $F(u(\tau, \theta), \tau, \theta) \equiv 0$. Then

$$\alpha = \mu(c_{10} + u) = c_{10}\mu + \sum_{n+m \geq 1, n, m \geq 0} \tilde{c}_{n,m} \mu^{n+1} \theta^m.$$

Theorem is proved. □

Proof of Theorem 2.5 From (3.3) it follows that the map $p \in \mathbb{T}^d \mapsto |v|^2(\vec{\pi} + p)$ is even. Now the expansions of $e(\mu)$ at $\mu = -\mu_o$ can be proven along the same lines of Theorem 2.4 using Proposition A.2 with $f = |v|^2$. □

Remark 3.3 Let $\widehat{v} : \mathbb{Z}^d \rightarrow \mathbb{C}$ satisfy (H1). Since $\epsilon(\cdot)$ is even,

$$\int_{\mathbb{T}^d} \frac{|v(p)|^2 dp}{\epsilon(p) - z} = \frac{1}{(2\pi)^d} \int_{\mathbb{T}^d} \frac{f(p) dp}{\epsilon(p) - z},$$

where

$$\begin{aligned} f(p) := & \left(\sum_{x \in \mathbb{Z}^d} \widehat{v}_1(x) \cos p \cdot x \right)^2 + \left(\sum_{x \in \mathbb{Z}^d} \widehat{v}_2(x) \cos p \cdot x \right)^2 \\ & + \left(\sum_{x \in \mathbb{Z}^d} \widehat{v}_1(x) \sin p \cdot x \right)^2 + \left(\sum_{x \in \mathbb{Z}^d} \widehat{v}_2(x) \sin p \cdot x \right)^2 \end{aligned}$$

and $\widehat{v} = \widehat{v}_1 + i\widehat{v}_2$ for some $\widehat{v}_1, \widehat{v}_2 : \mathbb{Z}^d \rightarrow \mathbb{R}$. By Lemma 3.1, the unique eigenvalue $e(\mu)$ of H_μ solves

$$1 - \mu \int_{\mathbb{T}^d} \frac{f(p)dp}{e(p) - e(\mu)} = 0.$$

Since both $p \in \mathbb{T}^d \mapsto f(p)$ and $p \in \mathbb{T}^d \mapsto f(\vec{\pi} + p)$ are even analytic functions, we can still apply Propositions A.1 and A.2 to find the expansions of $z \mapsto \int_{\mathbb{T}^d} \frac{f(p)dp}{e(p) - z}$ and thus, repeating the same arguments of the proofs of Theorems 2.4 and 2.5 one can obtain the corresponding expansions of $e(\mu)$.

Remark 3.4 When

$$|\widehat{v}(x)| = O(|x|^{2n_0+d+1}) \quad \text{as } |x| \rightarrow \infty$$

for some $n_0 \geq 1$, in view of Remark A.3, we need to solve equation (3.4) with respect to μ using only that left-hand side is an asymptotic sum (not a convergent series). This still can be done using appropriate modification of the Implicit Function Theorem for differentiable functions. As a result, we obtain only (Taylor-type) asymptotics of $e(\mu)$.

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Appendix A. Asymptotics of some integrals

In this section we study the behaviour of the integral

$$I_f(z) := \int_{\mathbb{T}^d} \frac{f(q)dq}{e(q) - z}, \quad z \in \mathbb{C} \setminus [0, 4d^2], \tag{A.1}$$

as $z \rightarrow 0$ and $z \rightarrow 4d^2$, where $f : \mathbb{T}^d \rightarrow \mathbb{R}$ is a real-analytic even function on \mathbb{T}^d . Further we denote by $W_r(\xi) \subset \mathbb{C}$ the complex disc of radius $r > 0$ centered at $\xi \in \mathbb{C}$.

Proposition A.1 *Let $f : \mathbb{T}^d \rightarrow \mathbb{R}$ be a real-analytic even function such that*

$$f(0) = D^2 f(0) = \dots = D^{2n_0-2} f(0) = 0, \quad D^{2n_0}(0) \neq 0 \tag{A.2}$$

for some $n_o \geq 0$. Then:

- ι_f is continuous at 0 if and only if $2n_o + d \geq 5$;
- ι_f is continuously differentiable at 0 if and only if $2n_o + d \geq 9$, in this case,

$$\iota'_f(0) := \int_{\mathbb{T}^d} \frac{f(q)dq}{(\epsilon(q))^2} = \lim_{z \searrow 0} \int_{\mathbb{T}^d} \frac{f(q)dq}{(\epsilon(q) - z)^2}.$$

Moreover, for any $z \in (-\frac{1}{64}, 0)$:

(a) if d is odd, then

$$\iota_f(z) = \begin{cases} \frac{\pi}{4(-z)^{3/4}} \left(c_f + \sum_{n \geq 1} a_n^d (-z)^{n/4} \right), & 2n_o + d = 1, \\ \frac{\pi}{8(-z)^{1/4}} \left(c_f + \sum_{n \geq 1} a_n^d (-z)^{n/4} \right), & 2n_o + d = 3, \\ \iota_f(0) - \frac{\pi(-z)^{1/4}}{8} \left(c_f + \sum_{n \geq 1} a_n^d (-z)^{n/4} \right), & 2n_o + d = 5, \\ \iota_f(0) - \frac{\pi(-z)^{3/4}}{8} \left(c_f + \sum_{n \geq 1} a_n^d (-z)^{n/4} \right), & 2n_o + d = 7, \\ \iota_f(0) + z \left(\iota'_f(0) + \sum_{n \geq 1} a_n^d (-z)^{n/4} \right), & 2n_o + d \geq 9, \end{cases} \quad (\text{A.3})$$

(b) if d is even, then

$$\iota_f(z) = \begin{cases} \frac{\pi}{8(-z)^{1/2}} \left(c_f + \sum_{n \geq 1} b_n^d (-z)^{n/2} \right) - \frac{1}{16} \ln(-z) \sum_{n \geq 0} c_n^d z^n, & 2n_o + d = 2, \\ -\frac{1}{16} \ln(-z) \left(c_f + \sum_{n \geq 1} c_n^d z^n \right) + \sum_{n \geq 0} b_n^d (-z)^{n/2}, & 2n_o + d = 4, \\ \iota_f(0) - \frac{\pi(-z)^{1/2}}{8} \left(c_f + \sum_{n \geq 1} b_n^d (-z)^{n/2} \right) + z \ln(-z) \sum_{n \geq 0} c_n^d z^n, & 2n_o + d = 6, \\ \iota_f(0) - \frac{z}{16} \ln(-z) \left(c_f + \sum_{n \geq 1} c_n^d z^n \right) + \sum_{n \geq 2} b_n^d (-z)^{n/2}, & 2n_o + d = 8, \\ \iota_f(0) + z \left(\iota'_f(0) + \sum_{n \geq 1} b_n^d (-z)^{n/2} \right) + z^2 \ln(-z) \sum_{n \geq 0} c_n^d z^n, & 2n_o + d \geq 10, \end{cases} \quad (\text{A.4})$$

where $\{a_n^d\}$, $\{b_n^d\}$ and $\{c_n^d\}$ are some real coefficients,

$$c_f := \frac{2^{2n_o+d}}{(2n_o)!} \int_{\mathbb{S}^{d-1}} D^{2n_o} f(0)[w, \dots, w] d\mathcal{H}^{d-1}; \quad (\text{A.5})$$

and all series in (A.3) and (A.4) converge absolutely for $z \in W_{1/64}(0) \subset \mathbb{C}$.

Proof Given $\gamma \in (0, \frac{1}{\sqrt{2}}]$, let $\varphi : B_\gamma(0) \subset \mathbb{R}^d \rightarrow \varphi(B_\gamma(0)) \subset \mathbb{R}^d$ be the smooth diffeomorphism

$$\varphi_i(y) = 2 \arcsin y_i, \quad i = 1, \dots, d.$$

Note that

$$\epsilon(\varphi(y)) = \left(\sum_{i=1}^d (1 - \cos(2 \arcsin(y_i))) \right)^2 = 4 \left(\sum_{i=1}^d y_i^2 \right)^2 = 4y^4, \quad (\text{A.6})$$

therefore,

$$\epsilon(q) \geq 4\gamma^4 \quad \text{for any } q \in \mathbb{T}^d \setminus \varphi(B_\gamma). \quad (\text{A.7})$$

We rewrite $\mathfrak{I}_f(z)$ as

$$\mathfrak{I}_f(z) := \int_{\varphi(B_\gamma(0))} \frac{f(q) dq}{\epsilon(q) - z} + \int_{\mathbb{T}^d \setminus \varphi(B_\gamma(0))} \frac{f(q) dq}{\epsilon(q) - z} := \mathfrak{I}^*(z) + \mathfrak{I}^{**}(z).$$

By virtue of (A.7),

$$\mathfrak{I}^{**}(z) = \int_{\mathbb{T}^d \setminus \varphi(B_\gamma(0))} \frac{f(q)}{\epsilon(q)} \left(1 - \frac{z}{\epsilon(q)} \right)^{-1} dq = \sum_{n \geq 0} z^n \int_{\mathbb{T}^d \setminus \varphi(B_\gamma(0))} \frac{f(q) dq}{(\epsilon(q))^{n+1}}, \quad (\text{A.8})$$

i.e. $\mathfrak{I}^{**}(\cdot)$ is analytic in $W_{2\gamma^4}(0)$. In \mathfrak{I}^* making the change of variables $q = \varphi(y)$ and using (A.6) we get

$$\mathfrak{I}^*(z) = \int_{B_\gamma(0)} \frac{f(\varphi(y)) J(\varphi(y)) dy}{4y^4 - z}, \quad (\text{A.9})$$

where $y^4 := (y^2)^2$ with $y^2 := \sum_{i=1}^d y_i^2$, and

$$J(\varphi(y)) = \prod_{i=1}^d \frac{2}{\sqrt{1 - y_i^2}} \quad (\text{A.10})$$

is the Jacobian of φ . Since f is an even analytic function satisfying (A.2), even each coordinate, from the Taylor series for f it follows that

$$f(p) = \sum_{n \geq n_o} \frac{1}{(2n)!} D^{2n} f(0) \underbrace{[p, \dots, p]}_{2n\text{-times}}, \quad (\text{A.11})$$

and by the analyticity of f in $B_\pi(0) \subset \mathbb{R}^d$, the series converges absolutely in $p \in B_\pi(0)$. By the definition of φ , $\varphi(rw) \subset B_\pi(0)$ for any $r \in (0, \gamma)$ and $w = (w_1, \dots, w_d) \in \mathbb{S}^{d-1}$, where \mathbb{S}^{d-1} is the unit sphere in \mathbb{R}^d . Then letting $p = \varphi(rw)$ and using the Taylor series

$$\varphi_i(rw) = 2rw_i + \frac{r^3 w_i^3}{3} + \sum_{n \geq 3} \tilde{c}_n r^{2n-1} w_i^{2n-1}$$

of $2 \arcsin(\cdot)$, which is absolutely convergent for $|rw_i| < 1$, from (A.11) we obtain

$$f(\varphi(rw)) = \sum_{n \geq n_o} \tilde{C}_n(w) r^{2n}, \tag{A.12}$$

where $\tilde{C}_n : \mathbb{S}^{d-1} \rightarrow \mathbb{R}$ is a homogeneous polynomial of $w \in \mathbb{S}^{d-1}$ of degree $2n$, and

$$\tilde{C}_{n_o}(w) = \frac{2^{2n_o}}{(2n_o)!} D^{2n_o} f(0) \left[\underbrace{w, \dots, w}_{2n_o \text{ - times}} \right]$$

Next consider $J(\varphi(y))$. Inserting the Taylor series of $(1 - t)^{-1/2}$ into (A.10) we obtain

$$J(\varphi(rw)) = 2^d \left(1 + \sum_{n \geq 1} \hat{C}_n(w) r^{2n} \right), \tag{A.13}$$

where $\hat{C}_n : \mathbb{S}^{d-1} \rightarrow \mathbb{R}$ is a homogeneous symmetric polynomial of $w \in \mathbb{S}^{d-1}$ of degree $2n$, and the series converges absolutely.

Now passing to polar coordinates by $y = rw$ in (A.9) and using (A.12) and (A.13) as well as the absolute convergence of the series we get

$$\mathfrak{l}^*(z) = 2^d \int_0^\gamma \frac{r^{d-1}}{4r^4 - z} \left(\sum_{n \geq n_o} \int_{\mathbb{S}^{d-1}} C_n(w) r^{2n} \right) d\mathcal{H}^{d-1} dr = \sum_{n \geq n_o} \hat{c}_n \int_0^\gamma \frac{r^{2n+d-1} dr}{4r^4 - z}, \tag{A.14}$$

where $C_n : \mathbb{S}^{d-1} \rightarrow \mathbb{R}$ is a homogeneous polynomial of $w \in \mathbb{S}^{d-1}$ of degree $2n$ and

$$\hat{c}_n := 2^d \int_{\mathbb{S}^{d-1}} C_n(w) d\mathcal{H}^{d-1}.$$

Note that $\hat{c}_{n_o} = c_f$, where c_f is given by (A.5) and the last series in (A.14) uniformly converges in any compact subset of $\mathbb{C} \setminus [0, 4]$ since \mathfrak{l}^* and

$$z \in \mathbb{C} \setminus [0, 4] \mapsto \mathfrak{j}_{2n+d-1}(z) := \int_0^\gamma \frac{r^{2n+d-1} dr}{4r^4 - z}$$

are analytic functions in $\mathbb{C} \setminus [0, 4]$ and all series in (A.14) converge pointwise¹. Note that for any $m \geq 0$, there exist $c_m \in \mathbb{R}$ and an analytic function f_m in the ball $W_{\gamma^4}(0) \subset \mathbb{C}$ such that for any $z \in (-\gamma^4, 0)$,

$$j_m(z) = z^n j_l^o(z) + c_m + z^\nu f_m((-z)^{1/2}), \tag{A.15}$$

where $n := \lceil \frac{m}{4} \rceil$, $l := m - 4n \in \{0, 1, 2, 3\}$, $\nu = \frac{1}{2}$ for $m = 0, 2$ and $\nu = 1$ for $m = 1, 3$ or $m \geq 4$, and

$$j_l^o(z) := \begin{cases} \frac{\pi}{4} (-z)^{-3/4} & \text{if } l = 0, \\ \frac{\pi}{8} (-z)^{-1/2} & \text{if } l = 1, \\ \frac{\pi}{8} (-z)^{-1/4} & \text{if } l = 2, \\ -\frac{1}{16} \ln(-z) & \text{if } l = 3. \end{cases}$$

Inserting (A.15) into (A.14) we obtain

$$\begin{aligned} \mathfrak{l}^*(z) = & \sum_{n \geq n_o} \widehat{c}_n \left(z^{\lceil \frac{2n+d-1}{4} \rceil} j_{2n+d-1-4\lceil \frac{2n+d-1}{4} \rceil}^o(z) \right. \\ & \left. + c_{2n+d-1} + \widehat{c}_n (-z)^{\nu_n} f_{2n+d-1}((-z)^{1/2}) \right), \end{aligned}$$

where $\{c_{2n+d-1}\} \subset \mathbb{R}$ and $\{f_{2n+d-1}\}$ is a sequence of analytic functions in $W_{\gamma^4}(0)$ and

$$\nu_n := \begin{cases} \frac{1}{2}, & 2n + d = 1, 3, \\ 1, & \text{otherwise.} \end{cases}$$

Since (A.14) converges locally uniformly in $\mathbb{C} \setminus [0, 4]$, $C := \sum_{n \geq n_o} \widehat{c}_n c_{2n+d-1}$ is finite and

$$\sum_{n \geq n_o} \widehat{c}_n (-z)^{\nu_n} f_{2n+d-1}((-z)^{1/2}) = (-z)^\nu g((-z)^{1/2}),$$

where g is analytic in $W_{\gamma^2}(0)$ and $\nu = \frac{1}{2}$ for $2n_o + d = 1, 3$ and $\nu = 1$ otherwise. Hence,

$$\mathfrak{l}^*(z) = C + (-z)^\nu g((-z)^{1/2}) + \sum_{n \geq n_o} \widehat{c}_n z^{\lceil \frac{2n+d-1}{4} \rceil} j_{2n+d-1-4\lceil \frac{2n+d-1}{4} \rceil}^o(z), \tag{A.16}$$

¹ If $\{h_n\}$ is an equi-bounded sequence of analytic functions in a connected open set $\Omega \subset \mathbb{C}$ converging pointwise to a function $h : \Omega \rightarrow \mathbb{C}$, then h is analytic and h_n converges uniformly to h in compact subsets of Ω .

If $0 \leq 2n_o + d - 1 \leq 3$, then by (A.16),

$$\begin{aligned}
 \Gamma^*(z) &= C + (-z)^\nu g((-z)^{1/2}) + \widehat{c}_{n_o} j_{2n_o+d-1}^o(z) \\
 &+ \sum_{n \geq n_o+1} \widehat{c}_n z^{\lfloor \frac{2n+d-1}{4} \rfloor} j_{2n+d-1-4\lfloor \frac{2n+d-1}{4} \rfloor}^o(z).
 \end{aligned}
 \tag{A.17}$$

In view of (A.8) and the definition of j_l^o , from (A.17) we obtain the expansions (A.3) and (A.4) of Γ_f for $2n_o + d \leq 4$. In particular, since $\lfloor \frac{2n+d-1}{4} \rfloor \geq 1$ for $n \geq n_o + 1$, letting $z \rightarrow 0$ in (A.17) we get

$$\lim_{z \rightarrow 0} \Gamma^*(z) = +\infty.
 \tag{A.18}$$

If $2n_o + d - 1 \geq 4$, then $\lfloor \frac{2n+d-1}{4} \rfloor \geq 1$ for any $n \geq n_o$. Therefore, by (A.16), $\Gamma^*(0) := \lim_{z \rightarrow 0} \Gamma^*(z)$ exists and equals to C . In particular, for $2n_o + d - 1 \leq 7$, one has

$$\begin{aligned}
 \Gamma^*(z) &= \Gamma^*(0) - zg((-z)^{1/2}) + \widehat{c}_{n_o} z j_{2n_o+d-1}^o(z) \\
 &+ \sum_{n \geq n_o+1} \widehat{c}_n z^{\lfloor \frac{2n+d-1}{4} \rfloor} j_{2n+d-1-4\lfloor \frac{2n+d-1}{4} \rfloor}^o(z),
 \end{aligned}
 \tag{A.19}$$

from which and (A.8) we deduce the expansions (A.3) and (A.4) of Γ_f for $5 \leq 2n_o + d \leq 8$. In particular, by virtue of (A.18) and analyticity of Γ^{**} at $z = 0$, Γ_f is continuous at 0 if and only if $2n_o + d \geq 5$. Notice also by (A.19)

$$\lim_{z \rightarrow 0} \frac{\Gamma^*(z) - \Gamma^*(0)}{z} = +\infty,
 \tag{A.20}$$

i.e. Γ^* (and hence Γ_f) is not differentiable at $z = 0$.

Finally, if $2n_o + d - 1 \geq 8$, then $\lfloor \frac{2n+d-1}{4} \rfloor \geq 2$ for any $n \geq n_o$. Therefore, by (A.16) there exists

$$\Gamma'^*(0) := \lim_{z \rightarrow 0} \frac{\Gamma^*(z) - \Gamma^*(0)}{z} = -g(0).$$

Now using the Taylor series of g at 0 we get

$$zg((-z)^{1/2}) = \Gamma'^*(0)z + z \sum_{n \geq 1} \frac{g^{(n)}(0)}{n!} (-z)^{n/2}.$$

Inserting this in (A.16), using the definition of j_l^o and the analyticity of Γ^{**} we get the expansions (A.3) and (A.4) of Γ_f for $2n_o + d \geq 9$.

By (A.18) and (A.20), Γ_f is continuously differentiable at 0 if and only if $2n_o + d \geq 9$.

Now the choice $\gamma = \frac{1}{\sqrt{2}}$ completes the proof. □

Proposition A.2 *Let $f : \mathbb{T}^d \rightarrow \mathbb{R}$ be a real-analytic function such that $q \in \mathbb{T}^d \mapsto f(\vec{\pi} + q)$ is even and*

$$f(\vec{\pi}) = D^2 f(\vec{\pi}) = \dots = D^{2n_o-2} f(\vec{\pi}) = 0, \quad D^{2n_o}(\vec{\pi}) \neq 0$$

for some $n_o \in \mathbb{N}_0$. Then:

- ι_f is continuous at $z = 4d^2$ if and only if for $2n_o + d \geq 3$,
- ι_f is continuously differentiable at $z = 4d^2$ if and only if for $2n_o + d \geq 5$, in this case

$$\iota'_f(4d^2) := \int_{\mathbb{T}^d} \frac{f(q)dq}{(\epsilon(q) - 4d^2)^2} = \lim_{z \searrow 4d^2} \int_{\mathbb{T}^d} \frac{f(q)dq}{(\epsilon(q) - z)^2}$$

exists.

Moreover, if $z - 4d^2 \in (0, \frac{1}{16})$, $\iota_f(z)$ is represented as:

(a) if d is odd, then

$$\iota_f(z) = \begin{cases} -\frac{\pi C_f}{\sqrt{z-4d^2}} + \sum_{k \geq 0} a_k^d (z - 4d^2)^{k/2}, & 2n_o + d = 1, \\ \iota_f(4d^2) + \pi C_f \sqrt{z - 4d^2} + \sum_{k \geq 2} a_k^d (z - 4d^2)^{k/2}, & 2n_o + d = 3, \\ \iota_f(4d^2) + \iota'_f(4d^2) (z - 4d^2) + \sum_{k \geq 3} a_k^d (z - 4d^2)^{k/2}, & 2n_o + d \geq 5; \end{cases} \tag{A.21}$$

(b) if d is even, then

$$\iota_f(z) = \begin{cases} C_f \ln \alpha + \ln \alpha \sum_{k \geq 1} b_k^d \alpha^k + \sum_{k \geq 0} c_k^d \alpha^k, & 2n_o + d = 2, \\ \iota_f(4d^2) - C_f \alpha \ln \alpha + \ln \alpha \sum_{k \geq 2} b_k^d \alpha^k + \sum_{k \geq 1} c_k^d \alpha^k, & 2n_o + d = 4, \\ \iota_f(4d^2) + \iota'_f(4d^2) \alpha + \ln \alpha \sum_{k \geq 2} b_k^d \alpha^k + \sum_{k \geq 2} c_k^d \alpha^k, & 2n_o + d \geq 6, \end{cases} \tag{A.22}$$

where $\alpha := z - 4d^2$, $\{a_k^d\}, \{b_k^d\}, \{c_k^d\} \subset \mathbb{R}$ and

$$C_f := \frac{2^{2n_o+d-1}}{(8d)^{n_o+d/2} (2n_o)!} \int_{\mathbb{S}^{d-1}} D^{2n_o} f(\vec{\pi})[w, \dots, w] d\mathcal{H}^{d-1}.$$

Proof Since $4d^2 - \epsilon(\cdot)$ has a unique non-degenerate minimum at $\vec{\pi}$, the asymptotics of $\iota_f(z)$ as $z \searrow 4d^2$ can be done along the lines of, for instance, [22, Lemma 4.1], hence, we skip the proof. □

Remark A.3 When

$$|\widehat{v}(x)| = O(|x|^{2n_0+d+1}) \quad \text{as } |x| \rightarrow \infty$$

for some $n_0 \geq 1$, one has $v \in C^{2n_0}(\mathbb{T}^d)$. In this case the Taylor series of f becomes only asymptotics of order $2n_0 - 1$ and thus, instead of expansions (A.3)-(A.4) and (A.21)-(A.22) of \mathfrak{I}_f one has only asymptotics up to order $2n_0 - 1$.

References

1. Albeverio, S., Lakaev, S., Muminov, Z.: Schrödinger operators on lattices. The Efimov effect and discrete spectrum asymptotics. *Ann. Inst. H. Poincaré Phys. Theor.* **5**, 743–772 (2004)
2. Albeverio, S., Lakaev, S., Makarov, K., Muminov, Z.: The threshold effects for the two-particle Hamiltonians on lattices. *Commun. Math. Phys.* **262**, 91–115 (2006)
3. Andrew, A., Paine, J.: Correction of finite element estimates for Sturm-Liouville eigenvalues. *Numer. Math.* **50**, 205–215 (1986)
4. Basti, G., Teta, A.: Efimov effect for a three-particle system with two identical fermions. *Ann. Henri Poincaré* **18**, 3975–4003 (2017)
5. Ben-Artzi, M., Katriel, G.: Spline functions, the biharmonic operator and approximate eigenvalues. *Numer. Math.* **141**, 839–879 (2019)
6. Boumenir, A.: Sampling for the fourth-order Sturm-Liouville differential operator. *J. Math. Anal. Appl.* **278**, 542–550 (2003)
7. Damanik, D., Hundertmark, D., Killip, R., Simon, B.: Variational estimates for discrete Schrödinger operators with potentials of indefinite sign. *Comm. Math. Phys.* **238**, 545–562 (2003)
8. Damanik, D., Teschl, G.: Bound states of discrete Schrödinger operators with super-critical inverse square potentials. *Proc. Amer. Math. Soc.* **135**, 1123–1127 (2007)
9. Dipierro, S., Karakhanyan, A., Valdinoci, E.: A free boundary problem driven by the biharmonic operator. [arXiv:1808.07696v2](https://arxiv.org/abs/1808.07696v2) [math.AP]
10. Egorova, I., Kopylova, E., Teschl, G.: Dispersion estimates for one-dimensional discrete Schrödinger and wave equations. *J. Spectr. Theory* **5**, 663–696 (2015)
11. Graf, G., Schenker, D.: 2-magnon scattering in the Heisenberg model. *Ann. Inst. Henri Poincaré Phys. Théor.* **67**, 91–107 (1997)
12. Graef, J., Heidarkhani, Sh., Kong, L., Wang, M.: Existence of solutions to a discrete fourth order boundary value problem. *J. Difference Equ. Appl.* **24**, 849–858 (2018)
13. Gridnev, D.: Three resonating fermions in flatland: proof of the super Efimov effect and the exact discrete spectrum asymptotics. *J. Phys. A: Math. Theor.* **47** (2014)
14. Hiroshima, F., Lőrinczi, J.: The spectrum of non-local discrete Schrödinger operators with a δ -potential. *Pacific J. Math. Industry* **6**, 1–6 (2014)
15. Hoffmann, S., Plonka, G., Weickert, J.: Discrete green’s functions for harmonic and biharmonic inpainting with sparse atoms. In: X. Tai *et al* (eds) *Energy Minimization Methods in Computer Vision and Pattern Recognition. EMCCVPR 2015. Lecture Notes in Computer Science*, vol 8932 (2015). Springer, Cham
16. Jaksch, D., *et al.*: Cold bosonic atoms in optical lattices. *Phys. Rev. Lett.* **81**, 3108–3111 (1998)
17. Kholmatov, Sh., Pardabaev, M.: On spectrum of the discrete bilaplacian with zero-range perturbation. *Lobachevskii J. Math.* **42**, 1286–1293 (2021)
18. Klaus, M., Simon, B.: Coupling constant thresholds in nonrelativistic quantum mechanics. I. Short-range two-body case. *Ann. Phys.* **130**, 251–281 (1980)
19. Lakaev, S.: The Efimov effect of a system of three identical quantum lattice particles. *Funkcional. Anal. Prilozhen.* **27**, 15–28 (1993)
20. Lakaev, S., Khalkhuzhaev, A., Lakaev, Sh.: Asymptotic behavior of an eigenvalue of the two-particle discrete Schrödinger operator. *Theoret. Math. Phys.* **171**, 800–811 (2012)
21. Lakaev, S., Kholmatov, Sh.: Asymptotics of eigenvalues of two-particle Schrödinger operators on lattices with zero range interaction. *J. Phys. A: Math. Theor.* **44** (2011)

22. Lakaev, S., Kholmatov, S.: Asymptotics of the eigenvalues of a discrete Schrödinger operator with zero-range potential. *Izv. Math.* **76**, 946–966 (2012)
23. Luef, F., Teschl, G.: On the finiteness of the number of eigenvalues of Jacobi operators below the essential spectrum. *J. Difference Equ. Appl.* **10**, 299–307 (2004)
24. Lewenstein, M., Sanpera, A., Ahufinger, A.: *Ultracold Atoms in Optical Lattices. Simulating Quantum Many-Body Systems.* Oxford University Press, Oxford (2012)
25. Mardanov, R., Zaripov, S.: Solution of Stokes flow problem using biharmonic equation formulation and multiquadratics method. *Lobachevskii J. Math.* **37**, 268–273 (2016)
26. Mattis, D.: The few-body problem on a lattice. *Rev. Mod. Phys.* **58**(2), 361–379 (1986)
27. McKenna, P., Walter, W.: Nonlinear oscillations in a suspension bridge. *Arch. Rational Mech. Anal.* **98**, 167–177 (1987)
28. Mogilner, A.: Hamiltonians in solid-state physics as multiparticle discrete Schrödinger operators: problems and results. *Adv. Sov. Math.* **5**, 139–194 (1991)
29. Naidon, P., Endo, S.: Efimov physics: a review. *Rep. Prog. Phys.* **80** (2017)
30. Rattana, A., Böckmann, C.: Matrix methods for computing eigenvalues of Sturm-Liouville problems of order four. *J. Comput. Appl. Math.* **249**, 144–156 (2013)
31. Sobolev, A.: The Efimov effect. *Discret. Spectr. Asymptotics. Commun. Math. Phys.* **156**, 127–168 (1993)
32. Tamura, H.: The Efimov effect of three-body Schrödinger operator. *J. Funct. Anal.* **95**, 433–459 (1991)
33. Tee, G.: A novel finite-difference approximation to the biharmonic operator. *Comput. J.* **6**, 177–192 (1963)
34. Yafaev, D.: On the theory of the discrete spectrum of the three-particle Schrödinger operator. *Math. USSR-Sb.* **23**, 535–559 (1974)
35. Wall, M.: *Quantum many-body physics of ultracold molecules in optical lattices. Models and simulation models.* Springer Theses, Cham-Heidelberg-New York (2015)
36. Winkler, K., et al.: Repulsively bound atom pairs in an optical lattice. *Nature* **441**, 853–856 (2006)

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