# The Dirichlet problem for the Monge-Ampère equation on Hermitian manifolds with boundary 

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#### Abstract

We study weak quasi-plurisubharmonic solutions to the Dirichlet problem for the complex Monge-Ampère equation on a general Hermitian manifold with non-empty boundary. We prove optimal subsolution theorems: for bounded and Hölder continuous quasiplurisubharmonic functions. The continuity of the solution is proved for measures that are well dominated by capacity, for example measures with $L^{p}, p>1$ densities, or moderate measures in the sense of Dinh-Nguyen-Sibony.


## 1 Introduction

Background. The complex Monge-Ampère equation in a strictly pseudoconvex bounded domain has been extensively studied since 1970's. Bedford and Taylor [1] proved the fundamental result on the existence of (weak) continuous plurisubharmonic solutions to the Dirichlet problem with continuous datum. The classical solutions were obtained later by Caffarelli et al. [8] for the smooth boundary condition and smooth positive right hand side. The comprehensive book [27] contains results, theirs applications and references for both the Dirichlet problem and the Monge-Ampère equation on compact Kähler manifolds without boundary.

On the other hand, it was discovered independently by Semmes [45] and Donaldson [20] that the geodesic equation in the space of Kähler potentials on a compact Kähler manifold is equivalent to the homogeneous Monge-Ampère equation on a compact Kähler manifold with boundary, the product of the Kähler manifold and an annulus in the complex plane. In general the solution is at most $C^{1,1}$-smooth, due to an example of Gamelin and Sibony

[^0][21]. Thus it leads to the study of weak solutions to the equation. The (unique) weak $C^{1, \overline{1}_{-}}$ solution was obtained by Chen [12] (with a complement in [5]) via a sequence of solutions to non degenerate equations-the so called $\varepsilon$-geodesic equations. To solve the non-degenerate Monge-Ampère equation without pseudoconvexity assumption on the boundary one needs to use ideas from Guan's [22]. We refer to the expository paper of Boucksom [7] for more detailed discussion of those developments.

Guan and Li [23] generalized the smooth subsolution theorem in [22] and [5] to the Hermitian setting. This work occurred amid renewed interest in the complex Monge-Ampère equation on Hermitian manifolds which has been studied earlier by Cherrier [13]. It culminated in the resolution of the Monge-Ampère equation on compact Hermitian manifolds by Tosatti and Weinkove [47]. Weak solutions theory have been developed in [18, 35-37, 39, 40], and recent advances for semi-positive Hermitian forms have been made in [24, 25].

Our goal is to study weak solutions to the Dirichlet problem on a smooth compact Hermitian manifold with non-empty boundary. We consider here very general right hand sides, which are positive Radon measures well dominated by capacity considered in [31, 32]. Notice the Käher case is included as a special one, not yet available in the literature.

Results. Let $(\bar{M}, \omega)$ be a $C^{\infty}$ smooth compact Hermitian manifold of dimension $n$ and with non-empty boundary $\partial M$. Thus $\bar{M}=M \cup \partial M$. Let $\mu$ be a positive Radon measure on $M$. Let us denote by $\operatorname{PSH}(M, \omega)$ the set of all $\omega$-plurisubharmonic ( $\omega$-psh for short) functions on $M$. Consider $\varphi \in C^{0}(\partial M, \mathbb{R})$. We study the Dirichlet problem

$$
\left\{\begin{array}{l}
u \in P S H(M, \omega) \cap L^{\infty}(\bar{M}),  \tag{1.1}\\
\left(\omega+d d^{c} u\right)^{n}=\mu, \\
\lim _{z \rightarrow x} u(z)=\varphi(x) \quad \forall x \in \partial M .
\end{array}\right.
$$

Since there is no convexity condition on the boundary, to solve the Dirichlet problem a necessary condition is the existence of a subsolution.
Definition 1.1 (Subsolution) Let $\underline{u} \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that

$$
\lim _{z \rightarrow x} \underline{u}(z)=\varphi(x) \text { for every } x \in \partial M .
$$

(a) It is called a bounded subsolution for the measure $\mu$ if it satisfies:

$$
\left(\omega+d d^{c} \underline{u}\right)^{n} \geq \mu \quad \text { on } M .
$$

(b) If a bounded subsolution is also continuous (resp. Hölder continuous) on $\bar{M}$, then we call it a continuous subsolution (resp. Hölder continuous subsolution) to $\mu$.
Theorem 1.2 Suppose that there exists a bounded subsolution $\underline{u} \in P S H(M, \omega) \cap L^{\infty}(\bar{M})$ for $\mu$. Then, the Dirichlet problem (1.1) has a solution.

This result is a generalization of the bounded subsolution theorem in a bounded strictly pseudoconvex domain due to the first author [29]. Furthermore, it can be considered as the weak solution version of $[5,7,22]$. Note however that our proof does not use those smooth solutions as approximants of the weak solution. For a Hermitian form $\omega$ one needs to improve the stability estimates in $[29,32]$ and apply them to get the statement.

Next we turn to the study of the continuity of solutions if we further assume that the right hand side is also well dominated by capacity, as in the first author's [31].

Recall the Bedford-Taylor capacity defined in our context as follows. For a Borel subset $E \subset M$,

$$
\operatorname{cap}_{\omega}(E):=\sup \left\{\int_{E}\left(\omega+d d^{c} w\right)^{n}: w \in \operatorname{PSH}(M, \omega), 0 \leq w \leq 1\right\} .
$$

Let $h: \mathbb{R}_{+} \rightarrow(0, \infty)$ be an increasing function such that

$$
\begin{equation*}
\int_{1}^{\infty} \frac{1}{x[h(x)]^{\frac{1}{n}}} d x<+\infty . \tag{1.2}
\end{equation*}
$$

In particular, $\lim _{x \rightarrow \infty} h(x)=+\infty$. Such a function $h$ is called admissible. If $h$ is admissible, then so is $A_{2} h\left(A_{1} x\right)$ for every $A_{1}, A_{2}>0$. Define

$$
\begin{equation*}
F_{h}(x)=\frac{x}{h\left(x^{-\frac{1}{n}}\right)} \tag{1.3}
\end{equation*}
$$

Let $\mu$ be a positive Radon measure satisfying for some admissible $h$ :

$$
\begin{equation*}
\mu(E) \leq F_{h}\left(\operatorname{cap}_{\omega}(E)\right), \tag{1.4}
\end{equation*}
$$

for any Borel set $E \subset X$. The set of all measures satisfying this inequality for some admissible $h$ is denoted by $\mathcal{F}(X, h)$. In what follows we often omit to mention that $h$ is admissible.

Corollary 1.3 Consider the subsolution from Theorem 1.2. Assume that $\mu \in \mathcal{F}(M, h)$ for an admissible function $h$. Then, the solution is continuous on $\bar{M}$.

An important class of such measures are the ones with $L^{p}$-density with respect to the Lebesgue measure ( $p>1$ ) (Lemma 5.7). Another class is the Monge-Ampère measures of Hölder continuous quasi-plurisubharmonic functions on $\bar{M}$ (Theorem 5.9). It still remains an open problem, even in a strictly pseudoconvex set in $\mathbb{C}^{n}$, whether a continuous subsolution leads to the continuous solution.

The best regularity of solutions of the Dirichlet problem in our considerations is the Hölder one (see e.g. [26]). In the compact Kähler manifold case it was proved in [34] for $L^{p},(p>1)$ right hand side. Then the Hölder continuous subsolution theorem was proved in [15]. We prove here a significant generalization of [43, 44], which answered positively a question by Zeriahi in the local setting. That is if the subsolution is Hölder continuous, then the solution has this property too.

Theorem 1.4 Assume that the subsolution $\underline{u} \in \operatorname{PSH}(M, \omega) \cap C^{0, \alpha}(\bar{M})$ for some $\alpha>0$. Then, the solution u obtained in Theorem 1.2 is Hölder continuous on $\bar{M}$.

Notice that the Hölder exponent of the solution depends only on the dimension and the Hölder exponent of the subsolution, as in $[15,36]$ for compact complex manifolds without boundary.

The uniqueness of a weak solution is still an open problem in general on compact Hermitian manifolds without boundary (see [37] for a partial result). We are able to prove this property under some extra assumption on either the metric $\omega$ or the manifold.

Corollary 1.5 (Uniqueness of bounded solution) Suppose that $M$ is Stein or $\omega$ is Kähler. Let $u$, $v$ be bounded $\omega$-psh on $M$ such that $\liminf _{z \rightarrow \partial M}(u-v)(z) \geq 0$. Assume that $(\omega+$ $\left.d d^{c} u\right)^{n} \leq\left(\omega+d d^{c} v\right)^{n}$ in $M$. Then, $u \geq v$ on $\bar{M}$. In particular, there is at most one bounded solution to the Dirichlet problem (1.1) in this setting.

Organization. In Sect. 2 we prove various convergence theorems in the Cegrell class of plurisubharmonic functions. These results combined with the Perron method allow to derive the bounded subsolution theorem (Theorem 1.2) in Sect. 3. In Sect. 4 we prove the stability estimates for Hermitian manifolds with boundary for measures that are well-dominated by capacity. This is done by adopting the proofs of stability estimates from the setting of compact

Hermitian manifold without boundary. Section 5 contains local estimates on interior and boundary charts. The key technical result is the bound on volumes of sublevel sets of quasiplurisubharmonic functions in a certain Cegrell class in a boundary chart (Lemma 5.7). In Sect. 6 we prove the Hölder continuous subsolution theorem. One needs to consider a smoothing of the bounded solution obtained in Theorem 1.2, via the geodesic convolution, due to Demailly. Then we obtain the global stability estimate to control the modulus of continuity of the solution. For this we use a rather delicate construction choosing carefully two exhaustive sequences $M_{\varepsilon} \subsetneq M_{\delta}$ of the manifold $M$, and keeping track of the dependence of the modulus of continuity of the solution in collar sets on both parameters $\varepsilon>\delta>0$ (Proposition 6.11). Finally, we give the proof of the uniqueness of solution when either the manifold is Stein or $\omega$ is Kähler.

Notation. The uniform constants $C, C_{0}, C_{1}, \ldots$ may differ from line to line. For simplicity we denote $\|\cdot\|_{\infty}$ to be the supremum norm of functions in the considered domain. We often write $\omega_{v}:=\omega+d d^{c} v$ for a quasi-plurisubharmonic function $v$.

## 2 Convergence in the Cegrell class

Let $\Omega$ be a bounded strictly pseudoconvex domain in $\mathbb{C}^{n}$. Let us recall the Cegrell class introduced in [9]:

$$
\mathcal{E}_{0}(\Omega)=\left\{u \in \operatorname{PSH}(\Omega) \cap L^{\infty}(\Omega): \lim _{z \rightarrow \partial \Omega} u(z)=0, \int_{\Omega}\left(d d^{c} u\right)^{n}<+\infty\right\} .
$$

First we prove a bunch of convergence results to be used in the following sections. In what follows, when the domain of integration $\Omega$ is fixed and no confusion arises we often write

$$
\int g d \lambda:=\int_{\Omega} g d \lambda
$$

for a Borel function $g$ on $\Omega$. For a Borel set $E \subset \Omega$, we denote by $\operatorname{cap}(E):=\operatorname{cap}(E, \Omega)$ its Bedford-Taylor capacity.

The following is implicitly contained in the last part of the proof of [9, Lemma 5.2]. Cegrell dealt with sequences from $\mathcal{E}_{0}(\Omega)$ but the proof works for sequences of negative plurisubharmonic functions as well.

Lemma 2.1 Let $d \lambda$ be a finite positive Radon measure on $\Omega$ which vanishes on pluripolar sets. Suppose that $u_{j} \in \mathcal{E}_{0}(\Omega)$ is a uniformly bounded sequence that converges a.e. with respect to the Lebesgue measure $d V_{2 n}$ to $u \in \mathcal{\mathcal { E } _ { 0 }}(\Omega)$. Then there exists a subsequence $\left\{u_{j_{s}}\right\} \subset\left\{u_{j}\right\}$ such that

$$
\lim _{s \rightarrow+\infty} \int_{\Omega} u_{j_{s}} d \lambda=\int_{\Omega} u d \lambda
$$

Proof Since $d \lambda$ is a finite measure it follows that $\sup _{j} \int_{\Omega}\left|u_{j}\right|^{2} d \lambda<+\infty$. So there exists a subsequence $\left\{u_{j}\right\}$ weakly converging to $v \in L^{2}(d \lambda)$. By the Banach-Saks theorem we can find a subsequence $u_{j_{k}}$ such that

$$
F_{k}=\frac{1}{k}\left(u_{j_{1}}+\cdots+u_{j_{k}}\right) \rightarrow v \text { in } L^{2}(d \lambda)
$$

as $k \rightarrow+\infty$. Extracting a subsequence $\left\{F_{k_{s}}\right\}_{s}$ of $\left\{F_{k}\right\}$ we get $F_{k_{s}} \rightarrow v$ a.e in $d \lambda$, and also that $F_{k_{s}}$ converges a.e to $u$ with respect to the Lebesgue measure. Therefore, $\left(\sup _{s>t} F_{k_{s}}\right)^{*} \searrow u$
everywhere as $t \rightarrow+\infty$. It follows that there is a subsequence which we still denote by $\left\{u_{j}\right\}$ such that

$$
\lim _{j \rightarrow \infty} \int u_{j} d \lambda=\int v d \lambda=\lim _{s \rightarrow \infty} \int F_{k_{s}} d \lambda=\lim _{t \rightarrow \infty} \int \sup _{s>t} F_{k_{s}} d \lambda=\int u d \lambda
$$

where the first identity used the decreasing convergence property; the second one used the a.e- $d \lambda$ convergence, and the last used the fact that $d \lambda$ does not charge pluripolar sets. This completes the proof.

Corollary 2.2 We keep the assumptions of Lemma 2.1. Assume moreover that $d \lambda(E) \leq$ $C_{0} \operatorname{cap}(E)$ for every Borel set $E \subset \Omega$ with a uniform constant $C_{0}$. Then there exists a subsequence, which is still denoted by $\left\{u_{j}\right\}$, such that

$$
\lim _{j \rightarrow \infty} \int\left|u_{j}-u\right| d \lambda=0
$$

Proof By Lemma 2.1 we know that

$$
\lim _{j \rightarrow \infty} \int u_{j} d \lambda=\int u d \lambda, \quad \lim _{j \rightarrow \infty} \int \max \left\{u_{j}, u\right\} d \lambda=\int_{\Omega} u d \lambda
$$

Fix $a>0$. We have $\left\{\left|u-u_{j}\right|>a\right\}=\left\{u-u_{j}>a\right\} \cup\left\{u-u_{j}<-a\right\}$. Therefore

$$
\int_{\left\{u-u_{j}>a\right\}} d \lambda=\int_{\left\{\max \left\{u_{j}, u\right\}-u_{j}>a\right\}} d \lambda \leq \frac{1}{a} \int_{\Omega}\left(\max \left\{u_{j}, u\right\}-u_{j}\right) d \lambda \rightarrow 0 .
$$

Next note that $\max \left\{u, u_{j}\right\} \rightarrow u$ in capacity, by the Hartogs lemma and the quasi-continuity of $u$. It follows that after using the last assumption

$$
\int_{\left\{u-u_{j}<-a\right\}} d \lambda \leq C_{0} \operatorname{cap}\left(\left|\max \left\{u_{j}, u\right\}-u\right|>a\right) \rightarrow 0 .
$$

In conclusion we get that $u_{j} \rightarrow u$ with respect to the measure $d \lambda$ and $\lim \int\left|u_{j}\right| d \lambda=\int|u| d \lambda$. As a byproduct we also get that $u_{j} \rightarrow u$ in $L^{1}(d \lambda)$.

Lemma 2.3 Still under the assumptions of Lemma 2.1 we also suppose that sup ${ }_{j} \int\left(d d^{c} u_{j}\right)^{n} \leq$ $C_{1}$ for some $C_{1}>0$. Let $w_{j} \in \mathcal{E}_{0}(\Omega)$ be a uniformly bounded sequence of plurisubharmonic functions in $\Omega$ satisfying sup ${ }_{j} \int\left(d d^{c} w_{j}\right)^{n} \leq C_{2}$ for some $C_{2}>0$. Assume that $w_{j}$ converges in capacity to $w \in \mathcal{E}_{0}(\Omega)$. Then,

$$
\lim _{j \rightarrow \infty} \int\left|u-u_{j}\right|\left(d d^{c} w_{j}\right)^{n}=0
$$

Proof Note that $\left|u-u_{j}\right|=\left(\max \left\{u, u_{j}\right\}-u_{j}\right)+\left(\max \left\{u, u_{j}\right\}-u\right)$. First, as in the proof of Corollary 2.2 we have $\phi_{j}:=\max \left\{u, u_{j}\right\} \rightarrow u$ in capacity. Fix $\varepsilon>0$. Then, when $j$ is large,

$$
\begin{aligned}
\int_{\Omega}\left(\max \left\{u, u_{j}\right\}-u\right)\left(d d^{c} w_{j}\right)^{n} & \leq \int_{\left\{\left|\phi_{j}-u\right|>\varepsilon\right\}}\left(d d^{c} w_{j}\right)^{n}+\varepsilon \int_{\Omega}\left(d d^{c} w_{j}\right)^{n} \\
& \leq C_{0} \operatorname{cap}\left(\left|\phi_{j}-u\right|>\varepsilon\right)+C_{2} \varepsilon .
\end{aligned}
$$

Therefore, $\lim _{j \rightarrow \infty} \int\left(\phi_{j}-u\right)\left(d d^{c} w_{j}\right)^{n}=0$. Next, we consider for $j>k$,

$$
\int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{j}\right)^{n}-\int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{k}\right)^{n}=\int\left(\phi_{j}-u_{j}\right) d d^{c}\left(w_{j}-w_{k}\right) \wedge T
$$

where $T=T(j, k)=\sum_{s=1}^{n-1}\left(d d^{c} w_{j}\right)^{s} \wedge\left(d d^{c} w_{k}\right)^{n-1-s}$. By integration by parts

$$
\begin{aligned}
\int\left(\phi_{j}-u_{j}\right) d d^{c}\left(w_{j}-w_{k}\right) \wedge T & =\int\left(w_{j}-w_{k}\right) d d^{c}\left(\phi_{j}-u_{j}\right) \wedge T \\
& \leq \int\left|w_{j}-w_{k}\right| d d^{c}\left(\phi_{j}+u_{j}\right) \wedge T
\end{aligned}
$$

Since $\left\|w_{j}\right\|_{\infty},\left\|u_{j}\right\|_{\infty} \leq A$ in $\Omega$ it follows that

$$
\begin{aligned}
\int_{\Omega}\left|w_{j}-w_{k}\right| d d^{c}\left(\phi_{j}+u_{j}\right) \wedge T \leq & A \int_{\left\{\left|w_{j}-w_{k}\right|>\varepsilon\right\}} d d^{c}\left(\phi_{j}+u_{j}\right) \wedge T \\
& +\varepsilon \int_{\left\{\left|w_{j}-w_{k}\right| \leq \varepsilon\right\}} d d^{c}\left(\phi_{j}+u_{j}\right) \wedge T \\
\leq & A^{n+1} \operatorname{cap}\left(\left|w_{j}-w_{k}\right|>\varepsilon\right)+C \varepsilon
\end{aligned}
$$

where the uniform bound for the second integral on the right hand side followed from [10] (see also Corollary 3.4 below). It means that the left hand side is less than $2 C \varepsilon$ for some $k_{0}$ and every $j>k \geq k_{0}$. Thus,

$$
\begin{aligned}
\int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{j}\right)^{n} \leq & \int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{k}\right)^{n} \\
& +\left|\int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{j}\right)^{n}-\int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{k}\right)^{n}\right| \\
\leq & \int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{k}\right)^{n}+2 C \varepsilon \\
\leq & \int\left|u-u_{j}\right|\left(d d^{c} w_{k}\right)^{n}+2 C \varepsilon .
\end{aligned}
$$

Fix $k=k_{0}$ and apply Corollary 2.2 for $d \lambda=\left(d d^{c} w_{k_{0}}\right)^{n}$ to get that for $j \geq k_{1} \geq k_{0}$

$$
\int\left(\phi_{j}-u_{j}\right)\left(d d^{c} w_{j}\right)^{n} \leq(2 C+1) \varepsilon
$$

Since $\varepsilon>0$ was arbitrary, the proof of the lemma is completed.

## 3 Bounded subsolution theorems

Our goal in this section is the proof of Theorem 1.2, but first we shall prove it in the special case of $M \equiv \Omega \subset \mathbb{C}^{n}$ —a strictly pseudoconvex bounded domain. Then the general statement will follow from this and the balayage procedure. Let $\mu$ be a positive Radon measure in $\Omega$ and let $\varphi$ be a continuous function on the boundary $\partial \Omega$. Assume that $\omega$ is a Hermitian form in a neighborhood of $\bar{\Omega}$.

Theorem 3.1 Suppose that $d \mu \leq\left(d d^{c} v\right)^{n}$ for some bounded plurisubharmonic function $v$ in $\Omega$ with $\lim _{z \rightarrow \partial \Omega} v(z)=0$. Then exists a unique $\omega$-plurisubharmonic function $u \in$ $\operatorname{PSH}(\Omega, \omega) \cap L^{\infty}(\Omega)$ solving

$$
\begin{align*}
& \left(\omega+d d^{c} u\right)^{n}=d \mu, \\
& \lim _{\zeta \rightarrow z} u(\zeta)=\varphi(z) \quad \text { for } z \in \partial \Omega . \tag{3.1}
\end{align*}
$$

We begin with showing that it is enough to prove the statement under additional hypothesis on $\varphi, d \mu$ and $\left(d d^{c} v\right)^{n}$. This reduction is done in three steps:

Step 1: One can assume that supp $\mu$ is compact in $\Omega$. Indeed, let $u_{0} \in \operatorname{PSH}(\Omega, \omega) \cap C^{0}(\bar{\Omega})$ be the solution satisfying $u_{0}=\varphi$ on $\partial \Omega$ and $\left(\omega+d d^{c} u_{0}\right)^{n}=0$ in $\Omega$ (it exists thanks to [35, Corollary 4.1]). Let $\eta_{j}$ be a non-decreasing sequence of cut-off functions such that $\eta_{j} \uparrow 1$ on $\Omega$. Then, the sequence of solution $u_{j}$ 's corresponding to $\mu_{j}=\eta_{j} \mu$ is uniformly bounded by $u_{0}+v \leq u_{j} \leq u_{0}$. Hence, by the comparison principle and the convergence theorem ([18], [35, Corollary 3.4]) they will decrease to the solution for $\mu$.

Step 2: We may assume further that the boundary data is in $C^{2}(\partial \Omega)$. Indeed, suppose that the problem is solvable for $\mu$ with compact support and let $\varphi_{k} \in C^{2}(\partial \Omega)$ be a sequence that decreases to $\varphi \in C^{0}(\partial \Omega)$. Then the sequence of solutions $u_{k}$ to

$$
\left(\omega+d d^{c} u_{k}\right)^{n}=\mu, \quad u_{k}=\varphi_{k} \quad \text { on } \partial \Omega
$$

is decreasing and uniformly bounded. The limit $u=\lim u_{k}$ is the required solution for the continuous boundary data.

Step 3: Reduction to the case of $v$ defined in a neighborhood of $\bar{\Omega}$ with $\lim _{z \rightarrow \partial \Omega} v(z)=0$, and the support of $\left(d d^{c} v\right)^{n}$ compact in $\Omega$. We already suppose that $\mu$ has a compact support in $\Omega$. Then we can modify the subsolution $v$ so that $v$ is defined in a neighborhood of $\bar{\Omega}$ and it satisfies for any $z \in \partial \Omega$,

$$
\begin{equation*}
\lim _{\zeta \rightarrow z} v(\zeta)=0 \tag{3.2}
\end{equation*}
$$

By the balayage procedure we may further assume that the support of $\left(d d^{c} v\right)^{n}$ is compact in $\Omega$.

With the above additional assumptions we proceed to define the expected solution. For $v$ as in Step 3, we consider the standard regularizing sequence $v_{j} \downarrow v$. Then, we write $\left(d d^{c} v_{j}\right)^{n}=f_{j} d V_{2 n}$. By [30] there exists $\check{v}_{j} \in \operatorname{PSH}(\Omega) \cap C^{0}(\bar{\Omega})$ such that $\check{v}_{j}=0$ on $\partial \Omega$, and

$$
\begin{equation*}
\left(d d^{c} \check{v}_{j}\right)^{n}=f_{j} d V_{2 n} \quad \text { in } \Omega \tag{3.3}
\end{equation*}
$$

We observe that by the Dini theorem $v_{j}$ converges to $v$ uniformly on compact sets, where the restriction of $v$ is continuous. Consequently, $\check{v}_{j}$ converges to $v$ on those compact sets because by the stability estimate for the Monge-Ampère equation

$$
\sup _{\Omega}\left|\check{v}_{j}-v_{j}\right| \leq \sup _{\partial \Omega}\left|\check{v}_{j}-v_{j}\right|=\sup _{\partial \Omega}\left|v_{j}-v\right| .
$$

Note also that

$$
\begin{equation*}
\int_{\Omega}\left(d d^{c} \check{v}_{j}\right)^{n} \leq C_{1} \tag{3.4}
\end{equation*}
$$

This follows from the compactness of the support of $v=\left(d d^{c} v\right)^{n}$.
Let $0 \leq h \leq 1$ be a continuous function with compact support in $\Omega$. Notice that

$$
\begin{equation*}
h f_{j} d V_{2 n} \rightarrow h\left(d d^{c} v\right)^{n} \quad \text { weakly. } \tag{3.5}
\end{equation*}
$$

We first show the existence of a solution for the measure $h\left(d d^{c} v\right)^{n}$ obtained as the limit of solutions of $h f_{j} d V_{2 n}$ for a certain subsequence of $\left\{f_{j}\right\}$. Applying [35, Corollary 0.4] we solve the Dirichlet problem

$$
\begin{align*}
& u_{j} \in \operatorname{PSH}(\Omega, \omega) \cap C^{0}(\bar{\Omega}), \\
& \left(\omega+d d^{c} u_{j}\right)^{n}=h f_{j} d V_{2 n},  \tag{3.6}\\
& u_{j}(z)=\varphi(z) \text { for } z \in \partial \Omega
\end{align*}
$$

We define

$$
u=\left(\limsup _{j \rightarrow \infty} u_{j}\right)^{*}=\lim _{j \rightarrow \infty}\left(\sup _{\ell \geq j} u_{\ell}\right)^{*} .
$$

By passing to a subsequence we may assume that $u_{j} \rightarrow u$ in $L^{1}(\Omega)$ and $u_{j} \rightarrow u$ a.e. with respect to the Lebesgue measure. Note also that $u \in \operatorname{PSH}(\Omega, \omega) \cap L^{\infty}(\Omega)$ and

$$
\lim _{z \rightarrow \zeta} u(z)=\varphi(\zeta) \text { for all } \zeta \in \partial \Omega
$$

This $u$ will be shown to be the solution we are seeking. To do this we need to prove several lemmas.

First observe that $\varphi$ can be extended to a $C^{2}$ smooth function in a neighborhood of $\bar{\Omega}$. This allows to produce a strictly plurisubharmonic function $g$ in a neighborhood of $\bar{\Omega}$ such that $d d^{c} g \geq \omega$ and $g=-\varphi$ on $\partial \Omega$. Let us use the notation

$$
\widehat{u}=u+g, \quad \widehat{u}_{j}=u_{j}+g .
$$

Using an idea of Cegrell [9, page 210] we first show that
Lemma $3.2 \widehat{u}_{j} \in \mathcal{E}_{0}(\Omega)$.
Proof By [23, Theorem 1.1] there exists $\Phi \in \operatorname{PSH}(\Omega, \omega)$ a $C^{2}$-smooth function on $\bar{\Omega}$ that solves $\left(\omega+d d^{c} \Phi\right)^{n} \equiv 1$ and $\Phi=\varphi$ on $\partial \Omega$. Then,

$$
\left(\omega+d d^{c} u_{j}\right)^{n} \leq\left(d d^{c} \check{v}_{j}\right)^{n} \leq\left(\omega+d d^{c} \check{v}_{j}+d d^{c} \Phi\right)^{n} .
$$

By the comparison principle [35, Corollary 3.4] we have $u_{j} \geq \check{v}_{j}+\Phi$. So, $\widehat{u}_{j}=u_{j}+g \geq$ $\check{v}_{j}+\Phi+g$. Since $\Phi+g \in \operatorname{PSH}(\Omega) \cap C^{2}(\bar{\Omega})$, and equals zero on $\partial \Omega$ it belongs to $\mathcal{E}_{0}(\Omega)$. Thus $\check{v}_{j}+\Phi+g \in \mathcal{E}_{0}(\Omega)$, and so the same is true about $\widehat{u}_{j}$.
Lemma 3.3 There exists a uniform constant $C_{0}$ such that

$$
\int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{n} \leq C_{0} .
$$

Proof Set $\gamma:=d d^{c} g-\omega$. Then $\omega+d d^{c} u_{j}=d d^{c} \widehat{u}_{j}-\gamma$. It follows that

$$
\begin{equation*}
\int_{\Omega}\left(d d^{c} \widehat{u}_{j}-\gamma\right)^{n}=\int_{\Omega}\left(\omega+d d^{c} u_{j}\right)^{n} \leq \int_{\Omega}\left(d d^{c} \check{v}_{j}\right)^{n} \leq C_{1} . \tag{3.7}
\end{equation*}
$$

Using the Newton expansion for the integrand on the left hand side we get that

$$
\begin{equation*}
\int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{n}-\binom{n}{1} \int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{n-1} \wedge \gamma+\cdots+(-1)^{n}\binom{n}{n} \int_{\Omega} \gamma^{n} \leq C_{1} . \tag{3.8}
\end{equation*}
$$

We are going to show that for $k=1, \ldots, n$,

$$
\begin{equation*}
\int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{k} \wedge \gamma^{n-k} \leq C_{2} \tag{3.9}
\end{equation*}
$$

for a uniform constant $C_{2}$. Indeed, since $\gamma$ is a smooth $(1,1)$ form in $\bar{\Omega}$, there is a defining function $\psi \in \mathcal{E}_{0}(\Omega)$ of $\Omega$ such that $\omega+\gamma \leq d d^{c} \psi$ on $\bar{\Omega}$. Using the Cegrell inequality [10, Lemma 5.4] we get for every $k \geq 1$

$$
\begin{align*}
\int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{k} \wedge \gamma^{n-k} & \leq \int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{k} \wedge\left(d d^{c} \psi\right)^{n-k} \\
& \leq\left(\int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{n}\right)^{\frac{k}{n}}\left(\int_{\Omega}\left(d d^{c} \psi\right)^{n}\right)^{\frac{n-k}{n}} . \tag{3.10}
\end{align*}
$$

If we write $m_{j}^{n}=\int\left(d d^{c} \widehat{u}_{j}\right)^{n}$, it follows from (3.8) and (3.10) that

$$
m_{j}^{n}-\text { const. }\binom{n}{k} \sum_{n>k \text { odd }} m_{j}^{k} \leq C_{3} \text { for all } j .
$$

Therefore, the total mass of $\left(d d^{c} \widehat{u}_{j}\right)^{n}$ is bounded by a uniform constant independent of $j$. Consequently,

$$
\int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{k} \wedge \gamma^{n-k} \leq C_{4}
$$

follows by (3.10).
We have also a more general statement.
Corollary 3.4 There exists a uniform constant $C$ such that

$$
\begin{equation*}
\int_{\Omega} T \leq C \tag{3.11}
\end{equation*}
$$

where $T$ are wedge products of $d d^{c} \widehat{u}_{j}$ and $d d^{c} \check{v}_{k}$.
Proof By an application of Cegrell's inequality [10] for every $k \geq 1$,

$$
\begin{equation*}
\int_{\Omega}\left(d d^{c} \check{v}_{j}\right)^{k} \wedge\left(d d^{c} \psi\right)^{n-k} \leq C_{4} \tag{3.12}
\end{equation*}
$$

where $\psi$ is a strictly plurisubharmonic function defining function of $\Omega$ as in the proof of the above lemma. Now using Cegrell's inequality one more time,

$$
\begin{aligned}
\int_{\Omega} T & =\int_{\Omega}\left(d d^{c} \widehat{u}_{j}\right)^{p} \wedge\left(d d^{c} \check{v}_{k}\right)^{q} \wedge\left(d d^{c} \psi\right)^{n-p-q} \\
& \leq\left[I\left(\widehat{u}_{j}\right)\right]^{\frac{p}{n}}\left[I\left(\check{v}_{k}\right)\right]^{\frac{q}{n}}[I(\psi)]^{\frac{n-p-q}{n}},
\end{aligned}
$$

where $I(w)=\int_{\Omega}\left(d d^{c} w\right)^{n}$. Finally, all three factors on the right hand side are bounded by (3.4), Lemma 3.3 and the smoothness of $\psi$ on $\bar{\Omega}$.

Lemma 3.5 Let $\left\{u_{j}\right\} \subset \operatorname{PSH}(\Omega, \omega)$ be a uniformly bounded subsequence of functions satisfying $u_{j}(z)=\varphi(z)$ for $z \in \partial \Omega$ and $u_{j} \rightarrow u$ in $L^{1}(\Omega)$ and $u_{j} \rightarrow u$ a.e. with respect to the Lebesgue measure. Then one can pick a subsequence $\left\{u_{j_{s}}\right\}$ such that for

$$
\begin{equation*}
w_{s}=\max \left\{u_{j_{s}}, u-1 / s\right\} . \tag{3.13}
\end{equation*}
$$

the following equalities hold
(a) $\lim _{s \rightarrow+\infty} \int_{\Omega}\left|u_{j_{s}}-u\right|\left(\omega+d d^{c} u\right)^{n}=0$.
(b) $\lim _{s \rightarrow+\infty} \int_{\Omega}\left|u_{j_{s}}-u\right|\left(\omega+d d^{c} w_{s}\right)^{n}=0$.
(c) $\lim _{s \rightarrow+\infty} \int_{\Omega}\left|u_{j_{s}}-u\right|\left(\omega+d d^{c} u_{j_{s}}\right)^{n}=0$.

Proof Since $u_{j}-u=\widehat{u}_{j}-\widehat{u}$, where $\widehat{u}_{j}, \widehat{u}$ and $\omega_{u}^{n}$ satisfy the assumption of Corollary 2.2, the proof of (a) follows.

By the Hartogs lemma $w_{s}$ converges to $u$ uniformly on any compact set $E$ such that $u_{\mid E}$ is continuous. Combining this and the quasi-continuity of $u$ it follows that $w_{s}$ converges to $u$ in capacity. Therefore, by the convergence theorems in [2, 18],

$$
\left(\omega+d d^{c} u\right)^{n}=\lim _{s \rightarrow+\infty}\left(\omega+d d^{c} w_{s}\right)^{n} .
$$

We observe that $\omega_{u}^{n}$ is a finite Radon measure in $\Omega$. Recall the notation $\widehat{w}_{s}=w_{s}+g$ and, $\widehat{u}=u+g$, where $g$ is a strictly plurisubharmonic defining function for $\Omega$ such that $d d^{c} g \geq \omega$ in a neighborhood of $\bar{\Omega}$. Since $w_{s}=u_{j_{s}}$ in a neighborhood of $\partial \Omega$, it follows from Stokes' theorem and Lemma 3.3 that

$$
\int_{\Omega}\left(\omega+d d^{c} w_{s}\right)^{n} \leq \int_{\Omega}\left(d d^{c} \widehat{w}_{s}\right)^{n}=\int_{\Omega}\left(d d^{c} \widehat{u}_{j_{s}}\right)^{n} \leq C_{0} .
$$

Letting $s \rightarrow+\infty$ we get that $\int \omega_{u}^{n}$ is finite, and thus $\widehat{u} \in \mathcal{E}_{0}(\Omega)$. Summarizing, $(\omega+$ $\left.d d^{c} w_{s}\right)^{n} \leq\left(d d^{c} \widehat{w}_{s}\right)^{n}$ and $\widehat{w}_{s} \rightarrow \widehat{u} \in \mathcal{E}_{0}(\Omega)$ in capacity. Hence (b) follows from Lemma 2.3.

Similarly, with the notation from Lemma 3.2, $\left(\omega+d d^{c} u_{j}\right)^{n} \leq\left(d d^{c} \check{v}_{j}\right)^{n}$ and $\check{v}_{j}$ converges to $v$ in capacity, thus the proof of (c) follows.

Lemma 3.6 Consider $\left\{u_{j}\right\}$ from the previous lemma. Then for a suitably chosen subsequence $\left\{u_{j_{s}}\right\} \subset\left\{u_{j}\right\}$ we have

$$
\left(\omega+d d^{c} u_{j_{s}}\right)^{n} \rightarrow\left(\omega+d d^{c} u\right)^{n} \text { weakly. }
$$

Proof By Lemma 3.5 we can choose a subsequence $\left\{u_{j_{s}}\right\} \subset\left\{u_{j}\right\}$ so that

$$
\int\left|u-u_{j_{s}}\right|\left(\omega+d d^{c} u_{j_{s}}\right)^{n}+\int\left|u-u_{j_{s}}\right|\left(\omega+d d^{c} w_{s}\right)^{n}<1 / s^{2} .
$$

Recall from (3.13) that $w_{s}=\max \left\{u_{j_{s}}, u-1 / s\right\}$. Then

$$
\mathbf{1}_{\left\{u_{j_{s}}>u-1 / s\right\}}\left(\omega+d d^{c} w_{s}\right)^{n}=\mathbf{1}_{\left\{u_{j_{s}}>u-1 / s\right\}}\left(\omega+d d^{c} u_{j_{s}}\right)^{n} .
$$

Therefore, for $\eta \in C_{c}^{\infty}(\Omega)$,

$$
\begin{aligned}
\left|\int \eta \omega_{u}^{n}-\int \eta \omega_{u_{j_{s}}}^{n}\right| & \leq\left|\int \eta \omega_{u}^{n}-\int \eta \omega_{w_{s}}^{n}\right|+\left|\int \eta \omega_{w_{s}}^{n}-\int \eta \omega_{u_{j_{s}}}^{n}\right| \\
& \leq\left|\int \eta \omega_{u}^{n}-\int \eta \omega_{w_{s}}^{n}\right|+\left|\int_{\left\{u_{\left.j_{s} \leq u-1 / s\right\}}\right.} \eta \omega_{w_{s}}^{n}-\eta \omega_{u_{j_{s}}}^{n}\right|
\end{aligned}
$$

The first term on the right hand side goes to zero as $\omega_{w_{s}}^{n} \rightarrow \omega_{u}^{n}$. It remains to estimate the second term. Firstly, by the choice of $\left\{u_{j_{s}}\right\}$ at the begining of this proof,

$$
\begin{aligned}
\left|\int_{\left\{u_{j_{s}} \leq u-1 / s\right\}} \eta \omega_{u_{j_{s}}}^{n}\right| & \leq\|\eta\|_{L^{\infty}} \int_{\left\{u_{\left.j_{s} \leq u-1 / s\right\}}\right.} \omega_{u_{j_{s}}}^{n} \\
& \leq s\|\eta\|_{L^{\infty}} \int\left|u-u_{j_{s}}\right| \omega_{u_{j_{s}}}^{n} \leq \frac{1}{s}\|\eta\|_{L^{\infty}} \rightarrow 0 \text { as } s \rightarrow+\infty
\end{aligned}
$$

Similarly,

$$
\begin{aligned}
\left|\int_{\left\{u_{j_{s}} \leq u-1 / s\right\}} \eta \omega_{w_{s}}^{n}\right| & \leq\|\eta\|_{L^{\infty}}\left|\int_{\left\{u_{\left.j_{s} \leq u-1 / s\right\}}\right.} \eta \omega_{w_{s}}^{n}\right| \\
& \leq s\|\eta\|_{L^{\infty}} \int\left|u-u_{j_{s}}\right| \omega_{w_{s}}^{n} \rightarrow 0 \text { as } s \rightarrow+\infty .
\end{aligned}
$$

The last two estimates complete the proof.
End of proof of Theorem 3.1 Now we come back to the sequence defined in (3.6) and its limit $u$. Let $\left\{u_{j_{s}}\right\}$ be a subsequence of $\left\{u_{j}\right\}$ whose $L^{1}$-limit and pointwise almost everywhere limit is equal $u$.

Let $0 \leq h \leq 1$ be a continuous function with compact support in $\Omega$. Then, applying the last lemma and (3.5), there exists a unique solution $u \in \operatorname{PSH}(\Omega, \omega) \cap L^{\infty}(\bar{\Omega})$ to

$$
\left(\omega+d d^{c} u\right)^{n}=h\left(d d^{c} v\right)^{n}, \quad u=\varphi \quad \text { on } \partial \Omega .
$$

By the Radon-Nikodym theorem $d \mu=h d \nu$ for some Borel function $0 \leq h \leq 1$. Since $h \in L^{1}(d \nu)$ and $C_{c}(\Omega)$ is dense in $L^{1}(d \nu)$, there exists a sequence of continuous functions $0 \leq h_{k} \leq 1$, whose supports are compact in $\Omega$, such that

$$
\lim _{k \rightarrow+\infty} \int\left|h_{k}-h\right| d v=0
$$

In particular, $h_{k} d v \rightarrow h d v$ weakly. Applying the argument above for continuous $h$, we can find $u_{k} \in \operatorname{PSH}(\Omega, \omega) \cap L^{\infty}(\Omega)$ such that $\lim _{z \rightarrow \zeta} u_{k}(z)=\varphi(\zeta)$ for every $\zeta \in \partial \Omega$ and

$$
\left(\omega+d d^{c} u_{k}\right)^{n}=h_{k}\left(d d^{c} v\right)^{n}=h_{k} d v .
$$

Define

$$
u=\left(\lim \sup u_{k}\right)^{*} .
$$

Passing to a subsequence we may assume that $u_{k} \rightarrow u$ in $L^{1}(\Omega)$ and converging $a . e$ to $u$ with respect to the Lebesgue measure. Again by Lemma 3.6 there exists a subsequence $\left\{u_{k_{s}}\right\}$ of $\left\{u_{k}\right\}$ such that

$$
\left(\omega+d d^{c} u_{k_{s}}\right)^{n} \rightarrow\left(\omega+d d^{c} u\right)^{n} \quad \text { weakly. }
$$

Hence,

$$
\left(\omega+d d^{c} u\right)^{n}=\lim _{k_{s} \rightarrow+\infty} h_{k_{s}}\left(d d^{c} v\right)^{n}=h d v .
$$

The proof is completed.
Proof of Theorem 1.2 Let us proceed with the proof of the subsolution theorem on $\bar{M}$ a smooth compact Hermitian manifold with boundary. Consider the following set of functions

$$
\begin{equation*}
\mathcal{B}(\varphi, \mu):=\left\{w \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(M):\left(\omega+d d^{c} w\right)^{n} \geq \mu, w_{\text {Іे }}^{*} \leq \varphi\right\}, \tag{3.14}
\end{equation*}
$$

where $w^{*}(x)=\lim \sup _{M \ni z \rightarrow x} w(z)$ for every $x \in \partial M$. Clearly, $\underline{u} \in \mathcal{B}(\varphi, \mu)$. Let us solve the linear PDE finding $h_{1} \in C^{0}(\bar{M}, \mathbb{R})$ such that

$$
\begin{array}{r}
\left(\omega+d d^{c} h_{1}\right) \wedge \omega^{n-1}=0, \\
h_{1}=\varphi \text { on } \partial M . \tag{3.15}
\end{array}
$$

Since $\left(\omega+d d^{c} w\right) \wedge \omega^{n-1} \geq 0$ for $w \in P S H(M, \omega)$, the maximum principle for the Laplace operator with respect to $\omega$ gives

$$
w \leq h_{1} \quad \text { for all } w \in \mathcal{B}(\varphi, \mu)
$$

Set

$$
\begin{equation*}
u(z)=\sup _{w \in \mathcal{B}(\varphi, \mu)} w(z) \text { for every } z \in M . \tag{3.16}
\end{equation*}
$$

Then, by Choquet's lemma and the fact that $\mathcal{B}(\varphi, \mu)$ satisfies the lattice property, $u=u^{*} \in$ $\mathcal{B}(\varphi, \mu)$. Again by the definition of $u$, we have $\underline{u} \leq u \leq h_{1}$. It follows that

$$
\begin{equation*}
\lim _{z \rightarrow x} u(z)=\varphi(x) \text { for every } x \in \partial M . \tag{3.17}
\end{equation*}
$$

Lemma 3.7 (Lift) Let $v \in \mathcal{B}(\varphi, \mu)$. Let $B \subset \subset M$ be a small coordinate ball (a chart biholomorphic to a ball in $\left.\mathbb{C}^{n}\right)$. Then, there exists $\tilde{v} \in \mathcal{B}(\varphi, \mu)$ such that $v \leq \widetilde{v}$ and $(\omega+$ $\left.d d^{c} \widetilde{v}\right)^{n}=\mu$ on $B$.

Proof We implicitly identify $B^{\prime} \subset \subset B$ with a small ball in $\mathbb{C}^{n}$, with $B^{\prime}$ also fixed. Note that on $B$ we have $\mathbf{1}_{B^{\prime}} d \mu \leq\left(\omega+d d^{c} v\right)^{n}$. Since $\overline{B^{\prime}}$ is compact in $B$, we can easily get a bounded plurisubharmonic subsolution with zero boundary value on $\partial B$ for $\mathbf{1}_{B^{\prime}} d \mu$. Let $\phi_{j} \searrow v$ on $\partial B$ be a uniformly bounded sequence of continuous functions. Theorem 3.1 gives the existence of solutions:

$$
\left\{\begin{array}{l}
v_{j} \in P S H(B, \omega) \cap L^{\infty}(B), \\
\left(\omega+d d^{c} v_{j}\right)^{n}=\mathbf{1}_{B^{\prime}} d \mu, \\
\lim _{z \rightarrow x \in \partial B} v_{j}(z)=\phi_{j}(x), \quad \forall x \in \partial B .
\end{array}\right.
$$

By the comparison principle $v_{j}$ is a decreasing sequence as $j \rightarrow+\infty$ and $v_{j} \geq v$ on $B$. Set $w=\lim _{j} v_{j}$. Then, $w \in \operatorname{PSH}(B, \omega)$ and $w \geq v$. By the convergence theorem $(\omega+$ $\left.d d^{c} w\right)^{n}=\mathbf{1}_{B^{\prime}} d \mu$. Note that $\lim _{j} v_{j}(x)=v(x)$ for every $x \in \partial B$. Hence, $\lim _{\sup _{z \rightarrow x}} w(z) \leq$ $\phi_{j}(x)$. Letting $j \rightarrow+\infty$ we get that

$$
\limsup _{z \rightarrow x} w(z) \leq v(x)
$$

for every $x \in \partial B$
Now we define

$$
\tilde{w}=\left\{\begin{array}{l}
\max \{w, v\} \quad \text { on } B, \\
v \text { on } M \backslash B .
\end{array}\right.
$$

Then, $\widetilde{w} \in P S H(M, \omega) \cap L^{\infty}(M)$ and it satisfies $\widetilde{w}^{*} \leq \varphi$ on $\partial M$. Moreover,

$$
\left(\omega+d d^{c} \widetilde{w}\right)^{n} \geq d \mu \quad \text { on }(M \backslash B) \cup B^{\prime} .
$$

Finally, let $B_{j} \nearrow B$ be a sequence of open balls increasing to $B$, then by the above construction we get a decreasing sequence $\widetilde{w}_{j} \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(M)$ such that $\widetilde{w}_{j}^{*} \leq \varphi$ on $\partial M$ and

$$
\left(\omega+d d^{c} w_{j}\right)^{n} \geq d \mu \quad \text { on }(M \backslash B) \cup B_{j} .
$$

Set $\tilde{v}=\lim w_{j} \geq v$. Then, $\tilde{v}$ is the required lift function.
End of Proof of the bounded subsolution theorem. By (3.17) it remains to show that the function $u$ above satisfies $\left(\omega+d d^{c} u\right)^{n}=\mu$. Let $B \subset \subset M$ be a small coordinate ball. It is enough to check $\left(\omega+d d^{c} u\right)^{n}=\mu$ on $B$. Let $\tilde{u}$ be the lift of $u$ as in Lemma 3.7. It follows that $\widetilde{u} \geq u$ and $\left(\omega+d d^{c} \widetilde{u}\right)^{n}=\mu$ on $B$. However, by the definition $\widetilde{u} \leq u$ on $M$. Thus, $\widetilde{u}=u$ on $B$, in particular on $B$ we have $\left(\omega+d d^{c} \widetilde{u}\right)^{n}=\left(\omega+d d^{c} u\right)^{n}=\mu$.

## 4 Continuity of the solution

First we recall some facts in pluripotential theory for a Hermitian background form. The proofs from [35-37] can be easily adapted to our setting.

Let $\mathbf{B}>0$ be a constant such that on $\bar{M}$ we have

$$
-\mathbf{B} \omega^{2} \leq 2 n d d^{c} \omega \leq \mathbf{B} \omega^{2}, \quad-\mathbf{B} \omega^{3} \leq 4 n^{2} d \omega \wedge d^{c} \omega \leq \mathbf{B} \omega^{3} .
$$

Then we have the following version (see [35, Theorem 0.2]) of the classical comparison principle [2].

Lemma 4.1 Let $u, v \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that $u, v \leq 0$. Assume $\lim \inf _{z \rightarrow \partial M}(u-v)(z) \geq 0$ and $-s_{0}=\sup _{M}(v-u)>0$. Fix $0<\theta<1$ and set $m(\theta)=\inf _{M}[u-(1-\theta) v]$. Then for any $0<s<\min \left\{\frac{\theta^{3}}{16 \mathbf{B}},\left|s_{0}\right|\right\}$,

$$
\begin{equation*}
\int_{\{u<(1-\theta) v+m(\theta)+s\}} \omega_{(1-\theta) v}^{n} \leq\left(1+\frac{s \mathbf{B}}{\theta^{n}} C\right) \int_{\{u<(1-\theta) v+m(\theta)+s\}} \omega_{u}^{n} \tag{4.1}
\end{equation*}
$$

where $C$ is a uniform constant depending only on $n$.
Corollary 4.2 (Domination principle) Let $u, v \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that $\liminf _{z \rightarrow \partial M}(u-v)(z) \geq 0$. Suppose that $u \geq v$ almost everywhere with respect to $\omega_{u}^{n}$. Then $u \geq v$ in $M$.

Proof See [42, Lemma 2.3] and [40, Proposition 2.2].
The next one is a generalization of [35, Theorem 5.3] to Hermitian manifolds with boundary.

Theorem 4.3 Fix $0<\theta<1$. Let $u, v \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that $u \leq 0$ and $-1 \leq v \leq 0$. Assume that $\lim \inf _{z \rightarrow \partial M}(u-v)(z) \geq 0$ and $-s_{0}=\sup _{M}(v-u)>0$. Denote by $m(\theta)=\inf _{M}[u-(1-\theta) v]$, and put

$$
\theta_{0}:=\frac{1}{3} \min \left\{\theta^{n}, \frac{\theta^{3}}{16 \mathbf{B}}, 4(1-\theta) \theta^{n}, 4(1-\theta) \frac{\theta^{3}}{16 \mathbf{B}},\left|s_{0}\right|\right\} .
$$

Suppose that $\omega_{u}^{n} \in \mathcal{F}(M, h)$. Then, for $0<t<\theta_{0}$,

$$
\begin{equation*}
t \leq \kappa\left[\operatorname{cap}_{\omega}(U(\theta, t))\right], \tag{4.2}
\end{equation*}
$$

where $U(\theta, t)=\{u<(1-\theta) v+m(\theta)+t\}$, and the function $\kappa$ is defined on the interval $(0,1)$ by the formula

$$
\begin{equation*}
\kappa\left(s^{-n}\right)=4 C_{n}\left\{\frac{1}{[h(s)]^{\frac{1}{n}}}+\int_{s}^{\infty} \frac{d x}{x[h(x)]^{\frac{1}{n}}}\right\}, \tag{4.3}
\end{equation*}
$$

with a dimensional constant $C_{n}$.
We use it to to generalize the stability estimate for manifolds with boundary. Let $\hbar(s)$ be the inverse function of $\kappa(s)$ and

$$
\begin{equation*}
\Gamma(s) \text { the inverse function of } s^{n(n+2)+1} \hbar\left(s^{n+2}\right) \text {. } \tag{4.4}
\end{equation*}
$$

Notice that $\Gamma(s) \rightarrow 0$ as $s \rightarrow 0^{+}$.
Proposition 4.4 (Stability of solutions) Let $u, v \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that $u, v \leq 0$. Let $\mu \in \mathcal{F}(M, h)$. Assume that $\liminf _{z \rightarrow \partial M}(u-v)(z) \geq 0$ and

$$
\left(\omega+d d^{c} u\right)^{n}=\mu .
$$

Then, there exists a constant $C>0$ depending only on $\omega$ and $\|v\|_{\infty}$ such that

$$
\sup _{M}(v-u) \leq C \Gamma\left(\left\|(v-u)_{+}\right\|_{L^{1}(d \mu)}\right) .
$$

Proof It is identical to the one in [36, Proposition 2.4].
Now we are going to prove Corollary 1.3. So we assume the existence of a subsolution for the measure $\mu$ and the boundary data $\varphi$ as in Theorem 1.2 and on top of that suppose $\mu \in \mathcal{F}(M, h)$ for an admissible function $h$. Then we claim that the solution is continuous.

We argue by contradiction. Suppose $u$ were not continuous. Then the discontinuity of $u$ would occur at an interior point of $M$. Hence

$$
d=\sup _{\bar{M}}\left(u-u_{*}\right)>0,
$$

where $u_{*}(z)=\lim _{\epsilon \rightarrow 0} \inf _{w \in B(z, \epsilon)} u(w)$ is the lower-semicontinuous regularization of $u$. Consider the closed nonempty set

$$
F=\left\{u-u_{*}=d\right\} \subset \subset M .
$$

The inclusion follows from the boundary condition. One can extend the boundary data $\varphi$ to a continuous function on $\bar{M}$. Then one can use the boundary condition again and the compactness of $\partial M$ to choose $M^{\prime} \subset \subset M^{\prime \prime} \subset \subset M$ with $\overline{M^{\prime}}$ so close to $\bar{M}$ (in the sense of the distance induced by the metric $\omega$ ) that

$$
|u-\varphi|<d / 4 \text { on } M \backslash M^{\prime} .
$$

By the approximation property of quasi-plurisubharmonic functions [6], one can find a sequence

$$
\begin{equation*}
\operatorname{PSH}\left(M^{\prime \prime}, \omega\right) \cap C^{\infty}\left(M^{\prime \prime}\right) \ni u_{j} \searrow u \quad \text { in } M^{\prime \prime} . \tag{4.5}
\end{equation*}
$$

By the Hartogs lemma there exists $j_{0}>0$ such that for $j>j_{0}$

$$
u_{j} \leq \varphi+d / 2 \text { on } M^{\prime \prime} \backslash M^{\prime} .
$$

Then the sets $\left\{u<u_{j}-d / 4\right\}$ are nonempty and relatively compact in $M^{\prime}$ for $j>j_{0}$. Moreover, by subtracting a uniform constant we may assume that

$$
-C_{0} \leq u, u_{j} \leq 0 \quad \text { on } M .
$$

Note that for a Borel set $E \subset M^{\prime}, \operatorname{cap}_{\omega}(E, M) \leq \operatorname{cap}_{\omega}\left(E, M^{\prime}\right)$. It follows that $\omega_{u}^{n} \in$ $\mathcal{F}\left(M^{\prime}, h\right)$. Now we apply Theorem 4.3 for $u$ and $v:=u_{j}$ on $M^{\prime}$ to get a contradiction. In fact, $-m_{j}=\sup _{M^{\prime}}\left(u_{j}-u\right) \geq d / 4$ for $j>j_{0}$. Let $0<\varepsilon<d / 12\left(1+A_{0}\right)$ and take $\theta_{0}$ from Theorem 4.3. For $0<s<\theta_{0}$ we have

$$
\begin{aligned}
U_{j}(\varepsilon, s) & :=\left\{u<(1-\varepsilon) u_{j}+\inf _{M^{\prime}}\left(u-(1-\varepsilon) u_{j}\right)\right\} \\
& \subset\left\{u<u_{j}+m_{j}+\varepsilon\left\|u_{j}\right\|_{\infty}+s\right\} \\
& \subset\left\{u<u_{j}\right\},
\end{aligned}
$$

where for the last inclusion we used the fact $\left\|u_{j}\right\| \leq A_{0}$ and $0<s \leq-m_{j} / 3$. Fix $0<s<\theta_{0}$. Then that theorem gives

$$
s \leq \kappa\left[\operatorname{cap}_{\omega}(U(\varepsilon, s))\right] \leq \kappa\left[\operatorname{cap}_{\omega}\left(u<u_{j}\right)\right] .
$$

This leads to a contradiction since $\operatorname{cap}_{\omega}\left(u<u_{j}\right) \rightarrow 0$ as $j \rightarrow+\infty$.
Remark 4.5 We only need to assume that for any $M^{\prime} \subset \subset M$, there exists an admissible function $h^{\prime}$ (it may depend on $M^{\prime}$ ) such that $\mu \in \mathcal{F}\left(M^{\prime}, h^{\prime}\right)$. Then, the solution $u$ is continuous on $\bar{M}$. Therefore, if $\mu$ is locally dominated by the Monge-Ampère measure of a Hölder
continuous plurisubharmonic functions, then we have for a fixed compact set $M^{\prime} \subset \subset M$, $\mu \in \mathcal{F}\left(M^{\prime}, h\right)$ for some admissible function $h$. This is a simple consequence of a result due to Dinh-Nguyen-Sibony [19] (Remark 5.10). More generally, if the modulus of continuity of $\underline{u}$ satisfies a Dini-type condition then we also get the same conclusion (see [38]).

We have an immediate consequence of this remark.
Corollary 4.6 The bounded solution obtained in Theorem 1.2 is continuous if the subsolution $\underline{u}$ is Hölder continuous.

## 5 Estimates on boundary charts

For the proof of the subsolution theorem in the Hölder continuous class we need rather delicate estimates close to the boundary of $\bar{M}$. In this section we prove estimates in local boundary charts of $\bar{M}$. The estimates on interior charts are easier and they are only mentioned as remarks after the corresponding estimates for boundary charts.

Let $q \in \partial M$ be a point on the boundary. Let $\Omega$ be a boundary chart centered at $q$. Fix $\rho$ a defining function for $\partial M$ in $\Omega$. Identifying this chart with a subset of $\mathbb{C}^{n}$, we can take as $\Omega$ the coordinate "half-ball" of radius $2 R>0$ centered at $q$. More precisely let $B_{r} \subset \mathbb{C}^{n}$ for $r>0$ denote the ball of radius $r$ centered at 0 . Let $\rho: B_{2 R} \rightarrow \mathbb{R}$ be a smooth function, with $\rho(0)=0$ and $d \rho \neq 0$ along $\{\rho=0\}$. Then,

$$
\Omega=\left\{z \in B_{2 R}: \rho(z)<0\right\}, \quad \partial M \cap \Omega=\left\{z: B_{2 R}: \rho(z)=0\right\} .
$$

The Hermitian form $\omega$ is smooth up to the boundary $\partial M$, so we can extend it to a smooth Hermitian form on $\overline{B_{2 R}}$. Multiplying $\rho$ by a small positive constant we assume also that

$$
\omega+d d^{c} \rho>0 \quad \text { and } \quad-1 \leq \rho \leq 0 \quad \text { on } \overline{B_{2 R}} .
$$

Moreover, the estimates are local, so we can fix a Kähler form $\Theta=d d^{c} g_{0}$, where $g_{0}=$ $C\left(|z|^{2}-(3 R)^{2}\right)$ and $C>0$ is large enough so that

$$
\Theta \geq \omega \text { on } \overline{B_{2 R}} .
$$

Suppose that

$$
\begin{equation*}
\psi \in P S H(\Omega) \cap C^{0, \alpha}(\bar{\Omega}) \text { and } \psi \leq 0, \tag{5.1}
\end{equation*}
$$

where $0<\alpha \leq 1$. Our goal is to get the stability estimates for $\omega$-psh functions with respect to the Monge-Ampère measure of $\psi$. This will provide the stability estimates on boundary charts under the assumption that a Hölder continuous subsolution exists.

Let us recall the following class of domains defined in [11, Definition 3.2]: A domain $D \subset \mathbb{C}^{n}$ is called quasi-hyperconvex if $D$ admits a continuous negative $\omega$-psh exhaustion function $\rho: D \rightarrow[-1,0)$.

We now give several estimates on the quasi-hyperconvex domain

$$
\Omega_{R}=\Omega \cap B_{R} .
$$

The first observation is that the continuous exhaustion function for $\Omega_{R}$,

$$
\begin{equation*}
\rho_{R}=\max \left\{\rho,|z|^{2}-R^{2}\right\} \in \operatorname{PSH}\left(\Omega_{R}, \omega\right) \tag{5.2}
\end{equation*}
$$

has finite Monge-Ampère mass on $\Omega_{R}$, i.e.,

$$
\int_{\Omega_{R}}\left(\Theta+d d^{c} \rho_{R}\right)^{n} \leq C_{0} .
$$

In fact, it follows from smoothness of $\rho$ and $|z|^{2}$ on $B_{2 R}$ and the Chern-Levin-Nirenberg inequality (see e.g. [33, page 8]) that $\rho_{R}$ has finite Monge-Ampère mass. More generally, for $1 \leq k \leq n$,

$$
\begin{equation*}
\int_{\Omega_{R}}\left(\Theta+d d^{c} \rho_{R}\right)^{k} \wedge \Theta^{n-k} \leq C_{0}\left\|g_{0}+\rho_{R}\right\|_{L^{\infty}\left(B_{2 R}\right)}^{k} \tag{5.3}
\end{equation*}
$$

Let us consider a class corresponding to the Cegrell class $\mathcal{E}_{0}$ in Sect. 2.

$$
\begin{equation*}
\mathcal{P}_{0}(\Theta)=\left\{v \in \operatorname{PSH}\left(\Omega_{R}, \Theta\right): v \leq 0, \lim _{z \rightarrow \partial \Omega_{R}} v(z)=0, \int_{\Omega_{R}} \Theta_{v}^{n}<+\infty\right\}, \tag{5.4}
\end{equation*}
$$

where $\operatorname{PSH}\left(\Omega_{R}, \Theta\right)$ denotes the set of all $\Theta$-psh functions in $\Omega_{R}$. We first adopt an inequality due to Błocki [4] to our setting.

Lemma 5.1 Let $v \in \mathcal{P}_{0}(\Theta)$ and $1 \leq k \leq n$ an integer. Let $\phi_{i} \leq 0, i=1, \ldots, k$, be plurisubharmonic functions in $\Omega_{R}$. Then

$$
\int_{\Omega_{R}}(-v)^{k} d d^{c} \phi_{1} \wedge \cdots \wedge d d^{c} \phi_{k} \wedge \Theta^{n-k} \leq k!\left\|\phi_{1}\right\|_{\infty} \cdots\left\|\phi_{k}\right\|_{\infty} \int_{\Omega_{R}} \Theta_{v}^{k} \wedge \Theta^{n-k}
$$

Proof We give the proof in a particular case (which is used below) $\phi_{1}=\cdots=\phi_{k}$. The general case is similar. Denote by $w_{\varepsilon}$ the function $\max \{v,-\varepsilon\}$ for $\varepsilon>0$. Then, $w_{\epsilon} \nearrow 0$ as $\varepsilon \searrow 0$. By the Lebesgue convergence theorem

$$
\lim _{\varepsilon \rightarrow 0} \int_{\Omega_{R}}\left(w_{\varepsilon}-v\right)^{k}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k}=\int_{\Omega_{R}}(-v)^{k}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k}
$$

Fix small $\varepsilon>0$ and write $w$ instead of $w_{\varepsilon}$. Then $v \leq w \in \operatorname{PSH}\left(\Omega_{R}, \Theta\right)$ and $w=v$ in a neighborhood of $\partial \Omega_{R}$. We first show that
$\int_{\Omega_{R}}(w-v)^{k}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k} \leq k\|\phi\|_{\infty} \int_{\Omega_{R}}(w-v)^{k-1}\left(\Theta+d d^{c} v\right) \wedge\left(d d^{c} \phi\right)^{k-1} \wedge \Theta^{n-k}$.
We have

$$
\begin{aligned}
-d d^{c}(w-v)^{k} & =-k(k-1) d(w-v) \wedge d^{c}(w-v)-k(w-v)^{k-1} d d^{c}(w-v) \\
& =-k(k-1) d(w-v) \wedge d^{c}(w-v)+k(w-v)^{k-1}\left(\Theta_{v}-\Theta_{w}\right) \\
& \leq k(w-v)^{k-1} \Theta_{v}
\end{aligned}
$$

By integration by parts

$$
\int_{\Omega_{R}}(w-v)^{k}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k}=\int_{\Omega_{R}} \phi d d^{c}(w-v)^{k} \wedge\left(d d^{c} \phi\right)^{k-1} \wedge \Theta^{n-k}
$$

Since $\phi$ is negative in $\Omega_{R}$, it follows from the previous inequality that

$$
\begin{aligned}
\int_{\Omega_{R}}(-\phi)\left[-d d^{c}(w-v)^{k}\right] & \wedge\left(d d^{c} \phi\right)^{k-1} \wedge \Theta^{n-k} \\
& \leq k \int_{\Omega_{R}}(-\phi)(w-v)^{k-1} \Theta_{v} \wedge\left(d d^{c} \phi\right)^{k-1} \wedge \Theta^{n-k} \\
& \leq k\|\phi\|_{\infty} \int_{\Omega_{R}}(w-v)^{k-1}\left(d d^{c} \phi\right)^{k-1} \wedge \Theta_{v} \wedge \Theta^{n-k}
\end{aligned}
$$

So the claim above follows.

Repeating this process $k$-times we obtain

$$
\int_{\Omega_{R}}(w-v)^{k}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k} \leq k!\|\phi\|_{\infty}^{k} \int_{\Omega_{R}} \Theta_{v}^{k} \wedge \Theta^{n-k} .
$$

The lemma follows by letting $\varepsilon \rightarrow 0$.
Thanks to this lemma and the quasi-hyperconvexity of $\Omega_{R}$ (every bounded smooth domain in $\mathbb{C}^{n}$ is quasi-hyperconvex) we are able to estimate the Monge-Ampère mass of a bounded plurisubharmonic function on

$$
\begin{equation*}
D_{\varepsilon}(R)=\left\{z \in \Omega_{R}: \rho_{R}(z)<-\varepsilon\right\}, \quad \varepsilon>0 . \tag{5.5}
\end{equation*}
$$

Corollary 5.2 Let $1 \leq k \leq n$ be an integer. Let $\phi \in \operatorname{PSH}\left(\Omega_{R}\right) \cap L^{\infty}\left(\Omega_{R}\right)$ be such that $\phi \leq 0$. Then,

$$
\int_{D_{\varepsilon}(R)}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k} \leq \frac{C\|\phi\|_{\infty}^{k}}{\varepsilon^{k}},
$$

where $C$ is a uniform constant independent of $\varepsilon$.
Proof $\operatorname{On} D_{\varepsilon}(R)$ we have $\max \left\{\rho_{R},-\varepsilon / 2\right\}-\rho_{R} \geq \varepsilon / 2$. It follows from Lemma 5.1 that

$$
\begin{aligned}
\int_{D_{\varepsilon}(R)}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k} & \left.\leq\left(\frac{2}{\varepsilon}\right)^{k} \int_{\Omega_{R}}\left(\max \left\{\rho_{R},-\varepsilon / 2\right\}\right)-\rho_{R}\right)^{k}\left(d d^{c} \phi\right)^{k} \wedge \Theta^{n-k} \\
& \leq\left(\frac{2}{\varepsilon}\right)^{k}\|\phi\|_{\infty}^{k} \int_{\Omega_{R}} \Theta_{\rho_{R}}^{k} \wedge \Theta^{n-k}
\end{aligned}
$$

The last integral is bounded thanks to (5.3).
Remark 5.3 We observe that for a bounded plurisubharmonic function

$$
2 d w \wedge d^{c} w=d d^{c} w^{2}-2 w d d^{c} w .
$$

Then applying the above corollary to both terms on the right hand side one obtains

$$
\begin{equation*}
\int_{D_{\varepsilon}(R)} d w \wedge d^{c} w \wedge\left(d d^{c} \phi\right)^{k-1} \wedge \Theta^{n-k} \leq \frac{C\|w\|_{\infty}^{2}\|\phi\|_{\infty}^{k-1}}{\varepsilon^{k}} \tag{5.6}
\end{equation*}
$$

The Hölder continuity of a function can be detected by checking the speed of convergence of regularizing sequences which we now define. Let us use the following the standard smoothing kernel $\tilde{\chi}(z)=\chi\left(|z|^{2}\right)$ in $\mathbb{C}^{n}$, where

$$
\chi(t)= \begin{cases}\frac{c_{n}}{(1-t)^{2}} \exp \left(\frac{1}{t-1}\right) & \text { if } 0 \leq t \leq 1,  \tag{5.7}\\ 0 & \text { if } t>1\end{cases}
$$

with the constant $c_{n}$ chosen so that

$$
\begin{equation*}
\int_{\mathbb{C}^{n}} \chi\left(\|z\|^{2}\right) d V_{2 n}(z)=1, \tag{5.8}
\end{equation*}
$$

and $d V_{2 n}$ denoting the Lebesgue measure in $\mathbb{C}^{n}$.

Let $u \in \operatorname{PSH}(\Omega) \cap L^{\infty}(\Omega)$ and $\delta>0$. For $z \in \Omega_{\delta}:=\{z \in \Omega: \operatorname{dist}(z, \partial \Omega)>\delta\}$ define

$$
\begin{align*}
& u * \chi_{\delta}(z)=\int_{|x|<1} u(z+\delta x) \widetilde{\chi}(x) d V_{2 n}(x),  \tag{5.9}\\
& \check{u}_{\delta}(z)=\frac{1}{v_{2 n} \delta^{n}} \int_{B_{\delta}(z)} u(x) d V_{2 n}(x), \tag{5.10}
\end{align*}
$$

where $B_{\delta}(z)=\left\{x \in \mathbb{C}^{n}:|z-x|<\delta\right\}$. Let us denote by $\sigma$ the surface measure on a sphere $S(z, r)$ and by $\bar{\sigma}_{2 n-1}$ the area of the unit sphere. Consider the averages

$$
\mu_{S}(u ; z, \delta)=\frac{1}{\bar{\sigma}_{2 n-1} \delta^{2 n-1}} \int_{S(z, \delta)} u(x) d \sigma_{2 n-1},
$$

which is are increasing in $\delta$. Therefore, for any $z \in \Omega_{\delta}$,

$$
\begin{aligned}
u * \chi_{\delta}(z)-u(z) & =\bar{\sigma}_{2 n-1} \int_{0}^{1}\left[\mu_{S}(u ; z, \delta t)-u(z)\right] \chi\left(t^{2}\right) t^{2 n-1} d t \\
& \leq\|\chi\|_{L^{\infty}} \int_{0}^{1}\left[\mu_{S}(u ; z, \delta t)-u(z)\right] t^{2 n-1} d t \\
& =\|\chi\|_{L^{\infty}} \frac{\left(\check{u}_{\delta}-u\right)(z)}{2 n} .
\end{aligned}
$$

On the other hand, since $\chi \geq 1 / C$ on [0, 1/2],

$$
\begin{aligned}
\int_{0}^{1}\left[\mu_{S}(u ; z, \delta t)-u(z)\right] \chi\left(t^{2}\right) t^{2 n-1} d t & \geq \frac{1}{C} \int_{0}^{1 / 2}\left[\mu_{S}(u ; z, \delta t)-u(z)\right] t^{2 n-1} d t \\
& =\frac{1}{2^{n} C} \int_{0}^{1}\left[\mu_{S}\left(u ; z, \frac{\delta s}{2}\right)-u(z)\right] s^{2 n-1} d s \\
& =\frac{1}{2^{n} C} \frac{\left(\check{u}_{\delta / 2}-u\right)(z)}{2 n} .
\end{aligned}
$$

We conclude that there exist a uniform constant $C>0$ such that in $\Omega_{\delta}$

$$
\begin{align*}
u * \chi_{\delta}-u & \leq C\left(\check{u}_{\delta}-u\right),  \tag{5.11}\\
\check{u}_{\delta / 2}-u & \leq C\left(u * \chi_{\delta}-u\right) . \tag{5.12}
\end{align*}
$$

Lemma 5.4 Let $u \in \operatorname{PSH}(\Omega) \cap L^{\infty}(\Omega)$. For $0<\delta<\delta_{0} \leq R / 4$,

$$
\int_{B_{R / 2} \cap \Omega_{2 \delta}}\left(\check{u}_{\delta}(z)-u(z)\right) d V_{2 n} \leq C \delta^{2} \int_{B_{3 R / 4} \cap \Omega_{\delta}} \Delta u(z) d V_{2 n} .
$$

Consequently,

$$
\int_{B_{R / 2} \cap \Omega_{\delta}}\left(\check{u}_{\delta}(z)-u(z)\right) d V_{2 n} \leq C \delta .
$$

Proof The first inequality follows from the classical Jensen formula (see e.g. [26, Lemma 4.3]. For the second one we observe

$$
\begin{aligned}
\int_{B_{R / 2} \cap \Omega_{\delta}}\left(\check{u}_{\delta}(z)-u(z)\right) d V_{2 n} \leq & \int_{B_{R / 2} \cap \Omega_{2 \delta}}\left(\check{u}_{\delta}(z)-u(z)\right) d V_{2 n} \\
& +\int_{\Omega_{\delta} \backslash \Omega_{2 \delta}}\left(\check{u}_{\delta}(z)-u(z)\right) d V_{2 n} \\
\leq & C \delta^{2} \int_{B_{3 R / 4} \cap \Omega_{\delta}} \Delta u(z) d V_{2 n}+C \delta .
\end{aligned}
$$

Since $\partial \Omega$ is a smooth manifold defined by $\rho$ in $B_{2 R}$, there exists a uniform constant $c_{0}>0$ such that $|\rho(z)| \geq c_{0} \operatorname{dist}(z, \partial \Omega)$ for every $z \in B_{R} \cap \Omega$. Hence, there exists a uniform constant $c_{1}=c_{1}\left(c_{0}, R\right)>0$ such that $B_{3 R / 4} \cap \Omega_{\delta} \subset\left\{\rho_{R}<-c_{1} \delta\right\}$ for every $0<\delta<\delta_{0}$. Now, we apply Corollary 5.2 with $\varepsilon=c_{1} \delta$ and $k=1$ for the first term on the right and the proof is completed.

Remark 5.5 In the interior chart the corresponding (stronger) inequality is given by [26, Lemma 4.3]. Let $u \in P S H \cap L^{\infty}\left(B_{2 R}\right)$. Then, for $0<\delta<R / 2$,

$$
\begin{equation*}
\int_{B_{R}}\left(\check{u}_{\delta}-u\right) d V_{2 n} \leq c_{n}\left(\int_{B_{3 R / 2}} \Delta u d V_{2 n}\right) \delta^{2} . \tag{5.13}
\end{equation*}
$$

We also need various capacities in the estimates. Consider the local $\Theta$-capacity:

$$
\operatorname{cap}_{\Theta}\left(E, \Omega_{R}\right)=\sup \left\{\int_{E}\left(\Theta+d d^{c} v\right)^{n}: v \in \operatorname{PSH}\left(\Omega_{R}, \Theta\right), \quad-1 \leq v \leq 0\right\}
$$

This capacity is equivalent to the Bedford-Taylor capacity.
Lemma 5.6 There exists a constant $A_{1}>0$ such that for every Borel set $E \subset \Omega$,

$$
\frac{1}{A_{1}} \operatorname{cap}\left(E, \Omega_{R}\right) \leq \operatorname{cap}_{\Theta}\left(E, \Omega_{R}\right) \leq A_{1} \operatorname{cap}\left(E, \Omega_{R}\right)
$$

Proof The first inequality is straightforward while the second inequality follows from the fact that $\Theta=d d^{c} g_{0}=d d^{c}\left[C\left(|z|^{2}-(3 R)^{2}\right]\right.$ on $\bar{\Omega}$ for some $C>0$ large enough depending only on $\omega$ and $\Omega$.

The following estimate of volumes of sublevel sets (comp. [33, Lemma 4.1]) will be crucial for proving the volume-capacity inequality on a smoothly bounded domain (without any pseudoconvexity assumption).

Lemma 5.7 There exist constants $C_{0}, \tau_{0}>0$ such that for every $v \in \mathcal{P}_{0}(\Theta)$ with $a^{n}:=$ $\int_{\Omega_{R}} \Theta_{v}^{n}$ and $s \geq 0$,

$$
V_{2 n}(v<-s) \leq C_{0} e^{-\tau_{0} s / a} .
$$

Proof Since the domain $\Omega_{R}$ is fixed, for any Borel set $E \subset \Omega_{R}$ we will only write $\operatorname{cap}_{\Theta}(E)=$ $\operatorname{cap}_{\Theta}\left(E, \Omega_{R}\right)$ and $\operatorname{cap}(E)=\operatorname{cap}\left(E, \Omega_{R}\right)$ in the proofs.

First we show that there exists a uniform constant $C>0$ independent of $v$ such that for $s \geq 0$,

$$
\begin{equation*}
\operatorname{cap}_{\Theta}(v<-s) \leq \frac{C a^{n}}{s^{n}} . \tag{5.14}
\end{equation*}
$$

In fact, let $h \in \operatorname{PSH}\left(\Omega_{R}, \Theta\right)$ be such that $-1 \leq h \leq 0$. Note that $\Theta=d d^{c} g_{0}$ for a strictly smooth plurisubharmonic function $g_{0}$ in a neighborhood of $\bar{\Omega}$ and $g_{0} \leq 0$. Then for $\phi=h+g_{0}$,

$$
\int_{\Omega_{R}}(-v)^{n}\left(\omega+d d^{c} h\right)^{n} \leq \int_{\Omega_{R}}(-v)^{n}\left(d d^{c} \phi\right)^{n} .
$$

By Lemma 5.1 it follows that

$$
\int_{\Omega_{R}}(-v)^{n}\left(d d^{c} \phi\right)^{n} \leq n!\|\phi\|_{\infty}^{n} \int_{\Omega_{R}} \Theta_{v}^{n} .
$$

Since $\|\phi\|_{\infty} \leq 1+\left\|g_{0}\right\|_{\infty}$,

$$
\int_{\{v<-s\}}\left(\Theta+d d^{c} h\right)^{n} \leq \frac{1}{s^{n}} \int_{\Omega_{R}}(-v)^{n}\left(\Theta+d d^{c} h\right)^{n} \leq \frac{C a^{n}}{s^{n}} .
$$

Taking supremum on the left hand side over $h$ we get the desired inequality.
Next, denoting $\Omega_{s}:=\{v<-s\} \subset \Omega_{R}$ by an observation in [41, Proposition 3.5] we know that

$$
V_{2 n}\left(\Omega_{s}\right) \leq C_{1} \exp \left(\frac{-\tau_{1}}{\left[\operatorname{cap}\left(\Omega_{s}\right)\right]^{\frac{1}{n}}}\right)
$$

for uniform constants $C_{1}, \tau_{1}>0$. From this and the equivalence between $\operatorname{cap}(\bullet)$ and $\operatorname{cap}_{\Theta}(\bullet)$ (Lemma 5.6) we have

$$
V_{2 n}\left(\Omega_{s}\right) \leq C_{1} \exp \left(\frac{-\tau_{1}}{\left[A_{1} \operatorname{cap}_{\Theta}\left(\Omega_{s}\right)\right]^{\frac{1}{n}}}\right) .
$$

Combining this and (5.14) we get $V_{2 n}\left(\Omega_{s}\right) \leq C_{1} e^{-\tau_{0} s / a}$ with $\tau_{0}=\tau_{1} /\left(C A_{1}\right)^{\frac{1}{n}}$.
By comparing the Monge-Ampère measure of a Hölder continuous plurisubharmonic function and the one of its convolution we show that the estimate of measures of sublevel sets also holds.

Proposition 5.8 Denote by $\mu=\left(d d^{c} \psi\right)^{n}$, where $\psi \in \operatorname{PSH}(\Omega) \cap C^{0, \alpha}(\bar{\Omega})$ as in (5.1). Let $\tau_{0}>0$ be the uniform exponent in Lemma 5.7. There exists $\widetilde{\tau}_{0}=\widetilde{\tau}_{0}\left(n, \alpha, \tau_{0}\right)>0$ such that for $v \in \mathcal{P}_{0}(\Theta)$ with $a^{n}=\int_{\Omega_{R}} \Theta_{v}^{n} \leq 1$ and for every $s>0$,

$$
\mu(v<-s) \leq \frac{C_{0} e^{-\tilde{\tau}_{0} s / a}}{s^{n+1}}
$$

Proof Let $v_{s}=\max \{v,-s\}$. Then, $v_{s} \in \mathcal{P}_{0}(\Theta)$, and by [11, Lemma 3.4] we also know that $\int_{\Omega_{R}} \Theta_{v_{s}}^{n} \leq \int_{\Omega_{R}} \Theta_{v}^{n}=a^{n}$. We are going to show that there are uniform constant $0<\alpha_{n} \leq 1$ and $C$ independent of $s$ and $v$ such that

$$
\begin{equation*}
\int_{\Omega_{R}}\left(v_{s}-v\right)\left(d d^{c} \psi\right)^{n} \leq \frac{C}{s^{n}}\left(\int_{\Omega_{R}}\left(v_{s}-v\right) d V_{2 n}\right)^{\alpha_{n}} . \tag{5.15}
\end{equation*}
$$

Suppose this is true for a moment, and let us finish the proof of the proposition. By the inequality $0 \leq v_{s}-v \leq \mathbf{1}_{\{v<-s\}}|v| \leq \mathbf{1}_{\{v<-s\}} e^{-\tau v} / \tau$ for every $\tau>0$ (to be determined later), it follows that

$$
\begin{align*}
\int_{\Omega_{R}}\left|v_{s}-v\right| d V_{2 n} & \leq \frac{1}{\tau} \int_{\{v<-s\}} e^{-\tau v} d V_{2 n} \\
& \leq \frac{1}{\tau} \int_{\Omega_{R}} e^{-\tau v-\tau(v+s)} d V_{2 n}  \tag{5.16}\\
& \leq \frac{e^{-\tau s}}{\tau} \int_{\Omega_{R}} e^{-2 \tau v} d V_{2 n}
\end{align*}
$$

Using Lemma 5.7, for $\tau=\tau_{0} /(4 a)$, the last integral can be bounded by a uniform constant which is independent of $v$. Moreover, $v_{s / 2}-v \geq \frac{s}{2}$ on $\{v<-s\} \subset \subset \Omega_{R}$ for every $s>0$.

Therefore,

$$
\begin{aligned}
\mu(v<-s) & \leq \frac{2}{s} \int_{\Omega_{R}}\left(v_{s / 2}-v\right)\left(d d^{c} \psi\right)^{n} \\
& \leq \frac{2 C}{s^{n+1}}\left(\int_{\Omega_{R}}\left(v_{s / 2}-v\right) d V_{2 n}\right)^{\alpha_{n}},
\end{aligned}
$$

where we used (5.15) for the second inequality. Combining this and (5.16) we get for every $s>0$,

$$
\mu(v<-s) \leq \frac{C}{s^{n+1}}\left(\frac{a e^{-\tau_{0} s /(8 a)}}{\tau_{0}}\right)^{\alpha_{n}}
$$

Let $\widetilde{\tau}_{0}=\tau_{0} \alpha_{n} / 8$. Using $0<a, \alpha_{n} \leq 1$, we get

$$
\mu(v<-s) \leq \frac{C e^{-\tilde{\tau}_{0} s / a}}{s^{n+1}}
$$

This is the desired estimate.
Now let us prove the promised inequality (5.15). For $0 \leq k \leq n$ we write $S_{k}:=\left(d d^{c} \psi\right)^{k} \wedge$ $\Theta^{n-k}$. We show by induction over $k$ that

$$
\begin{equation*}
\int_{\Omega_{R}}\left(v_{s}-v\right)\left(d d^{c} \psi\right)^{k} \wedge \Theta^{n-k} \leq \frac{C}{s^{k}}\left(\int_{\Omega_{R}}\left(v_{s}-v\right) d V_{2 n}\right)^{\alpha_{k}} \tag{5.17}
\end{equation*}
$$

For $k=0$, the inequality is obviously true. Suppose it is true for $0 \leq k<n$. Denote $T=\left(d d^{c} \psi\right)^{k} \wedge \Theta^{n-k-1}$. This means that

$$
\int_{\Omega_{R}}\left(v_{s}-v\right) T \wedge \Theta \leq \frac{C}{s^{k}}\left\|v_{s}-v\right\|_{L^{1}\left(\Omega_{R}\right)}^{\alpha_{k}}
$$

for some $0<\alpha_{k} \leq 1$. We need to show that the inequality holds for $k+1 \leq n$ with possibly smaller $\alpha_{k+1}>0$ and larger $C$. Write $\psi_{\epsilon}=\psi * \chi_{\epsilon}$ on $B_{2 R}$ (where we still write $\psi$ for its $\alpha$-Hölder continuous extension onto $\left.B_{2 R}\right)$ and take $\chi_{\epsilon}(z)=\chi\left(|z|^{2} / \epsilon^{2}\right) / \epsilon^{2 n}$ the standard smoothing family defined in (5.7).

Firstly,

$$
\begin{align*}
\int_{\Omega_{R}}\left(v_{s}-v\right) d d^{c} \psi \wedge T \leq & \left|\int_{\Omega_{R}}\left(v_{s}-v\right) d d^{c} \psi_{\epsilon} \wedge T\right| \\
& +\left|\int_{\Omega_{R}}\left(v_{s}-v\right) d d^{c}\left(\psi_{\epsilon}-\psi\right) \wedge T\right|  \tag{5.18}\\
= & :\left|I_{1}\right|+\left|I_{2}\right| .
\end{align*}
$$

Since $\|\psi\|_{\infty} \leq 1$,

$$
d d^{c} \psi_{\epsilon} \leq \frac{C\|\psi\|_{\infty}}{\epsilon^{2}} \Theta \leq \frac{C \Theta}{\epsilon^{2}} \quad \text { on } B_{2 R} .
$$

Using this and the induction hypothesis we get

$$
\begin{equation*}
\left|I_{1}\right| \leq \frac{C\|\psi\|_{\infty}}{\epsilon^{2}} \int_{\Omega_{R}}\left(v_{s}-v\right) T \wedge \Theta \leq \frac{C\left\|v_{s}-v\right\|_{L^{1}\left(\Omega_{R}\right)}^{\tau_{k}}}{s^{k} \epsilon^{2}} \tag{5.19}
\end{equation*}
$$

Next, we estimate $I_{2}$. By integration by parts we rewrite it as

$$
\begin{aligned}
I_{2} & =\int_{\Omega_{R}}\left(\psi_{\epsilon}-\psi\right) d d^{c}\left(v_{s}-v\right) \wedge T \\
& =\int_{\{v \leq-s\}}\left(\psi_{\epsilon}-\psi\right)\left(\Theta_{v_{s}}-\Theta_{v}\right) \wedge T \\
& \leq \int_{\{v \leq-s\}}\left|\psi_{\epsilon}-\psi\right|\left(\Theta_{v_{s}}+\Theta_{v}\right) \wedge T
\end{aligned}
$$

Notice that $v_{s}-v=0$ outside $\{v \leq-s\}$. By the Hölder continuity of $\psi$ we have $\left|\psi_{\epsilon}-\psi\right| \leq$ $C \epsilon^{\alpha}$. Hence,

$$
\begin{equation*}
\left|I_{2}\right| \leq C \epsilon^{\alpha} \int_{\{v \leq-s\}}\left(\Theta_{v_{s}}+\Theta_{v}\right) \wedge T . \tag{5.20}
\end{equation*}
$$

It follows from Lemma 5.1 that

$$
\begin{align*}
\int_{\{v \leq-s\}} \Theta_{v_{s}} \wedge T & =\int_{\{v \leq-s\}}\left(d d^{c} \psi\right)^{k} \wedge \Theta_{v_{s}} \wedge \Theta^{n-k-1} \\
& \leq \frac{1}{s^{k}} \int_{\Omega_{R}}(-v)^{k}\left(d d^{c} \psi\right)^{k} \wedge \Theta_{v_{s}} \wedge \Theta^{n-k-1}  \tag{5.21}\\
& \leq \frac{k!\|\psi\|_{\infty}^{k}}{s^{k}} \int_{\Omega_{R}} \Theta_{v}^{k} \wedge \Theta_{v_{s}} \wedge \Theta^{n-k-1}
\end{align*}
$$

Similarly,

$$
\begin{equation*}
\int_{\{v \leq-s\}} \Theta_{v} \wedge T \leq \frac{k!\|\psi\|_{\infty}^{k}}{s^{k}} \int_{\Omega_{R}} \Theta_{v}^{k+1} \wedge \Theta^{n-k-1} \tag{5.22}
\end{equation*}
$$

Using [11, Corollary 3.5-(3)] we obtain

$$
\begin{align*}
\int_{\Omega_{R}} \Theta_{v}^{k} & \wedge \Theta_{v_{s}} \wedge \Theta^{n-k-1} \\
& \leq 2^{n-1}\left(k \int_{\Omega_{R}} \Theta_{v_{s}}^{n}+\int_{\Omega_{R}} \Theta_{v}^{n}+(n-k-1) \int_{\Omega_{R}} \Theta^{n}\right)  \tag{5.23}\\
& \leq 2^{n-1} n\left(a+C_{0}\right)
\end{align*}
$$

where $C_{0}=\int_{\Omega_{R}} \Theta^{n}$. Similarly

$$
\begin{equation*}
\int_{\Omega_{R}} \Theta_{v}^{k+1} \wedge \Theta^{n-k-1} \leq 2^{n-1} n\left(a+C_{0}\right) \tag{5.24}
\end{equation*}
$$

Combining (5.21), (5.22), (5.23) (5.24) and the assumption $0<a \leq 1$ we get that

$$
\left|I_{2}\right| \leq \frac{C \epsilon^{\alpha}}{s^{k}}
$$

It follows from the estimates for $I_{1}$ and $I_{2}$ that

$$
\begin{aligned}
& \int_{\Omega_{R}}\left(v_{s}-v\right)\left(d d^{c} \psi\right)^{k+1} \wedge \Theta^{n-k-1} \\
& \quad=\int_{\Omega_{R}}\left(v_{s}-v\right) d d^{c} \psi \wedge T \\
& \quad \leq\left|I_{1}\right|+\left|I_{2}\right| \leq \frac{C\left\|v_{s}-v\right\|_{L^{1}\left(\Omega_{R}\right)}^{\alpha_{k}}}{s^{k} \epsilon^{2}}+\frac{C \epsilon^{\alpha}}{s^{k}} .
\end{aligned}
$$

Finally, we can choose

$$
\epsilon=\left\|v_{s}-v\right\|_{L^{1}\left(\Omega_{R}\right)}^{\alpha_{k} / 3}>0, \quad \alpha_{k+1}=\alpha \alpha_{k} / 3
$$

(otherwise $v_{s}=v$ and the inequality is obvious). Then

$$
\begin{aligned}
\int_{\Omega_{R}}\left(v_{s}-v\right) S_{k+1} & =\int_{\Omega_{R}}\left(v_{s}-v\right)\left(d d^{c} \psi\right)^{k+1} \wedge \Theta^{n-k-1} \\
& \leq \frac{C}{s^{k}}\left\|v_{s}-v\right\|_{L^{1}\left(\Omega_{R}\right)}^{\alpha_{k+1}} .
\end{aligned}
$$

The proof of the step $(k+1)$ is finished, and so is the proof of the proposition.
We state a volume-capacity inequality between the Monge-Ampère measure of a Hölder continuous plurisubharmonic function and the Bedford-Taylor capacity on quasihyperconvex domains, where subsets are of $\varepsilon$-distance from the boundary.

Theorem 5.9 Let $\psi \in \operatorname{PSH}(\Omega) \cap C^{0, \alpha}(\bar{\Omega})$ as in (5.1). Suppose $0<\varepsilon<R / 4$. Then there exist uniform constants $C, \tau_{0}>0$ such that for every compact set $K \subset D_{\varepsilon}(R) \cap B_{R / 2}$,

$$
\begin{equation*}
\int_{K}\left(d d^{c} \psi\right)^{n} \leq \frac{C}{\varepsilon^{n+1}} \exp \left(\frac{-\tau_{0} \varepsilon}{\left[\operatorname{cap}\left(K, \Omega_{R}\right)\right]^{\frac{1}{n}}}\right) . \tag{5.25}
\end{equation*}
$$

In particular, for any $\tau>0$,

$$
\int_{K}\left(d d^{c} \psi\right)^{n} \leq \frac{C_{\tau}}{\varepsilon^{n+1}\left(\tau_{0} \varepsilon\right)^{n(1+\tau)}}\left[\operatorname{cap}\left(K, \Omega_{R}\right)\right]^{1+\tau} .
$$

Proof Our first observation is that we only need to consider subsets satisfying $\operatorname{cap}_{\Theta}(K) \leq 1$. Otherwise, if $\operatorname{cap}_{\Theta}(K)>1$, then by Lemma 5.6, it follows that $\operatorname{cap}\left(K, \Omega_{R}\right) \geq 1 / A_{1}$. This implies that for $0<\tau_{0} \leq 2 n$,

$$
\exp \left(\frac{-\tau_{0} \varepsilon}{\left[\operatorname{cap}\left(K, \Omega_{R}\right)\right]^{\frac{1}{n}}}\right) \geq \exp \left(\frac{-\tau_{0} \varepsilon}{A_{1}^{\frac{1}{n}}}\right) \geq \exp \left(\frac{-n R / 2}{A_{1}^{\frac{1}{n}}}\right) .
$$

Then, the inequality follows from Corollary 5.2 with some uniform constant $C$. In what follows we work with a subset $E \subset \Omega_{R}$ satisfying

$$
\operatorname{cap}_{\Theta}\left(E, \Omega_{R}\right) \leq 1
$$

Since the domain $\Omega_{R}$ is fixed, we omit it in capacities formulae. We consider yet another capacity which takes into account the geometry of the domain. Let $E \subset \Omega_{R}$ be a Borel subset,

$$
\begin{equation*}
\operatorname{cap}_{\rho}(E):=\sup \left\{\int_{E} \Theta_{w}^{n}: w \in \operatorname{PSH}\left(\Omega_{R}, \Theta\right), \rho_{R} \leq w \leq 0\right\} . \tag{5.26}
\end{equation*}
$$

Recall that $-1 \leq \rho \leq 0$ is the defining function for $\partial M$ on $B_{2 R}$. By Definition (5.2) we have $-1 \leq \rho_{R} \leq 0$ in $\Omega$. Hence,

$$
\begin{equation*}
\operatorname{cap}_{\rho}(E) \leq \operatorname{cap}_{\Theta}(E) \tag{5.27}
\end{equation*}
$$

So $\operatorname{cap}_{\rho}(\bullet)$ does not charge pluripolar sets. Thus, without loss of generality we may assume $K$ is a compact regular subset, in the sense that $h_{\rho, K}$ is continuous in $\Omega_{R}$. The relative extremal function of this $\rho$-capacity is given by

$$
h_{\rho, K}(x)=\sup \left\{w(x): w \in \operatorname{PSH}\left(\Omega_{R}, \Theta\right), w_{\left.\right|_{K}} \leq \rho_{R}, w \leq 0\right\} .
$$

The desired property of $h_{\rho, K}$ is that $\rho_{R} \leq h_{\rho, K} \leq 0$, which implies that $h_{\rho, K}$ has zero boundary value and $h_{\rho, K} \in \mathcal{P}_{0}(\Theta)$. Moreover, for every compact set $K \subset \Omega_{R}$ the balayage argument shows $\left(\Theta+d d^{c} h_{\rho, K}\right)^{n} \equiv 0$ on $\Omega_{R} \backslash K$. Therefore,

$$
\begin{equation*}
\operatorname{cap}_{\rho}(K) \geq \int_{K}\left(\Theta+d d^{c} h_{\rho, K}\right)^{n}=\int_{\Omega_{R}}\left(\Theta+d d^{c} h_{\rho, K}\right)^{n} . \tag{5.28}
\end{equation*}
$$

(This inequality is indeed an identity, we refer the readers to [2, Proposition 6.5] and its generalization in $[16,17])$. Note that

$$
K \subset\left\{h_{\rho, K}=\rho_{R}\right\}=\left\{h_{\rho, K} \leq \rho_{R}\right\} .
$$

Hence

$$
K \subset\left\{h_{\rho, K} \leq \sup _{K} \rho_{R}=: \delta_{K}\right\} .
$$

Write $\mu:=\left(d d^{c} \psi\right)^{n}$. Applying Proposition 5.8 for $h_{\rho, K} \in \mathcal{P}_{0}(\Theta)$ we get

$$
\mu(K) \leq \mu\left(h_{\rho, K} \leq \delta_{K}\right) \leq \frac{C}{\left|\delta_{K}\right|^{n+1}} \exp \left(\frac{-\widetilde{\tau}_{0}\left|\delta_{K}\right|}{a}\right),
$$

where

$$
a^{n}=\int_{\Omega_{R}}\left(\Theta+d d^{c} h_{\rho, K}\right)^{n} \leq \operatorname{cap}_{\rho}(K) \leq \operatorname{cap}_{\Theta}(K) \leq 1 .
$$

Thus

$$
\mu(K) \leq \frac{C}{\left|\delta_{K}\right|^{n+1}} \exp \left(\frac{-\tilde{\tau}_{0}\left|\delta_{K}\right|}{\left[\operatorname{cap}_{\Theta}(K)\right]^{\frac{1}{n}}}\right) .
$$

If $K \subset\left\{z \in \Omega_{R}: \rho_{R}(z)<-\varepsilon\right\} \cap B_{R / 2}$, then $\left|\delta_{K}\right|=\left|\sup _{K} \rho_{R}\right| \geq \varepsilon$; hence for such compact sets

$$
\begin{equation*}
\mu(K) \leq \frac{C}{\varepsilon^{n+1}} \exp \left(\frac{-\widetilde{\tau}_{0} \varepsilon}{\left[\operatorname{cap}_{\Theta}(K)\right]^{\frac{1}{n}}}\right) . \tag{5.29}
\end{equation*}
$$

By the equivalence of the capacities in Lemma 5.6 the proof of the theorem follows with $\tau_{0}=\widetilde{\tau}_{0} / A_{1}^{\frac{1}{n}}$.

Remark 5.10 In an interior coordinate chart of $\bar{M}$ we have the corresponding inequality due to Dinh-Nguyen-Sibony [19]. Let us identify a fixed holomorphic coordinate ball with $B_{2 R} \subset \mathbb{C}^{n}$. Suppose that

$$
\psi \in P S H\left(B_{2 R}\right) \cap C^{0, \alpha}\left(\bar{B}_{2 R}\right)
$$

with $0<\alpha \leq 1$. Then, there exist uniform constants $C=C(R, \psi)$ and $\tau_{0}=\tau_{0}(n, \alpha)>0$ such that for every compact subset $K \subset B_{R / 2}$,

$$
\begin{equation*}
\int_{K}\left(d d^{c} \psi\right)^{n} \leq C \exp \left(\frac{-\tau_{0}}{\left[\operatorname{cap}\left(K, B_{R}\right)\right]^{\frac{1}{n}}}\right) . \tag{5.30}
\end{equation*}
$$

We are in the position to state the main stability estimate of this section.

Proposition 5.11 Take $\psi \in \operatorname{PSH}(\Omega) \cap C^{0, \alpha}(\bar{\Omega})$ as in (5.1). Assume that $u, v \in \operatorname{PSH}(\Omega) \cap$ $L^{\infty}(\Omega)$ and $u=v$ on $\{z \in \Omega: \rho(z) \geq-\varepsilon\}$. Then,

$$
\int_{\Omega \cap B_{R / 2}}|u-v|\left(d d^{c} \psi\right)^{n} \leq \frac{C}{\varepsilon^{n}}\|u-v\|_{L^{1}\left(\Omega_{R}\right)}^{\alpha_{1}},
$$

where $C=C\left(\Omega, R, \psi,\|u\|_{\infty},\|v\|_{\infty}\right)>0$ and $\alpha_{1}=\alpha_{1}(n, \alpha)>0$ are uniform constants.
Proof Note that

$$
\begin{aligned}
\partial\left(\Omega \cap B_{R / 2}\right) & =\left(\partial \Omega \cap B_{R / 2}\right) \cup\left(\Omega \cap \partial B_{R / 2}\right) \\
& =\left(\{\rho=0\} \cap B_{R / 2}\right) \cup\left(\Omega \cap \partial B_{R / 2}\right) .
\end{aligned}
$$

By the assumption $u=v$ on $\{z \in \Omega: \rho(z) \geq-\varepsilon\}$ the integrand on the left hand side is zero near the first portion of the boundary. Furthermore $\psi$ is Hölder continuous plurisubharmonic function on $\Omega$ a neighborhood of the second boundary portion, so it has finite Monge-Ampère mass by the Chern-Levine-Nirenberg inequality.

By subtracting from $u, v$ a constant we may assume that $u, v \leq 0$. Also dividing both sides by $\left(1+\|u\|_{\infty}\right)\left(1+\|v\|_{\infty}\right)\left(1+\|\psi\|_{\infty}\right)^{n}$ we may assume that

$$
-1 \leq u, v, \psi \leq 0 \quad \text { on } \Omega .
$$

First we suppose $v \geq u$. Let $0 \leq \eta \leq 1$ be a cut-off function in $\Omega$ such that $\eta \equiv 1$ on $\Omega \cap B_{R / 2}$ and supp $\eta \subset \Omega \cap B_{3 R / 4}$. Then

$$
\operatorname{supp}(\eta(v-u)) \subset D_{\varepsilon}(R) \subset \subset \Omega_{R}
$$

Thus it is enough to estimate

$$
\int_{\Omega_{R}} \eta(v-u)\left(d d^{c} \psi\right)^{n} .
$$

Let us still write $\psi$ for its Hölder continuous extension of $\psi \in C^{0, \alpha}$ onto $B_{2 R}$. We are going to prove, by induction over $0 \leq k \leq n$, the inequalities

$$
\begin{equation*}
\int_{\Omega_{R}} \eta(v-u)\left(d d^{c} \psi\right)^{k} \wedge \Theta^{n-k} \leq \frac{C}{\varepsilon^{k}}\|u-v\|_{L^{1}\left(\Omega_{R}\right)}^{\tau_{k}} \tag{5.31}
\end{equation*}
$$

For $k=n$ it is our statement, for $k=0$ the inequality holds with $\alpha_{1}=\tau_{0}=1$. Assume that the inequality is true for $k \in[0, n)$. We need to show it for $k+1$ with possibly smaller $\tau_{k+1}>0$ and larger $C$.

Consider the standard regularizing family $\chi_{t}(z)=\chi\left(|z|^{2} / t^{2}\right)$ for $0<t<R / 4$. Define

$$
T_{k}=\left(d d^{c} \psi\right)^{k} \wedge \Theta^{n-k}
$$

Write $\psi_{t}=\psi * \chi_{t}$ on $B_{2 R}$. Then, for $T=\left(d d^{c} \psi\right)^{k} \wedge \Theta^{n-k-1}$ we have

$$
\begin{align*}
\int_{\Omega_{R}} \eta(v-u) d d^{c} \psi \wedge T \leq & \left|\int_{\Omega_{R}} \eta(v-u) d d^{c} \psi_{t} \wedge T\right| \\
& +\left|\int_{\Omega_{R}} \eta(v-u) d d^{c}\left(\psi_{t}-\psi\right) \wedge T\right|  \tag{5.32}\\
= & :\left|I_{1}\right|+\left|I_{2}\right| .
\end{align*}
$$

Since $\|\psi\|_{\infty} \leq 1$,

$$
d d^{c} \psi_{t} \leq \frac{C\|\psi\|_{\infty}}{t^{2}} \Theta \leq \frac{C \Theta}{t^{2}} \quad \text { on } B_{2 R} .
$$

Using this and the induction hypothesis we get

$$
\begin{equation*}
\left|I_{1}\right| \leq \frac{C\|\psi\|_{\infty}}{t^{2}} \int_{\Omega_{R}} \eta(v-u) T \wedge \Theta \leq \frac{C\|v-u\|_{L^{1}\left(\Omega_{R}\right)}^{\tau_{k}}}{t^{2}} \tag{5.33}
\end{equation*}
$$

Next, by integration by parts we rewrite the second integral in (5.32) as

$$
I_{2}=\int_{\Omega_{R}}\left(\psi_{t}-\psi\right) d d^{c}(\eta(v-u)) \wedge T .
$$

Compute

$$
\begin{align*}
d d^{c}[\eta(v-u)] \wedge T= & (v-u) d d^{c} \eta \wedge T+2 d \eta \wedge d^{c}(v-u) \wedge T  \tag{5.34}\\
& +\eta d d^{c}(v-u) \wedge T .
\end{align*}
$$

Note that $\eta$ is smooth on $\Omega$, so $d d^{c} \eta \leq C \Theta$. By the Cauchy-Schwarz inequality

$$
\begin{aligned}
& \left|\int\left(\psi_{t}-\psi\right) d \eta \wedge d^{c}(v-u) \wedge T\right|^{2} \\
& \quad \leq \int_{D_{\varepsilon}(R)}\left(\psi_{t}-\psi\right)^{2} d \eta \wedge d^{c} \eta \wedge T \int_{D_{\varepsilon}(R)} d(v-u) \wedge d^{c}(v-u) \wedge T
\end{aligned}
$$

Observe that

$$
\begin{aligned}
\left|\psi_{t}(z)-\psi(z)\right| & \leq \int_{B(0,1)}|\psi(z-t w)-\psi(z)| \chi\left(|w|^{2}\right) d V_{2 n}(w) \\
& \leq C_{\alpha} t^{\alpha}
\end{aligned}
$$

with $C_{\alpha}$ the Hölder norm of $\psi$ on $\bar{\Omega}$, and $d \eta \wedge d^{c} \eta \leq C_{1} \Theta$,

$$
\int_{D_{\varepsilon}(R)}\left(\psi_{t}-\psi\right)^{2} d \eta \wedge d^{c} \eta \wedge T \leq C_{\alpha} C_{1} t^{2 \alpha} \int_{D_{\varepsilon}(R)} T \wedge \Theta \leq \frac{C t^{2 \alpha}\|\psi\|_{\infty}^{k}}{\varepsilon^{k}}
$$

where we used Corollary 5.2 for the last inequality. Similarly by Remark 5.3,

$$
\int_{D_{\varepsilon}(R)} d(v-u) \wedge d^{c}(v-u) \wedge T \leq \frac{C\|\psi\|_{\infty}^{k}\|u\|_{\infty}^{2}\|v\|_{\infty}^{2}}{\varepsilon^{k+1}}
$$

For the last term in (5.34) using Corollary 5.2 again and the Hölder continuity of $\psi$ we have

$$
\begin{aligned}
\left|\int_{D_{\varepsilon}(R)}\left(\psi_{t}-\psi\right) \eta d d^{c}(v-u) \wedge T\right| & \leq \int_{D_{\varepsilon}(R)}\left|\psi_{t}-\psi\right| \eta\left(d d^{c} u+d d^{c} v\right) \wedge T \\
& \leq \frac{C t^{\alpha}\left(\|u\|_{\infty}+\|v\|_{\infty}\right)\|\psi\|_{\infty}^{k}}{\varepsilon^{k+1}}
\end{aligned}
$$

Combining the above estimates we conclude

$$
\begin{equation*}
\left|I_{2}\right| \leq \frac{C t^{\alpha}}{\varepsilon^{k+1}} \tag{5.35}
\end{equation*}
$$

From (5.33) and (5.35) we get

$$
\int_{\Omega_{R}} \eta(v-u)\left(d d^{c} \psi\right)^{k+1} \leq\left|I_{1}\right|+\left|I_{2}\right| \leq \frac{C\|u-v\|_{L^{1}\left(\Omega_{R}\right)}^{\tau_{k}}}{t^{2}}+\frac{C t^{\alpha}}{\varepsilon^{k+1}} .
$$

If $\|u-v\|_{L^{1}\left(\Omega_{R}\right)}^{\tau_{k} / 4} \geq R / 4$, then the inequality of step $(k+1)$ holds for a fixed $t=R / 8$ and $\tau_{k+1}=\alpha \tau_{k} / 4$. On the other hand, we can choose $t=\|u-v\|_{L^{1}\left(\Omega_{R}\right)}^{\tau_{k} / 4}$ and this implies for $\tau_{k+1}=\alpha \tau_{k} / 4>0$,

$$
\int_{\Omega_{R}} \eta(v-u)\left(d d^{c} \psi\right)^{k+1} \leq C\|u-v\|_{L^{1}\left(\Omega_{R}\right)}^{\tau_{k+1}} .
$$

The induction proof is completed under extra hypothesis $v \geq u$. For the general case use the identity

$$
|u-v|=(\max \{u, v\}-u)+(\max \{u, v\}-v),
$$

and apply the above proof for the pairs $(\max \{u, v\}, u)$ and $(\max \{u, v\}, v)$.
By a similar (easier) argument for an interior chart of $\bar{M}$, which we identify with the ball $B_{2 R}$ of radius $2 R>0$ centered at 0 in $\mathbb{C}^{n}$, we get the following stability estimate.

Lemma 5.12 Suppose that $\psi \in \operatorname{PSH}\left(B_{2 R}\right) \cap C^{0, \alpha}\left(\bar{B}_{2 R}\right)$ with $0<\alpha \leq 1$. Assume that $u, v \in \operatorname{PSH}\left(B_{2 R}\right) \cap L^{\infty}\left(B_{2 R}\right)$. Then,

$$
\int_{B_{R / 2}}|v-u|\left(d d^{c} \psi\right)^{n} \leq C\|u-v\|_{L^{1}\left(B_{R}\right)}^{\alpha_{1}},
$$

where $C=C\left(R, \psi,\|u\|_{\infty},\|v\|_{\infty}\right)$ and $\alpha_{1}=\alpha_{1}(n, \alpha)>0$ are uniform constants.
Proof Let $T=\left(d d^{c} u\right)^{k} \wedge\left(d d^{c} v\right)^{\ell} \wedge\left(d d^{c} \psi\right)^{m} \wedge \Theta^{n-k-\ell-m}$. Since these functions are plurisubharmonic on $B_{2 R}$, the Chern-Levine-Nirenberg inequality gives

$$
\int_{B_{R}} T \leq C(R)\|u\|_{\infty}^{k}\|v\|_{\infty}^{\ell}\|\psi\|_{\infty}^{m}
$$

Then, the proof goes exactly along the lines of the proof of Proposition 5.11.

## 6 Hölder continuous subsolution theorems

In this section we prove Theorem 1.4. We first show that the global capacity is equivalent to the Bedford-Taylor capacity defined via a finite covering. We fix a finite covering of $\bar{M}$ $-\left\{B_{i}(s)\right\}_{i \in I} \cup\left\{U_{j}(s)\right\}_{j \in J}$, where $B_{i}(s)=B_{i}\left(x_{i}, s\right)$ and $U_{j}(s)=U_{j}\left(y_{j}, s\right)$ are coordinate balls and coordinate half-balls centered at $x_{i}$ and $y_{j}$ respectively, and of radius $0<s<1$. We choose $s>0$ so small that $B_{i}(2 s)$ and $U_{j}(2 s)$ are still contained in holomorphic charts. For any Borel set $E \subset M$ we can define another capacity

$$
\begin{equation*}
\operatorname{cap}^{\prime}(E)=\sum_{i \in I} \operatorname{cap}\left(E \cap B_{i}(s), B_{i}(2 s)\right)+\sum_{j \in J} \operatorname{cap}\left(E \cap U_{j}(s), U_{j}(2 s)\right), \tag{6.1}
\end{equation*}
$$

where $\operatorname{cap}(\bullet, \bullet)$ on the right hand side is just the Bedford-Taylor capacity.
Proposition 6.1 Two capacities cap ${ }_{\omega}$ and cap' are equivalent. More precisely, there exists a uniform constant $A_{0}>0$ such that for any Borel set $E \subset M$,

$$
\begin{equation*}
\frac{1}{A_{0}} \operatorname{cap}^{\prime}(E) \leq \operatorname{cap}_{\omega}(E) \leq A_{0} \operatorname{cap}^{\prime}(E) \tag{6.2}
\end{equation*}
$$

Proof The proof is an adaptation from [32]. We first prove for a uniform $C>0$,

$$
\begin{equation*}
\operatorname{cap}_{\omega}(E) \leq \operatorname{Ccap}^{\prime}(E) \tag{6.3}
\end{equation*}
$$

Let $i \in I \cup J$, and let $U(s)$ be either $B_{i}(s)$ or $U_{i}(s)$. Assume that $\omega \leq d d^{c} g$ for a strictly plurisubharmonic function $g \leq 0$ on a neighborhood of $\bar{U}(2 s)$. Consider $v \in \operatorname{PSH}(M, \omega)$ with $-1 \leq v \leq 0$.

$$
\begin{aligned}
\int_{E \cap U(s)}\left(\omega+d d^{c} v\right)^{n} & \leq \int_{E \cap U(s)}\left(d d^{c}(g+v)\right)^{n} \\
& \leq\left(\|g\|_{\infty}+1\right)^{n} \operatorname{cap}(E \cap U(s), U(2 s)) .
\end{aligned}
$$

Since $|I \cup J|$ is finite, the first inequality follows from the sub-additivity of $\operatorname{cap}_{\omega}(\cdot)$. The inverse inequality will follow if one can show that for a fixed $i \in I \cup J$ and $U(s)$ as above there is a uniform $C>0$ such that

$$
\begin{equation*}
\operatorname{cap}(E \cap U(s), U(2 s)) \leq \operatorname{Ccap}_{\omega}(E) \tag{6.4}
\end{equation*}
$$

In fact, let $v \in \operatorname{PSH}(U(2 s))$ and $-1 \leq v \leq 0$ in $U(2 s)$. We wish to find $\tilde{v}$ such that $\tilde{v}=a v-a$ on $U(s),-1 \leq \tilde{v} \leq 0$ and $\tilde{v} \in \operatorname{PSH}(M, \omega)$, where $0<a<1 / 2$ is a uniform constant depending only on $M, \omega$.

First we take a smooth function $\eta$ such that $\eta=0$ on $\bar{M} \backslash U(2 s)$ and $\eta<0$ in $U(2 s)$. Consider the function $\epsilon \eta$ for $\epsilon>0$ small. Since $\omega$ is a Hermitian metric on $\bar{M}$, we can choose $\epsilon>0$ depending on $\eta$ and $\omega$ such that $\epsilon \eta \in \operatorname{PSH}(M, \omega)$. Choose $0<a<1 / 2$ so that $\epsilon \eta \leq-3 a$ on $U(s)$. Writing $\eta$ for $\epsilon \eta$, we conclude that there exists a smooth $\omega$-psh function $\eta=0$ on $M \backslash U(2 s)$ and $\eta \leq-3 a$ on $\bar{U}(s)$ for $0<a<1 / 2$.

The function $\tilde{v}$ is defined as follows:

$$
\tilde{v}:= \begin{cases}\max \{a v-a, \eta\} & \text { on } \quad U(2 s),  \tag{6.5}\\ 0 & \text { on } \quad M \backslash U(2 s) .\end{cases}
$$

As $\lim \sup _{z \rightarrow \zeta} v(z) \leq-a<\eta(\zeta)=0$ for $\zeta \in \partial U(2 s) \cap M, \tilde{v} \in P S H(M, \omega)$. It is easy to see that $\tilde{v}$ satisfies all requirements. Thus,

$$
\begin{aligned}
\int_{E}\left(\omega+d d^{c} \tilde{v}\right)^{n} & \geq \int_{E \cap U(s)}\left(\omega+a d d^{c} v\right)^{n} \\
& \geq a^{n} \int_{E \cap U(s)}\left(d d^{c} v\right)^{n} .
\end{aligned}
$$

By taking supremum over $v$ it implies that $\operatorname{cap}(E \cap U(s), U(2 s)) \leq a^{-n} \operatorname{cap}_{\omega}(E)$.
Using the equivalence of capacities above and local volume-capacity inequalities on boundary and interior charts (Theorem 5.9 and Remark 5.10) we derive the global measurecapacity estimate on $M$. For $\varepsilon>0$ small let us denote

$$
\begin{equation*}
M_{\varepsilon}=\left\{z \in M: \operatorname{dist}_{\omega}(z, \partial M)>\varepsilon\right\}, \tag{6.6}
\end{equation*}
$$

where $\operatorname{dist}_{\omega}(\bullet, \partial M)$ is the distance function on $M$ with respect to the Riemannian metric induced by $\omega$.

Lemma 6.2 Let $\underline{u} \in \operatorname{PSH}(M, \omega) \cap C^{0, \alpha}(\bar{M})$ for some $0<\alpha \leq 1$. Let $\mu$ be a positive Borel measure on $M$. Suppose $\mu \leq\left(\omega+d d^{c} \underline{u}\right)^{n}$ in $M$. Then there exist uniform constants
$C, \alpha_{0}>0$ (also independent of $\varepsilon$ ) such that for every compact set $K \subset M_{\varepsilon}$,

$$
\mu(K) \leq \frac{C}{\varepsilon^{n+1}} \exp \left(-\frac{\alpha_{0} \varepsilon}{\left[\operatorname{cap}_{\omega}(K)\right]^{\frac{1}{n}}}\right)
$$

In particular, for any $\tau>0$,

$$
\mu(K) \leq \frac{C_{\tau}}{\varepsilon^{n+1}\left(\alpha_{0} \varepsilon\right)^{n(1+\tau)}}\left[\operatorname{cap}_{\omega}(K)\right]^{1+\tau},
$$

where $C_{\tau}$ depends additionally on $\tau$.

Proof Cover $\bar{M}$ by finitely many coordinate balls $B_{i}(R / 2)$ and coordinate half-balls $U_{j}(R / 2)$ with $R>0$ (fixed) so that $B_{i}(2 R)$ and $U_{j}(2 R)$ are still contained in holomorphic charts of $\bar{M}$. Let $K \subset M_{\varepsilon}$ be a compact set.

Consider its part $K_{i}=K \cap B_{i}(R / 2)$ which is contained in an interior chart $B_{i}(2 R)$. On this chart we can choose a strictly plurisubharmonic function $g$ such that $d d^{c} g \geq \omega$. Set $\psi=g+\underline{u}$. Then, $\mu \leq\left(d d^{c} \psi\right)^{n}$ on $B_{i}(2 R)$. Subtracting a constant we may assume that $\psi \leq 0$ on $\bar{B}_{i}(2 R)$. Applying Remark 5.10 and Proposition 6.1 we get

$$
\begin{aligned}
\mu\left(K_{i}\right) & \leq C \exp \left(\frac{-\tau_{0}}{\left[\operatorname{cap}\left(K_{i}, B_{R}\right)\right]^{\frac{1}{n}}}\right) \\
& \leq C \exp \left(\frac{-\tau_{0}}{\left[A_{0} \operatorname{cap}_{\omega}\left(K_{i}\right)\right]^{\frac{1}{n}}}\right) \\
& \leq C \exp \left(\frac{-\alpha_{0}}{\left[\operatorname{cap}_{\omega}(K)\right]^{\frac{1}{n}}}\right),
\end{aligned}
$$

where $\alpha_{0}=\tau_{0} / A_{0}^{\frac{1}{n}}$ is a uniform constant and we used $\operatorname{cap}_{\omega}\left(K_{i}\right) \leq \operatorname{cap}_{\omega}(K)$ in the last inequality.

Next consider $K_{j}=K \cap U_{j}(R / 2)$ which is contained in a boundary chart $\Omega=U_{j}(2 R)$. Similarly as above $\mu \leq\left(d d^{c} \psi\right)^{n}$ on $\Omega$ for a negative Hölder continuous $\psi$ on $\bar{\Omega}$ which is plurisubharmonic in $\Omega$. Now Theorem 5.9 and Proposition 6.1 give

$$
\begin{aligned}
\mu\left(K_{j}\right) & \leq \frac{C}{\varepsilon^{n+1}} \exp \left(\frac{-\tau_{0} \varepsilon}{\left[\operatorname{cap}\left(K_{j}, \Omega_{R}\right)\right]^{\frac{1}{n}}}\right) \\
& \leq \frac{C}{\varepsilon^{n+1}} \exp \left(\frac{-\tau_{0} \varepsilon}{\left[\operatorname{cap}^{\prime}(K)\right]^{\frac{1}{n}}}\right) \\
& \leq \frac{C}{\varepsilon^{n+1}} \exp \left(\frac{-\tau_{0} \varepsilon}{A_{0}^{\frac{1}{n}}\left[\operatorname{cap}_{\omega}(K)\right]^{\frac{1}{n}}}\right) .
\end{aligned}
$$

Since $|I \cup J|$ is finite, we conclude

$$
\begin{aligned}
\mu(K) & \leq \sum_{i \in I} \mu\left(K_{i}\right)+\sum_{j \in J} \mu\left(K_{j}\right) \\
& \leq C_{1} \exp \left(\frac{-\alpha_{0}}{\left[\operatorname{cap}_{\omega}(K)\right]^{\frac{1}{n}}}\right)+\frac{C_{2}}{\varepsilon^{n+1}} \exp \left(\frac{-\alpha_{0} \varepsilon}{\left[\operatorname{cap}_{\omega}(K)\right]^{\frac{1}{n}}}\right) \\
& \leq \frac{C}{\varepsilon^{n+1}} \exp \left(\frac{-\alpha_{0} \varepsilon}{\left[\operatorname{cap}_{\omega}(K)\right]^{\frac{1}{n}}}\right)
\end{aligned}
$$

This is the desired estimate.
We now fix the notation to finish the proof of Theorem 1.4. Let $\mu$ be a positive Borel measure on $M$. Suppose that there exists $\underline{u} \in \operatorname{PSH}(M, \omega) \cap C^{0, \alpha}(\bar{M})$ with $0<\alpha \leq 1$, a Hölder continuous subsolution for $\mu$ on $M$, satisfying

$$
\underline{u}_{\left.\right|_{\partial M}}=\varphi \in C^{0, \alpha}(\partial M) .
$$

By Theorem 1.2 and Corollary 4.6 there exists a solution $u \in P S H(M, \omega) \cap C^{0}(\bar{M})$ solving

$$
\begin{equation*}
\left(\omega+d d^{c} u\right)^{n}=\mu, \quad \lim _{z \rightarrow q} u(z)=\varphi(q) \quad \text { for every } q \in \partial M \tag{6.7}
\end{equation*}
$$

Proposition 6.3 Let $u \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be the solution to (6.7). Let $v \in$ $\operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that $v=u$ on $M \backslash M_{\varepsilon}$. Then there is $0<\alpha_{2} \leq 1$ such that

$$
\sup _{M}(v-u) \leq \frac{C}{\varepsilon^{3 n+1}}\left(\int_{M} \max \{v-u, 0\} d \mu\right)^{\alpha_{2}}
$$

where $C=C\left(M, \omega, \underline{u},\|u\|_{\infty},\|v\|_{\infty}\right)$ and $\alpha_{2}=\alpha_{2}(n, \alpha)>0$ are uniform constants.
Proof Subtracting a constant and then dividing both sides of the inequality by $\left(1+\|u\|_{\infty}+\right.$ $\|v\|_{\infty}$ ) we may assume that

$$
-1 \leq u, v \leq 0 \quad \text { on } \bar{M}
$$

Let us assume also that $-s_{0}=\sup _{M}(v-u)>0$, otherwise the statement trivially follows.
We will make use of Theorem 4.3. There for a given number $0<\theta<1$ we defined

$$
\theta_{0}:=\frac{1}{3} \min \left\{\theta^{n}, \frac{\theta^{3}}{16 \mathbf{B}}, 4(1-\theta) \theta^{n}, 4(1-\theta) \frac{\theta^{3}}{16 \mathbf{B}},\left|s_{0}\right|\right\},
$$

$m(\theta)=\inf _{M}[u-(1-\theta) v]$ and $U(\theta, t):=\{u<(1-\theta) v+m(\theta)+t\}$ for $0<t<\theta_{0}$. Then, for $0<\theta<\left|s_{0}\right| / 3$ and $0<t<\theta_{0}$ we have

$$
s_{0}+\theta+2 t \leq 0
$$

Since $-1 \leq u, v \leq 0$, it is clear that

$$
s_{0}-\theta \leq m(\theta) \leq s_{0}
$$

It follows that

$$
U(\theta, 2 t) \subset\left\{u<v+s_{0}+\theta+2 t\right\} \subset M_{\varepsilon}
$$

where in the last inclusion the assumption $v=u$ on $M \backslash M_{\varepsilon}$ is used.

Now the proof is identical to the one in [36, Proposition 2.4] except that we need to replace [36, Lemma 2.6] there by Lemma 6.2 at the cost of extra factor $\varepsilon^{-n(2+\tau)-1}$ in the uniform constants. If we fix $\tau=1$ in that lemma then

$$
\mu(K) \leq \frac{C_{\alpha}[\operatorname{cap}(K, \Omega)]^{2}}{\left(\alpha_{0} \varepsilon\right)^{3 n+1}}
$$

Therefore as in [36, Proposition 2.4] we can now take

$$
\alpha_{2}=\frac{1}{1+(n+2)(n+1)} .
$$

The next step is to estimate $L^{1}(d \mu)$-norm in terms of $L^{1}\left(d V_{2 n}\right)$-norm. Again this estimate is obtained for functions that are equal outside $M_{\varepsilon}$.

Lemma 6.4 Let $u$ be the solution to the Eq. (6.7). Let $v \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that $v=u$ on $M \backslash M_{\varepsilon}$. Then there is $0<\alpha_{3} \leq 1$ such that

$$
\int_{M}|u-v| d \mu \leq \frac{C}{\varepsilon^{n}}\left(\int_{M}|u-v| d V_{2 n}\right)^{\alpha_{3}}
$$

where $C=C\left(M, \omega, \underline{u},\|u\|_{\infty},\|v\|_{\infty}\right)$ and $\alpha_{3}=\alpha_{3}(n, \alpha)>0$ are uniform constants.
Proof As usual we cover $\bar{M}$ by a finite number of coordinate balls and half-balls of radius $R / 2$ so that the ones with radius $2 R$ are still contained in holomorphic charts. On a local coordinate chart $V$ consider a strictly plurisubharmonic function $g \leq 0$ such that

$$
d d^{c} g \geq \omega .
$$

Then we write $u^{\prime}=u+g, v^{\prime}=v+g$ and $\psi=\underline{u}+g$. They are plurisubharmonic functions on $V$. Moreover,

$$
\mu \leq\left(\omega+d d^{c} \underline{u}\right)^{n} \leq\left(d d^{c} \psi\right)^{n} \quad \text { on } V,
$$

where $\psi \in C^{0, \alpha}(\bar{V})$ is a Hölder continuous plurisubharmonic function on $V$.
On an interior chart $B_{2 R}$ of $\bar{M}$ by Lemma 5.12 we have

$$
\int_{B_{R / 2}}|u-v| d \mu \leq \int_{B_{R / 2}}\left|u^{\prime}-v^{\prime}\right|\left(d d^{c} \psi\right)^{n} \leq C\left(\int_{B_{R}}\left|u^{\prime}-v^{\prime}\right| d V_{2 n}\right)^{\alpha^{\prime}}
$$

This is bounded by $C\left(\int_{M}|u-v| d V_{2 n}\right)^{\alpha^{\prime}}$.
We now consider the case of a boundary chart $\Omega$. Let $\rho$ be is the defining function of $\partial M$ on $\Omega$ as in Sect. 5. Then, there exists a uniform constant $c_{0}>0$ such that

$$
|\rho(z)| \geq c_{0} \operatorname{dist}_{\omega}(z, \partial M) \quad \text { for all } z \in \Omega
$$

(shrinking $\Omega$ if necessary so that dist ${ }_{\omega}$ is a smooth function in $\Omega$ ). Hence $u=v$ on $\{|\rho(z)| \leq$ $\left.c_{0} \varepsilon\right\}$ by the assumption. Without loss of generality we may assume $c_{0}=1$. The conclusion follows from Proposition 5.11.

Since the covering is finite, the proof of the lemma follows.
We now proceed to find the Hölder exponent of the solution $u$ of the Eq. (6.7) over $\bar{M}$. By subtracting a uniform constant we may assume that

$$
u \leq 0 \quad \text { on } \bar{M} .
$$

Let $\delta>0$ be small. For $z \in M_{\delta}$ we define

$$
\begin{equation*}
\widehat{u}_{\delta}(z)=\sup \left\{u(x): x \in M \quad \text { and } \operatorname{dist}_{\omega}(x, z) \leq \delta\right\} . \tag{6.8}
\end{equation*}
$$

where $\operatorname{dist}_{\omega}$ is the Riemannian distance induced by the metric $\omega$. We wish to show that there exist constants $c_{0}>0, \delta_{0}>0$ and an exponent $0<\tau \leq 1$ such that for every $0<\delta \leq \delta_{0}$,

$$
\begin{equation*}
\sup _{\bar{M}_{\delta}}\left(\widehat{u}_{\delta}-u\right)=\sup _{M_{\delta}}\left(\widehat{u}_{\delta}-u\right) \leq c_{0} \delta^{\tau} . \tag{6.9}
\end{equation*}
$$

To do it let us fix a small constant $\delta_{0}>0$ such that for every $0<\delta<\delta_{0}, \sup _{M_{\delta}} \underline{u}_{\delta}-$ $\underline{u}) \leq C \delta^{\alpha}$ (for which we use Hölder continuity of the subsolution). Now we consider two parameters $\delta, \varepsilon$ such that

$$
\begin{equation*}
0<\delta \leq \varepsilon<\delta_{0} \tag{6.10}
\end{equation*}
$$

By a classical argument we obtain the following estimate on the $(\varepsilon, \delta)$-collars near the boundary $\partial M$.

Lemma 6.5 Consider $0<\delta \leq \varepsilon<\delta_{0}$. There is $0<\tau_{1} \leq 1$ such that for $z \in \overline{M_{\delta}} \backslash M_{\varepsilon}$,

$$
\widehat{u}_{\delta}(z) \leq u(z)+c_{1} \varepsilon^{\tau_{1}},
$$

where $\tau_{1}$ depends only on $M, \omega$ and $\alpha$ the Hölder exponent of $\underline{u}$. Moreover, for $z \in M$ and $q \in \partial M$ with $\operatorname{dist}_{\omega}(z, q) \leq \delta$ we have

$$
|u(z)-u(q)| \leq c_{1} \delta^{\tau_{1}} .
$$

Proof Let $h_{1} \in C^{0}(\bar{M}, \mathbb{R})$ be the unique solution to the linear PDE:

$$
\begin{array}{r}
\left(\omega+d d^{c} h_{1}\right) \wedge \omega^{n-1}=0,  \tag{6.11}\\
h_{1}=\varphi \text { on } \partial M .
\end{array}
$$

Note that $\left(\omega+d d^{c} u\right) \wedge \omega^{n-1} \geq 0$. The maximum principle for the Laplace operator with respect to $\omega$ gives $u \leq h_{1}$. It is classical fact from [28, Theorem 6] (see also [46, Theorem 5.3]) that $h_{1}$ is also Hölder continuous on $\bar{M}$ :

$$
h_{1} \in C^{0, \tau_{1}}(\bar{M}),
$$

with $0<\tau_{1} \leq \alpha$ (decreasing $\tau_{1}$ if necessary). Fix $z \in \overline{M_{\delta}} \backslash M_{\varepsilon}$. Since $u$ is continuous, there exists a point $z_{\delta} \in M$ with $\operatorname{dist}_{\omega}\left(z_{\delta}, z\right) \leq \delta$ such that $u\left(z_{\delta}\right)=\widehat{u}_{\delta}(z)$. Let $q \in \partial M$ be the point that is closest to $z$, i.e., $\operatorname{dist}_{\omega}(z, q) \leq \varepsilon$. By the fact that $\underline{u} \leq u \leq h_{1}$ on $M$ we have

$$
\widehat{u}_{\delta}(z)-u(z)=u\left(z_{\delta}\right)-u(z) \leq h_{1}\left(z_{\delta}\right)-\underline{u}(z) .
$$

Since $h_{1}(q)=\varphi(q)=\underline{u}(q)$, it follows that

$$
\begin{aligned}
h_{1}\left(z_{\delta}\right)-\underline{u}(z) & =\left(h_{1}\left(z_{\delta}\right)-h_{1}(q)\right)+(\underline{u}(q)-\underline{u}(z)) \\
& \leq\left(h_{1}\left(z_{\delta}\right)-h_{1}(z)\right)+\left(h_{1}(z)-h_{1}(q)\right)+(\underline{u}(q)-\underline{u}(z)) \\
& \leq C\left(\delta^{\tau_{1}}+\varepsilon^{\tau_{1}}\right),
\end{aligned}
$$

where $C$ is maximum of the Hölder norms of $h_{1}$ and $\underline{u}$ on $\bar{M}$. Because $\delta \leq \varepsilon$, the proof of the first inequality is completed. The remaining inequality follows from a similar argument.

Remark 6.6 The constants $c_{1}, \tau_{1}$ in the last lemma are independent of parameters $\delta, \varepsilon$. This applies also to all uniform constants appearing in the estimates that follow.

Corollary 6.7 The Hölder continuity of u over $\bar{M}$ will follow once we prove (6.9).
Proof We need to justify that for every $x, y \in \bar{M}$,

$$
|u(x)-u(y)| \leq C \operatorname{dist}_{\omega}(x, y)^{\tau},
$$

where $\tau=\tau_{1}$ in Lemma 6.5. We may assume that $0<\operatorname{dist}_{\omega}(x, y)=\delta \leq \delta_{0}$, and $\operatorname{dist}_{\omega}(x, \partial M)=\delta_{x} \leq \delta_{y}=\operatorname{dist}_{\omega}(y, \partial M)$.

Case 1: $\delta_{y} \leq \delta$. Then there exist $q_{x}, q_{y} \in \partial M$ such that $\operatorname{dist}_{\omega}\left(x, q_{x}\right) \leq \delta$ and $\operatorname{dist}_{\omega}\left(y, q_{y}\right) \leq \delta$. Since $\operatorname{dist}_{\omega}\left(q_{x}, q_{y}\right) \leq 3 \delta$, by the Hölder continuity of $\underline{u}$ we have $\left|u\left(q_{x}\right)-u\left(q_{y}\right)\right|=\left|\underline{u}\left(q_{x}\right)-\underline{u}\left(q_{y}\right)\right| \leq C \delta^{\tau}$. It follows from Lemma 6.5 that

$$
\begin{aligned}
|u(x)-u(y)| & \leq\left|u(x)-u\left(q_{x}\right)\right|+\left|u(y)-u\left(q_{y}\right)\right|+\left|u\left(q_{x}\right)-u\left(q_{y}\right)\right| \\
& \leq C \operatorname{dist}_{\omega}(x, y)^{\tau} .
\end{aligned}
$$

Case 2: $\delta_{x} \geq \delta$. Then, $u(y) \leq \widehat{u}_{\delta}(x)$ and $u(x) \leq \widehat{u}_{\delta}(y)$. Therefore, $u(x)-u(y) \leq$ $\widehat{u}_{\delta}(y)-u(y) \leq C \delta^{\tau}$ by (6.9). Similarly, $u(y)-u(x) \leq C \delta^{\tau}$.

Case 3: $\delta_{x}<\delta<\delta_{y}$. Without loss of generality we may assume that $\delta$ is a regular value of $\operatorname{dist}_{\omega}(\bullet, \partial M)$; that means: there exists a point $p \in \partial M_{\delta}$, which is the intersection of the shortest path joining $x, y$ and $\partial M_{\delta}$, such that

$$
\max \left\{\operatorname{dist}_{\omega}(x, p), \operatorname{dist}_{\omega}(y, p)\right\} \leq \delta
$$

Hence,

$$
|u(x)-u(y)| \leq|u(x)-u(p)|+|u(y)-u(p)| \leq C \operatorname{dist}_{\omega}(x, y)^{\tau},
$$

where we used Case 1 and Case 2 for the last inequality.
Now we use the global regularization of a quasi-plurisubharmonic function due to Demailly [14] which provides a lower bound on the complex Hessian.

Let $u$ be the continuous solution to the Eq. (6.7). Consider

$$
\begin{equation*}
\rho_{\delta} u(z)=\frac{1}{\delta^{2 n}} \int_{\zeta \in T_{z} M} u\left(\exp h_{z}(\zeta)\right) \chi\left(\frac{|\zeta|_{\omega}^{2}}{\delta^{2}}\right) d V_{\omega}(\zeta), \quad z \in M_{\delta}, \delta>0 \tag{6.12}
\end{equation*}
$$

where $\zeta \rightarrow \exp h_{z}(\zeta)$ is the (formal) holomorphic part of the Taylor expansion of the exponential map of the Chern connection on the tangent bundle of $M$ associated to $\omega$, and the mollifier $\chi: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$is given by (5.7) above.

By [14, Remark 4.6] we have for $0<t<\delta_{0}$ small

$$
\begin{align*}
\rho_{t} u(z) & =\int_{|x|<1} u(z+t x) \chi\left(|x|^{2}\right) d V_{2 n}(x)+O\left(t^{2}\right)  \tag{6.13}\\
& =u * \chi_{t}(z)+O\left(t^{2}\right)
\end{align*}
$$

in the normal coordinate system centered at $z$ (see [14, Proposition 2.9]). Since the metric $\omega$ is smooth on $\bar{M}$, the distance function satisfies

$$
\operatorname{dist}_{\omega}(z, x)=|z-x|+O\left(|z-x|^{2}\right)
$$

as $x \rightarrow z$. It follows that for $0<\delta \leq \delta_{0}\left(\delta_{0}\right.$ small enough)

$$
\begin{equation*}
\rho_{\delta} u(z) \leq \sup _{|z-x| \leq \delta} u(x)+C \delta^{2} \leq \widehat{u}_{2 \delta}(z)+c_{2} \delta \quad \text { for every } z \in \overline{M_{\delta}} . \tag{6.14}
\end{equation*}
$$

We observe that to prove the Hölder continuity it is enough to work with this geodesic convolution.

Lemma 6.8 Suppose that there exist constants $0<\tau \leq 1$ and $C>0$ such that for $0<\delta \leq$ $\delta_{0}$,

$$
\begin{equation*}
\rho_{\delta} u(z)-u(z) \leq C \delta^{\tau} \quad \text { for every } z \in M_{\delta} . \tag{6.15}
\end{equation*}
$$

Then $u$ is Hölder continuous on $\bar{M}$.
Proof By Corollary 6.7 it is enough to verify the inequality (6.9), i.e,

$$
\sup _{M_{\delta}}\left(\widehat{u}_{\delta}-u\right) \leq C^{\prime} \delta^{\tau} .
$$

First, it follows from (6.13) that in normal coordinates containing ball of radius $\delta$ we have $u * \chi_{\delta}-u \leq C \delta^{\tau}$. Therefore, (5.12) implies $\check{u}_{\delta / 2}-u \leq C \delta^{\tau}$. Now, the proof of the required inequality follows the lines of argument given in [26, Lemma 4.2] (see also [48, Theorem 3.2]).

Now let us state the important estimate for the complex Hessian of $\rho_{\delta} u$. The proof of the following variation of [14, Proposition 3.8] and [3, Lemma 1.12] was given in [37, Lemma 4.1].

Lemma 6.9 Let $0<\delta<\delta_{0}$ and $\rho_{t} u$ be as in (6.12). Define the Kiselman-Legendre transform with level $b>0$ by

$$
\begin{equation*}
u_{\delta, b}(z)=\inf _{t \in[0, \delta]}\left(\rho_{t} u(z)+K t^{2}+K t-b \log \frac{t}{\delta}\right), \tag{6.16}
\end{equation*}
$$

Then for some positive constant $K$ depending on the curvature, the function $\rho_{t} u+K t^{2}$ is increasing in $t$ and the following estimate holds:

$$
\begin{equation*}
\omega+d d^{c} u_{\delta, b} \geq-(A b+2 K \delta) \omega \text { on } \overline{M_{\delta}} . \tag{6.17}
\end{equation*}
$$

where $A$ is a lower bound of the negative part of the Chern curvature of $\omega$.
Thanks to this lemma we construct an $\omega$-psh function $U_{\delta}$ which approximates the solution $u$.

Lemma 6.10 Let $0<\tau \leq 1$ and $b=\left(\delta^{\tau}-2 K \delta\right) / A=O\left(\delta^{\tau}\right)$, where $K, A, \delta$ are parameters in Lemma 6.9. Define $U_{\delta}:=\left(1-\delta^{\tau}\right) u_{\delta, b}$. Then, $U_{\delta} \in \operatorname{PSH}\left(M_{\delta}, \omega\right)$ satisfies

$$
U_{\delta} \leq u+c_{3} \delta^{\tau}+c_{1} \varepsilon^{\tau_{1}} \quad \text { on } \partial M_{\varepsilon} .
$$

Moreover,

$$
U_{\delta} \leq \rho_{\delta} u+c_{3} \delta^{\tau} \quad \text { on } \overline{M_{\delta}},
$$

where $c_{1}, c_{3}$ are uniform constants independent of $\varepsilon$ and $\delta$.
Proof By Lemma 6.9 we have $\omega+d d^{c} U_{\delta} \geq \delta^{2 \tau} \omega$. The monotonicity of $\rho_{t} u+K t^{2}$ implies

$$
\begin{equation*}
u \leq u_{\delta, b} \leq \rho_{\delta} u+2 K \delta \quad \text { on } \overline{M_{\delta}} . \tag{6.18}
\end{equation*}
$$

On the boundary $\partial M_{\varepsilon}$, using (6.14) and Lemma 6.5 (with $\delta^{\prime}=2 \delta<\varepsilon$ ), we have

$$
\begin{equation*}
\rho_{\delta} u \leq \widehat{u}_{2 \delta}+c_{2} \delta \leq u+c_{1} \varepsilon^{\tau_{1}}+c_{2} \delta . \tag{6.19}
\end{equation*}
$$

Since $u_{\delta, b} \leq 0$ is uniformly bounded, it is easy to see that

$$
\begin{equation*}
\left(1-\delta^{\tau}\right) u_{\delta, b} \leq u_{\delta, b}+C \delta^{\tau} . \tag{6.20}
\end{equation*}
$$

By (6.18), (6.19) and (6.20) we have on $\partial M_{\varepsilon}$,

$$
\begin{align*}
U_{\delta} & \leq \rho_{\delta} u+2 K \delta+C \delta^{\tau} \\
& \leq u+c_{1} \varepsilon^{\tau_{1}}+\left(2 K+c_{2}\right) \delta+C \delta^{\tau} . \tag{6.21}
\end{align*}
$$

The proof of on the boundary part is completed by taking $c_{3}=2 K+c_{2}+C$. Furthermore, combining (6.18) and (6.20) we get $U_{\delta} \leq \rho_{\delta} u+c_{3} \delta^{\tau}$ on $\overline{M_{\delta}}$.

By this lemma and the stability estimates we get the following bound.
Proposition 6.11 Let $0<\delta \leq \varepsilon \leq \delta_{0}$. Let $U_{\delta}$ be the function defined in Lemma 6.10. Then,

$$
\sup _{M_{\delta}}\left(U_{\delta}-u\right) \leq \frac{C \delta^{\alpha_{2} \alpha_{3}}}{\varepsilon^{4 n+1}}+c_{1} \varepsilon^{\tau_{1}}
$$

where the exponents $\alpha_{2}, \alpha_{3}>0$ and $\tau_{1}$ are from Proposition 6.3, Lemmas 6.4, and 6.5, respectively.

Proof Using Lemma 6.10 we can produce:

$$
\tilde{U}_{\delta}= \begin{cases}\max \left\{U_{\delta}-c_{1} \varepsilon^{\tau_{1}}-c_{3} \delta^{\tau}, u\right\} & \text { on } M_{\varepsilon},  \tag{6.22}\\ u & \text { on } M \backslash M_{\varepsilon},\end{cases}
$$

which is an $\omega$-psh function on $M$. Next we apply Proposition 6.3 to get that for some $0<\alpha_{2} \leq 1$,

$$
\begin{equation*}
\sup _{M}\left(\widetilde{U}_{\delta}-u\right) \leq \frac{C}{\varepsilon^{3 n+1}}\left\|\left(\widetilde{U}_{\delta}-u\right)_{+}\right\|_{L^{1}(d \mu)}^{\alpha_{2}} . \tag{6.23}
\end{equation*}
$$

Note that $0<\delta \leq \varepsilon$, so $M_{\varepsilon} \subset M_{\delta}$. Then,

$$
\begin{aligned}
\left(\widetilde{U}_{\delta}-u\right)_{+} & =\max \left\{U_{\delta}-c_{1} \varepsilon^{\tau_{1}}-c_{3} \delta^{\tau}-u, 0\right\} \cdot \mathbf{1}_{M_{\varepsilon}} \\
& \leq \max \left\{\rho_{\delta} u-u, 0\right\} \cdot \mathbf{1}_{M_{\varepsilon}},
\end{aligned}
$$

where we used $U_{\delta}-c_{3} \delta^{\tau}-u \leq \rho_{\delta} u-u$ on $\bar{M}_{\delta}$ (Lemma 6.10) for the second inequality. Hence,

$$
\left(\widetilde{U}_{\delta}-u\right)_{+}=\mathbf{1}_{M_{\delta}} \cdot\left(\widetilde{U}_{\delta}-u\right)_{+} \leq \mathbf{1}_{M_{\delta}} \cdot\left(\rho_{\delta}-u\right)_{+} .
$$

This combined with Lemma 6.4 gives for some $0<\alpha_{3} \leq 1$,

$$
\begin{align*}
\left\|\left(\widetilde{U}_{\delta}-u\right)_{+}\right\|_{L^{1}(d \mu)} & \leq \frac{C}{\varepsilon^{n}}\left\|\left(\widetilde{U}_{\delta}-u\right)_{+}\right\|_{L^{1}\left(d V_{2 n}\right)}^{\alpha_{3}} \\
& \leq \frac{C}{\varepsilon^{n}}\left\|\mathbf{1}_{M_{\delta}} \cdot\left(\rho_{\delta}-u\right)_{+}\right\|_{L^{1}\left(d V_{2 n}\right)}^{\alpha_{3}} \tag{6.24}
\end{align*}
$$

The next step is to show that $L^{1}$-norm on the right hand side has a desired bound. Covering $M_{2 \delta}$ be finitely many normal charts (contained in coordinates balls or coordinate half-balls) and invoking the inequalities (6.13), (5.11) and Lemma 5.4 one obtains that for $0<\delta<\delta_{0}$,

$$
\int_{M_{2 \delta}}\left(\rho_{\delta} u-u\right)_{+} d V_{2 n} \leq C \delta^{2}+C \delta .
$$

This implies

$$
\begin{equation*}
\int_{M_{\delta}}\left(\rho_{\delta}-u\right)_{+} d V_{2 n} \leq C \int_{M_{\delta} \backslash M_{2 \delta}} d V_{2 n}+\int_{M_{2 \delta}}\left(\rho_{\delta}-u\right)_{+} d V_{2 n} \leq C \delta, \tag{6.25}
\end{equation*}
$$

where we used the compactness of $\bar{M}$ to get $V_{2 n}\left(M_{\delta} \backslash M_{2 \delta}\right) \leq C \delta$.
We conclude from (6.23), (6.24), (6.25) and $0<\alpha_{2}, \alpha_{3} \leq 1$ that

$$
\begin{equation*}
\sup _{M}\left(\widetilde{U}_{\delta}-u\right) \leq \frac{C \delta^{\alpha_{2} \alpha_{3}}}{\varepsilon^{4 n+1}} \tag{6.26}
\end{equation*}
$$

Notice that

$$
\begin{aligned}
\sup _{M_{\delta}}\left(U_{\delta}-u\right) & \leq \max \left\{\sup _{M_{\varepsilon}}\left(U_{\delta}-u\right), \sup _{\bar{M}_{\delta} \backslash M_{\varepsilon}}\left(U_{\delta}-u\right)\right\} \\
& \leq \max \left\{\sup _{M}\left(\widetilde{U}_{\delta}-u\right), \sup _{\bar{M}_{\delta} \backslash M_{\varepsilon}}\left(U_{\delta}-u\right)\right\} .
\end{aligned}
$$

Combining this with (6.26) and Lemma 6.5 we get

$$
\begin{aligned}
\sup _{M_{\delta}}\left(U_{\delta}-u\right) & \leq \sup _{M}\left(\widetilde{U}_{\delta}-u\right)+c_{1} \varepsilon^{\tau_{1}} \\
& \leq \frac{C \delta^{\alpha_{2} \alpha_{3}}}{\varepsilon^{4 n+1}}+c_{1} \varepsilon^{\tau_{1}} .
\end{aligned}
$$

This is the desired inequality.
We are ready to verify the hypothesis of Lemma 6.8.
End of the proof of the Hölder continuity of $u$ Let us choose $\alpha_{4}=\alpha_{2} \alpha_{3} / 2(4 n+1)>0$ and $\varepsilon=\delta^{\alpha_{4}}$, then Proposition 6.11 implies

$$
\sup _{M_{\delta}}\left(U_{\delta}-u\right) \leq C \delta^{\alpha_{4}}+c_{1} \delta^{\tau_{1} \alpha_{4}} \quad \text { on } \bar{M}_{\delta} .
$$

Therefore it follows from $u_{\delta, b} \leq 0$ and $0<\tau_{1}, \alpha_{4} \leq 1$ that

$$
\begin{equation*}
u_{\delta, b}-u \leq U_{\delta}-u \leq C \delta^{\tau_{1} \alpha_{4}} \quad \text { on } \bar{M}_{\delta} . \tag{6.27}
\end{equation*}
$$

Let us fix a point $z \in \bar{M}_{\delta}$. The minimum in the definition of $u_{\delta, b}$ is realized at $t_{0}=t_{0}(z)$. So, at this point we have

$$
\rho_{t_{0}} u+K t_{0}^{2}+K t_{0}-b \log \left(t_{0} / \delta\right)-u \leq C \delta^{\tau_{1} \alpha_{4}} .
$$

Since $\rho_{t} u+K t^{2}+K t-u \geq 0$, we have $b \log \left(t_{0} / \delta\right) \geq-C \delta^{\tau_{1} \alpha_{4}}$. Hence,

$$
t_{0} \geq e^{-C \delta^{\tau_{1} \alpha_{4}} / b} \delta,
$$

where

$$
\frac{\delta^{\tau_{1} \alpha_{4}}}{b}=\frac{A \delta^{\tau_{1} \alpha_{4}}}{\delta^{\tau}-2 K \delta} .
$$

Now we choose $\tau=\tau_{1} \alpha_{4}>0$, so $b \geq \delta^{\alpha_{1} \alpha_{4}} / 2 A$, which implies

$$
t_{0} \geq c_{4} \delta
$$

with $c_{4}=e^{-2 C A}$ as a uniform constant.
Since $t \mapsto \rho_{t} u+K t^{2}$ is increasing in $t$,

$$
\begin{aligned}
& \rho_{c_{4} \delta} u(z)+K\left(c_{4} \delta\right)^{2}+K\left(c_{4} \delta\right)-u(z) \\
& \quad \leq \rho_{t_{0}} u(z)+K t_{0}^{2}+K t_{0}-b \log \left(t_{0} / \delta\right)-u(z) \\
& \quad=u_{\delta, b}(z)-u(z) \\
& \quad \leq C \delta^{\tau},
\end{aligned}
$$

where we used the definition of $t_{0}$ and (6.27) for the third identity and the last inequality, respectively. Hence, $\rho_{c_{4} \delta} u(z)-u(z) \leq C \delta^{\tau} \leq C \delta^{\tau}$. The desired estimate is obtained by rescaling $\delta:=c_{4} \delta$.

Let us give the proof of uniqueness of solutions on Stein manifolds or when $\omega$ is Kähler. We will need a variation of Lemma 4.1 for the proof of uniqueness on Stein manifolds. In this setting there exists a strictly plurisubharmonic function on the whole manifold. This is a straightforward generalization of [35, Theorem 3.1].

Lemma 6.12 Fix $\theta>0$ small. Let $u, v \in \operatorname{PSH}(M, \omega) \cap L^{\infty}(\bar{M})$ be such that $\liminf _{z \rightarrow \partial M}(u-v) \geq 0$. Suppose that $-s_{0}=\sup _{M}(v-u)>0$ and $\omega+d d^{c} v \geq \theta \omega$ in $M$. Then for any $0<s<\theta_{0}=\min \left\{\frac{\theta^{n}}{16 \mathbf{B}},\left|s_{0}\right|\right\}$,

$$
\int_{\left\{u<v+s_{0}+s\right\}} \omega_{v}^{n} \leq\left(1+\frac{C_{n} \mathbf{B} s}{\theta^{n}}\right) \int_{\left\{u<v+s_{0}+s\right\}} \omega_{u}^{n} .
$$

Proof of Corollary 1.5 We first assume that $\omega$ is Kähler. Let $\varepsilon>0$ and define $u_{\varepsilon}=\max \{u+$ $\varepsilon, v\}$. Then, $u_{\varepsilon} \geq u+\varepsilon$ on $M$ and by the assumption we have $u_{\varepsilon}=u+\varepsilon$ near $\partial M$. Moreover, since $\omega_{v}^{n} \geq \omega_{u}^{n}$, it follows from a well-known inequality of Demailly that

$$
\left(\omega+d d^{c} \max \{u+\varepsilon, v\}\right)^{n} \geq \mathbf{1}_{\{u+\varepsilon \geq v\}} \omega_{u}^{n}+\mathbf{1}_{\{u+\varepsilon<v\}} \omega_{v}^{n} \geq \omega_{u}^{n}
$$

Applying a result of Błocki [5, Theorem 2.3] for $u_{\varepsilon}$ and $u+\varepsilon$ we get $u_{\varepsilon}=u+\varepsilon$ on $M$ (strictly speaking he only stated for continuous functions, but the proof works for bounded functions). Thus, $u+\varepsilon \geq v$ on $M$. Letting $\varepsilon$ to zero we get the proof of the corollary in the first case.

Next, suppose that $M$ is a Stein manifold and $\omega$ is a Hermitian metric on $\bar{M}$. Let $\rho \in$ $C^{\infty}(M)$ be a strictly plurisubharmonic exhaustive function for $M$. As in the previous case we only need to prove $u+\varepsilon \geq v$ on $M$ for a fixed $\varepsilon>0$, then let $\varepsilon$ to zero. Hence we may assume that $\liminf _{z \rightarrow \partial M}(u-v)(z) \geq \varepsilon$.

Let $M^{\prime} \subset \subset M$ be such that $u \geq v+\varepsilon$ on $M \backslash M^{\prime}$. By subtracting a uniform constant we may assume that $u, v \leq 0$ and $-C_{0} \leq \rho \leq 0$ on $\overline{M^{\prime}}$ for some constant $C_{0}>0$. Under these assumptions the proof follows the lines of the one in [35, Corollary 3.4] after replacing $\Omega$ by $M^{\prime}$ and the comparison principle by Lemma 6.12.

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