On the maximal operator of a general Ornstein–Uhlenbeck semigroup

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Abstract

If Q is a real, symmetric and positive definite $n \times n$ matrix, and B a real $n \times n$ matrix whose eigenvalues have negative real parts, we consider the Ornstein–Uhlenbeck semigroup on $\mathbb{R}^n$ with covariance Q and drift matrix B . Our main result says that the associated maximal operator is of weak type $(1, 1)$ with respect to the invariant measure. The proof has a geometric gist and hinges on the “forbidden zones method” previously introduced by the third author.

Keywords Ornstein–Uhlenbeck semigroup · Maximal operator · Gaussian measure · Mehler kernel · Weak type $(1, 1)$

Mathematics Subject Classification 47D03 · 42B25

1 Introduction

In this paper we prove a weak type $(1, 1)$ theorem for the maximal operator associated to a general Ornstein–Uhlenbeck semigroup. We extend the proof given by the third author in 1983 in a symmetric context. Our setting is the following.

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In $\mathbb{R}^n$ we will consider the semigroup generated by the elliptic operator

$$\mathcal{L} = \frac{1}{2} \sum_{i,j=1}^n q_{ij} \frac{\partial^2}{\partial x_i \partial x_j} + \sum_{i,j=1}^n b_{ij} x_i \frac{\partial}{\partial x_j},$$

or, equivalently,

$$\mathcal{L} = \frac{1}{2} \operatorname{tr}(Q \nabla^2) + \langle Bx, \nabla \rangle,$$

where ∇ is the gradient and ∇^2 the Hessian. Here $Q = (q_{ij})$ is a real, symmetric and positive definite $n \times n$ matrix, indicating the covariance of $\mathcal{L}$. The real $n \times n$ matrix $B = (b_{ij})$ is negative in the sense that all its eigenvalues have negative real parts, and it gives the drift of $\mathcal{L}$.

The semigroup is formally $\mathcal{H}_t = e^{t\mathcal{L}}$, $t > 0$, but to write it more explicitly we first introduce the positive definite, symmetric matrices

$$Q_t = \int_0^t e^{sB} Q e^{sB^*} ds, \quad 0 < t \leq +\infty, \quad (1.1)$$

and the normalized Gaussian measures γ_t in $\mathbb{R}^n$, with $t \in (0, +\infty]$, having density

$$y \mapsto (2\pi)^{-\frac{n}{2}} (\det Q_t)^{-\frac{1}{2}} \exp\left(-\frac{1}{2} \langle Q_t^{-1} y, y \rangle\right)$$

with respect to Lebesgue measure. Then for functions f in the space of bounded continuous functions in $\mathbb{R}^n$ one has

$$\mathcal{H}_t f(x) = \int f(e^{tB} x - y) d\gamma_t(y), \quad x \in \mathbb{R}^n, \quad (1.2)$$

a formula due to Kolmogorov. The measure γ_∞ is invariant under the action of $\mathcal{H}_t$; it will be our basic measure, replacing Lebesgue measure.

We remark that $(\mathcal{H}_t)_{t>0}$ is the transition semigroup of the stochastic process

$$\chi(x, t) = e^{tB} x + \int_0^t e^{(t-s)B} dW(s),$$

where W is a Brownian motion in $\mathbb{R}^n$ with covariance Q .

We are interested in the maximal operator defined as

$$\mathcal{H}_* f(x) = \sup_{t>0} |\mathcal{H}_t f(x)|.$$

Under the above assumptions on Q and B , our main result is the following.

Theorem 1.1 *The Ornstein–Uhlenbeck maximal operator $\mathcal{H}_*$ is of weak type $(1, 1)$ with respect to the invariant measure γ_∞ , with an operator quasinorm that depends only on the dimension and the matrices Q and B .*

In other words, the inequality

$$\gamma_\infty\{x \in \mathbb{R}^n : \mathcal{H}_* f(x) > \alpha\} \leq \frac{C}{\alpha} \|f\|_{L^1(\gamma_\infty)}, \quad \alpha > 0, \quad (1.3)$$

holds for all functions $f \in L^1(\gamma_\infty)$, with $C = C(n, Q, B)$.

For large values of the time parameter, we also obtain a refinement of this result. Indeed, we prove in Proposition 6.1 that

$$\gamma_\infty \left\{ x \in \mathbb{R}^n : \sup_{t>1} |\mathcal{H}_t f(x)| > \alpha \right\} \leq \frac{C}{\alpha \sqrt{\log \alpha}} \quad (1.4)$$

for large $\alpha > 0$ and all normalized functions $f \in L^1(\gamma_\infty)$. Here $C = C(n, Q, B)$, and this estimate is shown to be sharp. It cannot be extended to $\mathcal{H}_*$, since the maximal operator corresponding to small values of t only satisfies the ordinary weak type inequality. This sharpening is not surprising, in the light of some recent results for the standard case $Q = I$ and $B = -I$ by Lehec [8]. He proved the following conjecture, proposed by Ball, Barthe, Bednorz, Oleszkiewicz and Wolff [2]:

For each fixed $t > 0$, there exists a function $\psi_t = \psi_t(\alpha)$, with $\lim_{\alpha \rightarrow +\infty} \psi_t(\alpha) = 0$, satisfying

$$\gamma_\infty \{x \in \mathbb{R}^n : |\mathcal{H}_t f(x)| > \alpha\} \leq \frac{\psi_t(\alpha)}{\alpha}$$

for all large $\alpha > 0$ and all $f \in L^1(\gamma_\infty)$ such that $\|f\|_{L^1(\gamma_\infty)} = 1$. Lehec proved this conjecture with $\psi_t(\alpha) = C(t)/\sqrt{\log \alpha}$ independent of the dimension, and this ψ_t is sharp. Our estimates depend strongly on the dimension n , but on the other hand we estimate the supremum over large t .

The history of $\mathcal{H}_*$ is quite long and started with the first attempts to prove L^p estimates. When $(\mathcal{H}_t)_{t>0}$ is symmetric, i.e., when each operator $\mathcal{H}_t$ is self-adjoint on $L^2(\gamma_\infty)$, then $\mathcal{H}_*$ is bounded on $L^p(\gamma_\infty)$ for $1 < p \leq \infty$, as a consequence of the general Littlewood–Paley–Stein theory for symmetric semigroups of contractions on L^p spaces [16, Ch. III].

It is easy to see that the maximal operator is unbounded on $L^1(\gamma_\infty)$. This led, about fifty years ago, to the study of the weak type $(1, 1)$ of $\mathcal{H}_*$ with respect to γ_∞ . The first positive result is due to B. Muckenhoupt [13], who proved the estimate (1.3) in the one-dimensional case with $Q = I$ and $B = -I$. The analogous question in the higher-dimensional case was an open problem until 1983, when the third author [15] proved the weak type $(1, 1)$ in any finite dimension. Other proofs are due to Menárguez, Pérez and Soria [11] (see also [10, 14]) and to García-Cuerva, Mauceri, Meda, Sjögren and Torrea [7]. Moreover, a different proof of the weak type $(1, 1)$ of $\mathcal{H}_*$, based on a covering lemma halfway between covering results by Besicovitch and Wiener, was given by Aimar, Forzani and Scotto [1]. A nice overview of the literature may be found in [17, Ch.4].

In [4] the present authors recently considered a normal Ornstein–Uhlenbeck semigroup in $\mathbb{R}^n$, that is, we assumed that $\mathcal{H}_t$ is for each $t > 0$ a normal operator on $L^2(\gamma_\infty)$. Under this extra assumption, we proved that the associated maximal operator is of weak type $(1, 1)$ with respect to the invariant measure γ_∞ . This extends earlier work in the non-symmetric framework by Mauceri and Noselli [9], who proved that if $Q = I$ and $B = \lambda(R - I)$ for some positive λ and a real skew-symmetric matrix R generating a periodic group, then the maximal operator $\mathcal{H}_*$ is of weak type $(1, 1)$.

In Theorem 1.1 we go beyond the hypothesis of normality. The proof has a geometric core and relies on the *ad hoc* technique developed by the third author in [15]. It is worth noticing that, while the proof in [4] required an analysis of the special case when $Q = I$ and $B = (-\lambda_1, \dots, -\lambda_n)$, with $\lambda_j > 0$ for $j = 1, \dots, n$, and then the application of factorization results, we apply here directly, avoiding many intermediate steps, the "forbidden zones" technique introduced in [15].

Since the maximal operator $\mathcal{H}_*$ is trivially bounded from L^∞ to L^∞ , we obtain by interpolation the following corollary.

Corollary 1.2 *The Ornstein–Uhlenbeck maximal operator $\mathcal{H}_*$ is bounded on $L^p(\gamma_\infty)$ for all $p > 1$.*

This result improves Theorem 4.2 in [9], where the L^p boundedness of $\mathcal{H}_*$ is proved for all $p > 1$ in the normal framework, under the additional assumption that the infinitesimal generator of $(\mathcal{H}_t)_{t>0}$ is a sectorial operator of angle less than $\pi/2$.

In this paper we focus our attention on the Ornstein–Uhlenbeck semigroup in $\mathbb{R}^n$. In view of possible applications to stochastic analysis and to SPDE's, it would be very interesting to investigate the case of the infinite-dimensional Ornstein–Uhlenbeck maximal operator as well (see [3, 6, 18] for an introduction to the infinite-dimensional setting). The Riesz transforms associated to a general Ornstein–Uhlenbeck semigroup in $\mathbb{R}^n$ have been studied in the authors' paper [5].

The scheme of the paper is as follows. In Sect. 2 we introduce the Mehler kernel $K_t(x, u)$, that is, the integral kernel of $\mathcal{H}_t$. Some estimates for the norm and the determinant of Q_t and related matrices are provided in Sect. 3. As a consequence, we obtain bounds for the Mehler kernel. In Sect. 4 we consider the relevant geometric features of the problem, and introduce in Sect. 4.1 a system of polar-like coordinates. We also express Lebesgue measure in terms of these coordinates. Sections 5, 6, 7 and 8 are devoted to the proof of Theorem 1.1. First, Sect. 5 introduces some preliminary simplifications of the proof; in particular, we restrict the variable x to an ellipsoidal annulus. In Sect. 6 we consider the supremum in the definition of the maximal operator taken only over $t > 1$ and prove the sharp estimate (1.4). Section 7 is devoted to the case of small t under an additional local condition. Finally, in Sect. 8 we treat the remaining case and conclude the proof of Theorem 1.1, by proving the estimate (1.3) for small t under a global assumption.

In the following, we use the “variable constant convention”, according to which the symbols $c > 0$ and $C < \infty$ will denote constants which are not necessarily equal at different occurrences. They all depend only on the dimension and on Q and B . For any two nonnegative quantities a and b we write $a \lesssim b$ instead of $a \leq Cb$ and $a \gtrsim b$ instead of $a \geq cb$. The symbol $a \simeq b$ means that both $a \lesssim b$ and $a \gtrsim b$ hold.

By $\mathbb{N}$ we mean the set of all nonnegative integers. If A is an $n \times n$ matrix, we write $\|A\|$ for its operator norm on $\mathbb{R}^n$ with the Euclidean norm $|\cdot|$.

2 The Mehler kernel

For $t > 0$, the difference

$$Q_\infty - Q_t = \int_t^\infty e^{sB} Q e^{sB^*} ds \quad (2.1)$$

is a symmetric and strictly positive definite matrix. So is the matrix

$$Q_t^{-1} - Q_\infty^{-1} = Q_t^{-1}(Q_\infty - Q_t)Q_\infty^{-1}, \quad (2.2)$$

and we can define

$$D_t = (Q_t^{-1} - Q_\infty^{-1})^{-1} Q_t^{-1} e^{tB}, \quad t > 0. \quad (2.3)$$

Then formula (1.2), the definition of the Gaussian measure and some elementary computations yield

$$\begin{aligned}\mathcal{H}_t f(x) &= (2\pi)^{-\frac{n}{2}} (\det Q_t)^{-\frac{1}{2}} \int f(e^{tB}x - y) \exp\left[-\frac{1}{2}\langle Q_t^{-1}y, y\rangle\right] dy \\ &= \left(\frac{\det Q_\infty}{\det Q_t}\right)^{1/2} \exp\left[\frac{1}{2}\langle Q_t^{-1}e^{tB}x, D_t x - e^{tB}x\rangle\right] \\ &\quad \times \int f(u) \exp\left[\frac{1}{2}\langle (Q_\infty^{-1} - Q_t^{-1})(u - D_t x), u - D_t x\rangle\right] d\gamma_\infty(u),\end{aligned}\quad (2.4)$$

where we repeatedly used the fact that $Q_\infty^{-1} - Q_t^{-1}$ is symmetric. We now express the matrix D_t in various ways.

Lemma 2.1 *For all $x \in \mathbb{R}^n$ and $t > 0$ we have*

- (i) $D_t = Q_\infty e^{-tB^*} Q_\infty^{-1}$;
- (ii) $D_t = e^{tB} + Q_t e^{-tB^*} Q_\infty^{-1}$.

Proof (i) The formulae (2.1) and (1.1) imply

$$Q_\infty - Q_t = e^{tB} Q_\infty e^{tB^*} \quad (2.5)$$

(see also [12, formula (2.1)]). From (2.3) and (2.2) it follows that

$$D_t = Q_\infty (Q_\infty - Q_t)^{-1} e^{tB},$$

and combining this with (2.5) we arrive at (i).

- (ii) Multiplying (2.5) by $e^{-tB^*} Q_\infty^{-1}$ from the right, we obtain

$$Q_\infty e^{-tB^*} Q_\infty^{-1} - Q_t e^{-tB^*} Q_\infty^{-1} = e^{tB},$$

and (ii) now follows from (i). □

By means of (i) in this lemma, we can define D_t for all $t \in \mathbb{R}$, and they will form a one-parameter group of matrices.

Now (ii) in Lemma 2.1 yields

$$\langle Q_t^{-1} e^{tB} x, D_t x - e^{tB} x \rangle = \langle Q_t^{-1} e^{tB} x, Q_t e^{-tB^*} Q_\infty^{-1} x \rangle = \langle Q_\infty^{-1} x, x \rangle.$$

Thus (2.4) may be rewritten as

$$\mathcal{H}_t f(x) = \int K_t(x, u) f(u) d\gamma_\infty(u),$$

where K_t denotes the Mehler kernel, given by

$$\begin{aligned}K_t(x, u) &= \left(\frac{\det Q_\infty}{\det Q_t}\right)^{1/2} \exp(R(x)) \exp\left[-\frac{1}{2}\langle (Q_t^{-1} - Q_\infty^{-1})(u - D_t x), u - D_t x\rangle\right]\end{aligned}\quad (2.6)$$

for $x, u \in \mathbb{R}^n$. Here we introduced the quadratic form

$$R(x) = \frac{1}{2} \langle Q_\infty^{-1} x, x \rangle, \quad x \in \mathbb{R}^n.$$

3 Some auxiliary results

In this section we collect some preliminary bounds, which will be essential for the sequel.

Lemma 3.1 For $s > 0$ and for all $x \in \mathbb{R}^n$ the matrices D_s and $D_{-s} = D_s^{-1}$ satisfy

$$e^{Cs}|x| \lesssim |D_s x| \lesssim e^{Cs}|x|,$$

and

$$e^{-Cs}|x| \lesssim |D_{-s} x| \lesssim e^{-Cs}|x|.$$

This also holds with D_s replaced by e^{-sB} and e^{-sB^*} .

Proof We make a Jordan decomposition of B^* , thus writing it as the sum of a complex diagonal matrix and a triangular, nilpotent matrix, which commute with each other. This leads to expressions for e^{-sB^*} and e^{sB^*} , and since B^* like B has only eigenvalues with negative real parts, we see that

$$\|e^{-sB^*}\| \lesssim e^{Cs} \quad \text{and} \quad \|e^{sB^*}\| \lesssim e^{-Cs}. \quad (3.1)$$

From (i) in Lemma 2.1, we now get the claimed upper estimates for $D_{\pm s}$. To prove the lower estimate for D_s , we write

$$|x| = |D_{-s} D_s x| \lesssim e^{-Cs} |D_s x|.$$

The other parts of the lemma are completely analogous. $\square$

In the following lemma, we collect estimates of some basic quantities related to the matrices Q_t .

Lemma 3.2 For all $t > 0$ we have

- (i) $\det Q_t \simeq (\min(1, t))^n$;
- (ii) $\|Q_t^{-1}\| \simeq (\min(1, t))^{-1}$;
- (iii) $\|Q_\infty - Q_t\| \lesssim e^{-ct}$;
- (iv) $\|Q_t^{-1} - Q_\infty^{-1}\| \lesssim t^{-1} e^{-ct}$;
- (v) $\|(Q_t^{-1} - Q_\infty^{-1})^{-1/2}\| \lesssim t^{1/2} e^{Ct}$.

Proof (i) and (ii) Using (3.1), we see that for each $t > 0$ and for all $v \in \mathbb{R}^n$

$$\begin{aligned} \langle Q_t v, v \rangle &= \left\langle \int_0^t e^{sB} Q e^{sB^*} v \, ds, v \right\rangle = \int_0^t \langle Q^{1/2} e^{sB^*} v, Q^{1/2} e^{sB^*} v \rangle \, ds \\ &= \int_0^t |Q^{1/2} e^{sB^*} v|^2 \, ds \simeq \int_0^t |e^{sB^*} v|^2 \, ds \\ &\lesssim \int_0^t e^{-Cs} \, ds |v|^2 \simeq \min(1, t) |v|^2. \end{aligned}$$

Since $\|(e^{sB^*})^{-1}\| = \|e^{-sB^*}\| \lesssim e^{Cs}$, there is also a lower estimate

$$\int_0^t |e^{sB^*} v|^2 \, ds \gtrsim \int_0^t e^{-Cs} \, ds |v|^2 \simeq \min(1, t) |v|^2.$$

Thus any eigenvalue of Q_t has order of magnitude $\min(1, t)$, and (i) and (ii) follow.

(iii) From the definition of Q_t and (3.1), we get

$$\|Q_\infty - Q_t\| = \left\| \int_t^\infty e^{sB} Q e^{sB^*} ds \right\| \lesssim e^{-ct}.$$

(iv) Using now (ii) and (iii), we have

$$\begin{aligned} \|Q_t^{-1} - Q_\infty^{-1}\| &= \|Q_t^{-1}(Q_\infty - Q_t)Q_\infty^{-1}\| \lesssim \|Q_t^{-1}\| \|Q_\infty - Q_t\| \\ &\lesssim (\min(1, t))^{-1} e^{-ct} \lesssim t^{-1} e^{-ct}. \end{aligned}$$

(v) Since $\|A^{1/2}\| = \|A\|^{1/2}$ for any symmetric positive definite matrix A , we consider $(Q_t^{-1} - Q_\infty^{-1})^{-1}$, which can be rewritten as

$$(Q_t^{-1} - Q_\infty^{-1})^{-1} = (Q_\infty^{-1}(Q_\infty - Q_t)Q_t^{-1})^{-1} = Q_t(Q_\infty - Q_t)^{-1}Q_\infty. \quad (3.2)$$

It follows from (2.5) that $(Q_\infty - Q_t)^{-1} = e^{-tB^*} Q_\infty^{-1} e^{-tB}$, so that

$$\|(Q_\infty - Q_t)^{-1}\| \lesssim e^{Ct}$$

as a consequence of (3.2). Inserting this and the simple estimate $\|Q_t\| \lesssim t$ in (3.2), we obtain $\|(Q_t^{-1} - Q_\infty^{-1})^{-1}\| \lesssim t e^{Ct}$, and (v) follows. $\square$

Proposition 3.3 For $t \geq 1$ and $w \in \mathbb{R}^n$, we have

$$\langle (Q_t^{-1} - Q_\infty^{-1})D_t w, D_t w \rangle \simeq |w|^2.$$

Proof By (2.3) and Lemma 2.1 (i) we have

$$\begin{aligned} \langle (Q_t^{-1} - Q_\infty^{-1})D_t w, D_t w \rangle &= \langle Q_t^{-1} e^{tB} w, Q_\infty e^{-tB^*} Q_\infty^{-1} w \rangle \\ &= \langle Q_\infty Q_t^{-1} e^{tB} w, e^{-tB^*} Q_\infty^{-1} w \rangle. \end{aligned}$$

Since $Q_\infty Q_t^{-1} = I + (Q_\infty - Q_t)Q_t^{-1}$, this leads to

$$\begin{aligned} &\langle (Q_t^{-1} - Q_\infty^{-1})D_t w, D_t w \rangle \\ &= \langle e^{tB} w, e^{-tB^*} Q_\infty^{-1} w \rangle + \langle (Q_\infty - Q_t)Q_t^{-1} e^{tB} w, e^{-tB^*} Q_\infty^{-1} w \rangle \\ &= \langle Q_\infty^{-1} w, w \rangle + \langle e^{-tB} (Q_\infty - Q_t)Q_t^{-1} e^{tB} w, Q_\infty^{-1} w \rangle. \end{aligned}$$

Here $\langle Q_\infty^{-1} w, w \rangle \simeq |w|^2$. Using (2.1) and then the definition of Q_∞ , we observe that the last term can be written as

$$\begin{aligned} &\left\langle \int_t^\infty e^{(s-t)B} Q e^{(s-t)B^*} ds e^{tB^*} Q_t^{-1} e^{tB} w, Q_\infty^{-1} w \right\rangle \\ &= \langle Q_\infty e^{tB^*} Q_t^{-1} e^{tB} w, Q_\infty^{-1} w \rangle \\ &= \langle e^{tB^*} Q_t^{-1} e^{tB} w, w \rangle \\ &= |Q_t^{-1/2} e^{tB} w|^2. \end{aligned} \quad (3.3)$$

Since $|Q_t^{-1/2} e^{tB} w|^2 \lesssim |w|^2$ for $t \geq 1$ by Lemmata 3.1 and 3.2(ii), the proposition follows. $\square$

We finally give estimates of the kernel K_t , for small and large values of t . When $t \leq 1$, one has $\|(Q_t^{-1} - Q_\infty^{-1})^{1/2}\| \simeq t^{-1/2}$ and $\|(Q_t^{-1} - Q_\infty^{-1})^{-1/2}\| \simeq t^{1/2}$, by (iv) and (v) in Lemma 3.2. Combined with (2.6), this implies

$$\frac{e^{R(x)}}{t^{n/2}} \exp\left(-C \frac{|u - D_t x|^2}{t}\right) \lesssim K_t(x, u) \lesssim \frac{e^{R(x)}}{t^{n/2}} \exp\left(-c \frac{|u - D_t x|^2}{t}\right), \quad 0 < t \leq 1. \quad (3.4)$$

Lemma 3.4 For $t \geq 1$ and $x, u \in \mathbb{R}^n$, we have

$$e^{R(x)} \exp\left[-C |D_{-t} u - x|^2\right] \lesssim K_t(x, u) \lesssim e^{R(x)} \exp\left[-c |D_{-t} u - x|^2\right]. \quad (3.5)$$

Proof This follows from (2.6), if we write $u - D_t x = D_t(D_{-t} u - x)$ and apply Proposition 3.3 with $w = D_{-t} u - x$. $\square$

4 Geometric aspects of the problem

4.1 A system of adapted polar coordinates

We first need a technical lemma.

Lemma 4.1 For all $x \in \mathbb{R}^n$ and $s \in \mathbb{R}$, we have

$$\langle B^* Q_\infty^{-1} x, x \rangle = -\frac{1}{2} |Q^{1/2} Q_\infty^{-1} x|^2; \quad (4.1)$$

$$\frac{\partial}{\partial s} D_s x = -Q_\infty B^* Q_\infty^{-1} D_s x = -Q_\infty e^{-sB^*} B^* Q_\infty^{-1} x; \quad (4.2)$$

$$\frac{\partial}{\partial s} R(D_s x) = \frac{1}{2} |Q^{1/2} Q_\infty^{-1} D_s x|^2 \simeq |D_s x|^2. \quad (4.3)$$

Proof To prove (4.1), we use the definition of Q_∞ to write for any $z \in \mathbb{R}^n$

$$\begin{aligned} \langle B^* z, Q_\infty z \rangle &= \int_0^\infty \langle B^* z, e^{sB} Q e^{sB^*} z \rangle ds \\ &= \int_0^\infty \langle e^{sB^*} B^* z, Q e^{sB^*} z \rangle ds \\ &= \frac{1}{2} \int_0^\infty \frac{d}{ds} \langle e^{sB^*} z, Q e^{sB^*} z \rangle ds \\ &= -\frac{1}{2} |Q^{1/2} z|^2. \end{aligned}$$

Setting $z = Q_\infty^{-1} x$, we get (4.1).

Further, (4.2) easily follows if we observe that

$$\frac{\partial}{\partial s} D_s x = \frac{\partial}{\partial s} \left(Q_\infty e^{-sB^*} Q_\infty^{-1} x \right) = -Q_\infty B^* Q_\infty^{-1} Q_\infty e^{-sB^*} Q_\infty^{-1} x = -Q_\infty B^* Q_\infty^{-1} D_s x.$$

Finally, we get by means of (4.2) and (4.1)

$$\begin{aligned} \frac{\partial}{\partial s} R(D_s x) &= \frac{1}{2} \frac{\partial}{\partial s} \langle Q_\infty^{-1/2} D_s x, Q_\infty^{-1/2} D_s x \rangle \\ &= -\langle Q_\infty^{-1/2} Q_\infty B^* Q_\infty^{-1} D_s x, Q_\infty^{-1/2} D_s x \rangle \\ &= \frac{1}{2} |Q^{1/2} Q_\infty^{-1} D_s x|^2, \end{aligned}$$

and (4.3) is verified. $\square$

We observe here that an integration of (4.2) leads to

$$|x - D_t x| \lesssim t |x|, \quad 0 \leq t \leq 1. \quad (4.4)$$

Fix now $\beta > 0$ and consider the ellipsoid

$$E_\beta = \{x \in \mathbb{R}^n : R(x) = \beta\}.$$

As a consequence of (4.3), the map $s \mapsto R(D_s z)$ is strictly increasing for each $0 \neq z \in \mathbb{R}^n$. Hence any $x \in \mathbb{R}^n$, $x \neq 0$, can be written uniquely as

$$x = D_s \tilde{x}, \quad (4.5)$$

for some $\tilde{x} \in E_\beta$ and $s \in \mathbb{R}$. We consider s and $\tilde{x}$ as the polar coordinates of x . Our estimates in what follows will be uniform in β .

Next, we shall write Lebesgue measure in terms of these polar coordinates. A normal vector to the surface E_β at the point $\tilde{x} \in E_\beta$ is $\mathbf{N}(\tilde{x}) = Q_\infty^{-1} \tilde{x}$, and the tangent hyperplane at $\tilde{x}$ is $\mathbf{N}(\tilde{x})^\perp$. For $s > 0$ the tangent hyperplane of the surface $D_s E_\beta = \{D_s \tilde{x} : \tilde{x} \in E_\beta\}$ at the point $D_s \tilde{x}$ is $D_s(\mathbf{N}(\tilde{x})^\perp)$, and a normal to $D_s E_\beta$ at the same point is $w = (D_s^{-1})^*(\mathbf{N}(\tilde{x})) = D_{-s}^* Q_\infty^{-1} \tilde{x} = Q_\infty^{-1} e^{sB} \tilde{x}$.

The scalar product of w and the tangent of the curve $s \mapsto D_s \tilde{x}$ at the point $D_s \tilde{x}$ is, because of (4.2) and (4.1),

$$\begin{aligned} & \left\langle \frac{\partial}{\partial s} D_s \tilde{x}, w \right\rangle \\ &= -\langle Q_\infty e^{-sB^*} B^* Q_\infty^{-1} \tilde{x}, Q_\infty^{-1} e^{sB} \tilde{x} \rangle = -\langle B^* Q_\infty^{-1} \tilde{x}, \tilde{x} \rangle = \frac{1}{2} |Q_\infty^{1/2} Q_\infty^{-1} \tilde{x}|^2 > 0. \end{aligned} \quad (4.6)$$

Thus the curve $s \mapsto D_s \tilde{x}$ is transversal to each surface $D_s E_\beta$. Let dS_s denote the area measure of $D_s E_\beta$. Then Lebesgue measure is given in terms of our polar coordinates by

$$dx = H(s, \tilde{x}) dS_s(D_s \tilde{x}) ds, \quad (4.7)$$

where

$$H(s, \tilde{x}) = \left\langle \frac{\partial}{\partial s} D_s \tilde{x}, \frac{w}{|w|} \right\rangle = \frac{|Q_\infty^{1/2} Q_\infty^{-1} \tilde{x}|^2}{2 |Q_\infty^{-1} e^{sB} \tilde{x}|}.$$

To see how dS_s varies with s , we take a continuous function $\varphi = \varphi(\tilde{x})$ on E_β and extend it to $\mathbb{R}^n \setminus \{0\}$ by writing $\varphi(D_s \tilde{x}) = \varphi(\tilde{x})$. For any $t > 0$ and small $\varepsilon > 0$, we define the shell

$$\Omega_{t,\varepsilon} = \{D_s \tilde{x} : t < s < t + \varepsilon, \tilde{x} \in E_\beta\}.$$

Then $\Omega_{t,\varepsilon}$ is the image under D_t of $\Omega_{0,\varepsilon}$, and the Jacobian of this map is $\det D_t = e^{-t \operatorname{tr} B}$. Thus

$$\int_{\Omega_{t,\varepsilon}} \varphi(x) dx = e^{-t \operatorname{tr} B} \int_{\Omega_{0,\varepsilon}} \varphi(D_t x) dx,$$

which we can rewrite as

$$\begin{aligned} & \int_{t < s < t + \varepsilon} \int_{\tilde{x} \in E_\beta} \varphi(\tilde{x}) H(s, \tilde{x}) dS_s(D_s \tilde{x}) ds \\ &= e^{-t \operatorname{tr} B} \int_{0 < s < \varepsilon} \int_{\tilde{x} \in E_\beta} \varphi(\tilde{x}) H(s, \tilde{x}) dS_s(D_s \tilde{x}) ds. \end{aligned}$$

Now we divide by ε and let $\varepsilon \rightarrow 0$, getting

$$\int_{E_\beta} \varphi(\tilde{x}) H(t, \tilde{x}) dS_t(D_t \tilde{x}) = e^{-t \operatorname{tr} B} \int_{E_\beta} \varphi(\tilde{x}) H(0, \tilde{x}) dS_0(\tilde{x}).$$

Since this holds for any φ , it follows that

$$dS_t(D_t \tilde{x}) = e^{-t \operatorname{tr} B} \frac{H(0, \tilde{x})}{H(t, \tilde{x})} dS_0(\tilde{x}).$$

Together with (4.7), this implies the following result.

Proposition 4.2 *The Lebesgue measure in $\mathbb{R}^n$ is given in terms of polar coordinates $(t, \tilde{x})$ by*

$$dx = e^{-t \operatorname{tr} B} \frac{|Q^{1/2} Q_\infty^{-1} \tilde{x}|^2}{2 |Q_\infty^{-1} \tilde{x}|} dS_0(\tilde{x}) dt.$$

We also need estimates of the distance between two points in terms of the polar coordinates. The following result is a generalization of Lemma 4.2 in [4], and its proof is analogous.

Lemma 4.3 *Fix $\beta > 0$. Let $x^{(0)}, x^{(1)} \in \mathbb{R}^n \setminus \{0\}$ and assume $R(x^{(0)}) > \beta/2$. Write*

$$x^{(0)} = D_{s^{(0)}}(\tilde{x}^{(0)}) \quad \text{and} \quad x^{(1)} = D_{s^{(1)}}(\tilde{x}^{(1)})$$

with $s^{(0)}, s^{(1)} \in \mathbb{R}$ and $\tilde{x}^{(0)}, \tilde{x}^{(1)} \in E_\beta$.

(i) *Then*

$$|x^{(0)} - x^{(1)}| \gtrsim c |\tilde{x}^{(0)} - \tilde{x}^{(1)}|. \quad (4.8)$$

(ii) *If also $s^{(1)} \geq 0$, then*

$$|x^{(0)} - x^{(1)}| \gtrsim c \sqrt{\beta} |s^{(0)} - s^{(1)}|. \quad (4.9)$$

Proof Let $\Gamma : [0, 1] \rightarrow \mathbb{R}^n \setminus \{0\}$ be a differentiable curve with $\Gamma(0) = x^{(0)}$ and $\Gamma(1) = x^{(1)}$. It suffices to bound the length of any such curve from below by the right-hand sides of (4.8) and (4.9).

For each $\tau \in [0, 1]$, we write

$$\Gamma(\tau) = D_{s(\tau)} \tilde{x}(\tau),$$

with $\tilde{x}(\tau) \in E_\beta$ and $\tilde{x}(i) = \tilde{x}^{(i)}, s(i) = s^{(i)}$ for $i = 0, 1$. Thus

$$\Gamma'(\tau) = -s'(\tau) \frac{\partial}{\partial s} D_s \Big|_{s=s(\tau)} \tilde{x}(\tau) + D_{s(\tau)} \tilde{x}'(\tau).$$

The group property of D_s implies that

$$\frac{\partial}{\partial s} D_s \Big|_{s=s(\tau)} = D_{s(\tau)} \frac{\partial}{\partial s} D_s \Big|_{s=0},$$

and so

$$\Gamma'(\tau) = D_{s(\tau)} v,$$

with

$$v = -s'(\tau) \frac{\partial}{\partial s} D_s \Big|_{s=0} \tilde{x}(\tau) + \tilde{x}'(\tau).$$

The vector $\tilde{x}'(\tau)$ is tangent to E_β and thus orthogonal to $\mathbf{N}(\tilde{x})$. Then (4.6) (with $s = 0$) implies that the angle between $\frac{\partial}{\partial s} D_s|_{s=0} \tilde{x}(\tau)$ and $\tilde{x}'(\tau)$ is larger than some positive constant. It follows that

$$|v|^2 \gtrsim |s'(\tau)|^2 \left| \frac{\partial}{\partial s} D_s|_{s=0} \tilde{x}(\tau) \right|^2 + |\tilde{x}'(\tau)|^2 \gtrsim |s'(\tau)|^2 \beta + |\tilde{x}'(\tau)|^2, \quad (4.10)$$

where we also used the fact that, by (4.2),

$$\left| \frac{\partial}{\partial s} D_s|_{s=0} \tilde{x}(\tau) \right| \simeq |\tilde{x}(\tau)| \simeq \sqrt{\beta}.$$

Since

$$|v| = |D_{-s(\tau)} \Gamma'(\tau)| \leq \|D_{-s(\tau)}\| |\Gamma'(\tau)| \lesssim e^{-C \min(s(\tau), 0)} |\Gamma'(\tau)|$$

because of Lemma 3.1, we obtain from (4.10)

$$|\Gamma'(\tau)| \gtrsim e^{C \min(s(\tau), 0)} (\sqrt{\beta} |s'(\tau)| + |\tilde{x}'(\tau)|). \quad (4.11)$$

Next, we derive a lower bound for $s(0)$; assume first that $s(0) < 0$. The assumption $R(x^{(0)}) > \beta/2$ implies, together with Lemma 3.1,

$$\beta/2 \leq R(D_{s(0)} \tilde{x}^{(0)}) \lesssim |D_{s(0)} \tilde{x}^{(0)}|^2 \lesssim e^{c s(0)} |\tilde{x}^{(0)}|^2 \simeq e^{c s(0)} \beta.$$

It follows that

$$s(0) > -\tilde{s},$$

for some $\tilde{s}$ with $0 < \tilde{s} < C$, and this obviously holds also without the assumption $s(0) < 0$.

Assume now that $s(\tau) > -\tilde{s} - 1$ for all $\tau \in [0, 1]$. Then (4.11) implies

$$|\Gamma'(\tau)| \gtrsim \sqrt{\beta} |s'(\tau)|$$

and

$$|\Gamma'(\tau)| \gtrsim |\tilde{x}'(\tau)|.$$

Integrating these estimates with respect to τ in $[0, 1]$, we immediately see that one can control the length of Γ from below by the right-hand sides of (4.8) and (4.9).

If instead $s(\tau) \leq -\tilde{s} - 1$ for some $\tau \in [0, 1]$, we can proceed as in the proof of Lemma 4.2 in [4]. More precisely, since the image $s([0, 1])$ contains the interval $[-\tilde{s} - 1, \max(s(0), s(1))]$, we can find a closed subinterval I of $[0, 1]$ whose image $s(I)$ is exactly the interval $[-\tilde{s} - 1, \max(s(0), s(1))]$. Thus we may use (4.11) to control the length of Γ by

$$\int_0^1 |\Gamma'(\tau)| d\tau \geq \int_I |\Gamma'(\tau)| d\tau \gtrsim \sqrt{\beta} \int_I |s'(\tau)| d\tau \geq \sqrt{\beta} (\max(s(0), s(1)) + \tilde{s} + 1).$$

Here

$$\sqrt{\beta} (\max(s(0), s(1)) + \tilde{s} + 1) \gtrsim \sqrt{\beta} \gtrsim \text{diam } E_\beta \geq |\tilde{x}^{(0)} - \tilde{x}^{(1)}|,$$

and (4.8) follows. Under the additional hypothesis $s(1) \geq 0$ of (ii), we have

$$\tilde{s} \geq \max(-s(0), -s(1)) = -\min(s(0), s(1)).$$

Then

$$\begin{aligned}\sqrt{\beta} \left(\max(s(0), s(1)) + \tilde{s} + 1 \right) &\gtrsim \sqrt{\beta} \left(\max(s(0), s(1)) - \min(s(0), s(1)) \right) \\ &= \sqrt{\beta} |s(0) - s(1)|,\end{aligned}$$

and (4.9) follows. $\square$

4.2 The Gaussian measure of a tube

We fix a large $\beta > 0$. Define for $x^{(1)} \in E_\beta$ and $a > 0$ the set

$$\Omega = \left\{ x \in E_\beta : |x - x^{(1)}| < a \right\}.$$

This is a spherical cap of the ellipsoid E_β , centered at $x^{(1)}$. Observe that $|x| \simeq \sqrt{\beta}$ for $x \in \Omega$, and that the area of Ω is $|\Omega| \simeq \min(a^{n-1}, \beta^{(n-1)/2})$. Then consider the tube

$$Z = \{D_s \tilde{x} : s \geq 0, \tilde{x} \in \Omega\}. \quad (4.12)$$

Lemma 4.4 *There exists a constant C such that $\beta > C$ implies that the Gaussian measure of the tube Z fulfills*

$$\gamma_\infty(Z) \lesssim \frac{a^{n-1}}{\sqrt{\beta}} e^{-\beta}.$$

Proof Proposition 4.2 yields, since $H(0, \tilde{x}) \simeq |\tilde{x}| \simeq \sqrt{\beta}$,

$$\gamma_\infty(Z) \simeq \int_0^\infty e^{-s \operatorname{tr} B} e^{-R(D_s \tilde{x})} \int_\Omega H(0, \tilde{x}) dS(\tilde{x}) ds \lesssim \sqrt{\beta} a^{n-1} \int_0^\infty e^{-s \operatorname{tr} B} e^{-R(D_s \tilde{x})} ds.$$

By (4.3) we have

$$R(D_s \tilde{x}) - R(\tilde{x}) \simeq \int_0^s |D_{s'} \tilde{x}|^2 ds' \gtrsim s |\tilde{x}|^2 \simeq s\beta,$$

which implies

$$\gamma_\infty(Z) \lesssim \sqrt{\beta} a^{n-1} e^{-\beta} \int_0^\infty e^{-s \operatorname{tr} B} e^{-cs\beta} ds.$$

Assuming β large enough, one has $c\beta > -2 \operatorname{tr} B$, and then the last integral is finite and no larger than C/β . The lemma follows. $\square$

5 Simplifications

In this section, we introduce some preliminary simplifications and reductions for the proof of (1.3), i.e., of Theorem 1.1.

(1) We may assume that f is nonnegative and normalized in the sense that

$$\|f\|_{L^1(\gamma_\infty)} = 1,$$

since this involves no loss of generality.

(2) We may assume that α is large, $\alpha > C$, since otherwise (1.3) and (1.4) are trivial.

(3) In many cases, we may restrict x in (1.3) and (1.4) to the ellipsoidal annulus

$$\mathcal{E}_\alpha = \left\{ x \in \mathbb{R}^n : \frac{1}{2} \log \alpha \leq R(x) \leq 2 \log \alpha \right\}.$$

To begin with, we can always forget the unbounded component of the complement of $\mathcal{E}_\alpha$, since

$$\gamma_\infty \{x \in \mathbb{R}^n : R(x) > 2 \log \alpha\} \lesssim \int_{R(x) > 2 \log \alpha} \exp(-R(x)) dx \lesssim (\log \alpha)^{(n-2)/2} \exp(-2 \log \alpha) \lesssim \frac{1}{\alpha}. \quad (5.1)$$

(4) When $t > 1$, we may forget also the inner region where $R(x) < \frac{1}{2} \log \alpha$. Indeed, from (3.5) we get, if $(x, u) \in \mathbb{R}^n \times \mathbb{R}^n$ with $R(x) < \frac{1}{2} \log \alpha$,

$$K_t(x, u) \lesssim e^{R(x)} < \sqrt{\alpha} < \alpha,$$

since α is large. In other words, for any $(x, u) \in \mathbb{R}^n \times \mathbb{R}^n$

$$R(x) < \frac{1}{2} \log \alpha \quad \Rightarrow \quad K_t(x, u) \lesssim \alpha, \quad (5.2)$$

for all $t > 1$.

Replacing α by $C\alpha$ for some C , we see from (3) and (4) that we can assume $x \in \mathcal{E}_\alpha$ in the proof of (1.3) and (1.4), when the supremum in the maximal operator is taken only over $t > 1$.

Before introducing the last simplification, we need to define a global region

$$G = \left\{ (x, u) \in \mathbb{R}^n \times \mathbb{R}^n : |x - u| > \frac{1}{1 + |x|} \right\}$$

and a local region

$$L = \left\{ (x, u) \in \mathbb{R}^n \times \mathbb{R}^n : |x - u| \leq \frac{1}{1 + |x|} \right\}.$$

Notice that the definition of G and L does not depend on Q and B .

(5) When $t \leq 1$ and $(x, u) \in G$, we shall see that (5.2) is still valid, and it is again enough to consider $x \in \mathcal{E}_\alpha$.

To prove this, we need a lemma which will also be useful later.

Lemma 5.1 *If $(x, u) \in G$ and $0 < t \leq 1$, then*

$$\frac{1}{(1 + |x|)^2} \lesssim t^2 |x|^2 + |u - D_t x|^2.$$

Proof From the definition of G and (4.4) we get

$$\frac{1}{1 + |x|} \leq |x - u| \leq |x - D_t x| + |D_t x - u| \lesssim t|x| + |u - D_t x|.$$

The lemma follows. $\square$

To verify now (5.2) in the global region with $t \leq 1$, we recall from (3.4) that

$$K_t(x, u) \lesssim \frac{e^{R(x)}}{t^{n/2}} \exp\left(-c \frac{|u - D_t x|^2}{t}\right).$$

It follows from Lemma 5.1 that

$$t^2 \gtrsim \frac{1}{(1 + |x|)^4} \quad \text{or} \quad \frac{|u - D_t x|^2}{t} \gtrsim \frac{1}{(1 + |x|)^2 t}. \quad (5.3)$$

The first inequality here implies that

$$K_t(x, u) \lesssim e^{R(x)} (1 + |x|)^n \lesssim e^{2R(x)},$$

and (5.2) follows. If the second inequality of (5.3) holds, we have

$$K_t(x, u) \lesssim \frac{e^{R(x)}}{t^{n/2}} \exp\left(-\frac{c}{(1 + |x|)^2 t}\right) \lesssim e^{R(x)} (1 + |x|)^n,$$

and we get the same estimate. Thus (5.2) is verified.

Finally, let

$$\mathcal{H}_*^G f(x) = \sup_{0 < t \leq 1} \left| \int K_t(x, u) \chi_G(x, u) f(u) d\gamma_\infty(u) \right|,$$

and

$$\mathcal{H}_*^L f(x) = \sup_{0 < t \leq 1} \left| \int K_t(x, u) \chi_L(x, u) f(u) d\gamma_\infty(u) \right|.$$

6 The case of large t

In this section, we consider the supremum in the definition of the maximal operator taken only over $t > 1$, and we prove (1.4).

Proposition 6.1 *For all functions $f \in L^1(\gamma_\infty)$ such that $\|f\|_{L^1(\gamma_\infty)} = 1$,*

$$\gamma_\infty \left\{ x : \sup_{t>1} |\mathcal{H}_t f(x)| > \alpha \right\} \lesssim \frac{1}{\alpha \sqrt{\log \alpha}}, \quad \alpha > 2. \quad (6.1)$$

In particular, the maximal operator

$$\sup_{t>1} |\mathcal{H}_t f(x)|$$

is of weak type $(1, 1)$ with respect to the invariant measure γ_∞ .

Proof We can assume that $f \geq 0$. Looking at the arguments in Sect. 5, items (3) and (4), we see that it suffices to consider points $x \in \mathcal{E}_\alpha$. For both x and u we use the coordinates introduced in (4.5) with $\beta = \log \alpha$, that is,

$$x = D_s \tilde{x}, \quad u = D_{s'} \tilde{u},$$

where $\tilde{x}, \tilde{u} \in E_{\log \alpha}$ and $s, s' \in \mathbb{R}$.

From (3.5) we have

$$K_t(x, u) \lesssim \exp(R(x)) \exp(-c |D_{-t} u - x|^2)$$

for $t > 1$ and $x, u \in \mathbb{R}^n$. Since $x \in \mathcal{E}_\alpha$ and $D_{-t} u = D_{s'-t} \tilde{u}$, we can apply Lemma 4.3 (i), getting

$$|D_{-t} u - x| \gtrsim |\tilde{x} - \tilde{u}|,$$

so that

$$\int K_t(x, u) f(u) d\gamma_\infty(u) \lesssim \exp(R(D_s \tilde{x})) \int \exp(-c |\tilde{x} - \tilde{u}|^2) f(u) d\gamma_\infty(u).$$

In view of (4.3), the right-hand side here is strictly increasing in s , and therefore the inequality

$$\exp(R(D_s \tilde{x})) \int \exp(-c |\tilde{x} - \tilde{u}|^2) f(u) d\gamma_\infty(u) > \alpha \quad (6.2)$$

holds if and only if $s > s_\alpha(\tilde{x})$ for some function $\tilde{x} \mapsto s_\alpha(\tilde{x})$, with equality for $s = s_\alpha(\tilde{x})$. Since $\alpha > 2$ and $\|f\|_{L^1(\gamma_\infty)} = 1$, it follows that $s_\alpha(\tilde{x}) > 0$.

For some C , the set of points $x \in \mathcal{E}_\alpha$ where the supremum in (6.1) is larger than $C\alpha$ is contained in the set $\mathcal{A}(\alpha)$ of points $D_s \tilde{x} \in \mathcal{E}_\alpha$ fulfilling (6.2). We use Proposition 4.2 to estimate the γ_∞ measure of $\mathcal{A}(\alpha)$. Observe that $H(0, \tilde{x}) \simeq |\tilde{x}| \simeq \sqrt{\log \alpha}$ and that $D_s \tilde{x} \in \mathcal{E}_\alpha$ implies $s \lesssim 1$, so that also $e^{-s \operatorname{tr} B} \lesssim 1$. We get

$$\begin{aligned} \gamma_\infty(\mathcal{A}(\alpha)) &= \int_{\mathcal{A}(\alpha) \cap \mathcal{E}_\alpha} e^{-R(x)} dx \\ &\lesssim \sqrt{\log \alpha} \int_{E_{\log \alpha}} \int_{s_\alpha(\tilde{x})}^C e^{-R(D_s \tilde{x})} ds dS(\tilde{x}) \\ &\lesssim \sqrt{\log \alpha} \int_{E_{\log \alpha}} \int_{s_\alpha(\tilde{x})}^{+\infty} \exp(-R(D_{s_\alpha(\tilde{x})} \tilde{x}) - c \log \alpha (s - s_\alpha(\tilde{x}))) ds dS(\tilde{x}), \end{aligned}$$

where the last inequality follows from (4.3), since $|D_s \tilde{x}|^2 \gtrsim |\tilde{x}|^2 \simeq \log \alpha$. Integrating in s , we obtain

$$\gamma_\infty(\mathcal{A}(\alpha)) \lesssim \frac{1}{\sqrt{\log \alpha}} \int_{E_{\log \alpha}} \exp(-R(D_{s_\alpha(\tilde{x})} \tilde{x})) dS(\tilde{x}).$$

Now combine this estimate with the case of equality in (6.2) and change the order of integration, to get

$$\begin{aligned} \gamma_\infty(\mathcal{A}(\alpha)) &\lesssim \frac{1}{\alpha \sqrt{\log \alpha}} \int \int_{E_{\log \alpha}} \exp(-c |\tilde{x} - \tilde{u}|^2) dS(\tilde{x}) f(u) d\gamma_\infty(u) \\ &\lesssim \frac{1}{\alpha \sqrt{\log \alpha}} \int f(u) d\gamma_\infty(u), \end{aligned}$$

which proves Proposition 6.1. $\square$

Finally, we show that the factor $1/\sqrt{\log \alpha}$ in (6.1) is sharp.

Proposition 6.2 *For any $t > 1$ and any large α , there exists a function f normalized in $L^1(\gamma_\infty)$ and such that*

$$\gamma_\infty \{x : |\mathcal{H}_t f(x)| > \alpha\} \simeq \frac{1}{\alpha \sqrt{\log \alpha}}.$$

Proof Take a point z with $R(z) = \log \alpha$, and let f be (an approximation of) a Dirac measure at the point $u = D_t z$. Then, as a consequence of (3.5), $K_t(x, u) \simeq \exp(R(x))$ when x is in the ball $B(D_{-t} u, 1) = B(z, 1)$. We then have $\mathcal{H}_t f(x) = K_t(x, u) \gtrsim \alpha$ in the set $\mathcal{B} = \{x \in B(z, 1) : R(x) > R(z)\}$, whose measure is

$$\gamma_\infty(\mathcal{B}) \simeq e^{-R(z)} \frac{1}{\sqrt{R(z)}} = \frac{1}{\alpha \sqrt{\log \alpha}}.$$

□

7 The local case for small t

Proposition 7.1 *If $(x, u) \in L$ and $0 < t \leq 1$, then*

$$|K_t(x, u)| \lesssim \frac{\exp(R(x))}{t^{n/2}} \exp\left(-c \frac{|u - x|^2}{t}\right).$$

Proof In view of (3.4), it is enough to show that

$$\frac{|u - D_t x|^2}{t} \geq \frac{|u - x|^2}{t} - C. \quad (7.1)$$

We write

$$\begin{aligned} |u - D_t x|^2 &= |u - x + x - D_t x|^2 = |u - x|^2 + 2\langle u - x, x - D_t x \rangle + |x - D_t x|^2 \\ &\geq |u - x|^2 - 2|u - x| |x - D_t x|. \end{aligned}$$

By (4.4),

$$|u - x| |x - D_t x| \lesssim |u - x| t |x| \leq t$$

since $(x, u) \in L$, and (7.1) follows. □

Proposition 7.2 *The maximal operator $\mathcal{H}_*^L$ is of weak type $(1, 1)$ with respect to the invariant measure γ_∞ .*

Proof The proof is standard, since Proposition 7.1 implies

$$\mathcal{H}_*^L f(x) \lesssim \sup_{0 < t \leq 1} \frac{\exp(R(x))}{t^{n/2}} \int \exp\left(-c \frac{|x - u|^2}{t}\right) \chi_L(x, u) f(u) d\gamma_\infty(u).$$

The supremum here defines an operator of weak type $(1, 1)$ with respect to Lebesgue measure in $\mathbb{R}^n$. From this the proposition follows, cf. [7, Section 3]. □

8 The global case for small t

In this section, we conclude the proof of Theorem 1.1.

Proposition 8.1 *The maximal operator $\mathcal{H}_*^G$ is of weak type $(1, 1)$ with respect to the invariant measure γ_∞ .*

Proof We take f and α as in items (1) and (2) of Sect. 5. Then item (5) tells us that we need only consider $\mathcal{H}_*^G f(x)$ for $x \in \mathcal{E}_\alpha$.

For $m \in \mathbb{N}$ and $0 < t \leq 1$, we introduce regions S_t^m . If $m > 0$, we let

$$S_t^m = \left\{ (x, u) \in G : 2^{m-1}\sqrt{t} < |u - D_t x| \leq 2^m\sqrt{t} \right\}.$$

If $m = 0$, we replace the condition $2^{m-1}\sqrt{t} < |u - D_t x| \leq 2^m\sqrt{t}$ by $|u - D_t x| \leq \sqrt{t}$. Note that for any fixed $t \in (0, 1]$ these sets form a partition of G .

In the set S_t^m we have, because of (3.4),

$$K_t(x, u) \lesssim \frac{\exp(R(x))}{t^{n/2}} \exp(-c2^{2m}).$$

Then setting

$$\mathcal{K}_t^m(x, u) = \frac{\exp(R(x))}{t^{n/2}} \chi_{S_t^m}(x, u),$$

one has, for all $(x, u) \in G$ and $0 < t < 1$,

$$K_t(x, u) \lesssim \sum_{m=0}^{\infty} \exp(-c2^{2m}) \mathcal{K}_t^m(x, u).$$

Hence, it suffices to prove that for $m = 0, 1, \dots$

$$\gamma_{\infty} \left\{ x \in \mathcal{E}_{\alpha} : \sup_{0 < t \leq 1} \int \mathcal{K}_t^m(x, u) f(u) d\gamma_{\infty}(u) > \alpha \right\} \lesssim \frac{2^{Cm}}{\alpha}, \quad (8.1)$$

for large α and some C , since this will allow summing in m in the space $L^{1,\infty}(\gamma_{\infty})$.

Fix $m \in \mathbb{N}$ and assume that $(x, u) \in S_t^m$ for some $t \in (0, 1]$, so that $|u - D_t x| \leq 2^m\sqrt{t}$. Then Lemma 5.1 leads to

$$1 \lesssim (1 + |x|)^4 t^2 + (1 + |x|)^2 2^{2m} t \leq ((1 + |x|)^2 2^{2m} t)^2 + (1 + |x|)^2 2^{2m} t.$$

Consequently, a point $x \in \mathcal{E}_{\alpha}$ satisfies

$$(1 + |x|)^2 2^{2m} t \gtrsim 1 \quad (8.2)$$

as soon as there exists a point u with $\mathcal{K}_t^m(x, u) \neq 0$, and then $t \geq \varepsilon > 0$ for some $\varepsilon = \varepsilon(\alpha, m) > 0$. Hence the supremum in (8.1) will be the same if taken only over $\varepsilon \leq t \leq 1$, and it follows that this supremum is a continuous function of $x \in \mathcal{E}_{\alpha}$.

To prove (8.1), the idea, which goes back to [15], is to construct a finite sequence of pairwise disjoint balls $(\mathcal{B}^{(\ell)})_{\ell=1}^{\ell_0}$ in $\mathbb{R}^n$ and a finite sequence of sets $(\mathcal{Z}^{(\ell)})_{\ell=1}^{\ell_0}$ in $\mathbb{R}^n$, called forbidden zones. These zones will together cover the level set in (8.1). We will then verify that

$$\left\{ x \in \mathcal{E}_{\alpha} : \sup_{\varepsilon \leq t \leq 1} \int \mathcal{K}_t^m(x, u) f(u) d\gamma_{\infty}(u) \geq \alpha \right\} \subset \bigcup_{\ell=1}^{\ell_0} \mathcal{Z}^{(\ell)}, \quad (8.3)$$

that for each ℓ

$$\gamma_{\infty}(\mathcal{Z}^{(\ell)}) \lesssim \frac{2^{Cm}}{\alpha} \int_{\mathcal{B}^{(\ell)}} f(u) d\gamma_{\infty}(u), \quad (8.4)$$

and that the $\mathcal{B}^{(\ell)}$ are pairwise disjoint. This would imply

$$\gamma_{\infty}\left(\bigcup_{\ell=1}^{\ell_0} \mathcal{Z}^{(\ell)}\right) \lesssim \frac{2^{C_m}}{\alpha} \sum_{\ell=1}^{\ell_0} \int_{\mathcal{B}^{(\ell)}} f(u) d\gamma_{\infty}(u) \lesssim \frac{2^{C_m}}{\alpha},$$

and thus also (8.1) and Proposition 8.1.

The sets $\mathcal{B}^{(\ell)}$ and $\mathcal{Z}^{(\ell)}$ will be introduced by means of a sequence of points $x^{(\ell)}$, $\ell = 1, \dots, \ell_0$, which we define by recursion. To start, we choose as $x^{(1)}$ a point where the quadratic form $R(x)$ takes its minimal value in the compact set

$$\mathcal{A}_1(\alpha) = \left\{ x \in \mathcal{E}_{\alpha} : \sup_{\varepsilon \leq t \leq 1} \int \mathcal{K}_t^m(x, u) f(u) d\gamma_{\infty} \geq \alpha \right\}.$$

However, should this set be empty, (8.1) is immediate.

We now describe the recursion to construct $x^{(\ell)}$ for $\ell \geq 2$. Like $x^{(1)}$, these points will satisfy

$$\sup_{\varepsilon \leq t \leq 1} \int \mathcal{K}_t^m(x^{(\ell)}, u) f(u) d\gamma_{\infty} \geq \alpha.$$

Once an $x^{(\ell)}$, $\ell \geq 1$, is defined, we can thus by continuity choose $t_{\ell} \in [\varepsilon, 1]$ such that

$$\int \mathcal{K}_{t_{\ell}}^m(x^{(\ell)}, u) f(u) d\gamma_{\infty} \geq \alpha. \quad (8.5)$$

Using this t_{ℓ} , we associate with $x^{(\ell)}$ the tube

$$\mathcal{Z}^{(\ell)} = \left\{ D_s \eta \in \mathbb{R}^n : s \geq 0, R(\eta) = R(x^{(\ell)}), |\eta - x^{(\ell)}| < A 2^{3m} \sqrt{t_{\ell}} \right\},$$

Here the constant $A > 0$ is to be determined, depending only on n , Q and B .

All the $x^{(\ell)}$ will be minimizing points of $R(x)$. To avoid having them too close to one another, we will not allow $x^{(\ell)}$ to be in any $\mathcal{Z}^{(\ell')}$ with $\ell' < \ell$. More precisely, assuming $x^{(1)}, \dots, x^{(\ell)}$ already defined, we will choose $x^{(\ell+1)}$ as a minimizing point of $R(x)$ in the set

$$\mathcal{A}_{\ell+1}(\alpha) = \left\{ x \in \mathcal{E}_{\alpha} \setminus \bigcup_{\ell'=1}^{\ell} \mathcal{Z}^{(\ell')} : \sup_{\varepsilon \leq t \leq 1} \int \mathcal{K}_t^m(x, u) f(u) d\gamma_{\infty}(u) \geq \alpha \right\}, \quad (8.6)$$

provided this set is nonempty. But if $\mathcal{A}_{\ell+1}(\alpha)$ is empty, the process stops with $\ell_0 = \ell$ and (8.3) follows. We will see that this actually occurs for some finite ℓ .

Now assume that $\mathcal{A}_{\ell+1}(\alpha) \neq \emptyset$. In order to assure that a minimizing point exists, we must verify that $\mathcal{A}_{\ell+1}(\alpha)$ is closed and thus compact, although the $\mathcal{Z}^{(\ell')}$ are not open. To do so, observe that for $1 \leq \ell' \leq \ell$, the minimizing property of $x^{(\ell')}$ means that there is no point x in $\mathcal{A}_{\ell'}(\alpha)$ with $R(x) < R(x^{(\ell')})$. Thus we have the inclusions

$$\mathcal{A}_{\ell+1}(\alpha) \subset \mathcal{A}_{\ell'}(\alpha) \subset \left\{ x : R(x) \geq R(x^{(\ell')}) \right\}, \quad 1 \leq \ell' \leq \ell.$$

It follows that

$$\begin{aligned} \mathcal{A}_{\ell+1}(\alpha) &= \mathcal{A}_{\ell+1}(\alpha) \cap \bigcap_{1 \leq \ell' \leq \ell} \{x : R(x) \geq R(x^{(\ell')})\} \\ &= \bigcap_{\ell'=1}^{\ell} \left\{ x \in \mathcal{E}_{\alpha} \setminus \mathcal{Z}^{(\ell')} : R(x) \geq R(x^{(\ell')}), \sup_{\varepsilon \leq t \leq 1} \int \mathcal{K}_t^m(x, u) f(u) d\gamma_{\infty}(u) \geq \alpha \right\}. \end{aligned}$$

For each $\ell' = 1, \dots, \ell$ we have

$$\begin{aligned} & \{x \in \mathcal{E}_\alpha \setminus \mathcal{Z}^{(\ell')} : R(x) \geq R(x^{(\ell')})\} \\ &= \left\{ D_s \eta \in \mathcal{E}_\alpha : s \geq 0, R(\eta) = R(x^{(\ell')}), |\eta - x^{(\ell')}| \geq A 2^{3m} \sqrt{t_{\ell'}} \right\}, \end{aligned}$$

and this set is closed. It follows that $\mathcal{A}_{\ell+1}(\alpha)$ is compact, and a minimizing point $x^{(\ell+1)}$ can be chosen. Thus the recursion is well defined.

We observe that (8.2) applies to t_ℓ and $x^{(\ell)}$, and $|x^{(\ell)}|$ is large, so

$$|x^{(\ell)}|^2 2^{2m} t_\ell \gtrsim 1. \quad (8.7)$$

Further, we define balls

$$\mathcal{B}^{(\ell)} = \{u \in \mathbb{R}^n : |u - D_{t_\ell} x^{(\ell)}| \leq 2^m \sqrt{t_\ell}\}.$$

Because of the definitions of $\mathcal{K}_t^m$ and $\mathcal{S}_t^m$, the inequality (8.5) implies

$$\alpha \leq \frac{\exp(R(x^{(\ell)}))}{t_\ell^{n/2}} \int_{\mathcal{B}^{(\ell)}} f(u) d\gamma_\infty(u). \quad (8.8)$$

It remains to verify the claimed properties of $\mathcal{B}^{(\ell)}$ and $\mathcal{Z}^{(\ell)}$. The arguments below follow the lines of the proof of Lemma 6.2 in [4], with only slight modifications.

Lemma 8.2 *The balls $\mathcal{B}^{(\ell)}$ are pairwise disjoint.*

Proof Two balls $\mathcal{B}^{(\ell)}$ and $\mathcal{B}^{(\ell')}$ with $\ell < \ell'$ will be disjoint if

$$|D_{t_\ell} x^{(\ell)} - D_{t_{\ell'}} x^{(\ell')}| > 2^m (\sqrt{t_\ell} + \sqrt{t_{\ell'}}). \quad (8.9)$$

By means of our polar coordinates with $\beta = R(x^{(\ell)})$, we write

$$x^{(\ell')} = D_s \tilde{x}^{(\ell')}$$

for some $\tilde{x}^{(\ell')}$ with $R(\tilde{x}^{(\ell')}) = R(x^{(\ell)})$ and some $s \in \mathbb{R}$. Note that $s \geq 0$, because $R(x^{(\ell')}) \geq R(x^{(\ell)})$. Since $x^{(\ell')}$ does not belong to the forbidden zone $\mathcal{Z}^{(\ell)}$, we must have

$$|\tilde{x}^{(\ell')} - x^{(\ell)}| \geq A 2^{3m} \sqrt{t_\ell}. \quad (8.10)$$

We first assume that $t_{\ell'} \geq M 2^{4m} t_\ell$, for some $M = M(n, Q, B) \geq 2$ to be chosen. Lemma 4.3 (ii) implies

$$|D_{t_\ell} x^{(\ell)} - D_{t_{\ell'}} x^{(\ell')}| = |D_{t_\ell} x^{(\ell)} - D_{t_{\ell'}+s} \tilde{x}^{(\ell')}| \gtrsim |x^{(\ell)}| (t_{\ell'} + s - t_\ell) \gtrsim |x^{(\ell)}| t_{\ell'},$$

the last step by our assumption. Using again the assumption and then (8.7), we get

$$|x^{(\ell)}| t_{\ell'} \gtrsim |x^{(\ell)}| \sqrt{M} 2^{2m} \sqrt{t_\ell} \sqrt{t_{\ell'}} \gtrsim \sqrt{M} 2^m \sqrt{t_{\ell'}} \simeq \sqrt{M} 2^m (\sqrt{t_{\ell'}} + \sqrt{t_\ell}).$$

Fixing M suitably large, we obtain (8.9) from the last two formulae.

It remains to consider the case when $t_{\ell'} < M 2^{4m} t_\ell$. Then

$$\sqrt{t_\ell} > \frac{2^{-2m-1}}{\sqrt{M}} (\sqrt{t_{\ell'}} + \sqrt{t_\ell}).$$

Applying this to (8.10), we obtain (8.9) by choosing A so that $A/\sqrt{M}$ is large enough. $\square$

We next verify that the sequence $(x^{(\ell)})$ is finite. For $\ell < \ell'$, we have (8.10), and Lemma 4.3 (i) implies

$$|x^{(\ell')} - x^{(\ell)}| \gtrsim A 2^{3m} \sqrt{t_\ell}.$$

Since $t_\ell \geq \varepsilon$, we see that the distance $|x^{(\ell')} - x^{(\ell)}|$ is bounded below by a positive constant. But all the $x^{(\ell)}$ are contained in the bounded set $\mathcal{E}_\alpha$, so they are finite in number. Thus the set considered in (8.6) must be empty for some ℓ , and the recursion stops. This implies (8.3).

We finally prove (8.4). Observe that the forbidden zone $\mathcal{Z}^{(\ell)}$ is a tube as defined in (4.12), with $a = A 2^{3m} \sqrt{t_\ell}$ and $\beta = R(x^{(\ell)})$. This value of β is large since $x^{(\ell)} \in \mathcal{E}_\alpha$, and thus we can apply Lemma 4.4 to obtain

$$\gamma_\infty(\mathcal{Z}^{(\ell)}) \lesssim \frac{(A 2^{3m} \sqrt{t_\ell})^{n-1}}{\sqrt{R(x^{(\ell)})}} \exp\left(-R(x^{(\ell)})\right).$$

We bound the exponential here by means of (8.8) and observe that $R(x^{(\ell)}) \sim |x^{(\ell)}|^2$, getting

$$\gamma_\infty(\mathcal{Z}^{(\ell)}) \lesssim \frac{1}{\alpha |x^{(\ell)}| \sqrt{t_\ell}} (A 2^{3m})^{n-1} \int_{\mathcal{B}^{(\ell)}} f(u) d\gamma_\infty(u).$$

As a consequence of (8.7), we obtain

$$\gamma_\infty(\mathcal{Z}^{(\ell)}) \lesssim \frac{2^m}{\alpha} (A 2^{3m})^{n-1} \int_{\mathcal{B}^{(\ell)}} f(u) d\gamma_\infty(u) \lesssim \frac{2^{C_m}}{\alpha} \int_{\mathcal{B}^{(\ell)}} f(u) d\gamma_\infty(u),$$

proving (8.4). This concludes the proof of Proposition 8.1. $\square$

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