# Metrisability of projective surfaces and pseudo-holomorphic curves 

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#### Abstract

We show that the metrisability of an oriented projective surface is equivalent to the existence of pseudo-holomorphic curves. A projective structure $\mathfrak{p}$ and a volume form $\sigma$ on an oriented surface $M$ equip the total space of a certain disk bundle $Z \rightarrow M$ with a pair ( $J_{\mathfrak{p}}, \mathfrak{J}_{\mathfrak{p}, \sigma}$ ) of almost complex structures. A conformal structure on $M$ corresponds to a section of $Z \rightarrow M$ and $\mathfrak{p}$ is metrisable by the metric $g$ if and only if $[g]: M \rightarrow Z$ is a pseudo-holomorphic curve with respect to $J_{\mathfrak{p}}$ and $\mathfrak{J}_{\mathfrak{p}, d A_{g}}$.


## 1 Introduction

A projective structure on a smooth manifold consists of an equivalence class $\mathfrak{p}$ of torsion-free connections on its tangent bundle, where two such connections are called equivalent if they have the same geodesics up to parametrisation. A projective structure $\mathfrak{p}$ is called metrisable if it contains the Levi-Civita connection of some Riemannian metric. The problem of (locally) characterising the projective structures that are metrisable was first studied in the work of R. Liouville [17] in 1889, but was solved only relatively recently by Bryant, Dunajski and Eastwood for the case of two dimensions [2]. Since then, there has been renewed interest in the problem, see $[5,6,8,10,11,13,14,25,27]$ for related recent work.

The purpose of this short note is to show that in the case of an oriented projective surface $(M, \mathfrak{p})$, the metrisability of $\mathfrak{p}$ is equivalent to the existence of certain pseudo-holomorphic curves.

An orientation compatible complex structure on $M$ corresponds to a section of the bundle $\pi: Z \rightarrow M$ whose fibre at $x \in M$ consists of the orientation compatible linear complex structures on $T_{x} M$. The choice of a torsion-free connection $\nabla$ on $T M$ equips $Z$ with an almost complex structure $J[7,26]$. Namely, at $j \in Z$ we lift $j$ horizontally and take a natural complex structure on each fibre vertically. It turns out that $J$ is always integrable and does only depend on the projective equivalence class $\mathfrak{p}$ of $\nabla$, we thus denote it by $J_{\mathfrak{p}}$. Reversing the orientation on each fibre yields another almost complex structure $\mathfrak{J}$ which is however never integrable and is not projectively invariant. Fixing a volume form $\sigma$ on the projective

[^0]surface $(M, \mathfrak{p})$ determines a unique representative connection ${ }^{\sigma} \nabla \in \mathfrak{p}$ which preserves $\sigma$. We will write $\mathfrak{J}_{\mathfrak{p}, \sigma}$ for the non-integrable almost complex structure arising from ${ }^{\sigma} \nabla \in \mathfrak{p}$.

The choice of a conformal structure $[g]$ on an oriented surface $M$ defines an orientation compatible complex structure by rotating a tangent vector counterclockwise by $\pi / 2$ with respect to $[g]$. Thus, we may think of a conformal structure as a section $[g]: M \rightarrow Z$. Denoting the area form of a Riemannian metric $g$ by $d A_{g}$, we show:

Theorem 1 An oriented projective surface $(M, \mathfrak{p})$ is metrisable by the metric $g$ on $M$ if and only if $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$ is a holomorphic curve and $[g]: M \rightarrow\left(Z, \mathfrak{J}_{\mathfrak{p}, d A_{g}}\right)$ is a pseudo-holomorphic curve.

Applying a general existence result for pseudo-holomorphic curves [24, Theorem III] it follows that locally we can always find a Riemannian metric $g$ so that $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$ is a holomorphic curve or so that $[g]: M \rightarrow\left(Z, \mathfrak{J}_{\mathfrak{p}, d A_{g}}\right)$ is a pseudo-holomorphic curve. The geometric significance of the existence of such (pseudo-)holomorphic curves is given in Proposition 9 below.

The construction of the (integrable) almost complex structure $J_{\mathfrak{p}}$ on $Z$ given in $[7,26]$ is adapted from the construction of an almost complex structure $J$ on the twistor space $Y \rightarrow N$ of an oriented Riemannian 4-manifold ( $N, g$ ), see [1]. In the Riemannian setting the almost complex structure $J$ is integrable if and only if $g$ is self-dual. In [12], Eells-Salamon observe that reversing the orientation on each fibre of $Y \rightarrow N$ associates another almost complex structure $\mathfrak{J}$ on $Y$ to $(N, g)$ which is never integrable. Thus, the non-integrable almost complex structure $\mathfrak{J}$ used here may be thought of as the affine analogue of the non-integrable almost complex structure in oriented Riemannian 4-manifold geometry.

## 2 Pseudo-holomorphic curves and metrisability

Recall that the set of torsion-free connections on the tangent bundle of a surface $M$ is an affine space modelled on the smooth sections of the vector bundle $V=S^{2}\left(T^{*} M\right) \otimes T M$. We have a natural trace mapping $\operatorname{tr}: V \rightarrow T^{*} M$, given in abstract index notation by $A_{j k}^{i} \mapsto A_{i k}^{k}$, as well as an inclusion Sym : $T^{*} M \rightarrow V$, given by $b_{i} \mapsto \delta_{j}^{i} b_{k}+\delta_{k}^{i} b_{j}$. The bundle $V$ thus decomposes as $V=V_{0} \oplus T^{*} M$, where $V_{0}$ denotes the trace-free part of $V$. We have (Cartan, Eisenhart, Weyl)-the reader may also consult [9] for a modern reference:

Lemma 2 Two torsion-free connections $\nabla$ and $\nabla^{\prime}$ on $T M$ are projectively equivalent if and only if there exists a 1 -form $\xi$ on $M$ so that $\nabla-\nabla^{\prime}=\operatorname{Sym}(\xi)$.

This gives immediately:
Lemma 3 Let $(M, \mathfrak{p})$ be an oriented projective surface and $\sigma$ a volume form on $M$. Then there exists a unique representative connection ${ }^{\sigma} \nabla \in \mathfrak{p}$ preserving $\sigma$.

Proof Let $\nabla \in \mathfrak{p}$ be a representative connection. Since $\sigma$ is a volume form there exists a unique 1 -form $\alpha$ on $M$ such that $\nabla \sigma=\alpha \otimes \sigma$. An elementary computation shows that the connection $\nabla+\operatorname{Sym}(\xi)$ satisfies

$$
(\nabla+\operatorname{Sym}(\xi)) \sigma=\nabla \sigma-3 \xi \otimes \sigma,
$$

for all $\xi \in \Omega^{1}(M)$. Thus the connection ${ }^{\sigma} \nabla=\nabla+\frac{1}{3} \operatorname{Sym}(\alpha)$ preserves $\sigma$ and clearly is the only connection in $\mathfrak{p}$ doing so.

We also have:

Lemma 4 Let $\varphi \in \Gamma\left(V_{0}\right)$ and $\nabla$ be a torsion-free connection on $T M$. Then $\nabla+\varphi$ preserves a volume form $\sigma$ on $M$ if and only if $\nabla$ preserves the volume form $\sigma$.

Proof Since $\varphi \in \Gamma\left(V_{0}\right)$, an elementary computation shows that the connections $\nabla$ and $\nabla+\varphi$ induce the same connection on the bundle $\Lambda^{2}\left(T^{*} M\right)$ whose non-vanishing sections are the volume forms.

For our purposes it is convenient to construct the almost complex structures $(J, \mathfrak{J})$ associated to $\nabla$ in terms of the connection form $\theta$ on the oriented frame bundle of $M$. The oriented frame bundle $F$ of the oriented surface $M$ is the bundle $v: F \rightarrow M$ whose fibre at $x \in M$ consists of the linear isomorphisms $u: \mathbb{R}^{2} \rightarrow T_{x} M$ that are orientation preserving with respect to the standard orientation on $\mathbb{R}^{2}$ and the given orientation on $T_{x} M$. The group $\mathrm{GL}^{+}(2, \mathbb{R})$ acts transitively from the right on each fibre by the rule $R_{a}(u)=u \circ a$ for all $a \in \mathrm{GL}^{+}(2, \mathbb{R}), u \in F$ and this action turns $v: F \rightarrow M$ into a principal right $\mathrm{GL}^{+}(2, \mathbb{R})$-bundle. The total space $F$ carries a tautological $\mathbb{R}^{2}$-valued 1 -form $\omega$ defined by $\omega_{u}=u^{-1} \circ v_{u}^{\prime}$ and $\omega$ satisfies the equivariance property

$$
\begin{equation*}
R_{a}^{*} \omega=a^{-1} \omega \tag{1}
\end{equation*}
$$

for all $a \in \mathrm{GL}^{+}(2, \mathbb{R})$. We may embed $\mathrm{GL}(1, \mathbb{C})$ as the subgroup of $\mathrm{GL}^{+}(2, \mathbb{R})$ consisting of matrices that commute with the standard linear complex structure on $\mathbb{R}^{2}$. Note that may think of the oriented frame bundle $v: F \rightarrow M$ as a principal $\mathrm{GL}(1, \mathbb{C})$-bundle over $Z=$ $F / \mathrm{GL}(1, \mathbb{C})$. We may describe an almost complex structure on $Z$ by describing the pullback of its ( 1,0 )-forms to $F$. The pullback of a 1 -form on $Z$ to $F$ is semi-basic for the projection $v: F \rightarrow Z$, that is, it vanishes when evaluated on vector fields that are tangent to the fibres of $v$. For $y \in \mathfrak{g l}(2, \mathbb{R})$ we denote by $Y_{y}$ the vector field on $F$ that is generated by the flow $R_{\exp (t y)}$. Clearly, the vector fields $Y_{y}$ for $y \in \mathfrak{g l}(1, \mathbb{C})$ span the vector fields on $F$ that are tangent to the fibres of $v$.

Let $\nabla$ be a torsion-free connection on $T M$ with connection form $\theta=\left(\theta_{j}^{i}\right)$ on $F$. Recall that $\theta$ satisfies the equivariance property

$$
\begin{equation*}
R_{a}^{*} \theta=a^{-1} \theta a \tag{2}
\end{equation*}
$$

for all $a \in \mathrm{GL}^{+}(2, \mathbb{R})$ and the structure equations

$$
\begin{align*}
\mathrm{d} \omega^{i} & =-\theta_{j}^{i} \wedge \omega^{j}, \\
\mathrm{~d} \theta_{j}^{i} & =-\theta_{k}^{i} \wedge \theta_{j}^{k}+\Theta_{j}^{i}, \tag{3}
\end{align*}
$$

where $\Theta=\left(\Theta_{j}^{i}\right)$ denotes the curvature form of $\theta$. Since $\theta$ is a principal connection on $F$ it also satisfies $\theta\left(Y_{y}\right)=y$ for all $y \in \mathfrak{g l}(2, \mathbb{R})$. Since the Lie algebra of $\operatorname{GL}(1, \mathbb{C})$ is spanned by the matrices of the form

$$
\left(\begin{array}{cc}
z & -w \\
w & z
\end{array}\right)
$$

for $(z, w) \in \mathbb{R}^{2}$, the complex-valued 1-forms on $F$ that are semi-basic for the projection $v: F \rightarrow Z$ are spanned by the forms $\omega=\omega^{1}+\mathrm{i} \omega^{2}$ and

$$
\zeta=\left(\theta_{1}^{1}-\theta_{2}^{2}\right)+\mathrm{i}\left(\theta_{2}^{1}+\theta_{1}^{2}\right)
$$

and their complex conjugates. We now have:

Proposition 5 Let $\nabla$ be a torsion-free connection on $T M$ with connection form $\theta=\left(\theta_{j}^{i}\right)$ on $F$. Then there exists a unique pair $(J, \mathfrak{J})$ of almost complex structures on $Z$ whose $(1,0)$ forms pull back to become linear combinations of the forms $(\omega, \zeta)$ in the case of $J$ and to $(\omega, \bar{\zeta})$ in the case of $\mathfrak{J}$. Moreover, the almost complex structure $J$ is always integrable, whereas $\mathfrak{J}$ is never integrable.

Proof Writing

$$
r \mathrm{e}^{\mathrm{i} \phi} \simeq\left(\begin{array}{cc}
r \cos \phi & -r \sin \phi \\
r \sin \phi & r \cos \phi
\end{array}\right)
$$

for the elements of $\operatorname{GL}(1, \mathbb{C})$, the equivariance property (1) of $\omega$ and (2) of $\theta$ implies

$$
\begin{equation*}
\left(R_{r \mathrm{e}^{\mathrm{i} \phi}}\right)^{*} \omega=\frac{1}{r} \mathrm{e}^{\mathrm{i} \phi} \omega \text { and }\left(R_{r \mathrm{e}}\right)^{*} \zeta=\mathrm{e}^{-2 \mathrm{i} \phi} \zeta . \tag{4}
\end{equation*}
$$

It follows that there exists a unique almost complex structure $J$ on $Z$ whose ( 1,0 )-forms pull back to $F$ to become linear combinations of the forms $\omega, \zeta$. Likewise there exists a unique almost complex structure $\mathfrak{J}$ on $Z$ whose ( 1,0 )-forms pull back to $F$ to become linear combinations of the forms $\omega, \bar{\zeta}$. Furthermore, simple computations using the structure equations (3) imply that

$$
0=\mathrm{d} \zeta \wedge \omega \wedge \zeta=\mathrm{d} \omega \wedge \omega \wedge \zeta
$$

Consequently, the Newlander-Nirenberg theorem [23] implies that $J$ is integrable. On the other hand, we get

$$
\mathrm{d} \omega \wedge \omega \wedge \bar{\zeta}=\frac{1}{2} \omega \wedge \bar{\omega} \wedge \zeta \wedge \bar{\zeta}
$$

so that $\mathfrak{J}$ is never integrable.
Remark 6 The equivariance properties (4) imply that the bundles

$$
H=v^{\prime}\{\operatorname{Re}(\zeta)=0, \operatorname{Im}(\zeta)=0\} \quad \text { and } \quad V=v^{\prime}\{\operatorname{Re}(\omega)=0, \operatorname{Im}(\omega)=0\}
$$

are well-defined distributions on $Z$ that are invariant with respect to $J$ (and $\mathfrak{J}$ ). Hence we have $T Z=H \oplus V$.

For the convenience of the reader, we also show [7,26]:
Proposition 7 Suppose the torsion-free connections $\nabla$ and $\nabla^{\prime}$ on $T M$ are projectively equivalent, then they induce the same integrable almost complex structure $J$ on $Z$.

Proof The connections $\nabla$ and $\nabla^{\prime}$ are projectively equivalent if and only if there exists a 1-form $\xi$ on $M$ such that $\nabla^{\prime}=\nabla+\operatorname{Sym}(\xi)$. Writing $\theta=\left(\theta_{j}^{i}\right)$ for the connection form of $\nabla$ on $F$ and $v^{*} \xi=x_{i} \omega^{i}$ for real-valued functions $x_{i}$ on $F$, the connection form $\theta^{\prime}$ of $\nabla^{\prime}$ becomes

$$
\theta^{\prime}=\theta+\left(\begin{array}{cc}
2 x_{1} \omega^{1}+x_{2} \omega^{2} & x_{2} \omega^{1} \\
x_{1} \omega^{2} & x_{1} \omega^{1}+2 x_{2} \omega^{2}
\end{array}\right) .
$$

Consequently, we obtain

$$
\zeta^{\prime}=\zeta+\left(x_{1} \omega^{1}-x_{2} \omega^{2}\right)+\mathrm{i}\left(x_{2} \omega^{1}+x_{1} \omega^{2}\right)=\zeta+\left(x_{1}+\mathrm{i} x_{2}\right) \omega
$$

which shows that the complex span of $\omega, \zeta$ is the same as the one of $\omega, \zeta^{\prime}$ and hence the two integrable almost complex structures are the same.

Remark 8 For a projective structure $\mathfrak{p}$ on $M$ we will write $J_{\mathfrak{p}}$ for the integrable almost complex structure defined by any representative connection $\nabla \in \mathfrak{p}$. For a projective structure $\mathfrak{p}$ and a volume form $\sigma$ on $M$ we will write $\mathfrak{J}_{\mathfrak{p}, \sigma}$ for the non-integrable almost complex structure defined by the representative connection ${ }^{\sigma} \nabla \in \mathfrak{p}$. Note that the non-integrable almost complex structure is not projectively invariant.

Recall that a Weyl connection for a conformal structure [g] is a torsion-free connection ${ }^{[g]} \nabla$ on $T M$ which preserves $[g]$. Fixing a Riemannian metric $g \in[g]$, the Weyl connections for $[g]$ can be written as ${ }^{[g]} \nabla={ }^{g} \nabla+g \otimes B-\operatorname{Sym}(\beta)$ for some 1 -form $\beta$ on $M$ and where $B$ denotes the $g$-dual vector field to $\beta$. In [20] and in the language of thermostats in [22], it was observed that for every choice of a conformal structure $[g]$ on a projective surface $(M, \mathfrak{p})$, there exists a unique Weyl connection ${ }^{[g]} \nabla$ for $[g]$ and a unique 1-form $\varphi \in \Gamma\left(V_{0}\right)$ so that ${ }^{[g]} \nabla+\varphi$ is a representative connection of $\mathfrak{p}$. Moreover the endomorphism $\varphi(X)$ is symmetric with respect to $[g]$ for every vector field $X$ on $M$. We call ${ }^{[g]} \nabla$ the Weyl connection determined by $[g]$. Explicitly, if $\nabla$ is any representative connection of $\mathfrak{p}, g \in[g]$ and if we define a vector field $B=\frac{3}{4} \operatorname{tr}\left(g^{\sharp} \otimes\left(\nabla-{ }^{g} \nabla\right)_{0}\right)$, then

$$
\varphi=\left(\nabla-{ }^{g} \nabla-g \otimes B\right)_{0} \quad \text { and } \quad{ }^{[g]} \nabla={ }^{g} \nabla+g \otimes B-\operatorname{Sym}(\beta),
$$

where $A_{0}$ denotes the trace-free part of a tensor field $A \in \Gamma\left(S^{2}\left(T^{*} M\right) \otimes T M\right)$. We refer the reader to $[20,22]$ for a proof that ${ }^{[g]} \nabla$ and $\varphi$ do satisfy the claimed properties.

Proposition 9 Let $(M, \mathfrak{p})$ be an oriented projective surface and $g$ a Riemannian metric on M. Then we have:
(i) $\mathfrak{p}$ contains a Weyl connection for $[g]$ if and only if $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$ is a holomorphic curve;
(ii) the Weyl connection determined by $[g]$ is the Levi-Civita connection of $g$ if and only if $[g]: M \rightarrow\left(Z, \mathfrak{J}_{\mathfrak{p}, d A_{g}}\right)$ is a pseudo-holomorphic curve.

Remark 10 Here we say $[g]: M \rightarrow(Z, \mathfrak{J})$ is a (pseudo-)holomorphic curve if the image $\Sigma=[g](M) \subset Z$ admits the structure of a (pseudo-)holomorphic curve. By admitting the structure of (pseudo-)holomorphic curve, we mean that $\Sigma$ can be equipped with a complex structure $J$, so that the inclusion $\iota: \Sigma \rightarrow Z$ is $(J, \mathfrak{J})$-linear, that is, satisfies $\mathfrak{J} \circ \iota^{\prime}=\iota^{\prime} \circ J$.

As an immediate consequence, we obtain the Theorem 1:
Proof of Theorem 1 The projective structure $\mathfrak{p}$ is metrisable by $g$ if and only if the Weyl connection determined by [g] is the Levi-Civita connection of $g$ and the 1 -form $\varphi$ vanishes identically. The claim follows by applying Proposition 9.

For the proof of Proposition 9 we also need the following Lemma:
Lemma 11 Let $(Z, \mathfrak{J})$ be an almost complex four-manifold and $\omega, \chi \in \Omega^{1}(Z, \mathbb{C})$ a basis for the $(1,0)$-forms of $Z$. Suppose $\iota: \Sigma \rightarrow Z$ is an immersed surface so that $\iota^{*}(\omega \wedge \bar{\omega})$ is non-vanishing on $\Sigma$. Then $\Sigma$ admits the structure of a pseudo-holomorphic curve if and only if $\iota^{*}(\omega \wedge \chi)$ vanishes identically on $\Sigma$.

Proof Since $\iota^{*}(\omega \wedge \bar{\omega})$ is non-vanishing on $\Sigma$, the forms $\iota^{*} \omega$ and $\iota^{*} \bar{\omega}$ span the complex-valued 1-forms on $\Sigma$. Recall that $\iota: \Sigma \rightarrow Z$ is $(j, \mathfrak{J})$-linear if and only if the pullback of every ( 1,0 )-form on $Z$ is a $(1,0)$-form on $\Sigma$, the claim follows.

Proof of Proposition 9 Let $g$ be a Riemannian metric on the oriented projective surface ( $M, \mathfrak{p}$ ). Without losing generality we can assume that the projective structure $\mathfrak{p}$ arises from a connection of the form ${ }^{[g]} \nabla+\varphi$. The Weyl connection ${ }^{[g]} \nabla$ satisfies

$$
{ }^{[g]} \nabla d A_{g}=2 \beta \otimes d A_{g}
$$

for some 1-form $\beta$ on $M$ and hence can be written as ${ }^{[g]} \nabla={ }^{g} \nabla+g \otimes \beta^{\sharp}-\operatorname{Sym}(\beta)$.
Now suppose $\nabla \in \mathfrak{p}$ preserves the volume form $d A_{g}$ of $g$. Then, by Lemma 4 it must be of the form

$$
\begin{equation*}
\nabla={ }^{[g]} \nabla+\varphi+\frac{2}{3} \operatorname{Sym}(\beta)={ }^{g} \nabla+g \otimes \beta^{\sharp}-\frac{1}{3} \operatorname{Sym}(\beta)+\varphi . \tag{5}
\end{equation*}
$$

Proposition 5 and Lemma 11 imply that the condition that $[g]: M \rightarrow Z$ defines a pseudoholomorphic curve with respect to $J_{\mathfrak{p}}$ respectively $\mathfrak{J}_{\mathfrak{p}, d A_{g}}$ is equivalent to the condition that on the pullback bundle $[g]^{*} F \rightarrow M$ the form $\omega \wedge \zeta$, respectively $\omega \wedge \bar{\zeta}$ vanishes identically, where $\zeta$ is computed from the connection form of $\nabla$ and where we think of $F$ as fibering over $Z$. Keeping this in mind we now compute the pullback of the forms $\zeta$ and $\bar{\zeta}$ to $[g]^{*} F$. Recall that the semi-basic 1 -forms on $F$ are spanned by the components of $\omega$, hence there exist unique real-valued functions $g_{i j}=g_{j i}$ on $F$ so that $v^{*} g=g_{i j} \omega^{i} \otimes \omega^{j}$. Likewise, there exist unique real-valued functions $b_{i}$ on $F$ so that $v^{*} \beta=b_{i} \omega^{i}$ and unique real-valued function $A_{j k}^{i}=A_{k j}^{i}$ on $F$ so that $\left(v^{*} \varphi\right)_{j}^{i}=A_{j k}^{i} \omega^{k}$. The functions $A_{j k}^{i}$ satisfy furthermore $A_{k i}^{k}=0$ and $g_{i k} A_{j l}^{k}=g_{j k} A_{i l}^{k}$ since $\varphi$ takes values in the endomorphisms of $T M$ that are trace-free and symmetric with respect to $g$. The Levi-Civita connection $\left(\psi_{j}^{i}\right)$ of $g$ is the unique principal $\mathrm{GL}^{+}(2, \mathbb{R})$-connection on $F$ that satisfies

$$
\begin{aligned}
\mathrm{d} \omega^{i} & =-\psi_{j}^{i} \wedge \omega^{j}, \\
\mathrm{~d} g_{i j} & =g_{i k} \psi_{j}^{k}+g_{k j} \psi_{i}^{k} .
\end{aligned}
$$

The pullback bundle $P:=[g]^{*} F$ is cut out by the equations $g_{11}=g_{22}$ and $g_{12}=0$. On $P$ we have

$$
\begin{gathered}
0=\mathrm{d} g_{12}=g_{11} \psi_{2}^{1}+g_{22} \psi_{1}^{2}=g_{11}\left(\psi_{2}^{1}+\psi_{1}^{2}\right) \\
0=\mathrm{d} g_{11}-\mathrm{d} g_{22}=2 g_{11} \psi_{1}^{1}-2 g_{22} \psi_{2}^{2}=g_{11}\left(\psi_{1}^{1}-\psi_{2}^{2}\right)
\end{gathered}
$$

On $P$ the condition $g_{i k} A_{j l}^{k}=g_{j k} A_{i l}^{k}$ implies $A_{11}^{2}=-A_{22}^{2}$ and $A_{22}^{1}=-A_{11}^{1}$. Writing $A_{11}^{1}=a_{1}$ and $A_{22}^{2}=a_{2}$ and using (5), the connection form $\theta$ of $\nabla$ thus becomes

$$
\begin{aligned}
\theta= & \left(\begin{array}{cc}
\psi_{1}^{1} & -\psi_{1}^{2} \\
\psi_{1}^{2} & \psi_{1}^{1}
\end{array}\right)+\left(\begin{array}{ll}
b_{1} \omega^{1} & b_{1} \omega^{2} \\
b_{2} \omega^{1} & b_{2} \omega^{2}
\end{array}\right)-\frac{1}{3}\left(\begin{array}{cc}
2 b_{1} \omega^{1}+b_{2} \omega^{2} & b_{2} \omega^{1} \\
b_{1} \omega^{2} & b_{1} \omega^{1}+2 b_{2} \omega^{2}
\end{array}\right) \\
& +\left(\begin{array}{cc}
a_{1} \omega^{1}-a_{2} \omega^{2} & -a_{2} \omega^{1}-a_{1} \omega^{2} \\
-a_{2} \omega^{1}-a_{1} \omega^{2} & -a_{1} \omega^{1}+a_{2} \omega^{2}
\end{array}\right)
\end{aligned}
$$

Introducing the complex notation $a=a_{1}+\mathrm{i} a_{2}$ and $b=\frac{1}{2}\left(b_{1}-\mathrm{i} b_{2}\right)$, we obtain from a simple calculation

$$
\zeta=\left(\theta_{1}^{1}-\theta_{2}^{2}\right)+\mathrm{i}\left(\theta_{2}^{1}+\theta_{1}^{2}\right)=\frac{4}{3} \bar{b} \omega+2 \bar{a} \bar{\omega},
$$

where we write $\omega=\omega^{1}+\mathrm{i} \omega^{2}$.
Finally, since $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$ is a holomorphic curve if and only if $\omega \wedge \zeta$ vanishes identically on $P$, it follows that $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$ is a holomorphic curve if and only if

$$
0=\omega \wedge \zeta=2 \bar{a} \omega \wedge \bar{\omega}
$$

which is equivalent to $\varphi$ vanishing identically. This shows (i).
Likewise $[g]: M \rightarrow\left(Z, \mathfrak{J}_{\mathfrak{p}, d A_{g}}\right)$ is a pseudo-holomorphic curve if and only if

$$
0=\omega \wedge \bar{\zeta}=\frac{4}{3} b \omega \wedge \bar{\omega}
$$

on $P$. This is equivalent to $\beta$ vanishing identically. This shows (ii).
As a corollary we obtain:
Corollary 12 Let $(M, \mathfrak{p})$ be a projective surface. Then locally $\mathfrak{p}$ contains
(i) a Weyl connection ${ }^{[g]} \nabla$ for some conformal structure $[g]$;
(ii) a connection of the form ${ }^{\tilde{g}} \nabla+\varphi$ for some Riemannian metric $\tilde{g}$ and some $\varphi \in \Gamma\left(V_{0}\right)$ with $\varphi$ taking values in the endomorphisms that are $\tilde{g}$-symmetric.

Remark 13 The first statement of Proposition 9 and Corollary 12 was previously obtained in [19].

Proof of Corollary 12 We first consider the case (ii). We fix a volume form $\sigma$ on $M$. We need to show that in a neighbourhood $U_{x}$ of every point $x \in M$ there exists a conformal structure $[g]$ which is a pseudo-holomorphic curve into the total space of the bundle $\pi: Z \rightarrow M$, where we equip $Z$ with the almost complex structure $\mathfrak{J}_{p, \sigma}$. Choose $j \in Z$ with $\pi(j)=x$. Recall from Remark 6 that the subspace $H_{j} \subset T_{j} Z$ is invariant under $\mathfrak{J}_{\mathfrak{p}, \sigma}$. Now [24, Theorem III] implies that there exists a pseudo-holomorphic curve $\Sigma \subset\left(Z, \mathfrak{J}_{\mathfrak{p}, \sigma}\right)$ which contains $j$ and has $H_{j}$ as its tangent space at $j$. Since $H_{j} \subset T_{j} Z$ is horizontal, the restriction $\left.\pi_{j}^{\prime}\right|_{H_{j}}: H_{j} \rightarrow T_{x} M$ is an isomorphism. Therefore, the restriction of $\pi$ to $\Sigma$ is a local diffeomorphism in some neighbourhood of $j$. Hence there exists a neighbourhood $U_{x}$ of $x \in M$ and a section $[g]: U_{x} \rightarrow Z$ so that $[g]\left(U_{x}\right) \subset \Sigma$. Thus, $[g]: U_{x} \rightarrow\left(Z, \mathfrak{J}_{\mathfrak{p}, \sigma}\right)$ is a pseudo-holomorphic curve in the sense of Remark 10. Taking $\tilde{g}$ to be the unique metric in [g] with volume form $\sigma$ and applying Proposition 9 shows the claim. The case (i) follows in the same fashion, except that [24] is not needed, as $J_{\mathfrak{p}}$ is integrable and hence the construction of a holomorphic curve realising a prescribed $J_{\mathfrak{p}}$-invariant tangent plane is an elementary exercise.

Remark 14 Locally we can always find a holomorphic curve $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$, but globally this is not always possible. A properly convex projective structure $\mathfrak{p}$ on a closed surface $M$ with $\chi(M)<0$ admits a holomorphic curve $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$ if and only if $\mathfrak{p}$ is hyperbolic [22]. One would expect that a corresponding global non-existence result should also hold in the pseudo-holomorphic setting for a suitable class of projective surfaces.

Remark 15 If $(M, \mathfrak{p})$ is a closed oriented projective surface of with $\chi(M)<0$, then there exists at most one holomorphic curve $[g]: M \rightarrow\left(Z, J_{\mathfrak{p}}\right)$, see [21].
Remark 16 Hitchin [15] gave a twistorial construction of (complex) two-dimensional holomorphic projective structures. In the holomorphic category such a projective structure corresponds to a complex surface $Z$ having a family of rational curves with self-intersection number one. Denoting the canonical bundle of $Z$ by $K_{Z}$, such a holomorphic projective surface is metrisable if and only if $K_{Z}^{-2 / 3}$ admits a holomorphic section which intersects each rational curve in $Z$ at two points $[2,3,16]$.

Remark 17 The notion of a projective structure also makes sense in the complex setting and such structures are referred to as c-projective, see [4]. Correspondingly, there is a Kähler metrisability problem of c-projective structures. Some obstructions to Kähler metrisability of a (complex) two-dimensional c-projective structure have been obtained in [18].

We conclude by describing the holomorphic curves for the standard projective structure $\mathfrak{p}_{0}$ on the 2 -sphere whose geodesics are the great circles.

Example 18 Let $S^{2}$ denote the sphere of radius 1 centered at the origin in $\mathbb{R}^{3}$ and $g$ its induced round metric of constant Gauss curvature 1 whose geodesics are the great circles. We equip $S^{2}$ with its standard orientation.

Recall that the unit tangent bundle $\lambda: T_{1} S^{2} \rightarrow S^{2}$ of ( $S^{2}, g$ ) carries a canonical coframing $\left(\omega_{1}, \omega_{2}, \psi\right)$, where $\omega_{1}, \omega_{2}$ span the 1-forms on $T_{1} S^{2}$ that are semi-basic for the projection $\lambda$ and $\psi$ denotes the Levi-Civita connection form of $g$. The 1 -forms $\left(\omega_{1}, \omega_{2}, \psi\right)$ satisfy the structure equations

$$
\begin{equation*}
\mathrm{d} \omega_{1}=-\omega_{2} \wedge \psi \quad \text { and } \quad \mathrm{d} \omega_{2}=-\psi \wedge \omega_{1} \quad \text { and } \quad \mathrm{d} \psi=-\omega_{1} \wedge \omega_{2} . \tag{6}
\end{equation*}
$$

Let $\hat{g}$ be a Riemannian metric on $S^{2}$ and write $\lambda^{*} \hat{g}=\hat{g}_{i j} \omega_{i} \otimes \omega_{j}$ for unique real-valued functions $\hat{g}_{i j}=\hat{g}_{j i}$ on $T_{1} S^{2}$. Phrased in modern language (c.f. [2]) and applied to the case of the 2 -sphere, Liouville's result [17] implies that if the metrics $\hat{g}$ and $g$ have the same unparametrised geodesics then the functions $h_{i j}:=\hat{g}_{i j}\left(\hat{g}_{11} \hat{g}_{22}-\hat{g}_{12}^{2}\right)^{-2 / 3}$ satisfy the linear differential equations

$$
\begin{align*}
& \mathrm{d} h_{11}=-2 h_{1} \omega_{2}+2 h_{12} \psi, \\
& \mathrm{~d} h_{12}=h_{1} \omega_{1}-h_{2} \omega_{2}-\left(h_{11}-h_{22}\right) \psi,  \tag{7}\\
& \mathrm{d} h_{22}=2 h_{2} \omega_{1}-2 h_{12} \psi,
\end{align*}
$$

for some smooth real-valued functions $h_{i}$ on $T_{1} S^{2}$. Conversely, a solution to (7) on $T_{1} S^{2}$ satisfying $h_{11} h_{22}-h_{12}^{2} \neq 0$ gives a Riemannian metric $\hat{g}$ on $S^{2}$ with $\lambda^{*} \hat{g}=\left(h_{i j}\left(h_{11} h_{22}-\right.\right.$ $\left.\left.h_{12}^{2}\right)^{-2}\right) \omega_{i} \otimes \omega_{j}$ and that has the same unparametrised geodesics as $g$.

Applying the exterior derivative to the above system of equations implies the existence of a unique real-valued function $h$ on $T_{1} S^{2}$ such that

$$
\begin{aligned}
& \mathrm{d} h_{1}=-h_{12} \omega_{1}+\left(h_{11}+h\right) \omega_{2}+h_{2} \psi, \\
& \mathrm{~d} h_{2}=-\left(h_{22}+h\right) \omega_{1}+h_{12} \omega_{2}-h_{1} \psi .
\end{aligned}
$$

Taking yet another exterior derivative gives that

$$
\mathrm{d} h=-2 h_{1} \omega_{1}+2 h_{2} \omega_{2} .
$$

Writing

$$
\vartheta=\left(\begin{array}{ccc}
0 & -\omega_{1} & -\omega_{2} \\
\omega_{1} & 0 & -\psi \\
\omega_{2} & \psi & 0
\end{array}\right) \quad \text { and } \quad H=\left(\begin{array}{ccc}
h & h_{2} & -h_{1} \\
h_{2} & -h_{22} & h_{12} \\
-h_{1} & h_{12} & -h_{11}
\end{array}\right)
$$

the above system of differential equations can be expressed as

$$
\mathrm{d} H+\vartheta H+H \vartheta^{t}=0 .
$$

The structure equations (6) imply that $\mathrm{d} \vartheta+\vartheta \wedge \vartheta=0$, hence we may write $\vartheta=\Xi^{-1} \mathrm{~d} \Xi$ for some diffeomorphism $\Xi: T_{1} S^{2} \rightarrow \mathrm{SO}(3)$. It follows that the solutions are of the form $H=\Xi^{-1} C\left(\Xi^{-1}\right)^{t}$ for some constant symmetric 3-by-3 matrix $C$. In particular, taking $C=A A^{t}$ for some $A \in \operatorname{SL}(3, \mathbb{R})$, we obtain a solution $H_{A}$ providing a metric $\hat{g}_{A}$ on $S^{2}$ having the great circles as its geodesics.

Finally, in order to construct the holomorphic curve $\left[\hat{g}_{A}\right]: S^{2} \rightarrow Z$ from $H_{A}$, we interpret $Z$ as an associated bundle to $T_{1} S^{2}$. We will only give a sketch of the construction and refer the reader to [22, §4] for additional details. The orientation and metric turn $S^{2}$ into
a Riemann surface and hence a conformal structure on $S^{2}$ is given in terms of a Beltrami differential. Denoting the canonical bundle of $S^{2}$ by $K_{S^{2}}$, a Beltrami differential is a section $\mu$ of $\overline{K_{S^{2}}} \otimes K_{S^{2}}^{-1}$ satisfying $|\mu(x)|<1$ for all $x \in S^{2}$, where $|\cdot|$ denotes the norm induced by the natural Hermitian bundle metric on $\overline{K_{S^{2}}} \otimes K_{S^{2}}^{-1}$. The Riemannian metric $g$ gives an isomorphism $\overline{K_{S^{2}}} \otimes K_{S^{2}}^{-1} \simeq K_{S^{2}}^{-2}$ and thus $Z$ may be identified with $T_{1} S^{2} \times{ }_{S^{1}} \mathbb{D}$, where $S^{1}$ acts by usual rotation on $T_{1} S^{2}$ and by $z \cdot \mathrm{e}^{\mathrm{i} \phi}=z \mathrm{e}^{-2 \mathrm{i} \phi}$ on the open unit disk $\mathbb{D} \subset \mathbb{C}$. A holomorphic curve $[\hat{g}]: S^{2} \rightarrow Z$ is therefore represented by a map $\mu: T_{1} S^{2} \rightarrow \mathbb{D}$. Explicitly, the conformal structure arising from a Riemannian metric $\hat{g}$ on $S^{2}$ is represented by the map

$$
\mu=\frac{p-q+2 \mathrm{i} r}{p+q+2 \sqrt{p q-r^{2}}},
$$

where we write $\lambda^{*} \hat{g}=p \omega_{1} \otimes \omega_{1}+2 r \omega_{1} \circ \omega_{2}+q \omega_{2} \otimes \omega_{2}$ for unique real-valued functions $p, q, r$ on $T_{1} S^{2}$. In our case, the holomorphic curve $\left[\hat{g}_{A}\right]: S^{2} \rightarrow Z$ is thus represented by $\mu$ with

$$
p=\frac{h_{11}}{\left(h_{11} h_{22}-h_{12}^{2}\right)^{2}}, \quad r=\frac{h_{12}}{\left(h_{11} h_{22}-h_{12}^{2}\right)^{2}}, \quad q=\frac{h_{22}}{\left(h_{11} h_{22}-h_{12}^{2}\right)^{2}}
$$

and where the functions $h_{i j}$ arise from $H_{A}$ as above.
Remark 19 In the case of the standard projective structure on $S^{2}$ the complex surface ( $Z, J_{\mathfrak{p}_{0}}$ ) is biholomorphic to $\mathbb{C P}^{2} \backslash \mathbb{R} \mathbb{P}^{2}$ and moreover, the image of a holomorphic curve $[g]: S^{2} \rightarrow Z$ is a smooth quadric, see [19]. Trying to explicitly relate the holomorphic curve $\left[\hat{g}_{A}\right]$ to its image quadric does in general however not seem to give manageable expressions.

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## References

1. Atiyah, M.F., Hitchin, N.J., Singer, I.M.: Self-duality in four-dimensional Riemannian geometry. Proc. R. Soc. Lond. Ser. A 36, 425-461 (1978). https://doi.org/10.1098/rspa.1978.0143
2. Bryant, R.L., Dunajski, M., Eastwood, M.: Metrisability of two-dimensional projective structures. J. Differ. Geom. 83, 465-499 (2009)
3. Calderbank, D.M.J.: Selfdual 4-manifolds, projective surfaces, and the Dunajski-West construction, SIGMA Symmetry Integrability Geom. Methods Appl. 10, Paper 035, 18 (2014). https://doi.org/10. 3842/SIGMA.2014.035
4. Calderbank, D.M., Eastwood, M., Matveev, V.S., Neusser, K.: C-projective geometry (2015). arxiv:1512.04516
5. Čap, A., Gover, A.R., Macbeth, H.R.: Einstein metrics in projective geometry. Geom. Dedicata 168, 235-244 (2014). https://doi.org/10.1007/s10711-013-9828-3
6. Contatto, F., Dunajski, M.: Metrisability of Painlevé equations. J. Math. Phys. 59(9), 023507 (2018). https://doi.org/10.1063/1.4998147
7. Dubois-Violette, M.: Structures complexes au-dessus des variétés, applications. In: Mathematics and physics (Paris, 1979/1982), Progr. Math., vol. 37. Birkhäuser, Boston, pp. 1-42 (1983). https://doi.org/ 10.1007/BF02591680
8. Dunajski, M., Eastwood, M.: Metrisability of three-dimensional path geometries. Eur. J. Math. 2, 809-834 (2016). https://doi.org/10.1007/s40879-016-0095-3
9. Eastwood, M.: Notes on projective differential geometry, in Symmetries and overdetermined systems of partial differential equations, IMA Vol. Math. Appl., vol. 144. Springer, New York, pp. 41-60 (2008). https://doi.org/10.1007/978-0-387-73831-4_3
10. Eastwood, M.: Metrisability of three-dimensional projective structures (2017). arxiv:1712.06191
11. Eastwood, M., Matveev, V.: Metric connections in projective differential geometry. In: Symmetries and overdetermined systems of partial differential equations, IMA Vol. Math. Appl., vol. 144. Springer, New York, pp. 339-350 (2008). https://doi.org/10.1007/978-0-387-73831-4_16
12. Eells, J., Salamon, S.: Twistorial construction of harmonic maps of surfaces into four-manifolds. Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4) 12 (1985), 589-640 (1986)
13. Flood, K., Gover, A.R.: Metrics in projective differential geometry: the geometry of solutions to the metrizability equation. J. Geom. Anal. 29, 2492-2525 (2019). https://doi.org/10.1007/s12220-018-00845
14. Gover, A.R., Macbeth, H.R.: Detecting Einstein geodesics: Einstein metrics in projective and conformal geometry. Differ. Geom. Appl. 33, 44-69 (2014). https://doi.org/10.1016/j.difgeo.2013.10.011
15. Hitchin, N.: Complex manifolds and Einstein's equations. In: Twistor Geometry and Nonlinear Systems. Lecture Notes in Mathematics, vol. 970. Springer, Berlin (1982)
16. Lebrun, C., Mason, L.J.: Zoll manifolds and complex surfaces. J. Differ. Geom. 61, 453-535 (2002)
17. Liouville, R.: Sur les invariants de certaines équations différentielles et sur leurs applications. C. R. Acad. Sci. Paris 109, 560-563 (1889)
18. Mettler, T.: On Kähler metrisability of two-dimensional complex projective structures. Monatsh. Math. 174, 599-616 (2014). https://doi.org/10.1007/s00605-014-0636-0
19. Mettler, T.: Weyl metrisability of two-dimensional projective structures. Math. Proc. Camb. Philos. Soc. 156, 99-113 (2014). https://doi.org/10.1017/S0305004113000522
20. Mettler, T.: Extremal conformal structures on projective surfaces. Ann. Scuola Norm. Sup. Pisa Cl. Sci. (5) (2015). https://doi.org/10.2422/2036-2145.201802_008. arxiv:1510.01043 (to appear)
21. Mettler, T.: Geodesic rigidity of conformal connections on surfaces. Math. Z. 281, 379-393 (2015). https:// doi.org/10.1007/s00209-015-1489-5
22. Mettler, T., Paternain, G.P.: Convex projective surfaces with compatible Weyl connection are hyperbolic. Anal. PDE 13, 1073-1097 (2020). https://doi.org/10.2140/apde.2020.13.1073
23. Newlander, A., Nirenberg, L.: Complex analytic coordinates in almost complex manifolds. Ann. Math. 65(2), 391-404 (1957)
24. Nijenhuis, A., Woolf, W.B.: Some integration problems in almost-complex and complex manifolds. Ann. Math. 77(2), 424-489 (1963). https://doi.org/10.2307/1970126
25. Nurowski, P.: Projective versus metric structures. J. Geom. Phys. 62, 657-674 (2012). https://doi.org/10. 1016/j.geomphys.2011.04.011
26. O’Brian, N.R., Rawnsley, J.H.: Twistor spaces. Ann. Global Anal. Geom. 3, 29-58 (1985). https://doi. org/10.1007/BF00054490
27. Vollmer, A.: Projectively related superintegrable systems. J. Phys. A: Math. Theor. 53, 095202 (2020). https://doi.org/10.1088/1751-8121/ab6fc5

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