

Eigenvalue estimates for Fourier concentration operators on two domains

FELIPE MARCECA, JOSÉ LUIS ROMERO® & MICHAEL SPECKBACHER

Communicated by A. FIGALLI

Abstract

We study concentration operators associated with either the discrete or the continuous Fourier transform, that is, operators that incorporate a spatial cut-off and a subsequent frequency cut-off to the Fourier inversion formula. The spectral profiles of these operators describe the number of prominent degrees of freedom in problems where functions are assumed to be supported on a certain domain and their Fourier transforms are known or measured on a second domain. We derive eigenvalue estimates that quantify the extent to which Fourier concentration operators deviate from orthogonal projectors, by bounding the number of eigenvalues that are away from 0 and 1 in terms of the geometry of the spatial and frequency domains, and a factor that grows at most poly-logarithmically on the inverse of the spectral margin. The estimates are non-asymptotic in the sense that they are applicable to concrete domains and spectral thresholds, and almost match asymptotic benchmarks. Our work covers, for the first time, non-convex and non-symmetric spatial and frequency concentration domains, as demanded by numerous applications that exploit the expected approximate low dimensionality of the modeled phenomena. The proofs build on Israel's work on one dimensional intervals arXiv:1502.04404v1. The new ingredients are the use of redundant wave-packet expansions and a dyadic decomposition argument to obtain Schatten norm estimates for Hankel operators.

1. Introduction and Results

Fourier concentration operators act by incorporating a spatial cut-off and a subsequent frequency cut-off to the Fourier inversion formula. The chief example concerns the Fourier transform on the Euclidean space $\mathcal{F}:L^2(\mathbb{R}^d)\to L^2(\mathbb{R}^d)$, the cut-offs are then given by the indicator functions of two compact domains $E,F\subseteq\mathbb{R}^d$, and the concentration operator is

$$Sf = \chi_F \mathcal{F}^{-1} \chi_E \mathcal{F} \chi_F f, \qquad f \in L^2(\mathbb{R}^d). \tag{1.1}$$

Published online: 12 April 2024

These operators and their analogues, defined with respect to the discrete Fourier transform $L^2([-1/2, 1/2]^d) \to \ell^2(\mathbb{Z}^d)$, play a crucial role in many analysis problems and fields of application where the shapes of E, F are dictated by various physical constraints or measurement characteristics [16,24,25,37].

The basic intuition, stemming from the Fourier uncertainty principle [11,16], is that the concentration operator (1.1) is approximately a projection with rank $tr(S) = |E| \cdot |F|$. The error of such heuristic is encoded by the so-called *plunge* region

$$M_{\varepsilon}(S) = \{\lambda \in \sigma(S) : \varepsilon < \lambda < 1 - \varepsilon\}, \qquad \varepsilon \in (0, 1/2),$$
 (1.2)

consisting of intermediate eigenvalues. Asymptotics for the cardinality of $M_{\varepsilon}(S)$ go back to Landau and Widom [23,26] for the case of one dimensional intervals E = [-a, a], F = [-b, b] and read as

$$\#M_{\varepsilon}(S) = c \cdot \log(ab) \cdot \log(\frac{1-\varepsilon}{\varepsilon}) + o(\log(ab)), \quad \text{as } ab \to \infty,$$
 (1.3)

for an explicit constant c that depends on the normalization of the Fourier transform. The modern spectral theory of Wiener-Hopf operators gives similar asymptotics for concentration operators associated to rather general multi-dimensional domains subject to increasing isotropic dilations.

While (1.3) precisely describes the cardinality of the set $M_{\varepsilon}(S)$ in the limit $ab \to \infty$, the asymptotic is often insufficient for many purposes because of the quality of the error terms. Indeed, the error term in (1.3) depends in an unspecified way on the spectral threshold ε , which precludes applications where ε is let to vary with the domains E, F. Such limitations have motivated a great amount of work aimed at deriving upper bounds for $\#M_{\varepsilon}(S)$ that are threshold robust, that is, bounds that are effective for concrete concentration domains and explicit in their dependence on the spectral threshold [6,17,19,20,31,39], significantly improving on more classical results in this spirit [38].

With the exception of [17], the mentioned articles on threshold-robust spectral bounds for Fourier concentration operators concern only the one dimensional case, because they exploit a connection with a Sturm-Liouville equation which is specific to that setting. Such methods can be applied to some extent to higher dimensional domains enjoying special symmetries [13,35]. On the other hand, while [17] studies Fourier concentration operators associated with one dimensional intervals, the technique introduced by Israel is very general, as it relies on an explicit almost diagonalization of the concentration operator. In fact, as we were finishing this work, the preprint version of [18] provided an extension of [17] to higher dimensions (see Sects. 1.1 and 1.6).

In this article we derive upper bounds for the number of intermediate eigenvalues (1.2) associated with either the continuous or discrete Fourier transforms. In contrast to other results in the literature, we obtain estimates that apply to two suitably regular multi-dimensional spatial and frequency domains, which do not need to exhibit special symmetries. In this way, our results cover for the first time many setups of practical relevance, see Sect. 1.5.

Our proofs build on Israel's technique [17] and incorporate novel arguments to treat non-convex domains and their discrete counterparts. Instead of the orthonormal

wave-packet basis from [17], we use more versatile redundant expansions (frames). Second, we introduce a dyadic decomposition method implemented by means of Schatten norm estimates for Hankel operators; see Sect. 1.6.

1.1. The Euclidean Space

Given two compact sets $E, F \subseteq \mathbb{R}^d$, the *Fourier concentration operator S*: $L^2(\mathbb{R}^d) \to L^2(\mathbb{R}^d)$ is defined by (1.1) where \mathcal{F} denotes the Fourier transform

$$\mathcal{F}f(\xi) = \int_{\mathbb{R}^d} f(x)e^{-2\pi ix\xi} dx. \tag{1.4}$$

A set $E \subseteq \mathbb{R}^d$ is said to have a *maximally Ahlfors regular boundary* if there exists a constant $\kappa_{\partial E} > 0$ such that

$$\mathcal{H}^{d-1}\big(\partial E \cap B_r(x)\big) \geq \kappa_{\partial E} \cdot r^{d-1}, \qquad 0 < r \leq \mathcal{H}^{d-1}(\partial E)^{1/(d-1)}, \quad x \in \partial E.$$

Here, \mathcal{H}^{d-1} denotes the (d-1)-dimensional Hausdorff measure. The term maximal in the definition refers to the range of r for which the estimate is required to hold. See Sect. 2 for more context on Ahlfors regularity. In what follows, we denote for short $|\partial E| = \mathcal{H}^{d-1}(\partial E)$.

In this article we prove the following:

Theorem 1.1. Let $E, F \subseteq \mathbb{R}^d, d \ge 2$, be compact domains with maximally Ahlfors regular boundaries with constants $\kappa_{\partial E}, \kappa_{\partial F}$ respectively, and assume that that $|\partial E||\partial F| \ge 1$. Consider the concentration operator (1.1) and its eigenvalues $\{\lambda_n : n \in \mathbb{N}\}$.

Then for every $\alpha \in (0, 1/2)$, there exists $A_{\alpha,d} \ge 1$ such that for $\varepsilon \in (0, 1/2)$:

$$\#\{n \in \mathbb{N} : \lambda_n \in (\varepsilon, 1 - \varepsilon)\} \le A_{\alpha, d} \cdot \frac{|\partial E|}{\kappa_{\partial E}} \cdot \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log\left(\frac{|\partial E||\partial F|}{\kappa_{\partial E} \varepsilon}\right)^{2d(1 + \alpha) + 1}.$$
(1.5)

The strength of Theorem 1.1 lies in the fact that the right-hand side of (1.5) grows only mildly on ε , in agreement with the Landau-Widom asymptotic formula for one dimensional intervals (1.3). In contrast, cruder estimates based on computing first and second moments of concentration operators, as done often in sampling theory [22], lead to error bounds of the order $O(1/\varepsilon)$.

A result closely related to Theorem 1.1 is presented in the recent article [18]. For $F = [0, 1]^d$ and E = rK, where $r \ge 1$ is a dilation parameter and $K \subset B_1(0) \subset \mathbb{R}^d$ is a *convex, coordinate symmetric domain* [18, Theorem 1.1] gives the following bound for $\varepsilon \in (0, 1/2)$:

$$\#\{n \in \mathbb{N} : \lambda_n \in (\varepsilon, 1 - \varepsilon)\} \le C_d \cdot \max\{r^{d-1} \log(r/\varepsilon)^{5/2}, \log(r/\varepsilon)^{5d/2}\}. \quad (1.6)$$

For large r, the right-hand side of (1.6) becomes $O_d(r^{d-1}\log(r/\varepsilon)^{5/2})$ while Theorem 1.1 gives the weaker bound $O_{\alpha,d}(r^{d-1}\log(r/\varepsilon)^{2d(\alpha+1)+1})$, or, at best, the

slightly better $O_{\alpha,d}(r^{d-1}\log(r/\varepsilon)^{2d(\alpha+1)})$ by applying other technical results presented below. On the other hand, Theorem 1.1 applies to possibly non-convex, non-coordinate-symmetric and non-dilated domains E, and other regular domains F besides cubes.

When E and F are both cubes, even slightly stronger estimates hold, as follows by tensoring sharp bounds for one dimensional intervals (that match the Landau-Widom asymptotic (1.3)); see, e.g. [18, Theorem 1.4]. This can be used to argue that in (1.5) the factors $|\partial E|$ and $|\partial F|$ cannot be replaced by smaller powers. Similarly, while the power of the logarithm in (1.5) can conceivably be reduced, the logarithmic factor cannot be completely removed.

Our work is in great part motivated by applications where concentration domains may be non-convex, such as the complement of a disk within a two dimensional square; see Sect. 1.5. Such a domain E is allowed by Theorem 1.1 (and Theorems 1.2 and 1.3 below) and has moreover a favorable regularity constant $\kappa_{\partial E}$.

1.2. Discretization of Continuous Domains

Theorem 1.1 is obtained by taking a limit on a more precise result concerning a discrete setting, which is our main focus.

We consider a resolution parameter L>0 and define the discrete Fourier transform $\mathcal{F}_L: L^2((-L/2,L/2)^d) \to \ell^2(L^{-1}\mathbb{Z}^d)$ by

$$\mathcal{F}_{L}f(k/L) = \int_{(-L/2, L/2)^{d}} f(x)e^{-2\pi i x k/L} dx, \quad k \in \mathbb{Z}^{d}.$$
 (1.7)

We think of L as a discretization parameter for an underlying continuous problem. Let us define the *discretization at resolution* L > 0 of a domain $E \subseteq \mathbb{R}^d$ by

$$E_L = L^{-1} \mathbb{Z}^d \cap E. \tag{1.8}$$

Given two compact domains $E \subseteq \mathbb{R}^d$ and $F \subseteq (-L/2, L/2)^d$, consider the *discretized concentration operator* $T: L^2(F) \to L^2(F)$ given by

$$T = \chi_F \mathcal{F}_L^{-1} \chi_{E_L} \mathcal{F}_L. \tag{1.9}$$

Our second result is as follows:

Theorem 1.2. Let $E, F \subseteq \mathbb{R}^d$, $d \ge 2$, be compact domains with maximally Ahlfors regular boundaries with constants $\kappa_{\partial E}, \kappa_{\partial F}$ respectively, and assume that that $|\partial E| |\partial F| \ge 1$.

Fix a discretization resolution $L \ge |\partial E|^{-1/(d-1)}$ such that $F \subseteq (-L/2, L/2)^d$ and consider the discretized concentration operator (1.9) and its eigenvalues $\{\lambda_n : n \in \mathbb{N}\}$.

Indeed, the bound $O_{\alpha,d}(r^{d-1}\log(r/\varepsilon)^{2d(\alpha+1)})$ follows from Theorem 4.1, presented below, which is applicable to the domains in question, together with the discretization argument in the proof of Theorem 1.1.

Then for every $\alpha \in (0, 1/2)$ there exists $A_{\alpha,d} \ge 1$ such that for $\varepsilon \in (0, 1/2)$:

$$\#\{n \in \mathbb{N} : \lambda_n \in (\varepsilon, 1 - \varepsilon)\} \le A_{\alpha, d} \cdot \frac{|\partial E|}{\kappa_{\partial E}} \cdot \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log\left(\frac{|\partial E||\partial F|}{\kappa_{\partial E} \varepsilon}\right)^{2d(1 + \alpha) + 1}.$$
(1.10)

The eigenvalue sequence on the left-hand side of (1.10) is finitely supported (with a bound that depends on the resolution parameter L). In contrast, the right-hand side of (1.10) is independent of L. In applications, this helps capture the transition between analog models and their finite computational counterparts, rigorously showing that the latter remain faithful to the former.

1.3. The Discrete Fourier Transform

Finally, we consider a discrete concentration problem associated with the usual *discrete Fourier transform*, denoted as

$$\mathcal{F}_1: L^2((-1/2, 1/2)^d) \to \ell^2(\mathbb{Z}^d),$$

for consistency with (1.7).

Given a finite set $\Omega \subseteq \mathbb{Z}^d$ and $F \subseteq (-1/2, 1/2)^d$, the *discrete Fourier concentration operator* $T: L^2(F) \to L^2(F)$ is defined as

$$T = \chi_F \mathcal{F}_1^{-1} \chi_\Omega \mathcal{F}_1. \tag{1.11}$$

The discrete boundary of a set $\Omega \subseteq \mathbb{Z}^d$ is given by

$$\partial \Omega = \{ k \in \Omega : \min\{ |j - k| : j \in \mathbb{Z}^d \setminus \Omega \} = 1 \}. \tag{1.12}$$

We say that $\Omega \subseteq \mathbb{Z}^d$ has a *maximally Ahlfors regular boundary* if there exists a constant $\kappa_{\partial\Omega}$ such that

$$\inf_{k \in \partial \Omega} \# \left(\partial \Omega \cap k + [-n/2, n/2)^d \right) \ge \kappa_{\partial \Omega} \cdot n^{d-1}, \quad 1 \le n \le (\# \partial \Omega)^{1/(d-1)}, \ k \in \partial \Omega.$$

(Note the slight notational abuse: though $\Omega \subseteq \mathbb{Z}^d \subseteq \mathbb{R}^d$, the notions of boundary and boundary regularity are to be understood in the discrete sense.)

Our last result is.

Theorem 1.3. Let $d \geq 2$, $\Omega \subseteq \mathbb{Z}^d$ a finite set with maximally Ahlfors regular boundary and constant $\kappa_{\partial\Omega}$. Let $F \subseteq (-1/2, 1/2)^d$ be compact with maximally Ahlfors regular boundary and constant $\kappa_{\partial F}$. Assume that $\#\partial\Omega \cdot |\partial F| \geq 1$, and consider the concentration operator (1.11) and its eigenvalues $\{\lambda_n : n \in \mathbb{N}\}$.

Then, for every $\alpha \in (0, 1/2)$, there exists $A_{\alpha,d} \ge 1$ such that, for $\varepsilon \in (0, 1/2)$

$$\# \big\{ n \in \mathbb{N} : \lambda_n \in (\varepsilon, 1 - \varepsilon) \big\} \le A_{\alpha, d} \cdot \frac{\# \partial \Omega}{\kappa_{\partial \Omega}} \cdot \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log \left(\frac{\# \partial \Omega \cdot |\partial F|}{\kappa_{\partial \Omega} \varepsilon} \right)^{2d(1 + \alpha) + 1}.$$

1.4. One Sided Estimates

We remark that bounds on the number of intermediate eigenvalues, as in Theorems 1.1, 1.2 and 1.3, can be equivalently formulated in terms of the distribution function

$$N_{\varepsilon} := \{ n \in \mathbb{N} : \lambda_n > \varepsilon \}, \quad \varepsilon \in (0, 1).$$

Remark 1.4. For example, for $\varepsilon \in (0, 1)$ under the assumptions of Theorem 1.1 we have

$$\left| \# N_{\varepsilon}(S) - |E| \cdot |F| \right| \leq C_{\alpha,d} \cdot \frac{|\partial E|}{\kappa_{\partial E}} \cdot \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log \left(\frac{|\partial E| |\partial F|}{\kappa_{\partial E} \min\{\varepsilon, 1 - \varepsilon\}} \right)^{2d(1+\alpha)+1}; \tag{1.13}$$

see Sect. 8 for details.

1.5. Significance

Fourier concentration operators arise in problems where functions are assumed to be supported on a certain domain F and their Fourier transforms are known or measured on a second domain E. The insight that the class of such functions is approximately a vector space of dimension $|E| \cdot |F|$ is at the core of many classical and modern physical and signal models, and measurement and estimation methods [16] [12, "A Historical View"].

While in classical applications, such as telecommunications, the concentration domains are rectangles or unions thereof, the increasingly complex geometric nature of data and physical models has sparked great interest in spatial and frequency concentration domains with possibly intricate shapes. To name a few: (a) In geophysics and astronomy the power sprectrum of various quantities of interest is often assumed to be bandlimited to a disk or annulus and needs to be estimated from measurements taken on a domain as irregular as a geological continent [14,32–34]; (b) The Fourier extension algorithm approximates a function on an arbitrary domain by a Fourier series on an enclosing box and crucially exploits the expected moderate size of the plunge region (1.2) [2,29]; (c) Noise statistics are often estimated from those pixels of a square image located outside a central disk, which is assumed to contain the signal of interest — thus, the need to sample pure noise leads one to consider the complement of a disk within a two dimensional square as concentration domain (or, more realistically, a set of grid points within that domain) [3–5,21].

The expected low dimensionality of physical and signal models based on spatiotemporal constraints is often exploited without direct computation of eigendecompositions of Fourier concentration operators (which may in fact be ill-posed in the absence of symmetries). Rather, the expected asymptotic spectral profile of such operators informs strategies based on randomized linear algebra and information theory. The quest to analyze such models and methods has motivated a great amount

of recent research [6,17-20,31,39] which led to deep and far reaching improvements over more classical non-asymptotic results [38]. However, the mentioned literature covers only rectangular or other convex and symmetric domains, which precludes applications involving complex geometries. Due to this limitation, signal and measurement models with a complex geometric nature are often analyzed based on estimates much cruder than (1.5), which have error factors of the order $1/\varepsilon$ instead of $\log(1/\varepsilon)^{\alpha}$ [3,29,33], and thus poorly reflect the remarkable practical effectiveness of low dimensional models. Estimates on measurement/reconstruction complexity, estimation confidence, or approximation/stability trade-offs are as a consequence orders of magnitude too conservative. Our work bridges this theory to practice gap by providing the first spectral deviation estimates for Fourier concentration operators valid without simplifying symmetry or convexity assumptions that also match the precision of what is rigorously known for intervals [6,17– 20,31,36,37,39]. In addition, while Euclidean Fourier concentration operators help analyze computational schemes in their asymptotic continuous limit, our results for the discrete Fourier transform apply, more quantitatively, to finite settings, as they occur in many applications.

1.6. Methods and Related Literature

We work for the most part with the discrete Fourier transform and then obtain consequences for the continuous one by a limiting argument. Theorem 1.2 is thus a more precise and quantitative version of Theorem 1.1, and is proved in two steps. We first revisit Israel's argument [17] and adapt it to prove eigenvalue estimates when one of the domains is a rectangle and the other one is a general multi-dimensional discrete domain (Theorem 4.1 below). These estimates are slightly stronger than those in Theorem 1.2, and the extra precision is exploited in the subsequent step. We follow the method of almost diagonalization with wave-packets, which we achieve, unlike [17], through a redundant system (frame) instead of an orthonormal basis. The versatility of redundant expansions helps us avoid requiring symmetries from the concentration domain.

The second step is a decomposition, rescaling, and dyadic approximation argument, implemented by means of p-Schatten norm estimates for certain Hankel operators, and especially by quantifying those estimates as a function of p, as $p \to 0^+$. In this way we reduce the problem to the case where one of the domains is a rectangle, while relying on the refined estimates derived in the first step.

Our intermediate result, Theorem 4.1, is close in spirit to Theorem 1.1 in [18] (which appeared as we were finishing this article). The estimates in [18], formulated in the context of the continuous Fourier transform and concerning dilated convex domains, are stronger than what follows from Theorem 4.1 in that regime, as [18, Theorem 1.1] involves smaller powers of a certain logarithmic factor (see also Sect. 1.1 and (1.6)). On the other hand, Theorem 4.1 concerns a sufficiently regular possibly non-convex and non-symmetric domain, and covers the discrete Fourier transform (while Theorem 1.2 concerns two such domains).

We also mention our recent work on concentration operators for the short-time Fourier transform [27], that also makes use of Ahlfors regularity and Schatten norm estimates. Though the goals and results are philosophically similar to those in the present article, the settings are rather different from the technical point of view. Indeed, the arguments used in [27] rely on the rapid off-diagonal decay of the reproducing kernel of the range of the short-time Fourier transform, and do not seem to be applicable to Fourier concentration operators.

Finally, we comment on the notion of maximal Ahlfors regularity introduced in this article, which is a variant of more common notions (see Sect. 2.2). While this condition is sufficient to treat all the applications that we are aware of, it is probably non-optimal. Indeed, the eigenvalue estimate (1.5) is robust under certain geometric operations on the concentration domains that may not preserve Ahlfors regularity (see for example the dyadic decomposition and approximation in Sect. 5.2). However, we do not presently know how to formulate a simple and elegant regularity assumption that leads to a significant refinement of our results.

The remainder of the article is organized as follows. Section 2 sets up the notation and provides background on boundary regularity. Section 3 revisits and aptly adapts the technique from [17]. This is used in Sect. 4 to prove Theorem 4.1. Theorem 1.2 is proved in Sect. 5, Theorem 1.1 is proved in Sect. 6, and Theorem 1.3 is proved in Sect. 7. Remark 1.4 is proved in Sect. 8.

2. Preliminaries

2.1. Notation

We shall focus on Theorem 1.2 and set up the notation accordingly. Theorems 1.1 and 1.3 will be obtained afterwards as an application of Theorem 1.2.

We denote cubes by $Q_a = [-a/2, a/2)^d$. The Euclidean norm on \mathbb{R}^d is denoted $|\cdot|$. For two non-negative functions f, g we write $f \lesssim g$ if there exist a constant C such that $f(x) \leq Cg(x)$, and write $f \approx g$ if $f \lesssim g$ and $g \lesssim f$. The implied constant is allowed to depend on the dimension d and the parameter α from Theorems 1.1, 1.2 and 1.3, but not on other parameters.

We enumerate the eigenvalues of a compact self adjoint operator $S: \mathcal{H} \to \mathcal{H}$ acting on a Hilbert space \mathcal{H} as follows:

$$\lambda_k = \inf\{\|S - S_0\| : S_0 \in \mathcal{L}(L^2(\mathbb{R}^d)), \dim(\text{Range}(S_0)) < k\}, \quad k \ge 1.$$
 (2.1)

Then $\{\lambda_k : k \ge 1\} \setminus \{0\} = \sigma(S) \setminus \{0\}$ as sets with multiplicities — see, e.g., [8, Lemma 4.3].

Recall that the *discretization at resolution* L > 0 of a set $E \subseteq \mathbb{R}^d$ is defined by (1.8). We also write

$$E_L^c = L^{-1} \mathbb{Z}^d \smallsetminus E_L$$

and ∂E_L for the points in E_L which are at distance L^{-1} of E_L^c :

$$\partial E_L = \{k/L \in E_L : \min\{|k/L - j/L| : j/L \in E_L^c\} = L^{-1}\}.$$

For L = 1 this is consistent with (1.12).

We will work with the discrete Fourier transform $\mathcal{F}_L: L^2((-L/2, L/2)^d) \to \ell^2(L^{-1}\mathbb{Z}^d)$ given by (1.7) and reserve the notation $\mathcal{F}f$ or \widehat{f} for the continuous Fourier transform (1.4). Note that if $\operatorname{supp}(f) \subseteq (-L/2, L/2)^d$, then $\mathcal{F}_L f(k/L) = \mathcal{F}f(k/L)$ for every $k \in \mathbb{Z}^d$.

We also write $P_{E,L}=\mathcal{F}_L^{-1}\chi_{E_L}\mathcal{F}_L$. For $F\subseteq (-L/2,L/2)^d$ we define the operator $T=T_{E,F,L}:L^2(F)\to L^2(F)$ by

$$T = T_{E,F,L} = \chi_F P_{E,L}$$

and let $\lambda_n = \lambda_n(T)$ denote its eigenvalues as in (2.1). An easy computation shows that

$$T_{t^{-1}E,tF,tL} = \mathcal{M}_{t^{-1}}T_{E,F,L}\mathcal{M}_t, \qquad t > 0,$$

where \mathcal{M}_t denotes the *dilation operator*

$$\mathcal{M}_t f(x) = f(tx).$$

In particular,

$$\lambda_n(T_{t^{-1}E,tF,tL}) = \lambda_n(T_{E,F,L}), \qquad n \in \mathbb{N}. \tag{2.2}$$

2.2. Boundary Regularity

Let us introduce regularity of sets in more generality and discuss a few properties.

An \mathcal{H}^{d-1} -measurable set $X \subseteq \mathbb{R}^d$ is said to be *lower Ahlfors* (d-1)-regular (regular for short) at scale $\eta_X > 0$ if there exists a constant $\kappa_X > 0$ such that

$$\mathcal{H}^{d-1}(X \cap B_r(x)) \ge \kappa_X \cdot r^{d-1}, \quad 0 < r \le \eta_X, \quad x \in X.$$

(See for example [28, Definitions 2.1 and 2.2]). In the literature, Ahlfors regularity is usually stated for $\eta_X \sim \text{diam}(X)$. In contrast, we work with various scales and introduce the term *maximally Ahlfors regular* whenever the regularity condition holds at scale $\eta_X = \mathcal{H}^{d-1}(X)^{1/(d-1)}$. The use of this scale instead of the usual diam(X) allows us to include disconnected sets with distant connected components in our analysis.

Note that if $X \subseteq \mathbb{R}^d$ is regular at scale $\eta_X > 0$ with constant $\kappa_X > 0$ and t > 0, then $tX \subseteq \mathbb{R}^d$ is regular at scale $\eta_{tX} = t\eta_X$ with constant $\kappa_{tX} = \kappa_X$. By differentiation around a point of positive \mathcal{H}^{d-1} -density,

$$\kappa_X \le c_d,$$
 (2.3)

for any regular X of finite \mathcal{H}^{d-1} -measure. We also mention that if X is regular with parameters η_X and κ_X , then choosing an arbitrary $x \in X$ gives

$$\mathcal{H}^{d-1}(X) \ge \mathcal{H}^{d-1}(X \cap B_{\eta_X}(x)) \ge \kappa_X \cdot \eta_X^{d-1}. \tag{2.4}$$

We shall use the following basic result, derived from [7]:

Lemma 2.1. There exists a universal constant $C_d > 0$ such that for every compact set $X \subseteq \mathbb{R}^d$ that is regular at scale $\eta_X > 0$ with constant κ_X and every s > 0,

$$|X + B_s(0)| \le \frac{C_d}{\kappa_X} \cdot \mathcal{H}^{d-1}(X) \cdot s \cdot \left(1 + \frac{s^{d-1}}{\eta_X^{d-1}}\right).$$

Proof. From [7, Theorems 5 and 6] it follows that

$$\mathcal{H}^{d-1}\big(\{x\in\mathbb{R}^d:\ d(x,X)=r\}\big)\leq \frac{C_d}{\kappa_X}\cdot\mathcal{H}^{d-1}(X)\cdot\Big(1+\frac{r^{d-1}}{\eta_Y^{d-1}}\Big),$$

for almost every r > 0, and in addition, $|\nabla d(x, X)| = 1$, for almost every $x \in \mathbb{R}^d$. From this and the coarea formula — see, e.g., [10, Theorem 3.11] — it follows that

$$\begin{aligned} |X + B_s(0)| &= \int_{\mathbb{R}^d} \chi_{[0,s)}(d(x,X)) dx = \int_{\mathbb{R}^d} \chi_{[0,s)}(d(x,X)) |\nabla d(x,X)| dx \\ &= \int_0^s \mathcal{H}^{d-1} \Big(\{x : d(x,X) = r\} \Big) dr \le \frac{C_d}{\kappa_X} \mathcal{H}^{d-1}(X) \int_0^s \Big(1 + \frac{r^{d-1}}{\eta_X^{d-1}} \Big) dr \\ &\le \frac{C_d}{\kappa_X} \mathcal{H}^{d-1}(X) s \Big(1 + \frac{s^{d-1}}{\eta_X^{d-1}} \Big). \end{aligned}$$

Corollary 2.2. For $E \subseteq \mathbb{R}^d$ a compact domain with regular boundary at scale $\eta_{\partial E} \geq 1$ with constant $\kappa_{\partial E}$ and a discretization resolution $L \geq 1$, we have

$$L^{-d} \# E_L \lesssim |E| + \frac{|\partial E|}{\kappa_{2E} L}.$$

In particular, for $d \geq 2$,

$$L^{-d} \# E_L \lesssim \frac{\max\{|\partial E|^{d/(d-1)}, 1\}}{\kappa_{\partial E}}.$$

Proof. Recall that $Q_{L^{-1}}=L^{-1}[-1/2,1/2)^d$ and define $E'_L=\{m\in E_L:m+Q_{L^{-1}}\subseteq E\}$. From Lemma 2.1, we get

$$\begin{split} L^{-d} \# E_L &= \Big| \bigcup_{m \in E_L'} m + Q_{L^{-1}} \Big| + \Big| \bigcup_{m \in E_L \smallsetminus E_L'} m + Q_{L^{-1}} \Big| \leq |E| + |\partial E + B_{L^{-1}\sqrt{d}}(0)| \\ &\lesssim |E| + \frac{|\partial E|}{\kappa_{\partial E} L}. \end{split}$$

Finally, the second inequality follows from the isoperimetric inequality $|E| \lesssim |\partial E|^{d/(d-1)}$ and (2.3).

3. Israel's Argument Revisited

We now revisit the core argument of [17] and aptly adapt it so as to treat multidimensional and discrete domains.

3.1. Israel's Lemma

We need a slight generalization of Lemma 1 in [17], phrasing it in terms of frames rather than orthonormal bases. We include a proof for the sake of completeness

Recall that a frame for a Hilbert space \mathcal{H} is a subset of vectors $\{\phi_i\}_{i\in\mathcal{I}}$ for which the exist constants $0< A, B<\infty$ — called lower and upper frame bounds — such that

$$A \|f\|^2 \le \sum_{i \in \mathcal{I}} |\langle f, \phi_i \rangle|^2 \le B \|f\|^2, \quad f \in \mathcal{H}.$$

If, moreover, A = B, the say that the frame is *tight*.

Lemma 3.1. Let $T: \mathcal{H} \to \mathcal{H}$ be a positive, compact, self-adjoint operator on a Hilbert space \mathcal{H} with $||T|| \le 1$ and eigendecomposition $T = \sum_{n \ge 1} \lambda_n \langle \cdot, f_n \rangle f_n$.

Let $\{\phi_i\}_{i\in\mathcal{I}}$ be a frame of unit norm vectors for \mathcal{H} with lower frame bound A. If $\mathcal{I} = \mathcal{I}_1 \cup \mathcal{I}_2 \cup \mathcal{I}_3$, and

$$\sum_{i \in \mathcal{I}_1} \|T\phi_i\|^2 + \sum_{i \in \mathcal{I}_3} \|(I - T)\phi_i\|^2 \le \frac{A}{2}\varepsilon^2, \tag{3.1}$$

then $\#M_{\varepsilon}(T) \leq \frac{2}{A} \#\mathcal{I}_2$, where $M_{\varepsilon}(T)$ is defined as in (1.2).

Proof. Let $S_{\varepsilon} = \operatorname{span}\{f_n : \lambda_n \in (\varepsilon, 1 - \varepsilon)\}$ and let $P_{\varepsilon} : \mathcal{H} \to S_{\varepsilon}$ denote the orthogonal projection onto S_{ε} . Observe that

$$\varepsilon \|f\| \le \|Tf\|$$
, and $\varepsilon \|f\| \le \|(I-T)f\|$, $f \in S_{\varepsilon}$,

where the second inequality follows from the fact that for $f \in S_{\varepsilon}$ one has $||f|| - ||(I - T)f|| \le ||Tf|| \le (1 - \varepsilon)||f||$. Also note that T and P_{ε} commute since S_{ε} is spanned by a collection of eigenvectors of T. Therefore, by (3.1) we obtain

$$\sum_{i \in \mathcal{I}_1 \cup \mathcal{I}_3} \varepsilon^2 \|P_{\varepsilon} \phi_i\|^2 \le \sum_{i \in \mathcal{I}_1} \|T P_{\varepsilon} \phi_i\|^2 + \sum_{i \in \mathcal{I}_3} \|(I - T) P_{\varepsilon} \phi_i\|^2 \le \frac{A}{2} \varepsilon^2,$$

which implies

$$\sum_{i \in \mathcal{I}_1 \cup \mathcal{I}_3} \|P_{\varepsilon} \phi_i\|^2 \le \frac{A}{2}. \tag{3.2}$$

Using the frame property, we get, for $f \in S_{\varepsilon}$

$$A \|f\|^2 \le \sum_{i \in \mathcal{I}} |\langle f, \phi_i \rangle|^2 = \sum_{i \in \mathcal{I}} |\langle f, P_{\varepsilon} \phi_i \rangle|^2.$$

Now assume that $\dim(S_{\varepsilon}) \geq 1$ (otherwise the result is trivial), take an orthonormal basis $\{\psi_k\}_{k=1}^{\dim(S_{\varepsilon})}$ of S_{ε} , and sum the inequality above over all basis elements to derive

$$A \cdot \# M_{\varepsilon}(T) = A \cdot \dim(S_{\varepsilon}) = A \sum_{k=1}^{\dim(S_{\varepsilon})} \|\psi_k\|^2 \le \sum_{k=1}^{\dim(S_{\varepsilon})} \sum_{i \in \mathcal{I}} |\langle \psi_k, \phi_i \rangle|^2$$

$$= \sum_{i \in \mathcal{I}} \|P_{\varepsilon}\phi_i\|^2 \le \sum_{i \in \mathcal{I}_2} \|\phi_i\|^2 + \frac{A}{2} = \#\mathcal{I}_2 + \frac{A}{2} \le \#\mathcal{I}_2 + \frac{A}{2} \#M_{\varepsilon}(T),$$

where, in the second line, we used (3.2). This shows that $\#M_{\varepsilon}(T) \leq \frac{2}{4} \#\mathcal{I}_2$.

3.2. Local Trigonometric Frames

In this section, we construct a tight frame that allows us to apply Lemma 3.1. Let $\alpha > 0$, and $\theta \in \mathcal{C}^{\infty}(\mathbb{R})$ be such that

- (i) $\theta(x) = 1$, for $x \ge 1$, and $\theta(x) = 0$, for $x \le -1$,
- (ii) $\theta(-x)^2 + \theta(x)^2 = 1$, for every $x \in \mathbb{R}$, (iii) $|D^k \theta(x)| \le C_\alpha^k k^{(1+\alpha)k}$, for all $k \in \mathbb{N}_0$, all $x \in \mathbb{R}$, and a constant $C_\alpha > 0$.

See, for example, [17, Proposition 1] or [15, Chapter 1] for the existence of such a function.

Let W > 0. We decompose the interval $\left(-\frac{W}{2}, \frac{W}{2}\right)$ into disjoint intervals

$$I_j = x_j + \frac{W}{3 \cdot 2^{|j|+1}} [-1, 1), \quad j \in \mathbb{Z},$$

where

$$x_j = \frac{\operatorname{sign}(j)W}{2} \left(1 - \frac{1}{2^{|j|}}\right).$$

Note that $|I_j| = |I_{|j|}| = 2|I_{|j|+1}|$ for every $j \in \mathbb{Z}$. We will also denote $D_j =$ $I_i \cup I_{i+1}$. Now define

$$\theta_j(x) = \theta\left(\frac{2(x-x_j)}{|I_j|}\right)\theta\left(-\frac{2(x-x_{j+1})}{|I_{j+1}|}\right).$$

We have that $\theta_i(x) = 0$ for $x \notin D_i$, and furthermore by properties (i) and (ii)

$$\begin{split} \|\theta_j\|_2^2 &= \int_{I_j} \theta \left(\frac{2(x-x_j)}{|I_j|}\right)^2 dx + \int_{I_{j+1}} \theta \left(-\frac{2(x-x_{j+1})}{|I_{j+1}|}\right)^2 dx \\ &= \frac{|I_j|}{2} \int_{-1}^1 \theta(x)^2 dx + \frac{|I_{j+1}|}{2} \int_{-1}^1 \theta(x)^2 dx = \frac{|D_j|}{2}. \end{split}$$

We define the set of vectors

$$\phi_{j,k}(x) = \sqrt{\frac{2}{|D_j|}} \cdot \theta_j(x) \cdot \exp\left(2\pi i \frac{xk}{|D_j|}\right), \quad j, k \in \mathbb{Z},$$
 (3.3)

and note that $\|\phi_{i,k}\|_2 = 1$.

Lemma 3.2. The family $\{\phi_{j,k}\}_{j,k\in\mathbb{Z}}$ defined in (3.3) forms a tight frame for L^2 (-W/2, W/2) with frame constants A = B = 2.

Proof. Let $f \in L^2(-W/2, W/2)$ and set $f_j := f \chi_{I_j}$ so that $f = \sum_{i \in \mathbb{Z}} f_i$. Since $supp(\theta_i) \subseteq D_i = I_j \cup I_{j+1}$, we observe

$$\sum_{j,k\in\mathbb{Z}} |\langle f,\phi_{j,k}\rangle|^2 = \sum_{j,k\in\mathbb{Z}} |\langle f_j + f_{j+1},\phi_{j,k}\rangle|^2.$$

As $\{|D_j|^{-1/2} \exp(2\pi i kx/|D_j|)\}_{k \in \mathbb{Z}}$ is an orthonormal basis for $L^2(D_j)$, we find

$$\sum_{k \in \mathbb{Z}} |\langle f_j + f_{j+1}, \phi_{j,k} \rangle|^2 = 2 \|(f_j + f_{j+1})\theta_j\|_2^2 = 2 \|f_j\theta_j\|_2^2 + 2 \|f_{j+1}\theta_j\|_2^2.$$

Combining both identities and using property (ii), we conclude

$$\begin{split} \sum_{j,k\in\mathbb{Z}} |\langle f,\phi_{j,k}\rangle|^2 &= 2\sum_{j\in\mathbb{Z}} \left(\|f_j\theta_j\|_2^2 + \|f_{j+1}\theta_j\|_2^2\right) = 2\sum_{j\in\mathbb{Z}} \left(\|f_j\theta_j\|_2^2 + \|f_j\theta_{j-1}\|_2^2\right) \\ &= 2\sum_{j\in\mathbb{Z}} \int_{I_j} |f(x)|^2 \left(\theta\left(\frac{2(x-x_j)}{|I_j|}\right)^2 + \theta\left(-\frac{2(x-x_j)}{|I_j|}\right)^2\right) dx \\ &= 2\sum_{j\in\mathbb{Z}} \|f_j\|_2^2 = 2\|f\|_2^2. \end{split}$$

Let $0 < W_i \le L$, i = 1, ..., d, and consider the rectangle $\prod_{i=1}^{d} (-W_i/2, W_i/2)$. Also, set that

$$W_{\max} := \max_{i=1,\ldots,d} W_i.$$

We define a frame for $L^2(\prod_{i=1}^d (-W_i/2, W_i/2))$ via the tensor product

$$\Phi_{j,k}(x) = \Phi_{j_1,\dots,j_d,k_1,\dots,k_d}(x_1,\dots,x_d) = \phi_{j_1,k_1}(x_1) \cdot \dots \cdot \phi_{j_d,k_d}(x_d),$$

where each family $\{\phi_{j_i,k_i}(x_i)\}_{j_i,k_k\in\mathbb{Z}}$ is the frame for $L^2(-W_i/2,W_i/2)$ given by (3.3). This construction also yields a tight frame with frame bounds equal to 2^d .

3.3. Energy Estimates

Consider

$$\psi_j(x) = \theta_j \Big(|D_j|x + x_j - \frac{W}{3 \cdot 2^{\lfloor j \rfloor + 1}} \Big), \quad x \in \mathbb{R}, \ j \in \mathbb{Z}.$$

A straightforward computation shows that ψ_i is supported on [0, 1] and satisfies $|D^k\psi_j(x)| \leq \widetilde{C_\alpha}^k k^{(1+\alpha)k}$ by property (iii). As shown in [9] or [17, Lemma 4] it thus follows that $|\widehat{\psi_j}(\xi)| \leq A_\alpha \cdot \exp\left(-a_\alpha |\xi|^{(1+\alpha)^{-1}}\right)$. Since $1-\alpha \leq (1+\alpha)^{-1}$, we derive that $t^{(1+\alpha)^{-1}} \ge t^{1-\alpha} - 1$ for $t \ge 0$. Adjusting the constant A_{α} , we therefore

$$|\widehat{\theta_j}(\xi)| \le A_\alpha \cdot |D_j| \cdot \exp\left(-a_\alpha (|D_j| \cdot |\xi|)^{1-\alpha}\right), \quad \xi \in \mathbb{R}.$$
 (3.4)

With this at hand, we estimate the decay of

$$\mathcal{F}(\Phi_{j,k})(\xi) = 2^{d/2} \prod_{i=1}^d |D_{j_i}|^{-1/2} \cdot \widehat{\theta_{j_i}} \left(\xi_i - \frac{k_i}{|D_{j_i}|} \right), \quad \xi \in \mathbb{R}^d.$$

Define

$$M_i = \text{diag}(|D_{i_1}|, ..., |D_{i_d}|) \in \mathbb{R}^{d \times d}$$
.

By (3.4) (possibly enlarging A_{α}) and $|\xi|^{1-\alpha} \leq \sum_{i=1}^{d} |\xi_i|^{1-\alpha}$, it follows

$$|\mathcal{F}(\Phi_{j,k})(\xi)| \le A_{\alpha}^{d} \prod_{i=1}^{d} |D_{j_{i}}|^{1/2} \cdot \exp\left(-a_{\alpha} \left| |D_{j_{i}}| \xi_{i} - k_{i} \right|^{1-\alpha}\right)$$

$$\le A_{\alpha}^{d} \cdot \det(M_{j})^{1/2} \cdot \exp\left(-a_{\alpha} \left| M_{j}(\xi - \xi_{j,k}) \right|^{1-\alpha}\right), \tag{3.5}$$

where $(\xi_{j,k})_i = k_i |D_{j_i}|^{-1}$.

Consider now a compact domain $E \subseteq \mathbb{R}^d$. Let $s \ge 1$ be a parameter that will be determined later. For $j \in \mathbb{Z}^d$ fixed, we cover the index set \mathbb{Z}^d with three subsets as follows:

$$\mathcal{L}_{j}^{\text{low}} := \left\{ k \in \mathbb{Z}^{d} : \operatorname{dist}(k, M_{j} E_{L}^{c}) \geq s \right\};$$

$$\mathcal{L}_{j}^{\text{med}} := \left\{ k \in \mathbb{Z}^{d} : \operatorname{dist}(k, M_{j} E_{L}) < s, \text{ and } \operatorname{dist}(k, M_{j} E_{L}^{c}) < s \right\};$$

$$\mathcal{L}_{j}^{\text{high}} := \left\{ k \in \mathbb{Z}^{d} : \operatorname{dist}(k, M_{j} E_{L}) \geq s \right\}.$$

$$(3.6)$$

(Here, dist is associated with the usual Euclidean distance.) We claim that

$$\mathcal{L}_{j}^{\text{med}} \subseteq \left\{ k \in \mathbb{Z}^{d} : \operatorname{dist}(k, M_{j} \partial E_{L}) < s \right\}; \tag{3.7}$$

let us briefly sketch an argument. Fix $k \in \mathcal{L}_j^{\text{med}}$ and let $k_0 \in M_j L^{-1} \mathbb{Z}^d$ minimize the distance to k. For any point $x \in M_j L^{-1} \mathbb{Z}^d$ with |k-x| < s we can construct a path of adjacent points in $M_j L^{-1} \mathbb{Z}^d$ from x to k_0 whose distance to k is decreasing. In particular, choosing $x_1 \in M_j E_L$ and $x_2 \in M_j E_L^c$ at distance less than s from k, we can connect x_1 and x_2 through a path of adjacent points in $M_j L^{-1} \mathbb{Z}^d$ that stays at distance less than s from k. Necessarily, one of the points in the path must belong to $M_j \partial E_L$, which proves (3.7).

The indices (j, k) with $k \in \mathcal{L}_j^{\text{med}}$, and j satisfying a certain condition specified below (see (3.11)) will play the role of \mathcal{I}_2 in Lemma 3.1, so we need to estimate $\#(\mathcal{L}_j^{med})$.

Lemma 3.3. Let $E \subseteq \mathbb{R}^d$ be a compact domain with regular boundary at scale $\eta_{\partial E} \geq 1$ and constant $\kappa_{\partial E}$. Let $L \geq W_{\max}$ and $s \geq 1$. Then, for all $j \in \mathbb{Z}^d$, we have

$$\#\{k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j \partial E_L) < s\} \lesssim \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} \cdot \frac{|\partial E|}{\kappa_{\partial E}} \cdot s^d.$$

Proof. Since for every $x \in \mathbb{R}^d$ the cube $x + Q_1$ contains one point in \mathbb{Z}^d ,

$$\#\{k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j \partial E_L) < s\} \lesssim s^d \cdot \#\{k \in \mathbb{Z}^d : k \in M_j \partial E_L + Q_1\}.$$
 (3.8)

Second, we claim that for $0 < a \le 1$ and a set $X \subset \mathbb{R}$,

$$\#\{k \in \mathbb{Z} : k \in aX + [-1/2, 1/2)\} \le 2\#\{k \in \mathbb{Z} : k \in X + [-1/2, 1/2)\}.$$
 (3.9)

To see this, consider a map that sends k = ax + t, with $x \in X$ and $t \in [-1/2, 1/2)$ to the unique integer in x + [-1/2, 1/2), denoted k_* (the choice of x and t may be non-unique). Whenever $k_* = k_*'$, it follows that $|k - k'| \le 1$. Hence, at most two k's are mapped into the same k_* .

From (3.8), if we apply (3.9) componentwise (noting that $(M_j)_{i,i} \leq W_{\max}$), we obtain for $s \geq 1$

$$\#\{k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j \partial E_L) < s\} \lesssim s^d \cdot \#\{k \in \mathbb{Z}^d : k \in W_{\max} \partial E_L + Q_1\}.$$

Since for every $x \in \partial E_L$ there exists $x' \in \partial E$ such that $|x - x'| \leq L^{-1}$, and $W_{\text{max}}/L \leq 1$, it follows that

$$\#\{k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j \partial E_L) < s\} \lesssim s^d \cdot \#\{k \in \mathbb{Z}^d : k \in W_{\max} \partial E + Q_3\}$$
$$=: s^d \cdot \#(\mathcal{K}_{W_{\max}}). \tag{3.10}$$

Now let $k \in \mathcal{K}_{W_{\max}}$. There exists at least one point $x_k \in \partial E$ such that $k \in W_{\max} x_k + Q_3$. In particular, we have that $x_k \in W_{\max}^{-1} k + Q_{4/W_{\max}}$. Therefore, for every $k \in \mathcal{K}_{W_{\max}}$ we get by regularity of ∂E

$$\kappa_{\partial E} \cdot \min\{W_{\max}^{-1}, \eta_{\partial E}\}^{d-1} \leq \mathcal{H}^{d-1} \left(\partial E \cap B_{1/W_{\max}}(x_k)\right)$$

$$\leq \mathcal{H}^{d-1} \left(\partial E \cap x_k + Q_{2/W_{\max}}\right)$$

$$\leq \mathcal{H}^{d-1} \left(\partial E \cap W_{\max}^{-1} k + Q_{6/W_{\max}}\right).$$

So,

$$\begin{split} \kappa_{\partial E} \cdot \min \left\{ W_{\text{max}}^{-1}, \, \eta_{\partial E} \right\}^{d-1} \cdot \# (\mathcal{K}_{W_{\text{max}}}) & \leq \sum_{k \in \mathcal{K}_{W_{\text{max}}}} \mathcal{H}^{d-1} \big(\partial E \cap W_{\text{max}}^{-1} k + Q_{6/W_{\text{max}}} \big) \\ & \lesssim \sum_{k \in \mathbb{Z}^d} \mathcal{H}^{d-1} \big(\partial E \cap W_{\text{max}}^{-1} k + Q_{1/W_{\text{max}}} \big) \\ & = \mathcal{H}^{d-1} (\partial E). \end{split}$$

Plugging this estimate into (3.10) completes the proof.

Next, for a compact domain $E \subseteq \mathbb{R}^d$ and a parameter $s \ge 1$ we recall the sets (3.6), introduce a second auxiliary parameter $0 < \delta < 1$, and define the following covering of \mathbb{Z}^{2d} :

$$\Gamma^{\text{low}} := \left\{ (j, k) : \min_{i} |D_{j_{i}}| \geq \delta, \ k \in \mathcal{L}_{j}^{\text{low}} \right\};
\Gamma^{\text{med}} := \left\{ (j, k) : \min_{i} |D_{j_{i}}| \geq \delta, \ k \in \mathcal{L}_{j}^{\text{med}} \right\};
\Gamma^{\text{high}} := \left\{ (j, k) : \min_{i} |D_{j_{i}}| \geq \delta, \ k \in \mathcal{L}_{j}^{\text{high}} \right\} \cup \left\{ (j, k) : \min_{i} |D_{j_{i}}| < \delta, \ k \in \mathbb{Z}^{d} \right\}.$$
(3.11)

Lemma 3.4. Under the conditions of Lemma 3.3, let $0 < \delta < 1$ and consider the set Γ^{med} from (3.11). Then

$$\#(\Gamma^{\mathrm{med}}) \lesssim \max\{W_{\mathrm{max}}, 1/\eta_{\partial E}\}^{d-1} \cdot \frac{|\partial E|}{\kappa_{\partial E}} \cdot \max\{\log(W_{\mathrm{max}}/\delta), 1\}^d \cdot s^d.$$

Proof. By (3.7) and Lemma 3.3 it follows

$$\#(\Gamma^{\mathrm{med}}) = \sum_{\substack{j \in \mathbb{Z}^d \\ \min|D_{j_i}| \geq \delta}} \#(\mathcal{L}_j^{\mathrm{med}}) \lesssim \sum_{\substack{j \in \mathbb{Z}^d \\ \min|D_{j_i}| \geq \delta}} \max\{W_{\mathrm{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} s^d.$$

In each coordinate, we have that the number of intervals D_{j_i} for which $|D_{j_i}| \ge \delta$ is bounded by $C \max\{\log(W_{\max}/\delta), 1\}$. Hence, we arrive at

$$\#(\Gamma^{\text{med}}) \lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \max\{\log(W_{\text{max}}/\delta), 1\}^d s^d.$$

Lemma 3.5. Let $d \geq 2$, $L \geq \max\{W_{\max}, 1\}$, and $E \subseteq \mathbb{R}^d$ be a compact domain with regular boundary at scale $\eta_{\partial E} \geq 1$ with constant $\kappa_{\partial E}$ and such that $|\partial E| \geq 1$. Let $s \geq 1$ and $\delta \in (0, 1)$ be parameters and consider the sets from (3.11). Then there exists a constant $c = c_{\alpha} > 0$ such that

$$\begin{split} \sum_{(j,k)\in\Gamma^{\text{low}}} L^{-d} \sum_{m\in E_L^c} |\mathcal{F}(\Phi_{j,k})(m)|^2 \\ \lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \cdot \frac{|\partial E|}{\kappa_{\partial E}} \cdot \exp\left(-cs^{1-\alpha}\right) \max\{\log(W_{\text{max}}/\delta), 1\}^d, \end{split} \tag{3.12}$$

and

$$\begin{split} \sum_{(j,k)\in\Gamma^{\text{high}}} L^{-d} \sum_{m\in E_L} |\mathcal{F}(\Phi_{j,k})(m)|^2 \lesssim & \frac{\max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1}}{\kappa_{\partial E}} \cdot \left(|\partial E|^{\frac{d}{d-1}} \cdot \delta \right. \\ & + |\partial E| \cdot \exp\left(-cs^{1-\alpha}\right) \cdot \max\{\log(W_{\text{max}}/\delta), 1\}^d\right). \end{split} \tag{3.13}$$

Proof. For $j \in \mathbb{Z}^d$ and $l \in \mathbb{N}_0$ we set

$$\mathcal{L}_{i,l}^{\text{low}} = \left\{ k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j E_L^c) \in [s2^l, s2^{l+1}) \right\},\,$$

and

$$\mathcal{L}_{j,l}^{\text{high}} = \left\{ k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j E_L) \in [s2^l, s2^{l+1}) \right\}.$$

Notice that

$$\mathcal{L}_{j,l}^{\text{low}} \cup \mathcal{L}_{j,l}^{\text{high}} \subseteq \left\{ k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j E_L) < s2^{l+1}, \text{ and } \operatorname{dist}(k, M_j E_L^c) < s2^{l+1} \right\}$$
$$\subseteq \left\{ k \in \mathbb{Z}^d : \operatorname{dist}(k, M_j \partial E_L) < s2^{l+1} \right\},$$

where the last step follows as in (3.7). From Lemma 3.3 we get

$$\#(\mathcal{L}_{j,l}^{\text{low}}), \#(\mathcal{L}_{j,l}^{\text{high}}) \lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} s^d 2^{dl}. \tag{3.14}$$

From (3.5) it follows that if $k \in \mathcal{L}_{i,l}^{low}$

$$L^{-d} \sum_{m \in E_L^c} |\mathcal{F}(\Phi_{j,k})(m)|^2 \le L^{-d} \sum_{m \in E_L^c} A_\alpha^{2d} \det(M_j) \exp\left(-2a_\alpha |M_j(m - \xi_{j,k})|^{1-\alpha}\right)$$

$$\le A_\alpha^{2d} L^{-d} \det(M_j) \sum_{m' \in M_j E_L^c} \exp\left(-2a_\alpha |m' - k|^{1-\alpha}\right)$$

$$\lesssim \int_{\{|x| \ge s2^l\}} \exp\left(-2a_\alpha |x|^{1-\alpha}\right) dx$$

$$\le \exp\left(-c(s2^l)^{1-\alpha}\right), \tag{3.15}$$

where c can for example be chosen as a_{α} . A similar argument also shows that for $k \in \mathcal{L}_{i,l}^{\text{high}}$,

$$L^{-d} \sum_{m \in E_I} |\mathcal{F}(\Phi_{j,k})(m)|^2 \lesssim \exp\left(-c(s2^l)^{1-\alpha}\right).$$

As $\mathcal{L}_j^{low} = \bigcup_{l \in \mathbb{N}_0} \mathcal{L}_{j,l}^{low}$, it follows from (3.14) and (3.15) that

$$\begin{split} &\sum_{(j,k)\in\Gamma^{\text{low}}} L^{-d} \sum_{m\in E_L^c} |\mathcal{F}(\Phi_{j,k})(m)|^2 = \sum_{j\in\mathbb{Z}^d} \sum_{l\in\mathbb{N}_0} \sum_{k\in\mathcal{L}^{\text{low}}_{j,l}} L^{-d} \sum_{m\in E_L^c} |\mathcal{F}(\Phi_{j,k})(m)|^2 \\ &\lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \sum_{\substack{j\in\mathbb{Z}^d\\ \min|D_{j_i}|\geq \delta}} \sum_{l\in\mathbb{N}_0} (s2^l)^d \exp\left(-c(s2^l)^{1-\alpha}\right) \\ &\lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \sum_{\substack{j\in\mathbb{Z}^d\\ \min|D_{j_i}|\geq \delta}} \exp\left(-c's^{1-\alpha}\right) \\ &\lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \exp\left(-c's^{1-\alpha}\right) \max\{\log(W_{\text{max}}/\delta), 1\}^d, \end{split}$$

which completes the proof of (3.12). Again, we can use an analogous reasoning to show that

$$\sum_{\substack{j \in \mathbb{Z}^d \\ \min |D_{j_i}| \ge \delta}} \sum_{k \in \mathcal{L}_j^{\text{high}}} L^{-d} \sum_{m \in E_L} |\mathcal{F}(\Phi_{j,k})(m)|^2$$

$$\lesssim \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \exp\left(-c's^{1-\alpha}\right) \max\{\log(W_{\max}/\delta), 1\}^d.$$

Now suppose that $j \in \mathbb{Z}^d$ is such that $\min_{1 \le i \le d} |D_{j_i}| < \delta$. For every $m \in \mathbb{Z}^d$ we can uniformly bound the following series

$$\sum_{k\in\mathbb{Z}^d} \exp\left(-2a_\alpha |M_j(m-\xi_{j,k})|^{1-\alpha}\right) = \sum_{k\in\mathbb{Z}^d} \exp\left(-2a_\alpha |M_jm-k|^{1-\alpha}\right) \le C.$$

Since $det(M_j) = |D_j|$, where $D_j = D_{j_1} \times ... \times D_{j_d}$, we thus get by (3.5)

$$\begin{split} \sum_{j \in \mathbb{Z}^d \atop \min |D_{j_i}| < \delta} \sum_{k \in \mathbb{Z}^d} L^{-d} \sum_{m \in E_L} \left| \mathcal{F}(\Phi_{j,k})(m) \right|^2 &\leq C \sum_{j \in \mathbb{Z}^d \atop \min |D_{j_i}| < \delta} L^{-d} \sum_{m \in E_L} \det(M_j) \\ &\leq C L^{-d} \# E_L \sum_{j \in \mathbb{Z}^d \atop \min |D_{j_i}| < \delta} |D_j| \\ &\lesssim \frac{|\partial E|^{d/(d-1)}}{\kappa_{\partial E}} \sum_{j \in \mathbb{Z}^d \atop \min |D_{j_i}| < \delta} |D_j|, \end{split}$$

where in the last inequality we used Corollary 2.2. Finally,

$$\sum_{\substack{j \in \mathbb{Z}^d \\ \min|D_{j_i}| < \delta}} |D_j| \le \sum_{i=1}^d \sum_{\substack{j \in \mathbb{Z}^d \\ |D_{j_i}| < \delta}} |D_j| \le \sum_{i=1}^d W_{\max}^{d-1} 4\delta$$
$$\lesssim \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} \delta,$$

where we used that each interval D_{j_i} is at most at $|D_{j_i}| < \delta$ distance from the boundary of $(-W_i/2, W_i/2)$ and consequently the union of all such intervals lie in the complement of a rectangle of side-lengths $W_i - 4\delta$. This concludes the proof of (3.13).

4. General Domain Versus Rectangle

In this section, we prove the following variant of Theorem 1.2 for F a rectangle.

Theorem 4.1. Let $L \ge 1$ be a discretization resolution, $d \ge 2$, and $E \subseteq \mathbb{R}^d$ be a compact domain with regular boundary at scale $\eta_{\partial E} \ge 1$ with constant $\kappa_{\partial E}$ and such that $|\partial E| \ge 1$. For $0 < W_i \le L$, i = 1, ..., d, take $F = \prod_{i=1}^d (-W_i/2, W_i/2)$ and denote $W_{\text{max}} = \max_i W_i$.

Then for every $\alpha \in (0, 1/2)$ there exists $A_{\alpha,d} \ge 1$ such that for $\varepsilon \in (0, 1/2)$:

$$\begin{split} &\# \big\{ n \in \mathbb{N} : \ \lambda_n \in (\varepsilon, 1 - \varepsilon) \big\} \\ &\leq A_{\alpha, d} \cdot \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} \cdot \frac{|\partial E|}{\kappa_{\partial E}} \cdot \log \left(\frac{\max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} |\partial E|}{\kappa_{\partial E} \ \varepsilon} \right)^{2d(1 + \alpha)}. \end{split}$$

Proof. We adopt all the notation of Sect. 3. Fix parameters $s \ge 1$, $\delta \in (0, 1)$ and consider the sets from (3.11).

Observe that for $f \in L^2(F)$ one has

$$||Tf||_2^2 = ||\chi_F P_{E,L} f||_2^2 \le ||P_{E,L} f||_2^2 = L^{-d} \sum_{m \in E_I} |\widehat{f}(m)|^2,$$

and

$$\|f - Tf\|_2^2 = \|\chi_F f - \chi_F P_{E,L} f\|_2^2 \le \|(I - P_{E,L}) f\|_2^2 = L^{-d} \sum_{m \in E_T^c} |\widehat{f}(m)|^2.$$

By Lemma 3.5 it thus follows that

$$\begin{split} \sum_{(j,k)\in\Gamma^{\text{low}}} & \|(I-T)\Phi_{j,k}\|_2^2 + \sum_{(j,k)\in\Gamma^{\text{high}}} \|T\Phi_{j,k}\|_2^2 \\ & \leq C \frac{\max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1}}{\kappa_{\partial E}} \Big(|\partial E| \exp\big(-cs^{1-\alpha}\big) \max\{\log(W_{\text{max}}/\delta), 1\}^d \\ & + |\partial E|^{d/(d-1)}\delta \Big), \end{split} \tag{4.1}$$

where the constants depend only on α and d, and C can be taken ≥ 2 .

At last, we can now specify the parameters δ and s in order for the sets Γ^{low} , Γ^{med} and Γ^{high} to play the role of \mathcal{I}_1 , \mathcal{I}_2 and \mathcal{I}_3 in Lemma 3.1. We take

$$\delta = \frac{\kappa_{\partial E} \, \varepsilon^2}{C \, \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} |\partial E|^{d/(d-1)}}.$$

This ensures that

$$\frac{C \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} |\partial E|^{d/(d-1)}}{\kappa_{\partial E}} \delta \le \varepsilon^2. \tag{4.2}$$

while (2.4) shows that $\delta \in (0, 1)$.

In addition, we shall select s such that

$$C \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \exp\left(-cs^{1-\alpha}\right) \max\{\log(W_{\max}/\delta), 1\}^d \le \varepsilon^2.$$
(4.3)

This condition on s is equivalent to

$$s \ge \left(\frac{1}{c} \log \left(\frac{C \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} |\partial E| \max\{\log(W_{\max}/\delta), 1\}^d}{\kappa_{\partial E} \ \varepsilon^2}\right)\right)^{1/(1-\alpha)}. \tag{4.4}$$

Denote

$$u = \max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} |\partial E|/\kappa_{\partial E},$$

which satisfies $u \ge 1$ by (2.4), and note that

$$\log(W_{\text{max}}/\delta) \le \log\left(\frac{C\kappa_{\partial E}^{1/(d-1)}u^{d/(d-1)}}{\varepsilon^2}\right) \lesssim \log(u/\varepsilon),\tag{4.5}$$

where in the last step we used (2.3). Thus, we can bound the right-hand side of (4.4) by

$$\left(\frac{1}{c}\log\left(\frac{C'u\log(u/\varepsilon)^d}{\varepsilon^2}\right)\right)^{1/(1-\alpha)}\lesssim \log(u/\varepsilon)^{1/(1-\alpha)}.$$

In particular, (4.4) is satisfied if

$$s = A_{\alpha,d} \log \left(\frac{\max\{W_{\max}, 1/\eta_{\partial E}\}^{d-1} |\partial E|}{\kappa_{\partial E} \varepsilon} \right)^{1/(1-\alpha)}$$
(4.6)

for an adequate constant $A_{\alpha,d}$. Moreover, we can guarantee that $s \geq 1$, since, by (2.4), the term inside the logarithm in (4.6) is \geq 2. From (4.1), (4.2), (4.3), Lemma 3.1 and Lemma 3.4.

$$\begin{split} \#M_{\varepsilon}(T) &\leq 2^{1-d} \#(\Gamma^{\text{med}}) \\ &\lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \max\{\log(W_{\text{max}}/\delta), 1\}^{d} s^{d} \\ &\lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \log \left(\frac{\max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} |\partial E|}{\kappa_{\partial E} \varepsilon}\right)^{d/(1-\alpha)+d} \\ &\lesssim \max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} \frac{|\partial E|}{\kappa_{\partial E}} \log \left(\frac{\max\{W_{\text{max}}, 1/\eta_{\partial E}\}^{d-1} |\partial E|}{\kappa_{\partial E} \varepsilon}\right)^{2d(1+\alpha)}, \end{split}$$

where in the third step we used (4.5) and the definition of s.

5. Eigenvalue Estimates for Two Domains

5.1. Schatten Quasi-Norm Estimates

For $0 , and <math>\varepsilon > 0$, define the auxiliary function $g = g_{p,\varepsilon} : [0,1] \to \mathbb{R}$ given by

$$g(t) = \left(\frac{t - t^2}{s - s^2}\right)^p.$$

Note that since $\chi_{(\varepsilon,1-\varepsilon)} \leq g$, for a positive operator $0 \leq S \leq 1$,

$$#M_{\varepsilon}(S) = \operatorname{tr}(\chi_{(\varepsilon, 1-\varepsilon)}(S)) \le \operatorname{tr}(g(S)) = \frac{\|S - S^2\|_p^p}{(\varepsilon - \varepsilon^2)^p}, \tag{5.1}$$

where $\|\cdot\|_p$, 0 , denotes the Schatten quasi-norm. The next lemma showsthat upper bounds for the left-hand side of the last inequality can be transferred to the right-hand side without much loss.

Lemma 5.1. Suppose that for a positive operator $0 \le S \le 1$ there are constants C, D, a > 0 such that for every $\varepsilon \in (0, 1/2)$,

$$\#M_{\varepsilon}(S) \leq C(D + \log(\varepsilon^{-1}))^a$$
.

Then, for every $0 there is a constant <math>C_a > 0$ such that

$$||S - S^2||_p^p \le C_a C(D + p^{-1})^a$$
.

Page 21 of

Proof. By the symmetry of the function $h(x) = x - x^2$ around 1/2, for $0 \le x \le 1$,

$$h(x)^{p} = \int_{0}^{\min\{x, 1-x\}} (h^{p})'(t)dt = \int_{0}^{1/2} \chi_{(t, 1-t)}(x) p h^{p-1}(t) h'(t) dt$$

$$\leq \int_{0}^{1/2} \chi_{(t, 1-t)}(x) p t^{p-1} dt.$$

By a monotone convergence argument we get that

$$||S - S^{2}||_{p}^{p} \leq \int_{0}^{1/2} M_{t}(S) p t^{p-1} dt \leq C \int_{0}^{1} (D + \log(t^{-1}))^{a} p t^{p-1} dt$$

$$= C \int_{0}^{\infty} (D + u/p)^{a} e^{-u} du$$

$$\leq C (D + 1/p)^{a} + C \int_{1}^{\infty} (D + u/p)^{a} e^{-u} du$$

$$\leq C (D + 1/p)^{a} + C (D + 1/p)^{a} \int_{1}^{\infty} u^{a} e^{-u} du$$

$$\leq (1 + \Gamma(a+1)) C (D + 1/p)^{a}.$$

5.2. Decomposition of the Domain and Hankel Operators

In what follows, we let $F \subseteq (-L/2, L/2)^d$ be a compact domain with regular boundary at scale $\eta_{\partial F} = |\partial F|^{1/(d-1)} > 1$ with constant $\kappa_{\partial F}$. We also assume that $\mathcal{H}^d(\partial F) = 0$ (as otherwise the bounds that we shall derive are trivial). We construct two auxiliary sets $F^- \subseteq F \subseteq F^+$ which will be dyadic approximations of F from the inside and outside by cubes of side-length at least 1. More precisely, let

$$F = \bigcup_{k \in \mathbb{Z}} \bigcup_{j \in J_k} Q_{k,j}, \quad \text{(up to a null measure set)}$$

be a dyadic decomposition of F in pairwise disjoint cubes of the form $Q_{k,j}$ = $Q_{2^k} + 2^k j$ with $k \in \mathbb{Z}$ and $j \in J_k \subseteq \mathbb{Z}^d$, that are maximal (i.e., they are not contained in a larger dyadic cube included in F). We define

$$F^- = \bigcup_{k \ge 0} \bigcup_{j \in J_k} Q_{k,j}.$$

For F^+ we add cubes of length 1 to fully cover F and intersect them with $(-L/2, L/2)^d$. The result is a covering of F that combines the cubes from F^- with rectangles of maximal side-length < 1. More precisely, define

$$V = \{ v \in \mathbb{Z}^d : (F \setminus F^-) \cap (Q_1 + v) \neq \emptyset \},$$

and

$$F^{+} = F^{-} \cup \bigcup_{v \in V} \left((Q_{1} + v) \cap (-L/2, L/2)^{d} \right) =: F^{-} \cup \bigcup_{v \in V} R_{v}.$$

We write T^{\pm} for $T_{E,F^{\pm},L}$.

We shall invoke Theorem 4.1 and apply it to each rectangle in the decomposition of F^- and F^+ . This is allowed because translating the rectangles, or removing their null \mathcal{H}^d -measure boundaries, does not affect the eigenvalue estimates in question.

For a set $K \subseteq (-L/2, L/2)^d$ define the *Hankel operator* on $L^2((-L/2, L/2)^d)$ by

$$H_K = (I - P_{E,L})\chi_K P_{E,L}$$

and write $H^{\pm} = H_{F^{\pm}}$. The Hankel operator H_K is related to the *Toeplitz operator* $P_{E,L} \chi_K P_{E,L}$ in a similar way as the correspondingly named operators on Hardy spaces [30]. In particular, we note that

$$(H_K)^* H_K = P_{E,L} \chi_K P_{E,L} - P_{E,L} \chi_K P_{E,L} \chi_K P_{E,L}$$

= $P_{E,L} \chi_K P_{E,L} - (P_{E,L} \chi_K P_{E,L})^2$.

Since $P_{E,L}\chi_K P_{E,L}$ and T_K share the same non-zero eigenvalues, for p > 0,

$$||T_K - (T_K)^2||_p^p = ||H_K||_{2p}^{2p}.$$

Recall that for two operators S_1 , S_2 in the p-Schatten class, 0 , one has

$$||S_1 + S_2||_p^p \le ||S_1||_p^p + ||S_2||_p^p.$$

Lemma 5.2. Let $L \ge 1$, $d \ge 2$, and E, $F \subseteq \mathbb{R}^d$ be compact domains with regular boundaries at scales $\eta_{\partial E} \ge 1$, $\eta_{\partial F} = |\partial F|^{1/(d-1)} \ge 1$, with constants $\kappa_{\partial E}$, $\kappa_{\partial F}$ respectively. Assume also that $|\partial E| \ge 1$ and $F \subseteq (-L/2, L/2)^d$.

Then for $\varepsilon \in (0, 1/2)$, we have

$$\#M_{\varepsilon}(T^{\pm}) \lesssim \frac{|\partial E|}{\kappa_{\partial E}} \cdot \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log \left(\frac{|\partial E||\partial F|}{\kappa_{\partial E} \varepsilon}\right)^{2d(1+\alpha)+1}.$$

Proof. If $k \in \mathbb{Z}$ is such that $J_k \neq \emptyset$, then there is a cube of length 2^k included in F. In particular, the projection of ∂F onto the hyperplane $\{x_1 = 0\}$ contains a (d-1)-dimensional cube of length 2^k and therefore $2^{k(d-1)} \leq |\partial F|$. The maximality of the dyadic decomposition of F implies that $Q_{j,k} \subseteq \partial F + B_{\sqrt{d}2^{k+1}}(0)$ for $j \in J_k$. From Lemma 2.1 and the fact that $\eta_{\partial F} = |\partial F|^{1/(d-1)}$, we thus derive

$$2^{dk} \# J_k \le |\partial F + B_{\sqrt{d}2^{k+1}}(0)| \lesssim 2^k \frac{|\partial F|}{\kappa_{\partial F}} \left(1 + \frac{2^{k(d-1)}}{|\partial F|} \right) \lesssim 2^k \frac{|\partial F|}{\kappa_{\partial F}}. \tag{5.2}$$

Similarly,

$$\#V \le |\partial F + B_{\sqrt{d}}(0)| \lesssim \frac{|\partial F|}{\kappa_{\partial F}}.$$
 (5.3)

31

246.33 Fage 23 01

For $0 < 2p \le 1$, and $\varepsilon \in (0, 1/2)$, by (5.1), we get

$$\begin{split} \# M_{\varepsilon}(T^{+}) &\leq \frac{\|T^{+} - (T^{+})^{2}\|_{p}^{p}}{(\varepsilon - \varepsilon^{2})^{p}} = \frac{\|H^{+}\|_{2p}^{2p}}{(\varepsilon - \varepsilon^{2})^{p}} \\ &\leq (2/\varepsilon)^{p} \sum_{k \geq 0} \sum_{j \in J_{k}} \|H_{Q_{k,j}}\|_{2p}^{2p} + (2/\varepsilon)^{p} \sum_{v \in V} \|H_{R_{v}}\|_{2p}^{2p} \\ &\lesssim \varepsilon^{-p} \sum_{k \geq 0} \sum_{j \in J_{k}} \|T_{Q_{k,j}} - T_{Q_{k,j}}^{2}\|_{p}^{p} + \varepsilon^{-p} \sum_{v \in V} \|T_{R_{v}} - T_{R_{v}}^{2}\|_{p}^{p}. \end{split}$$

We invoke Theorem 4.1 to obtain spectral deviation estimates for each operator $T_{Q_{k,j}}$. These imply that we can apply Lemma 5.1 to $T_{Q_{k,j}}$ with

$$C \lesssim 2^{k(d-1)} \frac{|\partial E|}{\kappa_{\partial E}}, \quad \text{ and } \quad D = \log \left(2^{k(d-1)} \frac{|\partial E|}{\kappa_{\partial E}} \right).$$

Similarly, the same holds for T_{R_v} with k = 1. Choosing $p = \log(2) (2 \log(\varepsilon^{-1}))^{-1}$ (which ensures that $2p \le 1$ for every $\varepsilon \in (0, 1/2)$) thus yields

$$\begin{split} \# M_{\varepsilon}(T^+) &\lesssim \frac{|\partial E|}{\kappa_{\partial E}} \Biggl(\sum_{k \geq 0} \sum_{j \in J_k} 2^{k(d-1)} \log \Biggl(\frac{|\partial E| 2^{k(d-1)}}{\kappa_{\partial E} \; \varepsilon} \Biggr)^{2d(1+\alpha)} + \# V \cdot \log \Biggl(\frac{|\partial E|}{\kappa_{\partial E} \; \varepsilon} \Biggr)^{2d(1+\alpha)} \Biggr) \\ &\lesssim \frac{|\partial E|}{\kappa_{\partial E}} \frac{|\partial F|}{\kappa_{\partial F}} \sum_{\substack{k \geq 0 \\ 2^{k(d-1)} \leq |\partial F|}} \log \Biggl(\frac{|\partial E| 2^{k(d-1)}}{\kappa_{\partial E} \; \varepsilon} \Biggr)^{2d(1+\alpha)} \\ &= \frac{|\partial E|}{\kappa_{\partial E}} \frac{|\partial F|}{\kappa_{\partial F}} \sum_{\substack{0 \leq k \leq \left | \frac{\log(|\partial F|)}{\log(2)(d-1)} \right |}} \Biggl(\log \left(\frac{|\partial E|}{\kappa_{\partial E} \; \varepsilon} \right) + (d-1) \log(2)k \Biggr)^{2d(1+\alpha)} \;, \end{split}$$

where in the second-to-last step we used (5.2), (5.3), and the fact that $2^{k(d-1)} \le |\partial F|$ whenever $J_k \ne \emptyset$. Finally, noting that for C, D, a > 0,

$$\sum_{k=0}^{N} (C+Dk)^{a} \le \int_{0}^{N+1} (C+Dx)^{a} dx \le \frac{(C+D(N+1))^{a+1}}{D(a+1)},$$

we get,

$$\#M_{\varepsilon}(T^{+}) \lesssim \frac{|\partial E|}{\kappa_{\partial E}} \frac{|\partial F|}{\kappa_{\partial F}} \log \left(\frac{|\partial E||\partial F|}{\kappa_{\partial E} \varepsilon} \right)^{2d(1+\alpha)+1}.$$

The same argument applies to $\#M_{\varepsilon}(T^{-})$, and yields the desired conclusion.

5.3. The Transition Index

The following estimate is part of the proof of [1, Theorem 1.5] (see also [27, Lemma 4.3]) and allows us to find the index where eigenvalues cross the 1/2 threshold. We include a proof for the sake of completeness.

Lemma 5.3. For any trace class operator $0 \le S \le 1$,

(i)
$$\lambda_n \leq \frac{1}{2}$$
, for every $n \geq \lceil \operatorname{tr}(S) \rceil + \max\{2 \operatorname{tr}(S - S^2), 1\}$;
(ii) $\lambda_n \geq \frac{1}{2}$, for every $1 \leq n \leq \lceil \operatorname{tr}(S) \rceil - \max\{2 \operatorname{tr}(S - S^2), 1\}$.

Proof. First notice that if S is an orthogonal projection, then the result holds trivially, so we can assume otherwise. In particular, we have that $tr(S - S^2) > 0$.

Set $K = \lceil \operatorname{tr}(S) \rceil$ and write

$$\operatorname{tr}(S) - \operatorname{tr}(S^{2}) = \sum_{n=1}^{\infty} \lambda_{n} (1 - \lambda_{n})$$

$$= \sum_{n=1}^{K} \lambda_{n} (1 - \lambda_{n}) + \sum_{n=K+1}^{\infty} \lambda_{n} (1 - \lambda_{n})$$

$$\geq \lambda_{K} \sum_{n=1}^{K} (1 - \lambda_{n}) + (1 - \lambda_{K}) \sum_{n=K+1}^{\infty} \lambda_{n}$$

$$= \lambda_{K} K - \lambda_{K} \sum_{n=1}^{K} \lambda_{n} + (1 - \lambda_{K}) \left(\operatorname{tr}(S) - \sum_{n=1}^{K} \lambda_{n} \right)$$

$$= \lambda_{K} K + (1 - \lambda_{K}) \operatorname{tr}(S) - \sum_{n=1}^{K} \lambda_{n}$$

$$= \operatorname{tr}(S) - \sum_{n=1}^{K} \lambda_{n} + \lambda_{K} (K - \operatorname{tr}(S)).$$

Hence

$$\sum_{n=K+1}^{\infty} \lambda_n = \operatorname{tr}(S) - \sum_{n=1}^{K} \lambda_n \le \operatorname{tr}(S) - \operatorname{tr}(S^2), \tag{5.4}$$

and

$$\sum_{n=1}^{K-1} (1 - \lambda_n) = \operatorname{tr}(S) - \sum_{n=1}^{K} \lambda_n + \lambda_K (K - \operatorname{tr}(S)) - (1 - \lambda_K)(1 + \operatorname{tr}(S) - K)$$

$$\leq \operatorname{tr}(S) - \sum_{n=1}^{K} \lambda_n + \lambda_K (K - \operatorname{tr}(S)) \leq \operatorname{tr}(S) - \operatorname{tr}(S^2). \tag{5.5}$$

Now let $j \in \mathbb{N}$ such that $j > 2(\operatorname{tr}(S) - \operatorname{tr}(S^2))$ and consider k = K + j. It follows from (5.4) that

$$2(\operatorname{tr}(S) - \operatorname{tr}(S^2)) \cdot \lambda_k \le j \cdot \lambda_{K+j} \le \sum_{n=K+1}^{\infty} \lambda_n \le \operatorname{tr}(S) - \operatorname{tr}(S^2),$$

which shows $\lambda_k \le 1/2$ as $0 < \operatorname{tr}(S) - \operatorname{tr}(S^2) < \infty$; this proves part (i).

For part (ii), if $1 < k = K - i < K - 2(tr(S) - tr(S^2))$ for $i \in \mathbb{N}$, then (5.5) implies

$$2(\operatorname{tr}(S) - \operatorname{tr}(S^2)) \cdot (1 - \lambda_k) \le j \cdot (1 - \lambda_{K-j}) \le \sum_{n=1}^{K-1} (1 - \lambda_n) \le \operatorname{tr}(S) - \operatorname{tr}(S^2),$$

yielding $\lambda_k \geq 1/2$. This completes the proof.

5.4. Proof of the Main Result

With all the preparatory work at hand, we are ready to prove the main result.

Proof of Theorem 1.2. Recall from (2.2) that the eigenvalues of the concentration operator remain the same if we replace E, F and L with $t^{-1}E$, tF and tL respectively. We choose $t = |\partial E|^{1/(d-1)}$ and notice that $t^{-1}E$ satisfies $|\partial t^{-1}E| = 1$, $\eta_{\partial t^{-1}E} = t^{-1}\eta_{\partial E} = 1$, and $\kappa_{\partial t^{-1}E} = \kappa_{\partial E}$. Furthermore, we also have that tF has regular boundary at scale $\eta_{\partial t} = t \eta_{\partial F} = (|\partial E| |\partial F|)^{1/(d-1)} \ge 1$ with constant $\kappa_{\partial tF} = \kappa_{\partial F}$, and $tL \ge 1$ by assumption on L.

Note that for $F' \subseteq (-tL/2, tL/2)^d$, the operator T has integral kernel

$$K(x, y) = \chi_{F'}(x)\chi_{F'}(y)\frac{1}{(tL)^d}\sum_{k \in (t^{-1}E)_{t,t}} e^{-2\pi i k(x-y)}.$$

Thus,

$$\operatorname{tr}(T) = \int K(x, x) dx = \int_{F'} \frac{1}{(tL)^d} \sum_{k \in (t^{-1}E)_{tL}} 1 dx = \frac{\#(t^{-1}E)_{tL}}{(tL)^d} |F'|.$$

On the other hand, from Lemmas 5.1 and 5.2 we have that

$$\operatorname{tr}\left(T^{\pm} - (T^{\pm})^{2}\right) \leq C_{\alpha,d} \frac{|\partial t^{-1}E|}{\kappa_{\partial t^{-1}E}} \frac{|\partial tF|}{\kappa_{\partial t}F} \log \left(\frac{e|\partial t^{-1}E||\partial tF|}{\kappa_{\partial t^{-1}E}}\right)^{2d(1+\alpha)+1}$$
$$= C_{\alpha,d} \frac{|\partial E|}{\kappa_{\partial E}} \frac{|\partial F|}{\kappa_{\partial F}} \log \left(\frac{e|\partial E||\partial F|}{\kappa_{\partial E}}\right)^{2d(1+\alpha)+1} =: C_{E,F}.$$

So from Lemma 5.3,

$$\lambda_n(T^+) \le \frac{1}{2}, \qquad n \ge \left\lceil \frac{\#(t^{-1}E)_{tL}}{(tL)^d} | (tF)^+| \right\rceil + 2C_{E,F};$$

$$\lambda_n(T^-) \ge \frac{1}{2}, \qquad n \le \left\lceil \frac{\#(t^{-1}E)_{tL}}{(tL)^d} | (tF)^-| \right\rceil - 2C_{E,F}.$$

By Corollary 2.2 and $|\partial t^{-1}E| = 1$,

$$\#\{n \in \mathbb{N} : \lambda_n(T^-) < 1/2, \ \lambda_n(T^+) > 1/2\} \lesssim \frac{1}{\kappa_{\partial E}} |(tF)^+ \setminus (tF)^-| + C_{E,F}$$

$$\leq \frac{1}{\kappa_{\partial E}} |\partial tF + B_{\sqrt{d}}(0)| + C_{E,F}$$

$$\lesssim \frac{1}{\kappa_{\partial E}} \frac{t^{d-1} |\partial F|}{\kappa_{\partial F}} + C_{E,F} \lesssim C_{E,F},$$

where in the second to last step we used Lemma 2.1. Since $\lambda_n(T^-) \leq \lambda_n(T) \leq \lambda_n(T)$ $\lambda_n(T^+)$ for every $n \in \mathbb{N}$, again by Lemma 5.2,

$$\#M_{\varepsilon}(T) \leq \#\{n \in \mathbb{N} : 1/2 \leq \lambda_{n}(T^{-}) < 1 - \varepsilon\} + \#\{n \in \mathbb{N} : \varepsilon < \lambda_{n}(T^{+}) \leq 1/2\}$$

$$+ \#\{n \in \mathbb{N} : \lambda_{n}(T^{-}) < 1/2, \lambda_{n}(T^{+}) > 1/2\}$$

$$\lesssim \#M_{\varepsilon}(T^{-}) + \#M_{\varepsilon}(T^{+}) + \#\{n \in \mathbb{N} : \lambda_{n}(T^{-}) < 1/2, \lambda_{n}(T^{+}) > 1/2\}$$

$$\lesssim \frac{|\partial E|}{\kappa_{\partial E}} \frac{|\partial F|}{\kappa_{\partial E}} \log \left(\frac{|\partial E||\partial F|}{\kappa_{\partial E}}\right)^{2d(1+\alpha)+1} .$$

6. The Continuous Fourier Transform

In this section we deduce Theorem 1.1 from Theorem 1.2 by letting $L \to \infty$.

Proof of Theorem 1.1. Fix E and F as in the statement of Theorem 1.1. We consider a sufficiently large resolution such that $L \ge |\partial E|^{-1/(d-1)}$ and $F \subseteq (-L/2, L/2)^d$. Let $S_L: L^2(\mathbb{R}^d) \to L^2(\mathbb{R}^d)$ be the operator given by

$$t S_L : L^2(\mathbb{R}^n) \to L^2(\mathbb{R}^n)$$
 be the operator given by

$$S_L f = T_{E,F,L}(\chi_F f) = \chi_F \mathcal{F}_L^{-1} \chi_{E_L} \mathcal{F}_L \chi_F f, \quad f \in L^2(\mathbb{R}^d).$$

Note that S_L and $T_{E,F,L}$ share the same non-zero eigenvalues, and recall the operator *S* from (1.1).

Step 1. We show that

$$\lim_{L \to \infty} ||S_L - S|| = 0. \tag{6.1}$$

Recall that $Q_{L^{-1}} = L^{-1}[-1/2, 1/2)^d$ and define the auxiliary set

$$\Gamma_L = \bigcup_{m \in E_L} m + Q_{L^{-1}}.$$

Note that the symmetric difference $E\Delta\Gamma_L$ is included in $\partial E + B_{L^{-1}\sqrt{d}}(0)$. From Lemma 2.1,

$$|E\Delta\Gamma_L| \le |\partial E + B_{L^{-1}\sqrt{d}}(0)| \lesssim \frac{|\partial E|}{\kappa L} \left(1 + (L\eta_{\partial E})^{-(d-1)}\right) \xrightarrow{L\to\infty} 0.$$

Using this and setting $R_L = \chi_F \mathcal{F}^{-1} \chi_{\Gamma_L} \mathcal{F} \chi_F$, for $f \in L^2(\mathbb{R}^d)$ we have

$$\begin{aligned} \|(R_L - S)f\|_2^2 &\leq \|(\chi_{\Gamma_L} - \chi_E)\mathcal{F}(\chi_F f)\|_2^2 \leq |E\Delta\Gamma_L| \|\mathcal{F}(\chi_F f)\|_{\infty}^2 \\ &\leq |E\Delta\Gamma_L| \|\chi_F f\|_1^2 \leq |E\Delta\Gamma_L| |F| \|f\|_2^2 \xrightarrow{L \to \infty} 0. \end{aligned}$$

(2024) 248:35

To prove (6.1), it only remains to show that

$$||R_L - S_L|| \xrightarrow{L \to \infty} 0. \tag{6.2}$$

To this end, let $f \in L^2(\mathbb{R}^d)$ and estimate

$$||R_{L}f - S_{L}f||_{2}^{2} = \int_{F} \left| \int_{\Gamma_{L}} \mathcal{F}(\chi_{F}f)(w) e^{2\pi i w x} dw - L^{-d} \sum_{m \in E_{L}} \mathcal{F}(\chi_{F}f)(m) e^{2\pi i m x} \right|^{2} dx$$

$$= \int_{F} \left| \sum_{m \in E_{L}} \int_{m+Q_{L^{-1}}} \mathcal{F}(\chi_{F}f)(w) e^{2\pi i w x} - \mathcal{F}(\chi_{F}f)(m) e^{2\pi i m x} dw \right|^{2} dx$$

$$= \int_{F} \left| \sum_{m \in E_{L}} \int_{m+Q_{L^{-1}}} \int_{F} f(t) \left(e^{2\pi i w (x-t)} - e^{2\pi i m (x-t)} \right) dt dw \right|^{2} dx$$

$$\lesssim \int_{F} \left(\sum_{m \in E_{L}} \int_{m+Q_{L^{-1}}} \int_{F} |f(t)| |w - m| |x - t| dt dw \right)^{2} dx$$

$$\lesssim \int_{F} \left(\sum_{m \in E_{L}} L^{-(d+1)} \int_{F} |f(t)| |x - t| dt \right)^{2} dx$$

$$\lesssim (\#E_{L})^{2} L^{-2(d+1)} \int_{F} ||f||_{2}^{2} \int_{F} |x - t|^{2} dt dx$$

$$\lesssim L^{-2} \frac{\max\{|\partial E|^{2d/(d-1)}, 1\}}{\kappa_{F}^{2}} ||F|^{2} \operatorname{diam}(F)^{2} ||f||_{2}^{2},$$

where in the last estimate we used Corollary 2.2. Hence (6.2) holds.

Step 2. Since S_L and $T_{E,F,L}$ share the same non-zero eigenvalues, the estimates in Theorem 1.2 apply also to S_L for all sufficiently large L. By the Fischer-Courant formula, operator convergence of positive compact operators implies convergence of their eigenvalues. Hence, by (6.1), the estimate satisfied by the spectrum of S_L extends to the spectrum of S.

7. The Discrete Fourier Transform

Proof of Theorem 1.3. Let us define $E := \Omega + \overline{Q_1}$. Then $\Omega = E_L$ for L = 1. Let us apply Theorem 1.2 with L = 1 to E, F. We check the relevant hypotheses.

We first note that ∂E is an almost disjoint union of faces of cubes (by almost disjoint we mean that the intersection of any two faces has zero \mathcal{H}^{d-1} -measure). Moreover, each one is contained in $k + \overline{Q_1}$ for exactly one $k \in \Omega$ that must belong to $\partial \Omega$. In particular,

$$\partial E \subset \bigcup_{k \in \partial \Omega} k + \partial Q_1.$$

Conversely, for each point $k \in \partial \Omega$ at least one face of the cube $k + \overline{Q_1}$ lies in ∂E . Thus,

$$\#\partial\Omega \le |\partial E| \le 2d \cdot \#\partial\Omega,\tag{7.1}$$

and consequently,

$$|\partial E| |\partial F| \ge \# \partial \Omega \cdot |\partial F| \ge 1.$$

Moreover, (7.1) shows that the choice L=1 satisfies $L \ge |\partial E|^{-1/(d-1)}$ as $\partial \Omega$ contains at least one point.

Now fix $0 < r \le \sqrt{d} \cdot (\#\partial\Omega)^{1/(d-1)}$ and let us show that ∂E is regular at maximal scale. If $r \le 2\sqrt{d}$, and $x \in \partial E$ we clearly have $\mathcal{H}^{d-1}\big(\partial E \cap B_r(x)\big) \gtrsim r^{d-1}$ as E is a union of cubes of length 1. If $r > 2\sqrt{d}$, set $n = \lfloor r/\sqrt{d} \rfloor$ and let $x \in \partial E$. There exists $k_x \in \partial\Omega$ such that $|k_x - x| \le \sqrt{d}/2$. Note that for $y \in k_x + \overline{Q_{n+1}}$,

$$|y-x| \leq \frac{\sqrt{d}(n+1)}{2} + \frac{\sqrt{d}}{2} \leq \frac{r}{2} + \sqrt{d} < r.$$

Hence, $k_x + \overline{Q_{n+1}} \subseteq B_r(x)$. This together with the fact that for $k \in \partial \Omega$ at least one face of the cube $k + \overline{Q_1}$ lies in ∂E gives

$$\mathcal{H}^{d-1}(\partial E \cap B_r(x)) \ge \mathcal{H}^{d-1}(\partial E \cap k_x + \overline{Q_{n+1}}) \ge \mathcal{H}^{d-1}(\partial E \cap \bigcup_{k \in \partial \Omega \cap k_x + Q_n} k + \partial Q_1)$$

$$\ge \#(\partial \Omega \cap k_x + Q_n) \ge \kappa_{\partial \Omega} n^{d-1} \gtrsim \kappa_{\partial \Omega} r^{d-1}.$$

This shows that ∂E is regular at scale $\sqrt{d} \cdot (\#\partial\Omega)^{1/(d-1)}$ with constant $C_d \cdot \kappa_{\partial\Omega}$. Note that if a set X is regular at scale η_X and constant κ_X , then it is also regular at scale $\alpha\eta_X$ and constant $\min\{1, \alpha^{1-d}\}\kappa_X$, for every $\alpha > 0$. By (7.1) we therefore see that ∂E is regular at scale $\eta_{\partial E} = |\partial E|^{1/(d-1)}$ and constant $\kappa_{\partial E} \asymp \kappa_{\partial\Omega}$.

The desired estimates now follow by applying Theorem 1.2 to E and F with L=1, together with (7.1).

8. Proof of Remark 1.4

First we combine Lemma 5.3, Lemma 5.1 (for p=1) and Theorem 1.1 to conclude that there exist a constant $C=C_{\alpha,d}>0$ such that if

$$n \geq \lceil |E| \cdot |F| \rceil + C \frac{|\partial E|}{\kappa_{\partial E}} \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log \left(\frac{e|\partial E||\partial F|}{\kappa_{\partial E}} \right)^{2d(1+\alpha)+1} =: C_1,$$

then $\lambda_n \leq 1/2$, and if

$$n \leq \lceil |E| \cdot |F| \rceil - C \frac{|\partial E|}{\kappa_{\partial E}} \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log \left(\frac{e|\partial E||\partial F|}{\kappa_{\partial E}} \right)^{2d(1+\alpha)+1} =: C_2,$$

then $\lambda_n \geq 1/2$.

For $\varepsilon \in (0, 1)$, define $\varepsilon_0 := \min\{\varepsilon, 1 - \varepsilon\} \le 1/2$ and let $0 < \tau < \varepsilon_0$. Observe that

$$\{1, ..., \lfloor C_2 \rfloor\} \setminus M_{\tau}(S) \subseteq N_{1-\varepsilon_0}(S) \subseteq N_{\varepsilon}(S) \subseteq N_{\varepsilon_0}(S) \subseteq \{1, ..., \lceil C_1 \rceil\} \cup M_{\tau}(S),$$

where we understand $\{1, ..., \lfloor C_2 \rfloor\}$ to be \emptyset if $C_2 < 1$. Consequently,

$$C_2 - 1 - \# M_{\tau}(S) \le \# N_{\varepsilon}(S) \le C_1 + 1 + \# M_{\tau}(S).$$

Rearranging the last expression and using Theorem 1.1 for τ gives

$$\left| N_{\varepsilon}(S) - |E| \cdot |F| \right| \lesssim \frac{|\partial E|}{\kappa_{\partial E}} \cdot \frac{|\partial F|}{\kappa_{\partial F}} \cdot \log \left(\frac{|\partial E| |\partial F|}{\kappa_{\partial E} \tau} \right)^{2d(1+\alpha)+1}.$$

Letting $\tau \nearrow \varepsilon_0$ yields (1.13).

Acknowledgements. This research was funded in whole or in part by the Austrian Science Fund (FWF): 10.55776/Y1199. F. M. was also supported by the EPSRC: NIA EP/V002449/1.

Funding Open access funding provided by University of Vienna.

Data availability This research contains no data sets.

Declarations

Conflict of interest The authors are not aware of any competing financial or nonfinancial interests that are directly or indirectly related to this work.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http://creativecommons.org/ licenses/by/4.0/.

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

References

- 1. ABREU, L.D., PEREIRA, J.M., ROMERO, J.L.: Sharp rates of convergence for accumulated spectrograms. Inverse Probl. 33(11), 115008, 2017
- 2. ADCOCK, B., HUYBRECHS, D.: Frames and numerical approximation. SIAM Rev. 61(3), 443-473, 2019
- 3. ANDÉN, J., ROMERO, J.L.: Multitaper estimation on arbitrary domains. SIAM J. Imaging Sci. 13(3), 1565–1594, 2020
- 4. Andén, J., Singer, A.: Factor analysis for spectral estimation. In: *Proc. SampTA*, pp. 169-173. IEEE, 2017
- 5. BHAMRE, T., ZHANG, T., SINGER, A.: Denoising and covariance estimation of single particle cryo-EM images. J. Struct. Biol. 195(1), 72-81, 2016
- 6. BONAMI, A., JAMING, P., KAROUI, A.: Non-asymptotic behavior of the spectrum of the sinc-kernel operator and related applications. J. Math. Phys. 62(3), Paper No. 033511, 2021
- 7. CARABALLO, D.G.: Areas of level sets of distance functions induced by asymmetric norms. Pac. J. Math. 218(1), 37-52, 2005
- 8. DIESTEL, J., JARCHOW, H., TONGE, A.: Absolutely Summing Operators. Cambridge Studies in Advanced Mathematics, vol. 43. Cambridge University Press, Cambridge (1995)
- 9. DZIUBAŃSKI, J., HERNÁNDEZ, E.: Band-limited wavelets with subexponential decay. Canad. Math. Bull. 41(4), 398-403, 1998

- 10. EVANS, L.C., GARIEPY, R.F.: Measure Theory and Fine Properties of Functions. Studies in Advanced Mathematics. CRC Press, Boca Raton, FL (1992)
- 11. FEFFERMAN, C.L.: The uncertainty principle. Bull. Am. Math. Soc. (N.S.) 9(2), 129–206, 1983
- 12. GRÜNBAUM, F.A.: Serendipity strikes again. Proc. Natl. Acad. Sci. USA 119(26), e2207652119, 2022
- 13. Grünbaum, F.A., Longhi, L., Perlstadt, M.: Differential operators commuting with finite convolution integral operators: some nonabelian examples. SIAM J. Appl. Math. **42**(5), 941–955, 1982
- 14. HAN, S.-C., SIMONS, F.J.: Spatiospectral localization of global geopotential fields from the gravity recovery and climate experiment (grace) reveals the coseismic gravity change owing to the 2004 sumatra-andaman earthquake. J. Geophys. Res. 113, B01405, 2008
- 15. HERNÁNDEZ, E., WEISS, G.: A first course on wavelets Studies in Advanced Mathematics. CRC Press, Boca Raton, FL, 1996. With a foreword by Yves Meyer.
- 16. HOGAN, J. A., LAKEY, J. D.: Duration and bandwidth limiting. Applied and Numerical Harmonic Analysis. Birkhäuser/Springer, New York, 2012. Prolate functions, sampling, and applications.
- 17. ISRAEL, A.: The eigenvalue distribution of time-frequency localization operators. arXiv:1502.04404v1, 2015.
- 18. ISRAEL, A., MAYELI, A.: On the eigenvalue distribution of spatio-spectral limiting operators in higher dimensions. Appl. Comput. Harmon. Anal. 70, 101620, 2024
- 19. KARNIK, S., ROMBERG, J., DAVENPORT, M.A.: Improved bounds for the eigenvalues of prolate spheroidal wave functions and discrete prolate spheroidal sequences. Appl. Comput. Harmon. Anal. 55, 97-128, 2021
- 20. KARNIK, S., ZHU, Z., WAKIN, M.B., ROMBERG, J., DAVENPORT, M.A.: The fast Slepian transform. Appl. Comput. Harmon. Anal. 46(3), 624–652, 2019
- 21. LANDA, B., SHKOLNISKY, Y.: Steerable principal components for space-frequency localized images. SIAM J. Imaging Sci. 10(2), 508–534, 2017
- 22. LANDAU, H.J.: Necessary density conditions for sampling and interpolation of certain entire functions. Acta Math. 117, 37-52, 1967
- 23. Landau, H.J.: On Szego's eigenvalue distribution theorem and non-Hermitian kernels. J. Anal. Math. 28, 335-357, 1975
- 24. LANDAU, H.J., POLLAK, H.O.: Prolate spheroidal wave functions, Fourier analysis and uncertainty. II. Bell Syst. Tech. J. 40, 65–84, 1961
- 25. LANDAU, H.J., POLLAK, H.O.: Prolate spheroidal wave functions, Fourier analysis and uncertainty. III. The dimension of the space of essentially time- and band-limited signals. Bell Syst. Tech. J. 41, 1295–1336, 1962
- 26. LANDAU, H.J., WIDOM, H.: Eigenvalue distribution of time and frequency limiting. J. Math. Anal. Appl. 77(2), 469-481, 1980
- 27. MARCECA, F., ROMERO, J.L.: Spectral deviation of concentration operators for the shorttime Fourier transform. Stud. Math. 270(2), 145–173, 2023
- 28. Marín, J. J., Martell, J. M., Mitrea, D., Mitrea, I., Mitrea, M.: Singular integral operators, quantitative flatness, and boundary problems, volume 344 of Progress in Mathematics. Birkhäuser/Springer, Cham, [2022] 2022.
- 29. MATTHYSEN, R., HUYBRECHS, D.: Function approximation on arbitrary domains using Fourier extension frames. SIAM J. Numer. Anal. 56(3), 1360–1385, 2018
- 30. NIKOLSKI, N. K.: Operators, functions, and systems: an easy reading. Vol. 1, volume 92 of Mathematical Surveys and Monographs. American Mathematical Society, Providence, RI, 2002. Hardy, Hankel, and Toeplitz, Translated from the French by Andreas Hartmann.
- 31. Osipov, A.: Certain upper bounds on the eigenvalues associated with prolate spheroidal wave functions. Appl. Comput. Harmon. Anal. 35(2), 309-340, 2013
- 32. SIMONS, F.J., VAN DER HILST, R.D., ZUBER, M.T.: Spatiospectral localization of isostatic coherence anisotropy in Australia and its relation to seismic anisotropy: Implications for lithospheric deformation. J. Geophys. Res. 108(B5), 2250, 2003

- 33. SIMONS, F.J., WANG, D.V.: Spatiospectral concentration in the Cartesian plane. *GEM Int. J. Geomath.* **2**(1), 1–36, 2011
- 34. SIMONS, F.J., ZUBER, M.T., KORENAGA, J.: Isostatic response of the Australian lithosphere: estimation of effective elastic thickness and anisotropy using multitaper spectral analysis. *J. Geophys. Res. B* **105**(B8), 19163–19184, 2000
- 35. SLEPIAN, D.: Prolate spheroidal wave functions, Fourier analysis and uncertainty. IV. Extensions to many dimensions generalized prolate spheroidal functions. *Bell Syst. Tech. J.* **43**, 3009–3057, 1964
- 36. SLEPIAN, D.: Prolate spheroidal wave functions, Fourier analysis, and uncertainty-V: the discrete case. *Bell Syst. Tech. J.* **57**(5), 1371–1430, 1978
- 37. SLEPIAN, D., POLLAK, H.O.: Prolate spheroidal wave functions, Fourier analysis and uncertainty. *I. Bell Syst. Tech. J.* **40**, 43–63, 1961
- 38. WIDOM, H.: Asymptotic behavior of the eigenvalues of certain integral equations. II. *Arch. Ration. Mech. Anal.* 17, 215–229, 1964
- ZHU, Z., KARNIK, S., WAKIN, M.B., DAVENPORT, M.A., ROMBERG, J.: ROAST: rapid orthogonal approximate Slepian transform. *IEEE Trans. Signal Process.* 66(22), 5887– 5901, 2018

F. MARCECA
Department of Mathematics,
King's College London,
London
UK.

e-mail: felipe.marceca@kcl.ac.uk

and

J. L. ROMERO
Faculty of Mathematics,
University of Vienna,
Oskar-Morgenstern-Platz 1,
1090 Vienna
Austria.
e-mail: jose.luis.romero@univie.ac.at

and

Acoustics Research Institute, Austrian Academy of Sciences,
Dr. Ignaz Seipel-Platz 2,
1010 Vienna
Austria.

and

M. SPECKBACHER
Faculty of Mathematics,
University of Vienna,
Oskar-Morgenstern-Platz 1,
1090 Vienna
Austria.

e-mail: michael.speckbacher@univie.ac.at

(Received May 8, 2023 / Accepted March 14, 2024) © The Author(s) (2024)