



Hot workability behaviour of two P92 creep resistant steels: Constitutive analysis

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Abstract

This article reports the flow stress behaviour of two P92 steels at a temperature range of 850–1000°C and a strain rate of 0.1–10 s⁻¹ using the Gleeble® 3500 thermomechanical simulator. A physically-based constitutive model was used to analyse the effects of deformation conditions on the flow stress behaviour during deformation. This model incorporates the influence in the variation of Young's modulus and the self-diffusion coefficient as affected by temperature. The study developed constitutive equations that predict the flow stress behaviour of the two steels investigated. From the constitutive analysis of the results, the stress exponent n was: 9.8 (steel A) and 10.3 (steel B). The model used the self-diffusion activation energy of steel. The statistical parameters: correlation coefficient of 0.99 (for steel A and B), the absolute average relative error of 2.18% (steel A) and 2.20% (steel B) quantified the applicability of the model. The quantification results show that the constitutive equations developed have high accuracy in predicting the workability of the two P92 steels. The study has shown that this method is applicable in predicting the metal flow pattern of two P92 steels in the metalworking processes.

Keywords P92 steel · Hot deformation · Constitutive analysis · Stress exponent · Flow stress

1 Introduction

The creep-resistant steels of the family 9–12 wt% Cr steels are now used widely in modern power plant components [1]. Thermal efficiency and reduced CO₂ emissions from power plants require more research on these steels [2]. The most commonly used 9–12 wt% Cr steels for power plant components are: P91, P122, E911, and P92 [1, 3]. These steels have

become popular due to their excellent creep strength, weldability, creep resistance and fabricability [4, 5]. For example, P92 steel, which has ~30% higher creep strength than P91 steel [6], has replaced P91 steel and other older steel for power plant components such as turbine and boiler tubes and pipes. This steel can achieve a high steam temperature of up to 650°C than the 9–12 wt% Cr steels [7].

The processing route for most components after casting involves rolling, forging, drawing, and extrusion [8]. The forming process can eliminate defects like voids formed during casting and cause grain refinement of the microstructure. The metal flow behaviour during forming is of great concern to the engineers and the designers [9]. To improve the mechanical properties of these steels, forming parameters must be controlled [10]. Constitutive equations provide information on metal flow patterns [11]. These equations can accurately describe the workability of any material under different forming conditions [12]. These equations act as input in the Finite Element Method (FEM) simulation codes for studying forming process [13]. However, FEM accuracy depends on the constitutive equations developed for the material [14]. Computer simulation reduces production costs and time, especially in metal processing [15]. Moreover, FEM modules provide a

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cost-effective way of designing and optimising metal forming techniques commonly used in the production industry [16].

Constitutive modelling of flow stress during hot forming is of great importance. Most works on the hot workability of metal have used the conventional Arrhenius equations to determine the material constants [17–20]. The hyperbolic sine-law Arrhenius model, as given in Eq. 1.1 [21], has been accepted widely as the constitutive model for analysing the metal flow pattern of most materials, such as P91 steel [22] and P92 steel [17, 23, 24], and Titanium alloys [25]. This equation relates the deformation conditions and the flow stress as:

$$\dot{\epsilon} = A(\sinh(\alpha\sigma))^n \exp\left[\frac{-Q}{RT}\right] \quad (1.1)$$

where A , α and n are material constants, Q is the thermal activation energy, $\dot{\epsilon}$ is the strain rate, σ is the flow stress, T is the deformation temperature, and R is the universal gas constant (8.314 kJ. mol⁻¹). The constitutive constants obtained in Eq. 1.1 do not account for any internal microstructure evolution during deformation, thus referred to as apparent values [26]. The equation assumes that the microstructure remains constant during forming [27]. The obtained n and Q values are higher than 270 kJ.mol⁻¹ for iron [28]. The deviation of activation energy is due to a variation of Young's Modulus $E(T)$ and self-diffusion coefficient, which varies with temperature [26]. Therefore, to account for these factors, Eq. 1.1 for the general physically-based equation ($n=5$) is given as follows [27, 29, 30]:

$$\frac{\dot{\epsilon}}{D(T)} = B[\sinh(\alpha\sigma/E(T))]^5 \quad (1.2)$$

where $D(T)$ is determined using $D(T) = D_0 \exp(Q_{sd}/RT)$, D_0 is the pre-exponential constant. Q_{sd} and $E(T)$ are self-diffusion activation energy, and temperature affected Young's Modulus of the material, respectively. These values (D_0 , Q_{sd} , $E(T)$) are obtained from the Ashby table [31].

However, the stress exponent $n=5$ (Eq. 1.2) is an absolute value [32]. The equation assumes that no microstructure evolution such as DRX, DRV and dynamic precipitates occurs during forming, which affects the stress exponent. The n -value is, therefore, not a constant but a variable parameter [27, 33, 34], and the equation is as follows:

$$\frac{\dot{\epsilon}}{D(T)} = B[\sinh(\alpha\sigma/E(T))]^n \quad (1.3)$$

A few studies have reported results on using the physically-based model to analyse the flow stress behaviour of metals and alloys [27, 32, 33, 35, 36]. To this end, no study has reported on the applicability of the new physically-based constitutive equation in describing the deformation behaviour of creep-resistant steels. The novelty in this article refers to the effectiveness of the physically-based model in studying the workability of two P92 steels having variation in chromium and tungsten content, which is conspicuously missing in the literature. The present study aims to develop a suitable physically-based constitutive equation for evaluating and analysing the flow behaviour of P92 steel. The output of this study will provide more insight into this method for future applications in research and industrial application. Hot compression tests were conducted at different deformation conditions using Gleeble® 3500 equipment. The flow stress–strain curves were analysed to develop mathematical rate equations for predicting the flow stress behaviour of P92 steel under the investigated conditions.

2 Experimental procedure

Table 1 shows the chemical composition (in wt%) of the two P92 steel with variations in chromium and tungsten content studied. Test specimens measuring 8 mm diameter and 12 mm height underwent a uniaxial compression test using the Gleeble-3500 simulator under vacuum. Experimental test conditions were: a temperature range of 850–1000°C at an interval of 50°C and strain rates of 0.1, 1 and 10 s⁻¹. These deformation temperatures were in the austenite region. Hence, the deformation tests in this study were in a single phase. These forming conditions have been used previously for this steel [37]. Thermocouples at the mid-height of the specimen assisted in monitoring temperature during testing. Nickel paste and graphite foil are applied between the sample and the anvil to reduce friction. Before testing, the samples were heated at a rate of 5°C/s to 1100°C and held for 180 s. Then, cooled at 10°C/s and soaked for 60 s at the deformation temperature before compression to a strain of 0.6. After deformation, samples were air-cooled rapidly to room temperature. At the austenitisation temperature (1100°C), the complete dissolution of most carbides, especially M₂₃C₆, occurs. Dissolution temperature of chromium carbide starts at 900°C and completely dissolves at 1100°C [1]. At 1100°C, the P92 steel transforms into the austenitic phase. At this temperature, deformation resistance is low due to the absence of carbides [2]. Carbide dissolution occurs

Table 1 Chemical compositions (wt.%) of the two steels

Steel	C	Mn	Si	Cr	Mo	Ni	Cu	Al	V	Nb	W	Co	Fe
Steel-A	0.10	0.39	0.2	8.29	0.65	0.19	0.08	0.012	0.16	0.09	2.07	0.015	Bal
Steel-B	0.11	0.32	0.25	9.48	0.61	0.17	0.00	0.023	0.20	0.08	2.34	0.024	Bal

Table 2 Dissolution temperature of phase in the two steels investigated

Dissolution temperature (°C)			
Steel	M ₂₃ C ₆	MX	Laves phase
A	888	1250	720
B	911	1192	770

below 920°C according to ThermoCalc calculations, as shown in Fig. 2 (also, Table 2) for the two steels. The information above justifies the selection of 1100°C as the austenitisation temperature during forming in this study.

3 Results and discussion

3.1 Microstructure of as-received steels

Before being put into service, P92 steel undergoes normalisation and tempering conditions. The chemical composition of the two P92 steel (named steel A and B) are listed in Table 1 and have the following relative ‘amounts’ of Cr, W and Mo:

The optical micrographs (Fig. 1) show the tempered martensite microstructure of the as-received steels. The

Steel:	Steel A	Steel B
Cr	medium	high
W	medium	high
Mo	high	high
Cr _{eq}	13.10	13.92
Ni _{eq}	4.58	4.40

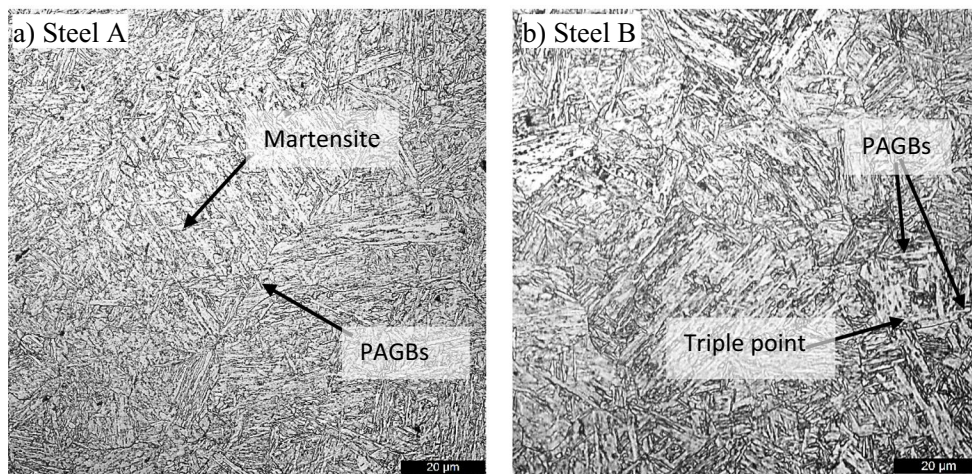
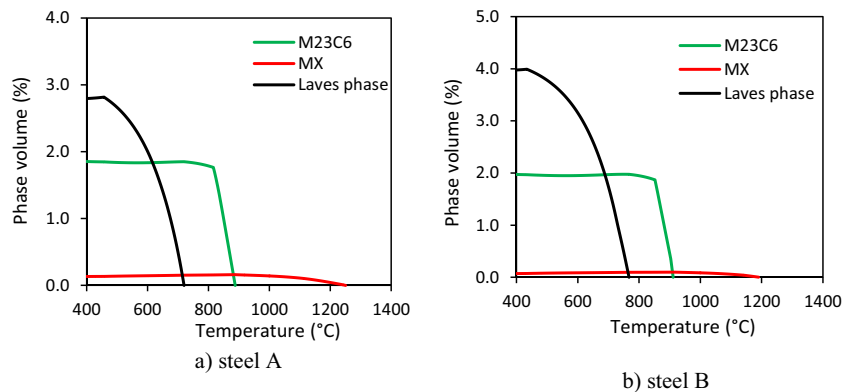


Fig. 1 Optical microstructures of as-received samples

Fig. 2 Thermo-Calc predicted volume % of M₂₃C₆, MX and Laves phases vs. temperature under equilibrium condition for the two steels



two steels had a tempered microstructure with well-defined prior austenite grain boundaries and triple points.

After heat treatment, precipitation occurs along the grain and lath boundaries and in the matrix. $M_{23}C_6$ ($M = Fe, Mo, W, Cr$) carbides precipitate along the prior austenite grain boundaries PAGBs [38], while MX ($M = V, Nb; X = C, N$) precipitates are usually randomly distributed in the matrix [39]. The carbides along the prior austenite grain boundary hinder the movement of the PAG boundaries and the sub-grain dislocations, thus improving the creep resistance [40]. $M_{23}C_6$ carbides improve the creep strength by pinning the grain boundary movement [41], while the MX precipitates impede dislocation movement during deformation [4].

Thermo-Calc software with the TCFE5 database predicted the equilibrium transformation and precipitate dissolution temperatures for the two steels investigated. Figure 2 shows the change in the phase volume (%) as a function of temperature predicted using ThermoCalc under equilibrium conditions for the commonly observed phases: $M_{23}C_6$ carbides, MX, and Laves phase. Dissolution of precipitates increases with an increase in temperature. The precipitate volume fraction increased as temperature decreased below the equilibrium dissolution temperature (Fig. 2).

The ThermoCalc results show that MX precipitates have higher dissolution temperatures, as shown in Fig. 2. At higher austenitising temperatures, the precipitates will dissolve into the solid solution and precipitate during tempering. The ThermoCalc results (Fig. 2) provided the basis for choosing 1100°C for austenitisation temperature during deformation.

3.2 Flow stress behaviour

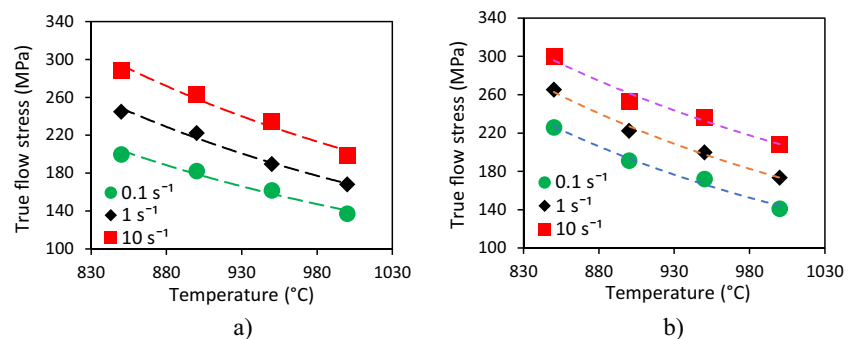
In this study, the flow stress–strain curves were friction corrected. During the uniaxial compression test, deformation at a lower strain rate caused “sticking” friction. Hence, the study accounted for interfacial friction correction on flow stress. However, the study did not consider adiabatic heating during deformation. The temperature variation between the pre-set and measured temperature in all deformation conditions was approximately below 20°C. The temperature

variation did not cause substantial variation in the flow stress values. Therefore, all the flow stress values used in this study were friction corrected. The relationship between the flow stress and temperature at a constant strain rate for the two steels is as given in Fig. 3. The plots show that the flow stress decreased as the deformation temperature increased at any strain rate. The decrease in flow stress is due to dynamic softening (dynamic recovery) [42], which is the dominant dynamic mechanism for P92 steel [18]. At higher deformation temperatures, the kinetic of atoms and dislocation movement increases, causing low flow stress [43]. The flow stress values are higher at a higher strain rate since the deformation time is shorter for dislocation and rearrangement to balance, resulting in work hardening and DRV as the dominant mechanisms. Table 3 shows the flow stress values obtained for the two steel at different deformation conditions. The flow stress results show that steel B had the highest flow stress values at 850°C and in all strain rates investigated. High flow stress values at 850°C were due to incomplete dissolution of carbides before deformation [18]. Hence, pinning dislocation results in an increase in flow stress values [44]. At higher temperatures (900–1000°C), the flow stress values for the two steels were relatively close. This deformation region (900–1000°C) lies in the austenitic zone for the two P92 steels. In this region, carbides dissolution of carbides occurs, causing low flow stress values. From the constitutive analysis, the stress exponent of the two steels was relatively the same, further suggesting less resistance to deformation under these conditions. The slight differences in flow stress values can be due to differences in Cr content.

Table 3 Flow stress values obtained during forming

Steel	Strain rate	850	900	950	1000
Steel-A	0.1	200	183	162	138
	1	245	222	189	168
	10	288	263	234	199
Steel-B	0.1	226	191	172	141
	1	265	222	200	172
	10	300	253	237	208

Fig. 3 Relationship between flow stress against temperature a) Steel-A, and b) Steel-B



Higher Cr content in steel B causes more carbides to form, thus hindering dislocation and causing high flow stress [44].

3.3 The physically-based constitutive equation

The Young’s modulus and the self-diffusion affected by temperature incorporated in the physically-based model can describe the metal flow pattern [28]. Therefore, Eq. 1.1 then becomes [27]:

$$\frac{\dot{\epsilon}}{D(T)} = B[\sinh(\alpha\sigma/E(T))]^n \tag{3.1}$$

$$D(T) = D_0 \exp(-Q_{sd}/RT) \tag{3.2}$$

$$E(T) = E_0 \left[1 - \frac{T_m}{G_0} \frac{dG(T-300)}{dT} \frac{1}{T_m} \right] \tag{3.3}$$

where T_m is steel absolute melting temperature, B , n , α are material constants and $\dot{\epsilon}$ is the strain rate (s^{-1}), E_0 and G_0 are the Young’s modulus and the shear modulus at temperature

of 300K respectively, the term $T_m/G_0 \cdot dG/dT$ is denoted as η representing modulus as affected by temperature and D is the self-diffusion coefficient. The constants: D_0 , E_0 and Q_{sd} values are obtained from Ashby [31] tables, provided in Table 4. Then, $\alpha \approx \beta/n'$.

The unknown material constants in Eq. 1.2 are obtained by plotting graphs using Eqs. 3.5 and 3.6.

$$\frac{\dot{\epsilon}}{D(T)} = B_1 \left(\frac{\sigma}{E(T)} \right)^{n'} \tag{3.5}$$

$$\frac{\dot{\epsilon}}{D(T)} = B_2 \exp\left(\frac{\beta'\sigma}{E(T)}\right) \tag{3.6}$$

n' and β' values were obtained from plots in Fig. 4 (Steel-A) and Fig. 5 (Steel-B) using linear regression analysis. Then, $\alpha \approx \beta'/n'$. From Eq. 1.2, the slope and the intercept of the plots in Fig. 4c (Steel-A) and Fig. 5c (Steel-B) determined the stress exponent n (slope of the graph) and the $\ln B$ (slope intercept). Table 5 gives the calculated material constants for the two steels.

By substituting material constants into Eq. 1.2, the resultant constitutive equations for the two steels is obtained as:

$$\text{Steel - A : } \dot{\epsilon} \exp\left(\frac{270000}{RT}\right) = Z = 1.33 \times 10^{11} (\sinh(576.68\sigma_{ss}/E(T)))^{9.8} \tag{3.7}$$

Table 4 Data values obtained for P92 steel in Ashby tables [31]

D_0 (m^2/s)	Q_{sd} (kJ/mol)	η	E_0 (MPa)	G_0 (MPa)
1.8×10^{-5}	270	-0.91	2.16×10^5	6.4×10^4

Fig. 4 Plots for determining a) $\ln B$, b) n' and c) n for steel A

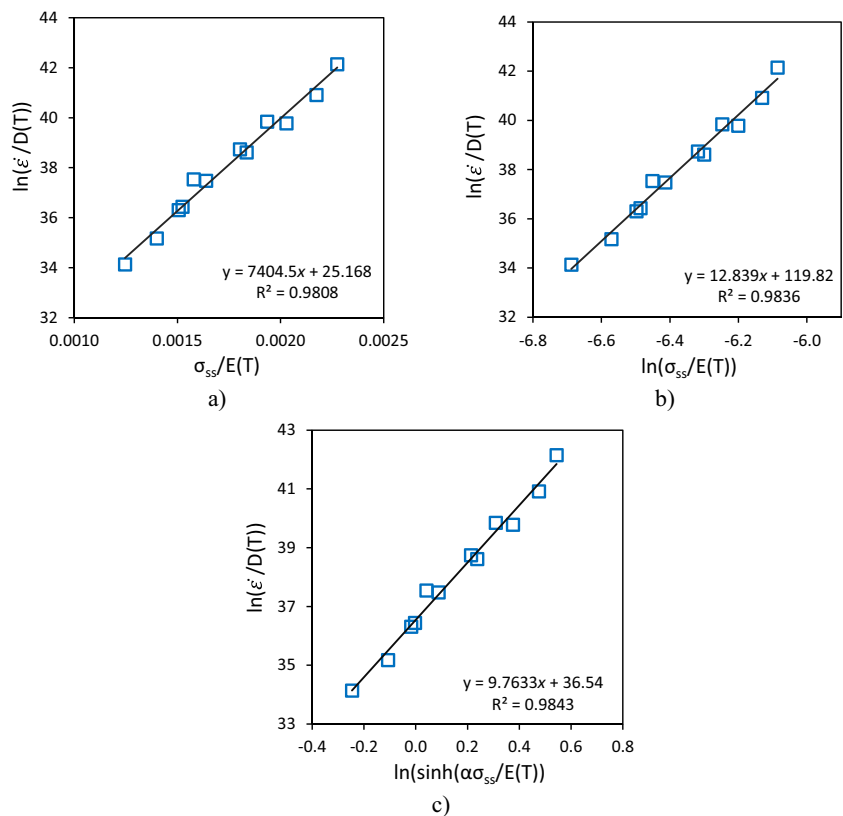
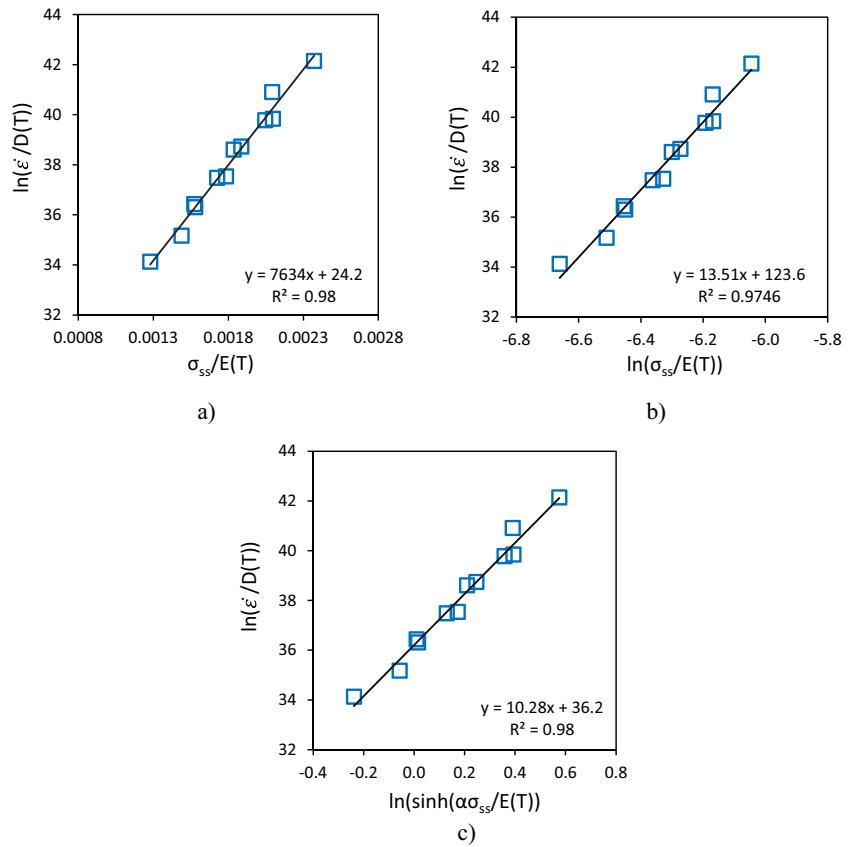


Fig. 5 Plots for determining **a)** $\ln B$, **b)** n' and **c)** n for steel-B:



$$\text{Steel - B} : \dot{\epsilon} \exp\left(\frac{270000}{RT}\right) = Z = 9.53 \times 10^{10} (\sinh(669.32\sigma_{ss}/E(T)))^{10.3} \quad (3.8)$$

Then, Eq. 1.2 becomes:

$$\left(\frac{Z}{B}\right)^{\frac{1}{n}} = \sinh\left(\alpha \frac{\sigma_{ss}}{E(T)}\right) = \frac{e^{\alpha\sigma_{ss}/E(T)}}{2} - \frac{-e^{-\alpha\sigma_{ss}/E(T)}}{2} \quad (3.9)$$

Solving Eq. 11 gives Eq. 12:

$$\sigma_{ss} = \frac{E(T)}{\alpha} \ln \left[\left(\frac{Z}{B}\right)^{\frac{1}{n}} + \left(\left(\frac{Z}{B}\right)^{\frac{2}{n}} + 1 \right) \right]^{\frac{1}{2}} \quad (3.10)$$

Equation 3.10 can predict the flow stress under different forming conditions using $Z = \dot{\epsilon} \exp\left(\frac{32475.34}{T}\right)$. Therefore, α , B and n values from Table 5, the following equations can describe the flow stress:

$$\text{For Steel - A} : \sigma_{ss} = \frac{E(T)}{576.7} \ln \left[\left(\frac{Z}{1.33 \times 10^{11}}\right)^{\frac{1}{9.8}} + \left(\left(\frac{Z}{1.33 \times 10^{11}}\right)^{\frac{2}{9.8}} + 1 \right) \right]^{\frac{1}{2}} \quad (3.11)$$

$$\text{For Steel - B} : \sigma_{ss} = \frac{E(T)}{565.0} \ln \left[\left(\frac{Z}{9.53 \times 10^{10}}\right)^{\frac{1}{10.3}} + \left(\left(\frac{Z}{9.53 \times 10^{10}}\right)^{\frac{2}{10.3}} + 1 \right) \right]^{\frac{1}{2}} \quad (3.12)$$

From the above analysis, the calculated material constants of the two steels did not show any differences. Even though the two steels have a slight variation in the chemical composition (Cr Content), they have relatively the same stress exponent n values (difference of 4.7%). A study reported that n -values are affected by flow stress, which depends on the interaction between precipitates and mobile dislocations [45]. The carbides pin dislocation hence, hindering the deformation process. The stress exponent is temperature-sensitive. The stress exponent n -value increases with a decrease in the deformation temperature and vice-versa [46]. However, deformation temperatures were slightly higher than the Ac_3 of the two steels. Therefore, the flow resistance during deformation should be relatively the same as most carbides might have dissolved. The results further show that the variation in the chemical composition had an insignificant effect on the stress exponent.

Table 5 Calculated constitutive constants of the two P92 steel

Material constant	Steel-A	Steel-B
α	576.68	565.00
B	1.33×10^{11}	9.53×10^{10}
n'	12.84	11.16
n	9.80	10.30

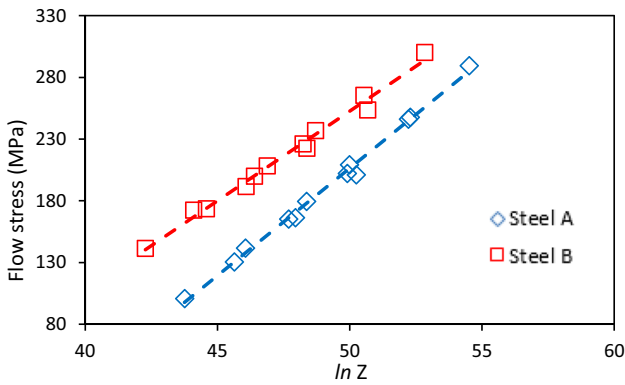


Fig. 6 Comparison of the flow stress and lnZ of the two steels

On the other hand, the B-value of steel A had one order of magnitude over steel B. Parameter *B* is proportional to the activation energy *Q*. Higher *Q*-value results in a higher *B*-value [47]. However, under the analysis method used, the *Q*-value was constant. Then, what might be the reason for this deference in *B*-value? *B*-value has no simple physical meaning [48]. Therefore, *B*-value is a function of forming conditions and microstructure orientation of the material [47]. The materials values given in Table 5 provide the necessary parameters to apply in the physically-based constitutive equation to determine precisely the flow stress behaviour during the deformation of this steel.

3.4 Comparison of Zener parameter and flow stress

The Zener parameter describes the combined effects of deformation temperature and strain rate on the flow stresses. Figure 6 shows the relationship between the flow stress and the lnZ. The plots showed a linear relationship between the flow stress and lnZ.

Figure 6 shows that at any given value of lnZ, the two steels experience different flow stress. This result shows the influence of the deformation conditions (temperature and strain rate) during deformation. Generally, the value of lnZ increases as the deformation temperature decreases with an increase in the strain rate. When comparing steel A and steel B, the latter had higher resistance to deformation. Steel B had the highest lnZ value, especially at higher strain rates, as shown in Fig. 6. Steel B also had higher flow stress than steel A at any given lnZ value (Eq. 3.12). These results indicate that steel B had higher resistance to deformation, which can be due to higher Cr content which contributes to a higher precipitation strengthening, hence hindering deformation. The Z-value also depends on the deformation temperature. Lower Z-values occur at higher forming temperatures. Higher Z-values show that work hardening is the dominant deformation mechanism, and the material undergoes severe plastic deformation. Lower Z-values indicate that DRX and DRV occurred due to an increased dislocation movement and reduced dislocation density. Work hardening and dynamic softening may simultaneously occur during forming, thus affecting the resultant dislocation density and influencing the flow stress behaviour [49]. Therefore, dynamic softening will occur depending on the initial thermal history of the sample subjected before deformation. For example, annealing before deformation can induce low dislocation density, which causes dislocation accumulation at the initial stages of deformation. As strain increases, a high generation of dislocation occurs, resulting in dislocation annihilation, hence initiating dynamic softening.

Equation 3.12 gives the constitutive equations that can be used to predict the flow stress using the Zener parameter for the two steels investigated.

$$\text{Steel A : } \sigma_{sat} = 17.41\ln Z - 664.25$$

$$\text{Steel B : } \sigma_{sat} = 14.55\ln Z - 475.07 \tag{3.12}$$

Fig. 7 Comparison of the calculated and the experimental flow stress a) Steel-A, and b) Steel-B

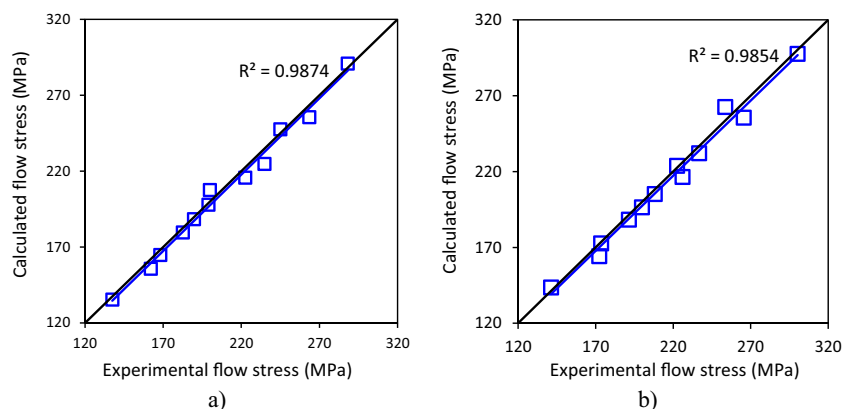


Table 6 Statistical parameters values

Parameter	Steel-A	Steel-B
R	0.99	0.99
$AARE$	2.18	2.20
R^2	0.98	0.99

4 Model verification

The validity of the developed constitutive models (Eqs. 3.11 and 3.12) to accurately predict the flow stress behaviour of the two steels investigated was done using statistical parameters: correlation coefficient R (Eq. 4.1) and the average absolute relative error $AARE$ (Eq. 4.2). These models are widely used to predict the linear relationship between two variables [14, 50, 51].

$$R = \frac{\sum_{i=1}^N (E_i - \bar{E})(C_i - \bar{C})}{\sqrt{\sum_{i=1}^N (E_i - \bar{E})^2 \sum_{i=1}^N (C_i - \bar{C})^2}} \quad (4.1)$$

$$AARE(\%) = \frac{1}{N} \sum_{i=1}^N \left[\frac{E_i - C_i}{E_i} \right] \quad (4.2)$$

where E is the flow stress values obtained from experimental data, C is the calculated flow stress and, \bar{E} is the mean values of experimental data and \bar{C} is the mean values calculated flow stress.

Figure 7 shows the correlation of experimental and calculated (predicted) flow stress data. Table 6 shows the calculated statistical parameter values for the two steels investigated. The results indicate that the developed physically-based constitutive model has higher accuracy in predicting the flow stress of the two steels. The calculated flow stress values had an excellent correlation compared to the experimental flow stress values. The $AARE$ values were 2.18% (P92-A) and 2.20% (P92-B), while the R values were the same for the two steels (0.99). These parameter values show that models have excellent predictability of flow stress behaviour. The analysis indicates that the physically-based model is applicable in determining the stress flow pattern of P92 steels for any given deformation conditions. Hence, it can be of use in most industrial metal-forming applications.

5 Conclusion

The study observed the following:

1. The flow stress values increase with a decrease in the temperature or an increase in strain rate and vice versa.

The results show that flow stress values depend on deformation conditions.

2. It is possible to obtain the material constants B , n , and α in the physically-based model using the self-diffusion activation energy of austenite iron during forming at different loading conditions.
3. The constitutive analysis results show that this model is accurate and reliable in analysing the flow stress behaviour of metals and alloys, hence an alternative technique to analyse metal flow patterns. From the results, the effect of the chemical composition of the two steels on the flow stress behaviour was insignificant.
4. The statistical analysis results show that the physically-based equations developed for the two steels exhibited high accuracy in predicting the metal flow pattern of the two P92 steel investigated. The predicted and experimental data had a good correlation.

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Data availability The authors confirm that the data supporting the findings of this study are available within the article.

Code availability Not applicable.

Declarations

Ethics approval This research follows all ethical standards.

Consent to participate All authors agree.

Consent for publication The authors gave consent for this publication.

Conflict of interest The authors declare no competing interests.

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References

1. Tabuchi M, Watanabe T, Kubo K, Matsui M, Kinugawa J, Abe F (2001) Creep crack growth behavior in the HAZ of weldments of W containing high Cr steel. Int J Press Vessel Pip

- 78(11–12):779–784. [https://doi.org/10.1016/S0308-0161\(01\)00090-4](https://doi.org/10.1016/S0308-0161(01)00090-4)
2. David SA, Siefert JA, Feng Z (2013) Welding and weldability of candidate ferritic alloys for future advanced ultrasupercritical fossil power plants. *Sci Technol Weld Joining* 18(8):631–651. <https://doi.org/10.1179/1362171813Y.0000000152>
 3. Czyska-filemonowicz A, Zielińska-lipiec A, Ennis PJ (2006) Modified 9 % Cr Steels for Advanced Power Generation : Microstructure and Properties. *J Achiev Mater Manuf Eng* 19(2):43–48 [Online]. Available: http://www.edu.ptnss.pb.journalamme.org/papers_vol19_2/1309.pdf
 4. Nagode A, Kosec L, Ule B, Kosec G (2011) Review of creep resistant alloys for power plant applications. *Metalurgija* 50(1):45–48
 5. Oñoro J (2006) Weld metal microstructure analysis of 9–12% Cr steels. *Int J Press Vessel Pip* 83(7):540–545. <https://doi.org/10.1016/j.ijpvp.2006.03.005>
 6. Peng YQ, Chen TC, Chung TJ, Jeng SL, Huang RT, Tsay LW (2017) Creep rupture of the simulated HAZ of T92 steel compared to that of a T91 steel. *Materials (Basel)* 10(2):139. <https://doi.org/10.3390/ma10020139>
 7. Ennis PJ, Czyska-Filemonowicz A (2003) Recent advances in creep resistant steels for power plant applications. *Sadhana* 28(3):709–730. <https://doi.org/10.1007/BF02706455>
 8. Ohgami M, Naoi H, Kinbara S, Mimura H, Ikemoto T, Fujita T (1997) Development of 9CrW tube, pipe and forging for ultra supercritical power plant boilers. *Nippon Steel Tech Rep* 72:59–64
 9. Samantaray D, Mandal S, Bhaduri AK (2011) Optimization of hot working parameters for thermo-mechanical processing of modified 9Cr-1Mo (P91) steel employing dynamic materials model. *Mater Sci Eng A* 528(15):5204–5211. <https://doi.org/10.1016/j.msea.2011.03.025>
 10. Rajput SK (2014) Physical Simulation of Hot Deformation of Low-Carbon Ti-Nb Microalloyed Steel and Microstructural Studies. *J Mater Eng and Perform* 8:2930. <https://doi.org/10.1007/s11665-014-1059-8>
 11. Mehtonen SV, Karjalainen LP, Porter DA (2014) Modeling of the high temperature flow behavior of stabilized 12–27wt% Cr ferritic stainless steels. *Mater Sci Eng A* 607:44–52. <https://doi.org/10.1016/j.msea.2014.03.124>
 12. Cai Z, Ji H, Pei W, Wang B, Huang X, Li Y (2019) Constitutive equation and model validation for 33Cr23Ni8Mn3N heat-resistant steel during hot compression. *Results Phys* 15(July):102633. <https://doi.org/10.1016/j.rinp.2019.102633>
 13. Nayan N, Singh G, Souza P, Murty S, Venkatesh M, Shivram B, Narayanan P, Mohan M, Jha S (2021) Hot workability and microstructure control in Monel®400 (Ni–30Cu) alloy: An approach using processing map, constitutive equation and deformation modeling. *Mater Sci Eng A* 825(July):141855. <https://doi.org/10.1016/j.msea.2021.141855>
 14. He A, Xie G, Zhang H, Wang X (2013) A comparative study on Johnson-Cook, modified Johnson-Cook and Arrhenius-type constitutive models to predict the high temperature flow stress in 20CrMo alloy steel. *Mater Des* 52:677–685. <https://doi.org/10.1016/j.matdes.2013.06.010>
 15. Zhang ZJ, Dai GZ, Wu SN, Dong LX, Liu LL (2009) Simulation of 42CrMo steel billet upsetting and its defects analyses during forming process based on the software DEFORM-3D. *Mater Sci Eng A* 499(1–2):49–52. <https://doi.org/10.1016/j.msea.2007.11.135>
 16. El Wahabi M, Cabrera JM, Prado JM (2003) Hot working of two AISI 304 steels: A comparative study. *Mater Sci Eng A* 343(1–2):116–125. [https://doi.org/10.1016/S0921-5093\(02\)00357-X](https://doi.org/10.1016/S0921-5093(02)00357-X)
 17. Liu CY, Zhang RJ, Yan YN (2011) Hot deformation behaviour and constitutive modelling of P92 heat resistant steel. *Mater Sci Technol* 27(8):1281–1286. <https://doi.org/10.1179/026708310X12683158443323>
 18. Obiko J, Chown LH, Whitefield DJ (2019) Warm deformation behaviour of P92 steel. *Mater Res Express* 6(12):1265j7. <https://doi.org/10.1088/2053-1591/ab5e9c>
 19. Li N, Zhao C, Jiang Z, Zhang H (2019) Flow behavior and processing maps of high-strength low-alloy steel during hot compression. *Mater Charact* 153(October 2018):224–233. <https://doi.org/10.1016/j.matchar.2019.05.009>
 20. Yamanaka K, Matsumoto H, Chiba A (2019) A Constitutive Model and Processing Maps Describing the High-Temperature Deformation Behavior of Ti-17 Alloy in the β -Phase Field. *Adv Eng Mater* 21(2):1–8. <https://doi.org/10.1002/adem.201800775>
 21. Sellars CM, McTegart WJ (1966) On the mechanism of hot deformation. *Acta Metall* 14(9):1136–1138. [https://doi.org/10.1016/0001-6160\(66\)90207-0](https://doi.org/10.1016/0001-6160(66)90207-0)
 22. Samantaray D, Mandal S, Bhaduri AK (2010) Constitutive analysis to predict high-temperature flow stress in modified 9Cr-1Mo (P91) steel. *Mater Des* 31(2):981–984. <https://doi.org/10.1016/j.matdes.2009.08.012>
 23. Shi RX, Liu ZD (2011) Hot deformation behavior of P92 steel used for ultra-super-critical power plants. *J Iron Steel Res Int* 18(7):53–58. [https://doi.org/10.1016/S1006-706X\(11\)60090-3](https://doi.org/10.1016/S1006-706X(11)60090-3)
 24. Sun SL, Zhang MG, He WW (2010) Hot deformation behavior and hot processing map of P92 steel. *Adv Mater Res* 37–101:290–295. <https://doi.org/10.4028/www.scientific.net/AMR.97-101.290>
 25. Ning Y, Fu MW, Hou H, Yao Z, Guo H (2011) Hot deformation behavior of Ti-5.0Al-2.40Sn-2.02Zr-3.86Mo-3.91Cr alloy with an initial lamellar microstructure in the α + β phase field. *Mater Sci Eng A* 528(3):1812–1818. <https://doi.org/10.1016/j.msea.2010.11.019>
 26. Wang L, Liu F, Cheng JJ, Zuo Q, Chen CF (2015) Hot deformation characteristics and processing map analysis for Nickel-based corrosion resistant alloy. *J Alloys Compd* 623:69–78. <https://doi.org/10.1016/j.jallcom.2014.10.034>
 27. Mirzadeh H, Cabrera JM, Najafizadeh A (2011) Constitutive relationships for hot deformation of austenite. *Acta Mater* 59(16):6441–6448. <https://doi.org/10.1016/j.actamat.2011.07.008>
 28. Cabrera JM, Al Omar A, Jonas JJ, Prado JM (1997) Modeling the flow behavior of a medium carbon microalloyed steel under hot working conditions. *Metall Mater Trans A Phys Metall Mater Sci* 28(11):2233–2243. <https://doi.org/10.1007/s11661-997-0181-8>
 29. Motallebi R, Savaedi Z, Mirzadeh H (2022) Additive manufacturing – A review of hot deformation behavior and constitutive modeling of flow stress. *Curr Opin Solid State Mater Sci* 26(3):100992. <https://doi.org/10.1016/j.cossms.2022.100992>
 30. Savaedi Z, Motallebi R, Mirzadeh H (2022) A review of hot deformation behavior and constitutive models to predict flow stress of high-entropy alloys. *J Alloys Compd* 903:163964. <https://doi.org/10.1016/j.jallcom.2022.163964>
 31. Ashby MF (1972) A first report on deformation-mechanism maps. *Acta Metall* 20(7):887–897. [https://doi.org/10.1016/0001-6160\(72\)90082-X](https://doi.org/10.1016/0001-6160(72)90082-X)
 32. Liu L, Wu Y, Gong H, Li S, Ahmad AS (2018) A physically based constitutive model and continuous dynamic recrystallization behavior analysis of 2219 aluminum alloy during hot deformation process. *Materials (Basel)* 11(8):1443. <https://doi.org/10.3390/ma11081443>
 33. Lian Wei H, Quan Liu G, He Zhang M (2014) Physically based constitutive analysis to predict flow stress of medium carbon and vanadium microalloyed steels. *Mater Sci Eng A* 602:127–133. <https://doi.org/10.1016/j.msea.2014.02.068>
 34. Lu C, Shi J, Wang J (2021) Physically based constitutive modeling for Ti17 alloy with original basketweave microstructure in β forging: A comparison of three approaches. *Mater Charact* 181(June):111455. <https://doi.org/10.1016/j.matchar.2021.111455>

35. Mirzadeh H (2015) Simple physically-based constitutive equations for hot deformation of 2024 and 7075 aluminum alloys. *Trans Nonferrous Met Soc China (English Ed)* 25(5):1614–1618
36. Feng R, Bao Y, Ding Y, Chen M, Ge Y, Xie L (2022) Three different mathematical models to predict the hot deformation behavior of TA32 titanium alloy. *J Mater Res* 37:1–14. <https://doi.org/10.1557/s43578-022-00532-2>
37. Li J, Yu R, Xu G, Chen C, Ha Y, Song L, Zhang H (2022) Research on Hot Deformation Behavior of F92 Steel Based on Stress Correction. *Metals (Basel)* 12(5):1–14. <https://doi.org/10.3390/met12050698>
38. Pandey C, Mahapatra M, Kumar P, Thakre JG, Saini N (2018) Role of evolving microstructure on the mechanical behaviour of P92 steel welded joint in as-welded and post weld heat treated state. *J Mater Process Technol* 263(June 2018):241–255. <https://doi.org/10.1016/j.jmatprotec.2018.08.032>
39. Barbadikar D, Deshmukh G, Maddi L, Laha K, Parameswaran P, Ballal A, Peshwe D, Paretkar R, Nandagopal M, Mathew M (2015) Effect of normalizing and tempering temperatures on microstructure and mechanical properties of P92 steel. *Int J Press Vessel Pip* 132–133:97–105. <https://doi.org/10.1016/j.ijpvp.2015.07.001>
40. Pandey C, Mahapatra M, Kumar P, Saini N (2018) Homogenization of P91 weldments using varying normalizing and tempering treatment. *Mater Sci Eng A* 710(June 2017):86–101. <https://doi.org/10.1016/j.msea.2017.10.086>
41. Abson DJ, Rothwell JS (2013) Review of type IV cracking of weldments in 9–12%Cr creep strength enhanced ferritic steels. *Int Mater Rev* 58(8):437–473. <https://doi.org/10.1179/1743280412Y.0000000016>
42. Shah N, Chakravartty J, Sarkar A (2019) Comparative study on hot deformation behaviour of P91, RAFM-CLAM and Low C RAFM-CLAM steels by processing maps. *Fusion Eng Des* 138:109–118. <https://doi.org/10.1016/j.fusengdes.2018.11.010>
43. Zhu L, He J, Zhang Y (2018) A two-stage constitutive model of X12CrMoWVNbN10–1–1 steel during elevated temperature. *Mater Res Express* 5(2):026505. <https://doi.org/10.1088/2053-1591/aaa911>
44. Carsí M, Peñalba F, Rieiro I, Ruano OA (2011) High temperature workability behavior of a modified P92 steel. *Int J Mater Res* 102(11):1378–1383. <https://doi.org/10.3139/146.110603>
45. Ennis P, Zielinska-Lipiec A, Wachter O, Czyska-Filemonowicz A (1997) Microstructural stability and creep rupture strength of the martensitic steel P92 for advanced power plant. *Acta Mater* 45(12):4901–4907. [https://doi.org/10.1016/S1359-6454\(97\)00176-6](https://doi.org/10.1016/S1359-6454(97)00176-6)
46. Sawada K, Kubo K, Abe F (2001) Creep behavior and stability of MX precipitates at high temperature in 9Cr-0.5Mo-1.8W-VNb steel. *Mater Sci Eng A* 319(321):784–787. [https://doi.org/10.1016/S0921-5093\(01\)00973-X](https://doi.org/10.1016/S0921-5093(01)00973-X)
47. Nkhoma R, Siyasiya C, Stumpf W (2014) Hot workability of AISI 321 and AISI 304 austenitic stainless steels. *J Alloys Compd* 595(April 2014):103–112. <https://doi.org/10.1016/j.jallcom.2014.01.157>
48. Rao K, Prasad Y (1986) High temperature deformation kinetics of Al–4Mg alloy. *J Mech Work Technol* 13:83–95
49. Yang X, Li W (2015) Flow Behavior and Processing Maps of a Low-Carbon Steel During Hot Deformation. *Metall Mater Trans A Phys Metall Mater Sci* 46(12):6052–6064
50. Yang Z, Dai L, Chu C, Zhang F, Wang L, Xiao A (2017) Effect of Aluminum Alloying on the Hot Deformation Behavior of Nanobainite Bearing Steel. *J Mater Eng Perform* 26(12):5954–5962. <https://doi.org/10.1007/s11665-017-3018-7>
51. Sun C, Zuo X, Xiang Y, Yang J (2016) Investigation on Hot Deformation Behavior and Hot Processing Map of BSTMU601 Super-Alloy. *Metals (Basel)* 6(3):70. <https://doi.org/10.3390/met6030070>

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