# The geometry of billiards in ellipses and their poncelet grids 

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#### Abstract

The goal of this paper is an analysis of the geometry of billiards in ellipses, based on properties of confocal central conics. The extended sides of the billiards meet at points which are located on confocal ellipses and hyperbolas. They define the associated Poncelet grid. If a billiard is periodic then it closes for any choice of the initial vertex on the ellipse. This gives rise to a continuous variation of billiards which is called billiard motion though it is neither a Euclidean nor a projective motion. The extension of this motion to the associated Poncelet grid leads to new insights and invariants.


Mathematics Subject Classification. Primary 51N35, Secondary 51N20, 52C30, 37D50.
Keywords. Billiard in ellipse, Caustic, Poncelet grid, Confocal conics, Billiard motion, Canonical parametrization.

## 1. Introduction

A billiard is the trajectory of a mass point within a domain with ideal physical reflections in the boundary. Already for two centuries, billiards in ellipses have attracted the attention of mathematicians, beginning with J.-V. Poncelet, C.G.J. Jacobi and A. Cayley. One basis for the investigations was the theory of confocal conics. In 2005 S . Tabachnikov published a book on various aspects of billiards, including their role as completely integrable systems [29]. In several publications and in the book [13], V. Dragović and M. Radnović studied billiards, also in higher dimensions, from the viewpoint of dynamical systems.

Computer animations of billiards in ellipses, which were carried out by Reznik [24], stimulated a new vivid interest on this well studied topic, where algebraic and analytic methods are meeting (see, e.g., $[2,3,11,21-23]$ and many further
references in [24]). These papers focus on invariants of periodic billiards when the vertices vary on the ellipse while the caustic remains fixed. This variation is called billiard motion though neither angles nor side lengths remain fixed; and it is not a projective motion preserving the circumscribed ellipse.
The goal of this paper is a geometric analysis of billiards in ellipses and their associated Poncelet grid, starting from properties of confocal conics. We concentrate on a certain symmetry between the vertices of any billiard and the contact points with the caustic, which can be an ellipse or hyperbola. Billiard motions induce motions of associated billiards with the same caustic and circumscribed confocal ellipses.

## 2. Metric properties of confocal conics

A family of confocal central conics (Fig. 1) is given by

$$
\begin{equation*}
\frac{x^{2}}{a^{2}+k}+\frac{y^{2}}{b^{2}+k}=1, \text { where } k \in \mathbb{R} \backslash\left\{-a^{2},-b^{2}\right\} \tag{2.1}
\end{equation*}
$$

serves as a parameter in the family. All these conics share the focal points

$$
\begin{equation*}
F_{1,2}=( \pm d, 0), \text { where } d^{2}:=a^{2}-b^{2} \tag{2.2}
\end{equation*}
$$

The confocal family sends through each point $P$ outside the common axes of symmetry two orthogonally intersecting conics, one ellipse and one hyperbola $\left[15\right.$, p. 38]. The parameters $\left(k_{e}, k_{h}\right)$ of these two conics define the elliptic coordinates of $P$ with

$$
-a^{2}<k_{h}<-b^{2}<k_{e}
$$

If $(x, y)$ are the cartesian coordinates of $P$, then $\left(k_{e}, k_{h}\right)$ are the roots of the quadratic equation

$$
\begin{equation*}
k^{2}+\left(a^{2}+b^{2}-x^{2}-y^{2}\right) k+\left(a^{2} b^{2}-b^{2} x^{2}-a^{2} y^{2}\right)=0 \tag{2.3}
\end{equation*}
$$

while conversely

$$
\begin{equation*}
x^{2}=\frac{\left(a^{2}+k_{e}\right)\left(a^{2}+k_{h}\right)}{d^{2}}, \quad y^{2}=-\frac{\left(b^{2}+k_{e}\right)\left(b^{2}+k_{h}\right)}{d^{2}} . \tag{2.4}
\end{equation*}
$$

Let $(a, b)=\left(a_{c}, b_{c}\right)$ be the semiaxes of the ellipse $c$ with $k=0$. Then, for points $P$ on a confocal ellipse $e$ with semiaxes $\left(a_{e}, b_{e}\right)$ and $k=k_{e}>0$, i.e., exterior to $c$, the standard parametrization yields

$$
\begin{align*}
P= & (x, y)=\left(a_{e} \cos t, b_{e} \sin t\right), 0 \leq t<2 \pi \\
& \text { with } a_{e}^{2}=a_{c}^{2}+k_{e}, b_{e}^{2}=b_{c}^{2}+k_{e} \tag{2.5}
\end{align*}
$$

For the elliptic coordinates $\left(k_{e}, k_{h}\right)$ of $P$ follows from (2.3) that

$$
k_{e}+k_{h}=a_{e}^{2} \cos ^{2} t+b_{e}^{2} \sin ^{2} t-a_{c}^{2}-b_{c}^{2} .
$$

After introducing the respective tangent vectors of $e$ and $c$, namely

$$
\begin{align*}
& \mathbf{t}_{e}(t):=\left(-a_{e} \sin t, b_{e} \cos t\right),  \tag{2.6}\\
& \mathbf{t}_{c}(t):=\left(-a_{c} \sin t, b_{c} \cos t\right),
\end{align*} \text { where }\left\|\mathbf{t}_{e}\right\|^{2}=\left\|\mathbf{t}_{c}\right\|^{2}+k_{e},
$$

we obtain ${ }^{1}$

$$
\begin{equation*}
k_{h}=k_{h}(t)=-\left(a_{c}^{2} \sin ^{2} t+b_{c}^{2} \cos ^{2} t\right)=-\left\|\mathbf{t}_{c}(t)\right\|^{2}=-\left\|\mathbf{t}_{e}(t)\right\|^{2}+k_{e} \tag{2.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\mathbf{t}_{e}(t)\right\|^{2}=k_{e}-k_{h}(t) \tag{2.8}
\end{equation*}
$$

Note that points on the confocal ellipses $e$ and $c$ with the same parameter $t$ have the same coordinate $k_{h}$. Consequently, they belong to the same confocal hyperbola (Figs. 5 and 8). Conversely, points of $e$ or $c$ on this hyperbola have a parameter out of $\{t,-t, \pi+t, \pi-t\}$.

Normal vectors of $e$ and $c$ can be defined respectively as

$$
\begin{align*}
& \mathbf{n}_{e}(t):=\left(\frac{\cos t}{a_{e}}, \frac{\sin t}{b_{e}}\right),  \tag{2.9}\\
& \mathbf{n}_{c}(t):=\left(\frac{\cos t}{a_{c}}, \frac{\sin t}{b_{c}}\right),
\end{align*} \quad \text { where }\left\|\mathbf{n}_{c}(t)\right\|=\frac{\left\|\mathbf{t}_{c}(t)\right\|}{a_{c} b_{c}} .
$$

We complete with two useful relations between the parameter $t$ and the second elliptic coordinate $k_{h}(t)$ :

$$
\begin{equation*}
\tan ^{2} t=-\frac{b_{c}^{2}+k_{h}(t)}{a_{c}^{2}+k_{h}(t)} \quad \text { and } \quad \sin t \cos t=\frac{a_{h} b_{h}}{d^{2}} \tag{2.10}
\end{equation*}
$$

with $a_{h}$ and $b_{h}$ as semiaxes of the hyperbola corresponding to the parameter $t$, i.e., $a_{h}^{2}=a_{c}^{2}+k_{h}$ and $b_{h}^{2}=-\left(b_{c}^{2}+k_{h}\right)$.

Proof. From (2.7) follows

$$
k_{h}=-\frac{a_{c}^{2} \tan ^{2} t+b_{c}^{2}}{1+\tan ^{2} t}, \quad \text { hence } \quad \tan ^{2} t\left(a_{c}^{2}+k_{h}\right)=-b_{c}^{2}-k_{h}
$$

and

$$
\sin t \cos t=\frac{\tan t}{1+\tan ^{2} t}=\frac{\sqrt{-\left(b_{c}^{2}+k_{h}\right)\left(a_{c}^{2}+k_{h}\right)}}{a_{c}^{2}-b_{c}^{2}}=\frac{a_{h} b_{h}}{d^{2}} .
$$

Referring to Fig. 1, the following lemma addresses an important property of confocal conics (note, e.g., [15, pp. 38 and 309]).

Lemma 2.1. The tangents drawn from any fixed point $P$ to the conics of a confocal family share the axes of symmetry, which are tangent to the two conics passing through $P$.

This means, if a ray is reflected at $P$ in one of the conics passing through, then the incoming and the outgoing ray contact the same confocal ellipse or hyperbola.
Below, we report about results concerning a pair of confocal conics. Due to their meaning for billiards in ellipses, we restrict ourselves to pairs $(e, c)$ of confocal ellipses with $c$ in the interior $e$, and we call $c$ the caustic (Fig. 2).

[^0]

Figure 1 The tangents from the point $P$ to the conics of a confocal family are symmetric w.r.t. the tangents at $P$ to the confocal conics passing through $P$


Figure 2 Periodic billiard $P_{1} P_{2} \ldots P_{5}$ inscribed in the ellipse $e$ with the caustic $c$

Lemma 2.2. Let $P=\left(a_{e} \cos t, b_{e} \sin t\right)$ with elliptic coordinates $\left(k_{e}, k_{h}\right)$ be a point on the ellipse $e$ with $k_{e}>0$ and $c$ be the confocal ellipse with $k=0$. Then, the angle $\theta(t) / 2$ between the tangent at $P$ to $e$ and any tangent from $P$ to $c$ satisfies

$$
\begin{align*}
\sin ^{2} \frac{\theta}{2} & =\frac{k_{e}}{\left\|\boldsymbol{t}_{e}(t)\right\|^{2}}=\frac{k_{e}}{k_{e}-k_{h}}, \tan \frac{\theta}{2}= \pm \sqrt{-\frac{k_{e}}{k_{h}}}  \tag{2.11}\\
\cos \theta & =1-\frac{2 k_{e}}{\left\|\boldsymbol{t}_{e}(t)\right\|^{2}}=\frac{k_{h}+k_{e}}{k_{h}-k_{e}}, \quad \sin \theta= \pm \frac{2 \sqrt{-k_{e} k_{h}}}{k_{e}-k_{h}} \tag{2.12}
\end{align*}
$$

Proof. The tangent $t_{P}$ to $e$ at $P=\left(a_{e} \cos t, b_{e} \sin t\right)$ in direction of $\mathbf{t}_{e}$ has the slope

$$
f:=\tan \alpha_{1}=\frac{-b_{e} \cos t}{a_{e} \sin t}
$$

If $s_{1}$ and $s_{2}$ denote the slopes of the tangents from $P$ to $c$, then they satisfy

$$
y-b_{e} \sin t=s_{i}\left(x-a_{e} \cos t\right), \quad i=1,2 .
$$

As tangents of $c$, their homogeneous line coordinates

$$
\left(u_{0}: u_{1}: u_{2}\right)=\left(\left(b_{e} \sin t-s_{i} a_{e} \cos t\right): s_{i}:-1\right)
$$

must satisfy the tangential equation $-u_{0}^{2}+a_{c}^{2} u_{1}^{2}+b_{c}^{2} u_{2}^{2}=0$ of $c$. This results in a quadratic equation for the unknown $s$, namely

$$
\left(a_{e}^{2} \sin ^{2} t-k_{e}\right) s^{2}+2 a_{e} b_{e} s \sin t \cos t+\left(b_{e}^{2} \cos ^{2} t-k_{e}\right)=0 .
$$

We conclude

$$
s_{1}+s_{2}=\frac{-2 a_{e} b_{e} \sin t \cos t}{a_{e}^{2} \sin ^{2} t-k_{e}} \text { and } s_{1} s_{2}=\frac{b_{e}^{2} \cos ^{2} t-k_{e}}{a_{e}^{2} \sin ^{2} t-k_{e}} .
$$

The slopes $f=\tan \alpha_{1}$ of $t_{P}$ and $\tan \alpha_{2}=s_{1}$ or $s_{2}$ of the tangents to $c$ imply for the enclosed signed angle $\theta(t) / 2$ (for brevity, we often suppress the parameter $t)$

$$
\tan \frac{\theta}{2}=\tan \left(\alpha_{1}-\alpha_{2}\right)=\frac{s_{1}-f}{1+s_{1} f}=\frac{f-s_{2}}{1+s_{2} f},
$$

hence

$$
\tan ^{2} \frac{\theta}{2}=\frac{\left(s_{1}-f\right)\left(f-s_{2}\right)}{\left(1+s_{1} f\right)\left(1+s_{2} f\right)}=\frac{f\left(s_{1}+s_{2}\right)-s_{1} s_{2}-f^{2}}{f\left(s_{1}+s_{2}\right)+1+f^{2} s_{1} s_{2}} .
$$

After some computations, we obtain

$$
\tan ^{2} \frac{\theta}{2}=\frac{k_{e}}{a_{e}^{2} \sin ^{2} t+b_{e}^{2} \cos ^{2} t-k_{e}}=\frac{k_{e}}{\left\|\mathbf{t}_{e}\right\|^{2}-k_{e}}=\frac{k_{e}}{\left\|\mathbf{t}_{c}\right\|^{2}},
$$

therefore

$$
\cot ^{2} \frac{\theta}{2}=\frac{\left\|\mathbf{t}_{e}\right\|^{2}}{k_{e}}-1 \text { and } \sin ^{2} \frac{\theta}{2}=\frac{1}{1+\cot ^{2} \frac{\theta}{2}}=\frac{k_{e}}{\left\|\mathbf{t}_{e}\right\|^{2}},
$$

where $k_{e}=a_{e}^{2}-a_{c}^{2}=b_{e}^{2}-b_{c}^{2}$, and finally

$$
\cos \theta=1-2 \sin ^{2} \frac{\theta}{2}=1-\frac{2 k_{e}}{\left\|\mathbf{t}_{c}\right\|^{2}+k_{e}}=\frac{\left\|\mathbf{t}_{c}\right\|^{2}-k_{e}}{\left\|\mathbf{t}_{c}\right\|^{2}+k_{e}} .
$$

Remark 2.3. A change of the origin $k=0$ for the elliptic coordinates in a family of confocal conics corresponds to a shift of the coordinates. Hence, if in Lemma 2.2 the ellipse $c$ is replaced by another confocal conic with the coordinate $k$, then the formulas (2.11) and (2.12) remain valid under the condition that we replace $k_{e}$ by $k_{e}-k$ and $k_{h}$ by $k_{h}-k$.

Lemma 2.4. Let $P_{1} P_{2}$ be a chord of the ellipse e, which contacts the caustic $c$ at the point $Q_{1}$. Then the signed distances of the line $\left[P_{1}, P_{2}\right]$ to the center $O$ and to the pole $R_{1}$ w.r.t. e have the constant product $-k_{e}$. The lines $\left[P_{1}, P_{2}\right]$ and $\left[Q_{1}, R_{1}\right]$ are orthogonal (Fig. 2).

Proof. Let the side $P_{1} P_{2}$ touch the caustic $c$ at the point $Q_{1}=\left(a_{c} \cos t_{1}^{\prime}, b_{c} \sin \right.$ $\left.t_{1}^{\prime}\right)$. Then the Hessian normal form of the spanned line $t_{Q}=\left[P_{1}, P_{2}\right]$ reads

$$
t_{Q}: \frac{b_{c} \cos t_{1}^{\prime} x+a_{c} \sin t_{1}^{\prime} y-a_{c} b_{c}}{\sqrt{b_{c}^{2} \cos ^{2} t_{1}^{\prime}+a_{c}^{2} \sin ^{2} t_{1}^{\prime}}}=0
$$

Its pole w.r.t. $e$ has the coordinates

$$
\begin{equation*}
R_{1}=\left(\frac{a_{e}^{2} \cos t_{1}^{\prime}}{a_{c}}, \frac{b_{e}^{2} \sin t_{1}^{\prime}}{b_{c}}\right) \tag{2.13}
\end{equation*}
$$

This yields for the signed distances to the line $t_{Q}$

$$
\begin{equation*}
\overline{O t_{Q}}=\frac{-a_{c} b_{c}}{\left\|\mathbf{t}_{c}\left(t_{1}^{\prime}\right)\right\|} \tag{2.14}
\end{equation*}
$$

and

$$
\begin{equation*}
\overline{R_{1} t_{Q}}=\overline{R_{1} Q_{1}}=\frac{k_{e}\left(b_{c}^{2} \cos ^{2} t_{1}^{\prime}+a_{c}^{2} \sin ^{2} t_{1}^{\prime}\right)}{a_{c} b_{c}\left\|\mathbf{t}_{c}\left(t_{1}^{\prime}\right)\right\|}=\frac{k_{e}\left\|\mathbf{t}_{c}\left(t_{1}^{\prime}\right)\right\|}{a_{c} b_{c}} . \tag{2.15}
\end{equation*}
$$

Thus, we obtain a constant product $\overline{O t_{Q}} \cdot \overline{R_{1} t_{Q}}=-k_{e}$, as stated in Lemma 2.4.
The last statement holds since $R_{1}$ and $Q_{1}$ are the poles of [ $P_{1}, P_{2}$ ] w.r.t. the confocal conics $e$ and $c$. It is wellknown that the poles of any line $\ell$ w.r.t. confocal conics lie on a line $\ell^{*}$ orthogonal to $\ell$ (see, e.g., [15, p. 340]).

## 3. Confocal conics and billiards

By virtue of Lemma 2.1, all sides of a billiard inscribed to the ellipse $e$ with parameter $k=k_{e}$ are tangent to a fixed conic $c$ confocal with $e$ (Fig. 2). The caustic $c$ with parameter $k_{c}$ can be a smaller ellipse with $-b^{2}<k_{c}<k_{e}$ or a hyperbola with $-a^{2}<k_{c}<-b^{2}$ or, in the limiting case with $k_{c}=-b^{2}$, consist of the pencils of lines with the focal points $F_{1}, F_{2}$ of $c$ as carriers. At the beginning, we confine ourselves to an ellipse with $k_{c}=0$ (Fig. 2), and we speak of an elliptic billiard. Only the Figs. 10, 11 and 12 show (periodic) billiards in $e$ with a hyperbola as caustic, called hyperbolic billards.

For billiards $\ldots P_{1} P_{2} P_{3} \ldots$ in the ellipse $e$ and with the ellipse $c$ as caustic, we assume from now on a counter-clockwise order and signed exterior angles $\theta_{1}, \theta_{2}, \theta_{3}, \ldots$ (see Fig. 2). The tangency points $Q_{1}, Q_{2}, \ldots$ of the billiard's sides $P_{1} P_{2}, P_{2} P_{3}, \ldots$ with the caustic $c$ subdivide the sides into two segments. We denote the lengths of the segments adjacent to $P_{i}$ as

$$
\begin{equation*}
l_{i}:=\overline{P_{i} Q_{i}} \text { and } r_{i}:=\overline{P_{i} Q_{i-1}} . \tag{3.1}
\end{equation*}
$$

Based on the parametrizations $\left(a_{e} \cos t, b_{e} \sin t\right)$ of $e$ and $\left(a_{c} \cos t^{\prime}, b_{c} \sin t^{\prime}\right)$ of $c$, we denote the respective parameters of $P_{1}, Q_{1}, P_{2}, Q_{2}, P_{3}, \ldots$ with $t_{1}, t_{1}^{\prime}$, $t_{2}, t_{2}^{\prime}, t_{3}, \ldots$ in strictly increasing order.
The following two lemmas deal with sides of billiards in the ellipse $e$.
Lemma 3.1. The connecting line $\left[P_{i}, P_{i+1}\right]$ of the vertices with respective parameters $t_{1}, t_{2}$ on $e$ contacts the caustic $c$ if and only if

$$
\frac{a_{c}^{2}}{a_{e}^{2}} \cos ^{2} \frac{t_{1}+t_{2}}{2}+\frac{b_{c}^{2}}{b_{e}^{2}} \sin ^{2} \frac{t_{1}+t_{2}}{2}=\cos ^{2} \frac{t_{1}-t_{2}}{2}
$$

This is equivalent to

$$
\sin ^{2} \frac{t_{1}-t_{2}}{2}=\frac{k_{e}}{a_{e} b_{e}}\left\|\boldsymbol{t}_{e}\left(\frac{t_{1}+t_{2}}{2}\right)\right\|^{2} .
$$

Proof. The line connecting the points $\left(a_{e} \cos t_{i}, b_{e} \sin t_{i}\right), i=1,2$, has homogeneous line coordinates ( $u_{0}: u_{1}: u_{2}$ ) equal to

$$
\left(a_{e} b_{e}\left(\cos t_{1} \sin t_{2}-\sin t_{1} \cos t_{2}\right): b_{e}\left(\sin t_{1}-\sin t_{2}\right): a_{e}\left(\cos t_{2}-\cos t_{1}\right)\right) .
$$

It contacts the caustic $c$ if $-u_{0}^{2}+a_{c}^{2} u_{1}^{2}+b_{c}^{2} u_{2}^{2}=0$, i.e.,

$$
\begin{aligned}
& a_{c}^{2} b_{e}^{2} \sin ^{2} \frac{t_{1}-t_{2}}{2} \cos ^{2} \frac{t_{1}+t_{2}}{2}+b_{c}^{2} a_{e}^{2} \sin ^{2} \frac{t_{1}-t_{2}}{2} \sin ^{2} \frac{t_{1}+t_{2}}{2} \\
& =a_{e}^{2} b_{e}^{2} \sin ^{2} \frac{t_{2}-t_{1}}{2} \cos ^{2} \frac{t_{2}-t_{1}}{2}
\end{aligned}
$$

Under the condition $\sin \left[\left(t_{1}-t_{2}\right) / 2\right] \neq 0$ we obtain the first claimed equation. The second follows after the substitutions $a_{c}^{2}=a_{e}^{2}-k_{e}$ and $b_{c}^{2}=b_{e}^{2}-k_{e}$ from

$$
1-\frac{k_{e}}{a_{e}^{2} b_{e}^{2}}\left(b_{e}^{2} \cos ^{2} \frac{t_{1}+t_{2}}{2}+a_{e}^{2} \sin ^{2} \frac{t_{1}+t_{2}}{2}\right)=\cos ^{2} \frac{t_{2}-t_{1}}{2} .
$$

by (2.9).
Lemma 3.2. Referring to the notation in Lemma 3.1, if the side $P_{i} P_{i+1}$ contacts the caustic c at $Q_{i}$ with parameter $t_{i}^{\prime}$, then

$$
\sin t_{i}^{\prime}=\frac{b_{c}}{b_{e}} \frac{\sin \frac{t_{i}+t_{i+1}}{2}}{\cos \frac{t_{i}-t_{i+1}}{2}}, \cos t_{i}^{\prime}=\frac{a_{c}}{a_{e}} \frac{\cos \frac{t_{i}+t_{i+1}}{2}}{\cos \frac{t_{i}-t_{i+1}}{2}}, \tan t_{i}^{\prime}=\frac{b_{c} a_{e}}{a_{c} b_{e}} \tan \frac{t_{i}+t_{i+1}}{2} .
$$

Proof. The tangent to $c$ at $Q_{1}$ has the line coordinates

$$
\left(u_{0}: u_{1}: u_{2}\right)=\left(-a_{c} b_{c}: b_{c} \cos t_{1}^{\prime}: a_{c} \sin t_{1}^{\prime}\right),
$$

which must be proportional to those in the proof of Lemma 3.1.
Remark 3.3. 1. The half-angle substitution

$$
\tau_{i}:=\tan \frac{t_{i}}{2} \quad \text { for } i=1,2
$$

allows to express the equation of Lemma 3.1 (for $i=1$ ) in projective coordinates on $e$. We obtain a symmetric biquadratic condition

$$
b_{e}^{2} k_{e} \tau_{1}^{2} \tau_{2}^{2}-b_{c}^{2} a_{e}^{2}\left(\tau_{1}^{2}+\tau_{2}^{2}\right)+2\left(a_{e}^{2} k_{e}+a_{c}^{2} b_{e}^{2}\right) \tau_{1} \tau_{2}+b_{e}^{2} k_{e}=0,
$$



Figure 3 The Joachimsthal integral $J_{e}:=-\left\langle\mathbf{u}_{i}, \mathbf{n}_{e \mid i}\right\rangle$ is constant along $e$
which defines a 2-2-correspondence on $e$ between the endpoints $P_{1}, P_{2}$ of a chord which contacts $c$. This remains valid after iteration, i.e., between the initial point $P_{1}$ and the endpoint $P_{N+1}$ of a billiard after $N$ reflections in $e$.
Now, we recall a classical argument for the underlying Poncelet porism (see also [18] and the references there): A 2-2-correspondence different from the identity keeps fixed at most four points. However, four fixed points on $e$ are already known as contact points between $e$ and the common complex conjugate (isotropic) tangents ${ }^{2}$ with the caustic $c$, since tangents of $e$ remain fixed under the reflection in $e$. If therefore one $N$-sided billiard in $e$ with caustic $c$ closes, then the correspondence is the identity and all billiards close.
2. With the aid of Jacobi's arguments in [20], Lemma 3.2 paves already the way to a representation of the billiard's vertices in terms of Jacobian elliptic functions (note [27]).

Given any billiard $P_{1} P_{2} \ldots$ in the ellipse $e$, let $\mathbf{p}_{i}=\left(x_{i}, y_{i}\right)$ denote the position vector of $P_{i}$ for $i=1,2, \ldots$, while $\mathbf{u}_{1}, \mathbf{u}_{2}, \ldots$ denote the unit vectors of the oriented sides $P_{1} P_{2}, P_{2} P_{3}, \ldots$ (Fig. 3). By (2.9), the vector $\mathbf{n}_{e \mid i}:=\left(x_{i} / a_{e}^{2}, y_{i} / b_{e}^{2}\right)$ is orthogonal to $e$ at $P_{i}$. According to [2, Proposition 2.1], the scalar product

$$
\begin{equation*}
J_{e}:=-\left\langle\mathbf{u}_{i}, \mathbf{n}_{e \mid i}\right\rangle \tag{3.2}
\end{equation*}
$$

is invariant along the billiard in $e$ and called Joachimsthal integral (note also [29, p. 54]).

The invariance of the Joachimsthal integral, which also holds in higher dimensions for billiards in quadrics, is the key result for the integrability of billiards, i.e., in the planar case for the existence of a caustic [2, p. 3]. In our approach, the invariance of $J_{e}$ follows from Lemma 2.2.

Lemma 3.4. The Joachimsthal integral $J_{e}:=-\left\langle\boldsymbol{u}_{i}, \boldsymbol{n}_{e \mid i}\right\rangle$ equals

$$
J_{e}=\frac{\sqrt{k_{e}}}{a_{e} b_{e}}
$$

[^1]with $k_{e}$ as elliptic coordinate of e w.r.t. $c$, i.e., $k_{e}=a_{e}^{2}-a_{c}^{2}=b_{e}^{2}-b_{c}^{2}$.
Proof. From (2.11) follows for the points $\left(a_{e} \cos t, b_{e} \sin t\right)$ of $e$
$$
J_{e}=-\left\langle\mathbf{u}, \mathbf{n}_{e}\right\rangle=-\cos \left(\frac{\pi}{2}+\frac{\theta}{2}\right)\left\|\mathbf{n}_{e}\right\|=\sin \frac{\theta}{2}\left\|\mathbf{n}_{e}\right\|=\sin \frac{\theta}{2} \frac{\left\|\mathbf{t}_{e}\right\|}{a_{e} b_{e}},
$$
hence by (2.11), (2.7), (2.6), and (2.9)
$$
J_{e}^{2}=\sin ^{2} \frac{\theta}{2}\left\|\mathbf{n}_{e}\right\|^{2}=\frac{k_{e}}{\left\|\mathbf{t}_{e}\right\|^{2}}\left\|\mathbf{n}_{e}\right\|^{2}=\frac{k_{e}}{a_{e}^{2} b_{e}^{2}} .
$$

This confirms the claim.

### 3.1. Poncelet grid

The following theorem is the basis for the Poncelet grid associated to each billiard. We formulate and prove a projective version. The special case dealing with confocal conics, has already been published by Chasles [9, p. 841] and later by Böhm in [8, p. 221]. The same theorem was studied in [21,25] and in [1]. In [19], the authors proved it in a differential-geometric way.

In the theorem and proof below, the term conic stands for regular dual conics, i.e., conics seen as the set of tangent lines, but also for pairs of line pencils and for single line pencils with multiplicity two. Expressed in terms of homogeneous line coordinates, the corresponding quadratic forms have rank 3,2 or 1, respectively. Moreover, we use the term range for a pencil of dual conics. The term net denotes a 2-parametric linear system of dual curves of degree 2 . Obviously, conics and ranges included in a net play the role of points and lines of a projective plane within the 5 -dimensional projective space of dual conics. Any two ranges in a net share a conic (compare with [10, Théorèmes I-IV]).

Theorem 3.5. Let c be a regular conic and $A_{1}, B_{1}$ two points such that the tangents $t_{1}, \ldots, t_{4}$ drawn from $A_{1}$ and $B_{1}$ to $c$ form a quadrilateral. Its remaining pairs of opposite vertices are denoted by $\left(A_{i}, B_{i}\right), i=2,3$. Then,

1. for each conic $c_{1}$ passing through $A_{1}$ and $B_{1}$, the range $\mathcal{R}_{c}$ spanned by $c$ and $c_{1}$ contains conics $c_{i}$ passing through $A_{i}$ and $B_{i}$, simultaneously. The tangents at $A_{j}$ and $B_{j}$ to $c_{j}$ for $j=1,2,3$ meet at a common point $R$. If $c_{i}$ has rank 2, then we obtain, as the limit of $c_{i}$, the diagonal $\left[A_{i}, B_{i}\right]$ of the quadrilateral $t_{1}, \ldots, t_{4}$.
2. This result holds also in the limiting case $t_{1}=t_{2}$, where the chord $A_{1} B_{1}$ of $c_{1}$ contacts $c$ at $B_{2}$.

In Fig. 4, the particular case is displayed where $c$ and $c_{1}=e$ span a range $\mathcal{R}_{c}$ of confocal conics (note also [6, Fig. 19]). Then by Lemma 2.1, the tangents at $A_{j}$ and $B_{j}$ to $c_{j}$ are angle bisectors of the quadrilateral. In case of a rank deficiency of $c_{i}$, either one axis of symmetry of the confocal family or the line at infinity shows up as $c_{i}$.


Figure 4 Left: Opposite vertices $A_{i}, B_{i}$ of the quadrilateral $t_{1} \ldots t_{4}$ of tangents to $c$ belong to a conic $c_{i}$ out of the range $\mathcal{R}_{c}$ (Theorem 3.5). Right: The net $\mathcal{N}$ and the ranges $\mathcal{R}_{t}, \mathcal{R}_{c}$ in the projective space of dual conics

Proof. The conics tangent to $t_{1}, \ldots, t_{4}$ define a range $\mathcal{R}_{t}$, which includes for $j=1,2,3$ the pairs of line pencils $\left(A_{j}, B_{j}\right)$ as well as the initial conic $c$. On the other hand, $c$ and $c_{1}$ span a range $\mathcal{R}_{c}$. Since both ranges share the conic $c$, they span a net $\mathcal{N}$ of conics.
The pair $\left(A_{1}, B_{1}\right)$ of line pencils spans together with $c_{1}$ the range of conics sharing the points $A_{1}, B_{1}$ and the tangents there, which meet at a point $R$. This range, which also belongs to $\mathcal{N}$, contains the rank- 1 conic with carrier $R$. Each pair of line pencils $\left(A_{i}, B_{i}\right), i=1,2$, spans with the pencil $R$ again a range within $\mathcal{N}$. This range shares with the range $\mathcal{R}_{c}$ a conic $c_{i}$ passing through $A_{i}$ and $B_{i}$ with respective tangent lines through $R .^{3}$

All these conclusions remain valid in the case, when $\mathcal{R}_{t}$ consists of conics which touch $c$ at $B_{2}$ and are tangent to $t_{3}$ and $t_{4}$.

As already indicated by the notation, we are interested in the particular case of Theorem 3.5 where the conics $c$ and $e$ in the range $\mathcal{R}_{c}$ are confocal. The following result follows directly from Theorem 3.5 and summarizes properties of the Poncelet grid. For the points of intersection between extended sides of a billiard $\ldots P_{0} P_{1} P_{2} \ldots$ we use the notation

$$
S_{i}^{(j)}:= \begin{cases}{\left[P_{i-k-1}, P_{i-k}\right] \cap\left[P_{i+k}, P_{i+k+1}\right] \text { for }} & j=2 k,  \tag{3.3}\\ {\left[P_{i-k}, P_{i-k+1}\right] \cap\left[P_{i+k}, P_{i+k+1}\right] \text { for }} & j=2 k-1\end{cases}
$$

where $i=\ldots, 0,1,2, \ldots$ and $j=1,2, \ldots$ Note that there are $j$ sides between those which intersect at $S_{i}^{(j)}$, and 'in the middle' of these $j$ sides there is for

[^2]even $j$ the vertex $P_{i}$ and otherwise the point of contact $Q_{i}$. At the same token, the point $S_{i}^{(j)}$ is the pole of the diagonal $\left[Q_{i-k-1}, Q_{i+k}\right]$ or $\left[Q_{i-k}, Q_{i+k}\right]$ of the polygon $\ldots Q_{1} Q_{2} Q_{3} \ldots$ of contact points w.r.t. the caustic.

Theorem 3.6. Let ... $P_{0} P_{1} P_{2} \ldots$ be a billiard in the ellipse e with sides $P_{i} P_{i+1}$ contacting the ellipse $c$ at the respective points $Q_{i}$ for all $i \in \mathbb{Z}$. Then the vertices $S_{i}^{(j)}$ of the associated Poncelet grid are distributed on the following conics.

1. The points $S_{i}^{(1)}, S_{i}^{(3)}, \ldots$ are located on the confocal hyperbola through $Q_{i}$, while the points $S_{i}^{(2)}, S_{i}^{(4)}, \ldots$ are located on the confocal hyperbola through $P_{i}$.
2. For each $j \in\{1,2, \ldots\}$, the points $\ldots S_{i}^{(j)} S_{i+(j+1)}^{(j)} S_{i+2(j+1)}^{(j)} \ldots$ are vertices of another billiard with the caustic c inscribed in a confocal ellipse $e^{(j)}$, provided that $e^{(j)}$ is regular. Otherwise $e^{(j)}$ coincides with an axis of symmetry or with the line at infinity. The locus $e^{(j)}$ is independent of the position of the initial vertex $P_{0} \in e$.

Proof. 1. The side lines $\left[P_{0}, P_{1}\right]\left(P_{0}=P_{7}\right.$ in Fig. 5) and $\left[P_{2}, P_{3}\right]$ meet at $S_{1}^{(1)}$, while $\left[P_{1}, P_{2}\right]$ contacts $c$ at $Q_{1}$. By Theorem $3.5,2$. the points $Q_{1}$ and $S_{1}^{(1)}$ belong to the same confocal hyperbola.
Now we go one step away from $Q_{1}$ : the tangents from $P_{0}$ and $P_{3}$ to $c$ intersect at $S_{1}^{(1)}$ and $S_{1}^{(3)}=\left[P_{-1}, P_{0}\right] \cap\left[P_{3}, P_{4}\right]$. The confocal conic through $S_{1}^{(1)}$ and $S_{1}^{(3)}$ must again be the hyperbola through $Q_{1}$. This follows by continuity after choosing $Q_{1}$ on one axis of symmetry. Iteration confirms the first claim in Theorem 3.6.
The tangents to $c$ from $P_{0}$ and $P_{2}$ form a quadrilateral with $P_{1}$ and $S_{1}^{(2)}$ as opposite vertices. Therefore, there exists a confocal conic passing through both points. This conic must be a hyperbola, as can be concluded by continuity: If $P_{1}$ is specified at a vertex of $e$, then due to symmetry the points $P_{1}$ and $S_{1}^{(2)}$ are located on an axis of symmetry.
The tangents to $c$ from $P_{-1}$ and $P_{3}$ form a quadrilateral with $S_{1}^{(2)}$ and $S_{1}^{(4)}$ as opposite vertices. Theorem 3.5 and continuity guarantee that this is again the confocal hyperbola through $P_{1}$. Iteration shows the same of $S_{1}^{(6)}$ etc. However, the points $P_{1}, S_{1}^{(2)}, S_{1}^{(4)}, \ldots$ need not belong to the same branch of the hyperbola.
2. The tangents through $P_{2}$ and $S_{2}^{(2)}$ (note Fig. 5) form a quadrilateral with $S_{1}^{(1)}$ and $S_{2}^{(1)}$ as opposite vertices. This time, continuity shows that the two points belong to the same confocal ellipse $e^{(1)}$. The same holds for the tangents through $P_{3}$ and $S_{3}^{(2)}$ etc.
Similarily, starting with the points $P_{0}$ and $P_{3}$, we find the ellipse $e^{(2)}$ through $S_{1}^{(2)}$ and $S_{2}^{(2)}$, and so on.
In order to prove that these ellipses $e^{(1)}, e^{(2)}, \ldots$ are independent of the choice of the initial point $P_{1} \in e$, we follow an argument from [2, proof of


Figure 5 Periodic billiard $(N=7)$ with extended sides

Corollary 2.2]: The claim holds for all confocal ellipses $e$ where billiards with the same caustic $c$ are aperiodic and traverse $e$ infinitely often. Since these ellipses form a dense set, the claim holds also for those with periodic billiards. The invariance of the ellipses $e^{(1)}, e^{(2)}, \ldots$ is already mentioned in [2, Theorem 7].
An alternative proof consists in demonstrating that the elliptic coordinate $k_{e}^{(j)}$ of $e^{(j)}$ for all $j$ does not depend on the parameter $t$. As one example, we present below in (3.7), (3.6) and (3.8) formulas for the semiaxes $a_{e \mid 1}$, $b_{e \mid 1}$ and the elliptic coordinate $k_{e \mid 1}$ of $e^{(1)}$.

Remark 3.7. Figure 5 reveals, that the polygons $P_{1} S_{1}^{(1)} P_{2} S_{2}^{(1)} \ldots$ as well as $P_{1} S_{2}^{(2)} P_{3} S_{4}^{(2)} \ldots$ and $S_{1}^{(2)} S_{1}^{(1)} S_{2}^{(2)} S_{2}^{(1)} \ldots$ are zigzag billiards in rings bounded by two confocal ellipses. However, we find also zigzag billiards between two confocal hyperbolas, e.g., $\ldots S_{1}^{(2)} P_{2} P_{1} S_{2}^{(2)} \ldots$ or the twofold covered $\ldots S_{1}^{(2)}$ $S_{1}^{(1)} P_{1} Q_{1} P_{1} S_{1}^{(1)} S_{1}^{(2)} \ldots$. Billiards between other pairs of confocal conics can be found in [12].

The coming lemma addresses invariants related to the incircles of quadrilaterals built from the tangents to $c$ from any two vertices $P_{i}$ and $P_{j}$ of a billiard in $e$ (note circle $d$ in Fig. 4 and $[1,5,19]$ ).

Lemma 3.8. Referring to Fig. 6, the power $w^{2}$ of the point $S_{i}^{(1)}$ w.r.t. the incircle of the triangle $P_{i} P_{i+1} S_{i}^{(1)}$ is the same for all $i$. Similarly, the power $w_{1}^{2}$ of $S_{i}^{(2)}$ w.r.t. the incircle of the quadrangle $P_{i} S_{i-1}^{(1)} S_{i}^{(2)} S_{i}^{(1)}$ is constant.

Proof. According to Graves's construction [15, p. 47], an ellipse $e$ can be constructed from a smaller ellipse $c$ in the following way: Let a closed piece of


Figure 6 The power $w^{2}$ of $S_{2}^{(1)}$ w.r.t. the incircle of the triangle $S_{2}^{(1)} P_{2} P_{3}$ shows up at all $S_{i}^{(1)}$ and equals the power of $P_{i}$ w.r.t. the incircle of the quadrangle $P_{i} S_{i-1}^{(1)} S_{i}^{(2)} S_{i}^{(1)}$
string strictly longer than the perimeter of $c$ be posed around $c$. If point $P$ is used to pull the string taut, then $P$ traces a confocal ellipse $e$. Consequently, for each vertex $P_{i}$ and neighboring tangency points $Q_{i-1}$ and $Q_{i}$ of a billiard in $e$ with caustic $c$, the sum of the lengths $\overline{Q_{i-1} P_{i}}$ and $\overline{P_{i} Q_{i}}$ minus the length of the elliptic arc between $Q_{i-1}$ and $Q_{i}$, i.e.,

$$
\begin{equation*}
D_{e}:=\overline{Q_{i-1} P_{i}}+\overline{P_{i} Q_{i}}-\stackrel{\frown}{Q_{i-1}} Q_{i} \tag{3.4}
\end{equation*}
$$

is constant (Fig. 6).
The incircle of $P_{2} P_{3} S_{2}^{(1)}$ has the center $R_{2}$ and the radius $\overline{Q_{2} R_{2}}$ by (2.15). The power of $P_{2}$ w.r.t. this circle is $l_{2}^{2}$, that of $P_{3}$ is $r_{3}^{2}$. From Graves' construction follows for the ellipse $e$ that

$$
D_{e}=r_{2}+l_{2}-\overparen{Q_{1} Q_{2}}=r_{3}+l_{3}-\overparen{Q_{2} Q_{3}}=\text { const }
$$

is the same for all $P_{i}$, provided that $\overparen{Q i}_{i}$ denotes the length of the (shorter) arc along $c$ between $Q_{i}$ and $Q_{j}$. For the analogue invariant at $e^{(1)}$ follows (Fig. 6)

$$
\begin{align*}
D_{e \mid 1} & :=\overline{Q_{1} S_{2}^{(1)}}+\overline{S_{2}^{(1)} Q_{3}}-\overparen{Q_{1}}{ }_{3} \\
& =\left(r_{2}+l_{2}+w\right)+\left(r_{3}+l_{3}+w\right)-\overparen{Q_{1}}{ }_{3}=2 D_{e}+2 w=\text { const. } \tag{3.5}
\end{align*}
$$

hence $w=$ const., where $w^{2}$ is the power of $S_{2}^{(1)}$ w.r.t. the said incircle.
Since the incircle of the quadrangle $P_{2} S_{1}^{(1)} S_{2}^{(2)} S_{2}^{(1)}$ is an excircle of the triangle $P_{2} P_{3} S_{2}^{(1)}$ (Fig. 6), the power of $P_{2}$ w.r.t. the excircle equals $w^{2}$, too. This follows from elementary geometry.
As an alternative, the constancy of $w$ can also be concluded from the fact, that neighboring circles with centers $R_{i}^{1}, R_{i}$ or $R_{i}, R_{i+1}^{1}$ share three tangents, and one circle is an incircle, the other an excircle of the triangle. Therefore, on the common side the same length $w$ shows up twice and also at the adjacent pairs of neighboring circles. Since the distance $w$ is constant for all aperiodic billiards, it reveals also for periodic billiards that $w$ is independent of the choice of the initial vertex. In a similar way follows the invariance of the length $w_{1}$, as shown in Fig. 6.

It needs to be noted that $S_{2}^{(1)}$ or $S_{2}^{(2)}$ can be located on the other branch of the related hyperbola. Then the said 'incircle' of the triangle $P_{2} P_{3} S_{2}^{(1)}$ has to be replaced by the excircle which contacts $c$ at $Q_{2}$ and is tangent to the side lines [ $P_{1}, P_{2}$ ] and $\left[P_{3}, P_{4}\right]$. Similarly, the said 'incircle' of the quadrangle becomes an excircle. In all these cases, Lemma 3.8 and the proof given above have to be adapted.

Remark 3.9. The Theorems 3.5 and 3.6 as well as the constancy of the length $w$ according to Lemma 3.8 are also valid in spherical geometry (note [26]) and in hyperbolic geometry. On the sphere (see Fig. 7), the caustic consists of a pair of opposite components, and for $N$-periodic billiards the confocal spherical ellipse $e^{(j)}$ coincides with $e^{(N-2-j)}$ w.r.t. the opposite caustic.
We obtain other families of incircles when we focus on pairs of consecutive sides of the billiards in $e^{(1)}, e^{(2)}$ and so on. However, these circles are not mutually disjoint. By the way, the centers $R_{i}^{(j)}$ of all these circles are the poles of diagonals of $\ldots P_{1} P_{2} P_{3} \ldots$ w.r.t. the ellipse $e$.

For the sake of completeness, we express below in (3.9) the distance $w$ in terms of the semiaxes of $e$ and $c$. For this purpose, we compute first the semiaxes $a_{e \mid 1}$ an $b_{e \mid 1}$ of $e^{(1)}$, since we need the coordinates of $S_{2}^{(1)}$. From (2.13) and (2.15) follows

$$
R_{2}=\left(\frac{a_{e}^{2} \cos t_{2}^{\prime}}{a_{c}}, \frac{b_{e}^{2} \sin t_{2}^{\prime}}{b_{c}}\right), \quad \overline{Q_{2} R_{2}}=\frac{k_{e}\left\|\mathbf{t}_{c}\left(t_{2}^{\prime}\right)\right\|}{a_{c} b_{c}}
$$

By virtue of Theorem 3.5, the tangents from $R_{2}$ to the confocal hyperbola through $Q_{2}$ contact at

$$
Q_{2}=\left(a_{c} \cos t_{2}^{\prime}, b_{c} \sin t_{2}^{\prime}\right) \in c \quad \text { and } \quad S_{2}^{(1)}=\left(a_{e \mid 1} \cos t_{2}^{\prime}, b_{e \mid 1} \sin t_{2}^{\prime}\right) \in e^{(1)}
$$

Both points lie on the polar of $R_{2}$ w.r.t. the hyperbola in question with the elliptic coordinate $k_{h}=-\left\|\mathbf{t}_{e}\left(t_{2}^{\prime}\right)\right\|^{2}$. This yields the condition

$$
\frac{a_{e}^{2} \cos t_{2}^{\prime}}{a_{c}\left(a_{c}^{2}+k_{h}\left(t_{2}^{\prime}\right)\right)} a_{e \mid 1} \cos t_{2}^{\prime}+\frac{b_{e}^{2} \sin t_{2}^{\prime}}{b_{c}\left(b_{c}^{2}+k_{h}\left(t_{2}^{\prime}\right)\right)} b_{e \mid 1} \sin t_{2}^{\prime}=1
$$



Figure 7 Periodic billiard $P_{1} P_{2} \ldots P_{7}$ on the sphere with extended sides and their contact points with the incircles. All circular arcs marked in black have the same length, as well as those marked in green, and both lengths are invariant against changes of $P_{1}$ on $e$
or, by virtue of (2.7),

$$
a_{e}^{2} b_{c} a_{e \mid 1}-b_{e}^{2} a_{c} b_{e \mid 1}=a_{c} b_{c} d^{2}, \quad \text { where } \quad a_{e \mid 1}^{2}-b_{e \mid 1}^{2}=d^{2} .
$$

We eliminate $a_{e \mid 1}$ and obtain after some computation the quadratic equation

$$
\left(b_{e}^{4}-2 b_{c}^{2} b_{e}^{2}-b_{c}^{2} d^{2}\right) b_{e \mid 1}^{2}+2 a_{c}^{2} b_{c} b_{e}^{2} b_{e \mid 1}+b_{c}^{2}\left(a_{c}^{2} d^{2}-a_{e}^{4}\right)=0
$$

The second solution besides $b_{e \mid 1}=b_{c}$ is

$$
\begin{equation*}
b_{e \mid 1}=\frac{b_{c}\left(a_{c}^{2} d^{2}-a_{e}^{4}\right)}{b_{e}^{4}-2 b_{c}^{2} b_{e}^{2}-b_{c}^{2} d^{2}}=\frac{b_{c}\left(a_{e}^{2} b_{e}^{2}+d^{2} k_{e}\right)}{a_{c}^{2} b_{c}^{2}-k_{e}^{2}} . \tag{3.6}
\end{equation*}
$$

This implies

$$
\begin{equation*}
a_{e \mid 1}=\frac{a_{c}}{a_{e}^{2} b_{c}}\left(b_{c} d^{2}+b_{e}^{2} b_{e \mid 1}\right)=\frac{a_{c}\left(a_{e}^{2} b_{e}^{2}-d^{2} k_{e}\right)}{a_{c}^{2} b_{c}^{2}-k_{e}^{2}} \tag{3.7}
\end{equation*}
$$

and

$$
\begin{equation*}
k_{e \mid 1}=b_{e \mid 1}^{2}-b_{c}^{2}=k_{e}\left(\frac{2 a_{c} b_{c} a_{e} b_{e}}{a_{c}^{2} b_{c}^{2}-k_{e}^{2}}\right)^{2} \tag{3.8}
\end{equation*}
$$

and yields finally

$$
\begin{equation*}
w:=\frac{2 a_{e} b_{e} \sqrt{k_{e}^{3}}}{a_{c}^{2} b_{c}^{2}-k_{e}^{2}} \tag{3.9}
\end{equation*}
$$

Negative semiaxes $a_{e \mid 1}, b_{e \mid 1}$ and a negative $w$ in the formulas above mean that the points $S_{i}^{(1)}$ are located on the respectively second branches of the hyperbolas and the incircles of the triangles $P_{i} P_{i+1} S_{i}^{(1)}$ become excircles. In the case of a vanishing denominator for $k_{e}=a_{c} b_{c}$ (periodic four-sided billiard) the ellipse $e^{(1)}$ is the line at infinity.
If on the right-hand side of the formulas (3.7), (3.6) and (3.8) we replace $a_{e}, b_{e}, k_{e}$ respectively by $a_{e \mid 1}, b_{e \mid 1}, k_{e \mid 1}$, then we obtain expressions for $a_{e \mid 3}, b_{e \mid 3}$, $k_{e \mid 3}$, i.e.,

$$
\begin{equation*}
k_{e \mid 3}=k_{e \mid 1}\left(\frac{2 a_{c} b_{c} a_{e \mid 1} b_{e \mid 1}}{a_{c}^{2} b_{c}^{2}-k_{e \mid 1}^{2}}\right)^{2} . \tag{3.10}
\end{equation*}
$$

### 3.2. Conjugate billiards

For two confocal ellipses $c$ and $e$, there exists an axial scaling

$$
\begin{equation*}
\alpha:(x, y) \mapsto\left(\frac{a_{e}}{a_{c}} x, \frac{b_{e}}{b_{c}} y\right) \quad \text { with } \quad c \rightarrow e . \tag{3.11}
\end{equation*}
$$

Corresponding points share the parameter $t$. Hence, they belong to the same confocal hyperbola (Fig. 8). The affine transformation $\alpha$ maps the tangency point $Q_{i} \in c$ of the side $P_{i} P_{i+1}$ to a point $P_{i}^{\prime} \in e$, while $\alpha^{-1}$ maps $P_{i}$ to the tangency point $Q_{i-1}^{\prime}$ of $P_{i-1}^{\prime} P_{i}^{\prime}$, i.e.,

$$
\alpha: Q_{i} \mapsto P_{i}^{\prime}, \quad Q_{i-1}^{\prime} \mapsto P_{i}
$$

This results from the symmetry between $t_{i}$ and $t_{i}^{\prime}$ in the equation

$$
\begin{equation*}
b_{c} a_{e} \cos t_{i} \cos t_{i}^{\prime}+a_{c} b_{e} \sin t_{i} \sin t_{i}^{\prime}=a_{c} b_{c} \tag{3.12}
\end{equation*}
$$

which expresses that $P_{i} \in e$ with parameter $t_{i}$ lies on the tangent to $c$ at $Q_{i}$ with parameter $t_{i}^{\prime}$. Referring to Fig. 8, $\alpha$ sends the tangent $\left[P_{i-1}^{\prime}, P_{i}^{\prime}\right]$ to $c$ at $Q_{i-1}^{\prime}$ to the tangent $\left[R_{i-1}, R_{i}\right]$ to $e$ at $P_{i}$. Hence, by $\alpha$ the polygon $Q_{1} Q_{2} \ldots$ is mapped to $P_{1}^{\prime} P_{2}^{\prime} \ldots$ and futhermore to that of the poles $R_{1} R_{2} \ldots$ of the billiard's sides $P_{1} P_{2}, P_{2} P_{3}, \ldots$.

Definition 3.10. Referring to Fig. 8, the billiard $\ldots P_{0}^{\prime} P_{1}^{\prime} P_{2}^{\prime} \ldots$ is called conjugate to the billiard $\ldots P_{0} P_{1} P_{2} \ldots$ in the ellipse $e$ with the ellipse $c$ as caustic, when the axial scaling $\alpha: c \rightarrow e$ defined in (3.11) maps the tangency point $Q_{i}$ of the side $P_{i} P_{i+1}$ to the vertex $P_{i}^{\prime}$.

Lemma 3.11. For each billiard $\ldots P_{0} P_{1} P_{2} \ldots$ in the ellipse $e$ with the ellipse $c$ as caustic, there exists a unique conjugate billiard $\ldots P_{0}^{\prime} P_{1}^{\prime} P_{2}^{\prime} \ldots$, and the relation between the two billiards in $e$ is symmetric. Moreover,

$$
\begin{equation*}
l_{i}=\overline{P_{i} Q_{i}}=\overline{P_{i}^{\prime} Q_{i-1}^{\prime}}=r_{i}^{\prime} \quad \text { and } \quad r_{i}=\overline{P_{i} Q_{i-1}}=\overline{P_{i-1}^{\prime} Q_{i-1}^{\prime}}=l_{i-1}^{\prime} \tag{3.13}
\end{equation*}
$$



Figure 8 The periodic billiard $P_{1} P_{2} \ldots P_{5}$ in $e$ with the caustic $c$ and the conjugate billiard $P_{1}^{\prime} P_{2}^{\prime} \ldots P_{5}^{\prime}$

Proof. From the symmetry in (3.12) follows for $\alpha: c \rightarrow e$ that $P_{i}$ is the preimage of $P_{i}$ is the tangency point $Q_{i-1}$ of $P_{i-1}^{\prime} P_{i}^{\prime}$. The congruences stated in (3.13) follow from Ivory's Theorem for the two diagonals in the curvilinear quadrangle $P_{i} P_{i}^{\prime} Q_{i} Q_{i-1}^{\prime}$. In view of the sequence of parameters $t_{1}, t_{1}^{\prime}, t_{2}, t_{2}^{\prime}, t_{3}, \ldots$ of the vertices $P_{1}, P_{1}^{\prime}, P_{2}, P_{2}^{\prime}, P_{3}, \ldots$ on $e$, the switch between the original billiard and its conjugate corresponds to the interchange of $t_{i}$ with $t_{i}^{\prime}$ for $i=1,2, \ldots$.

Finally we recall that, based on the Arnold-Liouville theorem from the theory of completely integrable systems, it is proved in [19,21] that there exist canonical coordinates $u$ on the ellipses $e$ and $c$ such that for any billiard the transitions from $P_{i} \rightarrow P_{i+1}$ and $Q_{i} \rightarrow Q_{i+1}$ correspond to shifts of the respective canonical coordinates $u_{i}$ and $u_{i+1}$ by $2 \Delta u$. Explicit formulas for the parameter transformation $t \mapsto u$ are provided in [27].

Figure 9 shows how on $c$ such coordinates can be constructed by iterated subdivision, provided that $Q_{1}$ and $Q_{3}$ get the respective canonical coordinates $u=0$ and 1 . A comparison with Fig. 8 reveals that, in the sense of a canonical parametrization, the contact point $Q_{i}$ is exactly halfway from $P_{i}$ to $P_{i+1}$, i.e.,

$$
\begin{equation*}
u_{i}^{\prime}=u_{i}+\Delta u, \quad u_{i+1}=u_{i}+2 \Delta u . \tag{3.14}
\end{equation*}
$$

Hence, the transition from a billiard to its conjugate is equivalent to a shift of canonical coordinates by $\Delta u$. An equivalent result can be found in [21, Sect. 4].

### 3.3. Billiards with a hyperbola as caustic

As illustrated in Fig. 10, billiards in ellipses $e$ with a confocal hyperbola $c$ as caustic are zig-zags between an upper and lower subarc of $e$. If the initial point $P_{1}$ is chosen at any point of intersection between $e$ and the hyperbola $c$, then the billiard is twofold covered, and the first side $P_{1} P_{2}$ is tangent to $c$ at $P_{1}$.


Figure 9 An example of canonical coordinates on $c, e$ and $e^{(1)}$, this time with origin $Q_{1}$ and unit point $Q_{3}$


Figure 10 Periodic billiard $P_{1} P_{2} \ldots P_{12}$ in the ellipse $e$ with the hyperbola $c$ as caustic, together with the hyperbolas $e^{(1)}$, $e^{(3)}$ and the ellipse $e^{(2)}$ with the inscribed billiard consisting of three quadrangles $S_{i}^{(2)} S_{i+3}^{(2)} S_{i+6}^{(2)} S_{i+9}^{(2)}$

Here we report briefly, in which way these billiard differ from those with an elliptic caustic. Proofs are left to the readers. In view of the associated Poncelet grid, we start with the analogue to Theorem 3.6 (see Figs. 10, 11 and 12).


Figure 11 The power $w^{2}$ of $P_{9}$ w.r.t. the circle tangent to the sides $\left[P_{7}, P_{8}\right],\left[P_{8}, P_{9}\right],\left[P_{9}, P_{10}\right]$ equals that of $P_{12}$ w.r.t. the circle tangent to [ $\left.P_{10}, P_{11}\right],\left[P_{11}, P_{12}\right],\left[P_{12}, P_{1}\right]$


Figure 12 Periodic billiard $P_{1} P \ldots P_{6}$ in the ellipse $e$ with the hyperbola $c$ as caustic, together with the conjugate billiard $P_{1}^{\prime} P_{2}^{\prime} \ldots P_{6}^{\prime}$. The associated polygon with vertices on $e^{(1)}$ splits into two triangles $S_{1}^{(1)} S_{3}^{(1)} S_{5}^{(1)}$ (green shaded) and $S_{2}^{(1)} S_{4}^{(1)} S_{6}^{(1)}$

Theorem 3.12. Let $\ldots P_{0} P_{1} P_{2} \ldots$ be a billiard in the ellipse $e$ with the hyperbola c as caustic.

1. Then the points $S_{i}^{(1)}, S_{i}^{(3)}, \ldots$ are located on confocal ellipses through the contact point $Q_{i}$ of $\left[P_{i}, P_{i+1}\right]$ with $c$, while the points $S_{i}^{(2)}, S_{i}^{(4)}, \ldots$ are located on the confocal hyperbola through $P_{i}$.
2. For even $j$, the points $\ldots S_{i}^{(j)} S_{i+(j+1)}^{(j)} S_{i+2(j+1)}^{(j)} \ldots$ are vertices of another billiard with the caustic $c$ inscribed in a confocal ellipse $e^{(j)}$, provided that $S_{i}^{(j)}$ is finite.
For odd $j$, the points $S_{i}^{(j)}$ are located on confocal hyperbolas $e^{(j)}$ or an axis of symmetry. At each vertex of $\ldots S_{i}^{(j)} S_{i+(j+1)}^{(j)} S_{i+2(j+1)}^{(j)} \ldots$, one angle bisector is tangent to $e^{(j)}$.
All conics $e^{(j)}$ are independent of the position of the initial vertex $P_{1} \in e$.
At the 12-periodic billiard depicted in Fig. 10, the billiard inscribed to $e^{(2)}$ splits into three quadrangles (dashed). Note that for odd $j$ there are some points $S_{i}^{(j)}$ where the tangent to $e^{(j)}$ is the interior bisector of the angle $\angle S_{i-(j+1)}^{(j)} S_{i}^{(j)} S_{i+(j+1)}^{(j)}$. Hence, we obtain no billiards inscribed to hyperbolas $e^{(j)}$ with the hyperbola $c$ as caustic. In Fig. 12, the 6 -periodic billiard $P_{1} P_{2} \ldots P_{6}$ yields two triangles $S_{i}^{(1)} S_{i+2}^{(1)} S_{i+4}^{(1)}$ inscribed to $e^{(1)}$; one of them is shaded in green.
Lemma 3.8 is also valid for hyperbolas as caustic. An example is depicted in Fig. 11: The power $w^{2}$ of $P_{9}$ w.r.t. the circle tangent to the four consecutive sides $\left[P_{7}, P_{8}\right],\left[P_{8}, P_{9}\right],\left[P_{9}, P_{10}\right]$, and $\left[P_{10}, P_{11}\right]$ equals that of $P_{12}$ w.r.t. the incircle of the quadrilateral with sides $\left[P_{10}, P_{11}\right],\left[P_{11}, P_{12}\right],\left[P_{12}, P_{1}\right]$, and $\left[P_{1}, P_{2}\right.$ ].
Also for billiards $P_{1} P_{2} \ldots$ in $e$ with a hyperbola $c$ as caustic, there exists a conjugate billiard $P_{1}^{\prime} P_{2}^{\prime} \ldots$, and the relation is symmetric. However, the definition is different. It uses the singular affine transformation

$$
\begin{equation*}
\alpha_{h}: e \rightarrow F_{1} F_{2} \text { with } P_{i} \mapsto T_{i}=\left[P_{i}, P_{i+1}\right] \cap\left[F_{1}, F_{2}\right] \tag{3.15}
\end{equation*}
$$

with $F_{1}$ and $F_{2}$ as the focal points of $e$ and $c$ (Fig. 12).
Definition 3.13. Referring to Fig. 12, the billiard $\ldots P_{0}^{\prime} P_{1}^{\prime} P_{2}^{\prime} \ldots$ in the ellipse $e$ with the hyperbola $c$ as caustic is called conjugate to the billiard $\ldots P_{0} P_{1} P_{2} \ldots$ in $e$ with the same caustic $c$ if the axial scaling $\alpha_{h}$ defined in (3.15) maps the point $P_{i}^{\prime}$ to the intersection $T_{i}$ of $P_{i} P_{i+1}$ with the principal axis.

Lemma 3.14. Let $\ldots P_{0} P_{1} P_{2} \ldots$ be a billiard in the ellipse $e$ with the hyperbola $c$ as caustic. Then to this billiard and to its mirror w.r.t. the principal axis exists a conjugate billiard $\ldots P_{0}^{\prime} P_{1}^{\prime} P_{2}^{\prime} \ldots$, and it is unique up to a reflection in the principal axis. The relation between two conjugate billiards in $e$ is symmetric. Moreover, if $T_{i}^{\prime}$ denotes the intersection of $P_{i}^{\prime} P_{i+1}^{\prime}$ with the principal axis,
then

$$
\begin{equation*}
\overline{P_{i} T_{i}}=\overline{P_{i}^{\prime} T_{i-1}^{\prime}} \text { and } \overline{P_{i} T_{i-1}}=\overline{P_{i-1}^{\prime} T_{i-1}^{\prime}} . \tag{3.16}
\end{equation*}
$$

Proof. The singular affine transformation $\alpha_{h}$ maps $P_{1}^{\prime}$ to $T_{1}$ and $P_{1}$ to a point $T^{\prime}\left(=T_{6}^{\prime}\right.$ in Fig. 12). We assume that $P_{1}$ and $P_{1}^{\prime}$ lie on the same side of the principal axis, since otherwise we apply a reflection in the axis. Then we obtain a curvilinear Ivory quadrangle $P_{1}^{\prime} T_{1} T^{\prime} P_{1}$ with diagonals of equal lengths. On the other hand, the lines $\left[P_{1}^{\prime}, T^{\prime}\right]$ and $\left[P_{1}, T_{1}\right]$ must contact the same confocal conic (see, e.g., [7, p. 153] or [26, Lemma 1]). Hence, the billiard through the point $P_{1}^{\prime} \in e$ with caustic $c$ contains one side on the line $\left[P_{1}^{\prime}, T^{\prime}\right]$. Iteration confirms the claim.

Let $\bar{P}_{1}$ and $\bar{P}_{2}$ be the images of $P_{1}$ and $P_{2}$ under reflection in the principal axis of $e$ (Fig. 12). Then, a comparison with Fig. 9 reveals that the confocal hyperbola through the intersection $T_{1}=\left[P_{1}, P_{2}\right] \cap\left[\bar{P}_{1}, \bar{P}_{1}\right]$ lies 'in the middle' between the hyperbolas through $P_{1}$ and $P_{2}$.

Remark 3.15. If $\ldots P_{0} P_{1} P_{2} \ldots$ and $\ldots P_{0}^{\prime} P_{1}^{\prime} P_{2}^{\prime} \ldots$ is a pair of conjugate billiards in an ellipse $e$, then the points of intersection $\left[P_{i}, P_{i+1}\right] \cap\left[P_{i}^{\prime}, P_{i-1}^{\prime}\right]$ are located on a confocal ellipse $e^{\prime}$ inside $e$. This holds for ellipses and hyperbolas as caustics. If the billiards are $N$-periodic, then in the elliptical case, the restriction of the two billiards to the interior of $e^{\prime}$ is $2 N$-periodic; conversely, $e$ plays the role of $e^{(1)}$ w.r.t. $e^{\prime}$ (Fig. 8). In the hyperbolic case, the restriction to the interior of $e^{\prime}$ gives two symmetric $2 N$-periodic billiards, provided that also the reflected billiards are involved (Fig. 12).

We conclude with citing a result from [28] about billiards in ellipses. It states that for each billiard $\ldots P_{1} P_{2} P_{3} \ldots$ in $e$ with a hyperbola as caustic there exists a billiard $\ldots P_{1}^{*} P_{2}^{*} P_{3}^{*} \ldots$ in $e^{*}$ with an ellipse as caustic such that corresponding sides $P_{i} P_{i+1}$ and $P_{i}^{*} P_{i+1}^{*}$ are congruent.

## 4. Periodic $N$-sided billiards

Let the billiard $P_{1} P_{2} \ldots P_{N}$ in the ellipse $e$ be periodic with an ellipse $c$ as caustic. Then, the sequence of parameters $t_{1}, t_{1}^{\prime}, t_{2}, \ldots, t_{N}, t_{N}^{\prime}$ of the vertices $P_{i}$ and the intermediate contact points $Q_{1}, \ldots, Q_{N}$ with $c$ is cyclic. Each side line intersects only a finite number of other side lines. Hence, the corresponding Poncelet grid contains a finite number of confocal ellipses $e^{(j)}$ through the points $S_{i}^{(j)}$, namely $\left[\frac{N-2}{2}\right.$ ] (including possibly the line at infinity), provided that $N \geq 5$ (Fig. 5). The sequence of ellipses $e, e^{(1)}, e^{(2)}, \ldots$ is cyclic, and

$$
\begin{equation*}
e^{(j)}=e^{(N-2-j)} \tag{4.1}
\end{equation*}
$$

For example, in the case $N=7$ (Fig. 6), the ellipse $e^{(2)}$ coincides with $e^{(3)}$.

 $l_{2}=\overline{P_{2} Q_{2}}=\overline{Q_{6} P_{7}}=r_{7}$ and $\overline{S_{2}^{(1)} P_{3}}=\overline{P_{7} S_{7}^{(1)}}$. The associated billiard in $e^{(2)}$ splits into three triangles

Definition 4.1. The sum of the oriented exterior angles $\theta_{i}$ of a periodic billiard in an ellipse $e$ is an integer multiple of $2 \pi$, namely $2 \tau \pi$. We call $\tau \in \mathbb{N}$ the turning number of the billiard. It counts the loops of the billiard around the center $O$ of $e$, anti-clockwise or clockwise.

If the periodic billiard $P_{1} P_{2} \ldots P_{N}$ has the turning number $\tau=1$ (Fig. 13), then the billiard $S_{1}^{(1)} S_{3}^{(1)} S_{5}^{(1)} \ldots$ in $e^{(1)}$ has $\tau=2$, that of $S_{1}^{(2)} S_{4}^{(2)} \ldots$ in $e^{(2)}$ the turning number $\tau=3$, and so on. In cases with $g=\operatorname{gcd}(N, \tau)>1$ the corresponding billiard splits into $g \frac{N}{g}$-sided billiards, each with turning number $\tau / g$ (note [25, Theorem 1.1]).

### 4.1. Symmetries of periodic billiards

The following is a corollary to Theorem 3.6.
Corollary 4.2. Let $P_{1} P_{2} \ldots P_{N}$ be an $N$-sided periodic billiard in the ellipse $e$ with the ellipse $c$ as caustic.
(i) For even $N$ and odd $\tau$, the billiard is centrally symmetric.
(ii) For odd $N=2 n+1$ and odd $\tau$, the billiard is centrally symmetric to the conjugate billiard, where $P_{i}$ corresponds to $P_{i+n}^{\prime}$. ${ }^{4}$
(iii) If $N$ is odd and $\tau$ is even, then the conjugate billiard coincides with the original one, and $P_{i}=P_{i+n}^{\prime}$.

Proof. By virtue of Theorem 3.6, the lines $\left[P_{i-j-1}, P_{i-j}\right.$, ] and $\left[P_{i+j}, P_{i+j+1}\right]$ for $j=1,2, \ldots$ meet at the point $S_{i}(j)$ on the confocal hyperbola through $P_{i}$.

[^3]

Figure 14 Periodic billiard with $N=7$ and $\tau=2$
(i) This means for even $N=2 n$, odd $\tau$ and $j=n-1$, that also the opposite vertex $S_{i}^{(n)}=P_{i-n}=P_{i+n}$ belongs to this hyperbola. If $P_{i}$ is specified at a vertex on the minor axis of the ellipse $e$, then $P_{i+n}$ is the opposite vertex. Continuity implies that the two points belong to different branches of the hyperbola and are symmetric w.r.t. the center $O$ of $e$.
(ii) , (iii): If $N$ is odd, say $N=2 n+1$, then for $j=n-1$ the sides [ $P_{i-n+1}, P_{i-n}$ ] and $\left[P_{i+n-1}, P_{i+n}\right]$ intersect at a point on the hyperbola through $Q_{i+n}$ and $P_{i+n}^{\prime}$. For odd $\tau$ (Figs. 8 and 13), the same continuity argument as before proves that $P_{i}$ and $P_{i+n}^{\prime}$ are opposite w.r.t. $O$. If $\tau$ is even (Fig. 14), then the choice of $P_{i} \in e$ on an axis of symmetry shows the coincidence with $P_{i+n}^{\prime} \in e$, and this must be preserved, when $P_{i}$ varies continuously on $e$. In the case of even $\tau$ and $N$ the billiard splits.

The billiards with a hyperbola $c$ as caustic (see Figs. 10 and 12) oscillate between the upper and lower section of $e$. Therefore, only billiards with an even $N$ can be periodic. Also for billiards of this type, it possible to define a turning number $\tau$ which counts how often the points $P_{1}, \bar{P}_{2}, P_{3}, \ldots, P_{N}$ (Fig. 12) run to and fro along the upper component of $e .{ }^{5}$ The symmetry properties of these periodic $N$-sided billiards differ from those in Corollary 4.2. They follow from Theorem 3.6, since opposite vertices $P_{i}$ and $P_{i+N / 2}$ belong to the same confocal hyperbola.

Corollary 4.3. Let $P_{1} P_{2} \ldots P_{N}$ be an $N$-sided periodic billiard in the ellipse $e$ with the hyperbola $c$ as caustic.

[^4](i) For $N \equiv 0(\bmod 4)$, the billiard is symmetric w.r.t. the secondary axis of $e$ and $c$.
(ii) For $N \equiv 2(\bmod 4)$ and odd turning number $\tau$, the billiards are centrally symmetric. For even $\tau$, each billiard is symmetric w.r.t. the principal axis of $e$ and $c$.

### 4.2. Some invariants

As a direct consequence of the results so far, we present new proofs for the invariants k101, k118 and k119 listed in [24, Table 2], though this table refers already to proofs for some of them in $[2,11]$.

We begin with a result that has first been proved for a much more general setting in [29, p. 103].

Lemma 4.4. The length $L_{e}$ of a periodic $N$-sided billiard in the ellipse $e$ with the ellipse $c$ as caustic is independent of the position of the initial vertex $P_{1} \in e$.

Proof. We refer to Graves's construction [15, p. 47]. According to (3.4) holds

$$
D_{e}:=\overline{Q_{i-1} P_{i}}+\overline{P_{i} Q_{i}}-\underset{Q_{i-1} Q_{i}}{\frown}
$$

This yields for an $N$-sided billiard with turning number $\tau$ the total length

$$
\begin{equation*}
L_{e}=N \cdot D_{e}-\tau \cdot P_{c} \tag{4.2}
\end{equation*}
$$

where $P_{c}$ denotes the perimeter of the caustic $c$. Thus, $L_{e}$ does not depend of the choice of the initial vertex $P_{1} \in e$.

If the billiard in $e$ has the turning number $\tau$, then its extension in $e^{(1)}$ has the turning number $2 \tau$, and from (3.5) and (3.9) follows

$$
\begin{equation*}
L_{e \mid 1}=N D_{e \mid 1}-2 \tau P_{c}=2 N\left(D_{e}+w\right)-2 \tau P_{c}=2 L_{e}+2 N \frac{2 a_{e} b_{e} \sqrt{k_{e}^{3}}}{a_{c}^{2} b_{c}^{2}-k_{e}^{2}} \tag{4.3}
\end{equation*}
$$

The following theorem on the invariant k118 in [24] deals with the lengths $r_{i}$ and $l_{i}$ of the segments $Q_{i-1} P_{i}$ and $P_{i} Q_{i}$, as defined in (3.1).

Theorem 4.5. In each $N$-sided periodic billiard opposite segments are congruent, i.e., if $N=2 n$, then $r_{i+n}=r_{i}$ and $l_{i+n}=l_{i}$, and if $N=2 n+1$, then $r_{i+n}=l_{i-1}$ and $l_{i+n}=r_{i}$. Thus, for odd $N$ holds

$$
\sum_{i=1}^{N} l_{i}=\sum_{i=1}^{N} r_{i}=\frac{L_{e}}{2}
$$

Proof. By Corollary 4.2, for even $N=2 n$ the central symmetry implies for opposite segments $r_{i}=r_{i+n}$ and $l_{i}=l_{i+n}$.
If $N=2 n+1$, then $l_{i}=\overline{P_{i} Q_{i}}$ shows up as $l_{i+n}^{\prime}=\overline{P_{i+n}^{\prime} Q_{i+n}^{\prime}}$ at the conjugate billiard and, by virtue of (3.13), this equals $r_{i+n+1}=\overline{P_{i+n+1} Q_{i+n}}$ (Fig. 13). Similarly follows $r_{i}=l_{i+n}$.


Figure 15 Billiard $P_{1} P_{2} \ldots$ with pedal points w.r.t. $O$

Remark 4.6. In the particular case $N=3$ the two segments adjacent to any side are congruent (note in Fig. 13 the triangular billiards in $e^{(2)}$ ). Therefore, the Cevians $\left[P_{i}, Q_{i+1}\right]$ are concurrent and meet at the Nagel point of the triangle. This has already been proved in [23] and agrees with the circles through $Q_{i}$ and centered at $R_{i}$ (see Figs. 6 and 14), which for $N=3$ are excircles of the triangle $P_{1} P_{2} P_{3}$.

The following theorem has first been proved in [2, p. 4]. Another proof can be found in [3, Cor. 3.2]. We give below a new proof.

Theorem 4.7. For the exterior angles $\theta_{1}, \ldots, \theta_{N}$ of the periodic $N$-sided elliptic billiard in the ellipse $e$, the sum of cosines is independent of the initial vertex, namely

$$
\sum_{i=1}^{N} \cos \theta_{i}=N-J_{e} L_{e}=N-\frac{\sqrt{k_{e}}}{a_{e} b_{e}} L_{e},
$$

where $L_{e}$ is the common perimeter of these billiards in $e$.

Proof. The pedal points $F_{i}$ and $F_{i-1}$ on the sides $P_{i} P_{i+1}$ and $P_{i-1} P_{i}$ w.r.t. the center $O$ (Fig. 15) have the position vectors

$$
\begin{gathered}
\mathbf{f}_{i, i-1}=\mathbf{p}+\frac{\lambda_{i, i-1}}{\left\|\mathbf{t}_{e}\right\|}\left(\cos \frac{\theta_{i}}{2} \mathbf{t} \pm \sin \frac{\theta_{i}}{2} \mathbf{t}^{\perp}\right), \\
\text { where } \quad 0=\left\langle\mathbf{f}_{i, i-1},\left(\cos \frac{\theta_{i}}{2} \mathbf{t} \pm \sin \frac{\theta_{i}}{2} \mathbf{t}^{\perp}\right)\right\rangle .
\end{gathered}
$$

Here, $\lambda_{i}$ and $\lambda_{i+1}$ denote the signed distances from the vertex $P_{i}$ in one case towards $P_{i+1}$, in the other opposite to $P_{i-1}$. From

$$
\left\langle\mathbf{p}_{i}, \mathbf{t}\right\rangle=\left(-a_{e}^{2}+b_{e}^{2}\right) \cos t \sin t \text { and }\left\langle\mathbf{p}_{i}, \mathbf{t}^{\perp}\right\rangle=-a_{e} b_{e}
$$

follows

$$
\lambda_{i, i-1}=\frac{1}{\left\|\mathbf{t}_{e}\right\|}\left(\left(-a_{e}^{2}+b_{e}^{2}\right) \cos t \sin t \cos \frac{\theta_{i}}{2} \pm a_{e} b_{e} \sin \frac{\theta_{i}}{2}\right) .
$$

This implies by virtue of (2.11) and after reversing the orientation for $\lambda_{i-1}$,

$$
\lambda_{i}-\lambda_{i-1}=\overline{P_{i} F_{i}}+\overline{P_{i} F_{i-1}}=\frac{2 a_{e} b_{e}}{\left\|\mathbf{t}_{e}\right\|^{2}} \sqrt{k_{e}} .
$$

Since the sum over all signed lengths between $P_{i}$ and the adjacent pedal points gives the total perimeter $L_{e}$ of the billiard, we obtain by (2.12)

$$
\begin{equation*}
L_{e}=\sum_{i=1}^{N}\left(\overline{P_{i} F_{i}}+\overline{P_{i} F_{i-1}}\right)=\frac{a_{e} b_{e}}{\sqrt{k_{e}}} \sum_{i=1}^{N} \frac{2 k_{e}}{\left\|\mathbf{t}_{e}\right\|^{2}}=\frac{a_{e} b_{e}}{\sqrt{k_{e}}} \sum_{i=1}^{N}\left(1-\cos \theta_{i}\right) \tag{4.4}
\end{equation*}
$$

hence

$$
\begin{equation*}
\sum_{i=1}^{N} \frac{1}{\left\|\mathbf{t}_{e}\right\|^{2}}=\frac{L_{e}}{2 a_{e} b_{e} \sqrt{k_{e}}} \text { and also } \sum_{i=1}^{N} \cos \theta_{i}=N-\frac{\sqrt{k_{e}}}{a_{e} b_{e}} L_{e} \tag{4.5}
\end{equation*}
$$

as stated.
Remark 4.8. Note that the result in [2] relates to the interior angles of the billiard. As already mentioned in [2, Theorem 7], the constant sum of cosines holds also for the 'extended' billiards in $e^{(j)}$, where the exterior angles are $\theta_{i}+\theta_{i+1}+\cdots+\theta_{i+j}$ (note Fig. 5).

The first equation in (4.5) gives rise to two invariants which are already known: Similar to (2.14), the distance of $O$ to the tangent $t_{P}$ to $e$ at $P$ equals

$$
\begin{equation*}
\overline{O t_{P}}=\frac{a_{e} b_{e}}{\left\|\mathbf{t}_{e}\right\|} \tag{4.6}
\end{equation*}
$$

and yields a result as stated in [3, Cor. 3.2, third equation]. The invariant k119, first proved by P. Roitmann, deals with the curvature of $e$, namely by [15, p. 79] with

$$
\kappa_{e}(t):=\frac{a_{e} b_{e}}{\left\|\mathbf{t}_{e}(t)\right\|^{3}} .
$$

Corollary 4.9. The squared distances from the center $O$ to the tangents $t_{P_{i}}$ at the vertices $P_{i}$ of the periodic $N$-sided elliptic billiard in the ellipse e have a constant sum, independent of the initial vertex, namely

$$
\sum_{i=1}^{N}{\overline{O t_{P}}}^{2}=\frac{a_{e} b_{e}}{2 \sqrt{k_{e}}} L_{e}
$$

The curvatures $\kappa_{i}$ of $e$ at the vertices $P_{i}$ give rise to an invariant sum

$$
\sum_{i=1}^{N} \kappa_{i}^{2 / 3}=\frac{L_{e}}{2 \sqrt{k_{e}}}\left(a_{e} b_{e}\right)^{-1 / 3}
$$

Remark 4.10. It is remarkable that the quantity $\kappa^{2 / 3}$ appears in the billiard setting also at the Lazutkin parameter which coincides with the Poritsky string length, i.e., a kind of canonical parameter, up to additive and multiplicative constants. For details see [16].

We conclude this section with a comment on (4.1) in connection with (3.8) and (3.10). For example, we obtain $k_{e \mid 1}=k_{e}$ for $N=3$ and $k_{e}=a_{c} b_{c}$ for $N=4$. The condition $k_{e \mid 2}=k_{e \mid 1}$ is valid for $N=5$, and $k_{e \mid 2}=\infty$ for $N=6$. This
yields algebraic conditions for the semiaxes of the ellipse $e$ with an inscribed $N$-periodic billiard when the ellipse $c$ is given as caustic. However, we need to recall that approximately 200 years ago N. Fuß and J. Steiner presented already equations for the projectively equivalent case of circles (see [20, pp. 378-380]), and A. Cayley published an explicit solution for general $N$ in a projective setting ( [17] or [15, Theorem 9.5.4]). Other approaches are provided in [4, Sect. VI] and [14, Sect. 11.2.3.9]. Equivalent conditions in terms of elliptic functions can be deduced from [27, Corollary 3].

## Acknowledgements

The author is grateful to Dan Reznik and Ronaldo Garcia for inspirations and interesting discussions and to the anonymous reviewer for important advice.

Funding Open access funding provided by TU Wien (TUW).

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Received: May 14, 2021.
Revised: August 24, 2021.
Accepted: October 2, 2021.


[^0]:    ${ }^{1}$ The norm $\left\|\mathbf{t}_{e}\right\|$ equals half length of the diameter of $e$ which is parallel to $\mathbf{t}_{e}$.

[^1]:    ${ }^{2}$ They follow from the second equation in Lemma 3.1 for $t_{1}=t_{2}$.

[^2]:    ${ }^{3}$ An extended version of this theorem in [26] addresses the symmetry between the ranges $\mathcal{R}_{c}$ and $\mathcal{R}_{t}$. This generalizes the statement that in the case of confocal conics $c$ and $c_{1}$ the quadrilateral $A_{1} A_{2} B_{1} B_{2}$ has an incircle $d$ (Figs. 4, 6 and 14, compare with $[1,19]$ ).

[^3]:    ${ }^{4}$ All subscripts in this section are understood modulo $N$.

[^4]:    5 The turning number of hyperbolic billiards becomes more intuitive when the billiard is seen as the limit of a focal billiard in the sense of [28, Theorem 2].

