

Endre Csáki · Miklós Csörgő · Antónia Földes · Zhan Shi

Increment sizes of the principal value of Brownian local time

Received: 1 April 1999 / Revised version: 27 September 1999 /
Published online: 14 June 2000 – © Springer-Verlag 2000

Abstract. Let W be a standard Brownian motion, and define $Y(t) = \int_0^t ds/W(s)$ as Cauchy's principal value related to local time. We determine: (a) the modulus of continuity of Y in the sense of P. Lévy; (b) the large increments of Y .

1. Introduction

Let $\{W(t); t \geq 0\}$ be a one-dimensional Brownian motion with $W(0) = 0$, and let $\{L(t, x); t \geq 0, x \in \mathbb{R}\}$ denote its local time process. That is, for any Borel function $f \geq 0$,

$$\int_0^t f(W(s)) ds = \int_{-\infty}^{\infty} f(x)L(t, x) dx, \quad t \geq 0.$$

We are interested in the process

$$Y(t) = \int_0^t \frac{ds}{W(s)}, \quad t \geq 0. \quad (1.1)$$

Rigorously speaking, the integral $\int_0^t ds/W(s)$ should be considered in the sense of Cauchy's principal value, i.e., $Y(t)$ is defined by

$$Y(t) = \lim_{\varepsilon \rightarrow 0^+} \int_0^t \frac{ds}{W(s)} \mathbb{1}_{\{|W(s)| \geq \varepsilon\}} = \int_0^{\infty} \frac{L(t, x) - L(t, -x)}{x} dx. \quad (1.2)$$

E. Csáki: Alfréd Rényi Institute of Mathematics, Hungarian Academy of Sciences, P.O.B. 127, H-1364, Budapest, Hungary. e-mail: csaki@math-inst.hu
Research supported by the Hungarian National Foundation for Scientific Research, Grant No. T 019346 and T 029621.

M. Csörgő: School of Mathematics and Statistics, Carleton University, Ottawa, Ontario K1S 5B6, Canada. e-mail: mcsorgo@math.carleton.ca
Research supported by an NSERC Canada Grant at Carleton University, Ottawa.

A. Földes: College of Staten Island, CUNY, 2800 Victory Blvd., Staten Island, New York 10314, USA. e-mail: afoldes@email.gc.cuny.edu
Research supported by a PSC CUNY Grant No. 6-67383.

Z. Shi: Laboratoire de Probabilités UMR 7599, Université Paris VI, 4 Place Jussieu, F-75252 Paris Cedex 05, France. e-mail: zhan@proba.jussieu.fr

Mathematics Subject Classification (1991): 60J65

Key words: Local time – modulus of continuity – large increment – Brownian motion

Since $x \mapsto L(t, x)$ is Hölder continuous of order ν , for any $\nu < 1/2$, the integral on the extreme right in (1.2) is almost surely absolutely convergent.

The study of Cauchy’s principal value of Brownian local time goes back at least to Itô and McKean [12], and has become very active since the late 70s, due to applications in various branches of stochastic analysis. For example, it is a natural example in Fukushima [10]’s theory for Dirichlet processes and zero-energy additive functionals (these processes cannot be treated in the frame of the usual Itô calculus techniques). Another important fact is that principal values of local times can be represented as the Hilbert transform, or more generally, fractional derivatives, of local times. The latter plays an important role in a class of limit theorems for occupation times of Brownian motion, discovered by Papanicolaou, Stroock and Varadhan [14]. Also, the principal values of Brownian local times are the key ingredient in establishing Bertoin [1]’s excursion theory for Bessel processes of small dimensions. For a detailed account of various motivations, historical facts and general properties of principal values of local times, we refer to the recent collection of research papers in Yor [18], to Chapter 10 of the lecture notes by Yor [19], and to the survey paper by Yamada [17].

The process $Y(\cdot)$ defined in (1.1)–(1.2) is continuous, having zero quadratic variation. Although it is not used in this paper, we mention an interesting property: stopped at some suitably chosen random times, the principal values give all the possible symmetric stable processes (cf. Biane and Yor [3], Fitzsimmons and Gettoor [9], Bertoin [2]).

It is easily seen that $Y(\cdot)$ inherits a scaling property from Brownian motion, namely, for any fixed $a > 0$, $t \mapsto a^{-1/2}Y(at)$ has the same law as $t \mapsto Y(t)$. Although the aforementioned zero quadratic variation property distinguishes $Y(\cdot)$ from Brownian motion (in particular, $Y(\cdot)$ is not a semimartingale), it is a kind of folklore that Y behaves somewhat like a Brownian motion. Let us first recall (cf. [11]) the global and local almost sure asymptotics of $Y(\cdot)$:

$$\limsup_{t \rightarrow \infty} \frac{Y(t)}{\sqrt{t \log \log t}} = \sqrt{8}, \quad \text{a.s.} \tag{1.3}$$

$$\limsup_{t \rightarrow 0} \frac{Y(t)}{\sqrt{t \log \log(1/t)}} = \sqrt{8}, \quad \text{a.s.} \tag{1.4}$$

Comparing (1.3)–(1.4) with the corresponding laws of the iterated logarithm (LIL’s) for Brownian motion, we see that $\frac{1}{2}Y(t)$ and $W(t)$ satisfy exactly the same global and local LIL’s.

The aim of this paper is to get a uniform version of (1.3)–(1.4) for the increments of $Y(t)$. Our first result characterizes its modulus of continuity in the sense of P. Lévy.

Theorem 1.1. *With probability one,*

$$\lim_{h \rightarrow 0} \sup_{0 \leq t \leq 1} \sup_{0 \leq s \leq h} \frac{|Y(t+s) - Y(t)|}{\sqrt{h \log(1/h)}} = 2.$$

Remark 1.1.1. So $\frac{1}{\sqrt{2}}Y(t)$ and $W(t)$ have the same moduli of continuity (and the same remark applies to Theorem 1.2 below). We have already seen that $\frac{1}{2}Y(t)$ and $W(t)$ satisfy the same LIL's. Heuristically speaking, that a factor $\sqrt{2}$ is missing in the modulus of continuity comes from the fact that the Hausdorff dimension of the zero set of W is $\frac{1}{2}$.

Our second result concerns the large increments of $Y(\cdot)$. The length of time window, in which the increments are considered, denoted by a_T , will be supposed to satisfy the following condition:

$$\left\{ \begin{array}{l} 0 < a_T \leq T, \\ T \mapsto a_T \text{ and } T \mapsto T/a_T \text{ are both non-decreasing,} \\ \lim_{T \rightarrow \infty} \frac{\log(T/a_T)}{\log \log T} = \infty. \end{array} \right. \tag{1.5}$$

Here is our main result concerning the large increments of $Y(\cdot)$.

Theorem 1.2. *Under (1.5),*

$$\lim_{T \rightarrow \infty} \sup_{0 \leq t \leq T - a_T} \sup_{0 \leq s \leq a_T} \frac{|Y(t + s) - Y(t)|}{\sqrt{a_T \log(T/a_T)}} = 2, \quad \text{a.s.}$$

We note that in Csáki et al. [5], we have already proved the upper bounds in Theorems 1.1 and 1.2 with a different constant, and assuming only the first two conditions of (1.5). In particular, under these conditions, we established

$$\begin{aligned} \limsup_{h \rightarrow 0} \sup_{0 \leq t \leq 1-h} \sup_{0 \leq s \leq h} \frac{|Y(t + s) - Y(t)|}{\sqrt{h \log(1/h)}} &\leq 3 \cdot 2^{7/6}, \quad \text{a.s.} \\ \limsup_{T \rightarrow \infty} \sup_{0 \leq t \leq T - a_T} \sup_{0 \leq s \leq a_T} \frac{|Y(t + s) - Y(t)|}{\sqrt{a_T \log(T/a_T)}} &\leq 3 \cdot 2^{7/6}, \quad \text{a.s.} \end{aligned}$$

These bounds were proved using estimates for local times. The reason for which we are able to get lower bounds and the exact constants is that we shall be using different techniques, based on fine analysis of Bessel processes. Also, due to third condition of (1.5), presently we have as well $a_T/T \rightarrow 0$ as $T \rightarrow \infty$. Hence we could have stated Theorem 1.2 with $\sup_{0 \leq t \leq T}$ instead of $\sup_{0 \leq t \leq T - a_T}$ in its present form. The proof of the latter version of Theorem 1.2 would require only a few slight changes in its current proof. We prefer to keep the present form of Theorem 1.2, for it could be still true as an exact $\limsup_{T \rightarrow \infty}$ statement under only the first two conditions of (1.5).

The rest of the paper is organized as follows. The upper bounds in Theorems 1.1 and 1.2 are proved in Section 2, and the lower bounds in Section 3. The proof of the upper bounds is the harder part, requiring careful analysis on path decompositions and deep properties of three-dimensional Bessel processes. The proof of the lower bounds mainly consists in choosing some “nice” random stopping times.

Notation. Throughout the paper, the letter c with subscripts denotes some finite and positive universal constants. When the constants depend on a parameter, say p , they are denoted by $c(p)$ with subscripts.

2. Upper bounds

The upper bounds in Theorems 1.1 and 1.2 are based on some key probability estimates which are stated below as Lemmas 2.7 and 2.8. Let us first recall some known results. Recall that $W(\cdot)$ is a standard Brownian motion, and that $Y(\cdot)$ is defined in (1.1).

The first useful result concerns the distribution of $Y(s)$ for any fixed s . This was evaluated by Biane and Yor [3].

Fact 2.1. *The density function of $Y(s)$ is given by: for $s > 0$ and $x > 0$,*

$$\frac{\mathbb{P}(Y(s) \in dx)}{dx} = \sqrt{\frac{2}{\pi^3 s}} \sum_{k=0}^{\infty} (-1)^k \exp\left(-\frac{(2k+1)^2 x^2}{8s}\right).$$

Comment 2.1.1. From Fact 2.1 it follows that

$$\lim_{\lambda \rightarrow \infty} \frac{1}{\lambda^2} \log \mathbb{P}(Y(1) > \lambda) = -\frac{1}{8}. \tag{2.1}$$

Moreover, for $s > 0$ and $\lambda > 0$,

$$\mathbb{P}(|Y(s)| > \lambda\sqrt{s}) \leq c_1 \exp\left(-\frac{\lambda^2}{8}\right), \tag{2.2}$$

where c_1 is a universal constant (cf. [11]). □

The second theorem we shall make use of in the proof of the upper bounds is time-reversal for Bessel processes, cf. Exercise XI.1.23 in Revuz and Yor [16]. We recall that a three-dimensional Bessel process is the Euclidean modulus of an \mathbb{R}^3 -valued Brownian motion.

Fact 2.2. *Let $T_0 = \inf\{t > 0 : W(t) = 0\}$, the first hitting time of W at 0. Then for any $x > 0$,*

$$\{W(t); 0 \leq t \leq T_0 \mid W(0) = x\} \stackrel{\text{law}}{=} \{R(\mathcal{L}_x - t); 0 \leq t \leq \mathcal{L}_x\}, \tag{2.3}$$

where $R(\cdot)$ is a three-dimensional Bessel process starting from 0, and $\mathcal{L}_x = \sup\{t > 0 : R(t) = x\}$, the last exit time from x .

Comment 2.2.1. Observe that (2.3) also guarantees the identity in law between T_0 (given $W(0) = x$) and \mathcal{L}_x . Thus

$$\mathbb{P}(u < \mathcal{L}_x \leq v) = \int_u^v \frac{x}{(2\pi y^3)^{1/2}} \exp\left(-\frac{x^2}{2y}\right) dy \leq \frac{(v-u)x}{u^{3/2}}, \tag{2.4}$$

for any $v > u > 0$ and $x > 0$. □

The next identity in law, due to Pitman and Yor [15], relates a particular additive functional of the three-dimensional Bessel process and the range of one-dimensional Brownian motion. This will be frequently applied to our situation.

Fact 2.3. *Let $R(\cdot)$ be as before a three-dimensional Bessel process starting from 0. The following identity in law holds:*

$$\int_0^1 \frac{ds}{R(s)} \stackrel{\text{law}}{=} \sup_{0 \leq s \leq 1} W(s) - \inf_{0 \leq s \leq 1} W(s).$$

Comment 2.3.1. It is an immediate consequence of Fact 2.3 and Feller’s exact distribution function of the range of Brownian motion (cf. [8]) that

$$\mathbb{P} \left(\int_0^1 \frac{ds}{R(s)} > \lambda \right) \leq c_2 \exp \left(-\frac{\lambda^2}{2} \right), \quad \lambda > 0, \tag{2.5}$$

with some absolute constant c_2 . By applying the diffusion comparison theorem stated in Theorem XI.3.7 of Revuz and Yor [16] to squared Bessel processes, we deduce the intuitively clear fact that a three-dimensional Bessel process starting from $x > 0$ is stochastically greater than a three-dimensional Bessel process starting from 0. Consequently, for any $x > 0$ and $\lambda > 0$,

$$\mathbb{P} \left(\int_0^1 \frac{ds}{R(s)} > \lambda \mid R(0) = x \right) \leq \mathbb{P} \left(\int_0^1 \frac{ds}{R(s)} > \lambda \right).$$

By means of the Markov and scaling properties we arrive at:

$$\mathbb{P} \left(\int_u^v \frac{ds}{R(s)} > \lambda \right) \leq c_2 \exp \left(-\frac{\lambda^2}{2(v-u)} \right), \tag{2.6}$$

for $v > u \geq 0$ and $\lambda > 0$. □

We start the proof of the upper bounds in Theorems 1.1 and 1.2 with an elementary estimate.

Lemma 2.4. *Let M and N be independent random variables such that for all $x \geq 0$,*

$$\mathbb{P}(M > x) \leq \mu \exp(-\alpha x^2), \quad \mathbb{P}(N > x) \leq \nu \exp(-\beta x^2), \tag{2.7}$$

for some positive constants μ, ν, α and β . Then for all $x \geq 0$,

$$\mathbb{P}(M + N > x) \leq c_3 (1 + x^2) \exp \left(-\frac{\alpha\beta x^2}{\alpha + \beta} \right),$$

where $c_3 = c_3(\mu, \nu, \beta) = \mu + 2\nu + 2\beta\mu\nu$.

Remark. We note that the term in the exponential here is sharp.

Proof. We have

$$\mathbb{P}(M + N > x) \leq \mathbb{P}(M > x) + \mathbb{P}(N > x) + \mathbb{P}(N > x - M, 0 \leq M \leq x).$$

Observe that the last probability term on the right hand side is

$$\begin{aligned} &\leq \mathbb{E} \left(v e^{-\beta(x-M)^2} \mathbb{1}_{\{0 \leq M \leq x\}} \right) \\ &= -v \int_{y \in [0, x]} e^{-\beta(x-y)^2} d_y \mathbb{P}(M > y) \\ &\leq v e^{-\beta x^2} + 2\beta v \int_{y \in [0, x]} (x - y) e^{-\beta(x-y)^2} \mathbb{P}(M > y) dy \\ &\leq v e^{-\beta x^2} + 2\beta \mu v x \int_{y \in [0, x]} e^{-\beta(x-y)^2 - \alpha y^2} dy. \end{aligned}$$

Since $\beta(x - y)^2 + \alpha y^2 \geq \alpha \beta x^2 / (\alpha + \beta)$, this yields,

$$\mathbb{P}(M + N > x) \leq \mu e^{-\alpha x^2} + 2v e^{-\beta x^2} + 2\beta \mu v x^2 \exp \left(-\frac{\alpha \beta x^2}{\alpha + \beta} \right),$$

as desired. □

It is intuitively clear that, if we want $Y(\cdot)$ to get extraordinarily large increments, the Brownian motion $W(\cdot)$ should be close to 0. However, due to the fact that $Y(\cdot)$ is defined only as a principal value, we have to treat it carefully. The next preliminary estimates (Lemmas 2.5 and 2.6) concern the two different situations: (i) $W(\cdot)$ is away from 0; and (ii) $W(\cdot)$ is close to 0. The tail probability of the increment of $Y(\cdot)$ in the second situation is greater, as expected.

Lemma 2.5. *For any positive δ, t, h , and any $\lambda \geq 1$ and $a \geq 1$,*

$$\begin{aligned} &\mathbb{P} \left(|Y(t+h) - Y(t)| > \lambda h^{1/2}, \inf_{s \in [t, t+h]} |W(s)| > 0 \right) \\ &\leq \frac{c_4(\delta) a^2 h^{1/2}}{(t+h)^{1/2}} \exp \left(-\frac{\lambda^2}{6(1+\delta)} \right) + 4 \exp \left(-\frac{(a - \lambda^{-1})^2}{2} \right). \end{aligned} \quad (2.8)$$

Proof. Let $t > 0, h > 0$ and $\lambda \geq 1$. Write $I_1 = I_1(t, h, \lambda)$ for the probability expression on the left hand side of (2.8). Then

$$\begin{aligned} I_1 &= \mathbb{P} \left(|Y(t+h) - Y(t)| > \lambda h^{1/2}, \inf_{s \in [t, t+h]} |W(s)| > 0 \right) \\ &= \mathbb{P} \left(|Y(t/h+1) - Y(t/h)| > \lambda, \inf_{s \in [t/h, t/h+1]} |W(s)| > 0 \right). \end{aligned}$$

Define

$$f(\lambda, x) = \mathbb{P} \left(\left| \int_0^1 \frac{ds}{W(s)} \right| > \lambda, T_0 > 1 \mid W(0) = x \right).$$

It follows from symmetry and the Markov property that

$$I_1 = 2 \int_0^\infty \mathbb{P}(W(t/h) \in dx) f(\lambda, x). \tag{2.9}$$

Let us now estimate $f(\lambda, x)$. According to (and in the notation of) Fact 2.2,

$$f(\lambda, x) = \mathbb{P} \left(\int_{\mathcal{L}_{x-1}}^{\mathcal{L}_x} \frac{ds}{R(s)} > \lambda, \mathcal{L}_x > 1 \right).$$

Fix a $\theta \in (0, 1]$. Then

$$\begin{aligned} f(\lambda, x) &= \sum_{k=1}^\infty \mathbb{P} \left(\int_{\mathcal{L}_{x-1}}^{\mathcal{L}_x} \frac{ds}{R(s)} > \lambda, (k-1)\theta + 1 < \mathcal{L}_x \leq k\theta + 1 \right) \\ &\leq \sum_{k=1}^\infty \mathbb{P} \left(\int_{(k-1)\theta}^{k\theta+1} \frac{ds}{R(s)} > \lambda, (k-1)\theta + 1 < \mathcal{L}_x \leq k\theta + 1 \right). \end{aligned} \tag{2.10}$$

By (2.6),

$$\mathbb{P} \left(\int_{(k-1)\theta}^{k\theta+1} \frac{ds}{R(s)} > \lambda \right) \leq c_2 \exp \left(-\frac{\lambda^2}{2(\theta + 1)} \right),$$

whereas by (2.4),

$$\mathbb{P} \left((k-1)\theta + 1 < \mathcal{L}_x \leq k\theta + 1 \right) \leq \frac{\theta x}{(1 + (k-1)\theta)^{3/2}} \leq \frac{x}{k^{3/2}\theta^{1/2}}.$$

Going back to (2.10) and using Hölder’s inequality, we get that, for any $p > 1$ and $q > 1$ with $1/p + 1/q = 1$,

$$\begin{aligned} f(\lambda, x) &\leq \sum_{k=1}^\infty c_2^{1/p} \exp \left(-\frac{\lambda^2}{2(\theta + 1)p} \right) \frac{x^{1/q}}{k^{3/(2q)}\theta^{1/(2q)}} \\ &= c_2^{1/p} \exp \left(-\frac{\lambda^2}{2(\theta + 1)p} \right) \frac{x^{1/q}}{\theta^{1/(2q)}} \zeta \left(\frac{3}{2q} \right), \end{aligned}$$

where $\zeta(\cdot)$ denotes the Riemann zeta function.

Now let $a \geq 1$ and $\delta > 0$. We first treat the situation when $x \leq a$. Note that $\zeta(3/(2q)) < \infty$ for $q < 3/2$. We can choose $\theta = \theta(\delta) \in (0, 1]$ and $p = p(\delta) > 3$ such that $q \in [1, 3/2)$ and that $(\theta + 1)p = 3(1 + \delta)$. Accordingly,

$$f(\lambda, x) \leq c_5(\delta) a \exp \left(-\frac{\lambda^2}{6(1 + \delta)} \right), \quad \text{for } x \leq a. \tag{2.11}$$

The other situation is $x > a$. By definition,

$$\begin{aligned} f(\lambda, x) &\leq \mathbb{P} \left(\inf_{s \in [0, 1]} W(s) < \lambda^{-1} \mid W(0) = x \right) \\ &\leq \mathbb{P} \left(\inf_{s \in [0, 1]} (W(s) - W(0)) < -(a - \lambda^{-1}) \right) \\ &\leq 2 \exp \left(-\frac{(a - \lambda^{-1})^2}{2} \right). \end{aligned} \tag{2.12}$$

Plugging (2.11)–(2.12) into (2.9), we see that

$$\begin{aligned}
 I_1 &\leq 2c_5(\delta) a \exp\left(-\frac{\lambda^2}{6(1+\delta)}\right) \mathbb{P}\left(0 < W\left(\frac{t}{h}\right) \leq a\right) \\
 &\quad + 4 \exp\left(-\frac{(a-\lambda^{-1})^2}{2}\right) \\
 &\leq \frac{2c_5(\delta) a^2 h^{1/2}}{(t+h)^{1/2}} \exp\left(-\frac{\lambda^2}{6(1+\delta)}\right) + 4 \exp\left(-\frac{(a-\lambda^{-1})^2}{2}\right),
 \end{aligned}$$

proving the lemma. □

Lemma 2.6. *For any positive δ, t, h, λ , and any $b \geq 1$,*

$$\begin{aligned}
 &\mathbb{P}\left(|Y(t+h) - Y(t)| > \lambda h^{1/2}, \inf_{s \in [t, t+h]} |W(s)| = 0\right) \\
 &\leq c_6 \frac{bh^{1/2}}{(t+h)^{1/2}} \exp\left(-\frac{\lambda^2}{8(1+\delta)}\right) + 4 \exp\left(-\frac{b^2}{2}\right). \tag{2.13}
 \end{aligned}$$

Proof. Write $I_2 = I_2(t, h, \lambda)$ for the probability on the left hand side of (2.13). It is easy to estimate I_2 when $t = 0$. Indeed, in this situation, we can use (2.2) to see that

$$\begin{aligned}
 I_2 &= \mathbb{P}\left(|Y(h)| > \lambda h^{1/2}, \inf_{s \in [0, h]} |W(s)| = 0\right) \\
 &\leq c_1 \exp\left(-\frac{\lambda^2}{8}\right), \quad \text{when } t = 0. \tag{2.14}
 \end{aligned}$$

Now assume $t > 0$. Note that

$$I_2 = 2 \int_0^\infty \mathbb{P}(W(t/h) \in dx) g(\lambda, x), \tag{2.15}$$

where

$$g(\lambda, x) = \mathbb{P}\left(\left|\int_0^1 \frac{ds}{W(s)}\right| > \lambda, \inf_{s \in [0, 1]} |W(s)| = 0 \mid W(0) = x\right).$$

Recall T_0 from Fact 2.2. Combining Fact 2.2 (in its notation) with the strong Markov and scaling properties, we obtain:

$$\begin{aligned}
 g(\lambda, x) &= \mathbb{P}\left(\left|\int_0^1 \frac{ds}{W(s)}\right| > \lambda, T_0 \leq 1 \mid W(0) = x\right) \\
 &\leq \mathbb{P}\left(\mathcal{L}_x \leq 1, \int_0^{\mathcal{L}_x} \frac{ds}{R(s)} + \sqrt{1 - \mathcal{L}_x} |Y(1)| > \lambda\right),
 \end{aligned}$$

where R and $Y(1)$ are assumed to be independent. Now fix an integer $n \geq 1$. Clearly,

$$\begin{aligned} g(\lambda, x) &\leq \sum_{k=1}^n \mathbb{P} \left(\int_0^{k/n} \frac{ds}{R(s)} + \sqrt{1 - \frac{k-1}{n}} |Y(1)| > \lambda \right) \\ &= \sum_{k=1}^n \mathbb{P} \left(\sqrt{\frac{k}{n}} \int_0^1 \frac{ds}{R(s)} + \sqrt{1 - \frac{k-1}{n}} |Y(1)| > \lambda \right). \end{aligned}$$

We want to apply Lemma 2.4 to $M = \sqrt{k/n} \int_0^1 ds/R(s)$, $N = \sqrt{1 - (k-1)/n} |Y(1)|$. In view of (2.5) and (2.2), it is seen that (2.7) is satisfied with $\mu = c_2$, $\alpha = n/(2k)$, $\nu = c_1$ and $\beta = 1/8(1 - (k-1)/n)$. It follows from Lemma 2.4 that

$$\begin{aligned} g(\lambda, x) &\leq c_7 \sum_{k=1}^n \frac{1 + \lambda^2}{1 - (k-1)/n} \exp \left(-\frac{\lambda^2}{8 + (8-6k)/n} \right) \\ &\leq c_7 n^2 (1 + \lambda^2) \exp \left(-\frac{\lambda^2}{8(1 + 1/n)} \right). \end{aligned}$$

Let $\delta > 0$. We can choose $n = n(\delta)$ such that $n^{-1} \leq \delta/2$. Therefore,

$$\begin{aligned} g(\lambda, x) &\leq c_8(\delta) (1 + \lambda^2) \exp \left(-\frac{\lambda^2}{8(1 + \delta/2)} \right) \\ &\leq c_9(\delta) \exp \left(-\frac{\lambda^2}{8(1 + \delta)} \right), \end{aligned} \tag{2.16}$$

the last inequality following from the fact that $(1 + \lambda^2) \exp \left(\frac{\lambda^2}{8(1+\delta)} - \frac{\lambda^2}{8(1+\delta/2)} \right)$ is uniformly bounded in $\lambda > 0$.

The estimate (2.16), which holds uniformly in x , is not accurate enough when x is large. Let $b \geq 1$. When $x \geq b$, we have

$$g(\lambda, x) \leq \mathbb{P} \left(\inf_{s \in [0,1]} (W(s) - W(0)) \leq -b \right) \leq 2 \exp \left(-\frac{b^2}{2} \right). \tag{2.17}$$

Using (2.16) for $x \in [0, b]$ and (2.17) for $x \in (b, \infty)$, and in view of (2.15), we obtain that, for $t > 0$,

$$\begin{aligned} I_2 &\leq c_9(\delta) \exp \left(-\frac{\lambda^2}{8(1 + \delta)} \right) \mathbb{P}(0 \leq W(t/h) \leq b) + 4 \exp \left(-\frac{b^2}{2} \right) \\ &\leq c_9(\delta) \frac{h^{1/2} b}{(t+h)^{1/2}} \exp \left(-\frac{\lambda^2}{8(1 + \delta)} \right) + 4 \exp \left(-\frac{b^2}{2} \right). \end{aligned}$$

This, together with (2.14), completes the proof of Lemma 2.6. □

The next two probability estimates (Lemmas 2.7 and 2.8) are the main ingredient in the proof of the upper bounds in Theorems 1.1 and 1.2. More precisely, Lemma 2.7 is our key probability estimate, which will be reinforced later in Lemma 2.8 into a maximal inequality.

Lemma 2.7. For all positive numbers δ, t and h , and all $\lambda \geq 1$,

$$\begin{aligned} & \mathbb{P} \left(|Y(t+h) - Y(t)| > \lambda h^{1/2} \right) \\ & \leq c_{10}(\delta) \left(\frac{h}{t+h} \right)^{1/2} \exp \left(-\frac{\lambda^2}{8(1+\delta)} \right) + 8 \exp \left(-\frac{\lambda^2}{2} \right). \end{aligned}$$

Proof. Take $\delta = 1/6$ and $a = \lambda + \lambda^{-1}$ in (2.8) to see that, for $\lambda \geq 1$,

$$\begin{aligned} & \mathbb{P} \left(|Y(t+h) - Y(t)| > \lambda h^{1/2}, \inf_{s \in [t, t+h]} |W(s)| > 0 \right) \\ & \leq \frac{c_{11} \lambda^2 h^{1/2}}{(t+h)^{1/2}} \exp \left(-\frac{\lambda^2}{7} \right) + 4 \exp \left(-\frac{\lambda^2}{2} \right) \\ & \leq \frac{c_{12} h^{1/2}}{(t+h)^{1/2}} \exp \left(-\frac{\lambda^2}{8} \right) + 4 \exp \left(-\frac{\lambda^2}{2} \right). \end{aligned} \tag{2.18}$$

On the other hand, replacing δ by 2δ and taking $b = \lambda$ in (2.13), we get that

$$\begin{aligned} & \mathbb{P} \left(|Y(t+h) - Y(t)| > \lambda h^{1/2}, \inf_{s \in [t, t+h]} |W(s)| = 0 \right) \\ & \leq \frac{c_6 \lambda h^{1/2}}{(t+h)^{1/2}} \exp \left(-\frac{\lambda^2}{8(1+2\delta)} \right) + 4 \exp \left(-\frac{\lambda^2}{2} \right) \\ & \leq \frac{c_{13}(\delta) h^{1/2}}{(t+h)^{1/2}} \exp \left(-\frac{\lambda^2}{8(1+\delta)} \right) + 4 \exp \left(-\frac{\lambda^2}{2} \right). \end{aligned} \tag{2.19}$$

Combining (2.18) with (2.19) yields the lemma, with $c_{10}(\delta) = c_{12} + c_{13}(\delta)$. \square

Lemma 2.8. For $\delta > 0, x > 0$ and $T > 0, h > 0$,

$$\begin{aligned} & \mathbb{P} \left(\sup_{0 \leq t \leq T} \sup_{0 \leq s \leq h} |Y(t+s) - Y(t)| > x\sqrt{h} \right) \\ & \leq c_{14}(\delta) \left(\sqrt{\frac{T+h}{h}} \exp \left(-\frac{x^2}{8(1+\delta)} \right) + \frac{T+h}{h} \exp \left(-\frac{x^2}{2(1+\delta)} \right) \right). \end{aligned}$$

Proof. Lemma 2.8 is a consequence of Lemma 2.7. The dyadic approximation argument we are using here is not new, and can be found for example in the proofs of Lemma 1.1.1 in Csörgő and Révész [7], and of Lemma 2.2 in Csáki et al. [4]. The main ideas go back to Lévy [13].

For positive real number s and integer n put $s_n = 2^{-n} \lfloor 2^n s \rfloor$. We have

$$\begin{aligned} |Y(t+s) - Y(t)| & \leq |Y((t+s)_n) - Y(t_n)| \\ & \quad + \sum_{j=0}^{\infty} |Y((t+s)_{n+j+1}) - Y((t+s)_{n+j})| \\ & \quad + \sum_{j=0}^{\infty} |Y(t_{n+j+1}) - Y(t_{n+j})|. \end{aligned} \tag{2.20}$$

Consider $t \in [0, T]$ and $s \in [0, h]$. Clearly $|(t + s)_n - t_n| \leq h + 2^{-n}$. Therefore, by Lemma 2.7, for any $\lambda \geq 1$,

$$\begin{aligned} & \mathbb{P} \left(|Y((t + s)_n) - Y(t_n)| > \lambda \sqrt{h + 2^{-n}} \right) \\ & \leq c_{10}(\delta) \sqrt{\frac{h + 2^{-n}}{(t + s)_n}} \exp \left(-\frac{\lambda^2}{8(1 + \delta)} \right) + 8 \exp \left(-\frac{\lambda^2}{2} \right). \end{aligned} \tag{2.21}$$

Note that there exists an i with $0 \leq i \leq (T + h)2^n$, such that $(t + s)_n = i2^{-n}$, i.e. $t + s \in [i2^{-n}, (i + 1)2^{-n})$. If $i = 0$, then the probability on the left hand side of (2.21) is simply 0. For each i with $1 \leq i \leq (T + h)2^n$, there are at most $2^nh + 3$ different values of $(t)_n$ such that $(t + s)_n = i2^{-n}$. Therefore, by (2.21),

$$\begin{aligned} & \mathbb{P} \left(\sup_{0 \leq t \leq T} \sup_{0 \leq s \leq h} |Y((t + s)_n) - Y(t_n)| > \lambda \sqrt{h + 2^{-n}} \right) \\ & \leq (2^nh + 3) \sum_{1 \leq i \leq (T+h)2^n} \left(\frac{c_{10}(\delta) \sqrt{h + 2^{-n}}}{\sqrt{i}2^{-n}} \exp \left(-\frac{\lambda^2}{8(1 + \delta)} \right) + 8 \exp \left(-\frac{\lambda^2}{2} \right) \right). \end{aligned}$$

Since $\sum_{1 \leq i \leq a} i^{-1/2} \leq 2\sqrt{a}$ for any $a \geq 1$, we arrive at:

$$\mathbb{P} \left(\sup_{0 \leq t \leq T} \sup_{0 \leq s \leq h} |Y((t + s)_n) - Y(t_n)| > \lambda \sqrt{h + 2^{-n}} \right) \leq \phi_1(n), \tag{2.22}$$

where

$$\begin{aligned} \phi_1(n) &= c_{15}(\delta) \sqrt{(T + h)2^n} (2^nh + 3)^{3/2} \exp \left(-\frac{\lambda^2}{8(1 + \delta)} \right) \\ & \quad + 8(2^nh + 3)(T + h)2^n \exp \left(-\frac{\lambda^2}{2} \right). \end{aligned}$$

Similarly, for any $u \in [0, T + h]$ and integer $j \geq 0$, since $|u_{n+j+1} - u_{n+j}| \leq 2^{-(n+j+1)}$, we have, by Lemma 2.7, for any $\lambda_j > 0$,

$$\begin{aligned} & \mathbb{P} \left(|Y(u_{n+j+1}) - Y(u_{n+j})| > \lambda_j 2^{-(n+j+1)/2} \right) \\ & \leq c_{10}(\delta) \sqrt{\frac{2^{-(n+j+1)}}{u_{n+j+1}}} \exp \left(-\frac{\lambda_j^2}{8(1 + \delta)} \right) + 8 \exp \left(-\frac{\lambda_j^2}{2} \right), \end{aligned}$$

which leads to:

$$\begin{aligned} & \mathbb{P} \left(\sup_{0 \leq u \leq T+h} |Y(u_{n+j+1}) - Y(u_{n+j})| > \lambda_j 2^{-(n+j+1)/2} \right) \\ & \leq \sum_{1 \leq i \leq (T+h)2^{n+j+1}} \left(\frac{c_{10}(\delta)}{\sqrt{i}} \exp \left(-\frac{\lambda_j^2}{8(1 + \delta)} \right) + 8 \exp \left(-\frac{\lambda_j^2}{2} \right) \right) \tag{2.23} \\ & \leq \phi_2(n, j), \end{aligned}$$

with

$$\begin{aligned} \phi_2(n, j) &= c_{16}(\delta) \sqrt{(T+h)2^{n+j+1}} \exp\left(-\frac{\lambda_j^2}{8(1+\delta)}\right) + 8(T+h)2^{n+j+1} \exp\left(-\frac{\lambda_j^2}{2}\right). \end{aligned}$$

Collect (2.20), (2.22) and (2.23) to see that,

$$\begin{aligned} \mathbb{P} &\left(\sup_{0 \leq t \leq T} \sup_{0 \leq s \leq h} |Y(t+s) - Y(t)| > \lambda \sqrt{h+2^{-n}} + 2 \sum_{j=0}^{\infty} \lambda_j 2^{-(n+j+1)/2} \right) \\ &\leq \phi_1(n) + 2 \sum_{j=0}^{\infty} \phi_2(n, j). \end{aligned}$$

Let $\mu = \mu(\delta) \in (0, 1)$ be such that $\sqrt{1+2\mu} + c_{20} \sqrt{8\mu} \leq 1 + \delta$, where c_{20} is the absolute constant defined in (2.24) below. Choose $\lambda_j = \sqrt{\lambda^2 + 4j}$, and let $n = n(h, \delta)$ be such that $2^{-n} \in [\mu h, 2\mu h]$. Then

$$\begin{aligned} \phi_1(n) + 2 \sum_{j=0}^{\infty} \phi_2(n, j) &\leq c_{17}(\delta) \left(\sqrt{\frac{T+h}{h}} \exp\left(-\frac{\lambda^2}{8(1+\delta)}\right) + \frac{T+h}{h} \exp\left(-\frac{\lambda^2}{2}\right) \right). \end{aligned}$$

On the other hand, since $\lambda_j \leq \lambda + 2j^{1/2}$ and $\lambda \geq 1$, we have,

$$\sum_{j=0}^{\infty} \lambda_j 2^{-(j+1)/2} \leq \sum_{j=0}^{\infty} (\lambda + 2j^{1/2}) 2^{-(j+1)/2} = c_{18} \lambda + c_{19} \leq c_{20} \lambda, \tag{2.24}$$

(with $c_{18} = \sum_{j=0}^{\infty} 2^{-(j+1)/2}$, $c_{19} = 2 \sum_{j=0}^{\infty} j^{1/2} 2^{-(j+1)/2}$ and $c_{20} = c_{18} + c_{19}$), which implies

$$\lambda \sqrt{h+2^{-n}} + 2 \sum_{j=0}^{\infty} \lambda_j 2^{-(n+j+1)/2} \leq \sqrt{1+2\mu} \lambda \sqrt{h} + c_{20} \sqrt{8\mu} \lambda \sqrt{h}.$$

Since $\sqrt{1+2\mu} + c_{20} \sqrt{8\mu} \leq 1 + \delta$, we have proved that

$$\begin{aligned} \mathbb{P} &\left(\sup_{0 \leq t \leq T} \sup_{0 \leq s \leq h} |Y(t+s) - Y(t)| > (1+\delta) \lambda \sqrt{h} \right) \\ &\leq c_{17}(\delta) \left(\sqrt{\frac{T+h}{h}} \exp\left(-\frac{\lambda^2}{8(1+\delta)}\right) + \frac{T+h}{h} \exp\left(-\frac{\lambda^2}{2}\right) \right). \end{aligned}$$

This holds actually for all $\lambda > 0$ (when $\lambda \in (0, 1)$, we only have to take an enlarged value of $c_{17}(\delta)$ if necessary). Taking $x = (1 + \delta)\lambda$, and since $\delta > 0$ is arbitrary, this yields the lemma. □

Proof of the upper bounds in Theorems 1.1 and 1.2. To check the upper bound in Theorem 1.1, we write

$$\Theta(h) = \sup_{0 \leq t \leq 1} \sup_{0 \leq s \leq h} |Y(t+s) - Y(t)|.$$

Let $\delta > 0$ and $h_k = h_k(\delta) = k^{-3/\delta}$. Applying Lemma 2.8 to $T = 1$, $h = h_k$ and $x = 2(1 + \delta)\sqrt{\log(1/h_k)}$ gives that

$$\begin{aligned} & \mathbb{P}\left(\Theta(h_k) > 2(1 + \delta)\sqrt{h_k \log(1/h_k)}\right) \\ & \leq 2c_{14}(\delta) \left(\frac{1}{\sqrt{h_k}} \exp\left(-\frac{1 + \delta}{2} \log \frac{1}{h_k}\right) + \frac{1}{h_k} \exp\left(-2(1 + \delta) \log \frac{1}{h_k}\right) \right) \\ & = 2c_{14}(\delta) (h_k^{\delta/2} + h_k^{1+2\delta}), \end{aligned}$$

which is summable in k . By the Borel–Cantelli lemma, almost surely for all large k , $\Theta(h_k) \leq 2(1 + \delta)\sqrt{h_k \log(1/h_k)}$. Now let $h \in [h_{k+1}, h_k]$. We have

$$\frac{\Theta(h)}{\sqrt{h \log(1/h)}} \leq \frac{\Theta(h_k)}{\sqrt{h_{k+1} \log(1/h_k)}} \leq 2(1 + \delta) \frac{\sqrt{h_k}}{\sqrt{h_{k+1}}}.$$

Since $\sqrt{h_{k+1}}/\sqrt{h_k} \rightarrow 1$ (as $k \rightarrow \infty$), we obtain that

$$\limsup_{h \rightarrow 0} \frac{\Theta(h)}{\sqrt{h \log(1/h)}} \leq 2(1 + \delta), \quad \text{a.s.}$$

This yields the upper bound in Theorem 1.1, as δ can be as close to 0 as possible.

The upper bound in Theorem 1.2 is proved using exactly the same argument, considering $\tilde{\Theta}(T) = \sup_{0 \leq t \leq T-a_T} \sup_{0 \leq s \leq a_T} |Y(t+s) - Y(t)|$ instead of $\Theta(h)$. □

3. Lower bounds

As before, $W(\cdot)$ is a standard Brownian motion, and $Y(\cdot)$ denotes the principal value defined in (1.1)–(1.2). The proof of the lower bounds in Theorems 1.1 and 1.2 relies on the following estimate.

Lemma 3.1. For $T \geq 2a > 0$, $\varepsilon \in (0, 1)$, $\delta > 0$ and $\lambda > 0$,

$$\begin{aligned} & \mathbb{P}\left(\sup_{0 \leq t \leq T-a} (Y(t+a) - Y(t)) \leq \lambda\sqrt{a}\right) \\ & \leq 5\left(\frac{a}{T}\right)^{\varepsilon/2} + \exp\left(-c_{21}(\delta) \left(\frac{T}{a}\right)^{(1-\varepsilon)/2} e^{-(1+\delta)\lambda^2/8}\right). \end{aligned}$$

Proof. Let us construct an increasing sequence of stopping times $\{\eta_k = \eta_k(a)\}_{k \geq 0}$ by: $\eta_0 = 0$ and

$$\eta_{k+1} = \inf\{t > \eta_k + a : W(t) = 0\}, \quad k \geq 0.$$

Let

$$v_T = v_T(a) = \max\{i \geq 0 : \eta_i \leq T - a\}.$$

Clearly,

$$\sup_{0 \leq t \leq T-a} (Y(t+a) - Y(t)) \geq \max_{0 \leq i \leq v_T} (Y(\eta_i + a) - Y(\eta_i)),$$

which yields

$$\begin{aligned} & \mathbb{P} \left(\sup_{0 \leq t \leq T-a} (Y(t+a) - Y(t)) \leq \lambda\sqrt{a} \right) \\ & \leq \mathbb{P} \left(\max_{0 \leq i \leq v_T} (Y(\eta_i + a) - Y(\eta_i)) \leq \lambda\sqrt{a} \right) \\ & \leq \mathbb{P} \left(v_T < \left(\frac{T}{a} \right)^{(1-\varepsilon)/2} \right) \\ & \quad + \mathbb{P} \left(\max_{0 \leq i \leq (T/a)^{(1-\varepsilon)/2}} (Y(\eta_i + a) - Y(\eta_i)) \leq \lambda\sqrt{a} \right). \end{aligned} \tag{3.1}$$

It was shown in Csáki and Földes [6] that for $T \geq 2a$ and $\varepsilon \in (0, 1)$,

$$\mathbb{P} \left(v_T < \left(\frac{T}{a} \right)^{(1-\varepsilon)/2} \right) \leq 3\sqrt{2} \left(\frac{a}{T} \right)^{\varepsilon/2}. \tag{3.2}$$

On the other hand, by the strong Markov property, $\{Y(\eta_i + a) - Y(\eta_i)\}_{i \geq 0}$ are iid variables, having the same law as $Y(a)$. Hence,

$$\begin{aligned} & \mathbb{P} \left(\max_{0 \leq i \leq (T/a)^{(1-\varepsilon)/2}} (Y(\eta_i + a) - Y(\eta_i)) \leq \lambda\sqrt{a} \right) \\ & = (1 - \mathbb{P}(Y(a) > \lambda\sqrt{a}))^{\lfloor (T/a)^{(1-\varepsilon)/2} \rfloor}. \end{aligned}$$

In view of (2.1), we have, for any $\delta > 0$ and $\lambda > 0$,

$$\mathbb{P}(Y(1) > \lambda) \geq c_{22}(\delta) e^{-(1+\delta)\lambda^2/8},$$

which implies that

$$\begin{aligned} & \mathbb{P} \left(\max_{0 \leq i \leq (T/a)^{(1-\varepsilon)/2}} (Y(\eta_i + a) - Y(\eta_i)) \leq \lambda\sqrt{a} \right) \\ & \leq \left(1 - c_{22}(\delta) e^{-(1+\delta)\lambda^2/8} \right)^{\lfloor (T/a)^{(1-\varepsilon)/2} \rfloor} \\ & \leq \exp \left(-\lfloor (T/a)^{(1-\varepsilon)/2} \rfloor c_{22}(\delta) e^{-(1+\delta)\lambda^2/8} \right), \end{aligned} \tag{3.3}$$

the last inequality following from the relation that $1 - x \leq e^{-x}$ (for $x \geq 0$). Assembling (3.1)–(3.3) yields the lemma. \square

Proof of the lower bounds in Theorems 1.1 and 1.2. Again, these bounds are proved using the same argument. For the upper bounds, we chose to prove Theorem 1.1 (cf. Section 2), so we give the proof of the lower bound in Theorem 1.2 here. The corresponding proof for Theorem 1.1 is similar and easier.

Let a_T be a function satisfying (1.5). Since $\frac{\log(T/a_T)}{\log \log T} \rightarrow \infty$, we have $T \geq 2a_T$ for all large T , say $T \geq n_0$. Consider

$$\Lambda(T) = \sup_{0 \leq t \leq T - a_T} (Y(t + a_T) - Y(t)), \quad T > 0.$$

Let $\delta > 0$ and $\varepsilon \in (0, 1/2)$. Define $T_k = T_k(\delta) = (1 + \delta)^k$. Apply Lemma 3.1 to $T = T_k$, $a = a_{T_k}$ and $\lambda = \frac{2\sqrt{1-2\varepsilon}}{\sqrt{1+\delta}} \sqrt{\log(T/a_{T_k})}$, to see that for $T \geq n_0$,

$$\begin{aligned} & \mathbb{P} \left(\Lambda(T_k) \leq \frac{2\sqrt{1-2\varepsilon}}{\sqrt{1+\delta}} \sqrt{a_{T_k} \log(T_k/a_{T_k})} \right) \\ & \leq 5 \left(\frac{a_{T_k}}{T_k} \right)^{\varepsilon/2} + \exp \left(-c_{21}(\delta) \left(\frac{T_k}{a_{T_k}} \right)^{(1-\varepsilon)/2} \exp \left(-\frac{(1-2\varepsilon) \log(T_k/a_{T_k})}{2} \right) \right) \\ & = 5 \left(\frac{a_{T_k}}{T_k} \right)^{\varepsilon/2} + \exp \left(-c_{21}(\delta) \left(\frac{T_k}{a_{T_k}} \right)^{\varepsilon/2} \right). \end{aligned}$$

Since $\frac{\log(T/a_T)}{\log \log T} \rightarrow \infty$, we have $T/a_T \geq (\log T)^{3/\varepsilon}$ for large T . Therefore, $\sum_k (a_{T_k}/T_k)^{\varepsilon/2} < \infty$ and $\sum_k \exp(-c_{21}(\delta) (T_k/a_{T_k})^{\varepsilon/2}) < \infty$. By the Borel–Cantelli lemma, almost surely for all large k ,

$$\Lambda(T_k) > \frac{2\sqrt{1-2\varepsilon}}{\sqrt{1+\delta}} \sqrt{a_{T_k} \log(T_k/a_{T_k})}.$$

Let

$$\Lambda_0(T) = \sup_{0 \leq t \leq T - a_T} \sup_{0 \leq s \leq a_T} (Y(t + s) - Y(t)).$$

Clearly $T \mapsto \Lambda_0(T)$ is non-decreasing, such that $\Lambda_0(T) \geq \Lambda(T)$. Therefore, for $T \in [T_k, T_{k+1}]$,

$$\frac{\Lambda_0(T)}{\sqrt{a_T \log(T/a_T)}} \geq \frac{\Lambda_0(T_k)}{\sqrt{a_{T_{k+1}} \log(T_{k+1}/a_{T_k})}} > \frac{2\sqrt{1-2\varepsilon}}{\sqrt{1+\delta}} \frac{\sqrt{a_{T_k} \log(T_k/a_{T_k})}}{\sqrt{a_{T_{k+1}} \log(T_{k+1}/a_{T_k})}}.$$

Since $\log(T/a_T) \rightarrow \infty$ (cf. (1.5)),

$$\frac{\log(T_k/a_{T_k})}{\log(T_{k+1}/a_{T_k})} = \frac{\log(T_k/a_{T_k})}{\log(T_k/a_{T_k}) + \log(1+\delta)} \rightarrow 1, \quad k \rightarrow \infty,$$

and since $T \mapsto T/a_T$ is non-decreasing,

$$\frac{a_{T_k}}{a_{T_{k+1}}} \geq \frac{T_k}{T_{k+1}} = \frac{1}{1+\delta},$$

we obtain:

$$\liminf_{T \rightarrow \infty} \frac{\Lambda_0(T)}{\sqrt{a_T} \log(T/a_T)} \geq \frac{2\sqrt{1-2\varepsilon}}{1+\delta}, \quad \text{a.s.}$$

Sending ε and δ to 0 yields the lower bound in Theorem 1.2. \square

Acknowledgements. The authors acknowledge support via fellowships from the Paul Erdős Summer Research Center of Mathematics, Budapest. Miklós Csörgő furthermore benefits from a Paul Erdős Visiting Professorship of the Center. Cooperation between the authors was also supported by the joint French–Hungarian Intergovernmental Grant “Balaton” (grant no. F25/97). We also wish to thank an anonymous referee and an associate editor for insightful comments and suggestions.

References

- Bertoin, J.: Excursions of a $BES_0(d)$ and its drift term ($0 < d < 1$), *Probab. Th. Rel. Fields*, **84**, 231–250 (1990)
- Bertoin, J.: Cauchy’s principal value of local times of Lévy processes with no negative jumps via continuous branching processes, *Electronic J. Probab.*, **2**, paper no. 6, pp. 1–12 (1997)
- Biane, P., Yor, M.: Valeurs principales associées aux temps locaux browniens, *Bull. Sci. Math.*, **111**, 23–101 (1987)
- Csáki, E., Csörgő, M., Földes, A., Révész, P.: How big are the increments of the local time of a Wiener process? *Ann. Probab.*, **11**, 593–608 (1983)
- Csáki, E., Csörgő, M., Földes, A., Shi, Z.: Path properties of Cauchy’s principal values related to local time. (Preprint)
- Csáki, E., Földes, A.: On the narrowest tube of a Wiener process, *Coll. Math. Soc. J. Bolyai*, **36**, 173–197. *Limit Theorems in Probability and Statistics*, (P. Révész, ed.) North-Holland, Amsterdam, 1984
- Csörgő, M., Révész, P.: *Strong Approximations in Probability and Statistics*, Academic Press, New York, 1981
- Feller, W.: The asymptotic distribution of the range of sums of independent random variables, *Ann. Math. Statist.*, **22**, 427–432 (1951)
- Fitzsimmons, P.J., Gettoor, R.K.: On the distribution of the Hilbert transform of the local time of a symmetric Lévy process, *Ann. Probab.*, **20**, 1484–1497 (1992)
- Fukushima, M.: *Dirichlet Forms and Markov Processes*, North-Holland, Amsterdam, 1980
- Hu, Y., Shi, Z.: An iterated logarithm law for Cauchy’s principal value of Brownian local times. In: *Exponential Functionals and Principal Values Related to Brownian Motion* (M. Yor, ed.), pp. 131–154, *Biblioteca de la Revista Matemática Iberoamericana*, Madrid, 1997
- Itô, K., McKean, H.P.: *Diffusion Processes and Their Sample Paths*, Springer, Berlin, 1965
- Lévy, P.: *Théorie de l’Addition des Variables Aléatoires*, Gauthier-Villars, Paris, 1937
- Papanicolaou, G., Stroock, D., Varadhan, S.R.S.: *Martingale approach to some limit theorems*, *Duke Univ. Maths. Series III*, (Statistical Mechanics Dynamical Systems), 1977
- Pitman, J.W., Yor, M.: Quelques identités en loi pour les processus de Bessel. In: *Homage to P.-A. Meyer et J. Neveu*, *Astérisque*, **236** pp. 249–276. *Société Mathématique de France*, Paris, 1996
- Revuz, D., Yor, M.: *Continuous Martingales and Brownian Motion*, Third Edition, Springer, Berlin, 1999

17. Yamada, T.: Principal values of Brownian local times and their related topics. In: Itô's Stochastic Calculus and Probability Theory, (N. Ikeda et al., eds.), pp. 413–422. Springer, Tokyo, 1996
18. Yor, M., editor: Exponential Functionals and Principal Values Related to Brownian Motion, Biblioteca de la Revista Matemática Iberoamericana, Madrid, 1997
19. Yor, M.: Some Aspects of Brownian Motion. Part II: Some Recent Martingale Problems, ETH Zürich Lectures in Mathematics, Birkhäuser, Basel, 1997