

Addendum: Updated constraints on non-standard interactions from global analysis of oscillation data

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In this addendum we re-assess the constraints on Non-Standard Interactions (NSI) from the global analysis of neutrino oscillation data after including the new results released since the publication of ref. [1], in particular those presented at the Neutrino2020 conference. The new data considered here includes the total energy spectrum and the day-night asymmetry of the 2970-day SK4 solar neutrino sample [2], as well as the latest results from long-baseline (LBL) experiments T2K [3, 4] and NOvA [5, 6]. In addition, we have updated the reactor experiments Double-Chooz [7, 8] to 1276/587 days of far/near detector data and RENO [9, 10] to 2908 days of exposure.

The main effect driven by the new results concerns the analysis of solar and KamLAND data discussed in section 3 of ref. [1]. As explained there, at the time of publication there was a tension of $\Delta\chi^2 \sim 7.4$ between these two data sets within the context of the 3ν oscillation analysis, arising from a combination of two effects: (a) the ${}^8\text{B}$ measurements performed by SNO, SK and Borexino did not show any evidence of the low energy spectrum

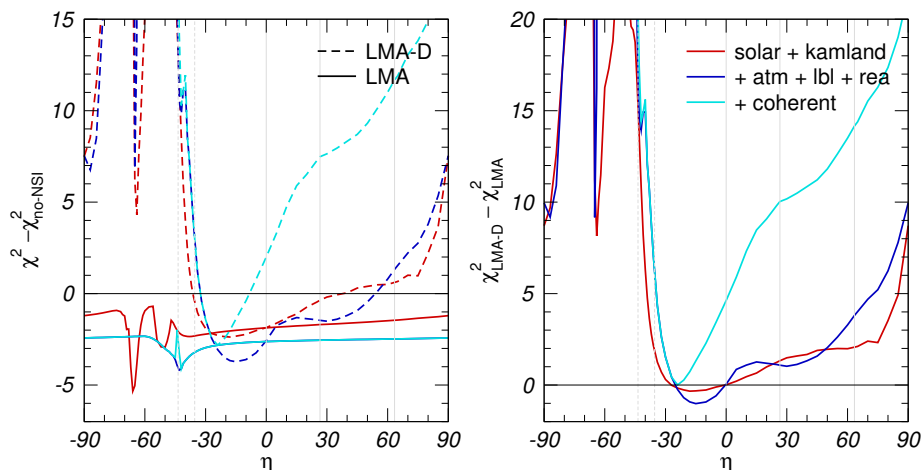


Figure 1. Left: $\chi^2_{\text{LMA}}(\eta) - \chi^2_{\text{no-NSI}}$ (full lines) and $\chi^2_{\text{LMA-D}}(\eta) - \chi^2_{\text{no-NSI}}$ (dashed lines) for the analysis of different data combinations (as labeled in the figure) as a function of the NSI quark coupling parameter η . The full dark blue and light blue curves lie on top of each other. Right: $\chi^2_{\text{dark}} - \chi^2_{\text{light}} \equiv \chi^2_{\text{LMA-D}}(\eta) - \chi^2_{\text{LMA}}(\eta)$ as a function of η . See text for details.

turn-up expected in the standard LMA-MSW [11, 12] solution for the value of Δm^2_{21} favored by KamLAND, and (b) the observation of a non-vanishing day-night asymmetry in SK, whose size was considerably larger than what predicted for the Δm^2_{21} value indicated by KamLAND. Such tension could be alleviated in presence of a non-standard matter potential, thus leading to a sizable decrease in the minimum χ^2 for the LMA solution for most values of η ($\Delta\chi^2 \sim -7 \rightarrow -11$), as could be observed in the left panel in figure 4 of ref. [1]. Correspondingly, in figure 2 of the same work, which showed the two-dimensional projections on the matter potential parameters ($\varepsilon_D^\eta, \varepsilon_N^\eta$) of the 1 σ , 90%, 2 σ , 99% and 3 σ CL (2 dof) allowed regions from the analysis of solar and KamLAND data in the presence of non-standard neutrino-matter interactions, the 3 ν standard LMA oscillation scenario ($\varepsilon_D^\eta = \varepsilon_N^\eta = 0$) was outside of such allowed regions for most values of η .

As discussed in ref. [13], with the updated SK4 solar data the tension between the best fit Δm^2_{21} of KamLAND and that of the solar results has decreased to $\Delta\chi^2_{\text{solar}} = 1.3$. This is due to both the smaller day-night asymmetry, and the slightly more pronounced turn-up in the low energy part of the spectrum. So now in the left panel in figure 1 we see that for the LMA solution the fit with NSI leads to a decrease of about 1 unit of χ^2 for most values of η . Correspondingly in figure 2 the 3 ν standard LMA oscillation scenario, $\varepsilon_D^\eta = \varepsilon_N^\eta = 0$ lies inside the 1 σ LMA allowed regions for most values of η . Concerning the status of the LMA-D solution, the right panel in figure 1 shows that now LMA-D is allowed below 3 σ for $\eta > -40^\circ$ in the analysis of solar+KamLAND, for $-38^\circ \lesssim \eta \lesssim 87^\circ$ in the global oscillation analysis, and for $-38^\circ \lesssim \eta \lesssim 20^\circ$ when including information from the total event rate at COHERENT.¹ From the left panel we read that the best fit for the global analysis of

¹We remind the reader that, as discussed in section 5 of ref. [1], while oscillation constraints apply to models where NSI are generated by arbitrarily light mediators, there is a minimum mediator mass for which the bounds of COHERENT are relevant, which we estimate to be $\mathcal{O}(10\text{--}50)$ MeV (see also refs. [14, 15]).

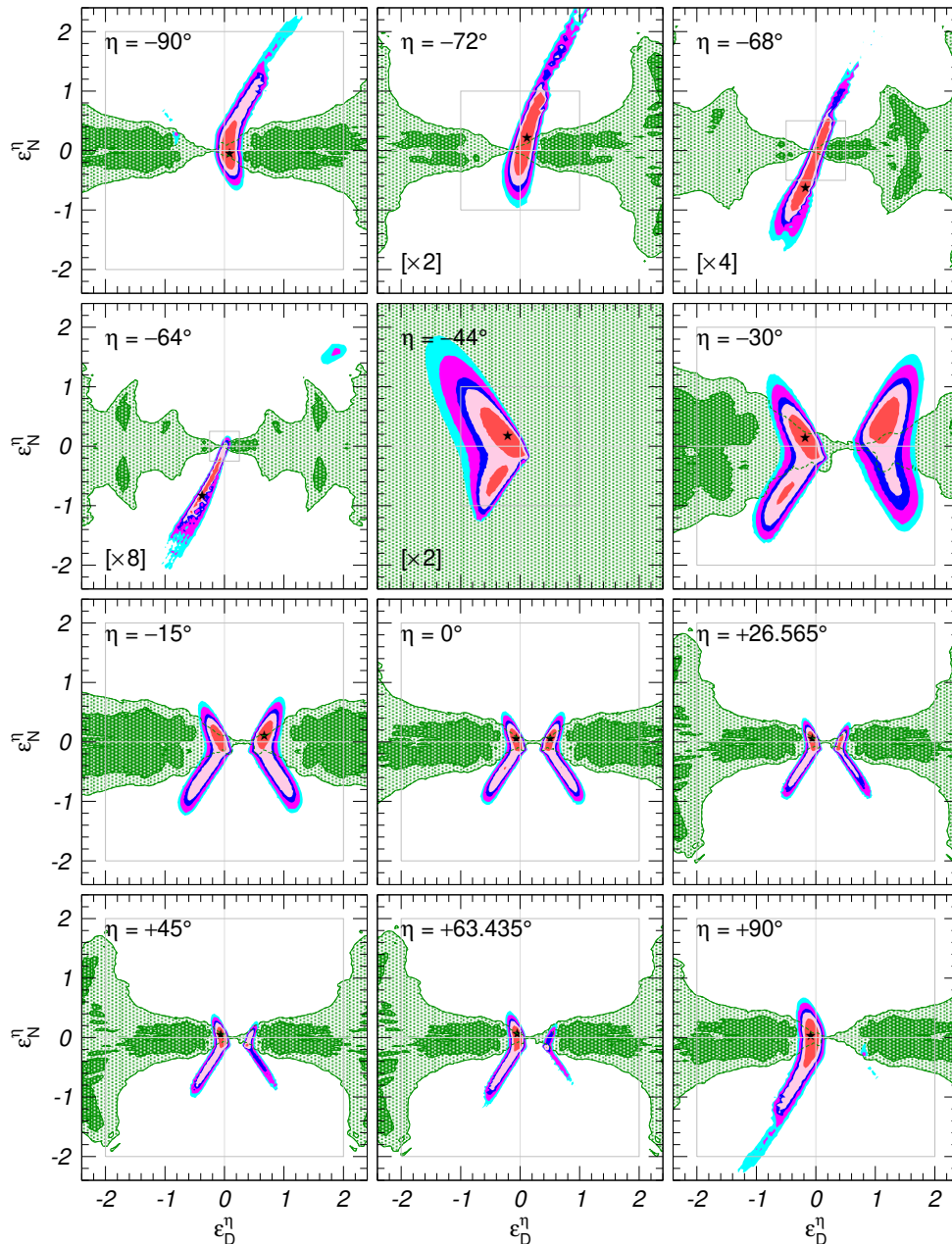


Figure 2. Two-dimensional projections of the 1σ , 90%, 2σ , 99% and 3σ CL (2 dof) allowed regions from the analysis of solar and KamLAND data in the presence of non-standard matter potential for the matter potential parameters $(\varepsilon_D^\eta, \varepsilon_N^\eta)$, for $\sin^2 \theta_{13} = 0.022$ and after marginalizing over the oscillation parameters. The best fit point is marked with a star. The results are shown for fixed values of the NSI quark coupling parameter η . The panels with a scale factor “[$\times N$]” in their lower-left corner have been “zoomed-out” by such factor with respect to the standard axis ranges, hence the grey square drawn in each panel always corresponds to $\max(|\varepsilon_D^\eta|, |\varepsilon_N^\eta|) = 2$ and has the same size in all the panels. For illustration we also show as shaded green areas the 90% and 3σ CL allowed regions from the analysis of the atmospheric and LBL data. Note that, as a consequence of the periodicity of η , the regions in the first ($\eta = -90^\circ$) and last ($\eta = +90^\circ$) panels are identical up to an overall sign flip.

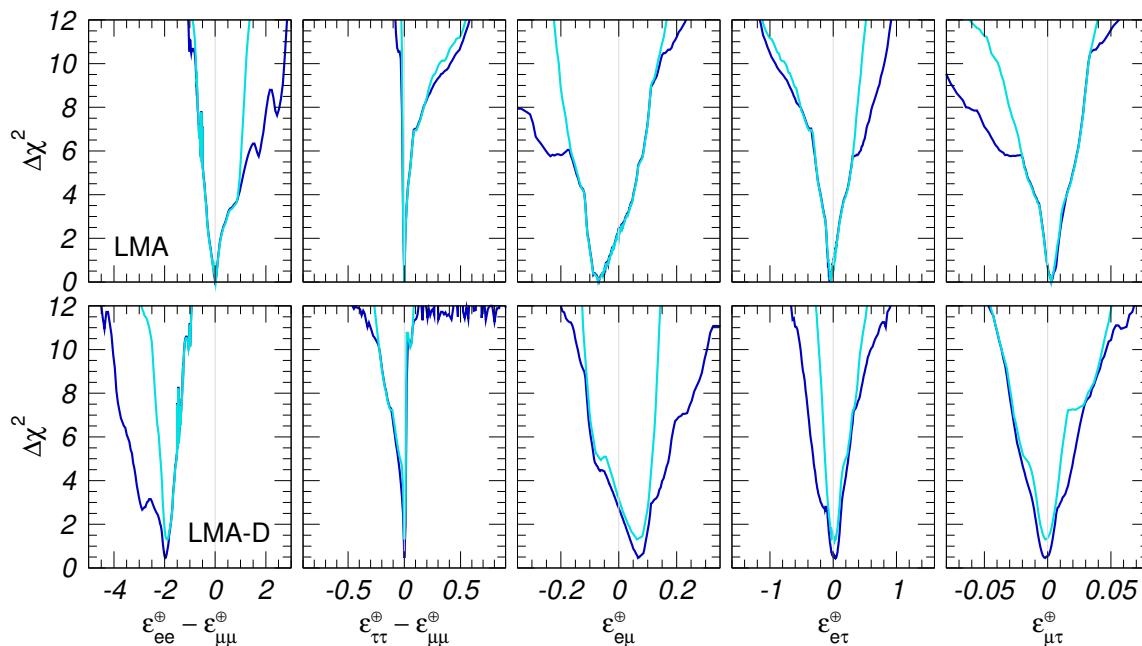


Figure 3. Dependence of the $\Delta\chi^2$ function on the effective NSI parameters relevant for matter effects in LBL experiments with arbitrary values of η , from the global analysis of solar, atmospheric, LBL-CPC and reactor data (blue lines) and including also COHERENT (cyan lines). The upper (lower) panels correspond to solutions within the LMA (LMA-D) subset of parameter space.

oscillations and also in combination with COHERENT corresponds to $\eta \sim -45^\circ$ for LMA. For LMA-D the best fit for OSC (OSC+COH) is obtained for $\eta \sim -15^\circ$ ($\eta \sim -20^\circ$).

In figure 3 we plot the dependence of the global χ^2 on each NSI effective coupling relevant for neutrino propagation in the Earth after marginalization over all other parameters including η , so that the $\Delta\chi^2$ functions plotted in the figure are defined with respect to the absolute minimum for any η . When compared with the corresponding figure in ref. [1] we observe that, following the discussion above, the minimum χ^2 within LMA and LMA-D are almost the same, while previously we had $\Delta\chi_{\min, \text{LMA-D}}^2 \sim 3$. The other observable difference is that including COHERENT has now a larger impact on the allowed ranges in LMA.

Finally, for the sake of convenience and comparison with previous results we list in the first columns in table 1 the 95% CL ranges for NSI with up-quarks only, down-quarks only, and protons. Generically the allowed ranges within LMA are slightly reduced and, as expected, the allowed ranges for $\varepsilon_{ee} - \varepsilon_{\mu\mu}$ are now more symmetric around zero.

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OSC			+ COHERENT	
	LMA	LMA \oplus LMA-D		LMA = LMA \oplus LMA-D
$\varepsilon_{ee}^u - \varepsilon_{\mu\mu}^u$	$[-0.072, +0.321]$	$\oplus[-1.042, -0.743]$	ε_{ee}^u	$[-0.067, +0.547]$
$\varepsilon_{\tau\tau}^u - \varepsilon_{\mu\mu}^u$	$[-0.001, +0.018]$	$[-0.016, +0.018]$	$\varepsilon_{\mu\mu}^u$	$[-0.076, +0.455]$
$\varepsilon_{e\mu}^u$	$[-0.050, +0.020]$	$[-0.050, +0.059]$	$\varepsilon_{\tau\tau}^u$	$[-0.076, +0.455]$
$\varepsilon_{e\tau}^u$	$[-0.077, +0.098]$	$[-0.111, +0.098]$	$\varepsilon_{e\mu}^u$	$[-0.050, +0.020]$
$\varepsilon_{\mu\tau}^u$	$[-0.006, +0.007]$	$[-0.006, +0.007]$	$\varepsilon_{e\tau}^u$	$[-0.077, +0.099]$
$\varepsilon_{ee}^d - \varepsilon_{\mu\mu}^d$	$[-0.084, +0.326]$	$\oplus[-1.081, -1.026]$	ε_{ee}^d	$[-0.063, +0.503]$
$\varepsilon_{\tau\tau}^d - \varepsilon_{\mu\mu}^d$	$[-0.001, +0.018]$	$[-0.001, +0.018]$	$\varepsilon_{\mu\mu}^d$	$[-0.072, +0.408]$
$\varepsilon_{e\mu}^d$	$[-0.051, +0.020]$	$[-0.051, +0.038]$	$\varepsilon_{\tau\tau}^d$	$[-0.072, +0.407]$
$\varepsilon_{e\tau}^d$	$[-0.077, +0.098]$	$[-0.077, -0.098]$	$\varepsilon_{e\mu}^d$	$[-0.050, +0.020]$
$\varepsilon_{\mu\tau}^d$	$[-0.006, +0.007]$	$[-0.006, +0.007]$	$\varepsilon_{e\tau}^d$	$[-0.078, +0.098]$
$\varepsilon_{ee}^p - \varepsilon_{\mu\mu}^p$	$[-0.190, +0.927]$	$\oplus[-2.927, -1.814]$	ε_{ee}^p	$[-0.222, +1.801]$
$\varepsilon_{\tau\tau}^p - \varepsilon_{\mu\mu}^p$	$[-0.001, +0.053]$	$[-0.052, +0.053]$	$\varepsilon_{\mu\mu}^p$	$[-0.248, +0.282] \oplus [+0.625, +1.551]$
$\varepsilon_{e\mu}^p$	$[-0.145, +0.058]$	$[-0.145, +0.145]$	$\varepsilon_{\tau\tau}^p$	$[-0.248, +0.281] \oplus [+0.646, +1.548]$
$\varepsilon_{e\tau}^p$	$[-0.238, +0.292]$	$[-0.292, +0.292]$	$\varepsilon_{e\mu}^p$	$[-0.145, +0.058]$
$\varepsilon_{\mu\tau}^p$	$[-0.019, +0.021]$	$[-0.021, +0.021]$	$\varepsilon_{e\tau}^p$	$[-0.239, +0.293]$
			$\varepsilon_{\mu\tau}^p$	$[-0.019, +0.021]$

Table 1. 2σ allowed ranges for the NSI couplings $\varepsilon_{\alpha\beta}^u$, $\varepsilon_{\alpha\beta}^d$ and $\varepsilon_{\alpha\beta}^p$ as obtained from the global analysis of oscillation data (left column) and also including COHERENT constraints. The results are obtained after marginalizing over oscillation and the other matter potential parameters either within the LMA only and within both LMA and LMA-D subspaces respectively (this second case is denoted as LMA \oplus LMA-D). Notice that once COHERENT data are included the two columns become identical in all cases since for NSI couplings to $f = u, d, p$ the LMA-D solution is only allowed above 95% CL.

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