

Resummed Drell-Yan cross-section at N^3LL

**A.H. Ajjath,^a Goutam Das,^{b,c} M.C. Kumar,^d Pooja Mukherjee,^a V. Ravindran^a
and Kajal Samanta^d**

^a*The Institute of Mathematical Sciences, HBNI,
IV Cross Road, Taramani, Chennai 600113, India*

^b*Theory Group, Deutsches Elektronen-Synchrotron (DESY),
Notkestrasse 85, D-22607 Hamburg, Germany*

^c*Theoretische Physik 1, Naturwissenschaftlich-Technische Fakultät, Universität Siegen,
Walter-Flex-Strasse 3, 57068 Siegen, Germany*

^d*Department of Physics, Indian Institute of Technology Guwahati,
Guwahati-781039, India*

E-mail: ajjathah@imsc.res.in, goutam.das@uni-siegen.de,
mckumar@iitg.ac.in, poojamukherjee@imsc.res.in, ravindra@imsc.res.in,
kajal.samanta@iitg.ac.in

ABSTRACT: We present the resummed predictions for inclusive cross-section for Drell-Yan (DY) production as well as onshell Z, W^\pm productions at next-to-next-to-next-to leading logarithmic (N^3LL) accuracy. Using the standard techniques, we derive the N -dependent coefficients in the Mellin- N space as well as the N -independent constants and match the resummed result through the minimal prescription procedure with the fixed order results. In addition to the standard $\ln N$ exponentiation, we study the numerical impacts of exponentiating N -independent part of the soft function and the complete \bar{g}_0 that appears in the resummed predictions in N space. All the analytical pieces needed in these different approaches are extracted from the soft-virtual part of the inclusive cross section known to next-to-next-to-next-to leading order (N^3LO). We perform a detailed analysis on the scale and parton distribution function (PDF) variations and present predictions for 13 TeV LHC for the neutral Drell-Yan process as well as onshell charged and neutral vector boson productions.

KEYWORDS: Perturbative QCD, Resummation

ARXIV EPRINT: [2001.11377](https://arxiv.org/abs/2001.11377)

Contents

1	Introduction	1
2	Theoretical framework	4
2.1	Drell-Yan and Z, W^\pm production	4
2.2	Soft-virtual cross-section	5
2.3	Threshold resummation	9
3	Numerical results	14
3.1	Soft-virtual correction for neutral DY invariant mass	14
3.2	Resummed prediction for neutral DY invariant mass	16
3.3	Resummed prediction for Z/W^\pm productions	21
4	Conclusions	23
A	Soft-virtual coefficient in N-space	24
B	Resummed coefficients	26
B.1	Resummation ingredients for the Standard \bar{N} exponentiation	26
B.2	Resummation ingredients for the Standard N exponentiation	29
B.3	Resummation ingredients for the Soft exponentiation	35
B.4	Resummation ingredients for the All exponentiation	38
C	Anomalous dimensions	38

1 Introduction

Standard Model(SM) has been very successful so far in describing the physics of elementary particles. Precision study has played an important role in establishing the SM through the latest discovery of Higgs boson at the Large Hadron Collider (LHC). The properties of the Higgs boson is being studied with higher accuracy. Recent observations at the LHC demonstrate that the systematic precision study is essential to look for any deviation from the SM in search of new physics beyond the SM (BSM). While there is no promising sign of new physics signature so far at the LHC, it is extremely important to know the SM predictions for the standard processes like Higgs and DY or Z, W^\pm productions to utmost accuracy. Not only this could help in BSM searches but also help to understand the perturbative structure of the underlying gauge theory.

Drell-Yan production has been a standard candle at the hadron colliders and is extremely important for luminosity monitoring. This is one of the hadronic processes which is well understood theoretically. For example, next to next to leading order (NNLO) quantum

chromodynamics (QCD) correction [1–3] to this process was computed three decades ago. DY is also an important process experimentally for several BSM searches. Experimentally, one has a very clean environment for precise measurements in terms of the kinematics of the final state lepton pairs. Higher order perturbative QCD corrections to DY provides ample opportunity to explore the structure of the perturbation series. Thus DY serves as an important process in collider experiments. At the LHC, the strong interaction dynamics dominates over the others and hence There have been attempts to go beyond NNLO accuracy in order to improve the precision from the theoretical side.

The calculation of complete N³LO cross-section is extremely difficult due to increasing number of subprocesses involved, however there have been significant progress to obtain third order contribution to this process in QCD. Very recently the first result at complete N³LO from only virtual photon mediator has been calculated in [4]. From the theory side, DY is seen to be extremely stable with respect to factorization and renormalisation scales already at NNLO. The scale uncertainty has been seen to be reduced to 2% for a canonical variation of factorization and renormalisation scales compared to NLO where uncertainty is about 9.2%, whereas the K-factor seem to improve marginally from 1.25 at NLO to 1.28 at NNLO. However keeping in mind the importance of this process, it is worth studying the results from next orders and devise methods to incorporate more and more higher order corrections. Since a complete calculation beyond NNLO level is difficult, the soft-virtual (SV) contributions is often computed as first step. In addition, the later constitutes a significant part of the cross-section in the region where the partonic scaling variable $z \rightarrow 1$, called the threshold region. The SV cross-sections are known for many SM processes e.g. Higgs production [5–11], associated production [12], bottom quark annihilation [13], pseudo-scalar Higgs production [14]. For DY production, using the three loop quark form factor [15], exploring the universal structure of the soft part [16] of SV cross section to Higgs production [5], the dominant soft-virtual (SV) corrections for DY at third order was obtained [17] and later it was confirmed in [18]. It is also worth noting that the same result was independently obtained in [19] using soft-gluon resummation.

The SV contributions dominate at every order in the perturbation theory through large logarithms spoiling the reliability of the fixed order predictions. The resolution to this is to resum these large logarithms to all orders. Resummation of these large logarithms is thus very important to correctly describe the cross section in the threshold region. In [20–22], a systematic approach was proposed to resum these logarithms to all orders. The large logarithms arise in the hard partonic cross section when the total available center-of-mass energy (\hat{s}) becomes close to the invariant mass (Q) of the final state, in other words the partonic scaling variable $z = Q^2/\hat{s} \rightarrow 1$. This results from the soft gluon emissions, as a consequence of which the cross-section is enhanced by the large logarithms that appear as distributions namely Dirac delta $\delta(1 - z)$ and $+$ distributions:

$$\mathcal{D}_j(z) = \left(\frac{\log^j(1 - z)}{1 - z} \right)_+ \tag{1.1}$$

In Mellin space these singular terms are transformed into powers of logarithms of the Mellin variable N . In Mellin N space, these contributions can be systematically resummed

to all orders and they display a nontrivial pattern of exponentiation. In the threshold region, the fixed order predictions often fail to describe the cross-section well and hence the resummation of these large logarithms becomes very important to correctly describe the region. Moreover, it has been very well established that the resummed contributions give a sizable contributions to the cross-section. In fact many SM fixed order calculations have been improved with the corresponding resummed results, for example, the inclusive scalar Higgs production in gluon fusion [6, 19, 23–25] (see also [26] for renormalisation group improved prediction) as well as in bottom quark annihilation [27], deep inelastic scattering [28, 29], DY production [6, 19, 30], pseudo-scalar Higgs production [31–33], spin-2 production [34, 35] etc. Threshold resummation not only improves the inclusive fixed order results but also differential observables like rapidity [21, 36–39] and in the context of LHC precision measurements, it is important to include these corrections and they are shown to improve the fixed order results.

In the resummed predictions for the cross-sections, there is an intrinsic ambiguity on what is exponentiated and what is not. In the standard approach, one exponentiates only large- N pieces coming from the soft function which are enhanced in the threshold region. However one can also define large logarithms in terms of a new variable $\bar{N} = N \exp(\gamma_E)$, γ_E being the *Euler-Mascheroni* (E-M) constant. Theoretically this is allowed, since γ_E arises as a mathematical artifact due to dimensional regularization in d -space time dimensions. Moreover, this does not spoil the fact that the large- N pieces are exponentiated in the threshold region. In this terminology, one exponentiates \bar{N} instead of N . Numerically, however this makes a difference already at the leading logarithmic accuracy. It has been already seen in [29] that the perturbative convergence is improved if one exponentiates the large- \bar{N} terms. Apart from the standard threshold exponentiation, one can in fact exponentiate the complete soft function i.e. all the large- \bar{N} terms as well as the $\delta(1-z)$ terms arising from the soft function. We call this ‘**Soft exponentiation**’ which renders some part of the N -independent constant (\bar{g}_0) for the **Standard \bar{N} exponentiation**. In addition to these procedures, one can also exponentiate the complete form factor along with the soft function. This was studied in the context of the SM Higgs production [24, 25] and was shown to improve the scale uncertainty better than the standard threshold approach. The form factor is process dependent and therefore is non-universal unlike the soft function. However, the form factor as well as the soft function both satisfy the similar Sudakov K+G type equation [8, 9, 40–43]. Hence the solution to K+G equation for the form factor is an exponential N independent constant justifying the exponentiation. The numerical impact of this has already been studied in the past for the DY production in [44] where the authors show that both in DIS scheme and in $\overline{\text{MS}}$ scheme the complete form factor exponentiates to the orders currently known.

The goal of the present article is to study the effect of threshold logarithms at N^3LL accuracy and match it to the known NNLO results. We perform this study for the neutral DY production as well as for onshell Z and W^\pm productions. The paper is organized as follows. In section 2.1, we collect the useful formulae required for the invariant mass distribution for DY and total cross-section for Z, W^\pm productions at the LHC. Next we discuss the theoretical set up in the context of resummation. Here we describe in detail the

factorization of soft-virtual coefficient in section 2.2. Next we set up in section 2.3 different resummation prescriptions as well as derive some useful formulas needed. Section 3 we study in detail the effect of threshold logarithms for different prescriptions and present our results along with the estimation of uncertainties. We finally conclude in section 4.

2 Theoretical framework

2.1 Drell-Yan and Z, W^\pm production

The hadronic cross-section for DY or onshell Z, W^\pm production at the LHC can be written as

$$\sigma = \sigma^{(0)} \sum_{ab=q,\bar{q},g} \int_0^1 dx_1 \int_0^1 dx_2 f_a(x_1, \mu_f^2) f_b(x_2, \mu_f^2) \int_0^1 dz \Delta_{ab}(z, Q^2, \mu_f^2) \delta(\tau - zx_1x_2), \quad (2.1)$$

where $\sigma = \frac{d\sigma}{dQ}(\tau, Q^2)$ for DY production, with Q being the invariant mass of the di-lepton pair. Here $f_a(x_1, \mu_f^2)$ and $f_b(x_2, \mu_f^2)$ are the non-perturbative parton distribution functions (PDFs) of the partons a, b carrying momentum fractions x_1, x_2 of the incoming protons at the factorization scale μ_f . These PDFs are appropriately convoluted with perturbatively calculable partonic coefficients $\Delta_{ab}(z, Q^2, \mu_f^2)$. For the on-shell Z, W^\pm production, $\sigma = \sigma_V, V = Z/W^\pm$ and $q = M_V$, the mass of the vector boson. The partonic coefficients are obtain from the partonic cross section using perturbation theory. For the DY production V we include contributions from γ and Z as well as their interference.

The partonic cross-section can be decomposed as

$$\Delta_{ab}(z, Q^2, \mu_f^2) = \Delta_{ab}^{(sv)}(z, Q^2, \mu_f^2) + \Delta_{ab}^{(reg)}(z, Q^2, \mu_f^2). \quad (2.2)$$

The first term $\Delta^{(sv)}$ is called the SV partonic coefficient and it contains distributions such as $\delta(1-z)$ and \mathcal{D}_+ , whereas the second term $\Delta^{(reg)}$ contains those terms that are regular in the scaling variable z . The prefactors for DY and Z, W^\pm production are given as below:

$$\begin{aligned} \sigma_{DY}^{(0)} &= \frac{2\pi}{n_c} \left[\frac{Q}{S} \mathcal{F}^{(0)} \right], \\ \sigma_Z^{(0)} &= \frac{2\pi}{n_c} \left[\frac{\pi\alpha}{8s_w^2 c_w^2 S} \right], \\ \sigma_{W^\pm}^{(0)} &= \frac{2\pi}{n_c} \left[\frac{\pi\alpha}{4s_w^2 S} \right], \end{aligned} \quad (2.3)$$

where S is the hadronic centre-of-mass energy and $n_c = 3$ in QCD. For DY production, the factor $\mathcal{F}^{(0)}$ is found to be,

$$\begin{aligned} \mathcal{F}^{(0)} &= \frac{4\alpha^2}{3Q^2} \left[Q_q^2 - \frac{2Q^2(Q^2 - M_Z^2)}{((Q^2 - M_Z^2)^2 + M_Z^2 \Gamma_Z^2)} c_w^2 s_w^2 Q_q g_e^V g_q^V \right. \\ &\quad \left. + \frac{Q^4}{((Q^2 - M_Z^2)^2 + M_Z^2 \Gamma_Z^2)} c_w^4 s_w^4 \left((g_e^V)^2 + (g_e^A)^2 \right) \left((g_q^V)^2 + (g_q^A)^2 \right) \right]. \end{aligned} \quad (2.4)$$

Here α is the fine structure constant, c_w, s_w are sine and cosine of Weinberg angle respectively. M_Z and Γ_Z are the mass and the decay width of the Z -boson.

$$g_a^A = -\frac{1}{2}T_a^3, \quad g_a^V = \frac{1}{2}T_a^3 - s_w^2 Q_a, \quad (2.5)$$

Q_a being electric charge and T_a^3 is the weak isospin of the electron or quarks.

In the threshold region, the SV terms which consist of distributions contribute significantly at the hadron level. After mass factorization, the partonic coefficient in the threshold region experiences further factorization in terms of the form factor and soft-collinear function. In the next section we will discuss in detail on the structure of distributions in the SV coefficient which will form the basis for the resummation.

2.2 Soft-virtual cross-section

In the following, we briefly describe the theoretical set up that is required to study the impact of threshold corrections within the framework of resummation *à la* Sterman, Catani and Trentadue [20, 21]. We do this in order to understand the role of various pieces that contribute to the resummed result. Exploiting the factorization of infrared sensitive contributions and gauge and renormalisation group invariances, inclusive cross section for the DY and on-shell Z/W^\pm productions in the threshold limit can be expressed in terms of form factor of the neutral/charged current, soft distribution function and Altarelli-Parisi kernels (see [8, 9]). The resulting expression expressed in z space is free of both ultraviolet and infrared divergences and captures the distributions \mathcal{D}_j with given logarithmic accuracy to all orders in perturbation theory. In the Mellin N space, we can achieve the same and in addition, one has the advantage to reorganize the series in such a way that order one contributions of the form $a_s \beta_0 \log(N)$ can be resummed systematically to all orders in the large N limit. Here, a_s is defined by $a_s = g_s^2(\mu_r^2)/16\pi^2$ with g_s begin the strong coupling constant and μ_r the renormalisation scale and β_0 is the first coefficient of the coupling constant beta function. Note that in Mellin N space, the convolutions in z space become simple products. The z space result can be used to compute the soft-virtual contributions in power series expansion of strong coupling constant a_s .

In $d = 4 + \epsilon$ space time dimensions, the threshold enhanced partonic soft-virtual cross-section to all orders in perturbation theory in z space can be written [8, 9] as

$$\Delta^{(sv)}(z, Q^2, \mu_r^2, \mu_f^2) = \mathcal{C} \exp\left(\Psi(z, Q^2, \mu_r^2, \mu_f^2, \epsilon)\right) \Big|_{\epsilon=0}. \quad (2.6)$$

Here Ψ is a distribution function which is finite in the limit $\epsilon \rightarrow 0$. The symbol \mathcal{C} denotes the Mellin convolution (denoted below as \otimes) which in the above expression should be treated as

$$\mathcal{C} \exp\left(f(z)\right) = \delta(1-z) + \frac{1}{1!}f(z) + \frac{1}{2!}f(z) \otimes f(z) + \dots, \quad (2.7)$$

with $f(z)$ being a function containing only $\delta(1-z)$ and plus distributions. The finite exponent in the above is re-factorized in the threshold limit and gets contribution from the form factor $(\hat{\mathcal{F}}(\hat{a}_s, q^2, \mu^2, \epsilon))$ with $q^2 = -Q^2$, soft-collinear function $(\Phi(\hat{a}_s, z, Q^2, \mu^2, \epsilon))$

(later called as soft function) as well as mass factorization kernels ($\Gamma(\hat{a}_s, z, \mu_f^2, \mu^2, \epsilon)$) and takes the following form in dimensional regularization:

$$\Psi\left(z, Q^2, \mu_r^2, \mu_f^2, \epsilon\right) = \left(\ln\left[\mathcal{Z}(\hat{a}_s, \mu_r^2, \mu^2, \epsilon)\right]^2 + \ln\left|\hat{\mathcal{F}}(\hat{a}_s, q^2, \mu^2, \epsilon)\right|^2\right)\delta(1-z) + 2\Phi(\hat{a}_s, z, Q^2, \mu^2, \epsilon) - 2\mathcal{C}\ln\Gamma(\hat{a}_s, z, \mu_f^2, \mu^2, \epsilon). \quad (2.8)$$

μ keeps the strong coupling (\hat{a}_s) dimensionless in the $d = 4 + \epsilon$ dimensions. $\mathcal{Z}(\hat{a}_s, \mu_r^2, \mu^2, \epsilon)$ denotes the overall UV renormalization constant which for the processes considered here is unity due to conserved current.

The bare quark form factor satisfies the Sudakov K+G equation [8, 9, 40–43] which follows as a consequence of the gauge invariance as well as renormalisation group invariance,

$$\frac{d\ln\hat{\mathcal{F}}}{d\ln q^2} = \frac{1}{2}\left[\mathcal{K}(\hat{a}_s, \frac{\mu_r^2}{\mu^2}, \epsilon) + \mathcal{G}(\hat{a}_s, \frac{q^2}{\mu_r^2}, \frac{\mu_r^2}{\mu^2}, \epsilon)\right]. \quad (2.9)$$

The function \mathcal{K} contains all the infrared poles in ϵ whereas the function \mathcal{G} is finite in the limit $\epsilon \rightarrow 0$. The renormalisation group invariance leads to the following solutions of these functions in terms of cusp anomalous dimensions (A):

$$\frac{d\mathcal{K}}{d\ln\mu_r^2} = -\frac{d\mathcal{G}}{d\ln\mu_r^2} = A(a_s(\mu_r)) = \sum_{i=1}^{\infty} a_s^i(\mu_r) A_i. \quad (2.10)$$

The cusp anomalous dimensions are known to fourth order [45–47, 47–49, 49–57, 57] and are collected in appendix C. The μ_r independent piece of the \mathcal{G} can be written in perturbative series as

$$\mathcal{G}(a_s(q), \epsilon) = \sum_{j=1}^{\infty} a_s^j(q) \mathcal{G}_j(\epsilon), \quad (2.11)$$

where the coefficients $\mathcal{G}^{(j)}(\epsilon)$ can be decomposed as

$$\mathcal{G}_i(\epsilon) = 2B_i + f_i + C_i + \sum_{k=1}^{\infty} \epsilon^k G_{ik}, \quad (2.12)$$

where

$$\begin{aligned} C_1 &= 0, \\ C_2 &= -2\beta_0 G_{11}, \\ C_3 &= -2\beta_1 G_{11} - 2\beta_0(G_{21} + 2\beta_0 G_{12}). \end{aligned} \quad (2.13)$$

The coefficients G_{ik} are the finite coefficients found in terms of QCD color factors and can be extracted from explicit calculation of quark form factor. Note that up to the third order one also needs coefficients G_{22}, G_{31} and thereby needs the three-loop calculation of the form factor [15]. We have collected them in the appendix C. Similar to the cusp anomalous

dimension, the coefficients f_i have been found to be maximally non-abelian to third order in strong coupling i.e. they satisfy

$$f_i^g = \frac{C_F}{C_A} f_i^q. \quad (2.14)$$

The initial state collinear singularities are removed using the Altarelli-Parisi (AP) splitting kernels $\Gamma(\hat{a}_s, \mu_f^2, \mu^2, z, \epsilon)$. They satisfy the well-known DGLAP evolution given as,

$$\frac{d\Gamma(z, \mu_f^2, \epsilon)}{d \ln \mu_f^2} = \frac{1}{2} P(z, \mu_f^2) \otimes \Gamma(z, \mu_f^2, \epsilon), \quad (2.15)$$

where $P(z, \mu_f^2)$ is the AP splitting functions. The perturbative expansion for these splitting functions has the following form:

$$P(z, \mu_f^2) = \sum_{i=0}^{\infty} a_s^{i+1}(\mu_f) P^{(i)}(z). \quad (2.16)$$

As already discussed, only the $q\bar{q}$ channel contributes to the SV cross-section and thus we find that, only the diagonal terms of the splitting functions contribute to the SV cross-section. The diagonal part of the splitting functions is known to contain the $\delta(1-z)$ and distributions and can be written as,

$$P_{II}^{(i)} = 2 \left[B_{i+1} \delta(1-z) + A_{i+1} \mathcal{D}_0 \right] + P_{II}^{(reg,i)}(z). \quad (2.17)$$

The splitting functions are known exactly to four loops [45, 58–60].

The finiteness of the soft-virtual cross-section demands that the soft-collinear function Φ will also satisfy similar Sudakov type equation like the form factor i.e. one can write

$$\frac{d\Phi}{d \ln Q^2} = \frac{1}{2} \left[\bar{\mathcal{K}} \left(\hat{a}_s, z, \frac{\mu_r^2}{\mu^2}, \epsilon \right) + \bar{\mathcal{G}} \left(\hat{a}_s, z, \frac{Q^2}{\mu_r^2}, \frac{\mu_r^2}{\mu^2}, \epsilon \right) \right], \quad (2.18)$$

where $\bar{\mathcal{K}}(\hat{a}_s, z, \frac{\mu_r^2}{\mu^2}, \epsilon)$ contains all the poles and $\bar{\mathcal{G}}(\hat{a}_s, z, \frac{Q^2}{\mu_r^2}, \frac{\mu_r^2}{\mu^2}, \epsilon)$ is finite in the dimensional regularization such that Ψ becomes finite as $\epsilon \rightarrow 0$. The solution to the above equation has been found [8, 9] to be

$$\Phi = \sum_{j=1}^{\infty} \hat{a}_s^j \frac{j\epsilon}{1-z} \left(\frac{Q^2(1-z)^2}{\mu^2} \right)^{j\epsilon/2} \mathcal{S}_\epsilon^j \hat{\Phi}^{(j)}(\epsilon). \quad (2.19)$$

$\hat{\Phi}^{(j)}$ can be found from the solution of the form factor by the replacement as $A \rightarrow -A, \mathcal{G}(\epsilon) \rightarrow \bar{\mathcal{G}}(\epsilon)$. Notice that $\bar{\mathcal{G}}(\epsilon)$ are now new finite z -independent coefficients coming from the soft function whereas the z dependence has been taken out in eq. (2.19). This can be found by comparing the poles and non-pole terms in $\hat{\Phi}^{(j)}$ with those coming from the form factors, overall renormalisation constants, splitting kernel and the lower order SV terms. The coefficient $\bar{\mathcal{G}}$ has same structure as the form factor in eq. (2.12) after setting $f_i \rightarrow -f_i, B_i \rightarrow 0, \gamma_i \rightarrow 0$,

$$\bar{\mathcal{G}}_i = -f_i + \tilde{C}_i + \sum_{k=1}^{\infty} \epsilon^k \tilde{G}_{ik}, \quad (2.20)$$

where

$$\begin{aligned}\tilde{C}_1 &= 0, \\ \tilde{C}_2 &= -2\beta_0\tilde{G}_{11}, \\ \tilde{C}_3 &= -2\beta_1\tilde{G}_{11} - 2\beta_0(\tilde{G}_{21} + 2\beta_0\tilde{G}_{12}).\end{aligned}\tag{2.21}$$

The coefficients f_i are same as those appear in the quark form factor in eq. (2.12). The coefficients \tilde{G}_{ij} required up to three loops have been extracted in [61] and also collected in the appendix C. Note that one has to perform the following expansion in eq. (2.19) in order to get all the distributions and delta function coming from the soft function,

$$\frac{1}{(1-z)}\left[(1-z)^2\right]^{j\epsilon/2} = \frac{1}{j\epsilon}\delta(1-z) + \sum_{k=0}^{\infty} \frac{(j\epsilon)^k}{k!} \mathcal{D}_k.\tag{2.22}$$

It is worth noting that $\bar{\mathcal{G}}$ as well as the complete soft function Φ_I satisfy the maximally non-abelian property up to three loops. Moreover Φ_I is also universal in the sense that it only depends on the initial legs and is completely unaware of the color neutral final state. Expanding $\Delta^{(sv)}$ in powers of a_s as

$$\Delta_{ab}^{(sv)} = \delta_{ab} \sum_{i=0}^{\infty} a_s^i \Delta^{(i)},\tag{2.23}$$

with the born contribution being $\Delta^{(0)} = \delta(1-z)$. The SV correction at the three loops are known [17] which we collect here for completeness in the appendix A.

In the following, we will study the numerical impact of resummed result resulting from $\Delta_{ab}^{(sv)}$ after performing the Mellin transformation in the large N limit. We start with Ψ which is finite while the individual contributions to it contain UV and IR singularities. Decomposing the later ones as sum of singular and finite parts as

$$\begin{aligned}\ln|\hat{\mathcal{F}}|^2(Q^2) &= \mathcal{L}_{\mathcal{F}}^{\text{sing}}(Q^2, \mu_r^2) + \mathcal{L}_{\mathcal{F}}^{\text{fin}}(Q^2, \mu_r^2) \\ \Phi(z, Q^2) &= \Phi^{\text{sing}}(z, Q^2, \mu_f^2, \mu_r^2) + \Phi_{\mathcal{D}}^{\text{fin}}(z, Q^2, \mu_f^2, \mu_r^2) + \Phi_{\delta}^{\text{fin}}(Q^2, \mu_f^2, \mu_r^2)\delta(1-z) \\ \mathcal{C} \ln \Gamma(z, \mu_f^2) &= \mathcal{L}_{\Gamma}^{\text{sing}}(z, \mu_f^2, \mu_r^2) + \mathcal{L}_{\Gamma\mathcal{D}_0}^{\text{fin}}(\mu_f^2, \mu_r^2)\mathcal{D}_0 + \mathcal{L}_{\Gamma\delta}^{\text{fin}}(\mu_f^2, \mu_r^2)\delta(1-z)\end{aligned}\tag{2.24}$$

Substituting the above equations in eq. (2.8), we can easily show that all the singular terms in the limit $\epsilon \rightarrow 0$ cancel among themselves. In addition, \mathcal{D}_0 terms in finite part of $\mathcal{C} \ln \Gamma$ go away when added to $\Phi_{\mathcal{D}}^{\text{fin}}$ resulting in a finite distribution. Substituting the Ψ in eq. (2.6), we obtain

$$\Delta^{\text{sv}}(z, Q^2) = C_0(Q^2) \otimes \mathcal{C}e^{G_+(z, Q^2)}\tag{2.25}$$

where (supressing dependence on μ_f and μ_r), the N independent constant C_0 is given by

$$C_0(z, Q^2) = \exp\left(\mathcal{L}_{\mathcal{F}}^{\text{fin}} + 2\Phi_{\delta}^{\text{fin}} - 2\mathcal{L}_{\Gamma\delta}^{\text{fin}}\right)\delta(1-z)\tag{2.26}$$

$$G_+(z, Q^2) = \left(\frac{1}{1-z} \left[\int_{\mu_f^2}^{Q^2(1-z)^2} \frac{d\mu^2}{\mu^2} 2 A(a_s(\mu^2)) + D(a_s(Q^2(1-z)^2)) \right] \right)_+\tag{2.27}$$

and D in G_+ is related to $\bar{\mathcal{G}}$ by $D = 2\bar{\mathcal{G}}$ and $\bar{\mathcal{G}}(a_s(Q^2(1-z)^2))$ is $\bar{\mathcal{G}}$ in eq. (2.18) evaluated at $\mu_r^2 = \mu_f^2 = Q^2$.

So far, we showed how various collinear soft gluon emissions as well as the wide angle soft emissions can be systematically summed to all orders to obtain eq. (2.25) in the z space when partonic variable $z \rightarrow 1$. Note that C_0 is obtained by first collecting those terms that are proportional to $\delta(1-z)$ terms of Ψ and then expanding the exponential of them in powers of a_s . The remaining function G_+ contains only distributions \mathcal{D}_j . Hence, one can predict the following structure for G_+ :

$$G_+(z, Q^2) = G_1(Q^2) \otimes \mathcal{D}_0 + G_2(z, Q^2) + a_s G_3(z, Q^2) + \dots \quad (2.28)$$

where each G_1 sums certain terms of the $a_s^i \mathcal{D}_{i-1}$ to all orders, and G_2 sums $a_s^i \mathcal{D}_{i-2}$ terms to all orders etc. The result Δ^{sv} expressed in terms of C_0 and the exponential of G_+ using eq. (2.28) systematically sums the distributions \mathcal{D}_j to all orders and hence can predict these distributions to all orders provided A and D are known to desired accuracy in a_s . For example, knowing A_1 , we can predict all the terms $a_s^i \mathcal{D}_i$ with $i = 1, 2, \dots, \infty$ in Φ , similarly given A_1 and D_1 , we can predict $a_s^i \mathcal{D}_{i-1}$ with $i = 1, 2, \dots, \infty$ etc. Hence, expression given in eq. (2.25) has the predictive power for Δ^{sv} to all orders in a_s given the logarithmic accuracy in z space, quantified by terms of the form $a_s^i \mathcal{D}_j$. Note that when the exponential of Φ is expanded using convolution rules given in eq. (2.7), we will get not only \mathcal{D}_j but also $\delta(1-z)$. In other words, $\delta(1-z)$ terms in Δ^{sv} can come from both $\exp(G_+)$ as well as C_0 .

Often in certain kinematic regions, these contributions can be enhanced when convoluted with the parton distribution functions spoiling the reliability of the perturbation theory. Hence we need to include these potentially large terms to all orders in perturbation theory for any sensible predictions. Such an exercise in the z space is technically challenging due to the complexity involved in computing the convolutions of \mathcal{D}_j . However, in the Mellin N space, the convolutions become simple products allowing us to study the impact of these large logarithms to all orders in a systematic fashion. In the following, we will describe how this can be done in Mellin N space.

2.3 Threshold resummation

In the last sub-section, we showed that threshold effects for partonic coefficients can be obtained near threshold as a product of well-defined functions, each organizing a class of infrared and collinear enhancements as can be seen from eq. (2.8). This re-factorization is valid up to corrections which are nonsingular at threshold when partonic $z \rightarrow 1$. While the z space result captures the entire underlying infrared dynamics in the threshold limit, it can be better described in the Mellin- N space where the threshold limit $z \rightarrow 1$ translates into $N \rightarrow \infty$. We found that the form in eq. (2.8) was already suitable for all order study, however complications arise in performing the convolution. On the other hand any such convolution becomes simple product in Mellin space and all the distributions coming from the soft function are thus translated into large logarithms in Mellin N .

Following [21], the resummed partonic SV coefficient function can be organized as follows:

$$\begin{aligned}\hat{\sigma}_N(Q^2) &= \int_0^1 dz z^{N-1} \Delta^{\text{sv}}(z, Q^2) \\ &= \bar{g}_0(Q^2) \exp\left(G_{\bar{N}}(Q^2)\right),\end{aligned}\tag{2.29}$$

where $G_{\bar{N}}$ is obtained by computing the large N limit of Mellin moment of G_+ and then by decomposing as

$$\lim_{N \rightarrow \infty} \int_0^1 dz z^{N-1} G_+(z, Q^2) = \bar{G}_0(Q^2) + G_{\bar{N}}(Q^2), \quad \text{with } G_{\bar{N}}(Q^2)|_{\bar{N}=1} = 0 \tag{2.30}$$

where $\bar{N} = N \exp(\gamma_E)$ and γ_E is E-M constant. The N independent constant \bar{g}_0 is given by

$$\bar{g}_0(Q^2) = \exp\left(\mathcal{L}_{\mathcal{F}}^{\text{fin}} + 2\Phi_{\delta}^{\text{fin}} - 2\mathcal{L}_{\Gamma\delta}^{\text{fin}} + \bar{G}_0(Q^2)\right) \tag{2.31}$$

$G_{\bar{N}}$ is function of the universal coefficients A which are known to fourth order and D known to third order in a_s . $G_{\bar{N}}$ collects and resums all the large- N logarithms to all orders and it can be expressed as a resummed perturbative series which takes the following form:

$$G_{\bar{N}}(Q^2) = \ln \bar{N} \bar{g}_1(\bar{N}, Q^2) + \bar{g}_2(\bar{N}, Q^2) + a_s \bar{g}_3(\bar{N}, Q^2) + a_s^2 \bar{g}_4(\bar{N}, Q^2) + \dots \tag{2.32}$$

Following [19, 23], we computed \bar{g}_i up to $i = 4$ (for \bar{g}_i up to $i = 3$, see [30]) and they are given in appendix B.1. Note that \bar{g}_i coefficients are universal in the sense that it depends only on whether the born process is $q\bar{q}$ channel or gg channel. In the Mellin N space, the $\delta(1-z)$ in z -space directly translates into N independent piece whereas the plus-distributions give rise to the $\ln(\bar{N})$ as well as N independent constants in the large N limit. Part of these constant pieces, namely \bar{G}_0 , is absorbed into the coefficients \bar{g}_0 in the standard resummation approach. Hence, \bar{g}_0 contains only N independent pieces which come from the form factor, soft distribution function, AP kernels and N independent part of the Mellin moment of $G_+(z, Q^2)$. The condition $G_{\bar{N}} = 0$ for $\bar{N} = 1$ allows the constants \bar{g}_i to contain N independent terms. Note that the expressions for \bar{g}_0 and \bar{g}_i obtained this way depend on the condition $G_{\bar{N}} = 0$ for $\bar{N} = 1$. In other words, there is an ambiguity in treating the N independent terms in the resummed results. Exploiting this, in [21], the N independent constants were defined by demanding $G_{\bar{N}} = 1$ when $\bar{N} = 1$. With this, \bar{g}_0 has the following perturbative expansion:

$$\bar{g}_0(Q^2) = 1 + \sum_{n=1}^{\infty} a_s^n \bar{g}_{0n}(Q^2). \tag{2.33}$$

The successive terms in the resummed exponent eq. (2.32) along with the corresponding terms in eq. (2.33) define the accuracies leading logarithmic (LL), next to LL (NLL), NNLL and N³LL etc. Terms independent of N can be treated, in principle, by the same methods that resum terms enhanced by logarithms of N .

In summary, the resummed result will differ depending on how we treat the N -independent constants. We define various schemes that differentiate how these constants

are treated in our numerical implementation for the phenomenological studies. This allows us to investigate numerical impact of the various resummed results in detail.

- **Standard \bar{N} exponentiation.** This is the case we have discussed so far where we define large logarithms are functions of $\bar{N} = N \exp(\gamma_E)$, where γ_E is E-M constant. The N dependent functions $G_{\bar{N}}$ in this case can be computed by simply performing the Mellin moment of $G_+(z, Q^2)$ in the large \bar{N} limit and keeping only those terms that vanish when $\bar{N} = 1$.
- **Standard N exponentiation.** This approach differs from the previous one in the definition of large- N variable. In this case the large logarithm is simply $\ln N$ and these terms are exponentiated to all orders through the resummed exponent. It is evident that this only accounts for reshuffling of γ_E between \bar{g}_0 and $G_{\bar{N}}$ in eq. (2.29) which now takes the following form:

$$\hat{\sigma}_N(Q^2) = g_0(Q^2) \exp\left(G_N(Q^2)\right). \quad (2.34)$$

The resummed exponent G_N also takes a different form compared to the standard \bar{N} exponent,

$$G_N(Q^2) = \ln N g_1(N, Q^2) + g_2(N, Q^2) + a_s g_3(N, Q^2) + a_s^2 g_4(N, Q^2) + \dots \quad (2.35)$$

The resummed coefficients g_i in the above equation which defines the resummed accuracy, differs from \bar{g}_i in eq. (2.32). The present scheme is defined by demanding $G_N = 0$ when $N = 1$.

$$\lim_{N \rightarrow \infty} \int_0^1 dz z^{N-1} G_+(z, Q^2) = G_0(Q^2) + G_N(Q^2), \quad \text{with } G_N(Q^2)|_{N=1} = 0 \quad (2.36)$$

With this definition, the rest of the N independent terms from the Mellin moment of G_+ is combined with finite parts of form factor, soft distribution function and the AP kernels as The N independent constant g_0 is given by

$$g_0(Q^2) = \exp\left(\mathcal{L}_{\mathcal{F}}^{\text{fin}} + 2\Phi_{\delta}^{\text{fin}} - 2\mathcal{L}_{\Gamma\delta}^{\text{fin}} + G_0(Q^2)\right) \quad (2.37)$$

and the above result is expanded in powers of a_s :

$$g_0(Q^2) = 1 + \sum_{n=1}^{\infty} a_s^n g_{0n}(Q^2). \quad (2.38)$$

Numerically this can make a difference and it was seen in the context of DIS previously. In case of DY also we find such differences which will be discussed in the next section. Up to N³LL accuracy, the resummed exponents $g_i, i = 1, \dots, 4$ for both quark as well as for gluon initiated process in N exponentiation scheme can be found in [23, 28] and we computed the results for the g_{0i} coefficients up to $i = 3$ which are listed in appendix B.2 along with g_i .

- **Soft exponentiation.** In the standard \bar{N} (N) exponentiation, one exponentiates $\ln \bar{N}$ ($\ln N$) and certain \bar{N} (N) independent terms which arise from G_+ , subjected to the condition $G_{\bar{N}} = 0$ ($G_N = 0$) when $\bar{N} = 1$ ($N = 1$). The remaining \bar{N} (N) independent terms in the Mellin moment of G_+ along with C_0 give the coefficient \bar{g}_0 (g_0). In principle, we can define a scheme wherein entire \bar{N} (N) independent terms of G_+ can be kept in the exponent. More specifically, we define the scheme (relaxing $G_{\bar{N}} = 1$ ($G_N = 0$) for $\bar{N} = 1$ ($N = 1$)) wherein we exponentiate all the terms coming from the finite part of soft distribution function and those from the AP kernels. That is, the exponential contains

$$G_{\bar{N}}^{\text{Soft}} = G_{\bar{N}} + 2\Psi_{\delta}^{\text{fin}} - 2\mathcal{L}_{\Gamma\delta}^{\text{fin}} \quad (2.39)$$

that is,

$$\hat{\sigma}_N(Q^2) = g_0^{\text{Soft}}(Q^2) \exp\left(G_{\bar{N}}^{\text{Soft}}(Q^2)\right). \quad (2.40)$$

with

$$G_{\bar{N}}^{\text{Soft}}(Q^2) = \ln \bar{N} g_1^{\text{Soft}}(\bar{N}, Q^2) + g_2^{\text{Soft}}(\bar{N}, Q^2) + a_s g_3^{\text{Soft}}(\bar{N}, Q^2) + a_s^2 g_4^{\text{Soft}}(\bar{N}, Q^2) + \dots. \quad (2.41)$$

The remaining N independent terms define g_0^{Soft} that is obtained by expanding $\exp(\mathcal{L}_{\mathcal{F}\delta}^{\text{fin}})$ in power series expansion in a_s :

$$g_0^{\text{Soft}}(Q^2) = 1 + \sum_{n=1}^{\infty} a_s^n g_{0n}^{\text{Soft}}(Q^2). \quad (2.42)$$

$G_{\bar{N}}^{\text{Soft}}$ and g_0^{Soft} have similar expansion as eq. (2.32) and eq. (2.33) respectively and the corresponding coefficients are calculated and presented in appendix B.3.

- **All exponentiation.** The soft function and the form factor satisfy K+G type Sudakov integro-differential equations given in eqs. (2.9), and (2.18) and the AP kernels satisfy renormalisation group equation eq. (2.15) governed by AP splitting functions. Hence, their solutions given the boundary conditions demonstrate exponential. The z space solutions that we obtained carry all order information on the distribution \mathcal{D}_j in terms of universal cusp A , soft f and collinear B anomalous dimensions and certain process dependent constants resulting from the form factor. Hence it is natural to study the numerical impact of the entire contribution in the Mellin space without imposing any condition on the N dependent terms. This can be easily achieved and the result for $\hat{\sigma}_N$ takes the following form

$$\hat{\sigma}_N = \exp\left(G_{\bar{N}}^{\text{All}}\right), \quad (2.43)$$

where

$$G_{\bar{N}}^{\text{All}}(Q^2) = \mathcal{L}_{\mathcal{F}}^{\text{fin}}(Q^2) + 2\Phi_{\delta}^{\text{fin}}(Q^2) - 2\mathcal{L}_{\Gamma\delta}^{\text{fin}} + \bar{G}_0(Q^2) + G_{\bar{N}}(Q^2) \quad (2.44)$$

where $G_{\bar{N}}^{\text{All}}$ is expanded as

$$G_{\bar{N}}^{\text{All}}(Q^2) = \ln \bar{N} g_1^{\text{All}}(Q^2) + g_2^{\text{All}}(Q^2) + a_s g_3^{\text{All}}(Q^2) + a_s^2 g_4^{\text{All}}(Q^2). \quad (2.45)$$

The present scheme was already explored in [24, 25] for studying inclusive cross section for the production of Higgs boson at the LHC. For similar study for the DY in DIS and $\overline{\text{MS}}$ schemes, see [44]. Here we will extend it to the N^3LL accuracy. The relevant resummed exponent has been provided in appendix B.4.

Note that a detailed comparison between the N -exponentiation and \bar{N} -exponentiation has been done in [29] for the charge and neutral DIS processes. There, one finds that the \bar{N} -exponentiation shows a faster convergence compared to the N -exponentiation. In fact, the convergence has already been achieved at NLO+NLL order in the threshold region in the case of \bar{N} -exponentiation, whereas in N -exponentiation, this occurs after the NLO+NLL order. Notice that the leading logarithmic term also differs between these two approaches. In the case of \bar{N} -exponentiation, all the γ_E terms are exponentiated through the variable $\bar{N} = N \exp(\gamma_E)$; but in the N -exponentiation these γ_E terms are distributed among the exponent and the N independent term g_0 . As a result the deviation starts already at the LL accuracy. In the next section, we will discuss how various schemes discussed so far can affect the predictions. Note that they all give same result at the LL accuracy, however from NLL they differ. At NNLO level, we have the contributions from all the channels and at N^3LO only SV contribution is known so far. Hence, our numerical predictions will be based on fixed order $\text{N}^3\text{LO}_{\text{sv}}$ results for the parton coefficients and on parton distribution functions known to NNLO accuracy. Note that the resummed result has to be matched to the fixed order result in order to avoid any double counting of threshold logarithms. Hence, the matched result which is usually denoted by N^nLL is computed by taking the difference between the resummed result and the same truncated up to order a_s^n . Hence, it contains contributions from the threshold logarithms to all orders in perturbation theory starting from a_s^{n+1} :

$$\begin{aligned} \sigma_V^{N^n\text{LO}+N^n\text{LL}} &= \sigma_V^{N^n\text{LO}} + \sigma_V^{(0)} \sum_{ab \in \{q, \bar{q}\}} \int_{c-i\infty}^{c+i\infty} \frac{dN}{2\pi i} (\tau)^{-N} \delta_{ab} f_{a,N}(\mu_f^2) f_{b,N}(\mu_f^2) \\ &\times \left(\hat{\sigma}_N \Big|_{N^n\text{LL}} - \hat{\sigma}_N \Big|_{\text{tr}N^n\text{LO}} \right). \end{aligned} \quad (2.46)$$

The Mellin space PDF ($f_{i,N}$) can be evolved using QCD-PEGASUS [62]. Alternatively they can be related to the derivative of z -space PDF as prescribed in [21, 23]. The contour c in the Mellin inverse integration can be chosen according to *Minimal prescription* [63] procedure. Notice that the second term in eq. (2.46) represents the resummed result truncated to N^nLO order, i.e. the same order to which singular SV results are available. In the next section we present the numerical results for the DY production as well as on-shell Z, W^\pm production for LHC where we match the existing N^3LO fixed order SV results with the N^3LL resummation derived in this article.

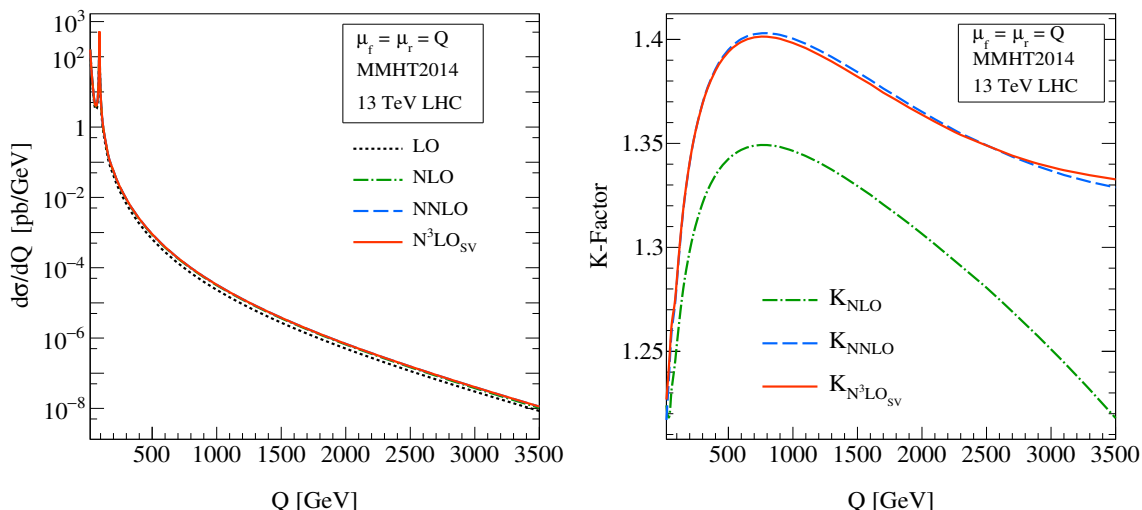


Figure 1. The di-lepton invariant mass distribution (left panel) and the corresponding K -factors (right panel) are presented to N^3LO_{sv} in QCD for 13 TeV LHC.

3 Numerical results

In this section, we present the numerical impact of resummed threshold corrections for neutral DY production as well as on-shell Z/W^\pm production at the LHC. For neutral DY production we consider all the partonic channels at the FO up to NNLO with off-shell γ^* , Z intermediate states. Detailed analysis is done for 13 TeV LHC, however it can be extended to other energies as well as to other colliders.

3.1 Soft-virtual correction for neutral DY invariant mass

We start our discussion by examining the SV corrections at N^3LO . For our numerical study, we use the following electro-weak parameters for the vector boson masses and widths, Weinberg angle (θ_w) and the fine structure constant (α):

$$\begin{aligned}
 m_Z &= 91.1876 \text{ GeV}, & \Gamma_Z &= 2.4952 \text{ GeV}, \\
 m_W &= 80.379 \text{ GeV}, & \sin^2\theta_w &= 0.22343 & \alpha &= 1/128.
 \end{aligned}
 \tag{3.1}$$

We present our results for the default choice of hadronic center of mass energy 13 TeV at the LHC. The parton distribution functions (PDFs) are directly taken from the `lhpdf` [64] routine. The coefficient functions are convoluted with the respective order by order PDFs, however for N^3LO_{sv} results we used NNLO PDFs. Except for studying the PDF uncertainties, we use MMHT2014 [65] parton densities throughout. The $(n+1)$ -loop strong coupling constant is used for computing N^nLO order cross sections with $\alpha_s(m_Z) = 0.120(0.118)$ at NLO(NNLO) respectively. Except for the study of scale uncertainties, the unphysical scale are set equal to the invariant mass of the di-lepton, $\mu_r = \mu_f = Q$.

In figure 1, we present the invariant mass distribution (left panel) of the di-lepton production for the neutral case to N^3LO_{sv} in QCD for 13 TeV LHC as well as the corresponding K -factors (right panel). It is worth noting here that at $\mathcal{O}(\alpha_s^3)$ level the $\delta(1-z)$ contribution is comparable but opposite in sign to the sum of logarithmic contributions

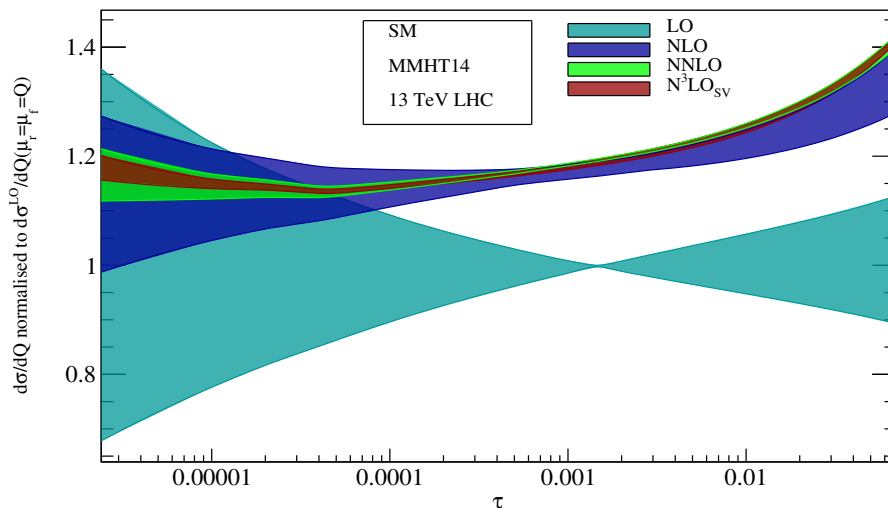


Figure 2. 7-point scale variation is plotted against the hadronic τ variable up to N^3LO_{sv} order. All the figures are normalized to LO contribution evaluated at the central scale $\mu_r = \mu_f = Q$.

as is mentioned in [17]. The 3-loop SV corrections are found to be positive up to around $Q = 400$ GeV and remain negative for $400 \text{ GeV} < Q < 2200$ GeV and become positive thereafter as threshold logarithms dominate in the high Q region. At around 3500 GeV, the 3-loop SV corrections contribute by about 2%. The observed values of Q where this change in the sign happens are not fixed but can change with the center of mass energy of the hadrons.

While the perturbation series is asymptotic and the higher orders terms are very small, the reliability of the theory predictions depends somewhat on the uncertainties due to the unphysical factorization (μ_f) and renormalization (μ_r) scales as well as those due to choice of PDFs. To this end, we estimate the 7-point scale uncertainties in the invariant mass distribution at various orders in the perturbation theory by varying the scales $\mu = \{\mu_f, \mu_r\}$ in the range $\frac{1}{2} \leq \frac{\mu}{Q} \leq 2$. The scale uncertainties are conveniently presented in terms of the invariant mass distribution at higher orders normalized with respect to LO ones. In figure 2 we present these normalized distributions up to N^3LO_{sv} as a function of $\tau = Q^2/S$. At LO, there is no dependence on μ_r , hence the observation that these scale uncertainties are minimum around $\tau = 0.001$ (corresponding to about $Q = 400$ GeV) can be directly related to the behavior of the corresponding quark fluxes. At higher orders, the dependence on μ_r and μ_f is known and the scale uncertainties are found to increase with Q in the region $Q > 400$ GeV. For $Q = 1500$ GeV, they are found to be 6.61%, 2.40%, 0.44% and 0.91% respectively at LO, NLO, NNLO and N^3LO_{sv} . For the 3-loop SV case, the scale uncertainties are expected to get further reduced only after including the regular terms that are yet to be computed in the fixed order perturbation theory. However, as we increase Q value, even N^3LO_{sv} show reasonable reduction in scale uncertainty as threshold logarithms dominate over the regular terms for larger Q values. For completeness, we note that the scale uncertainties for $Q = 3500$ GeV are found to be 12.90%, 4.64%, 1.21% and 0.68% at LO, NLO, NNLO and N^3LO_{sv} respectively.

It is interesting to compare the SV results obtained in this article with the full N³LO result recently calculated in [4]. For this comparison we chose the same parameters as in [4]. In particular at N³LO level, we use PDF4LHC15_nnlo_mc pdf and strong coupling through four-loop evolution. In order to see how much the SV terms are comparable to the full correction at the third order, we have adapted to DY production through only γ^* as in [4]. The full N³LO correction is negative up to $Q = 1$ TeV after that it becomes positive up to $Q = 2$ TeV (see figure (2) of [4]). On the contrary, the SV result is positive at lower invariant mass region and becomes negative at Q as low as 100 GeV. This indicates that the contributions from the other subprocesses are really important in the low invariant mass region $Q < 100$ GeV. In fact in the low- Q region, the $q\bar{q}$ channel really dominates over all other channels including the $q\bar{q}$ (see figure 4 of [4]) and thus completely controls the behavior. The $q\bar{q}$ channel on the other hand really becomes important in the higher Q region. Note that SV corrections are intrinsically ambiguous due to the choice of the function $g(z)$ in eq. (3.2). This essentially does not change the complete cross-section, however when expanded in the limit $z \rightarrow 1$, this basically includes a part of the subleading terms in the $q\bar{q}$ channel. It is thus interesting to study the subleading behavior of the $q\bar{q}$ channel itself by using the treatment performed in [66–68]. At N³LO we have modified the SV result by a proper weight function $g(z)$ which minimises the subleading regular effects,

$$\begin{aligned} \sigma &= \sigma^{(0)} \sum_{ab=q,\bar{q},g} \int_0^1 dx_1 \int_0^1 dx_2 \int_0^1 dz f_a(x_1, \mu_f^2) f_b(x_2, \mu_f^2) / g(z) \\ &\times [\Delta_{ab}(z, Q^2, \mu_f^2) g(z)] \delta(\tau - zx_1x_2). \end{aligned} \tag{3.2}$$

With the different choices of the weight function $g(z)$, we observe that $g(z) = z$ gives a result which is very close to the complete result from the $q\bar{q}$ channel at the lower orders. Notice that this choice correctly reproduces the leading collinear logarithm as well as a part of other subleading terms at each order. With this modified SV cross-section (denoted as SVM) we observe that the N³LO corrections are negative below $Q = 1$ TeV and becomes positive afterwards. This shows that the subleading pieces from the $q\bar{q}$ channels are also important in the region below 1 TeV. At $Q = 30$ GeV, the SV(SVM) corrections differ by as large as 5.2%(4.8%) from the complete N³LO. This is direct consequence of the large negative $q\bar{q}$ contribution in this region. As we approach the threshold region, the SV (SVM) terms dominate and around $Q = 1800$ GeV they differs only by 0.7% (0.2%) from the exact result, confirming the reliability of the threshold result in this region.

3.2 Resummed prediction for neutral DY invariant mass

We have studied the impact of different resummation schemes as described in the previous section. First we compare the resummed results between two approaches: the **Standard** \bar{N} and **Standard** N prescriptions. We find that the perturbative convergence is better in the case of \bar{N} exponentiation for the scale choice $\mu_r = \mu_f = Q$. This can be clearly seen from figure 3 where the convergence is already achieved at NLO+NLL whereas in N exponentiation it happens only after NLO+NLL order. At $Q = 2500$ GeV, we see the corrections received in **Standard** N exponentiation is 21.6% at NLO+NLL, 2.2%

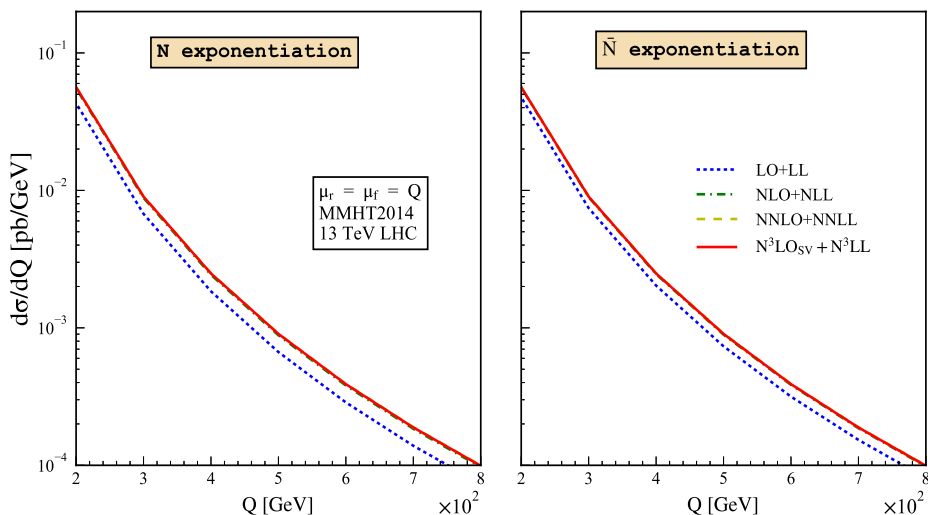


Figure 3. The comparison between **Standard** N and \bar{N} resummation approaches for di-lepton invariant mass distribution is presented up to N^3LL accuracy for 13 TeV LHC.

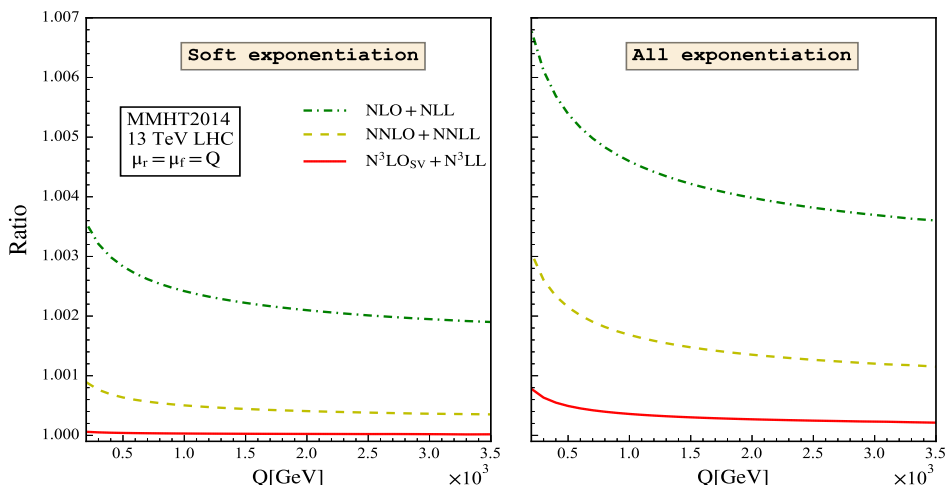


Figure 4. Comparison between the **Soft** (left panel) and **All exponentiation** (right panel) with **Standard** \bar{N} approach. Here, the ratio is taken over the **Standard** \bar{N} results.

at NNLO+NNLL whereas in the **Standard** \bar{N} exponentiation these are 6.7% and 2.3% respectively. This observation is also true for different scale choices. This is expected since naively one can expect that as we exponentiate more and more terms the convergence becomes better. In the rest of the discussion we will mention ‘**Standard**’ only in the context of \bar{N} exponentiation unless otherwise stated.

We now investigate the differences resulting from remaining two approaches *viz.* the **Soft exponentiation** and **All exponentiation** to study their perturbative behavior. To illustrate this, we show figure 4 where we took the ratio with respect to the **Standard** \bar{N} results at each order. Notice that LO+LL results are same for all these three approaches by construction. To this end one sees that at lower orders the resummed cross-sections deviate

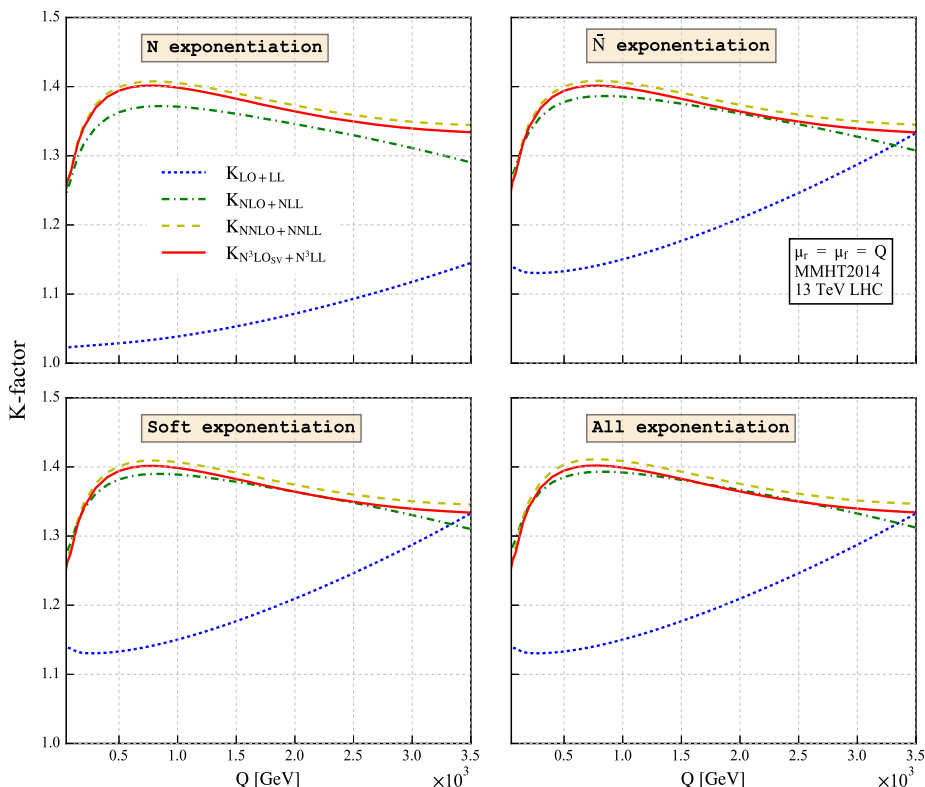


Figure 5. The K -factors are shown for resummed results up to N^3LL level for different threshold resummation approaches (discussed in the text).

more from those of \bar{N} exponentiations. At NNLL the **Soft exponentiation** gets additional 0.12% corrections compared to the **Standard \bar{N}** approach at $Q = 100$ GeV. However at N^3LL level the **Soft exponentiation** does not improve over the **Standard \bar{N}** results and both approaches provide almost same results. On the other hand, **All exponentiation** still gets some contribution from higher orders through the exponentiation of complete \bar{g}_0 even at N^3LL order. The increment is however very small giving only 0.12% corrections over the **Standard \bar{N}** scenario.

We have quantified the impact of resummed results through K -factor. In figure 5 we present the resummed K -factors ($K_{NLO+NLL}$, $K_{NNLO+NNLL}$, $K_{N^3LO_{sv}+N^3LL}$) up to order N^3LL . We define the K -factor as $\frac{d\sigma^{resum}}{dQ} / \frac{d\sigma^{LO}}{dQ}$, where *resum* represents all the resummed corrections up to $N^3LO_{sv} + N^3LL$. One observes that the perturbative convergence is improved in the case of **All exponentiation** compared to others although marginally. The K factor defined this way will be useful to directly compare against the experimental results. For **All exponentiation** case, we find that the K -factor is 1.294 at $Q = 100$ at NNLL which changes to 1.286 at N^3LL . The K -factor increases with Q . At higher $Q = 2500$ GeV the K -factors become 1.362 at NNLL and 1.350 at N^3LL .

Next we study the uncertainties resulting from unphysical scale in these approaches. We follow the canonical variation of μ_f and μ_r around the final state invariant mass Q within $[1/2, 2]Q$ imposing additional constraint $1/2 \leq \mu_r/\mu_f \leq 2$ as was done in the

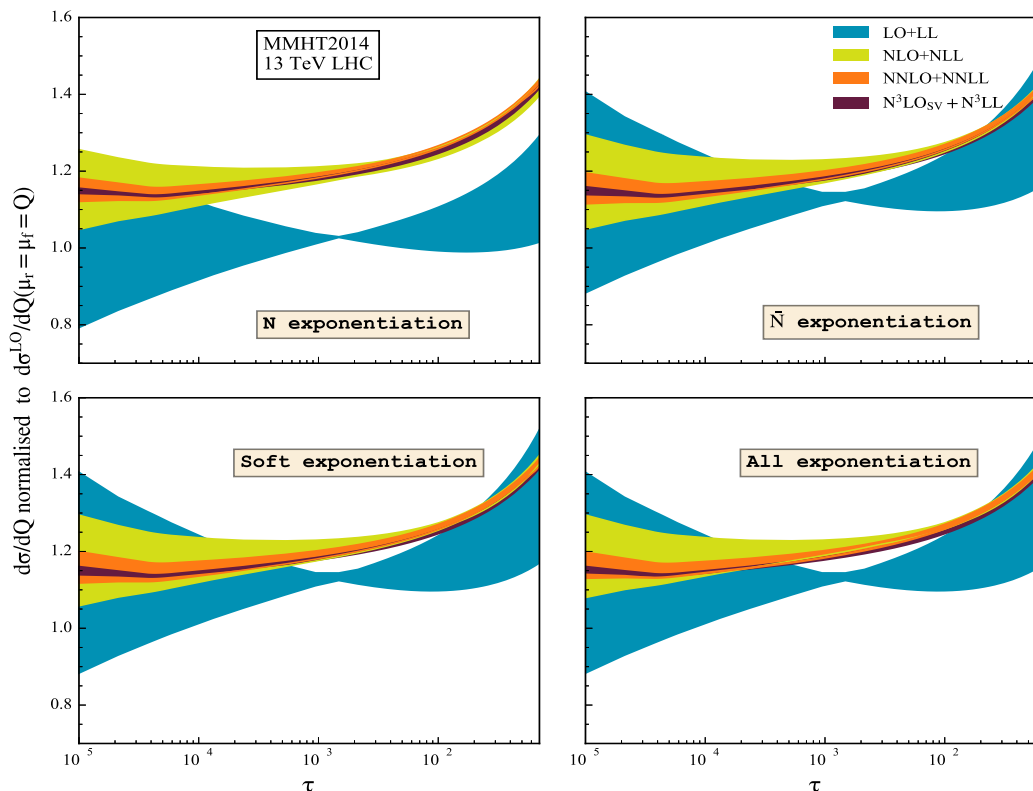


Figure 6. 7-point scale variations around the central scale choice ($\mu_r = Q, \mu_f = Q$) are presented as in figure 2 but for resummed predictions up to N^3LL accuracy.

third order SV prediction in the previous section. We notice that different approaches for resummation provide a systematic reduction in the scale uncertainties. For example, in the **Standard \bar{N}** case, the scale uncertainty reduces from 13.37% at LO+LL to 6.91% at NLO+NLL to 1.99% at NNLO+NNLL. For a higher invariant mass $Q = 1500$ GeV that we have considered in our analysis as we approach towards the threshold region, we notice that scale uncertainties got reduced significantly to 7.61% at LO+LL, 0.90% at NLO+NLL, 0.52% at NNLO+NNLL. In figure 6 a similar pattern of reduction in the scale uncertainties is seen for the **Soft** and **All exponentiations** as we go to higher logarithmic accuracy for invariant mass up to $Q = 1500$ GeV. However, when we compare among themselves, the scale uncertainties at LO+LL remain the same for all the approaches by construction. For higher logarithmic accuracy, we see that **All** has the smallest uncertainty up to NNLO+NNLL through $Q = 1500$ GeV. However, for the **Soft** case that we have considered here the NLO+NLL has larger uncertainty than that of **All** and at NNLO+NNLL the soft is the largest. At $N^3LO_{sv}+N^3LL$, the scale uncertainties for all the different approaches are found to be less than 1% for the entire Q region we have considered. However, these scale uncertainties at N^3LL level are largest for **All** case while \bar{N} and **Soft** have similar scale uncertainties. Though the behaviour of the scale uncertainties in all these three schemes are competing with each other, it can be seen that **All exponentiation** has the largest scale uncertainties at $N^3LO_{sv}+N^3LL$ and smallest at lower orders for a

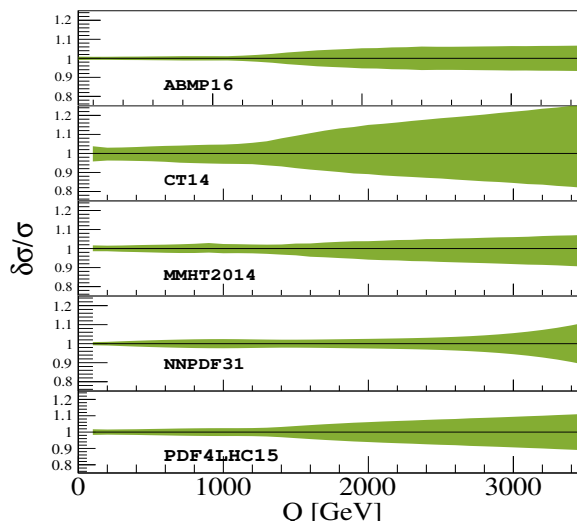


Figure 7. Intrinsic PDF uncertainties in di-lepton invariant mass distribution have been estimated at NNLO+NNLL level taking $\mu_r = \mu_f = Q$. Here σ is obtained from the central set $n = 0$ provided by the respective PDF group.

wide range of Q values. This shows that the sub-leading regular pieces are also important to capture the scale dependence properly. We will again come back on this discussion at the end of this section.

We have also estimated in our resummed predictions the uncertainties from the non-perturbative PDFs. We convolute the resummed coefficient at NNLL level with n different sets of a given PDF group and estimate the uncertainty from the `lhpdf` routines. We use the PDFs provided by ABMP16 ($n=30$) [69], CT14 ($n=57$) [70], MMHT2014 ($n=51$) [65], NNPDF31 ($n=101$) [71] and PDF4LHC15 ($n=31$) [72] groups. These results are shown in figure 7 in terms of $\delta\sigma/\sigma$ where $\delta\sigma$ is the difference between the extrema obtained from n different sets and σ is the one obtained from central set $n = 0$. These PDF uncertainties in general are found to increase with the invariant mass of the di-lepton pair and, for the range of Q considered here, we find that they are smallest in the low Q -region for AMP16 and are largest for CT14 case. These uncertainties for $Q = 1500$ GeV are found to be 6.14% (AMP16), 16.99% (CT14), 6.17% (MMHT2014), 4.21% (NNPDF31) and 7.43% (PDF4LHC15).

Finally, we discuss the matching relation presented in eq. (2.46). One can match the $N^3\text{LO}_{sv}$ fixed order results (with $n = 3$) with the resummed results subtracted up to $\mathcal{O}(a_s^3)$ (with $n = 3$) in order to avoid any double counting from the fixed order. So far, we have followed this approach. Instead we can match the complete NNLO fixed order result with the resummed result subtracted up to $\mathcal{O}(a_s^2)$, which also avoids double counting and retains the threshold terms at $\mathcal{O}(a_s^3)$ in N -space in the threshold limit $N \rightarrow \infty$. The difference in these two approaches is sub-leading and has to be related with the fact that N -space threshold results when transformed back into distribution space produces sub-leading logarithms in addition to the plus distributions. In figure 8 we compare these two approaches setting all the scales same as Q in `Standard \bar{N}` approach. We see that the threshold terms defined in Mellin- N space provide much better perturbative convergence

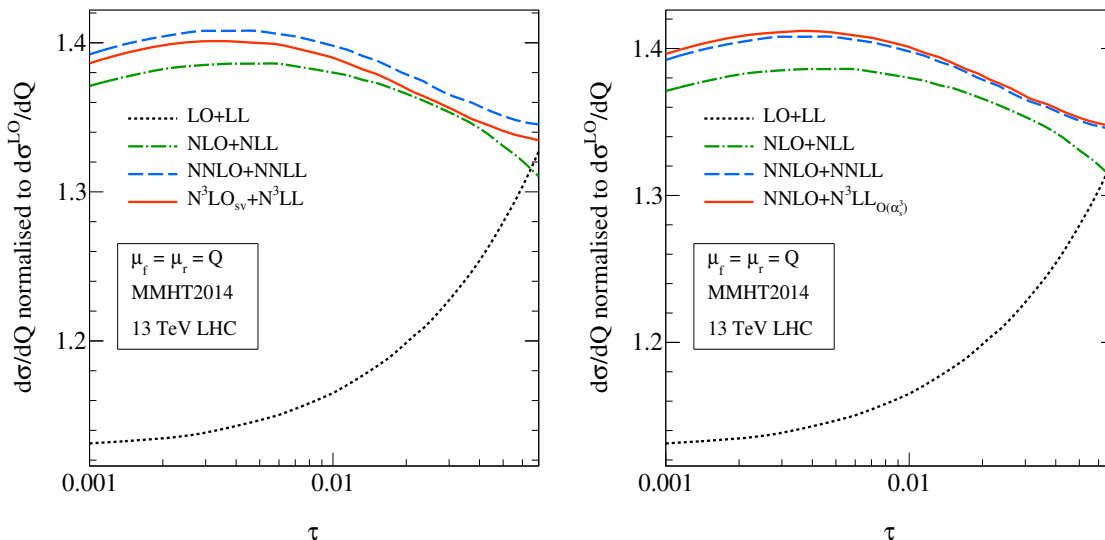


Figure 8. Comparison between two different ways of matching with the fixed order. In one case, the matching is done with threshold logarithms kept in the distribution (z) space (left) and in the other case the matching is done with threshold logarithms in the Mellin- N space (right).

\sqrt{S} (TeV)	LO	NLO	NNLO	N^3LO_{sv}	LO+LL	NLO+NLL	NNLO+NNLL	$N^3LO_{sv}+N^3LL$
7	22.286 ($\pm 9.99\%$)	29.041 ($\pm 3.47\%$)	29.994 ($\pm 0.87\%$)	30.246 ($\pm 2.28\%$)	25.466 ($\pm 10.51\%$)	29.996 ($\pm 5.99\%$)	30.148 ($\pm 2.04\%$)	30.242 ($\pm 2.27\%$)
8	26.202 ($\pm 10.85\%$)	33.846 ($\pm 3.78\%$)	34.905 ($\pm 0.92\%$)	35.197 ($\pm 2.22\%$)	29.904 ($\pm 11.36\%$)	34.946 ($\pm 6.28\%$)	35.083 ($\pm 2.10\%$)	35.192 ($\pm 2.20\%$)
13	46.465 ($\pm 13.84\%$)	57.957 ($\pm 4.91\%$)	59.379 ($\pm 1.17\%$)	59.840 ($\pm 2.04\%$)	52.829 ($\pm 14.31\%$)	59.774 ($\pm 7.35\%$)	59.666 ($\pm 2.29\%$)	59.834 ($\pm 2.02\%$)
14	50.610 ($\pm 14.27\%$)	62.770 ($\pm 5.08\%$)	64.231 ($\pm 1.12\%$)	64.723 ($\pm 2.02\%$)	57.512 ($\pm 14.73\%$)	64.627 ($\pm 7.51\%$)	64.540 ($\pm 2.32\%$)	64.717 ($\pm 1.99\%$)

Table 1. Fixed order (up to N^3LO_{sv}) and resummed (up to $N^3LO_{sv} + N^3LL$) cross section (in nb) for on-shell Z -boson production at different center of mass energy of LHC. The scale uncertainty has been estimated using seven-point scale variation around the central scale $(\mu_r, \mu_f) = (1, 1)M_Z$.

compared to the z -space definition. This is a well-known observation which shows that the sub-leading pieces are also important at this order. As far as scale uncertainty is concerned, this approach gives better estimate of scale uncertainty at N^3LL level reducing in some cases by a factor of two, however the general behavior does not change much.

3.3 Resummed prediction for Z/W^\pm productions

In this section we present the resummed results for on-shell Z and W^\pm productions to $N^3LO_{sv}+N^3LL$ accuracy. We set all the parameters same as the previous section. For pdf, we chose the central value from MMHT2014 set at the corresponding order. At the LHC, the underlying parton fluxes for W^+ production are larger than for W^- case, consequently the production cross sections for the former case are larger than the latter one. This is true also for higher centre of mass energies. In table 1, 2, 3, we present for different center of mass energies at the LHC, the central predictions for on-shell Z , W^+ and W^-

\sqrt{S} (TeV)	LO	NLO	NNLO	N ³ LO _{sv}	LO+LL	NLO+NLL	NNLO+NNLL	N ³ LO _{sv} +N ³ LL
7	43.758 (±10.68%)	56.962 (±3.25%)	58.717 (±0.74%)	59.354 (±3.41%)	49.761 (±11.19%)	58.761 (±5.76%)	59.006 (±1.00%)	59.347 (±3.38%)
8	50.820 (±11.50%)	65.525 (±3.50%)	67.400 (±0.88%)	68.140 (±3.48%)	57.742 (±12.00%)	67.578 (±6.00%)	67.728 (±0.95%)	68.131 (±3.46%)
13	86.542 (±14.34%)	107.427 (±4.41%)	109.454 (±1.43%)	110.700 (±3.86%)	98.044 (±14.8%)	110.700 (±6.86%)	109.967 (±0.94%)	110.688 (±3.83%)
14	93.726 (±14.75%)	115.635 (±4.55%)	117.616 (±1.52%)	118.961 (±3.93%)	106.139 (±15.21%)	119.145 (±6.99%)	118.163 (±0.95%)	118.948 (±3.91%)

Table 2. Same as table 1 but for W^+ production at the LHC.

\sqrt{S} (TeV)	LO	NLO	NNLO	N ³ LO _{sv}	LO+LL	NLO+NLL	NNLO+NNLL	N ³ LO _{sv} +N ³ LL
7	30.757 (±11.17%)	39.324 (±2.99%)	40.401 (±1.36%)	40.837 (±4.09%)	35.291 (±11.68%)	40.706 (±5.69%)	40.628 (±1.05%)	40.829 (±4.07%)
8	36.238 (±12.00%)	45.937 (±3.23%)	47.100 (±1.53%)	47.615 (±4.19%)	41.519 (±12.51%)	47.531 (±5.91%)	47.359 (±1.08%)	47.606 (±4.18%)
13	64.571 (±14.89%)	79.089 (±4.1%)	80.441 (±2.22%)	81.358 (±4.66%)	73.646 (±15.36%)	81.719 (±6.7%)	80.862 (±1.56%)	81.347 (±4.63%)
14	70.360 (±15.31%)	85.696 (±4.23%)	87.043 (±2.33%)	88.042 (±4.74%)	80.199 (±15.77%)	88.530 (±6.83%)	87.496 (±1.67%)	88.029 (±4.72%)

Table 3. Same as table 1 but for W^- production at the LHC.

respectively along with the corresponding percentage of scale uncertainties. Note that the scale uncertainties are calculated again using the same procedure i.e. the 7-point scale variation around the central scale which is now the vector boson mass i.e. the central scale has been chosen as $(\mu_r, \mu_f) = (1, 1)M_V$, with $V = Z$ for Z production and $V = W^\pm$ for W -boson production. In all the cases we observed that the fixed order scale uncertainties are systematically reduced while going to higher orders, however at N³LO_{sv}, it again increases which is due to the fact that, at this order there are still missing regular pieces as well as missing N³LO PDF which are essential to the scale uncertainty. Similar observation is also seen for the matched resummed prediction. We see that compared to the fixed order, the resummed results provide better perturbative convergence. To further investigate the source of this huge 7-point scale uncertainty at the N³LO_{sv} level, we have calculated the independent μ_r and μ_f uncertainties symmetrized around the central scale. We find that the μ_r scale uncertainty is indeed improved. For 14 TeV centre-of-mass energies we find the μ_r uncertainty at N³LO_{sv} reduces to $\pm 0.32\%$ for Z case, $\pm 0.49\%$ for W^+ and $\pm 0.51\%$ for W^- production. On the other hand the μ_f uncertainty deteriorates at the N³LO order and could lead to uncertainty as large as about $\pm 5\%$. This is expected since at N³LO level we are missing the correct PDF to cancel the residual μ_f part in the coefficient function. In case of resummed prediction we observe further reduction of μ_r scale uncertainties to $\pm 0.2\%$, $\pm 0.37\%$ and $\pm 0.38\%$ for Z , W^+ and W^- productions respectively. The resummed K -factors as defined before, increases from NNLO+NNLL to N³LO_{sv}+N³LL for all the cases. The absolute size of the perturbative corrections however decreases at N³LO_{sv}+N³LL compared to the previous orders confirming the reliability of perturbation theory.

Note that one can perform a similar study as we did at the end of section 3.1 to estimate the effect of subleading logarithms at the $N^3\text{LO}$ level. The subleading effects at this region are expected to be comparable to the SV corrections. We find that for on-shell Z production, the SVM cross-section decreases the SV corrections presented above at the third order by 0.1%, while for W case the respective contribution is about 0.2%. However similar to the photon mediated DY case, at this region the other subprocesses will be as important as the $q\bar{q}$ one. We have also compared the SVM result for W -production with the recently obtained exact result at $N^3\text{LO}$ [73], and we find that including the other subprocesses further bring down the cross-section additionally 2%.

4 Conclusions

We have studied the Drell-Yan production of di-lepton as well as on-shell Z and W^\pm productions in the context of threshold resummation and presented our results to $N^3\text{LL}$ accuracy for different resummation prescriptions. The threshold corrections are important in the large invariant mass region (above $Q = 1800\text{ GeV}$). However in the moderate invariant mass region also we find substantial contribution from the threshold terms. We have used all the necessary ingredients available to perform resummation, in particular the threshold enhanced large- N as well as the N -independent constants. The standard threshold resummation uses results of the SV cross-section at any given order. In particular we showed how the large N -independent constants can be found at $N^3\text{LL}$ level using the existing SV results. We have matched our resummed $N^3\text{LL}$ results with the existing NNLO($N^3\text{LO}_{sv}$) cross-section and presented results for 13 TeV LHC. First, we observed that the resummed results obtained by exponentiating $\ln\bar{N}$ pieces give faster convergence of the perturbation series compared to the conventional case where $\ln N$ terms have been exponentiated. Further, we explored other possibilities of doing resummation where we exponentiate complete soft pieces coming from the universal soft distribution function and notice that the perturbative convergence for the **Soft** case is bit faster than the **Standard** \bar{N} case. We also presented our results when the complete g_0 coefficients including the form factor have been exponentiated and found that the convergence rate of the perturbation series is competing with that of the **Soft** case. Over all, we observe that these different approaches show a systematic behavior of the resummed predictions where the convergence of the perturbation series gets better when more and more N -independent terms are exponentiated. For scale uncertainties up to NNLO+NNLL, **A11** has the lowest scale uncertainties for a wide range of Q values. We also note that at $N^3\text{LL}$ accuracy, however, the missing regular pieces are also important and so is $N^3\text{LO}$ PDFs to tame the overall scale uncertainty.

Acknowledgments

V.R. thanks C. Duhr for useful discussions. G.D. would like to thank G. Bell, Y. Ji, L. Magnea, S. Moch and A. Vogt for useful discussions. The algebraic computations have been done with the latest version of the symbolic manipulation system FORM [74, 75]. The

research of G.D. is supported partially by DESY postdoctoral fellowship and the DFG within the Collaborative Research Center TRR 257 (“*Particle Physics Phenomenology after the Higgs Discovery*”). M.C.K. and K.S. would like to thank the High Performance Computing Facility *Param Ishan* at IIT Guwahati, where a part of the computational work for this study has been carried out.

A Soft-virtual coefficient in N -space

The SV coefficient up to three loops are presented here (denoting $\bar{L} = \ln\bar{N}$),

$$\Delta_N^{\text{sv},(1)} = \bar{L}^2 \left(2A_1 \right) + \bar{L} \left(2A_1 L_{fr} - 2A_1 L_{qr} + 2f_1 \right) + \bar{g}_{01}, \quad (\text{A.1})$$

$$\begin{aligned} \Delta_N^{\text{sv},(2)} = & \bar{L}^4 \left(2A_1^2 \right) + \bar{L}^3 \left(4A_1^2 L_{fr} - 4A_1^2 L_{qr} + \frac{4}{3}A_1\beta_0 + 4A_1 f_1 \right) + \bar{L}^2 \left(10A_1^2 \zeta_2 + 2A_1^2 L_{fr}^2 \right. \\ & - 4A_1^2 L_{fr} L_{qr} + 2A_1^2 L_{qr}^2 - 2A_1\beta_0 L_{qr} - 4A_1 B_1 L_{fr} + 4A_1 B_1 L_{qr} + 4A_1 f_1 L_{fr} \\ & - 4A_1 f_1 L_{qr} + 4A_1 G_{11} + 4A_1 \tilde{G}_{11} + 2A_2 + 2\beta_0 f_1 + 2f_1^2 \left. \right) + \bar{L} \left(10A_1^2 \zeta_2 L_{fr} \right. \\ & - 10A_1^2 \zeta_2 L_{qr} + 4A_1\beta_0 \zeta_2 - A_1\beta_0 L_{fr}^2 + A_1\beta_0 L_{qr}^2 + 10A_1 \zeta_2 f_1 - 4A_1 B_1 L_{fr}^2 \\ & + 8A_1 B_1 L_{fr} L_{qr} - 4A_1 B_1 L_{qr}^2 + 4A_1 G_{11} L_{fr} - 4A_1 G_{11} L_{qr} + 4A_1 \tilde{G}_{11} L_{fr} \\ & - 4A_1 \tilde{G}_{11} L_{qr} + 2A_2 L_{fr} - 2A_2 L_{qr} - 2\beta_0 f_1 L_{qr} + 4\beta_0 \tilde{G}_{11} - 4B_1 f_1 L_{fr} + 4B_1 f_1 L_{qr} \\ & \left. + 4f_1 G_{11} + 4f_1 \tilde{G}_{11} + 2f_2 \right) + \bar{g}_{02}, \quad (\text{A.2}) \end{aligned}$$

$$\begin{aligned} \Delta_N^{\text{sv},(3)} = & \bar{L}^6 \left(\frac{4}{3}A_1^3 \right) + \bar{L}^5 \left(4A_1^3 L_{fr} - 4A_1^3 L_{qr} + \frac{8}{3}A_1^2\beta_0 + 4A_1^2 f_1 \right) + \bar{L}^4 \left(10A_1^3 \zeta_2 + 4A_1^3 L_{fr}^2 \right. \\ & - 8A_1^3 L_{fr} L_{qr} + 4A_1^3 L_{qr}^2 + \frac{8}{3}A_1^2\beta_0 L_{fr} - \frac{20}{3}A_1^2\beta_0 L_{qr} - 4A_1^2 B_1 L_{fr} + 4A_1^2 B_1 L_{qr} \\ & + 8A_1^2 f_1 L_{fr} - 8A_1^2 f_1 L_{qr} + 4A_1^2 G_{11} + 4A_1^2 \tilde{G}_{11} + 4A_1 A_2 + \frac{4}{3}A_1\beta_0^2 + \frac{20}{3}A_1\beta_0 f_1 \\ & \left. + 4A_1 f_1^2 \right) + \bar{L}^3 \left(20A_1^3 \zeta_2 L_{fr} - 20A_1^3 \zeta_2 L_{qr} + \frac{4}{3}A_1^3 L_{fr}^3 - 4A_1^3 L_{fr}^2 L_{qr} + 4A_1^3 L_{fr} L_{qr}^2 \right. \\ & - \frac{4}{3}A_1^3 L_{qr}^3 + \frac{44}{3}A_1^2\beta_0 \zeta_2 - 2A_1^2\beta_0 L_{fr}^2 - 4A_1^2\beta_0 L_{fr} L_{qr} + 6A_1^2\beta_0 L_{qr}^2 + 20A_1^2 \zeta_2 f_1 \\ & - 8A_1^2 B_1 L_{fr}^2 + 16A_1^2 B_1 L_{fr} L_{qr} - 8A_1^2 B_1 L_{qr}^2 + 4A_1^2 f_1 L_{fr}^2 - 8A_1^2 f_1 L_{fr} L_{qr} \\ & + 4A_1^2 f_1 L_{qr}^2 + 8A_1^2 G_{11} L_{fr} - 8A_1^2 G_{11} L_{qr} + 8A_1^2 \tilde{G}_{11} L_{fr} - 8A_1^2 \tilde{G}_{11} L_{qr} \\ & + 8A_1 A_2 L_{fr} - 8A_1 A_2 L_{qr} - \frac{8}{3}A_1\beta_0^2 L_{qr} - \frac{8}{3}A_1\beta_0 B_1 L_{fr} + \frac{8}{3}A_1\beta_0 B_1 L_{qr} \\ & + 4A_1\beta_0 f_1 L_{fr} - 12A_1\beta_0 f_1 L_{qr} + \frac{8}{3}A_1\beta_0 G_{11} + \frac{32}{3}A_1\beta_0 \tilde{G}_{11} + \frac{4}{3}A_1\beta_1 \\ & - 8A_1 B_1 f_1 L_{fr} + 8A_1 B_1 f_1 L_{qr} + 4A_1 f_1^2 L_{fr} - 4A_1 f_1^2 L_{qr} + 8A_1 f_1 G_{11} + 8A_1 f_1 \tilde{G}_{11} \\ & \left. + 4A_1 f_2 + \frac{8}{3}A_2\beta_0 + 4A_2 f_1 + \frac{8}{3}\beta_0^2 f_1 + 4\beta_0 f_1^2 + \frac{4}{3}f_1^3 \right) + \bar{L}^2 \left(25A_1^3 \zeta_2^2 + 10A_1^3 \zeta_2 L_{fr}^2 \right. \\ & - 20A_1^3 \zeta_2 L_{fr} L_{qr} + 10A_1^3 \zeta_2 L_{qr}^2 + 8A_1^2\beta_0 \zeta_2 L_{fr} - 28A_1^2\beta_0 \zeta_2 L_{qr} + \frac{16}{3}A_1^2\beta_0 \zeta_3 \end{aligned}$$

$$\begin{aligned}
& -2A_1^2\beta_0L_{fr}^3 + 2A_1^2\beta_0L_{fr}^2L_{qr} + 2A_1^2\beta_0L_{fr}L_{qr}^2 - 2A_1^2\beta_0L_{qr}^3 - 20A_1^2\zeta_2B_1L_{fr} \\
& + 20A_1^2\zeta_2B_1L_{qr} + 20A_1^2\zeta_2f_1L_{fr} - 20A_1^2\zeta_2f_1L_{qr} + 20A_1^2\zeta_2G_{11} + 20A_1^2\zeta_2\tilde{G}_{11} \\
& - 4A_1^2B_1L_{fr}^3 + 12A_1^2B_1L_{fr}^2L_{qr} - 12A_1^2B_1L_{fr}L_{qr}^2 + 4A_1^2B_1L_{qr}^3 + 4A_1^2G_{11}L_{fr}^2 \\
& - 8A_1^2G_{11}L_{fr}L_{qr} + 4A_1^2G_{11}L_{qr}^2 + 4A_1^2\tilde{G}_{11}L_{fr}^2 - 8A_1^2\tilde{G}_{11}L_{fr}L_{qr} + 4A_1^2\tilde{G}_{11}L_{qr}^2 \\
& + 20A_1A_2\zeta_2 + 4A_1A_2L_{fr}^2 - 8A_1A_2L_{fr}L_{qr} + 4A_1A_2L_{qr}^2 + 8A_1\beta_0^2\zeta_2 + 2A_1\beta_0^2L_{qr}^2 \\
& + 12A_1\beta_0\zeta_2B_1 + 28A_1\beta_0\zeta_2f_1 + 2A_1\beta_0B_1L_{fr}^2 + 4A_1\beta_0B_1L_{fr}L_{qr} - 6A_1\beta_0B_1L_{qr}^2 \\
& - 2A_1\beta_0f_1L_{fr}^2 - 4A_1\beta_0f_1L_{fr}L_{qr} + 6A_1\beta_0f_1L_{qr}^2 - 8A_1\beta_0G_{11}L_{qr} + 4A_1\beta_0G_{12} \\
& + 8A_1\beta_0\tilde{G}_{11}L_{fr} - 16A_1\beta_0\tilde{G}_{11}L_{qr} + 4A_1\beta_0\tilde{G}_{12} - 2A_1\beta_1L_{qr} + 10A_1\zeta_2f_1^2 \\
& + 4A_1B_1^2L_{fr}^2 - 8A_1B_1^2L_{fr}L_{qr} + 4A_1B_1^2L_{qr}^2 - 8A_1B_1f_1L_{fr}^2 + 16A_1B_1f_1L_{fr}L_{qr} \\
& - 8A_1B_1f_1L_{qr}^2 - 8A_1B_1G_{11}L_{fr} + 8A_1B_1G_{11}L_{qr} - 8A_1B_1\tilde{G}_{11}L_{fr} \\
& + 8A_1B_1\tilde{G}_{11}L_{qr} - 4A_1B_2L_{fr} + 4A_1B_2L_{qr} + 8A_1f_1G_{11}L_{fr} - 8A_1f_1G_{11}L_{qr} \\
& + 8A_1f_1\tilde{G}_{11}L_{fr} - 8A_1f_1\tilde{G}_{11}L_{qr} + 4A_1f_2L_{fr} - 4A_1f_2L_{qr} + 4A_1G_{11}^2 \\
& + 8A_1G_{11}\tilde{G}_{11} + 2A_1G_{21} + 4A_1\tilde{G}_{11}^2 + 2A_1\tilde{G}_{21} - 4A_2\beta_0L_{qr} - 4A_2B_1L_{fr} \\
& + 4A_2B_1L_{qr} + 4A_2f_1L_{fr} - 4A_2f_1L_{qr} + 4A_2G_{11} + 4A_2\tilde{G}_{11} + 2A_3 - 4\beta_0^2f_1L_{qr} \\
& + 8\beta_0^2\tilde{G}_{11} - 4\beta_0B_1f_1L_{fr} + 4\beta_0B_1f_1L_{qr} - 4\beta_0f_1^2L_{qr} + 4\beta_0f_1G_{11} + 12\beta_0f_1\tilde{G}_{11} \\
& + 4\beta_0f_2 + 2\beta_1f_1 - 4B_1f_1^2L_{fr} + 4B_1f_1^2L_{qr} + 4f_1^2G_{11} + 4f_1^2\tilde{G}_{11} + 4f_1f_2) \\
& + \bar{L} \left(25A_1^3\zeta_2^2L_{fr} - 25A_1^3\zeta_2^2L_{qr} + 20A_1^2\beta_0\zeta_2^2 - 5A_1^2\beta_0\zeta_2L_{fr}^2 - 10A_1^2\beta_0\zeta_2L_{fr}L_{qr} \right. \\
& + 15A_1^2\beta_0\zeta_2L_{qr}^2 + \frac{16}{3}A_1^2\beta_0\zeta_3L_{fr} - \frac{16}{3}A_1^2\beta_0\zeta_3L_{qr} + 25A_1^2\zeta_2^2f_1 - 20A_1^2\zeta_2B_1L_{fr}^2 \\
& + 40A_1^2\zeta_2B_1L_{fr}L_{qr} - 20A_1^2\zeta_2B_1L_{qr}^2 + 20A_1^2\zeta_2G_{11}L_{fr} - 20A_1^2\zeta_2G_{11}L_{qr} \\
& + 20A_1^2\zeta_2\tilde{G}_{11}L_{fr} - 20A_1^2\zeta_2\tilde{G}_{11}L_{qr} + 20A_1A_2\zeta_2L_{fr} - 20A_1A_2\zeta_2L_{qr} \\
& - 8A_1\beta_0^2\zeta_2L_{qr} + \frac{32}{3}A_1\beta_0^2\zeta_3 + \frac{2}{3}A_1\beta_0^2L_{fr}^3 - \frac{2}{3}A_1\beta_0^2L_{qr}^3 + 4A_1\beta_0\zeta_2B_1L_{fr} \\
& - 4A_1\beta_0\zeta_2B_1L_{qr} + 10A_1\beta_0\zeta_2f_1L_{fr} - 30A_1\beta_0\zeta_2f_1L_{qr} + 8A_1\beta_0\zeta_2G_{11} \\
& + 28A_1\beta_0\zeta_2\tilde{G}_{11} + \frac{16}{3}A_1\beta_0\zeta_3f_1 + 4A_1\beta_0B_1L_{fr}^3 - 4A_1\beta_0B_1L_{fr}^2L_{qr} \\
& - 4A_1\beta_0B_1L_{fr}L_{qr}^2 + 4A_1\beta_0B_1L_{qr}^3 - 2A_1\beta_0G_{11}L_{fr}^2 - 4A_1\beta_0G_{11}L_{fr}L_{qr} \\
& + 6A_1\beta_0G_{11}L_{qr}^2 + 4A_1\beta_0G_{12}L_{fr} - 4A_1\beta_0G_{12}L_{qr} - 2A_1\beta_0\tilde{G}_{11}L_{fr}^2 \\
& - 4A_1\beta_0\tilde{G}_{11}L_{fr}L_{qr} + 6A_1\beta_0\tilde{G}_{11}L_{qr}^2 + 4A_1\beta_0\tilde{G}_{12}L_{fr} - 4A_1\beta_0\tilde{G}_{12}L_{qr} + 4A_1\beta_1\zeta_2 \\
& - A_1\beta_1L_{fr}^2 + A_1\beta_1L_{qr}^2 - 20A_1\zeta_2B_1f_1L_{fr} + 20A_1\zeta_2B_1f_1L_{qr} + 20A_1\zeta_2f_1G_{11} \\
& + 20A_1\zeta_2f_1\tilde{G}_{11} + 10A_1\zeta_2f_2 + 4A_1B_1^2L_{fr}^3 - 12A_1B_1^2L_{fr}^2L_{qr} + 12A_1B_1^2L_{fr}L_{qr}^2 \\
& - 4A_1B_1^2L_{qr}^3 - 8A_1B_1G_{11}L_{fr}^2 + 16A_1B_1G_{11}L_{fr}L_{qr} - 8A_1B_1G_{11}L_{qr}^2 \\
& - 8A_1B_1\tilde{G}_{11}L_{fr}^2 + 16A_1B_1\tilde{G}_{11}L_{fr}L_{qr} - 8A_1B_1\tilde{G}_{11}L_{qr}^2 - 4A_1B_2L_{fr}^2 \\
& + 8A_1B_2L_{fr}L_{qr} - 4A_1B_2L_{qr}^2 + 4A_1G_{11}^2L_{fr} - 4A_1G_{11}^2L_{qr} + 8A_1G_{11}\tilde{G}_{11}L_{fr} \\
& - 8A_1G_{11}\tilde{G}_{11}L_{qr} + 2A_1G_{21}L_{fr} - 2A_1G_{21}L_{qr} + 4A_1\tilde{G}_{11}^2L_{fr} - 4A_1\tilde{G}_{11}^2L_{qr} \\
& \left. + 2A_1\tilde{G}_{21}L_{fr} - 2A_1\tilde{G}_{21}L_{qr} + 8A_2\beta_0\zeta_2 - 2A_2\beta_0L_{fr}^2 + 2A_2\beta_0L_{qr}^2 + 10A_2\zeta_2f_1 \right)
\end{aligned}$$

$$\begin{aligned}
 & -4A_2B_1L_{fr}^2 + 8A_2B_1L_{fr}L_{qr} - 4A_2B_1L_{qr}^2 + 4A_2G_{11}L_{fr} - 4A_2G_{11}L_{qr} \\
 & + 4A_2\tilde{G}_{11}L_{fr} - 4A_2\tilde{G}_{11}L_{qr} + 2A_3L_{fr} - 2A_3L_{qr} + 8\beta_0^2\zeta_2f_1 + 2\beta_0^2f_1L_{qr}^2 \\
 & - 8\beta_0^2\tilde{G}_{11}L_{qr} + 8\beta_0^2\tilde{G}_{12} + 12\beta_0\zeta_2B_1f_1 + 10\beta_0\zeta_2f_1^2 + 2\beta_0B_1f_1L_{fr}^2 \\
 & + 4\beta_0B_1f_1L_{fr}L_{qr} - 6\beta_0B_1f_1L_{qr}^2 - 8\beta_0B_1\tilde{G}_{11}L_{fr} + 8\beta_0B_1\tilde{G}_{11}L_{qr} - 8\beta_0f_1G_{11}L_{qr} \\
 & + 4\beta_0f_1G_{12} - 8\beta_0f_1\tilde{G}_{11}L_{qr} + 4\beta_0f_1\tilde{G}_{12} - 4\beta_0f_2L_{qr} + 8\beta_0G_{11}\tilde{G}_{11} + 8\beta_0\tilde{G}_{11}^2 \\
 & + 4\beta_0\tilde{G}_{21} - 2\beta_1f_1L_{qr} + 4\beta_1\tilde{G}_{11} + 4B_1^2f_1L_{fr}^2 - 8B_1^2f_1L_{fr}L_{qr} + 4B_1^2f_1L_{qr}^2 \\
 & - 8B_1f_1G_{11}L_{fr} + 8B_1f_1G_{11}L_{qr} - 8B_1f_1\tilde{G}_{11}L_{fr} + 8B_1f_1\tilde{G}_{11}L_{qr} - 4B_1f_2L_{fr} \\
 & + 4B_1f_2L_{qr} - 4B_2f_1L_{fr} + 4B_2f_1L_{qr} + 4f_1G_{11}^2 + 8f_1G_{11}\tilde{G}_{11} + 2f_1G_{21} + 4f_1\tilde{G}_{11}^2 \\
 & + 2f_1\tilde{G}_{21} + 4f_2G_{11} + 4f_2\tilde{G}_{11} + 2f_3) + \bar{g}_{03}. \tag{A.3}
 \end{aligned}$$

Here $L_{fr} = \ln\left(\frac{\mu_f^2}{\mu_r^2}\right)$, $L_{qr} = \ln\left(\frac{Q^2}{\mu_r^2}\right)$. The coefficients \bar{g}_{0i} are given in eq. (B.3).

B Resummed coefficients

Here we collect N -dependent and N -independent coefficients for all different prescriptions for resummation.

B.1 Resummation ingredients for the Standard \bar{N} exponentiation

For the Standard \bar{N} exponentiation we present here the \bar{N} independent coefficients g_0 to three loops in eq. (2.33) below

$$\bar{g}_{01} = \left[\tilde{G}_{11} \left(2 \right) + G_{11} \left(2 \right) + B_1 \left(2 L_{qr} - 2 L_{fr} \right) + A_1 \left(5 \zeta_2 \right) \right], \tag{B.1}$$

$$\begin{aligned}
 \bar{g}_{02} = & \left[\tilde{G}_{21} \left(1 \right) + \tilde{G}_{12} \left(2 \beta_0 \right) + \tilde{G}_{11} \left(- 2 \beta_0 L_{qr} \right) + \tilde{G}_{11}^2 \left(2 \right) + G_{21} \left(1 \right) \right. \\
 & + G_{12} \left(2 \beta_0 \right) + G_{11} \left(- 2 \beta_0 L_{qr} \right) + G_{11} \tilde{G}_{11} \left(4 \right) + G_{11}^2 \left(2 \right) + f_1 \left(5 \beta_0 \zeta_2 \right) \\
 & + B_2 \left(2 L_{qr} - 2 L_{fr} \right) + B_1 \left(- \beta_0 L_{qr}^2 + \beta_0 L_{fr}^2 + 6 \beta_0 \zeta_2 \right) + B_1 \tilde{G}_{11} \left(4 L_{qr} \right. \\
 & \left. - 4 L_{fr} \right) + B_1 G_{11} \left(4 L_{qr} - 4 L_{fr} \right) + B_1^2 \left(2 L_{qr}^2 - 4 L_{fr} L_{qr} + 2 L_{fr}^2 \right) \\
 & + A_2 \left(5 \zeta_2 \right) + A_1 \left(\frac{8}{3} \beta_0 \zeta_3 - 5 \beta_0 \zeta_2 L_{qr} \right) + A_1 \tilde{G}_{11} \left(10 \zeta_2 \right) + A_1 G_{11} \left(10 \zeta_2 \right) \\
 & \left. + A_1 B_1 \left(10 \zeta_2 L_{qr} - 10 \zeta_2 L_{fr} \right) + A_1^2 \left(\frac{25}{2} \zeta_2^2 \right) \right], \tag{B.2}
 \end{aligned}$$

$$\begin{aligned}
 \bar{g}_{03} = & \left[\tilde{G}_{31} \left(\frac{2}{3} \right) + \tilde{G}_{22} \left(\frac{4}{3} \beta_0 \right) + \tilde{G}_{21} \left(- 2 \beta_0 L_{qr} \right) + \tilde{G}_{13} \left(\frac{8}{3} \beta_0^2 \right) + \tilde{G}_{12} \left(\frac{4}{3} \beta_1 \right. \right. \\
 & \left. \left. - 4 \beta_0^2 L_{qr} \right) + \tilde{G}_{11} \left(- 2 \beta_1 L_{qr} + 2 \beta_0^2 L_{qr}^2 + 8 \beta_0^2 \zeta_2 \right) + \tilde{G}_{11} \tilde{G}_{21} \left(2 \right) \right. \\
 & \left. + \tilde{G}_{11} \tilde{G}_{12} \left(4 \beta_0 \right) + \tilde{G}_{11}^2 \left(- 4 \beta_0 L_{qr} \right) + \tilde{G}_{11}^3 \left(\frac{4}{3} \right) + G_{31} \left(\frac{2}{3} \right) + G_{22} \left(\frac{4}{3} \beta_0 \right) \right]
 \end{aligned}$$

$$\begin{aligned}
& + G_{21} \left(-2 \beta_0 L_{qr} \right) + G_{21} \tilde{G}_{11} \left(2 \right) + G_{13} \left(\frac{8}{3} \beta_0^2 \right) + G_{12} \left(\frac{4}{3} \beta_1 - 4 \beta_0^2 L_{qr} \right) \\
& + G_{12} \tilde{G}_{11} \left(4 \beta_0 \right) + G_{11} \left(-2 \beta_1 L_{qr} + 2 \beta_0^2 L_{qr}^2 - 12 \beta_0^2 \zeta_2 \right) + G_{11} \tilde{G}_{21} \left(2 \right) \\
& + G_{11} \tilde{G}_{12} \left(4 \beta_0 \right) + G_{11} \tilde{G}_{11} \left(-8 \beta_0 L_{qr} \right) + G_{11} \tilde{G}_{11}^2 \left(4 \right) + G_{11} G_{21} \left(2 \right) \\
& + G_{11} G_{12} \left(4 \beta_0 \right) + G_{11}^2 \left(-4 \beta_0 L_{qr} \right) + G_{11}^2 \tilde{G}_{11} \left(4 \right) + G_{11}^3 \left(\frac{4}{3} \right) \\
& + f_2 \left(10 \beta_0 \zeta_2 \right) + f_1 \left(5 \beta_1 \zeta_2 + \frac{16}{3} \beta_0^2 \zeta_3 - 10 \beta_0^2 \zeta_2 L_{qr} \right) + f_1 \tilde{G}_{11} \left(10 \beta_0 \zeta_2 \right) \\
& + f_1 G_{11} \left(10 \beta_0 \zeta_2 \right) + B_3 \left(2 L_{qr} - 2 L_{fr} \right) + B_2 \left(-2 \beta_0 L_{qr}^2 + 2 \beta_0 L_{fr}^2 \right. \\
& \left. + 12 \beta_0 \zeta_2 \right) + B_2 \tilde{G}_{11} \left(4 L_{qr} - 4 L_{fr} \right) + B_2 G_{11} \left(4 L_{qr} - 4 L_{fr} \right) + B_1 \left(\right. \\
& \left. - \beta_1 L_{qr}^2 + \beta_1 L_{fr}^2 + 6 \beta_1 \zeta_2 + \frac{2}{3} \beta_0^2 L_{qr}^3 - \frac{2}{3} \beta_0^2 L_{fr}^3 - 12 \beta_0^2 \zeta_2 L_{qr} \right) \\
& + B_1 \tilde{G}_{21} \left(2 L_{qr} - 2 L_{fr} \right) + B_1 \tilde{G}_{12} \left(4 \beta_0 L_{qr} - 4 \beta_0 L_{fr} \right) + B_1 \tilde{G}_{11} \left(\right. \\
& \left. - 6 \beta_0 L_{qr}^2 + 4 \beta_0 L_{fr} L_{qr} + 2 \beta_0 L_{fr}^2 + 12 \beta_0 \zeta_2 \right) + B_1 \tilde{G}_{11}^2 \left(4 L_{qr} - 4 L_{fr} \right) \\
& + B_1 G_{21} \left(2 L_{qr} - 2 L_{fr} \right) + B_1 G_{12} \left(4 \beta_0 L_{qr} - 4 \beta_0 L_{fr} \right) + B_1 G_{11} \left(\right. \\
& \left. - 6 \beta_0 L_{qr}^2 + 4 \beta_0 L_{fr} L_{qr} + 2 \beta_0 L_{fr}^2 + 12 \beta_0 \zeta_2 \right) + B_1 G_{11} \tilde{G}_{11} \left(8 L_{qr} - 8 L_{fr} \right) \\
& + B_1 G_{11}^2 \left(4 L_{qr} - 4 L_{fr} \right) + B_1 f_1 \left(10 \beta_0 \zeta_2 L_{qr} - 10 \beta_0 \zeta_2 L_{fr} \right) \\
& + B_1 B_2 \left(4 L_{qr}^2 - 8 L_{fr} L_{qr} + 4 L_{fr}^2 \right) + B_1^2 \left(-2 \beta_0 L_{qr}^3 + 2 \beta_0 L_{fr} L_{qr}^2 \right. \\
& \left. + 2 \beta_0 L_{fr}^2 L_{qr} - 2 \beta_0 L_{fr}^3 + 12 \beta_0 \zeta_2 L_{qr} - 12 \beta_0 \zeta_2 L_{fr} \right) + B_1^2 \tilde{G}_{11} \left(4 L_{qr}^2 \right. \\
& \left. - 8 L_{fr} L_{qr} + 4 L_{fr}^2 \right) + B_1^2 G_{11} \left(4 L_{qr}^2 - 8 L_{fr} L_{qr} + 4 L_{fr}^2 \right) + B_1^3 \left(\frac{4}{3} L_{qr}^3 \right. \\
& \left. - 4 L_{fr} L_{qr}^2 + 4 L_{fr}^2 L_{qr} - \frac{4}{3} L_{fr}^3 \right) + A_3 \left(5 \zeta_2 \right) + A_2 \left(\frac{16}{3} \beta_0 \zeta_3 - 10 \beta_0 \zeta_2 L_{qr} \right) \\
& + A_2 \tilde{G}_{11} \left(10 \zeta_2 \right) + A_2 G_{11} \left(10 \zeta_2 \right) + A_2 B_1 \left(10 \zeta_2 L_{qr} - 10 \zeta_2 L_{fr} \right) \\
& + A_1 \left(\frac{8}{3} \beta_1 \zeta_3 - 5 \beta_1 \zeta_2 L_{qr} - \frac{16}{3} \beta_0^2 \zeta_3 L_{qr} + 5 \beta_0^2 \zeta_2 L_{qr}^2 + \frac{21}{5} \beta_0^2 \zeta_2^2 \right) \\
& + A_1 \tilde{G}_{21} \left(5 \zeta_2 \right) + A_1 \tilde{G}_{12} \left(10 \beta_0 \zeta_2 \right) + A_1 \tilde{G}_{11} \left(\frac{16}{3} \beta_0 \zeta_3 - 20 \beta_0 \zeta_2 L_{qr} \right) \\
& + A_1 \tilde{G}_{11}^2 \left(10 \zeta_2 \right) + A_1 G_{21} \left(5 \zeta_2 \right) + A_1 G_{12} \left(10 \beta_0 \zeta_2 \right) + A_1 G_{11} \left(\frac{16}{3} \beta_0 \zeta_3 \right.
\end{aligned}$$

$$\begin{aligned}
 & - 20 \beta_0 \zeta_2 L_{qr} \Big) + A_1 G_{11} \tilde{G}_{11} \left(20 \zeta_2 \right) + A_1 G_{11}^2 \left(10 \zeta_2 \right) + A_1 f_1 \left(25 \beta_0 \zeta_2^2 \right) \\
 & + A_1 B_2 \left(10 \zeta_2 L_{qr} - 10 \zeta_2 L_{fr} \right) + A_1 B_1 \left(\frac{16}{3} \beta_0 \zeta_3 L_{qr} - \frac{16}{3} \beta_0 \zeta_3 L_{fr} \right. \\
 & \left. - 15 \beta_0 \zeta_2 L_{qr}^2 + 10 \beta_0 \zeta_2 L_{fr} L_{qr} + 5 \beta_0 \zeta_2 L_{fr}^2 + 30 \beta_0 \zeta_2^2 \right) \\
 & + A_1 B_1 \tilde{G}_{11} \left(20 \zeta_2 L_{qr} - 20 \zeta_2 L_{fr} \right) + A_1 B_1 G_{11} \left(20 \zeta_2 L_{qr} - 20 \zeta_2 L_{fr} \right) \\
 & + A_1 B_1^2 \left(10 \zeta_2 L_{qr}^2 - 20 \zeta_2 L_{fr} L_{qr} + 10 \zeta_2 L_{fr}^2 \right) + A_1 A_2 \left(25 \zeta_2^2 \right) \\
 & + A_1^2 \left(\frac{40}{3} \beta_0 \zeta_2 \zeta_3 - 25 \beta_0 \zeta_2^2 L_{qr} \right) + A_1^2 \tilde{G}_{11} \left(25 \zeta_2^2 \right) + A_1^2 G_{11} \left(25 \zeta_2^2 \right) \\
 & + A_1^2 B_1 \left(25 \zeta_2^2 L_{qr} - 25 \zeta_2^2 L_{fr} \right) + A_1^3 \left(\frac{125}{6} \zeta_2^3 \right) \Big]. \tag{B.3}
 \end{aligned}$$

The resummed exponent as in eq. (2.32) is calculated to the N³LL accuracy and collected below (with $\omega = 2\beta_0 a_s \ln \bar{N}$),

$$\bar{g}_1 = \left[\frac{A_1}{\beta_0} \left\{ 2 - 2 \ln(1 - \omega) + 2 \ln(1 - \omega) \omega^{-1} \right\} \right], \tag{B.4}$$

$$\begin{aligned}
 \bar{g}_2 = & \left[\frac{D_1}{\beta_0} \left\{ \frac{1}{2} \ln(1 - \omega) \right\} + \frac{A_2}{\beta_0^2} \left\{ - \ln(1 - \omega) - \omega \right\} + \frac{A_1}{\beta_0} \left\{ \left(\ln(1 - \omega) + \frac{1}{2} \ln(1 - \omega)^2 \right. \right. \right. \\
 & \left. \left. + \omega \right) \left(\frac{\beta_1}{\beta_0^2} \right) + \left(\omega \right) L_{fr} + \left(\ln(1 - \omega) \right) L_{qr} \right\} \Big], \tag{B.5}
 \end{aligned}$$

$$\begin{aligned}
 \bar{g}_3 = & \left[\frac{A_3}{\beta_0^2} \left\{ - \frac{\omega}{(1 - \omega)} + \omega \right\} + \frac{A_2}{\beta_0} \left\{ \left(2 \frac{\omega}{(1 - \omega)} \right) L_{qr} + \left(3 \frac{\omega}{(1 - \omega)} + 2 \frac{\ln(1 - \omega)}{(1 - \omega)} \right. \right. \right. \\
 & \left. \left. - \omega \right) \left(\frac{\beta_1}{\beta_0^2} \right) + \left(- 2 \omega \right) L_{fr} \right\} + A_1 \left\{ - 4 \zeta_2 \frac{\omega}{(1 - \omega)} + \left(- \frac{\ln(1 - \omega)^2}{(1 - \omega)} - \frac{\omega}{(1 - \omega)} \right. \right. \\
 & \left. \left. - 2 \frac{\ln(1 - \omega)}{(1 - \omega)} + 2 \ln(1 - \omega) + \omega \right) \left(\frac{\beta_1}{\beta_0^2} \right)^2 + \left(- \frac{\omega}{(1 - \omega)} \right) L_{qr}^2 + \left(- \frac{\omega}{(1 - \omega)} \right. \right. \\
 & \left. \left. - 2 \ln(1 - \omega) - \omega \right) \left(\frac{\beta_2}{\beta_0^3} \right) + \left(\left(- 2 \frac{\omega}{(1 - \omega)} - 2 \frac{\ln(1 - \omega)}{(1 - \omega)} \right) \left(\frac{\beta_1}{\beta_0^2} \right) \right) L_{qr} \right. \\
 & \left. + \left(\omega \right) L_{fr}^2 \right\} + \frac{D_2}{\beta_0} \left\{ \frac{\omega}{(1 - \omega)} \right\} + D_1 \left\{ \left(- \frac{\omega}{(1 - \omega)} \right) L_{qr} + \left(- \frac{\omega}{(1 - \omega)} \right. \right. \\
 & \left. \left. - \frac{\ln(1 - \omega)}{(1 - \omega)} \right) \left(\frac{\beta_1}{\beta_0^2} \right) \right\} \Big], \tag{B.6}
 \end{aligned}$$

$$\begin{aligned}
 \bar{g}_4 = & \left[\frac{A_4}{\beta_0^2} \left\{ \frac{1}{6} \frac{\omega(2 - \omega)}{(1 - \omega)^2} - \frac{1}{3} \omega \right\} + \frac{A_3}{\beta_0} \left\{ \left(- \frac{1}{2} \frac{\omega(2 - \omega)}{(1 - \omega)^2} \right) L_{qr} + \left(- \frac{5}{12} \frac{\omega(2 - \omega)}{(1 - \omega)^2} \right. \right. \right. \\
 & \left. \left. - \frac{1}{2} \frac{\ln(1 - \omega)}{(1 - \omega)^2} + \frac{1}{3} \omega \right) \left(\frac{\beta_1}{\beta_0^2} \right) + \left(\omega \right) L_{fr} \right\} + A_2 \left\{ 2 \zeta_2 \frac{\omega(2 - \omega)}{(1 - \omega)^2} + \left(\frac{1}{2} \frac{\ln(1 - \omega)}{(1 - \omega)^2} \right. \right. \\
 & \left. \left. - \frac{1}{12} \frac{\omega^2}{(1 - \omega)^2} + \frac{5}{6} \frac{\omega}{(1 - \omega)} + \frac{1}{2} \frac{\ln(1 - \omega)}{(1 - \omega)^2} - \frac{1}{3} \omega \right) \left(\frac{\beta_1}{\beta_0^2} \right)^2 + \left(\frac{1}{2} \frac{\omega(2 - \omega)}{(1 - \omega)^2} \right) L_{qr}^2 \right. \\
 & \left. + \left(- \frac{\omega}{(1 - \omega)} \right) L_{qr} \right\} \Big].
 \end{aligned}$$

$$\begin{aligned}
 & + \left(\frac{1}{3} \frac{\omega^2}{(1-\omega)^2} - \frac{1}{3} \frac{\omega}{(1-\omega)} + \frac{1}{3} \omega \right) \left(\frac{\beta_2}{\beta_0^3} \right) + \left(-\omega \right) L_{fr}^2 + \left(\left(\frac{1}{2} \frac{\omega(2-\omega)}{(1-\omega)^2} \right. \right. \\
 & + \left. \left. \frac{\ln(1-\omega)}{(1-\omega)^2} \right) \left(\frac{\beta_1}{\beta_0^2} \right) \right) L_{qr} \Big\} + \beta_0 A_1 \left\{ \frac{8}{3} \zeta_3 \frac{\omega(2-\omega)}{(1-\omega)^2} + \left(-\frac{1}{6} \frac{\ln(1-\omega)^3}{(1-\omega)^2} \right. \right. \\
 & + \left. \frac{1}{3} \frac{\omega^2}{(1-\omega)^2} - \frac{1}{3} \frac{\omega}{(1-\omega)} + \frac{1}{2} \frac{\ln(1-\omega)}{(1-\omega)^2} - \frac{\ln(1-\omega)}{(1-\omega)} + \frac{1}{2} \ln(1-\omega) \right. \\
 & + \left. \frac{1}{3} \omega \right) \left(\frac{\beta_1}{\beta_0^2} \right)^3 + \left(-\frac{1}{6} \frac{\omega(2-\omega)}{(1-\omega)^2} \right) L_{qr}^3 + \left(\frac{1}{12} \frac{\omega(2-\omega)}{(1-\omega)^2} + \frac{1}{2} \ln(1-\omega) \right. \\
 & + \left. \frac{1}{3} \omega \right) \left(\frac{\beta_3}{\beta_0^4} \right) + \left(-\frac{5}{12} \frac{\omega^2}{(1-\omega)^2} + \frac{1}{6} \frac{\omega}{(1-\omega)} - \frac{1}{2} \frac{\ln(1-\omega)}{(1-\omega)^2} + \frac{\ln(1-\omega)}{(1-\omega)} \right. \\
 & - \ln(1-\omega) - \frac{2}{3} \omega \Big) \frac{\beta_1 \beta_2}{\beta_0^5} + \left(\frac{1}{3} \omega \right) L_{fr}^3 + \left(-2 \zeta_2 \frac{\omega(2-\omega)}{(1-\omega)^2} + \left(-\frac{1}{2} \frac{\ln(1-\omega)^2}{(1-\omega)^2} \right. \right. \\
 & + \left. \frac{1}{2} \frac{\omega^2}{(1-\omega)^2} \right) \left(\frac{\beta_1}{\beta_0^2} \right)^2 + \left(-\frac{1}{2} \frac{\omega^2}{(1-\omega)^2} \right) \left(\frac{\beta_2}{\beta_0^3} \right) \Big\} L_{qr} + \left(\right. \\
 & - 2 \zeta_2 \frac{\ln(1-\omega)}{(1-\omega)^2} \Big) \left(\frac{\beta_1}{\beta_0^2} \right) + \left(\left(-\frac{1}{2} \frac{\ln(1-\omega)}{(1-\omega)^2} \right) \left(\frac{\beta_1}{\beta_0^2} \right) \right) L_{qr}^2 + \left(\left(\right. \right. \\
 & - \left. \frac{1}{2} \omega \right) \left(\frac{\beta_1}{\beta_0^2} \right) \Big) L_{fr}^2 \Big\} + \frac{D_3}{\beta_0} \left\{ -\frac{1}{4} \frac{\omega(2-\omega)}{(1-\omega)^2} \right\} + D_2 \left\{ \left(\frac{1}{4} \frac{\omega(2-\omega)}{(1-\omega)^2} \right. \right. \\
 & + \left. \frac{1}{2} \frac{\ln(1-\omega)}{(1-\omega)^2} \right) \left(\frac{\beta_1}{\beta_0^2} \right) + \left(\frac{1}{2} \frac{\omega(2-\omega)}{(1-\omega)^2} \right) L_{qr} \Big\} + \beta_0 D_1 \left\{ -\zeta_2 \frac{\omega(2-\omega)}{(1-\omega)^2} + \left(\right. \right. \\
 & - \left. \frac{1}{4} \frac{\ln(1-\omega)^2}{(1-\omega)^2} + \frac{1}{4} \frac{\omega^2}{(1-\omega)^2} \right) \left(\frac{\beta_1}{\beta_0^2} \right)^2 + \left(-\frac{1}{4} \frac{\omega(2-\omega)}{(1-\omega)^2} \right) L_{qr}^2 + \left(\right. \\
 & - \left. \frac{1}{4} \frac{\omega^2}{(1-\omega)^2} \right) \left(\frac{\beta_2}{\beta_0^3} \right) + \left(\left(-\frac{1}{2} \frac{\ln(1-\omega)}{(1-\omega)^2} \right) \left(\frac{\beta_1}{\beta_0^2} \right) \right) L_{qr} \Big\} \Big]. \tag{B.7}
 \end{aligned}$$

All the anomalous dimensions can be found in appendix C.

B.2 Resummation ingredients for the Standard N exponentiation

Below we present the resummed exponent for the Standard N -exponent as given in eq. (2.35).

$$g_1 = \left[\tilde{A}_1 \left\{ \lambda^{-1} (\ln \lambda') (2) + (\ln \lambda') (-2) + 2 \right\} \right], \tag{B.8}$$

$$\begin{aligned}
 g_2 = & \left[\tilde{D}_1 \left\{ (\ln \lambda') \left(\frac{1}{2} \right) \right\} + \tilde{A}_2 \left\{ (\ln \lambda') (-1) + \lambda (-1) \right\} + \tilde{A}_1 \left\{ (\ln \lambda') (\tilde{\beta}_1 \right. \right. \\
 & \left. \left. - 2 \gamma_E + L_{qr} \right) + (\ln \lambda')^2 \left(\frac{1}{2} \tilde{\beta}_1 \right) + \lambda \left(\tilde{\beta}_1 + L_{fr} \right) \right\} \Big], \tag{B.9}
 \end{aligned}$$

$$\begin{aligned}
 g_3 = & \beta_0 \left[\tilde{D}_2 \left\{ \frac{\lambda}{\lambda'} \left(-\frac{1}{2} \right) \right\} + \tilde{D}_1 \left\{ (\ln \lambda') \left(\frac{1}{\lambda'} \right) \left(\frac{1}{2} \tilde{\beta}_1 \right) + \frac{\lambda}{\lambda'} \left(\frac{1}{2} \tilde{\beta}_1 - \gamma_E + \frac{1}{2} L_{qr} \right) \right\} \right. \\
 & + \tilde{A}_3 \left\{ \frac{\lambda^2}{(\lambda')^2} \left(\frac{1}{2} \right) \right\} + \tilde{A}_2 \left\{ \frac{\lambda}{\lambda'} \left(-\frac{3}{2} \tilde{\beta}_1 + 2 \gamma_E - L_{qr} \right) + (\ln \lambda') \left(\frac{1}{\lambda'} \right) \left(-\tilde{\beta}_1 \right) \right. \\
 & \left. \left. + \lambda \left(\frac{1}{2} \tilde{\beta}_1 + L_{fr} \right) \right\} + \tilde{A}_1 \left\{ \frac{\lambda}{\lambda'} \left(\frac{1}{2} \tilde{\beta}_2 + \frac{1}{2} \tilde{\beta}_1^2 - 2 \gamma_E \tilde{\beta}_1 + 2 \gamma_E^2 + L_{qr} \tilde{\beta}_1 \right) \right\} \right]
 \end{aligned}$$

$$\begin{aligned}
 & -2 L_{qr} \gamma_E + \frac{1}{2} L_{qr}^2 + 2 \zeta_2) + (\ln \lambda') \left(\tilde{\beta}_2 - 2 \gamma_E \tilde{\beta}_1 + L_{qr} \tilde{\beta}_1 \right) + (\ln \lambda') \frac{\lambda}{\lambda'} \left(\tilde{\beta}_1^2 \right. \\
 & \left. - 2 \gamma_E \tilde{\beta}_1 + L_{qr} \tilde{\beta}_1 \right) + (\ln \lambda')^2 \left(\frac{1}{\lambda'} \right) \left(\frac{1}{2} \tilde{\beta}_1^2 \right) + \lambda \left(\frac{1}{2} \tilde{\beta}_2 - \frac{1}{2} \tilde{\beta}_1^2 - \frac{1}{2} L_{fr}^2 \right) \left. \right\}, \tag{B.10}
 \end{aligned}$$

$$\begin{aligned}
 g_4 = \beta_0^2 & \left[\tilde{D}_3 \left\{ \frac{\lambda(2-\lambda)}{(\lambda')^2} \left(-\frac{1}{4} \right) \right\} + \tilde{D}_2 \left\{ \frac{\lambda(2-\lambda)}{(\lambda')^2} \left(\frac{1}{4} \tilde{\beta}_1 - \gamma_E + \frac{1}{2} L_{qr} \right) \right. \right. \\
 & + (\ln \lambda') \left(\frac{1}{\lambda'} \right)^2 \left(\frac{1}{2} \tilde{\beta}_1 \right) \left. \right\} + \tilde{D}_1 \left\{ \frac{\lambda(2-\lambda)}{(\lambda')^2} \left(-\frac{1}{4} \tilde{\beta}_2 + \frac{1}{4} \tilde{\beta}_1^2 - \gamma_E^2 + L_{qr} \gamma_E \right. \right. \\
 & - \frac{1}{4} L_{qr}^2 - \zeta_2) + \frac{\lambda}{\lambda'} \left(\frac{1}{2} \tilde{\beta}_2 - \frac{1}{2} \tilde{\beta}_1^2 \right) + (\ln \lambda') \left(\frac{1}{\lambda'} \right)^2 \left(\gamma_E \tilde{\beta}_1 - \frac{1}{2} L_{qr} \tilde{\beta}_1 \right) \\
 & + (\ln \lambda')^2 \left(\frac{1}{\lambda'} \right)^2 \left(-\frac{1}{4} \tilde{\beta}_1^2 \right) \left. \right\} + \tilde{A}_4 \left\{ \frac{\lambda(2-\lambda)}{(\lambda')^2} \left(\frac{1}{6} \right) + \lambda \left(-\frac{1}{3} \right) \right\} \\
 & + \tilde{A}_3 \left\{ \frac{\lambda(2-\lambda)}{(\lambda')^2} \left(-\frac{5}{12} \tilde{\beta}_1 + \gamma_E - \frac{1}{2} L_{qr} \right) + (\ln \lambda') \left(\frac{1}{\lambda'} \right)^2 \left(-\frac{1}{2} \tilde{\beta}_1 \right) \right. \\
 & + \lambda \left(\frac{1}{3} \tilde{\beta}_1 + L_{fr} \right) \left. \right\} + \tilde{A}_2 \left\{ \frac{\lambda(2-\lambda)}{(\lambda')^2} \left(\frac{1}{3} \tilde{\beta}_2 - \frac{1}{12} \tilde{\beta}_1^2 - \gamma_E \tilde{\beta}_1 + 2 \gamma_E^2 \right. \right. \\
 & + \frac{1}{2} L_{qr} \tilde{\beta}_1 - 2 L_{qr} \gamma_E + \frac{1}{2} L_{qr}^2 + 2 \zeta_2) + \frac{\lambda}{\lambda'} \left(-\tilde{\beta}_2 + \tilde{\beta}_1^2 \right) \\
 & + (\ln \lambda') \left(\frac{1}{\lambda'} \right)^2 \left(\frac{1}{2} \tilde{\beta}_1^2 - 2 \gamma_E \tilde{\beta}_1 + L_{qr} \tilde{\beta}_1 \right) + (\ln \lambda')^2 \left(\frac{1}{\lambda'} \right)^2 \left(\frac{1}{2} \tilde{\beta}_1^2 \right) \\
 & + \lambda \left(\frac{1}{3} \tilde{\beta}_2 - \frac{1}{3} \tilde{\beta}_1^2 - L_{fr}^2 \right) \left. \right\} + \tilde{A}_1 \left\{ \frac{\lambda(2-\lambda)}{(\lambda')^2} \left(\frac{1}{12} \tilde{\beta}_3 - \frac{5}{12} \tilde{\beta}_1 \tilde{\beta}_2 + \frac{1}{3} \tilde{\beta}_1^3 \right. \right. \\
 & + \gamma_E \tilde{\beta}_2 - \gamma_E \tilde{\beta}_1^2 + \frac{4}{3} \gamma_E^3 - \frac{1}{2} L_{qr} \tilde{\beta}_2 + \frac{1}{2} L_{qr} \tilde{\beta}_1^2 - 2 L_{qr} \gamma_E^2 + L_{qr}^2 \gamma_E - \frac{1}{6} L_{qr}^3 \\
 & + \frac{8}{3} \zeta_3 + 4 \zeta_2 \gamma_E - 2 \zeta_2 L_{qr} \left. \right) + \frac{\lambda}{\lambda'} \left(\tilde{\beta}_1 \tilde{\beta}_2 - \tilde{\beta}_1^3 - 2 \gamma_E \tilde{\beta}_2 + 2 \gamma_E \tilde{\beta}_1^2 + L_{qr} \tilde{\beta}_2 \right. \\
 & - L_{qr} \tilde{\beta}_1^2) + (\ln \lambda') \left(\frac{1}{2} \tilde{\beta}_3 - \frac{1}{2} \tilde{\beta}_1^3 \right) + (\ln \lambda') \frac{\lambda}{\lambda'} \left(\tilde{\beta}_1 \tilde{\beta}_2 - \tilde{\beta}_1^3 \right) \\
 & + (\ln \lambda') \left(\frac{1}{\lambda'} \right)^2 \left(-\frac{1}{2} \tilde{\beta}_1 \tilde{\beta}_2 + \frac{1}{2} \tilde{\beta}_1^3 - 2 \gamma_E^2 \tilde{\beta}_1 + 2 L_{qr} \gamma_E \tilde{\beta}_1 - \frac{1}{2} L_{qr}^2 \tilde{\beta}_1 \right. \\
 & - 2 \zeta_2 \tilde{\beta}_1) + (\ln \lambda')^2 \left(\frac{1}{\lambda'} \right)^2 \left(\gamma_E \tilde{\beta}_1^2 - \frac{1}{2} L_{qr} \tilde{\beta}_1^2 \right) + (\ln \lambda')^3 \left(\frac{1}{\lambda'} \right)^2 \left(-\frac{1}{6} \tilde{\beta}_1^3 \right) \\
 & \left. + \lambda \left(\frac{1}{3} \tilde{\beta}_3 - \frac{2}{3} \tilde{\beta}_1 \tilde{\beta}_2 + \frac{1}{3} \tilde{\beta}_1^3 - \frac{1}{2} L_{fr}^2 \tilde{\beta}_1 + \frac{1}{3} L_{fr}^3 \right) \right\} \left. \right], \tag{B.11}
 \end{aligned}$$

with $\lambda' = (1 - \lambda)$, where $\lambda = 2a_s \beta_0 \ln N$. $\tilde{A}_i = A_i / \beta_0^i$, $\tilde{D}_i = D_i / \beta_0^i$, $\tilde{\beta}_i = \beta_i / \beta_0^{i+1}$. The constants g_0 are given by

$$\begin{aligned}
 g_{01} = & \left[\tilde{G}_{11} \left(2 \right) + G_{11} \left(2 \right) + f_1 \left(2 \gamma_E \right) + B_1 \left(2 L_{qr} - 2 L_{fr} \right) + A_1 \left(2 \gamma_E^2 - 2 L_{qr} \gamma_E \right. \right. \\
 & \left. \left. + 2 L_{fr} \gamma_E + 5 \zeta_2 \right) \right], \tag{B.12}
 \end{aligned}$$

$$\begin{aligned}
g_{02} = & \left[\tilde{G}_{21} \left(1 \right) + \tilde{G}_{12} \left(2 \beta_0 \right) + \tilde{G}_{11} \left(4 \beta_0 \gamma_E - 2 \beta_0 L_{qr} \right) + \tilde{G}_{11}^2 \left(2 \right) + G_{21} \left(1 \right) \right. \\
& + G_{12} \left(2 \beta_0 \right) + G_{11} \left(- 2 \beta_0 L_{qr} \right) + G_{11} \tilde{G}_{11} \left(4 \right) + G_{11}^2 \left(2 \right) + f_2 \left(2 \gamma_E \right) \\
& + f_1 \left(2 \beta_0 \gamma_E^2 - 2 \beta_0 L_{qr} \gamma_E + 5 \beta_0 \zeta_2 \right) + f_1 \tilde{G}_{11} \left(4 \gamma_E \right) + f_1 G_{11} \left(4 \gamma_E \right) \\
& + f_1^2 \left(2 \gamma_E^2 \right) + B_2 \left(2 L_{qr} - 2 L_{fr} \right) + B_1 \left(- \beta_0 L_{qr}^2 + \beta_0 L_{fr}^2 + 6 \beta_0 \zeta_2 \right) \\
& + B_1 \tilde{G}_{11} \left(4 L_{qr} - 4 L_{fr} \right) + B_1 G_{11} \left(4 L_{qr} - 4 L_{fr} \right) + B_1 f_1 \left(4 L_{qr} \gamma_E \right. \\
& \left. - 4 L_{fr} \gamma_E \right) + B_1^2 \left(2 L_{qr}^2 - 4 L_{fr} L_{qr} + 2 L_{fr}^2 \right) + A_2 \left(2 \gamma_E^2 - 2 L_{qr} \gamma_E \right. \\
& \left. + 2 L_{fr} \gamma_E + 5 \zeta_2 \right) + A_1 \left(\frac{4}{3} \beta_0 \gamma_E^3 - 2 \beta_0 L_{qr} \gamma_E^2 + \beta_0 L_{qr}^2 \gamma_E - \beta_0 L_{fr}^2 \gamma_E \right. \\
& \left. + \frac{8}{3} \beta_0 \zeta_3 + 4 \beta_0 \zeta_2 \gamma_E - 5 \beta_0 \zeta_2 L_{qr} \right) + A_1 \tilde{G}_{11} \left(4 \gamma_E^2 - 4 L_{qr} \gamma_E + 4 L_{fr} \gamma_E \right. \\
& \left. + 10 \zeta_2 \right) + A_1 G_{11} \left(4 \gamma_E^2 - 4 L_{qr} \gamma_E + 4 L_{fr} \gamma_E + 10 \zeta_2 \right) + A_1 f_1 \left(4 \gamma_E^3 \right. \\
& \left. - 4 L_{qr} \gamma_E^2 + 4 L_{fr} \gamma_E^2 + 10 \zeta_2 \gamma_E \right) + A_1 B_1 \left(4 L_{qr} \gamma_E^2 - 4 L_{qr}^2 \gamma_E - 4 L_{fr} \gamma_E^2 \right. \\
& \left. + 8 L_{fr} L_{qr} \gamma_E - 4 L_{fr}^2 \gamma_E + 10 \zeta_2 L_{qr} - 10 \zeta_2 L_{fr} \right) + A_1^2 \left(2 \gamma_E^4 - 4 L_{qr} \gamma_E^3 \right. \\
& \left. + 2 L_{qr}^2 \gamma_E^2 + 4 L_{fr} \gamma_E^3 - 4 L_{fr} L_{qr} \gamma_E^2 + 2 L_{fr}^2 \gamma_E^2 + 10 \zeta_2 \gamma_E^2 - 10 \zeta_2 L_{qr} \gamma_E \right. \\
& \left. + 10 \zeta_2 L_{fr} \gamma_E + \frac{25}{2} \zeta_2^2 \right) \left. \right], \tag{B.13}
\end{aligned}$$

$$\begin{aligned}
g_{03} = & \left[\tilde{G}_{31} \left(\frac{2}{3} \right) + \tilde{G}_{22} \left(\frac{4}{3} \beta_0 \right) + \tilde{G}_{21} \left(4 \beta_0 \gamma_E - 2 \beta_0 L_{qr} \right) + \tilde{G}_{13} \left(\frac{8}{3} \beta_0^2 \right) \right. \\
& + \tilde{G}_{12} \left(\frac{4}{3} \beta_1 + 8 \beta_0^2 \gamma_E - 4 \beta_0^2 L_{qr} \right) + \tilde{G}_{11} \left(4 \beta_1 \gamma_E - 2 \beta_1 L_{qr} + 8 \beta_0^2 \gamma_E^2 \right. \\
& \left. - 8 \beta_0^2 L_{qr} \gamma_E + 2 \beta_0^2 L_{qr}^2 + 8 \beta_0^2 \zeta_2 \right) + \tilde{G}_{11} \tilde{G}_{21} \left(2 \right) + \tilde{G}_{11} \tilde{G}_{12} \left(4 \beta_0 \right) \\
& + \tilde{G}_{11}^2 \left(8 \beta_0 \gamma_E - 4 \beta_0 L_{qr} \right) + \tilde{G}_{11}^3 \left(\frac{4}{3} \right) + G_{31} \left(\frac{2}{3} \right) + G_{22} \left(\frac{4}{3} \beta_0 \right) + G_{21} \left(\right. \\
& \left. - 2 \beta_0 L_{qr} \right) + G_{21} \tilde{G}_{11} \left(2 \right) + G_{13} \left(\frac{8}{3} \beta_0^2 \right) + G_{12} \left(\frac{4}{3} \beta_1 - 4 \beta_0^2 L_{qr} \right) \\
& + G_{12} \tilde{G}_{11} \left(4 \beta_0 \right) + G_{11} \left(- 2 \beta_1 L_{qr} + 2 \beta_0^2 L_{qr}^2 - 12 \beta_0^2 \zeta_2 \right) + G_{11} \tilde{G}_{21} \left(2 \right) \\
& + G_{11} \tilde{G}_{12} \left(4 \beta_0 \right) + G_{11} \tilde{G}_{11} \left(8 \beta_0 \gamma_E - 8 \beta_0 L_{qr} \right) + G_{11} \tilde{G}_{11}^2 \left(4 \right) \\
& \left. + G_{11} G_{21} \left(2 \right) + G_{11} G_{12} \left(4 \beta_0 \right) + G_{11}^2 \left(- 4 \beta_0 L_{qr} \right) + G_{11}^2 \tilde{G}_{11} \left(4 \right) \right]
\end{aligned}$$

$$\begin{aligned}
& + G_{11}^3 \left(\frac{4}{3} \right) + f_3 \left(2 \gamma_E \right) + f_2 \left(4 \beta_0 \gamma_E^2 - 4 \beta_0 L_{qr} \gamma_E + 10 \beta_0 \zeta_2 \right) \\
& + f_2 \tilde{G}_{11} \left(4 \gamma_E \right) + f_2 G_{11} \left(4 \gamma_E \right) + f_1 \left(2 \beta_1 \gamma_E^2 - 2 \beta_1 L_{qr} \gamma_E + 5 \beta_1 \zeta_2 \right. \\
& + \left. \frac{8}{3} \beta_0^2 \gamma_E^3 - 4 \beta_0^2 L_{qr} \gamma_E^2 + 2 \beta_0^2 L_{qr}^2 \gamma_E + \frac{16}{3} \beta_0^2 \zeta_3 + 8 \beta_0^2 \zeta_2 \gamma_E - 10 \beta_0^2 \zeta_2 L_{qr} \right) \\
& + f_1 \tilde{G}_{21} \left(2 \gamma_E \right) + f_1 \tilde{G}_{12} \left(4 \beta_0 \gamma_E \right) + f_1 \tilde{G}_{11} \left(12 \beta_0 \gamma_E^2 - 8 \beta_0 L_{qr} \gamma_E \right. \\
& + \left. 10 \beta_0 \zeta_2 \right) + f_1 \tilde{G}_{11}^2 \left(4 \gamma_E \right) + f_1 G_{21} \left(2 \gamma_E \right) + f_1 G_{12} \left(4 \beta_0 \gamma_E \right) \\
& + f_1 G_{11} \left(4 \beta_0 \gamma_E^2 - 8 \beta_0 L_{qr} \gamma_E + 10 \beta_0 \zeta_2 \right) + f_1 G_{11} \tilde{G}_{11} \left(8 \gamma_E \right) \\
& + f_1 G_{11}^2 \left(4 \gamma_E \right) + f_1 f_2 \left(4 \gamma_E^2 \right) + f_1^2 \left(4 \beta_0 \gamma_E^3 - 4 \beta_0 L_{qr} \gamma_E^2 + 10 \beta_0 \zeta_2 \gamma_E \right) \\
& + f_1^2 \tilde{G}_{11} \left(4 \gamma_E^2 \right) + f_1^2 G_{11} \left(4 \gamma_E^2 \right) + f_1^3 \left(\frac{4}{3} \gamma_E^3 \right) + B_3 \left(2 L_{qr} - 2 L_{fr} \right) + B_2 \left(\right. \\
& - \left. 2 \beta_0 L_{qr}^2 + 2 \beta_0 L_{fr}^2 + 12 \beta_0 \zeta_2 \right) + B_2 \tilde{G}_{11} \left(4 L_{qr} - 4 L_{fr} \right) + B_2 G_{11} \left(4 L_{qr} \right. \\
& - \left. 4 L_{fr} \right) + B_2 f_1 \left(4 L_{qr} \gamma_E - 4 L_{fr} \gamma_E \right) + B_1 \left(-\beta_1 L_{qr}^2 + \beta_1 L_{fr}^2 + 6 \beta_1 \zeta_2 \right. \\
& + \left. \frac{2}{3} \beta_0^2 L_{qr}^3 - \frac{2}{3} \beta_0^2 L_{fr}^3 - 12 \beta_0^2 \zeta_2 L_{qr} \right) + B_1 \tilde{G}_{21} \left(2 L_{qr} - 2 L_{fr} \right) \\
& + B_1 \tilde{G}_{12} \left(4 \beta_0 L_{qr} - 4 \beta_0 L_{fr} \right) + B_1 \tilde{G}_{11} \left(8 \beta_0 L_{qr} \gamma_E - 6 \beta_0 L_{qr}^2 - 8 \beta_0 L_{fr} \gamma_E \right. \\
& + \left. 4 \beta_0 L_{fr} L_{qr} + 2 \beta_0 L_{fr}^2 + 12 \beta_0 \zeta_2 \right) + B_1 \tilde{G}_{11}^2 \left(4 L_{qr} - 4 L_{fr} \right) \\
& + B_1 G_{21} \left(2 L_{qr} - 2 L_{fr} \right) + B_1 G_{12} \left(4 \beta_0 L_{qr} - 4 \beta_0 L_{fr} \right) + B_1 G_{11} \left(\right. \\
& - \left. 6 \beta_0 L_{qr}^2 + 4 \beta_0 L_{fr} L_{qr} + 2 \beta_0 L_{fr}^2 + 12 \beta_0 \zeta_2 \right) + B_1 G_{11} \tilde{G}_{11} \left(8 L_{qr} - 8 L_{fr} \right) \\
& + B_1 G_{11}^2 \left(4 L_{qr} - 4 L_{fr} \right) + B_1 f_2 \left(4 L_{qr} \gamma_E - 4 L_{fr} \gamma_E \right) + B_1 f_1 \left(4 \beta_0 L_{qr} \gamma_E^2 \right. \\
& - \left. 6 \beta_0 L_{qr}^2 \gamma_E - 4 \beta_0 L_{fr} \gamma_E^2 + 4 \beta_0 L_{fr} L_{qr} \gamma_E + 2 \beta_0 L_{fr}^2 \gamma_E + 12 \beta_0 \zeta_2 \gamma_E \right. \\
& + \left. 10 \beta_0 \zeta_2 L_{qr} - 10 \beta_0 \zeta_2 L_{fr} \right) + B_1 f_1 \tilde{G}_{11} \left(8 L_{qr} \gamma_E - 8 L_{fr} \gamma_E \right) \\
& + B_1 f_1 G_{11} \left(8 L_{qr} \gamma_E - 8 L_{fr} \gamma_E \right) + B_1 f_1^2 \left(4 L_{qr} \gamma_E^2 - 4 L_{fr} \gamma_E^2 \right) \\
& + B_1 B_2 \left(4 L_{qr}^2 - 8 L_{fr} L_{qr} + 4 L_{fr}^2 \right) + B_1^2 \left(-2 \beta_0 L_{qr}^3 + 2 \beta_0 L_{fr} L_{qr}^2 \right. \\
& + \left. 2 \beta_0 L_{fr}^2 L_{qr} - 2 \beta_0 L_{fr}^3 + 12 \beta_0 \zeta_2 L_{qr} - 12 \beta_0 \zeta_2 L_{fr} \right) + B_1^2 \tilde{G}_{11} \left(4 L_{qr}^2 \right.
\end{aligned}$$

$$\begin{aligned}
 & - 8 L_{fr} L_{qr} + 4 L_{fr}^2) + B_1^2 G_{11} \left(4 L_{qr}^2 - 8 L_{fr} L_{qr} + 4 L_{fr}^2 \right) + B_1^2 f_1 \left(4 L_{qr}^2 \gamma_E \right. \\
 & \left. - 8 L_{fr} L_{qr} \gamma_E + 4 L_{fr}^2 \gamma_E \right) + B_1^3 \left(\frac{4}{3} L_{qr}^3 - 4 L_{fr} L_{qr}^2 + 4 L_{fr}^2 L_{qr} - \frac{4}{3} L_{fr}^3 \right) \\
 & + A_3 \left(2 \gamma_E^2 - 2 L_{qr} \gamma_E + 2 L_{fr} \gamma_E + 5 \zeta_2 \right) + A_2 \left(\frac{8}{3} \beta_0 \gamma_E^3 - 4 \beta_0 L_{qr} \gamma_E^2 \right. \\
 & \left. + 2 \beta_0 L_{qr}^2 \gamma_E - 2 \beta_0 L_{fr}^2 \gamma_E + \frac{16}{3} \beta_0 \zeta_3 + 8 \beta_0 \zeta_2 \gamma_E - 10 \beta_0 \zeta_2 L_{qr} \right) \\
 & + A_2 \tilde{G}_{11} \left(4 \gamma_E^2 - 4 L_{qr} \gamma_E + 4 L_{fr} \gamma_E + 10 \zeta_2 \right) + A_2 G_{11} \left(4 \gamma_E^2 - 4 L_{qr} \gamma_E \right. \\
 & \left. + 4 L_{fr} \gamma_E + 10 \zeta_2 \right) + A_2 f_1 \left(4 \gamma_E^3 - 4 L_{qr} \gamma_E^2 + 4 L_{fr} \gamma_E^2 + 10 \zeta_2 \gamma_E \right) \\
 & + A_2 B_1 \left(4 L_{qr} \gamma_E^2 - 4 L_{qr}^2 \gamma_E - 4 L_{fr} \gamma_E^2 + 8 L_{fr} L_{qr} \gamma_E - 4 L_{fr}^2 \gamma_E + 10 \zeta_2 L_{qr} \right. \\
 & \left. - 10 \zeta_2 L_{fr} \right) + A_1 \left(\frac{4}{3} \beta_1 \gamma_E^3 - 2 \beta_1 L_{qr} \gamma_E^2 + \beta_1 L_{qr}^2 \gamma_E - \beta_1 L_{fr}^2 \gamma_E + \frac{8}{3} \beta_1 \zeta_3 \right. \\
 & \left. + 4 \beta_1 \zeta_2 \gamma_E - 5 \beta_1 \zeta_2 L_{qr} + \frac{4}{3} \beta_0^2 \gamma_E^4 - \frac{8}{3} \beta_0^2 L_{qr} \gamma_E^3 + 2 \beta_0^2 L_{qr}^2 \gamma_E^2 - \frac{2}{3} \beta_0^2 L_{qr}^3 \gamma_E \right. \\
 & \left. + \frac{2}{3} \beta_0^2 L_{fr}^3 \gamma_E + \frac{32}{3} \beta_0^2 \zeta_3 \gamma_E - \frac{16}{3} \beta_0^2 \zeta_3 L_{qr} + 8 \beta_0^2 \zeta_2 \gamma_E^2 - 8 \beta_0^2 \zeta_2 L_{qr} \gamma_E \right. \\
 & \left. + 5 \beta_0^2 \zeta_2 L_{qr}^2 + \frac{21}{5} \beta_0^2 \zeta_2^2 \right) + A_1 \tilde{G}_{21} \left(2 \gamma_E^2 - 2 L_{qr} \gamma_E + 2 L_{fr} \gamma_E + 5 \zeta_2 \right) \\
 & + A_1 \tilde{G}_{12} \left(4 \beta_0 \gamma_E^2 - 4 \beta_0 L_{qr} \gamma_E + 4 \beta_0 L_{fr} \gamma_E + 10 \beta_0 \zeta_2 \right) + A_1 \tilde{G}_{11} \left(\frac{32}{3} \beta_0 \gamma_E^3 \right. \\
 & \left. - 16 \beta_0 L_{qr} \gamma_E^2 + 6 \beta_0 L_{qr}^2 \gamma_E + 8 \beta_0 L_{fr} \gamma_E^2 - 4 \beta_0 L_{fr} L_{qr} \gamma_E - 2 \beta_0 L_{fr}^2 \gamma_E \right. \\
 & \left. + \frac{16}{3} \beta_0 \zeta_3 + 28 \beta_0 \zeta_2 \gamma_E - 20 \beta_0 \zeta_2 L_{qr} \right) + A_1 \tilde{G}_{11}^2 \left(4 \gamma_E^2 - 4 L_{qr} \gamma_E + 4 L_{fr} \gamma_E \right. \\
 & \left. + 10 \zeta_2 \right) + A_1 G_{21} \left(2 \gamma_E^2 - 2 L_{qr} \gamma_E + 2 L_{fr} \gamma_E + 5 \zeta_2 \right) + A_1 G_{12} \left(4 \beta_0 \gamma_E^2 \right. \\
 & \left. - 4 \beta_0 L_{qr} \gamma_E + 4 \beta_0 L_{fr} \gamma_E + 10 \beta_0 \zeta_2 \right) + A_1 G_{11} \left(\frac{8}{3} \beta_0 \gamma_E^3 - 8 \beta_0 L_{qr} \gamma_E^2 \right. \\
 & \left. + 6 \beta_0 L_{qr}^2 \gamma_E - 4 \beta_0 L_{fr} L_{qr} \gamma_E - 2 \beta_0 L_{fr}^2 \gamma_E + \frac{16}{3} \beta_0 \zeta_3 + 8 \beta_0 \zeta_2 \gamma_E \right. \\
 & \left. - 20 \beta_0 \zeta_2 L_{qr} \right) + A_1 G_{11} \tilde{G}_{11} \left(8 \gamma_E^2 - 8 L_{qr} \gamma_E + 8 L_{fr} \gamma_E + 20 \zeta_2 \right) \\
 & + A_1 G_{11}^2 \left(4 \gamma_E^2 - 4 L_{qr} \gamma_E + 4 L_{fr} \gamma_E + 10 \zeta_2 \right) + A_1 f_2 \left(4 \gamma_E^3 - 4 L_{qr} \gamma_E^2 \right. \\
 & \left. + 4 L_{fr} \gamma_E^2 + 10 \zeta_2 \gamma_E \right) + A_1 f_1 \left(\frac{20}{3} \beta_0 \gamma_E^4 - 12 \beta_0 L_{qr} \gamma_E^3 + 6 \beta_0 L_{qr}^2 \gamma_E^2 \right. \\
 & \left. + 4 \beta_0 L_{fr} \gamma_E^3 - 4 \beta_0 L_{fr} L_{qr} \gamma_E^2 - 2 \beta_0 L_{fr}^2 \gamma_E^2 + \frac{16}{3} \beta_0 \zeta_3 \gamma_E + 28 \beta_0 \zeta_2 \gamma_E^2 \right. \\
 & \left. - 30 \beta_0 \zeta_2 L_{qr} \gamma_E + 10 \beta_0 \zeta_2 L_{fr} \gamma_E + 25 \beta_0 \zeta_2^2 \right) + A_1 f_1 \tilde{G}_{11} \left(8 \gamma_E^3 - 8 L_{qr} \gamma_E^2 \right.
 \end{aligned}$$

$$\begin{aligned}
 & + 8 L_{fr} \gamma_E^2 + 20 \zeta_2 \gamma_E \Big) + A_1 f_1 G_{11} \left(8 \gamma_E^3 - 8 L_{qr} \gamma_E^2 + 8 L_{fr} \gamma_E^2 + 20 \zeta_2 \gamma_E \right) \\
 & + A_1 f_1^2 \left(4 \gamma_E^4 - 4 L_{qr} \gamma_E^3 + 4 L_{fr} \gamma_E^3 + 10 \zeta_2 \gamma_E^2 \right) + A_1 B_2 \left(4 L_{qr} \gamma_E^2 - 4 L_{qr}^2 \gamma_E \right. \\
 & \left. - 4 L_{fr} \gamma_E^2 + 8 L_{fr} L_{qr} \gamma_E - 4 L_{fr}^2 \gamma_E + 10 \zeta_2 L_{qr} - 10 \zeta_2 L_{fr} \right) \\
 & + A_1 B_1 \left(\frac{8}{3} \beta_0 L_{qr} \gamma_E^3 - 6 \beta_0 L_{qr}^2 \gamma_E^2 + 4 \beta_0 L_{qr}^3 \gamma_E - \frac{8}{3} \beta_0 L_{fr} \gamma_E^3 \right. \\
 & + 4 \beta_0 L_{fr} L_{qr} \gamma_E^2 - 4 \beta_0 L_{fr} L_{qr}^2 \gamma_E + 2 \beta_0 L_{fr}^2 \gamma_E^2 - 4 \beta_0 L_{fr}^2 L_{qr} \gamma_E \\
 & + 4 \beta_0 L_{fr}^3 \gamma_E + \frac{16}{3} \beta_0 \zeta_3 L_{qr} - \frac{16}{3} \beta_0 \zeta_3 L_{fr} + 12 \beta_0 \zeta_2 \gamma_E^2 - 4 \beta_0 \zeta_2 L_{qr} \gamma_E \\
 & \left. - 15 \beta_0 \zeta_2 L_{qr}^2 + 4 \beta_0 \zeta_2 L_{fr} \gamma_E + 10 \beta_0 \zeta_2 L_{fr} L_{qr} + 5 \beta_0 \zeta_2 L_{fr}^2 + 30 \beta_0 \zeta_2^2 \right) \\
 & + A_1 B_1 \tilde{G}_{11} \left(8 L_{qr} \gamma_E^2 - 8 L_{qr}^2 \gamma_E - 8 L_{fr} \gamma_E^2 + 16 L_{fr} L_{qr} \gamma_E - 8 L_{fr}^2 \gamma_E \right. \\
 & \left. + 20 \zeta_2 L_{qr} - 20 \zeta_2 L_{fr} \right) + A_1 B_1 G_{11} \left(8 L_{qr} \gamma_E^2 - 8 L_{qr}^2 \gamma_E - 8 L_{fr} \gamma_E^2 \right. \\
 & + 16 L_{fr} L_{qr} \gamma_E - 8 L_{fr}^2 \gamma_E + 20 \zeta_2 L_{qr} - 20 \zeta_2 L_{fr} \Big) + A_1 B_1 f_1 \left(8 L_{qr} \gamma_E^3 \right. \\
 & \left. - 8 L_{qr}^2 \gamma_E^2 - 8 L_{fr} \gamma_E^3 + 16 L_{fr} L_{qr} \gamma_E^2 - 8 L_{fr}^2 \gamma_E^2 + 20 \zeta_2 L_{qr} \gamma_E \right. \\
 & \left. - 20 \zeta_2 L_{fr} \gamma_E \right) + A_1 B_1^2 \left(4 L_{qr}^2 \gamma_E^2 - 4 L_{qr}^3 \gamma_E - 8 L_{fr} L_{qr} \gamma_E^2 + 12 L_{fr} L_{qr}^2 \gamma_E \right. \\
 & \left. + 4 L_{fr}^2 \gamma_E^2 - 12 L_{fr}^2 L_{qr} \gamma_E + 4 L_{fr}^3 \gamma_E + 10 \zeta_2 L_{qr}^2 - 20 \zeta_2 L_{fr} L_{qr} + 10 \zeta_2 L_{fr}^2 \right) \\
 & + A_1 A_2 \left(4 \gamma_E^4 - 8 L_{qr} \gamma_E^3 + 4 L_{qr}^2 \gamma_E^2 + 8 L_{fr} \gamma_E^3 - 8 L_{fr} L_{qr} \gamma_E^2 + 4 L_{fr}^2 \gamma_E^2 \right. \\
 & + 20 \zeta_2 \gamma_E^2 - 20 \zeta_2 L_{qr} \gamma_E + 20 \zeta_2 L_{fr} \gamma_E + 25 \zeta_2^2 \Big) + A_1^2 \left(\frac{8}{3} \beta_0 \gamma_E^5 \right. \\
 & \left. - \frac{20}{3} \beta_0 L_{qr} \gamma_E^4 + 6 \beta_0 L_{qr}^2 \gamma_E^3 - 2 \beta_0 L_{qr}^3 \gamma_E^2 + \frac{8}{3} \beta_0 L_{fr} \gamma_E^4 - 4 \beta_0 L_{fr} L_{qr} \gamma_E^3 \right. \\
 & + 2 \beta_0 L_{fr} L_{qr}^2 \gamma_E^2 - 2 \beta_0 L_{fr}^2 \gamma_E^3 + 2 \beta_0 L_{fr}^2 L_{qr} \gamma_E^2 - 2 \beta_0 L_{fr}^3 \gamma_E^2 + \frac{16}{3} \beta_0 \zeta_3 \gamma_E^2 \\
 & \left. - \frac{16}{3} \beta_0 \zeta_3 L_{qr} \gamma_E + \frac{16}{3} \beta_0 \zeta_3 L_{fr} \gamma_E + \frac{44}{3} \beta_0 \zeta_2 \gamma_E^3 - 28 \beta_0 \zeta_2 L_{qr} \gamma_E^2 \right. \\
 & + 15 \beta_0 \zeta_2 L_{qr}^2 \gamma_E + 8 \beta_0 \zeta_2 L_{fr} \gamma_E^2 - 10 \beta_0 \zeta_2 L_{fr} L_{qr} \gamma_E - 5 \beta_0 \zeta_2 L_{fr}^2 \gamma_E \\
 & \left. + \frac{40}{3} \beta_0 \zeta_2 \zeta_3 + 20 \beta_0 \zeta_2^2 \gamma_E - 25 \beta_0 \zeta_2^2 L_{qr} \right) + A_1^2 \tilde{G}_{11} \left(4 \gamma_E^4 - 8 L_{qr} \gamma_E^3 \right. \\
 & + 4 L_{qr}^2 \gamma_E^2 + 8 L_{fr} \gamma_E^3 - 8 L_{fr} L_{qr} \gamma_E^2 + 4 L_{fr}^2 \gamma_E^2 + 20 \zeta_2 \gamma_E^2 - 20 \zeta_2 L_{qr} \gamma_E \\
 & \left. + 20 \zeta_2 L_{fr} \gamma_E + 25 \zeta_2^2 \right) + A_1^2 G_{11} \left(4 \gamma_E^4 - 8 L_{qr} \gamma_E^3 + 4 L_{qr}^2 \gamma_E^2 + 8 L_{fr} \gamma_E^3 \right. \\
 & \left. - 8 L_{fr} L_{qr} \gamma_E^2 + 4 L_{fr}^2 \gamma_E^2 + 20 \zeta_2 \gamma_E^2 - 20 \zeta_2 L_{qr} \gamma_E + 20 \zeta_2 L_{fr} \gamma_E + 25 \zeta_2^2 \right)
 \end{aligned}$$

$$\begin{aligned}
 & + A_1^2 f_1 \left(4 \gamma_E^5 - 8 L_{qr} \gamma_E^4 + 4 L_{qr}^2 \gamma_E^3 + 8 L_{fr} \gamma_E^4 - 8 L_{fr} L_{qr} \gamma_E^3 + 4 L_{fr}^2 \gamma_E^3 \right. \\
 & + 20 \zeta_2 \gamma_E^3 - 20 \zeta_2 L_{qr} \gamma_E^2 + 20 \zeta_2 L_{fr} \gamma_E^2 + 25 \zeta_2^2 \gamma_E \left. \right) + A_1^2 B_1 \left(4 L_{qr} \gamma_E^4 \right. \\
 & - 8 L_{qr}^2 \gamma_E^3 + 4 L_{qr}^3 \gamma_E^2 - 4 L_{fr} \gamma_E^4 + 16 L_{fr} L_{qr} \gamma_E^3 - 12 L_{fr} L_{qr}^2 \gamma_E^2 - 8 L_{fr}^2 \gamma_E^3 \\
 & + 12 L_{fr}^2 L_{qr} \gamma_E^2 - 4 L_{fr}^3 \gamma_E^2 + 20 \zeta_2 L_{qr} \gamma_E^2 - 20 \zeta_2 L_{qr}^2 \gamma_E - 20 \zeta_2 L_{fr} \gamma_E^2 \\
 & + 40 \zeta_2 L_{fr} L_{qr} \gamma_E - 20 \zeta_2 L_{fr}^2 \gamma_E + 25 \zeta_2^2 L_{qr} - 25 \zeta_2^2 L_{fr} \left. \right) + A_1^3 \left(\frac{4}{3} \gamma_E^6 \right. \\
 & - 4 L_{qr} \gamma_E^5 + 4 L_{qr}^2 \gamma_E^4 - \frac{4}{3} L_{qr}^3 \gamma_E^3 + 4 L_{fr} \gamma_E^5 - 8 L_{fr} L_{qr} \gamma_E^4 + 4 L_{fr} L_{qr}^2 \gamma_E^3 \\
 & + 4 L_{fr}^2 \gamma_E^4 - 4 L_{fr}^2 L_{qr} \gamma_E^3 + \frac{4}{3} L_{fr}^3 \gamma_E^3 + 10 \zeta_2 \gamma_E^4 - 20 \zeta_2 L_{qr} \gamma_E^3 + 10 \zeta_2 L_{qr}^2 \gamma_E^2 \\
 & + 20 \zeta_2 L_{fr} \gamma_E^3 - 20 \zeta_2 L_{fr} L_{qr} \gamma_E^2 + 10 \zeta_2 L_{fr}^2 \gamma_E^2 + 25 \zeta_2^2 \gamma_E^2 - 25 \zeta_2^2 L_{qr} \gamma_E \\
 & \left. + 25 \zeta_2^2 L_{fr} \gamma_E + \frac{125}{6} \zeta_2^3 \right) \Big]. \tag{B.14}
 \end{aligned}$$

B.3 Resummation ingredients for the Soft exponentiation

In the case for **Soft exponentiation**, all the terms coming from the soft function are exponentiated and hence this means all the contribution to the finite (N-independent) piece from the soft function is also being exponentiated. This renders the g_0 coefficients of the **Standard** \bar{N} threshold and changes also the resummed exponent. We write these changes below in terms of the **Standard** \bar{N} threshold exponent and pre-factor,

$$\begin{aligned}
 g_1^{\text{Soft}} &= \bar{g}_1, \\
 g_2^{\text{Soft}} &= \bar{g}_2 + a_s \Delta_{g_2}^{\text{Soft}}, \\
 g_3^{\text{Soft}} &= \bar{g}_3 + a_s^2 \Delta_{g_3}^{\text{Soft}}, \\
 g_4^{\text{Soft}} &= \bar{g}_4 + a_s^3 \Delta_{g_4}^{\text{Soft}}, \tag{B.15}
 \end{aligned}$$

where the coefficients $\Delta_{g_i}^{\text{Soft}}$ are given as,

$$\Delta_{g_1}^{\text{Soft}} = \left[\tilde{G}_{11} \left(2 \right) + f_1 \left(-L_{qr} \right) + A_1 \left(\frac{1}{2} L_{qr}^2 + 2 \zeta_2 \right) \right], \tag{B.16}$$

$$\begin{aligned}
 \Delta_{g_2}^{\text{Soft}} &= \left[\tilde{G}_{21} \left(1 \right) + \tilde{G}_{12} \left(2 \beta_0 \right) + \tilde{G}_{11} \left(-2 \beta_0 L_{qr} \right) + f_2 \left(-L_{qr} \right) + f_1 \left(\frac{1}{2} \beta_0 L_{qr}^2 \right. \right. \\
 & \left. \left. + 2 \beta_0 \zeta_2 \right) + A_2 \left(\frac{1}{2} L_{qr}^2 + 2 \zeta_2 \right) + A_1 \left(-\frac{1}{6} \beta_0 L_{qr}^3 + \frac{8}{3} \beta_0 \zeta_3 - 2 \beta_0 \zeta_2 L_{qr} \right) \right], \tag{B.17}
 \end{aligned}$$

$$\begin{aligned}
 \Delta_{g_3}^{\text{Soft}} &= \left[\tilde{G}_{31} \left(\frac{2}{3} \right) + \tilde{G}_{22} \left(\frac{4}{3} \beta_0 \right) + \tilde{G}_{21} \left(-2 \beta_0 L_{qr} \right) + \tilde{G}_{13} \left(\frac{8}{3} \beta_0^2 \right) + \tilde{G}_{12} \left(\frac{4}{3} \beta_1 \right. \right. \\
 & \left. \left. - 4 \beta_0^2 L_{qr} \right) + \tilde{G}_{11} \left(-2 \beta_1 L_{qr} + 2 \beta_0^2 L_{qr}^2 + 8 \beta_0^2 \zeta_2 \right) + f_3 \left(-L_{qr} \right) \right. \\
 & \left. + f_2 \left(\beta_0 L_{qr}^2 + 4 \beta_0 \zeta_2 \right) + f_1 \left(\frac{1}{2} \beta_1 L_{qr}^2 + 2 \beta_1 \zeta_2 - \frac{1}{3} \beta_0^2 L_{qr}^3 + \frac{16}{3} \beta_0^2 \zeta_3 \right. \right. \\
 & \left. \left. - 4 \beta_0^2 \zeta_2 L_{qr} \right) + A_3 \left(\frac{1}{2} L_{qr}^2 + 2 \zeta_2 \right) + A_2 \left(-\frac{1}{3} \beta_0 L_{qr}^3 + \frac{16}{3} \beta_0 \zeta_3 \right) \right]
 \end{aligned}$$

$$\begin{aligned}
& -4 \beta_0 \zeta_2 L_{qr} \Big) + A_1 \left(-\frac{1}{6} \beta_1 L_{qr}^3 + \frac{8}{3} \beta_1 \zeta_3 - 2 \beta_1 \zeta_2 L_{qr} + \frac{1}{12} \beta_0^2 L_{qr}^4 \right. \\
& \left. - \frac{16}{3} \beta_0^2 \zeta_3 L_{qr} + 2 \beta_0^2 \zeta_2 L_{qr}^2 + \frac{36}{5} \beta_0^2 \zeta_2^2 \right) \Big]. \tag{B.18}
\end{aligned}$$

The N -independent constants in the case can be put in the following form:

$$\begin{aligned}
g_{01}^{\text{Soft}} &= \bar{g}_{01} + a_s \Delta_{g_{01}}^{\text{Soft}}, \\
g_{02}^{\text{Soft}} &= \bar{g}_{02} + a_s^2 \Delta_{g_{02}}^{\text{Soft}}, \\
g_{03}^{\text{Soft}} &= \bar{g}_{03} + a_s^3 \Delta_{g_{03}}^{\text{Soft}}, \tag{B.19}
\end{aligned}$$

where the coefficients $\Delta_{g_{0i}}^{\text{Soft}}$ are given by,

$$\Delta_{g_{01}}^{\text{Soft}} = \left[\tilde{G}_{11} \left(-2 \right) + f_1 \left(L_{qr} \right) + A_1 \left(-\frac{1}{2} L_{qr}^2 - 2 \zeta_2 \right) \right], \tag{B.20}$$

$$\begin{aligned}
\Delta_{g_{02}}^{\text{Soft}} &= \left[\tilde{G}_{21} \left(-1 \right) + \tilde{G}_{12} \left(-2 \beta_0 \right) + \tilde{G}_{11} \left(2 \beta_0 L_{qr} \right) + \tilde{G}_{11}^2 \left(-2 \right) + G_{11} \tilde{G}_{11} \left(-4 \right) \right. \\
& + f_2 \left(L_{qr} \right) + f_1 \left(-\frac{1}{2} \beta_0 L_{qr}^2 - 2 \beta_0 \zeta_2 \right) + f_1 G_{11} \left(2 L_{qr} \right) + f_1^2 \left(\frac{1}{2} L_{qr}^2 \right) \\
& + B_1 \tilde{G}_{11} \left(-4 L_{qr} + 4 L_{fr} \right) + B_1 f_1 \left(2 L_{qr}^2 - 2 L_{fr} L_{qr} \right) + A_2 \left(-\frac{1}{2} L_{qr}^2 - 2 \zeta_2 \right) \\
& + A_1 \left(\frac{1}{6} \beta_0 L_{qr}^3 - \frac{8}{3} \beta_0 \zeta_3 + 2 \beta_0 \zeta_2 L_{qr} \right) + A_1 \tilde{G}_{11} \left(-10 \zeta_2 \right) + A_1 G_{11} \left(-L_{qr}^2 \right. \\
& \left. - 4 \zeta_2 \right) + A_1 f_1 \left(-\frac{1}{2} L_{qr}^3 + 3 \zeta_2 L_{qr} \right) + A_1 B_1 \left(-L_{qr}^3 + L_{fr} L_{qr}^2 - 4 \zeta_2 L_{qr} \right. \\
& \left. + 4 \zeta_2 L_{fr} \right) + A_1^2 \left(\frac{1}{8} L_{qr}^4 - \frac{3}{2} \zeta_2 L_{qr}^2 - 8 \zeta_2^2 \right) \Big], \tag{B.21}
\end{aligned}$$

$$\begin{aligned}
\Delta_{g_{03}}^{\text{Soft}} &= \left[\tilde{G}_{31} \left(-\frac{2}{3} \right) + \tilde{G}_{22} \left(-\frac{4}{3} \beta_0 \right) + \tilde{G}_{21} \left(2 \beta_0 L_{qr} \right) + \tilde{G}_{13} \left(-\frac{8}{3} \beta_0^2 \right) + \tilde{G}_{12} \left(-\frac{4}{3} \beta_1 \right. \right. \\
& \left. + 4 \beta_0^2 L_{qr} \right) + \tilde{G}_{11} \left(2 \beta_1 L_{qr} - 2 \beta_0^2 L_{qr}^2 - 8 \beta_0^2 \zeta_2 \right) + \tilde{G}_{11} \tilde{G}_{21} \left(-2 \right) + \tilde{G}_{11} \tilde{G}_{12} \left(\right. \\
& \left. - 4 \beta_0 \right) + \tilde{G}_{11}^2 \left(4 \beta_0 L_{qr} \right) + \tilde{G}_{11}^3 \left(-\frac{4}{3} \right) + G_{21} \tilde{G}_{11} \left(-2 \right) + G_{12} \tilde{G}_{11} \left(-4 \beta_0 \right) \\
& + G_{11} \tilde{G}_{21} \left(-2 \right) + G_{11} \tilde{G}_{12} \left(-4 \beta_0 \right) + G_{11} \tilde{G}_{11} \left(8 \beta_0 L_{qr} \right) + G_{11} \tilde{G}_{11}^2 \left(-4 \right) \\
& + G_{11}^2 \tilde{G}_{11} \left(-4 \right) + f_3 \left(L_{qr} \right) + f_2 \left(-\beta_0 L_{qr}^2 - 4 \beta_0 \zeta_2 \right) + f_2 G_{11} \left(2 L_{qr} \right) + f_1 \left(\right. \\
& \left. -\frac{1}{2} \beta_1 L_{qr}^2 - 2 \beta_1 \zeta_2 + \frac{1}{3} \beta_0^2 L_{qr}^3 - \frac{16}{3} \beta_0^2 \zeta_3 + 4 \beta_0^2 \zeta_2 L_{qr} \right) + f_1 \tilde{G}_{11} \left(-10 \beta_0 \zeta_2 \right) \\
& + f_1 G_{21} \left(L_{qr} \right) + f_1 G_{12} \left(2 \beta_0 L_{qr} \right) + f_1 G_{11} \left(-3 \beta_0 L_{qr}^2 - 4 \beta_0 \zeta_2 \right) \\
& + f_1 G_{11}^2 \left(2 L_{qr} \right) + f_1 f_2 \left(L_{qr}^2 \right) + f_1^2 \left(-\frac{1}{2} \beta_0 L_{qr}^3 + 3 \beta_0 \zeta_2 L_{qr} \right) \\
& \left. + f_1^2 G_{11} \left(L_{qr}^2 \right) + f_1^3 \left(\frac{1}{6} L_{qr}^3 \right) + B_2 \tilde{G}_{11} \left(-4 L_{qr} + 4 L_{fr} \right) + B_2 f_1 \left(2 L_{qr}^2 \right. \right.
\end{aligned}$$

$$\begin{aligned}
& -2 L_{fr} L_{qr}) + B_1 \tilde{G}_{21} \left(-2 L_{qr} + 2 L_{fr} \right) + B_1 \tilde{G}_{12} \left(-4 \beta_0 L_{qr} + 4 \beta_0 L_{fr} \right) \\
& + B_1 \tilde{G}_{11} \left(6 \beta_0 L_{qr}^2 - 4 \beta_0 L_{fr} L_{qr} - 2 \beta_0 L_{fr}^2 - 12 \beta_0 \zeta_2 \right) + B_1 \tilde{G}_{11}^2 \left(-4 L_{qr} \right. \\
& \left. + 4 L_{fr} \right) + B_1 G_{11} \tilde{G}_{11} \left(-8 L_{qr} + 8 L_{fr} \right) + B_1 f_2 \left(2 L_{qr}^2 - 2 L_{fr} L_{qr} \right) \\
& + B_1 f_1 \left(-2 \beta_0 L_{qr}^3 + \beta_0 L_{fr} L_{qr}^2 + \beta_0 L_{fr}^2 L_{qr} + 2 \beta_0 \zeta_2 L_{qr} + 4 \beta_0 \zeta_2 L_{fr} \right) \\
& + B_1 f_1 G_{11} \left(4 L_{qr}^2 - 4 L_{fr} L_{qr} \right) + B_1 f_1^2 \left(L_{qr}^3 - L_{fr} L_{qr}^2 \right) + B_1^2 \tilde{G}_{11} \left(-4 L_{qr}^2 \right. \\
& \left. + 8 L_{fr} L_{qr} - 4 L_{fr}^2 \right) + B_1^2 f_1 \left(2 L_{qr}^3 - 4 L_{fr} L_{qr}^2 + 2 L_{fr}^2 L_{qr} \right) + A_3 \left(-\frac{1}{2} L_{qr}^2 \right. \\
& \left. - 2 \zeta_2 \right) + A_2 \left(\frac{1}{3} \beta_0 L_{qr}^3 - \frac{16}{3} \beta_0 \zeta_3 + 4 \beta_0 \zeta_2 L_{qr} \right) + A_2 \tilde{G}_{11} \left(-10 \zeta_2 \right) \\
& + A_2 G_{11} \left(-L_{qr}^2 - 4 \zeta_2 \right) + A_2 f_1 \left(-\frac{1}{2} L_{qr}^3 + 3 \zeta_2 L_{qr} \right) + A_2 B_1 \left(-L_{qr}^3 \right. \\
& \left. + L_{fr} L_{qr}^2 - 4 \zeta_2 L_{qr} + 4 \zeta_2 L_{fr} \right) + A_1 \left(\frac{1}{6} \beta_1 L_{qr}^3 - \frac{8}{3} \beta_1 \zeta_3 + 2 \beta_1 \zeta_2 L_{qr} \right. \\
& \left. - \frac{1}{12} \beta_0^2 L_{qr}^4 + \frac{16}{3} \beta_0^2 \zeta_3 L_{qr} - 2 \beta_0^2 \zeta_2 L_{qr}^2 - \frac{36}{5} \beta_0^2 \zeta_2^2 \right) + A_1 \tilde{G}_{21} \left(-5 \zeta_2 \right) \\
& + A_1 \tilde{G}_{12} \left(-10 \beta_0 \zeta_2 \right) + A_1 \tilde{G}_{11} \left(-\frac{16}{3} \beta_0 \zeta_3 + 20 \beta_0 \zeta_2 L_{qr} \right) + A_1 \tilde{G}_{11}^2 \left(\right. \\
& \left. -10 \zeta_2 \right) + A_1 G_{21} \left(-\frac{1}{2} L_{qr}^2 - 2 \zeta_2 \right) + A_1 G_{12} \left(-\beta_0 L_{qr}^2 - 4 \beta_0 \zeta_2 \right) \\
& + A_1 G_{11} \left(\frac{4}{3} \beta_0 L_{qr}^3 - \frac{16}{3} \beta_0 \zeta_3 + 8 \beta_0 \zeta_2 L_{qr} \right) + A_1 G_{11} \tilde{G}_{11} \left(-20 \zeta_2 \right) \\
& + A_1 G_{11}^2 \left(-L_{qr}^2 - 4 \zeta_2 \right) + A_1 f_2 \left(-\frac{1}{2} L_{qr}^3 + 3 \zeta_2 L_{qr} \right) + A_1 f_1 \left(\frac{5}{12} \beta_0 L_{qr}^4 \right. \\
& \left. - 6 \beta_0 \zeta_2 L_{qr}^2 - 16 \beta_0 \zeta_2^2 \right) + A_1 f_1 G_{11} \left(-L_{qr}^3 + 6 \zeta_2 L_{qr} \right) + A_1 f_1^2 \left(-\frac{1}{4} L_{qr}^4 \right. \\
& \left. + \frac{3}{2} \zeta_2 L_{qr}^2 \right) + A_1 B_2 \left(-L_{qr}^3 + L_{fr} L_{qr}^2 - 4 \zeta_2 L_{qr} + 4 \zeta_2 L_{fr} \right) + A_1 B_1 \left(\frac{5}{6} \beta_0 L_{qr}^4 \right. \\
& \left. - \frac{1}{3} \beta_0 L_{fr} L_{qr}^3 - \frac{1}{2} \beta_0 L_{fr}^2 L_{qr}^2 - \frac{16}{3} \beta_0 \zeta_3 L_{qr} + \frac{16}{3} \beta_0 \zeta_3 L_{fr} + 3 \beta_0 \zeta_2 L_{qr}^2 \right. \\
& \left. - 4 \beta_0 \zeta_2 L_{fr} L_{qr} - 2 \beta_0 \zeta_2 L_{fr}^2 - 12 \beta_0 \zeta_2^2 \right) + A_1 B_1 \tilde{G}_{11} \left(-20 \zeta_2 L_{qr} \right. \\
& \left. + 20 \zeta_2 L_{fr} \right) + A_1 B_1 G_{11} \left(-2 L_{qr}^3 + 2 L_{fr} L_{qr}^2 - 8 \zeta_2 L_{qr} + 8 \zeta_2 L_{fr} \right) \\
& + A_1 B_1 f_1 \left(-L_{qr}^4 + L_{fr} L_{qr}^3 + 6 \zeta_2 L_{qr}^2 - 6 \zeta_2 L_{fr} L_{qr} \right) + A_1 B_1^2 \left(-L_{qr}^4 \right. \\
& \left. + 2 L_{fr} L_{qr}^3 - L_{fr}^2 L_{qr}^2 - 4 \zeta_2 L_{qr}^2 + 8 \zeta_2 L_{fr} L_{qr} - 4 \zeta_2 L_{fr}^2 \right) + A_1 A_2 \left(\frac{1}{4} L_{qr}^4 \right.
\end{aligned}$$

$$\begin{aligned}
 & -3 \zeta_2 L_{qr}^2 - 16 \zeta_2^2) + A_1^2 \left(-\frac{1}{12} \beta_0 L_{qr}^5 + 2 \beta_0 \zeta_2 L_{qr}^3 - \frac{40}{3} \beta_0 \zeta_2 \zeta_3 \right. \\
 & \left. + 16 \beta_0 \zeta_2^2 L_{qr} \right) + A_1^2 \tilde{G}_{11} \left(-25 \zeta_2^2 \right) + A_1^2 G_{11} \left(\frac{1}{4} L_{qr}^4 - 3 \zeta_2 L_{qr}^2 - 16 \zeta_2^2 \right) \\
 & + A_1^2 f_1 \left(\frac{1}{8} L_{qr}^5 - \frac{3}{2} \zeta_2 L_{qr}^3 + \frac{9}{2} \zeta_2^2 L_{qr} \right) + A_1^2 B_1 \left(\frac{1}{4} L_{qr}^5 - \frac{1}{4} L_{fr} L_{qr}^4 - 3 \zeta_2 L_{qr}^3 \right. \\
 & \left. + 3 \zeta_2 L_{fr} L_{qr}^2 - 16 \zeta_2^2 L_{qr} + 16 \zeta_2^2 L_{fr} \right) + A_1^3 \left(-\frac{1}{48} L_{qr}^6 + \frac{3}{8} \zeta_2 L_{qr}^4 - \frac{9}{4} \zeta_2^2 L_{qr}^2 \right. \\
 & \left. - \frac{49}{3} \zeta_2^3 \right) \Big]. \tag{B.22}
 \end{aligned}$$

B.4 Resummation ingredients for the All exponentiation

In the case for All exponentiation, the complete g_0 is being exponentiated along with the large- N pieces. This brings into modification only for the resummed exponent compared to the Standard \bar{N} exponentiation. We write the resummed exponent in this case in terms of \bar{N} exponents as,

$$\begin{aligned}
 g_1^{\text{All}} &= \bar{g}_1, \\
 g_2^{\text{All}} &= \bar{g}_2 + a_s \Delta_{g_2}^{\text{All}}, \\
 g_3^{\text{All}} &= \bar{g}_3 + a_s^2 \Delta_{g_3}^{\text{All}}, \\
 g_4^{\text{All}} &= \bar{g}_4 + a_s^3 \Delta_{g_4}^{\text{All}},
 \end{aligned} \tag{B.23}$$

where $\Delta_{g_i}^{\text{All}}$ terms are found from exponentiating also the complete g_0 prefactor and they are given as,

$$\begin{aligned}
 \Delta_{g_2}^{\text{All}} &= \bar{g}_{01}, \\
 \Delta_{g_2}^{\text{All}} &= \left(-\frac{\bar{g}_{01}^2}{2} + \bar{g}_{02} \right), \\
 \Delta_{g_2}^{\text{All}} &= \left(\frac{\bar{g}_{01}^3}{3} - \bar{g}_{01} \bar{g}_{02} + \bar{g}_{03} \right),
 \end{aligned} \tag{B.24}$$

where the coefficients \bar{g}_{0i} are given in (B.3).

C Anomalous dimensions

Here we present all the anomalous dimensions used in performing the resummation.

The cusp anomalous dimensions are given as

$$A_1 = \left\{ C_F \left(4 \right) \right\}, \tag{C.1}$$

$$A_2 = \left\{ C_F n_f \left(-\frac{40}{9} \right) + C_F C_A \left(\frac{268}{9} - 8 \zeta_2 \right) \right\}, \tag{C.2}$$

$$\begin{aligned}
 A_3 = \left\{ C_F n_f^2 \left(-\frac{16}{27} \right) + C_F C_A n_f \left(-\frac{836}{27} - \frac{112}{3} \zeta_3 + \frac{160}{9} \zeta_2 \right) + C_F C_A^2 \left(\frac{490}{3} + \frac{88}{3} \zeta_3 \right. \right. \\
 \left. \left. - \frac{1072}{9} \zeta_2 + \frac{176}{5} \zeta_2^2 \right) + C_F^2 n_f \left(-\frac{110}{3} + 32 \zeta_3 \right) \right\}, \tag{C.3}
 \end{aligned}$$

$$\begin{aligned}
 A_4 = & \left\{ C_F \frac{d_F^{abcd} d_A^{abcd}}{N_A} \left(\frac{7040}{3} \zeta_5 + \frac{256}{3} \zeta_3 - 768 \zeta_3^2 - 256 \zeta_2 - \frac{15872}{35} \zeta_2^3 \right) + C_F n_f \frac{d_F^{abcd} d_A^{abcd}}{N_A} \left(\right. \right. \\
 & - \frac{2560}{3} \zeta_5 - \frac{512}{3} \zeta_3 + 512 \zeta_2 \left. \right) + C_F n_f^3 \left(-\frac{32}{81} + \frac{64}{27} \zeta_3 \right) + C_F^2 n_f^2 \left(\frac{2392}{81} - \frac{640}{9} \zeta_3 + \frac{64}{5} \zeta_2^2 \right) \\
 & + C_F^3 n_f \left(\frac{572}{9} - 320 \zeta_5 + \frac{592}{3} \zeta_3 \right) + C_A C_F n_f^2 \left(\frac{923}{81} + \frac{2240}{27} \zeta_3 - \frac{608}{81} \zeta_2 - \frac{224}{15} \zeta_2^2 \right) \\
 & + C_A C_F^2 n_f \left(-\frac{34066}{81} + 160 \zeta_5 + \frac{3712}{9} \zeta_3 + \frac{440}{3} \zeta_2 - 128 \zeta_2 \zeta_3 - \frac{352}{5} \zeta_2^2 \right) + C_A^2 C_F n_f \left(\right. \\
 & - \frac{24137}{81} + \frac{2096}{9} \zeta_5 - \frac{23104}{27} \zeta_3 + \frac{20320}{81} \zeta_2 + \frac{448}{3} \zeta_2 \zeta_3 - \frac{352}{15} \zeta_2^2 \left. \right) + C_A^3 C_F \left(\frac{84278}{81} \right. \\
 & \left. - \frac{3608}{9} \zeta_5 + \frac{20944}{27} \zeta_3 - 16 \zeta_3^2 - \frac{88400}{81} \zeta_2 - \frac{352}{3} \zeta_2 \zeta_3 + \frac{3608}{5} \zeta_2^2 - \frac{20032}{105} \zeta_2^3 \right) \left. \right\}. \quad (C.4)
 \end{aligned}$$

The quartic casimirs are given by

$$\frac{d_F^{abcd} d_A^{abcd}}{N_A} = \frac{n_c(n_c^2 + 6)}{48}, \quad \frac{d_F^{abcd} d_F^{abcd}}{N_A} = \frac{(n_c^4 - 6n_c^2 + 18)}{96n_c^2}, \quad (C.5)$$

with $N_A = n_c^2 - 1$ and $N_F = n_c$ where $n_c = 3$ for QCD.

The universal D coefficients are given as,

$$D_1 = C_F \left\{ 0 \right\}, \quad (C.6)$$

$$D_2 = C_F \left\{ n_f \left(\frac{224}{27} - \frac{32}{3} \zeta_2 \right) + C_A \left(-\frac{1616}{27} + 56 \zeta_3 + \frac{176}{3} \zeta_2 \right) \right\}, \quad (C.7)$$

$$\begin{aligned}
 D_3 = & C_F \left\{ n_f^2 \left(-\frac{3712}{729} + \frac{320}{27} \zeta_3 + \frac{640}{27} \zeta_2 \right) + C_F n_f \left(\frac{3422}{27} - \frac{608}{9} \zeta_3 - 32 \zeta_2 - \frac{64}{5} \zeta_2^2 \right) \right. \\
 & + C_A n_f \left(\frac{125252}{729} - \frac{2480}{9} \zeta_3 - \frac{29392}{81} \zeta_2 + \frac{736}{15} \zeta_2^2 \right) + C_A^2 \left(-\frac{594058}{729} - 384 \zeta_5 \right. \\
 & \left. + \frac{40144}{27} \zeta_3 + \frac{98224}{81} \zeta_2 - \frac{352}{3} \zeta_2 \zeta_3 - \frac{2992}{15} \zeta_2^2 \right) \left. \right\}. \quad (C.8)
 \end{aligned}$$

The coefficients B are given as

$$B_1 = \left\{ C_F \left(3 \right) \right\}, \quad (C.9)$$

$$B_2 = \left\{ C_F n_f \left(-\frac{1}{3} - \frac{8}{3} \zeta_2 \right) + C_F C_A \left(\frac{17}{6} - 12 \zeta_3 + \frac{44}{3} \zeta_2 \right) + C_F^2 \left(\frac{3}{2} + 24 \zeta_3 - 12 \zeta_2 \right) \right\}, \quad (C.10)$$

$$\begin{aligned}
 B_3 = & \left\{ C_F n_f^2 \left(-\frac{17}{9} - \frac{16}{9} \zeta_3 + \frac{80}{27} \zeta_2 \right) + C_F C_A n_f \left(20 + \frac{200}{9} \zeta_3 - \frac{1336}{27} \zeta_2 + \frac{4}{5} \zeta_2^2 \right) \right. \\
 & + C_F C_A^2 \left(-\frac{1657}{36} + 40 \zeta_5 - \frac{1552}{9} \zeta_3 + \frac{4496}{27} \zeta_2 - 2 \zeta_2^2 \right) + C_F^2 n_f \left(-23 - \frac{136}{3} \zeta_3 \right. \\
 & + \frac{20}{3} \zeta_2 + \frac{232}{15} \zeta_2^2 \left. \right) + C_F^2 C_A \left(\frac{151}{4} + 120 \zeta_5 + \frac{844}{3} \zeta_3 - \frac{410}{3} \zeta_2 + 16 \zeta_2 \zeta_3 - \frac{988}{15} \zeta_2^2 \right) \\
 & \left. + C_F^3 \left(\frac{29}{2} - 240 \zeta_5 + 68 \zeta_3 + 18 \zeta_2 - 32 \zeta_2 \zeta_3 + \frac{288}{5} \zeta_2^2 \right) \right\}. \quad (C.11)
 \end{aligned}$$

The anomalous dimensions f are given as

$$f_1 = \left\{ 0 \right\}, \tag{C.12}$$

$$f_2 = \left\{ C_F n_f \left(-\frac{112}{27} + \frac{4}{3} \zeta_2 \right) + C_F C_A \left(\frac{808}{27} - 28 \zeta_3 - \frac{22}{3} \zeta_2 \right) \right\}, \tag{C.13}$$

$$f_3 = \left\{ C_F n_f^2 \left(-\frac{2080}{729} + \frac{112}{27} \zeta_3 - \frac{40}{27} \zeta_2 \right) + C_F C_A n_f \left(-\frac{11842}{729} + \frac{728}{27} \zeta_3 + \frac{2828}{81} \zeta_2 - \frac{96}{5} \zeta_2^2 \right) + C_F C_A^2 \left(\frac{136781}{729} + 192 \zeta_5 - \frac{1316}{3} \zeta_3 - \frac{12650}{81} \zeta_2 + \frac{176}{3} \zeta_2 \zeta_3 + \frac{352}{5} \zeta_2^2 \right) + C_F^2 n_f \left(-\frac{1711}{27} + \frac{304}{9} \zeta_3 + 4 \zeta_2 + \frac{32}{5} \zeta_2^2 \right) \right\}. \tag{C.14}$$

The finite G coefficients coming from the explicit calculation of the form factor are given as

$$G_{11} = \left\{ C_F \left(-8 + \zeta_2 \right) \right\}, \tag{C.15}$$

$$G_{12} = \left\{ C_F \left(8 - \frac{7}{3} \zeta_3 - \frac{3}{4} \zeta_2 \right) \right\}, \tag{C.16}$$

$$G_{13} = \left\{ C_F \left(-8 + \frac{7}{4} \zeta_3 + \zeta_2 + \frac{47}{80} \zeta_2^2 \right) \right\}, \tag{C.17}$$

$$G_{21} = \left\{ C_F n_f \left(\frac{5813}{162} - \frac{8}{3} \zeta_3 + \frac{37}{9} \zeta_2 \right) + C_F C_A \left(-\frac{70165}{324} + \frac{260}{3} \zeta_3 - \frac{575}{18} \zeta_2 + \frac{88}{5} \zeta_2^2 \right) + C_F^2 \left(-\frac{1}{4} - 60 \zeta_3 + 58 \zeta_2 - \frac{88}{5} \zeta_2^2 \right) \right\}, \tag{C.18}$$

$$G_{22} = \left\{ C_F n_f \left(-\frac{129389}{1944} + \frac{301}{27} \zeta_3 - \frac{425}{54} \zeta_2 + \frac{7}{12} \zeta_2^2 \right) + C_F C_A \left(\frac{1547797}{3888} - 51 \zeta_5 - \frac{12479}{54} \zeta_3 + \frac{7297}{108} \zeta_2 + \frac{89}{3} \zeta_2 \zeta_3 - \frac{653}{24} \zeta_2^2 \right) + C_F^2 \left(-\frac{109}{16} + 12 \zeta_5 + 184 \zeta_3 - \frac{437}{4} \zeta_2 - 28 \zeta_2 \zeta_3 + \frac{108}{5} \zeta_2^2 \right) \right\}, \tag{C.19}$$

$$G_{31} = \left\{ C_F N_4 n_{fv} \left(12 - 80 \zeta_5 + 14 \zeta_3 + 30 \zeta_2 - \frac{6}{5} \zeta_2^2 \right) + C_F n_f^2 \left(-\frac{258445}{2187} + \frac{536}{81} \zeta_3 - \frac{3466}{81} \zeta_2 - \frac{40}{9} \zeta_2^2 \right) + C_F C_A n_f \left(\frac{3702974}{2187} - 72 \zeta_5 - \frac{68660}{81} \zeta_3 + \frac{155008}{243} \zeta_2 + \frac{392}{9} \zeta_2 \zeta_3 - \frac{1298}{45} \zeta_2^2 \right) + C_F C_A^2 \left(-\frac{48902713}{8748} + \frac{688}{3} \zeta_5 + \frac{85883}{18} \zeta_3 - \frac{1136}{3} \zeta_3^2 - \frac{1083305}{486} \zeta_2 + \frac{1786}{9} \zeta_2 \zeta_3 + \frac{37271}{90} \zeta_2^2 - \frac{6152}{63} \zeta_2^3 \right) + C_F^2 n_f \left(\frac{73271}{162} - \frac{368}{3} \zeta_5 + \frac{19700}{27} \zeta_3 - \frac{7541}{18} \zeta_2 - \frac{152}{3} \zeta_2 \zeta_3 - \frac{704}{45} \zeta_2^2 \right) + C_F^2 C_A \left(\frac{230}{3} - \frac{3020}{3} \zeta_5 - \frac{23402}{9} \zeta_3 + 296 \zeta_3^2 + \frac{55499}{18} \zeta_2 - \frac{3448}{3} \zeta_2 \zeta_3 + \frac{2432}{45} \zeta_2^2 - \frac{15448}{105} \zeta_2^3 \right) + C_F^3 \left(-\frac{1527}{4} + 1992 \zeta_5 - 2130 \zeta_3 + 48 \zeta_3^2 - 206 \zeta_2 + 840 \zeta_2 \zeta_3 - 534 \zeta_2^2 + \frac{21584}{105} \zeta_2^3 \right) \right\}. \tag{C.20}$$

Here $N_4 = (n_c^2 - 4)/n_c$ and n_{fv} is proportional to the charge weighted sum of the quark flavors [15]. The finite \tilde{G} coefficients are found to be

$$\tilde{G}_{11} = \left\{ C_F \left(-3 \zeta_2 \right) \right\}, \tag{C.21}$$

$$\tilde{G}_{12} = \left\{ C_F \left(\frac{7}{3} \zeta_3 \right) \right\}, \tag{C.22}$$

$$\tilde{G}_{13} = \left\{ C_F \left(-\frac{3}{16} \zeta_2^2 \right) \right\}, \tag{C.23}$$

$$\tilde{G}_{21} = \left\{ C_F n_f \left(-\frac{328}{81} + \frac{32}{3} \zeta_3 + \frac{70}{9} \zeta_2 \right) + C_F C_A \left(\frac{2428}{81} - \frac{176}{3} \zeta_3 - \frac{469}{9} \zeta_2 + 4 \zeta_2^2 \right) \right\}, \tag{C.24}$$

$$\begin{aligned} \tilde{G}_{22} = & \left\{ C_F n_f \left(\frac{976}{243} - \frac{310}{27} \zeta_3 - \frac{196}{27} \zeta_2 - \frac{1}{20} \zeta_2^2 \right) + C_F C_A \left(-\frac{7288}{243} + 43 \zeta_5 + \frac{2077}{27} \zeta_3 \right. \right. \\ & \left. \left. + \frac{1414}{27} \zeta_2 - \frac{203}{3} \zeta_2 \zeta_3 + \frac{11}{40} \zeta_2^2 \right) \right\}, \end{aligned} \tag{C.25}$$

$$\begin{aligned} \tilde{G}_{31} = & \left\{ C_F n_f^2 \left(\frac{11584}{2187} - \frac{2720}{81} \zeta_3 - \frac{1996}{81} \zeta_2 + \frac{32}{9} \zeta_2^2 \right) + C_F C_A n_f \left(-\frac{716509}{4374} \right. \right. \\ & \left. \left. + \frac{148}{3} \zeta_5 + \frac{45956}{81} \zeta_3 + \frac{105059}{243} \zeta_2 - \frac{1208}{9} \zeta_2 \zeta_3 - \frac{532}{9} \zeta_2^2 \right) + C_F C_A^2 \left(\frac{7135981}{8748} \right. \right. \\ & \left. \left. - \frac{1430}{3} \zeta_5 - \frac{59648}{27} \zeta_3 + \frac{536}{3} \zeta_3^2 - \frac{765127}{486} \zeta_2 + \frac{11000}{9} \zeta_2 \zeta_3 + \frac{1964}{9} \zeta_2^2 + \frac{152}{63} \zeta_2^3 \right) \right. \\ & \left. + C_F^2 n_f \left(-\frac{42727}{324} + \frac{112}{3} \zeta_5 + \frac{2536}{27} \zeta_3 + \frac{605}{6} \zeta_2 - 88 \zeta_2 \zeta_3 + \frac{152}{15} \zeta_2^2 \right) \right\}. \end{aligned} \tag{C.26}$$

Open Access. This article is distributed under the terms of the Creative Commons Attribution License ([CC-BY 4.0](https://creativecommons.org/licenses/by/4.0/)), which permits any use, distribution and reproduction in any medium, provided the original author(s) and source are credited.

References

- [1] R. Hamberg, W.L. van Neerven and T. Matsuura, *A complete calculation of the order $\alpha - s^2$ correction to the Drell-Yan K factor*, *Nucl. Phys. B* **359** (1991) 343 [Erratum *ibid.* **644** (2002) 403] [[INSPIRE](#)].
- [2] T. Matsuura, R. Hamberg and W.L. van Neerven, *The Contribution of the Gluon-gluon Subprocess to the Drell-Yan K Factor*, *Nucl. Phys. B* **345** (1990) 331 [[INSPIRE](#)].
- [3] R.V. Harlander and W.B. Kilgore, *Next-to-next-to-leading order Higgs production at hadron colliders*, *Phys. Rev. Lett.* **88** (2002) 201801 [[hep-ph/0201206](#)] [[INSPIRE](#)].
- [4] C. Duhr, F. Dulat and B. Mistlberger, *The Drell-Yan cross section to third order in the strong coupling constant*, [arXiv:2001.07717](#) [[INSPIRE](#)].
- [5] C. Anastasiou et al., *Higgs boson gluon-fusion production at threshold in N^3LO QCD*, *Phys. Lett. B* **737** (2014) 325 [[arXiv:1403.4616](#)] [[INSPIRE](#)].

- [6] S. Moch and A. Vogt, *Higher-order soft corrections to lepton pair and Higgs boson production*, *Phys. Lett. B* **631** (2005) 48 [[hep-ph/0508265](#)] [[INSPIRE](#)].
- [7] E. Laenen and L. Magnea, *Threshold resummation for electroweak annihilation from DIS data*, *Phys. Lett. B* **632** (2006) 270 [[hep-ph/0508284](#)] [[INSPIRE](#)].
- [8] V. Ravindran, *On Sudakov and soft resummations in QCD*, *Nucl. Phys. B* **746** (2006) 58 [[hep-ph/0512249](#)] [[INSPIRE](#)].
- [9] V. Ravindran, *Higher-order threshold effects to inclusive processes in QCD*, *Nucl. Phys. B* **752** (2006) 173 [[hep-ph/0603041](#)] [[INSPIRE](#)].
- [10] A. Idilbi, X.-d. Ji, J.-P. Ma and F. Yuan, *Threshold resummation for Higgs production in effective field theory*, *Phys. Rev. D* **73** (2006) 077501 [[hep-ph/0509294](#)] [[INSPIRE](#)].
- [11] Y. Li, A. von Manteuffel, R.M. Schabinger and H.X. Zhu, *Soft-virtual corrections to Higgs production at N^3 LO*, *Phys. Rev. D* **91** (2015) 036008 [[arXiv:1412.2771](#)] [[INSPIRE](#)].
- [12] M.C. Kumar, M.K. Mandal and V. Ravindran, *Associated production of Higgs boson with vector boson at threshold N^3 LO in QCD*, *JHEP* **03** (2015) 037 [[arXiv:1412.3357](#)] [[INSPIRE](#)].
- [13] T. Ahmed, N. Rana and V. Ravindran, *Higgs boson production through $b\bar{b}$ annihilation at threshold in N^3 LO QCD*, *JHEP* **10** (2014) 139 [[arXiv:1408.0787](#)] [[INSPIRE](#)].
- [14] T. Ahmed, M.C. Kumar, P. Mathews, N. Rana and V. Ravindran, *Pseudo-scalar Higgs boson production at threshold N^3 LO and N^3 LL QCD*, *Eur. Phys. J. C* **76** (2016) 355 [[arXiv:1510.02235](#)] [[INSPIRE](#)].
- [15] T. Gehrmann, E.W.N. Glover, T. Huber, N. Iqizlerli and C. Studerus, *Calculation of the quark and gluon form factors to three loops in QCD*, *JHEP* **06** (2010) 094 [[arXiv:1004.3653](#)] [[INSPIRE](#)].
- [16] V. Ravindran, J. Smith and W.L. van Neerven, *QCD threshold corrections to di-lepton and Higgs rapidity distributions beyond N^2 LO*, *Nucl. Phys. B* **767** (2007) 100 [[hep-ph/0608308](#)] [[INSPIRE](#)].
- [17] T. Ahmed, M. Mahakhud, N. Rana and V. Ravindran, *Drell-Yan Production at Threshold to Third Order in QCD*, *Phys. Rev. Lett.* **113** (2014) 112002 [[arXiv:1404.0366](#)] [[INSPIRE](#)].
- [18] Y. Li, A. von Manteuffel, R.M. Schabinger and H.X. Zhu, *N^3 LO Higgs boson and Drell-Yan production at threshold: The one-loop two-emission contribution*, *Phys. Rev. D* **90** (2014) 053006 [[arXiv:1404.5839](#)] [[INSPIRE](#)].
- [19] S. Catani, L. Cieri, D. de Florian, G. Ferrera and M. Grazzini, *Threshold resummation at N^3 LL accuracy and soft-virtual cross sections at N^3 LO*, *Nucl. Phys. B* **888** (2014) 75 [[arXiv:1405.4827](#)] [[INSPIRE](#)].
- [20] G.F. Sterman, *Summation of Large Corrections to Short Distance Hadronic Cross-Sections*, *Nucl. Phys. B* **281** (1987) 310 [[INSPIRE](#)].
- [21] S. Catani and L. Trentadue, *Resummation of the QCD Perturbative Series for Hard Processes*, *Nucl. Phys. B* **327** (1989) 323 [[INSPIRE](#)].
- [22] S. Catani and L. Trentadue, *Comment on QCD exponentiation at large x* , *Nucl. Phys. B* **353** (1991) 183 [[INSPIRE](#)].
- [23] S. Catani, D. de Florian, M. Grazzini and P. Nason, *Soft gluon resummation for Higgs boson production at hadron colliders*, *JHEP* **07** (2003) 028 [[hep-ph/0306211](#)] [[INSPIRE](#)].

- [24] M. Bonvini and S. Marzani, *Resummed Higgs cross section at N^3LL* , *JHEP* **09** (2014) 007 [[arXiv:1405.3654](#)] [[INSPIRE](#)].
- [25] M. Bonvini, S. Marzani, C. Muselli and L. Rottoli, *On the Higgs cross section at N^3LO+N^3LL and its uncertainty*, *JHEP* **08** (2016) 105 [[arXiv:1603.08000](#)] [[INSPIRE](#)].
- [26] T. Ahmed, G. Das, M.C. Kumar, N. Rana and V. Ravindran, *RG improved Higgs boson production to N^3LO in QCD*, [arXiv:1505.07422](#) [[INSPIRE](#)].
- [27] A.H. Ajjath, A. Chakraborty, G. Das, P. Mukherjee and V. Ravindran, *Resummed prediction for Higgs boson production through $b\bar{b}$ annihilation at N^3LL* , *JHEP* **11** (2019) 006 [[arXiv:1905.03771](#)] [[INSPIRE](#)].
- [28] S. Moch, J.A.M. Vermaseren and A. Vogt, *Higher-order corrections in threshold resummation*, *Nucl. Phys. B* **726** (2005) 317 [[hep-ph/0506288](#)] [[INSPIRE](#)].
- [29] G. Das, S.-O. Moch and A. Vogt, *Soft corrections to inclusive deep-inelastic scattering at four loops and beyond*, *JHEP* **03** (2020) 116 [[arXiv:1912.12920](#)] [[INSPIRE](#)].
- [30] A. Idilbi, X.-d. Ji and F. Yuan, *Resummation of threshold logarithms in effective field theory for DIS, Drell-Yan and Higgs production*, *Nucl. Phys. B* **753** (2006) 42 [[hep-ph/0605068](#)] [[INSPIRE](#)].
- [31] T. Ahmed et al., *Pseudo-scalar Higgs boson production at $N^3LO_A+N^3LL'$* , *Eur. Phys. J. C* **76** (2016) 663 [[arXiv:1606.00837](#)] [[INSPIRE](#)].
- [32] T. Schmidt and M. Spira, *Higgs Boson Production via Gluon Fusion: Soft-Gluon Resummation including Mass Effects*, *Phys. Rev. D* **93** (2016) 014022 [[arXiv:1509.00195](#)] [[INSPIRE](#)].
- [33] D. de Florian and J. Zurita, *Soft-gluon resummation for pseudoscalar Higgs boson production at hadron colliders*, *Phys. Lett. B* **659** (2008) 813 [[arXiv:0711.1916](#)] [[INSPIRE](#)].
- [34] G. Das, M.C. Kumar and K. Samanta, *Resummed inclusive cross-section in ADD model at $N^3LL+NNLO$* , [arXiv:1912.13039](#) [[INSPIRE](#)].
- [35] G. Das, M.C. Kumar and K. Samanta, *Resummed inclusive cross-section in Randall-Sundrum model at NNLO+NNLL*, *JHEP* **07** (2020) 040 [[arXiv:2004.03938](#)] [[INSPIRE](#)].
- [36] D. Westmark and J.F. Owens, *Enhanced threshold resummation formalism for lepton pair production and its effects in the determination of parton distribution functions*, *Phys. Rev. D* **95** (2017) 056024 [[arXiv:1701.06716](#)] [[INSPIRE](#)].
- [37] P. Banerjee, G. Das, P.K. Dhani and V. Ravindran, *Threshold resummation of the rapidity distribution for Higgs production at NNLO+NNLL*, *Phys. Rev. D* **97** (2018) 054024 [[arXiv:1708.05706](#)] [[INSPIRE](#)].
- [38] P. Banerjee, G. Das, P.K. Dhani and V. Ravindran, *Threshold resummation of the rapidity distribution for Drell-Yan production at NNLO+NNLL*, *Phys. Rev. D* **98** (2018) 054018 [[arXiv:1805.01186](#)] [[INSPIRE](#)].
- [39] G. Lusterians, J.K.L. Michel and F.J. Tackmann, *Generalized Threshold Factorization with Full Collinear Dynamics*, [arXiv:1908.00985](#) [[INSPIRE](#)].
- [40] V.V. Sudakov, *Vertex parts at very high-energies in quantum electrodynamics*, *Sov. Phys. JETP* **3** (1956) 65 [[INSPIRE](#)].
- [41] A.H. Mueller, *On the Asymptotic Behavior of the Sudakov Form-factor*, *Phys. Rev. D* **20** (1979) 2037 [[INSPIRE](#)].

- [42] J.C. Collins, *Algorithm to Compute Corrections to the Sudakov Form-factor*, *Phys. Rev. D* **22** (1980) 1478 [[INSPIRE](#)].
- [43] A. Sen, *Asymptotic Behavior of the Sudakov Form-Factor in QCD*, *Phys. Rev. D* **24** (1981) 3281 [[INSPIRE](#)].
- [44] T.O. Eynck, E. Laenen and L. Magnea, *Exponentiation of the Drell-Yan cross-section near partonic threshold in the DIS and MS-bar schemes*, *JHEP* **06** (2003) 057 [[hep-ph/0305179](#)] [[INSPIRE](#)].
- [45] S. Moch, J.A.M. Vermaseren and A. Vogt, *The three loop splitting functions in QCD: The nonsinglet case*, *Nucl. Phys. B* **688** (2004) 101 [[hep-ph/0403192](#)] [[INSPIRE](#)].
- [46] J. Henn, A.V. Smirnov, V.A. Smirnov, M. Steinhauser and R.N. Lee, *Four-loop photon quark form factor and cusp anomalous dimension in the large- N_c limit of QCD*, *JHEP* **03** (2017) 139 [[arXiv:1612.04389](#)] [[INSPIRE](#)].
- [47] S. Moch, B. Ruijl, T. Ueda, J.A.M. Vermaseren and A. Vogt, *Four-Loop Non-Singlet Splitting Functions in the Planar Limit and Beyond*, *JHEP* **10** (2017) 041 [[arXiv:1707.08315](#)] [[INSPIRE](#)].
- [48] A. Grozin, *Four-loop cusp anomalous dimension in QED*, *JHEP* **06** (2018) 073 [*Addendum ibid.* **01** (2019) 134] [[arXiv:1805.05050](#)] [[INSPIRE](#)].
- [49] J.M. Henn, T. Peraro, M. Stahlhofen and P. Wasser, *Matter dependence of the four-loop cusp anomalous dimension*, *Phys. Rev. Lett.* **122** (2019) 201602 [[arXiv:1901.03693](#)] [[INSPIRE](#)].
- [50] J. Davies, A. Vogt, B. Ruijl, T. Ueda and J.A.M. Vermaseren, *Large- n_f contributions to the four-loop splitting functions in QCD*, *Nucl. Phys. B* **915** (2017) 335 [[arXiv:1610.07477](#)] [[INSPIRE](#)].
- [51] R.N. Lee, A.V. Smirnov, V.A. Smirnov and M. Steinhauser, *The n_f^2 contributions to fermionic four-loop form factors*, *Phys. Rev. D* **96** (2017) 014008 [[arXiv:1705.06862](#)] [[INSPIRE](#)].
- [52] J.A. Gracey, *Anomalous dimension of nonsinglet Wilson operators at $O(1/N_f)$ in deep inelastic scattering*, *Phys. Lett. B* **322** (1994) 141 [[hep-ph/9401214](#)] [[INSPIRE](#)].
- [53] M. Beneke and V.M. Braun, *Power corrections and renormalons in Drell-Yan production*, *Nucl. Phys. B* **454** (1995) 253 [[hep-ph/9506452](#)] [[INSPIRE](#)].
- [54] S. Moch, B. Ruijl, T. Ueda, J.A.M. Vermaseren and A. Vogt, *On quartic colour factors in splitting functions and the gluon cusp anomalous dimension*, *Phys. Lett. B* **782** (2018) 627 [[arXiv:1805.09638](#)] [[INSPIRE](#)].
- [55] R. Brüser, A. Grozin, J.M. Henn and M. Stahlhofen, *Matter dependence of the four-loop QCD cusp anomalous dimension: from small angles to all angles*, *JHEP* **05** (2019) 186 [[arXiv:1902.05076](#)] [[INSPIRE](#)].
- [56] R.N. Lee, A.V. Smirnov, V.A. Smirnov and M. Steinhauser, *Four-loop quark form factor with quartic fundamental colour factor*, *JHEP* **02** (2019) 172 [[arXiv:1901.02898](#)] [[INSPIRE](#)].
- [57] A. von Manteuffel, E. Panzer and R.M. Schabinger, *Cusp and collinear anomalous dimensions in four-loop QCD from form factors*, *Phys. Rev. Lett.* **124** (2020) 162001 [[arXiv:2002.04617](#)] [[INSPIRE](#)].
- [58] A. Vogt, S. Moch and J.A.M. Vermaseren, *The three-loop splitting functions in QCD: The singlet case*, *Nucl. Phys. B* **691** (2004) 129 [[hep-ph/0404111](#)] [[INSPIRE](#)].

- [59] S. Moch, J.A.M. Vermaseren and A. Vogt, *The Three-Loop Splitting Functions in QCD: The Helicity-Dependent Case*, *Nucl. Phys. B* **889** (2014) 351 [[arXiv:1409.5131](#)] [[INSPIRE](#)].
- [60] J.M. Henn, G.P. Korchemsky and B. Mistlberger, *The full four-loop cusp anomalous dimension in $\mathcal{N} = 4$ super Yang-Mills and QCD*, *JHEP* **04** (2020) 018 [[arXiv:1911.10174](#)] [[INSPIRE](#)].
- [61] T. Ahmed, G. Das, P. Mathews, N. Rana and V. Ravindran, *Spin-2 Form Factors at Three Loop in QCD*, *JHEP* **12** (2015) 084 [[arXiv:1508.05043](#)] [[INSPIRE](#)].
- [62] A. Vogt, *Efficient evolution of unpolarized and polarized parton distributions with QCD-PEGASUS*, *Comput. Phys. Commun.* **170** (2005) 65 [[hep-ph/0408244](#)] [[INSPIRE](#)].
- [63] S. Catani, M.L. Mangano, P. Nason and L. Trentadue, *The resummation of soft gluons in hadronic collisions*, *Nucl. Phys. B* **478** (1996) 273 [[hep-ph/9604351](#)] [[INSPIRE](#)].
- [64] A. Buckley et al., *LHAPDF6: parton density access in the LHC precision era*, *Eur. Phys. J. C* **75** (2015) 132 [[arXiv:1412.7420](#)] [[INSPIRE](#)].
- [65] L.A. Harland-Lang, A.D. Martin, P. Motylinski and R.S. Thorne, *Parton distributions in the LHC era: MMHT 2014 PDFs*, *Eur. Phys. J. C* **75** (2015) 204 [[arXiv:1412.3989](#)] [[INSPIRE](#)].
- [66] S. Catani, D. de Florian and M. Grazzini, *Direct Higgs production and jet veto at the Tevatron and the LHC in NNLO QCD*, *JHEP* **01** (2002) 015 [[hep-ph/0111164](#)] [[INSPIRE](#)].
- [67] S. Catani, D. de Florian and M. Grazzini, *Higgs production in hadron collisions: Soft and virtual QCD corrections at NNLO*, *JHEP* **05** (2001) 025 [[hep-ph/0102227](#)] [[INSPIRE](#)].
- [68] C. Anastasiou et al., *High precision determination of the gluon fusion Higgs boson cross-section at the LHC*, *JHEP* **05** (2016) 058 [[arXiv:1602.00695](#)] [[INSPIRE](#)].
- [69] S. Alekhin, J. Bluemlein, S.-O. Moch and R. Placakyte, *The new ABMP16 PDF*, *PoS DIS2016* (2016) 016 [[arXiv:1609.03327](#)] [[INSPIRE](#)].
- [70] S. Dulat et al., *New parton distribution functions from a global analysis of quantum chromodynamics*, *Phys. Rev. D* **93** (2016) 033006 [[arXiv:1506.07443](#)] [[INSPIRE](#)].
- [71] NNPDF collaboration, *Parton distributions from high-precision collider data*, *Eur. Phys. J. C* **77** (2017) 663 [[arXiv:1706.00428](#)] [[INSPIRE](#)].
- [72] J. Butterworth et al., *PDF4LHC recommendations for LHC Run II*, *J. Phys. G* **43** (2016) 023001 [[arXiv:1510.03865](#)] [[INSPIRE](#)].
- [73] C. Duhr, F. Dulat and B. Mistlberger, *Charged Current Drell-Yan Production at N3LO*, [[arXiv:2007.13313](#)] [[INSPIRE](#)].
- [74] J.A.M. Vermaseren, *New features of FORM*, [[math-ph/0010025](#)] [[INSPIRE](#)].
- [75] B. Ruijl, T. Ueda and J. Vermaseren, *FORM version 4.2*, [[arXiv:1707.06453](#)] [[INSPIRE](#)].