

Explanation of the ATLAS Z-peaked excess by squark pair production in the NMSSM

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ABSTRACT: The ATLAS collaboration recently reported a 3σ excess in the leptonic- $Z + jets + E_T^{miss}$ channel. We intend to interpret this excess by squark pair production in the Next-to-Minimal Supersymmetric Standard Model (NMSSM). The decay chain we employ is $\tilde{q} \rightarrow q\tilde{\chi}_2^0 \rightarrow q\tilde{\chi}_1^0 Z$, where $\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0$ denote the lightest and the next-to-lightest neutralinos with singlino and bino as their dominant components respectively. Our simulations indicate that after considering the constraints from the ATLAS searches for $jets + E_T^{miss}$ signal the central value of the excess can be obtained for $m_{\tilde{q}} \lesssim 1.2\text{TeV}$, and if the constraint from the CMS on- Z search is further considered, more than 10 signal events are still attainable for $m_{\tilde{q}} \lesssim 750\text{GeV}$. Compared with the interpretation by gluino pair production, the squark explanation allows for a significantly wider range of $m_{\tilde{q}}$ as well as a less compressed SUSY mass spectrum. We also show that the squark explanation will be readily tested at the initial stage of the 14 TeV LHC.

KEYWORDS: Supersymmetry Phenomenology

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Contents

1	Introduction	1
2	Our scenario for the Z-excess	2
3	Z-peaked excess in our scenario	5
4	Distributions of the excess	10
5	Test of our explanation at the 14 TeV LHC	12
6	Conclusion	13

1 Introduction

Since the discovery of the Higgs-like particle by the ATLAS and CMS collaborations at the Large Hadron Collider (LHC) in 2012 [1, 2], the main task of the LHC program has shifted to the searches for new physics beyond the Standard Model (SM). These searches cover a wide range of possible signatures, notably various combinations of jets (with or without b-tagging), the missing transverse energy (E_T^{miss}) and/or leptons. In this direction, the ATLAS collaboration has recently reported an intriguing excess at 3σ significance in the leptonic- $Z + jets + E_T^{\text{miss}}$ channel [3]. Based on the full dataset in 2012, which corresponds to about $20.3 fb^{-1}$ integrated luminosity at the 8 TeV LHC, the collaboration observed 29 events for the on- Z electron and muon pair channels in contrast with the expected SM background number 10.6 ± 3.2 , and no excess over the SM background was observed in any other signal region (SR) [3].

So far there have been several attempts to explain the excess by the production of new physics particles, which are needed to decay with a sizeable rate into jets as well as at least one Z boson and one invisible particle [4–13]. In the context of supersymmetric theories (SUSY), in order to provide enough events after the rather tight cuts employed in [3], the production of a pair of moderately light gluinos was usually utilized for the explanation [4, 6–9, 12, 13]. The key point in doing this is to choose a lightest SUSY particle (LSP) with relatively suppressed couplings to squarks, so that the gluino prefers to decay first into the neutralino other than the LSP, and the neutralino subsequently decays into the LSP plus a Z boson. In the Minimal Supersymmetric standard model (MSSM), a higgsino-dominated neutralino has very weak couplings to light flavor squarks, so one can naturally imagine a higgsino-dominated LSP and assume the decay chain $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_i^0 \rightarrow q\bar{q}Z\tilde{\chi}_1^0$ (q represents a light flavor quark, and $\tilde{\chi}_i^0$ denotes a gaugino-dominated neutralino) to interpret the excess.¹ For this case, we note that the measured dark matter relic density is not easy

¹In this case, the higgsino-dominated LSP plays the same role as the gravitino in the ATLAS report [3] to interpret the excess.

to obtain if only the neutralino serves as the dark matter candidate (see for example, figure1 in [14]). In the Next-to Minimal Supersymmetric Standard Model (NMSSM) [15], however, a singlino-dominated neutralino also has the property, and meanwhile if it acts as the LSP, the correct relic density can be achieved by multiple annihilation channels [15, 16]. So in this work we are interested in the interpretation of the Z -excess in the NMSSM with a singlino-dominated LSP.

In the framework of the NMSSM with a singlino-dominated LSP, the gluino pair production with the three-body decay $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_2^0 \rightarrow q\bar{q}Z\tilde{\chi}_1^0$ has been studied for the Z -excess in [6, 9]. These works indicated that only with simple assumptions on the relevant model parameters can the NMSSM explain the excess quite well. Explicitly speaking, it was found that after considering the constraints from the ATLAS searches for $jets + E_T^{\text{miss}}$ signal, the NMSSM can reproduce the central value of the excess, and even if one further considers the constraint from the CMS search for the leptonic- $Z + jets + E_T^{\text{miss}}$ channel which observed no excess in all SRs, the event number of the ATLAS on- Z signal can still reach 11, which is about 1.2σ away from the measured central value [9]. Moreover, as illustrated in [6] the gluino explanation can reproduce well various distributions of the excess presented by the ATLAS collaboration. Albeit these advantages, we still think that it is necessary to seek for other explanations since in the gluino explanation, the gluino mass is limited in a narrow range, and at same time the relevant sparticle mass spectrum must be rather compressed to escape the constraints (see figure2 of [9]). In this work, we consider squark pair production as an explanation of the excess, and in order to compare it with the gluino explanation, we make similar assumptions on the model parameters to those of [9]. We find that in the squark explanation the range of the squark mass can be significantly extended in comparison with the gluino explanation, and moreover the relevant sparticle mass spectrum may become less compressed. We also find that, just like the gluino explanation, the distributions of the excess can also be reproduced well in the squark explanation.

This work is organized as follows. In section 2, we briefly introduce our scenario for the excess. In section 3, we perform a comprehensive analysis of the relevant parameter space and present the results of our simulations on the Z -peaked excess. In section 4, we choose some representative parameter points, and exhibit their predictions on various distributions of the excess in comparison with the corresponding data provided by the ATLAS collaboration. In section 5, we briefly discuss future test of our scenario at the 14 TeV LHC. At last, we draw our conclusions in section 6.

2 Our scenario for the Z -excess

As one of the most economical extensions of the MSSM, the NMSSM contains one gauge singlet Higgs superfield \hat{S} in its matter content. The superpotential of the general NMSSM is given by [15, 17]

$$W_{\text{NMSSM}} = W_{\text{MSSM}} + \lambda \hat{S} \hat{H}_u \cdot \hat{H}_d + \xi_F \hat{S} + \frac{\mu'}{2} \hat{S}^2 + \frac{\kappa}{3} \hat{S}^3, \quad (2.1)$$

where W_{MSSM} is the superpotential of the MSSM without the μ term, \hat{H}_u and \hat{H}_d are the $SU(2)_L$ doublet superfields, κ and λ are dimensionless coefficients, ξ_F parameterizes the tadpole term and μ' is a supersymmetric mass.

In this framework, the fermionic component field of \hat{S} which is usually called singlino \tilde{S} will mix with the gauginos and the higgsinos of the MSSM to form neutralinos. In the basis $(\psi_1 \equiv -i\tilde{B}, \psi_2 \equiv -i\tilde{W}^0, \psi_3 \equiv \tilde{H}_u^0, \psi_4 \equiv \tilde{H}_d^0, \psi_5 \equiv \tilde{S})$, the corresponding mass matrix is given by [15]

$$\mathcal{M} = \begin{pmatrix} M_1 & 0 & \frac{ev_u}{\sqrt{2}c_w} & -\frac{ev_d}{\sqrt{2}c_w} & 0 \\ 0 & M_2 & -\frac{ev_u}{\sqrt{2}s_w} & \frac{ev_d}{\sqrt{2}s_w} & 0 \\ \frac{ev_u}{\sqrt{2}c_w} & -\frac{ev_u}{\sqrt{2}s_w} & 0 & -\mu_{\text{eff}} & -\lambda v_d \\ -\frac{ev_d}{\sqrt{2}c_w} & \frac{ev_d}{\sqrt{2}s_w} & -\mu_{\text{eff}} & 0 & -\lambda v_u \\ 0 & 0 & -\lambda v_d & -\lambda v_u & 2\kappa s + \mu' \end{pmatrix}, \quad (2.2)$$

where M_1 and M_2 are soft gaugino masses, $v_u = v \sin \beta$ and $v_d = v \cos \beta$ are vacuum expectation values (vev) of the Higgs fields H_u and H_d respectively, $\mu_{\text{eff}} = \mu + \lambda s$ with s denoting the vev of the singlet scalar field S , and $c_w = \cos \theta_W$. This matrix can be diagonalized by a 5×5 unitary matrix N , and consequently neutralinos as mass eigenstates are defined by

$$\tilde{\chi}_i^0 = \sum_{j=1}^5 N_{ij} \psi_j,$$

where the mass order $m_{\tilde{\chi}_1^0} < m_{\tilde{\chi}_2^0} < \dots < m_{\tilde{\chi}_5^0}$ is assumed. Obviously, the matrix element N_{ij} measures the size of the ψ_j component in $\tilde{\chi}_i^0$ state, and for the singlino-dominated and bino-dominated neutralinos, their masses are mainly determined by the combination $2\kappa s + \mu'$ and M_1 respectively. Moreover, with the help of N_{ij} one can get the interactions of the neutralinos. As shown in [9], the $\tilde{q}\tilde{\chi}_i^0\tilde{q}$ coupling with q denoting a light flavor quark is determined by the gaugino components of $\tilde{\chi}_i^0$, and the $\tilde{\chi}_i^0\tilde{\chi}_j^0Z$ coupling is determined by the higgsino components of the neutralinos. By contrast, the $\tilde{\chi}_i^0\tilde{\chi}_j^0h$ coupling with h denoting the SM-like Higgs boson depends on all components of the neutralinos, and there may exist cancelations among different contributions. These characters are helpful to understand our explanation of the Z -excess.

In the following, we intend to interpret the ATLAS on- Z excess by squark pair production. To make our explanation as simple as possible, we have the following assumptions:

- Only the first and second generation squarks are responsible for the excess. In our analysis, we assume a common mass $m_{\tilde{q}}$ for the squarks, then the cross section for the squark pair production depends only on $m_{\tilde{q}}$ and $m_{\tilde{g}}$. We calculate the cross section at the NLO with the code Prospino [18], and show its dependence on $m_{\tilde{q}}$ at the 8 TeV LHC in figure 1.
- The leptonic- $Z + jets + E_T^{\text{miss}}$ signal is generated by the cascade decay $\tilde{q} \rightarrow q\tilde{\chi}_2^0 \rightarrow q\tilde{\chi}_1^0Z$. In order to maximize this signal rate, we require both $Br(\tilde{q} \rightarrow q\tilde{\chi}_2^0)$ and $Br(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0Z)$ at roughly 100%, where the former requirement can be satisfied if $\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0$ are singlino-dominated and bino-dominated respectively, and only the two particles in the neutralino and chargino sector are lighter than the squarks, and the latter condition can be realized if $m_Z < m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0} \leq m_h$ or if the $\tilde{\chi}_i^0\tilde{\chi}_j^0h$ interaction is significantly suppressed (see above discussion and also our previous work [9]).

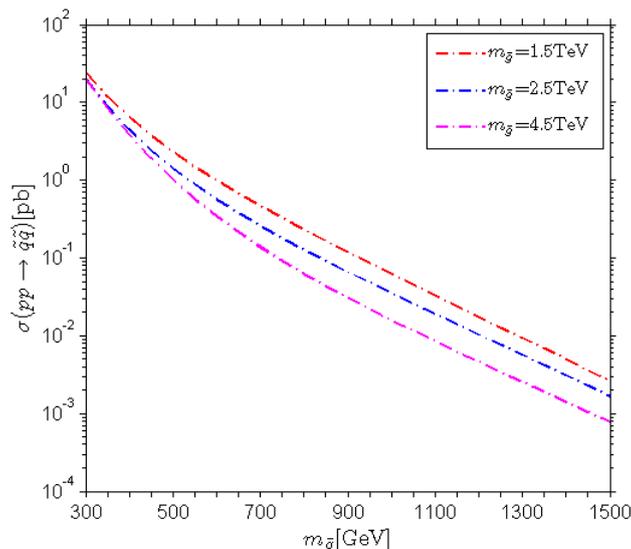


Figure 1. The cross section of the squark pair production at the 8 TeV LHC for different gluino mass.

- With the above assumptions, the parameters involved in our explanation are $m_{\tilde{q}}$, $m_{\tilde{g}}$, $\Delta m_1 \equiv m_{\tilde{q}} - m_{\tilde{\chi}_2^0}$ and $\Delta m_2 \equiv m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$. In our discussion, we vary these parameters freely, but noting that the process $pp \rightarrow \tilde{q}\tilde{q} \rightarrow \tilde{\chi}_1^0 Z q \tilde{\chi}_1^0 Z q$ can also generate multi-jets + E_T^{miss} signal, we limit these parameters by the ATLAS searches for the multi-jets signal, which were presented in [19, 20]. We also consider the CMS search for the leptonic- $Z + jets + E_T^{\text{miss}}$ signal [21] as an alternative constraint on the parameters.

About our scenario for the Z -excess, we have following additional remarks:

- We ad hoc require that only $\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0$ among the neutralinos are lighter than the squarks. This will simplify our analysis, but on the other hand, since the rate and various kinematic distributions of the process $pp \rightarrow \tilde{q}\tilde{q} \rightarrow \tilde{\chi}_1^0 Z q \tilde{\chi}_1^0 Z q$ are decided by few parameters, the capability of our scenario to interpret the excess is limited given that the scenario must satisfy the constraints mentioned above. In fact, as implied by the results of [11], allowing the squarks to decay in multiple ways facilitates SUSY to balance the ATLAS signal and the constraints, and is thus able to explain the excess in a better way. This, however, needs an intensive scan over a higher dimensional SUSY parameter space, and for each parameter point, simulations on various SUSY signals have to be done to compare with the corresponding experimental data. Such calculations are very time-consuming, and are beyond the capability of our cluster.
- Again for the sake of simplicity, we do not consider the effect of the third generation squarks in our analysis. These squarks have been tightly limited by the SM-like Higgs boson mass, and are preferred to be heavy [22–24]. For some optimized points in figure 2 for the Z -excess, we once included their contributions to the leptonic- $Z + jets + E_T^{\text{miss}}$ and the multi-jets + E_T^{miss} signals by assuming the degeneracy of all

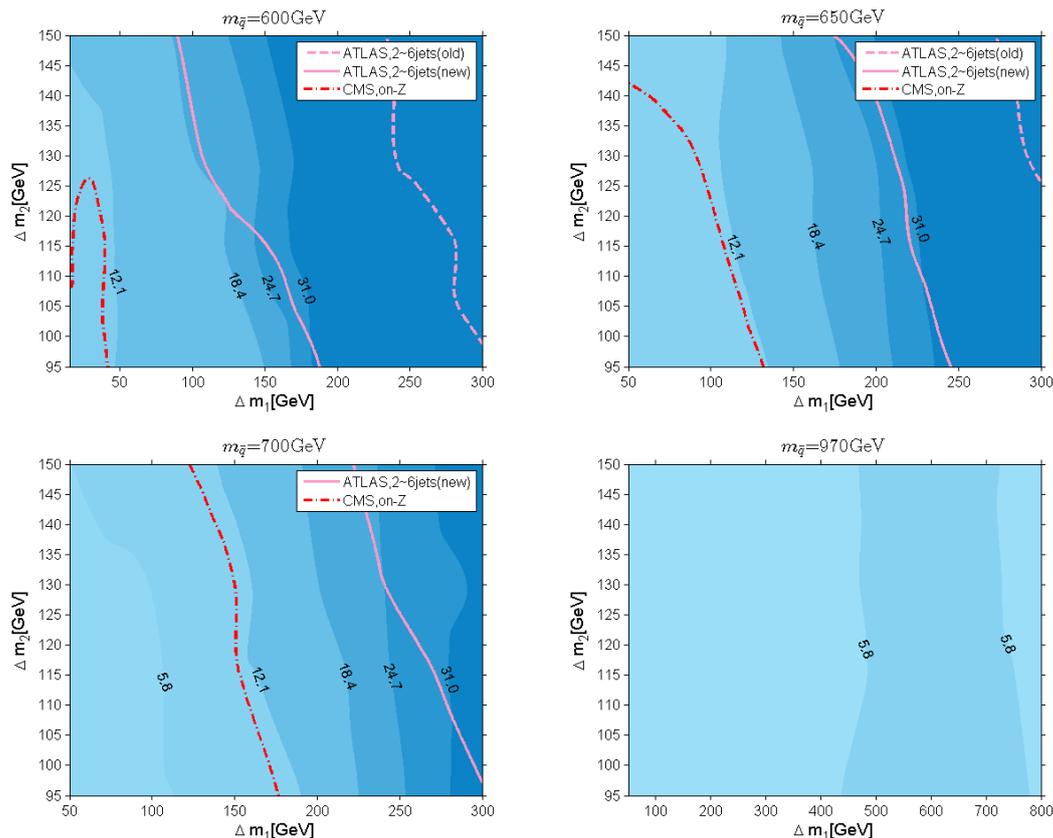


Figure 2. Constant contours of the event number for the ATLAS leptonic- $Z + jets + M_T^{\text{miss}}$ signal on the $\Delta m_1 - \Delta m_2$ plane in heavy gluino case. For each $m_{\tilde{g}}$, the region between the contour marked by the number 12.1 and that by 24.7 can explain the excess at 1σ level, and that between the lines marked by 5.8 and 31 respectively can account for the excess at 2σ level. The parameter spaces that coincide with different SUSY searches are also presented, which are right bounded by different types of lines. The dotted line and solid line are the boundaries coming from the ATLAS preliminary and updated searches for $2 \sim 6 jets + E_T^{\text{miss}}$ signal respectively, and the dash dotted line is for the CMS constraint. For the case $m_{\tilde{g}} = 700 \text{ GeV}$, the constraint from the ATLAS preliminary search for $2 \sim 6 jets + E_T^{\text{miss}}$ signal is too weak to be drawn on the plane, and for the case $m_{\tilde{g}} = 970 \text{ GeV}$, there are actually no boundaries on the plane.

squarks. However, we did not find any improvement on our explanation due to the constraints we considered.

- The assumptions on the properties of the LSP and NLSP in this work are same as those of our previous work [9], where gluino pair production was used for the excess. This enables us to compare directly the two explanations.

3 Z-peaked excess in our scenario

From the ATLAS analysis on the leptonic- $Z + jets + E_T^{\text{miss}}$ channel presented in [3], one can infer that the event number of the excess is 18.4 ± 6.3 after including the statistical and systematic uncertainties [9]. This means that, if one wants to explain the excess at

1σ and 2σ levels, the SUSY signal number after cuts should satisfy $12.1 \leq N_{ll} \leq 24.7$ and $5.8 \leq N_{ll} \leq 31$ respectively. In order to find the parameter space that can produce the required event number, we fix $m_{\tilde{g}} = 4.5\text{TeV}$ (heavy gluino case) and $m_{\tilde{g}} = 1.5\text{TeV}$ (light gluino case) separately, and for each case, we perform a grid scan over the parameters Δm_1 and Δm_2 by choosing a series of $m_{\tilde{q}}$. For each parameter point, we calculate the squark pair production rate at the 8 TeV LHC by the package Prospino [18], and generate the parton level events for the considered process with MG5_aMC [25], which includes Pythia [26] for parton showering and hadronization. Then we use the package CheckMATE-1.2.0 [27, 28] which contains fine-tuned fast detector simulation code Delphes3.0.10 [29] to repeat the analyses of various experiments. These experiments include the ATLAS on-Z search [3], the CMS on-Z search [21], and the ATLAS $2 \sim 6 \text{ jet} + E_T^{\text{miss}}$ searches [19, 20], among which the first one is used to generate the signal of the excess, and the other ones serve as constraints. In [9], we encoded the cuts for those experiments in the package CheckMATE-1.2.0, and the validation of them indicated that our calculations coincide with the corresponding experimental analyses at 20% level. In implementing the constraints of the SUSY searches on the parameters, we define for each search the ratio $R = \max(N_{S,i}/S_{obs,i}^{95\%})$, where $N_{S,i}$ is the event number of the SUSY signal in the i th SR of the search, $S_{obs,i}^{95\%}$ is its 95% upper limit usually provided in the experimental report, and the \max is over all SRs defined in the search. Obviously, only in case that $R < 1$, the corresponding parameter point is experimentally allowed at 95% C.L..

In figure 2 and figure 3, we present on $\Delta m_1 - \Delta m_2$ plane the constant contours of the event number for the ATLAS on-Z analysis in the heavy gluino and light gluino cases respectively. For each $m_{\tilde{q}}$, the region between the contour marked by the number 12.1 and that by 24.7 can explain the excess at 1σ level, and that between the contours marked by 5.8 and 31 respectively can account for the excess at 2σ level. The parameter spaces that coincide with various SUSY searches are also presented, which are right bounded by different types of lines (note that a compressed SUSY mass spectrum is helpful to evade the LHC constraints). The dotted line and solid line are the boundaries coming from the ATLAS preliminary and updated searches for $2 \sim 6 \text{ jets} + E_T^{\text{miss}}$ signal respectively, and they are obtained by setting the corresponding R values at 1. The constraint from the CMS on-Z search is obtained in a similar way, and shown by dash dotted lines.

From figure 2 for the heavy gluino case, one can learn following facts:

- For all the choices of $m_{\tilde{q}}$, the strongest constraint comes from the CMS dedicated on-Z counting experiment, and the weakest one is the ATLAS preliminary search for $2 \sim 6 \text{ jets} + E_T^{\text{miss}}$ signal.
- With the increase of the squark mass, Δm_1 is allowed to vary within a wider range. In this case, the improved cut efficiency due to the enlarged Δm_1 can compensate for the decrease of the squark pair production rate. As a result, even for $m_{\tilde{q}} \lesssim 800\text{GeV}$ the central value of the excess (18.4 events) can still be obtained if only the constraints from the ATLAS searches for the $\text{jets} + E_T^{\text{miss}}$ signal are considered.
- In the CMS dedicated on-Z counting experiment, six signal regions discriminated by the jet number n_j and E_T^{miss} were considered (see table 2 in [9]). We checked that

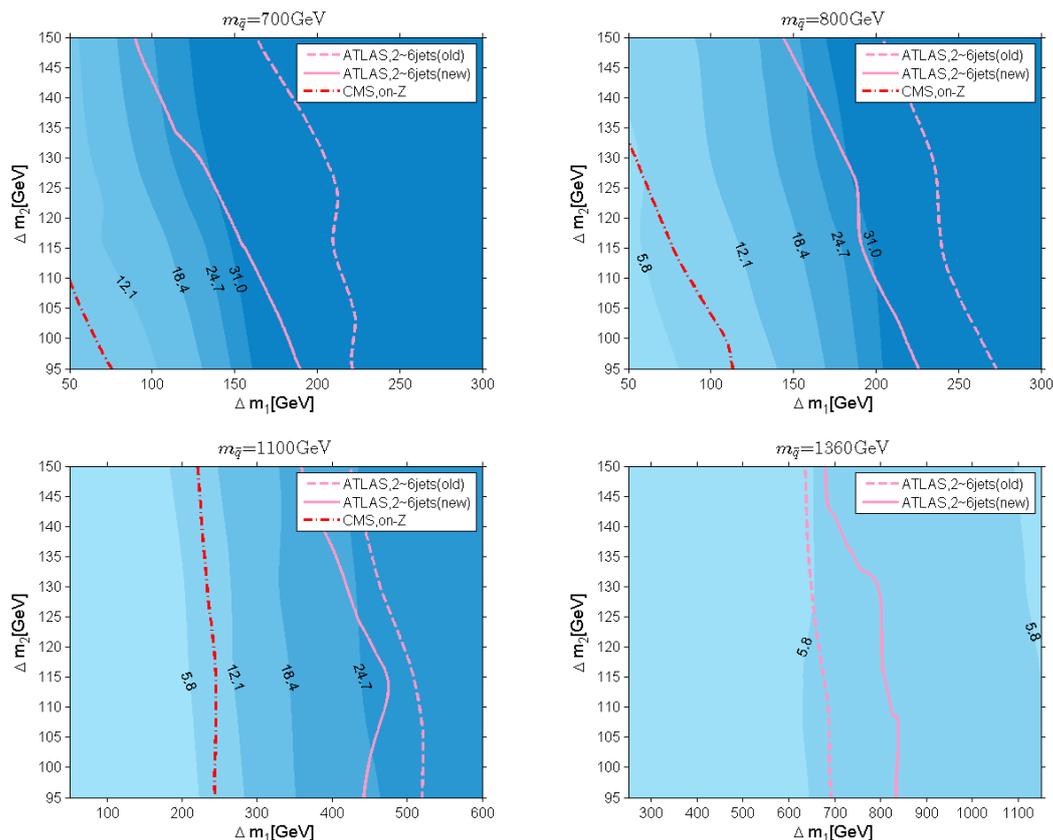


Figure 3. Similar to figure 2, but for the light gluino case, $m_{\tilde{g}} = 1.5\text{TeV}$. Note that for $m_{\tilde{q}} = 1360\text{GeV}$, the CMS constraint is too weak to be drawn on the plane.

for our scenario the tightest constraint of the experiment comes from $n_j \geq 2$ SRs with E_T^{miss} either satisfying $200\text{GeV} \leq E_T^{\text{miss}} \leq 300\text{GeV}$ (called SR-II hereafter) or satisfying $E_T^{\text{miss}} > 300\text{GeV}$ (called SR-III hereafter).² In either case, the signal of the ATLAS on- Z search has a large overlap with that of the CMS on- Z search, so due to the tension of the two search results the event number of the ATLAS experiment is always upper bounded by about 11 for $m_{\tilde{q}} \lesssim 750\text{GeV}$ after considering the CMS constraint. We also checked that, with the further increase of $m_{\tilde{q}}$ from about 750 GeV, the maximal reach of the event number drops either because the tension of the ATLAS and CMS data becomes stronger for moderately heavy \tilde{q} or because the squark pair production rate is sufficiently suppressed for heavy \tilde{q} .

- The lower right panel of figure 2 indicates that there are actually no boundaries on the $\Delta m_1 - \Delta m_2$ plane for $m_{\tilde{q}} = 970\text{GeV}$. In this case, the maximal reach of the event number is 6.8, which is still within the 2σ range of the excess.

²In more detail, from our calculations we learn that the SR-II is more powerful than the SR-III in limiting our scenario for the region defined by $100\text{GeV} \lesssim \Delta m_1 \lesssim 150\text{GeV}$ and $95\text{GeV} \lesssim \Delta m_2 \lesssim 150\text{GeV}$ as well as that defined by $150\text{GeV} \lesssim \Delta m_1 \lesssim 200\text{GeV}$ and $95\text{GeV} \lesssim \Delta m_2 \lesssim 120\text{GeV}$ for all panels in figure 2. Consequently, the boundaries of the CMS experiment in the cases of $m_{\tilde{q}} = 650, 700\text{GeV}$ are mainly determined by the SR-II, while in the case of $m_{\tilde{q}} \gtrsim 750\text{GeV}$, they are determined by the SR-III.

Point	$\sigma(\text{fb})$	CMS [21]				ATLAS(updated) [19]				ATLAS(preliminary) [20]			
		SR _{max}	ϵ	S _{obs} ^{95%}	R	SR _{max}	ϵ	S _{obs} ^{95%}	R	SR _{max}	ϵ	S _{obs} ^{95%}	R
S ₁	32.9	SR-III	3.3%	7.6	2.77	2jm	14.7%	90	1.09	CM	9.2%	81.2	0.75
S ₂	6.3	SR-III	4.2%	7.6	0.67	2jt	21.0%	38	0.71	BM	12.4%	14.9	1.06

Table 1. Details about the constraints in light gluino case for two points S_1 and S_2 , which are taken from the lower left and right panels in figure 3 respectively, and are defined by S_1 : $m_{\tilde{q}} = 1100\text{GeV}$, $\Delta m_1 = 458\text{GeV}$, $\Delta m_2 = 105\text{GeV}$ and S_2 : $m_{\tilde{q}} = 1360\text{GeV}$, $\Delta m_1 = 687\text{GeV}$, $\Delta m_2 = 105\text{GeV}$. In this table, SR_{max} represents the SR which provides the strongest constraint on the parameter for a certain experiment, and ϵ is the total cut efficiency of the SR.

Next we turn to the light gluino case. From the event contours in figure 3, one can get following information:

- The dependence of the ATLAS event number on the parameters Δm_1 and Δm_2 is quite similar to that in the heavy gluino case, and so are the dependencies of the constraints.
- For squarks lighter than about 1200GeV, the tightest constraint comes from the CMS experiment, which is similar to the heavy gluino case; but with the increase of the squark mass, the strongest constraint may come from the ATLAS preliminary search for $jets + E_T^{\text{miss}}$ signal, which is shown in the lower right panel of figure 3. The underlying reason for such a feature is that the capabilities of the SRs defined in the experiments in limiting SUSY depend on the configuration of the SUSY spectrum, and without the specification of the spectrum, there is no definite conclusion about which is the strongest. To illustrate this, we take one point from each of the two lower panels in figure 3 as an example, and show the details of the constraints in table 1. This table indicates that for the points S_1 and S_2 , the tightest constraints of the ATLAS preliminary search for $jets + E_T^{\text{miss}}$ are the SRs CM and BM respectively, and those of the ATLAS updated search correspond to the SRs 2jm and 2jt respectively. For the point S_1 , the constraint from the former experimental analysis is weaker, while for the point S_2 , the situation reverses.
- We checked that the ATLAS excess can be explained at 2σ level by squarks with mass up to about 1.4TeV. The wider mass range in comparison with the heavy gluino case is mainly due to the larger rate of the squark pair production in light gluino case. We also checked that the central value of the excess can be achieved for $m_{\tilde{q}} \lesssim 1.2\text{TeV}$ if only the constraints from the ATLAS searches for the $jets + E_T^{\text{miss}}$ signal are considered, and that if the constraint from the CMS experiment is further considered, the maximal event number is only 9.5. The latter fact implies that the heavy gluino case is able to provide a slightly better explanation.
- Note that for the lower right panel where $m_{\tilde{q}} = 1.36\text{TeV}$ and $m_{\tilde{g}} = 1.5\text{TeV}$, the effect from the squark-gluino associated production on the LHC searches is non-negligible. Discussing such an effect is beyond the scope of this work.

Before we end this section, we have four comments about our explanation. First, comparing figure 2 and figure 3 in this work with figure 2 in [9] where gluino pair production with the decay mode $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_2^0 \rightarrow q\bar{q}Z\tilde{\chi}_1^0$ was used to explain the excess, we conclude that the squark explanation allows for a significantly wider range of $m_{\tilde{q}}$ as well as a less compressed SUSY mass spectrum. One underlying reason we think is that the cut efficiency of the ATLAS on- Z search is usually larger for the squark explanation than for the gluino explanation, and consequently, moderately heavy squarks are still capable of explaining the excess. Second, we emphasize again that in our simple scenario, both the event number of the ATLAS signal and the constraints are determined by few SUSY parameters. Consequently, the capability of our scenario to interpret the excess is limited. As we mentioned at the end of section 2, a more complex scenario with higher dimensional SUSY parameters may improve this situation, but it needs tremendous calculation to search for the relevant parameter space [11]. Third, we note that so far there have appeared other simple SUSY scenarios to explain the excess, and the studies of these scenarios indicated that they can explain the excess at 1σ level in certain narrow SUSY parameter spaces [12, 13]. This conclusion is slightly better than ours, where the best explanation is about 1.2σ away from the central value of the excess. Three factors may contribute to the difference:

- The difference in the theoretical hypothesis on SUSY, which determines the kinematical distributions of the SUSY signals. For example, both [12] and [13] utilized the production $pp \rightarrow \tilde{g}\tilde{g}$ with the loop-induced decay $\tilde{g} \rightarrow g\tilde{\chi}_i^0 \rightarrow gZ\chi$ (χ denotes the lightest neutralino in [12] and the gravitino in [13]) to explain the excess. Comparing their interpretations with ours, one can learn that, although all of them considered the two-body decay of a strongly produced SUSY particle, due to the difference of the properties for the parent sparticles such as their spins and production channels, their kinematical distributions may differ greatly even when their predictions on the event number of the excess are same. One can get this point by comparing the E_{miss}^T distribution of the benchmark points P_1 and P_2 in this paper with that of the best point in [13], which are presented in figure 4 of this work and in figure 6 of [13] respectively. As a result of the difference, in principle there might exist SUSY points for which the ATLAS leptonic- Z signal is moderately enhanced and meanwhile the CMS signal is appropriately suppressed.
- The uncertainties induced by related simulations. For all scenarios to explain the excess, simulations of the experimental searches for SUSY have to be done. As shown in the appendices of [9] and [13] where the validations of the simulations were explicitly presented, the uncertainties of the simulations are at the level of 20%, and therefore, the calculation performed by different groups may result in a significant deviation. As far as our simulations are concerned, the computed efficiency for the ATLAS signal event is less than that presented by the ATLAS collaboration for the selected SUSY point by about 10%, and by contrast our efficiency for the CMS search is slightly larger than that in the CMS report.
- The treatment of the CMS constraint. From the CMS report presented in [21], one can only infer the approximate value of the S_{obs}^{95} for the SR-III. Confronted with such

a situation, we in our previous work [9] calculated the S_{obs}^{95} s for all the six SRs by the asymptotic CL_s prescription [30] (see table 2 of [9]). Furthermore, we pointed out that the SR-II may provide a stronger constraint on the parameters than the other SRs in discussing figure 2 of this work. This conclusion indicates that the calculation of all S_{obs}^{95} s is necessary; but on the hand, since the values of the S_{obs}^{95} s were not explicitly given in other previous literatures and meanwhile they depend on calculation method, there might exist deviations for different authors in considering the CMS constraint. Let's show this point by an explicit example. During the revision of this manuscript, the paper [31] appeared to interpret the excess in the NMSSM extension with a Dirac gluino, and the authors presented the details about their calculation of S_{obs}^{95} s. Briefly speaking, the calculation in [31] differs from ours in at least two aspects. One is the authors of [31] used the standard Bayesian procedure in the calculation, while we used the asymptotic CL_s method [30]. The other is that the work [31] had considered the theoretical uncertainty in calculating the signal, and by contrast, we ignored such an effect. As a result of the differences, the values of S_{obs}^{95} s in [31] are usually larger than our predictions by about 15%, and consequently, the CMS constraint is significantly relaxed in [31].

Finally, we note that in our scenario the singlino-dominated LSP is usually heavier than about 450GeV, and one may wonder how such a heavy dark matter (DM) achieves its measured relic density. In this case, the possible annihilation final states of the DM include $f\bar{f}$, VV , H_iH_j , A_iA_j and H_iA_j , where f (V) denotes any of the fermions (vector bosons) in the SM, and H_i (A_j) denotes a CP-even (CP-odd) Higgs boson (see [15] and references therein). The easiest way to achieve the density is through the s -channel annihilations mediated by a singlet-dominated Higgs boson, where just like the light DM case discussed in [16], the Higgs boson mass as well as the self-coupling coefficient κ for the singlet fields play an important role in tuning the annihilation rate.

4 Distributions of the excess

In this section, we investigate whether our explanation can reproduce the distributions of the excess reported by the ATLAS collaboration. For this end, we concentrate on three benchmark points P1, P2 and P3 with point P3 taken from our previous work [9]. Point P1 and P3 correspond to the best points after considering all the constraints in the squark explanation and the gluino explanation respectively, and contribute 11 and 10.5 events to the excess. By contrast, point P2 only satisfies the constraints from the ATLAS jets + E_T^{miss} searches, but it can reproduce the central value of the excess. Detailed information of these points is presented in table 2.

In order to compare our explanation with the experimental data for various distributions, we generate the distributions of E_T^{miss} , H_T (the scalar sum of the P_T s for the leptons and signal jets) and the jet multiplicity n_j in the electron and muon combined channel for each parameter point. In getting the distributions of E_T^{miss} and H_T , we include the overflow events into the last bin. The corresponding results are shown in figure 4, where the black

	P_1	P_2	P_3	SM
$m_{\tilde{g}}$ (GeV)	4500	4500	750	—
$m_{\tilde{q}}$ (GeV)	650	700	4500	—
$m_{\tilde{\chi}_2^0}$ (GeV)	530	500	650	—
$m_{\tilde{\chi}_1^0}$ (GeV)	430	370	565	—
Events in ATLAS on-Z SR($8\text{TeV}, 20.3\text{fb}^{-1}$)	11.0	18.2	10.5	10.6
$\chi^2_{E_T^{\text{miss}}}$ (d.o.f. = 9)	7.6	5.6	8.0	14.8
$\chi^2_{H_T}$ (d.o.f. = 7)	5.7	2.6	6.0	13.8
$\chi^2_{n_j}$ (d.o.f. = 5)	6.9	4.7	6.8	14.1
χ^2_{total} (d.o.f. = 21)	20.2	12.9	20.7	42.7
Events in ATLAS on-Z SR($14\text{TeV}, 10\text{fb}^{-1}$)	172.8	168.2	103.6	18.2
Expected significance($14\text{TeV}, 10\text{fb}^{-1}$)	25	24	14	—
Expected significance($8\text{TeV}, 20.3\text{fb}^{-1}$)	2.4	4.0	2.3	—

Table 2. Detailed information of the benchmark points P_1 , P_2 and P_3 , including their contributions to the ATLAS excess and behaviors at the LHC-14.

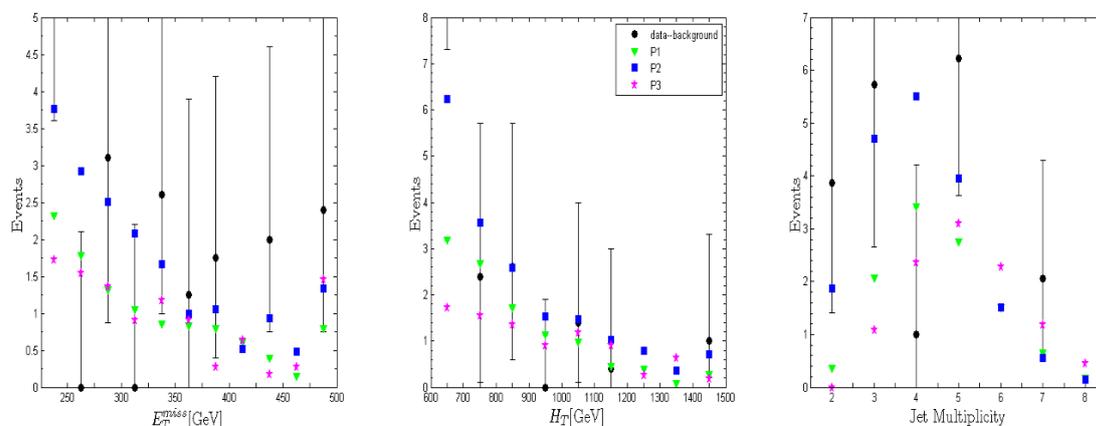


Figure 4. Distributions of E_T^{miss} , H_T and the jet multiplicity predicted by the three benchmark points. As a comparison, the experimental data for the ATLAS on-Z excess with the expected SM background extracted are also shown. Note that as did in [6], only the statistical uncertainties are included in the error bars of the data.

solid circle with the error bar stands for the data obtained by the ATLAS experiment with the expected SM background subtracted [3], and the predictions of the points P1, P2 and P3 are marked by triangles, squares and asterisks respectively. To quantize the difference between the theoretical predictions and the corresponding experimental data, we define a

χ^2 function for each distribution by a simple way

$$\chi^2 = \sum_i \frac{(s_i - \hat{s}_i)^2}{(\delta s_i)^2}, \tag{4.1}$$

where s_i is the theoretical prediction in the i th bin, \hat{s}_i is the corresponding experimental datum, and δs_i is the error of the datum. In table 2, we show the values of χ^2 for different distributions.

From figure 4 and table 2, one can learn that all points, especially point P2, can reproduce the distributions in an excellent way. The only significant difference between the squark explanation and the gluino explanation in generating the distributions comes from the jet multiplicity, i.e. for the former explanation, the n_j distribution peaks at 3, while for the latter explanation, it peaks at 4. Due to the large errors of the data at current time, we can not determine which explanation is preferred to account for the excess.

5 Test of our explanation at the 14 TeV LHC

Considering that the squark pair production rate at the LHC-14 can be greatly enhanced in comparison with that at the LHC-8, one may expect that the squark explanation will be tested very soon at the LHC. We investigate this issue by considering the lepton- Z + jets + E_T^{miss} signal of the production at the LHC-14. For simplicity, we assume the same cuts as those of the ATLAS on- Z search at the LHC-8, and estimate the SM background of the signal. In doing this, we suppose that the dominant background at the LHC-14 comes from the same processes as those at the LHC-8, which include flavor-symmetric backgrounds, Z + jets, rare top and diboson [3]. Since it is hard to get accurate background events by directly simulating the processes at the LHC-14, we simulate each background process at the LHC-14 and LHC-8 separately to get the ratio of their rates after cuts, then we scale the background at the LHC-8, which was given in the ATLAS report [3], by this ratio. We realize that the results obtained in this way may deviate significantly from their true values, but without any detailed information of the ATLAS detector at the LHC-14, our results may serve as a rough estimate of the background.

Once we know the signal and the total background after the cuts, we can calculate the expected significance by following formulae

$$S = \frac{N_s}{\sqrt{N_b + \epsilon^2 N_b^2}}, \tag{5.1}$$

where N_s and N_b denote the event numbers for the signal and the background respectively, and the coefficient ϵ parameterizes the effect induced by systematic errors. In our calculation, we set $\epsilon = 30\%$ for 14TeV LHC, which was adopted at the LHC-8 [3].

Assuming 10 fb^{-1} integrated luminosity at the LHC-14, we present in table 2 the predictions of the points on the event number of the signal as well as that of the background. This table indicates that the signals at the LHC-14 after cuts are enhanced by more than 9 times, and by contrast the background is enhance by only about 2 times. As a result,

either the squarks/gluino predicted by the points will be discovered, or in case of the non-observation of the leptonic- $Z + jets + E_T^{\text{miss}}$ signal, the points will be excluded. Moreover, comparing the squark explanation with the gluino explanation, we note that the former is more readily tested at 14TeV LHC.

6 Conclusion

In this paper, we intended to explain the 3σ excess recently reported by the ATLAS collaboration in the search for the leptonic- $Z + jets + E_T^{\text{miss}}$ signal. For this end, we considered the pair production of the first two generation squarks in the NMSSM with the decay chain $\tilde{q} \rightarrow q\tilde{\chi}_2^0 \rightarrow q\tilde{\chi}_1^0 Z$. In order to maximize the signal rate and also to simplify our analysis, we considered a singlino-dominated $\tilde{\chi}_1^0$ and a bino-dominated $\tilde{\chi}_2^0$, and assumed both $Br(\tilde{q} \rightarrow q\tilde{\chi}_2^0)$ and $Br(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 Z)$ at roughly 100%. With these assumptions, the parameters relevant to our analysis include the common squark mass $m_{\tilde{q}}$, the gluino mass $m_{\tilde{g}}$ as well as the mass splittings $\Delta m_1 \equiv m_{\tilde{q}} - m_{\tilde{\chi}_2^0}$ and $\Delta m_2 \equiv m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$.

In order to find the parameter space that can explain the excess, we fixed $m_{\tilde{g}} = 4.5\text{TeV}$ and $m_{\tilde{q}} = 1.5\text{TeV}$ separately, and for each case, we performed a grid scan over the parameters Δm_1 and Δm_2 by choosing a series of $m_{\tilde{q}}$. Then for each parameter point we encountered, we simulated the process $pp \rightarrow \tilde{q}\tilde{q} \rightarrow \tilde{\chi}_1^0 Z q\tilde{\chi}_1^0 Z q$ with the cuts adopted by the ATLAS on- Z search, the CMS on- Z search, and the ATLAS $2 \sim 6$ jets + E_T^{miss} searches respectively. Based on our simulations, we have following conclusions:

- After considering the constraints from the ATLAS searches for $jets + E_T^{\text{miss}}$ signals, the central value of the ATLAS Z -peaked excess can be obtained for $m_{\tilde{q}} \lesssim 1.2\text{TeV}$.
- If the constraint from the CMS on- Z search is further considered, more than 10 signal events are still attainable for $m_{\tilde{q}} \lesssim 750\text{GeV}$.
- For the squarks as heavy as about 1.4TeV, the squark pair production can still account for the excess at 2σ level without conflicting with any constraints.
- Compared with the explanation by gluino pair production, the squark explanation allows for a significantly wider range of $m_{\tilde{q}}$ as well as a less compressed SUSY mass spectrum.

Moreover, we also investigated whether the squark pair production can reproduce the distributions of the excess reported by the ATLAS collaboration. We found that, quite similar to the gluino pair production, the squark pair production can fit the data quite well. Probing the squark explanation at the 14TeV LHC was also investigated, and we concluded that only with 10fb^{-1} integrated luminosity, the squarks that are able to explain the excess will be either discovered or excluded.

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References

- [1] ATLAS collaboration, *Observation of a new particle in the search for the standard model Higgs boson with the ATLAS detector at the LHC*, *Phys. Lett. B* **716** (2013) 1 [[arXiv:1207.7214](https://arxiv.org/abs/1207.7214)] [[INSPIRE](#)].
- [2] CMS collaboration, *Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC*, *Phys. Lett. B* **716** (2012) 30 [[arXiv:1207.7235](https://arxiv.org/abs/1207.7235)] [[INSPIRE](#)].
- [3] ATLAS collaboration, *Search for supersymmetry in events containing a same-flavour opposite-sign dilepton pair, jets and large missing transverse momentum in $\sqrt{s} = 8$ TeV pp collisions with the ATLAS detector*, *Eur. Phys. J. C* **75** (2015) 318 [[arXiv:1503.03290](https://arxiv.org/abs/1503.03290)] [[INSPIRE](#)].
- [4] G. Barenboim, J. Bernabeu, V.A. Mitsou, E. Romero, E. Torro and O. Vives, *METing SUSY on the Z peak*, [arXiv:1503.04184](https://arxiv.org/abs/1503.04184) [[INSPIRE](#)].
- [5] N. Vignaroli, *Z-peaked excess from heavy gluon decays to vectorlike quarks*, *Phys. Rev. D* **91** (2015) 115009 [[arXiv:1504.01768](https://arxiv.org/abs/1504.01768)] [[INSPIRE](#)].
- [6] U. Ellwanger, *Possible explanation of excess events in the search for jets, missing transverse momentum and a Z boson in pp collisions*, *Eur. Phys. J. C* **75** (2015) 367 [[arXiv:1504.02244](https://arxiv.org/abs/1504.02244)] [[INSPIRE](#)].
- [7] B. Allanach, A. Raklev and A. Kvellestad, *Consistency of the recent ATLAS $Z + E_T^{\text{miss}}$ excess in a simplified GGM model*, *Phys. Rev. D* **91** (2015) 095016 [[arXiv:1504.02752](https://arxiv.org/abs/1504.02752)] [[INSPIRE](#)].
- [8] A. Kobakhidze, N. Liu, L. Wu and J.M. Yang, *ATLAS Z-peaked excess in the MSSM with a light sbottom or stop*, *Phys. Rev. D* **92** (2015) 075008 [[arXiv:1504.04390](https://arxiv.org/abs/1504.04390)] [[INSPIRE](#)].
- [9] J. Cao, L. Shang, J.M. Yang and Y. Zhang, *Explanation of the ATLAS Z-peaked excess in the NMSSM*, *JHEP* **06** (2015) 152 [[arXiv:1504.07869](https://arxiv.org/abs/1504.07869)] [[INSPIRE](#)].
- [10] B.A. Dobrescu, *Leptophobic boson signals with leptons, jets and missing energy*, [arXiv:1506.04435](https://arxiv.org/abs/1506.04435) [[INSPIRE](#)].
- [11] M. Cahill-Rowley, J.L. Hewett, A. Ismail and T.G. Rizzo, *The ATLAS $Z + MET$ excess in the MSSM*, [arXiv:1506.05799](https://arxiv.org/abs/1506.05799) [[INSPIRE](#)].
- [12] X. Lu, S. Shirai and T. Terada, *ATLAS Z excess in minimal supersymmetric standard model*, *JHEP* **09** (2015) 204 [[arXiv:1506.07161](https://arxiv.org/abs/1506.07161)] [[INSPIRE](#)].
- [13] S.P. Liew, A. Mariotti, K. Mawatari, K. Sakurai and M. Vereecken, *Z-peaked excess in goldstini scenarios*, *Phys. Lett. B* **750** (2015) 539 [[arXiv:1506.08803](https://arxiv.org/abs/1506.08803)] [[INSPIRE](#)].
- [14] H. Baer, V. Barger and D. Mickelson, *Direct and indirect detection of higgsino-like WIMPs: concluding the story of electroweak naturalness*, *Phys. Lett. B* **726** (2013) 330 [[arXiv:1303.3816](https://arxiv.org/abs/1303.3816)] [[INSPIRE](#)].

- [15] U. Ellwanger, C. Hugonie and A.M. Teixeira, *The next-to-minimal supersymmetric standard model*, *Phys. Rept.* **496** (2010) 1 [[arXiv:0910.1785](#)] [[INSPIRE](#)].
- [16] J. Cao, L. Shang, P. Wu, J.M. Yang and Y. Zhang, *Interpreting the galactic center gamma-ray excess in the NMSSM*, *JHEP* **10** (2015) 030 [[arXiv:1506.06471](#)] [[INSPIRE](#)].
- [17] F. Franke and H. Fraas, *Neutralinos and Higgs bosons in the next-to-minimal supersymmetric standard model*, *Int. J. Mod. Phys. A* **12** (1997) 479 [[hep-ph/9512366](#)] [[INSPIRE](#)].
- [18] W. Beenakker, R. Hopker and M. Spira, *PROSPINO: a program for the production of supersymmetric particles in next-to-leading order QCD*, [hep-ph/9611232](#) [[INSPIRE](#)].
- [19] ATLAS collaboration, *Search for squarks and gluinos with the ATLAS detector in final states with jets and missing transverse momentum using $\sqrt{s} = 8$ TeV proton-proton collision data*, *JHEP* **09** (2014) 176 [[arXiv:1405.7875](#)] [[INSPIRE](#)].
- [20] ATLAS collaboration, *Search for squarks and gluinos with the ATLAS detector in final states with jets and missing transverse momentum and 20.3 fb⁻¹ of $\sqrt{s} = 8$ TeV proton-proton collision data*, [ATLAS-CONF-2013-047](#) (2013).
- [21] CMS collaboration, *Search for physics beyond the standard model in events with two leptons, jets and missing transverse momentum in pp collisions at $\sqrt{s} = 8$ TeV*, *JHEP* **04** (2015) 124 [[arXiv:1502.06031](#)] [[INSPIRE](#)].
- [22] U. Ellwanger, *A Higgs boson near 125 GeV with enhanced di-photon signal in the NMSSM*, *JHEP* **03** (2012) 044 [[arXiv:1112.3548](#)] [[INSPIRE](#)].
- [23] J.-J. Cao, Z.-X. Heng, J.M. Yang, Y.-M. Zhang and J.-Y. Zhu, *A SM-like Higgs near 125 GeV in low energy SUSY: a comparative study for MSSM and NMSSM*, *JHEP* **03** (2012) 086 [[arXiv:1202.5821](#)] [[INSPIRE](#)].
- [24] S.F. King, M. Muhlleitner and R. Nevzorov, *NMSSM Higgs benchmarks near 125 GeV*, *Nucl. Phys. B* **860** (2012) 207 [[arXiv:1201.2671](#)] [[INSPIRE](#)].
- [25] J. Alwall, M. Herquet, F. Maltoni, O. Mattelaer and T. Stelzer, *MadGraph 5: going beyond*, *JHEP* **06** (2011) 128 [[arXiv:1106.0522](#)] [[INSPIRE](#)].
- [26] T. Sjöstrand, S. Mrenna and P.Z. Skands, *PYTHIA 6.4 physics and manual*, *JHEP* **05** (2006) 026 [[hep-ph/0603175](#)] [[INSPIRE](#)].
- [27] M. Drees, H. Dreiner, D. Schmeier, J. Tattersall and J.S. Kim, *CheckMATE: confronting your favourite new physics model with LHC data*, *Comput. Phys. Commun.* **187** (2014) 227 [[arXiv:1312.2591](#)] [[INSPIRE](#)].
- [28] J.S. Kim, D. Schmeier, J. Tattersall and K. Rolbiecki, *A framework to create customised LHC analyses within CheckMATE*, *Comput. Phys. Commun.* **196** (2015) 535 [[arXiv:1503.01123](#)] [[INSPIRE](#)].
- [29] DELPHES 3 collaboration, J. de Favereau et al., *DELPHES 3, a modular framework for fast simulation of a generic collider experiment*, *JHEP* **02** (2014) 057 [[arXiv:1307.6346](#)] [[INSPIRE](#)].
- [30] A.L. Read, *Presentation of search results: the CL(s) technique*, *J. Phys. G* **28** (2002) 2693 [[INSPIRE](#)].
- [31] R. Ding, Y. Fan, J. Li, T. Li and B. Zhu, *The ATLAS leptonic-Z excess from light squark productions in the NMSSM extension with a heavy Dirac gluino*, [arXiv:1508.07452](#) [[INSPIRE](#)].