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Biadjoint scalar tree amplitudes and intersecting dual associahedra

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ABSTRACT: We present a new formula for the biadjoint scalar tree amplitudes $m(\alpha|\beta)$ based on the combinatorics of dual associahedra. Our construction makes essential use of the cones in 'kinematic space' introduced by Arkani-Hamed, Bai, He, and Yan. We then consider dual associahedra in 'dual kinematic space.' If appropriately embedded, the intersections of these dual associahedra encode the amplitudes $m(\alpha|\beta)$. In fact, we encode all the partial amplitudes at *n*-points using a single object, a 'fan,' in dual kinematic space. Equivalently, as a corollary of our construction, all *n*-point partial amplitudes can be understood as coming from integrals over subvarieties in a toric variety. Explicit formulas for the amplitudes then follow by evaluating these integrals using the equivariant localisation formula. Finally, by introducing a lattice in kinematic space, we observe that our fan is also related to the inverse KLT kernel, sometimes denoted $m_{\alpha'}(\alpha|\beta)$.

KEYWORDS: Differential and Algebraic Geometry, Scattering Amplitudes, Gauge-gravity correspondence

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1 Introduction

Arkani-Hamed, Bai, He, and Yan (AHBHY) presented new formulas for the biadjoint scalar tree amplitudes in ref. [1]. Their formulas are based on constructions in 'kinematic space,' \mathcal{K}_n , which is the vector space of all Mandelstam variables $s_{ij} = 2k_i \cdot k_j$ subject to the momentum conservation relations. For example, at four points, the kinematic space \mathcal{K}_4 is the plane defined by

 $s_{12} + s_{13} + s_{23} = 0$

in \mathbb{R}^3 . AHBHY construct embeddings of the associahedron into kinematic space.¹ These are polytopes in \mathcal{K}_n of dimension n-3. For instance, at four points, the ordering $\alpha = 1234$ is associated to the line segment $\mathcal{A}(1234)$ between the two points

$$(s_{12}, s_{13}, s_{23}) = (c, -c, 0)$$
 and $(s_{12}, s_{13}, s_{23}) = (0, -c, c),$

for some constant c. This line segment is an embedding of the associahedron on three letters into kinematic space, regarded as a hyperplane in \mathbb{R}^3 . See figure 1 for an illustration of this. There are two other embeddings (also shown in figure 1). $\mathcal{A}(3124)$ is the line segment between

 $(s_{12}, s_{13}, s_{23}) = (c, 0, -c)$ and $(s_{12}, s_{13}, s_{23}) = (0, c, -c),$

and $\mathcal{A}(2314)$ is the segment between

$$(s_{12}, s_{13}, s_{23}) = (-c, 0, c)$$
 and $(s_{12}, s_{13}, s_{23}) = (-c, c, 0).$

Notice that the three line segments are not intersecting. In general, AHBHY's construction embeds (n-1)!/2 non-intersecting associahedra in \mathcal{K}_n . These associahedra can be realised as the intersection of hyperplanes with (n-1)!/2 cones that we call $\mathcal{C}(\alpha)$. For instance, at four points, there are three non-intersecting cones which we show in figure 3. An important observation for the present paper is that the 'dual cones' $\mathcal{C}(\alpha)^*$ in dual kinematic space do intersect. The purpose of this paper is to explain how the intersections of the dual cones are related to the biadjoint scalar tree amplitudes. For example, the cones $\mathcal{C}(1234)$ and $\mathcal{C}(2314)$ do not intersect in \mathcal{K}_4 . But their dual cones, $\mathcal{C}(1234)^*$ and $\mathcal{C}(2314)^*$, intersect in the span of a vector W_{23} . See figure 2. The vector W_{23} is 'dual' to the Mandelstam variable $-s_{23}$ in the sense that $W_{23} \cdot Z = -s_{23}$ for Z a vector in \mathcal{K}_4 . But $-s_{23}$ is the propagator for the partial amplitude

$$m(1234|2314) = \frac{1}{-s_{23}}.$$

So this partial amplitude can also be written as

$$m(1234|2314) = \frac{1}{W_{23} \cdot Z}.$$

This suggests that the amplitude m(1234|2314) is related to the intersection of the two dual cones $C(1234)^*$ and $C(2341)^*$. In this paper, we show how this observation can be made precise and generalised to all the partial tree amplitudes. Our main result is a generalisation of a formula given by AHBHY, which we now explain. Let α and β be orderings of $\{1, \ldots, n\}$. Then, following AHBHY, consider an associahedron $\mathcal{A}(\alpha)$ in kinematic space contained in some n-3 dimensional plane $H(\alpha) \subset \mathcal{K}_n$. The 'dual polytope' $\mathcal{A}(\alpha)^*$ is a polytope which lives in the dual vector space $H(\alpha)^*$. Duality on polytopes maps dimensionk faces in $\mathcal{A}(\alpha)$ to codimension-k faces in $\mathcal{A}(\alpha)^*$. AHBHY prove the following formula,

$$m(\alpha|\alpha) = \operatorname{Vol}(\mathcal{A}(\alpha)^*), \qquad (1.1)$$

¹This is carried out in section 3.2 of ref. [1].



Figure 1. The three associahedra $\mathcal{A}(\alpha)$ defined by AHBHY's construction at four points. The associahedra lie in a plane, kinematic space, inside \mathbb{R}^3 .

where Vol denotes the 'volume' of the polytope.² The 'volume' Vol is not a number, but rather a function on $H(\alpha)$. Or, if you prefer, the volume of the dual polytope depends on a choice of position (i.e. a choice of Mandelstam variables) $Z \in H(\alpha)$. The 'volume' Vol is similar to an ordinary volume in the sense that it satisfies

$$\operatorname{Vol}(A \cup B) = \operatorname{Vol}(A) + \operatorname{Vol}(B) - \operatorname{Vol}(A \cap B),$$

for any two polytopes A and B. We call anything which satisfies this relation a 'valuation' on polytopes. Our main result can be stated as (see corollary 4.2)

$$m(\alpha|\beta) = \operatorname{Val}(\partial \mathcal{A}(\alpha)^* \cap \partial \mathcal{A}(\beta)^*).$$
(1.2)

In this formula, $\partial \mathcal{A}(\alpha)^*$ is the boundary of $\mathcal{A}(\alpha)^*$ and 'Val' is a natural valuation on polytopes which we will introduce later. The novel aspect of this formula is that we must first define embeddings of the dual associahedra $\mathcal{A}(\alpha)^*$ into \mathcal{K}_n^* (just as AHBHY define embeddings of $\mathcal{A}(\alpha)$ into \mathcal{K}_n). We define a 'canonical embedding' in section 3. Given this canonical embedding, it turns out that

$$\operatorname{Val}(\partial \mathcal{A}(\alpha)^*) = \operatorname{Vol}(\mathcal{A}(\alpha)^*).$$

where, on the left-hand-side, we regard $\partial \mathcal{A}(\alpha)^*$ as embedded in \mathcal{K}_n^* and, on the right-handside, $\mathcal{A}(\alpha)^*$ is embedded only in $H(\alpha)^*$ (or some projective compactification of it, as done in AHBHY). For this reason, equation (1.2) is a mild generalisation of AHBHY's result, equation (1.1). Before saying more about our results, we will recall some of the reasons to be interested in the biadjoint scalar amplitudes and their various presentations.

The biadjoint scalar tree amplitudes $m(\alpha|\beta)$ can be computed in a Feynman diagram expansion by summing over all tri-valent graphs which are planar with respect to both α

²This formula is proposed in section 5.1 and explicitly proved in section 5.2 of ref. [1], though it also follows from a general theorem concerning 'positive geometries.'



Figure 2. The cones C(1234) and C(2314) do not intersect in kinematic space, \mathcal{K}_4 . However, the dual cones $C(1234)^*$ and $C(2314)^*$ do intersect, and their intersection can be regarded as encoding the amplitudes m(1234|2314).

and β . However, there are many other formulas for these amplitudes. A particularly novel formula for $m(\alpha|\beta)$ was given by Cachazo-He-Yuan, who expressed the amplitude as a residue pairing [2, 3]. In [4], Mizera pointed out that the CHY residue pairing is essentially the same thing as the intersection pairing of two cocycles in a certain cohomology theory.³ An attractive formula for these intersection pairings was computed by Matsumoto in [5]. The work by AHBHY is more combinatorial. In [1], the amplitudes $m(\alpha|\beta)$ are presented in terms of associahedra in kinematic space.⁴ One way to infer the amplitude from the associahedra is by studying 'canonical forms' that have logarithmic singularities on the boundaries of the associahedra.⁵ The biadjoint scalar tree amplitudes have aroused all of this interest because of their importance in understanding the relationship between tree-level gravity and Yang-Mills amplitudes. The partial amplitudes $m(\alpha|\beta)$ play an essential role in the 'double copy' relation (equation (6.2)) between gravity tree amplitudes and Yang-Mills partial tree amplitudes [8]. Indeed, Yang-Mills and gravity amplitudes can essentially be expanded as a sum over $m(\alpha|\beta)$ with certain numerators. Another source of interest is the ubiquitous presence of the biadjoint tree amplitudes in string theory. For instance, closed string tree amplitudes can be expanded in a basis of I-

³The appropriate cohomology theory is cohomology with values in a local system. The underlying space is the open string moduli space $\mathcal{M}_{0,n}(\mathbb{R})$ and the local system is defined by the monodromies of the Koba-Nielsen factor. The cocycles in the pairing are represented in (twisted) de Rham theory by the Parke-Taylor factors associated to the two orderings of $\{1, \ldots, n\}$. These are top-forms on the open string moduli space $\mathcal{M}_{0,n}(\mathbb{R})$.

⁴This realisation of the associahedron bears some resemblance to methods developed in the combinatorial literature, see [6] for a review. The exact prescription given by AHBHY does, however, appear to be new.

⁵This fits into a general paradigm. The associahedra are defined by finitely many linear inequalities and equalities. For this reason, they are 'linear semi-algebraic sets' or 'positive geometries.' A volume form with logarithmic singularities on the boundary of a positive geometry is called a 'canonical form' and a great deal can be said in general about these forms. See ref. [7] for a recent elaboration of these ideas.

integrals' $I(\alpha|\beta)$ whose leading term in the infinite tension limit $(\alpha' \to 0)$ is $m(\alpha|\beta)$.⁶ [9] Likewise, open string tree amplitudes can be expanded in a basis of 'Z-integrals' $Z_{\alpha}(\beta)$ whose leading term in the infinite tension limit is $m(\alpha|\beta)$.⁷ [10] The open and closed string tree amplitudes are related by the Kawai-Lewellen-Tai relations [11]. Following Mizera [12], we denote the inverse KLT kernel by $m_{\alpha'}(\alpha|\beta)$ since, in the infinite tension limit, the leading term in $m_{\alpha'}(\alpha|\beta)$ is given (up to a possible factor) by $m(\alpha|\beta)$. This means that we might regard the inverse KLT kernel $m_{\alpha'}(\alpha|\beta)$ as a the amplitude $m(\alpha|\beta)$ with α' corrections. Perhaps surprisingly, $m_{\alpha'}(\alpha|\beta)$ has a tidy description in terms of the open string moduli space $\mathcal{M}_{0,n}(\mathbb{R})$ and the monodromy factors picked up by the Koba-Nielsen factor when points collide. The compactified moduli space $\overline{\mathcal{M}_{0,n}}(\mathbb{R})$ is tesselated by (n-1)!/2 copies of the associahedron on n-1 letters. (See [13], theorem 3.1.3, for instance.) Mizera proved that the functions $m_{\alpha'}(\alpha|\beta)$ can be obtained as the intersections of these associahedra regarded as twisted cycles in $\mathcal{M}_{0,n}(\mathbb{R})$ [14]. Some of the resulting formulas for $m_{\alpha'}(\alpha|\beta)$ were encountered earlier in the mathematical literature on twisted de Rham theory (see [15], in particular). Mizera's result gives a new interpretation of the KLT relation as a purely cohomological statement, namely it is the analogue in twisted de Rham theory (or cohomology with values in a local system) of Riemann's period relations. This interpretation is presented and discussed in [14].

Our main result is a new formula for the amplitudes $m(\alpha|\beta)$. All *n*-point partial amplitudes are described by a single object, called a 'fan,' in dual kinematic space \mathcal{K}_n^* . Individual partial amplitudes $m(\alpha|\beta)$ are given by the 'volumes' of the intersections of dual associahedra in this fan. This result is stated precisely in section 4 as theorem 4.1 and corollary 4.2. Once the necessary definitions are given, the proof of this result is essentially tautological. After all, associahedra and dual associahedra are just a fancy way of encoding the Feynman diagram sum of trivalent graphs. Given this, why do we need another formula for $m(\alpha|\beta)$? Especially since so many formulas already exist? We think our formula has two intrinsic merits. First, it is a compact and symmetric way to present the relevant combinatorics. Second, it leads us to study a certain object (the fan) which encodes all of the partial amplitudes at once. In fact, as we point out at the end of section 6, our main result, restated in the language of toric geometry, is that all the *n*-point partial amplitudes $m(\alpha|\beta)$ are described by a single toric variety. Particular partial amplitudes $m(\alpha|\beta)$ correspond to toric subvarieties (or, really, cycles in the homology). This appears to us as a novel description of the amplitudes. We can only hope that these new geometric presentations of the amplitudes will suggest new approaches to the Yang-Mills and gravity amplitudes and to the double copy relations between them. We offer speculations to this end during the discussion in section 6. We also relegate several technical (but quite interesting) loose ends to this discussion. Before proving our results in section 4, our key 'canonical embedding' construction is described in 3. We explain all the essential definitions in section 2, where we also review the relevant results that we need from AHBHY (ref. [1]). Our second result is an observation concerning the inverse KLT kernel. The diagonal components of

⁶See equation (3.24) of [9] for the integral and equations (4.2), (4.3) and (4.4) for the expansion of the closed string amplitude.

⁷See equation (2.2) of [10] for the integral and (2.17) for the expansion of the open string amplitude.



Figure 3. The three cones $C(\alpha)$ in the kinematic space \mathcal{K}_4 for four points.

the inverse KLT kernel — that is, $m_{\alpha'}(\alpha|\alpha)$ — are also related to the fan in dual kinematic space which encodes the partial amplitudes $m(\alpha|\beta)$. To obtain the components $m_{\alpha'}(\alpha|\alpha)$ we have to introduce a lattice on kinematic space. This lattice amounts to imposing that the Mandelstam variables s_{ij} are integer multiples of $1/\alpha'$. (It seems natural to introduce a lattice here since, as we have just remarked, toric geometry is implicit throughout this paper and the construction of toric varieties is predicated on lattices.) Having introduced the lattice, we observe in section 5.3 that the components $m_{\alpha'}(\alpha|\alpha)$ can be given as a certain sum over lattice points. This curious, and potentially superficial observation, sparks a number of further speculations during the discussion in section 6. The main text of the paper is largely self contained. However, for the reader's convenience, several definitions and theorems mentioned briefly in the text are reviewed at greater length in appendix A, where we also give references to the many expositions that already exist.

2 Review and definitions

In this section we describe AHBHY's construction of associahedra in 'kinematic space.' At n points, their construction gives (n-1)!/2 distinct associahedra. This is reminiscent of the open string moduli space, $\mathcal{M}_{0,n}(\mathbb{R})$, which (after blow-ups) is tiled by (n-1)!/2associahedra. Unlike the tiles of $\mathcal{M}_{0,n}(\mathbb{R})$, the associahedra that AHBHY construct in \mathcal{K}_n do not intersect. However, we will see in section 3 that it is natural to regard the dual associahedra as intersecting in dual kinematic space, \mathcal{K}_n^* .

2.1 Cones in kinematic space

To begin, let us define 'kinematic space' at n points, \mathcal{K}_n , to be the space of all Mandelstam variables subject to momentum conservation.⁸ We can present \mathcal{K}_n explicitly as follows. Let $\{s_{ij}\}$, where $1 \leq i, j \leq n$, be the n(n-1)/2 Mandelstam variables. We impose the n

⁸See also section 2 of [1].

momentum conservation relations

$$\sum_{j=1}^{n} s_{ij} = 0.$$

We can use these relations to remove one index, say n, from our formulas. That is, we consider only those s_{ij} with $1 \le i, j \le n-1$. There are (n-1)(n-2)/2 such variables and we regard them as coordinates on a vector space

$$\mathcal{V}_n = \mathbb{R}^{\frac{n(n-3)}{2}+1}.$$

The remaining momentum conservation relation on these variables is

$$\sum_{i,j=1}^{n-1} s_{ij} = 0,$$

which presents the kinematic space \mathcal{K}_n as a hyperplane in \mathcal{V}_n . For example, at four points, the vector space \mathcal{V}_4 is \mathbb{R}^3 with coordinates (s_{12}, s_{23}, s_{13}) and the kinematic space \mathcal{K}_4 is the hyperplane

$$s_{12} + s_{23} + s_{13} = 0.$$

Now let α be an ordering of $\{1, \ldots, n\}$ such that $\alpha(n) = n$. Associated to each such ordering, AHBHY associate a cone in \mathcal{K}_n .⁹ The cone associated to α is cut out by the inequalities

$$X_S = -\frac{1}{2} \sum_{i,j \in S} s_{ij} \ge 0$$
 (2.1)

for all subsets $S \subset \{1, \ldots, n-1\}$ which are 'consecutive' with respect to α .¹⁰ There are

$$\binom{n}{2} - (n-1) = \frac{n(n-3)}{2} + 1$$

such subsets S, corresponding to the diagonals of the n-gon. However, $X_{123...n-1}$ is identically zero on the hyperplane \mathcal{K}_n . So only

$$\frac{n(n-3)}{2}$$

inequalities are implied by equation (2.1). We write

 $\mathcal{C}(\alpha) = \{(s_{ij}) \in \mathcal{K}_n | \text{ satisfying the inequalities, equation } (2.1)\}$

for this cone. The cone $C(\alpha)$ is 'polyhedral' in the sense that it has finitely many flat sides.¹¹ Moreover, the inequalities that define $C(\alpha)$ are linearly independent, which means that $C(\alpha)$ contains an interior region of full dimension, dim $\mathcal{K}_n = n(n-3)/2$. There are

⁹This is done in section 3.2 of [1].

¹⁰By 'consecutive' I mean that S can be written as $\{\alpha(a), \alpha(a+1), \ldots, \alpha(a+k)\}$ for some positive integers a and k. For instance $\{1, 2, 3\}$ is consecutive with respect to $\alpha = (1324)$, but $\{1, 2\}$ is not.

¹¹This is in contrast to a cone generated by a sphere, or any arbitrary convex set. See appendix A.1 for a review of definitions and results concerning polyhedral cones.

(n-1)! permutations α such that $\alpha(n) = n$. However, the inequalities, equation (2.1), do not define (n-1)! distinct cones $C(\alpha)$. Given a permutation α such that $\alpha(n) = n$, define the reversed permutation $\bar{\alpha}$ as

$$(\bar{\alpha}(1),\bar{\alpha}(2),\ldots,\bar{\alpha}(n-1),\bar{\alpha}(n))=(\alpha(n-1),\alpha(n-2),\ldots,\alpha(1),\alpha(n)).$$

The inequalities, equation (2.1), are identical for α and its reverse $\bar{\alpha}$. That is, $C(\alpha) = C(\bar{\alpha})$. For this reason, we obtain only (n-1)!/2 distinct cones. In the four point example, we obtain 3!/2 = 3 cones. These are

$$\mathcal{C}(1234) = \mathcal{C}(3214) = \{(s_{ij}) \mid -s_{12}, -s_{23} \ge 0\},\$$

$$\mathcal{C}(3124) = \mathcal{C}(2134) = \{(s_{ij}) \mid -s_{12}, -s_{13} \ge 0\},\$$

$$\mathcal{C}(2314) = \mathcal{C}(1324) = \{(s_{ij}) \mid -s_{13}, -s_{23} \ge 0\}.$$

We show these cones in figure 3. To do this, we employ a non-orthogonal basis

$$e_{12} = \begin{bmatrix} 0\\1 \end{bmatrix}$$
, and $e_{23} = \begin{bmatrix} \sqrt{3}/2\\-1/2 \end{bmatrix}$ (2.2)

for the plane, so that the point $(-s_{12}, -s_{23})$ in \mathcal{K}_4 is represented in the plane by $-s_{12}e_{12} - s_{23}e_{23}$. We choose this basis to make the symmetries of the construction more apparent in our drawings. Before we introduce associahedra in the following subsection, let us briefly comment on the physical significance of what we have done so far. Recall that the Mandelstam variables are $s_{ij} = 2k_i \cdot k_j$, where k_i and k_j are massless momentum vectors. The sign of s_{ij} then encodes whether k_i and k_j are time directed in the same way (i.e. both forward or both past directed) or in opposite ways. To see this, write $k_i = \omega_i(\epsilon_i, n_i)$ for a sign $\epsilon_i = \pm 1$, a positive energy ω_i , and a unit vector n_i . In Lorentzian signature we have

$$k_i \cdot k_j = \omega_i \omega_j (\epsilon_i \epsilon_j - \cos \theta),$$

where θ is the angle between n_i and n_j . So $s_{ij} \ge 0$ iff $\epsilon_i \epsilon_j = 1$ and $s_{ij} \le 0$ iff $\epsilon_i \epsilon_j = -1$. For example, at four points, both s_{12} and s_{23} are negative inside the cone C(1234). Momentum conservation implies that s_{13} is positive inside the cone. This means that k_1 and k_3 are both future-pointing (or both past-pointing) while k_2 is past-pointing (resp. future-pointing). So we see that some of the inequalities in equation (2.1) have a physical interpretation in terms of positive and negative energies.

2.2 Associahedra in kinematic space

AHBHY construct (n-1)!/2 associahedra in kinematic space \mathcal{K}_n by intersecting each of the cones $\mathcal{C}(\alpha)$ with an appropriate hyperplane.¹² For each ordering α , with $\alpha(n) = n$, define the hyperplane

 $H(\alpha) = \left\{ s_{\alpha(i),\alpha(j)} = \text{constant} \mid \text{for all non-adjacent pairs } i, j \text{ less than } n \right\}.$

¹²See again section 3.2 of [1].



Figure 4. The three realisations $\mathcal{A}(\alpha)$ of the associahedron in the kinematic space \mathcal{K}_4 for four points.

There are (n-2)(n-3)/2 pairs of non-adjacent (i, j) for i and j less than n. So the dimension of $H(\alpha)$ is

$$\frac{n(n-3)}{2} - \frac{(n-2)(n-3)}{2} = n - 3.$$

The intersection

$$\mathcal{A}(\alpha) = \mathcal{C}(\alpha) \cap H(\alpha) \subset \mathcal{K}_n$$

defines a polytope. For generic choices of the constants that define $H(\alpha)$, $\mathcal{A}(\alpha)$ is a realisation of the associahedron. The polytope $\mathcal{A}(\alpha)$ lies in the n-3 dimensional plane $H(\alpha)$. Nevertheless, the facets of $\mathcal{A}(\alpha)$ (these are the top dimension faces of $\mathcal{A}(\alpha)$) generate the cone $\mathcal{C}(\alpha)$, which bounds an interior region of dimension n(n-3)/2. Since all faces of $\mathcal{A}(\alpha)$ can be obtained as the intersections of facets, it follows that all faces of $\mathcal{A}(\alpha)$ are contained in the boundary of the cone $\mathcal{C}(\alpha)$. For example, at four points, consider the cone

$$\mathcal{C}(1234) = \{(s_{ij}) \mid -s_{12}, -s_{23} \ge 0\}.$$

We choose the hyperplane

$$H(1234) = \{(s_{ij}) \mid s_{13} = \sqrt{3/2}\}.$$

Then, using the basis introduced in equation (2.2), the associahedron is given by the convex hull

$$\mathcal{A}(1234) = \operatorname{Conv}\left(\begin{bmatrix}1\\0\end{bmatrix}, \begin{bmatrix}1/2\\\sqrt{3}/2\end{bmatrix}\right).$$

This is shown in figure 4, together with the associahedra $\mathcal{A}(3124)$ and $\mathcal{A}(2314)$.

2.3 Dual cones and associahedra

Let \mathcal{K}_n^* be the dual (as a vector space) of \mathcal{K}_n . There is a standard notion of a 'dual cone.' Define

$$\mathcal{C}(\alpha)^* = \{ W \in \mathcal{K}_n^* \,|\, W \cdot Z \ge 0 \text{ for all } Z \in \mathcal{C}(\alpha) \} \,. \tag{2.3}$$



Figure 5. The three dual cones $\mathcal{C}(\alpha)$ in \mathcal{K}_4^* .

The inequalities that define $\mathcal{C}(\alpha)$, equation (2.1), can be re-expressed as

 $W_S \cdot Z \ge 0,$

where the covectors $W_S \in \mathcal{K}_n^*$ are defined by

$$W_S \cdot Z = X_S. \tag{2.4}$$

Given this, the cone $\mathcal{C}(\alpha)^*$ is given by the 'conic hull' or 'positive span'¹³

 $\mathcal{C}(\alpha)^* = \operatorname{Cone} \{ W_S \mid \text{for } S \text{ consecutive with respect to } \alpha \}.$

In the four point example we have, for instance,

$$\mathcal{C}(1234)^* = \operatorname{Cone}(W_{12}, W_{23}).$$

In fact, the three dual cones $\mathcal{C}(1234)^*$, $\mathcal{C}(3124)^*$, and $\mathcal{C}(2314)^*$ fill the dual space \mathcal{K}_4^* , as shown in figure 5.¹⁴ What about dual associahedra? Unlike the cones, there is no natural notion of a dual associahedron in \mathcal{K}_n^* . Instead, we must restrict to the (n-3)-hyperplane $H(\alpha)$ that contains $\mathcal{A}(\alpha)$. Then the 'dual polytope' is given by

$$\mathcal{A}(\alpha)^* = \{ Y \in H(\alpha)^* \mid Y \cdot Z \ge -1 \text{ for all } Z \in \mathcal{A}(\alpha) \} \subset H(\alpha)^*.$$

The duality operation swaps dimension-k faces in $\mathcal{A}(\alpha)$ for codimension-(k + 1) faces in $\mathcal{A}(\alpha)^*$. Thus, vertices in $\mathcal{A}(\alpha)$ become facets in $\mathcal{A}(\alpha)^*$, and so on. To give a concrete example, let's return to the four point case and consider the line H(1234). Using $-s_{12}$ as a coordinate on this line, the associahedron is the interval

$$\mathcal{A}(1234) = \left[0, \sqrt{3}/2\right] \subset H(1234).$$

 $^{^{13}}$ See equation (A.1) in appendix A.1 for conic hulls.

¹⁴This is a consequence of the fact that the three cones C(1234), C(3124), and C(2314) in \mathcal{K}_4 are the tangent cones of an isosceles triangle (translated so that their apexes are all at the origin). In general, the sum of dual tangent cones of a convex polytope fills the dual vector space. See theorem A.3.



Figure 6. The dual polytope construction for a 1-simplex associahedron encountered in the four point example. Notice that duality for a polytope P in \mathcal{K}_n is not defined with respect to \mathcal{K}_n , but rather with respect to the smallest linear subspace containing P. This means that the dual associahedron $\mathcal{A}(1234)^*$ is not naturally regarded as living in \mathcal{K}_n^* until an embedding is chosen. We make a choice of embedding in section 3.

The dual polytope is then

$$\mathcal{A}(1234)^* = \left[-2/\sqrt{3},\infty\right) \subset H(1234)^*.$$

If we compactify at infinity, we can regard H(1234) as \mathbb{RP}^1 . Then the duality operation has sent a 1-simplex $\mathcal{A}(1234)$ to a dual 1-simplex $\mathcal{A}(1234)^*$. We illustrate this process in figure 6 so as to emphasise that the dual associahedron defined this way is not yet embedded into dual kinematic space. Such an embedding is an additional construction which inovlves making choices. This is analogous to the choices that are involved in AHBHY's embeddings of the associahedron into kinematic space. We choose a particular embedding in the next section.

3 Embedding dual associahedra

AHBHY construct (n-1)!/2 embeddings of the associahedron, which has dimension n-3, into a higher dimensional vector space, \mathcal{K}_n . We will now construct (n-1)!/2 embeddings of the dual associahedra $\mathcal{A}(\alpha)^*$ (or, in fact, their faces) into dual kinematic space \mathcal{K}_n^* . The idea is to embed the (faces of the) dual associahedron $\mathcal{A}(\alpha)^*$ into the (boundary of the) dual cone $\mathcal{C}(\alpha)^*$. This can be done in a canonical way such that the embedded associahedra $\mathcal{A}(\alpha)^* \subset \mathcal{K}_n^*$ have combinatorially meaningful intersections. In defining our embedding we will make use of the vectors

$$W_S \in \mathcal{K}_n^*, \qquad S \subset \{1, \dots, n-1\},$$

that we defined in equation (2.4). These vectors can be used as generators for the cones $C(\alpha)$. Now recall that there is a 1-1 correspondence between codimension k faces of

 $\mathcal{A}(\alpha) \subset H(\alpha)$ and dimension k-1 faces of $\mathcal{A}(\alpha)^* \subset H(\alpha)^*$. In particular, the facets (or codimension-1 faces) of $\mathcal{A}(\alpha) \subset H(\alpha)$ correspond to the vertices of $\mathcal{A}(\alpha)^* \subset H(\alpha)^*$. A facet of $\mathcal{A}(\alpha) \subset H(\alpha)$ is defined by a single inequality,

$$W_S \cdot Z \ge 0,$$

for some subset S. Let $Y_S \in H(\alpha)^*$ be the vertex of $\mathcal{A}(\alpha)^* \subset H(\alpha)^*$ corresponding to the facet $W_S \cdot Z = 0$ of $\mathcal{A}(\alpha)$. Then any dimension-k face of $\mathcal{A}(\alpha)^* \subset H(\alpha)^*$ is given by the convex hull

$$\operatorname{Conv}(Y_{S_1},\ldots,Y_{S_k}) \subset H(\alpha)^*$$

for some subsets S_i labelling the vertices of $\mathcal{A}(\alpha)^*$. We map this face to the convex hull

$$\operatorname{Conv}(W_{S_1},\ldots,W_{S_k}) \subset \mathcal{C}(\alpha)^*$$

in \mathcal{K}_n^* . In particular, we map the vertex Y_S to the vector W_S . In this way, we can embed all the faces of $\mathcal{A}(\alpha)^*$ into \mathcal{K}_n^* . In other words, we have an embedding of the boundary $\partial \mathcal{A}(\alpha)^*$ into \mathcal{K}_n^* . We call this the 'canonical embedding' of the faces and will typically abuse notation by denoting the embedded faces as $\partial \mathcal{A}(\alpha)^*$. In general, the canonically embedded faces $\partial \mathcal{A}(\alpha)^*$ do not bound a dimension n-3 polytope.¹⁵ This does happen for four points, but not at higher points. One can obtain an embedding of all $\mathcal{A}(\alpha)^*$ into \mathcal{K}_n^* by choosing a triangulation of $\mathcal{A}(\alpha)^*$ into (n-3)-simplices. However, we will not do this. The reason is that all the combinatorial data that we need for amplitudes is already contained in the faces of the dual associahedra. Finally, let us denote the union of all the canonically embedded faces by

$$FN_n = \bigcup_{\alpha} \partial \mathcal{A}(\alpha)^* \subset \mathcal{K}_n^*$$

(where we write 'FN' for 'face net'). We now give two examples. Consider first the four point example. The three dual vectors are

$$W_{12}, W_{13}, W_{23},$$

and the dual cones are given by

$$\mathcal{C}(1234)^* = \text{Cone}(W_{12}, W_{23}),$$

$$\mathcal{C}(2314)^* = \text{Cone}(W_{13}, W_{23}),$$

$$\mathcal{C}(3124)^* = \text{Cone}(W_{12}, W_{13}).$$

¹⁵It is possible to realize the whole dual associahedron $\mathcal{A}(\alpha)^*$ in dual kinematic space as follows. If $\mathcal{A}(\alpha)$ is the intersection of the cone $\mathcal{C}(\alpha)$ and the hyperplane $H(\alpha)$, then $\mathcal{A}(\alpha)^*$ can be regarded as the projection of $\mathcal{C}(\alpha)^*$ through the dual hyperplane $H(\alpha)^*$. (This is because, in projective geometry, restriction to hyperplanes is dual to projection through hyperplanes.) Any representative of this $\mathcal{A}(\alpha)^*$ in the preimage of the projection is then an embedding of $\mathcal{A}(\alpha)^*$ into \mathcal{K}_n^* . We do not use this here because we have not found a prescription which guarantees that these $\mathcal{A}(\alpha)^*$ intersect in the desired way. We thank Nima Arkani-Hamed for this comment.



Figure 7. The three dual associahedra $\mathcal{A}(\alpha)^*$ canonically embedded in \mathcal{K}_4^* form a triangle.

Then the canonically embedded dual associahedra are

$$\mathcal{A}(1234)^* = \text{Conv}(W_{12}, W_{23}), \tag{3.1}$$

$$\mathcal{A}(2314)^* = \text{Conv}(W_{13}, W_{23}), \tag{3.2}$$

$$\mathcal{A}(3124)^* = \text{Conv}(W_{12}, W_{13}). \tag{3.3}$$

That is, the dual associahedra tile an isosceles triangle in \mathcal{K}_4^* . See figure 7. So FN_4 gives the vertices of a triangle. The situation at five points is less straightforward. In this case, there are (5-1)!/2 = 12 associahedra, each being a pentagon. The dual of a pentagon is a pentagon. So FN_5 will be formed by embedding the boundaries of 12 pentagons into \mathcal{K}_5^* . The vertices of the pentagons will each be mapped under the canonical embedding to one of the following ten vectors,

$$\{W_{12}, W_{13}, W_{14}, W_{23}, W_{24}, W_{34}, W_{234}, W_{134}, W_{124}, W_{123}\}$$

Though we cannot easily sketch the result of this embedding, we can draw the 'net' associated to it, which is shown in figure 8. Vertices in the net with the same label are identified in \mathcal{K}_5^* under the embedding. Figure 8 can be arrived at from the analogous diagram showing $\overline{\mathcal{M}_{0,5}}(\mathbb{R})$ tiled by pentagons. This procedure is described in figure 9. Examining figure 8, we see that each vertex is contained in the boundary of 6 distinct faces. Moreover, there are 12 faces altogether and 10 vertices. This suggests that the embedded associahedra in \mathcal{K}_5^* tile a 'halved' or 'degenerate' dodecahedron, since a dodecahedron has 20 vertices, each of which meets 3 of its 12 faces. In general, FN_n contains

$$\sum_{k=2}^{n-2} \binom{n-1}{k} = 2^{n-1} - n - 1$$

vertices in \mathcal{K}_n^* . It is not clear whether FN_n (or perhaps some double cover of it) can be mapped to the face lattice of a convex polytope. See section 6 for a discussion of this and related problems.

4 The partial amplitudes $m(\alpha|\beta)$

In the previous section we constructed an embedding of the (faces of the) dual associahedra $\mathcal{A}(\alpha)^*$ into dual kinematic space, \mathcal{K}_n^* . Given this embedding, our main claim is that the



Figure 8. The net which corresponds to the embedded associahedra in \mathcal{K}_5^* . Vertices with the same label are identified in the embedding.



Figure 9. A portion of $\overline{\mathcal{M}_{0,5}}(\mathbb{R})$. By associating to every side of each pentagon a vertex of the dual pentagon, we arrive at a net of dual pentagons with shared vertices. Rearranging this net leads us to figure 8.

biadjoint scalar tree amplitudes are given by

$$m(\alpha|\beta) = \operatorname{Val}(\partial \mathcal{A}(\alpha)^* \cap \partial \mathcal{A}(\beta)^*),$$

where Val is a certain 'valuation' or 'volume' that we define below in equations (4.6) and (4.4). We prove this formula as theorem 4.1 at the end of this section. Before we define Val, we need to review a few preliminaries.¹⁶

4.1 Preliminaries

Let P be a polyhedron in \mathcal{K}_n^* . For instance, P could be a cone like $\mathcal{C}(\alpha)^*$. Then consider the integral

$$I_P(Z, \alpha') = \int_P e^{-\alpha' W \cdot Z} \mathrm{d}W_P, \qquad (4.1)$$

where dW_P is the Euclidean volume form of the appropriate dimension to be integrated over P. We have introduced a parameter, α' , which, at this stage, is not directly related to string theory in any way. Notice that W is dimensionful and so α' is dimensionful, too. The integral I_P is related to the volume of P. If P is a bounded polytope, then the $(\alpha')^0$ term in the Laurent expansion of $I_P(Z, \alpha')$,

$$I_P(Z,\alpha') = \dots + I_P^{(-1)}(Z)(\alpha')^{-1} + I_P^{(0)}(Z) + I_P^{(1)}(Z)\alpha' + \dots,$$

will coincide with the volume of P. For now, let us fix $\alpha' = 1$. Suppose P is a cone. Then, referring to equation (4.1), we see that the integral $I_P(V, \alpha')$ is well defined if $V \cdot W > 0$ for all $W \in P$. This means that $I_P(Z, 1)$ is well defined if Z is in the dual cone $P^* \subset \mathcal{K}_n$.¹⁷ If Z lies on the boundary of P^* , then $W \cdot Z = 0$ for W in some face of P. So $I_P(Z, 1)$ diverges to $+\infty$ as Z tends to the boundary of P^* and the integral is not defined for Zoutside P^* . Now consider any polyhedral cone

$$C = \operatorname{Cone}(W_1, \ldots, W_k),$$

where $k \leq \dim \mathcal{K}_n^*$. A standard calculation shows that for Z in the interior of C^* , the integral I_C is¹⁸

$$I_C(Z,1) = \frac{\langle W_1, \dots, W_k \rangle}{\prod_{i=1}^k (W_i \cdot Z)},$$
(4.2)

where $\langle W_1, \ldots, W_k \rangle$ is the Euclidean volume of the unit (open-closed) box

Box
$$(W_1, \ldots, W_k) = \left\{ \sum_{i=1}^k c_i W_i \, | \, 0 \le c_i < 1 \right\}.$$

¹⁷The dual cone was defined in equation (2.3) in section 2.3. See also equation (A.3) in appendix A.1.

$$\int_{0}^{\infty} \mathrm{d}x \, e^{-ax} = \frac{1}{a}$$

several times.

¹⁶See also appendix A for a more detailed review and references.

¹⁸See also theorem A.7 in appendix A.2 where I_C is computed for $k > \dim V$. The basic idea is to use the result that

We are not interested in evaluating the integral I for the cones $\mathcal{C}(\alpha)^* \subset \mathcal{K}_n^*$ introduced in section 2.3. For the purposes of writing amplitudes, the factors of $W_i \cdot Z$ appearing in equation (4.2) are propagators. The cone $\mathcal{C}(\alpha)^*$ has n(n-3)/2 vertices, and so $I_{\mathcal{C}(\alpha)^*}$ would contain too many propagators to be an *n*-point amplitude. An *n*-point amplitude should have n-3 propagators, since this is the number of internal lines in any *n*-point trivalent graph. All the combinatorial structure we need to write amplitudes is in the boundary $\partial \mathcal{C}(\alpha)^*$. This boundary is generated by the facets of the dual associahedron $\mathcal{A}(\alpha)^*$ canonically embedded in \mathcal{K}_n^* . Under the canonical embedding, $\mathcal{A}(\alpha)^*$ is not necessarily itself a convex polytope in \mathcal{K}_n^* (except for the n = 4 case). However, the faces of $\mathcal{A}(\alpha)^* \subset \mathcal{K}_n^*$ are all convex polytopes by construction. It turns out that the amplitude $m(\alpha|\alpha)$ is given by a sum of terms in $I_{\partial \mathcal{C}(\alpha)^*}$.

4.2 The formula

The integral I_P computes biadjoint scalar amplitudes when P is the cone generated by $\partial \mathcal{A}(\alpha)^*$. To see this, consider the boundary $\partial \mathcal{A}(\alpha)^*$ of the dual associahedron, canonically embedded in \mathcal{K}_n^* . This is the union of the embedded facets of $\mathcal{A}(\alpha)^*$, which we call $\{F_i\}$. So

$$\partial \mathcal{A}(\alpha)^* = \bigcup_i F_i. \tag{4.3}$$

Since n-3 is smaller than dim $\mathcal{K}_n^* = n(n-3)/2$, we can use equation (4.2) to evaluate I on $\partial \mathcal{A}(\alpha)^*$. If A and B are cones that intersect in a strictly lower dimensional cone, then the standard properties of integrals tell us that

$$I_{A\cup B} = I_A + I_B$$

It follows from this, and equation (4.3), that

$$I_{\text{Cone}(\partial \mathcal{A}(\alpha)^*)}(Z,1) = \sum_{\text{facets } F} \frac{1}{\prod_{i=1}^{n-3} (W_{F(i)} \cdot Z)},$$
(4.4)

where the facet F is given by the convex hull of $W_{F(1)}, \ldots, W_{F(n-3)}$ and the volume of the unit box spanned by these vectors is one.¹⁹ The right-hand-side of equation (4.4) is a biadjoint scalar tree amplitude. Indeed,

$$m(\alpha|\alpha) = I_{\text{Cone}(\partial \mathcal{A}(\alpha)^*)}(Z, 1).$$
(4.5)

This is a special case of the theorem, theorem 4.1, that we prove below. We emphasise here that equation (4.5) is merely an alternative form of the result presented for $m(\alpha|\alpha)$ by AHBHY. Their result is presented in sections 5.1 and 5.2 of [1], and the equivalence with equation (4.5) follows almost immediately from equation (5.7) in section 5.2 of their paper.

Theorem 4.1. For dual associahedra $\mathcal{A}(\alpha)^*$ and $\mathcal{A}(\beta)^*$ embedded in \mathcal{K}_n^* as described in section 3, the biadjoint scalar amplitudes are given by

$$m(\alpha|\beta) = I_{Cone(\partial \mathcal{A}(\alpha)^* \cap \partial \mathcal{A}(\beta)^*)}(Z, 1).$$

¹⁹In abstract terms, we could normalise the volume so that the W vectors span unit-volume cells. Concretely, however, taking coordinates $Z = (s_{ij})$ and defining W_S as in equation (2.4), the vectors $W_{F(1)}, \ldots, W_{F(n-3)}$ span a parallelogram formed by translations of a unit box.

Proof. Given the set up as we have described it, this result is almost tautological. Vertices of $\mathcal{A}(\alpha)$ correspond to α -planar trivalent graphs. So facets of $\mathcal{A}(\alpha)^*$ correspond to α -planar trivalent graphs. Suppose that a particular α -planar graph g has propogators $1/S_{g(1)}, \ldots, 1/S_{g(n-3)}$ where $-S_{g(i)} = W_{F(i)} \cdot V$ for some vectors $W_{F(i)} \in \mathcal{K}_n^*$. Then the corresponding facet F of $\mathcal{A}(\alpha)^*$ is canonically embedded in \mathcal{K}_n^* as

$$\operatorname{Conv}(W_{F(1)},\ldots,W_{F(n-3)}).$$

If g is also a β -planar graph, then $\operatorname{Conv}(W_{F(1)}, \ldots, W_{F(n-3)})$ will also be a facet of $\mathcal{A}(\beta)^*$, canonically embedded in \mathcal{K}_n^* . In this way, the α, β -planar graphs are in 1-1 correspondence with shared facets of the dual associahedra $\mathcal{A}(\alpha)^*, \mathcal{A}(\beta)^* \subset \mathcal{K}_n^*$ under the canonical embedding. The result then follows from

$$\sum_{\text{shared } F} \frac{1}{\prod_{i=1}^{n-3} (W_{F(i)} \cdot V)} = \sum_{\alpha, \beta - \text{planar } g} \frac{1}{\prod_{i=1}^{n-3} (-S_{g(i)})}.$$

This last expression is $m(\alpha|\beta)$.

The statement of theorem 4.1 emphasises the cones generated by the boundaries $\partial \mathcal{A}(\alpha)^*$. However, the integral $I_{\text{Cone}P}$ can also be regarded, if you prefer, as a volume or valuation of the polytope P itself. To emphasise the role of the dual associahedra, we might define

$$\operatorname{Val}(P)(Z) = I_{\operatorname{Cone}(P)}(Z, 1). \tag{4.6}$$

Notice that Val is a valuation (i.e. 'like' a volume). This follows from the valuation property for I_P ($I_{P\cup Q} = I_P + I_Q - I_{P\cap Q}$) and some other observations such as $\text{Cone}(P \cup Q) =$ $\text{Cone}(P) \cup \text{Cone}(Q)$. This said, we can write the following.

Corollary 4.2. The biadjoint scalar tree amplitudes are given by

$$m(\alpha|\beta) = Val(\partial \mathcal{A}(\alpha)^* \cap \partial \mathcal{A}(\beta)^*),$$

where Val is the valuation defined in equation (4.6).

4.3 Examples

We now illustrate theorem 4.1 at four and five points. No essentially new phenomena appear at higher points, so these examples suffice to illustrate the result.

Example 4.3. The kinematic associahedra at four points were presented explicitly in section 3, equations (3.1) to (3.3). For example, the dual associahedron $\mathcal{A}(1234)^*$ is given by

$$\mathcal{A}(1234)^* = Conv(W_{12}, W_{23}),$$

and it has two zero-dimensional facets, namely W_{12} and W_{23} . Then we compute

$$I_{Cone(W_{12})\cup Cone(W_{23})} = \frac{1}{W_{12} \cdot V} + \frac{1}{W_{23} \cdot V}.$$

It follows that the amplitude is

$$m(1234|1234) = \frac{1}{-s_{12}} + \frac{1}{-s_{23}}.$$



Figure 10. The shared facet of $\mathcal{A}(12345)^*$ and $\mathcal{A}(14235)^*$. Vertices in the net with the same label are identified under the canonical embedding into \mathcal{K}_5^* .

Example 4.4. Recall from equation (3.3) that

$$\mathcal{A}(3124)^* = Conv(W_{13}, W_{12}).$$

Then we evaluate the intersection

$$\mathcal{A}(1234)^* \cap \mathcal{A}(3124)^* = W_{12}.$$

This implies the amplitude

$$m(1234|3124) = Val(Cone(W_{12})) = \frac{1}{W_{12} \cdot V} = \frac{1}{-s_{12}}.$$

Example 4.5. At five points, the kinematic space \mathcal{K}_5 is defined by the hyperplane

$$s_{12} + s_{13} + s_{14} + s_{23} + s_{24} + s_{34} = 0$$

in $V = \mathbb{R}^6$. Consider the ordering $\alpha = 12345$. The cone $\mathcal{C}(\alpha)$ is defined by the inequalities

$$\begin{split} X_{13} &= -s_{12} \ge 0 \\ X_{14} &= -s_{123} \ge 0 \\ X_{24} &= -s_{23} \ge 0 \\ X_{25} &= -s_{234} \ge 0 \\ X_{35} &= -s_{34} \ge 0. \end{split}$$

The dual cone $\mathcal{C}(\alpha)^*$ is the conic hull

$$\mathcal{C}(\alpha)^* = Cone(W_{12}, W_{123}, W_{23}, W_{234}, W_{34}).$$

(For the definition of the dual vectors W_S , see equation (2.4) in section 2.3.) Similarly, the dual cone $C(\beta)^*$ for $\beta = 14235$ is the conic hull

$$\mathcal{C}(\beta)^* = Cone(W_{14}, W_{124}, W_{24}, W_{234}, W_{23}).$$

The associated dual associahedra, $\mathcal{A}(\alpha)^*$ and $\mathcal{A}(\beta)^*$, canonically embedded in \mathcal{K}_5^* , share one face:

$$\mathcal{A}(\alpha)^* \cap \mathcal{A}(\beta)^* = Conv(W_{23}, W_{234})$$

This can be read off from figure 8, or computed explicitly by listing the faces of the two dual associahedra. See figure 10 for a drawing that highlights the intersection. We compute

$$I_{Cone(W_{23}W_{234})}(Z,1) = \frac{1}{Z \cdot W_{23}Z \cdot W_{234}}$$

So the amplitude is

$$m(\alpha|\beta) = Val(Conv(W_{23}, W_{234})) = \frac{1}{s_{23}s_{234}}$$

Example 4.6. To give another example, the dual cone for $\gamma = 12435$ is the conic hull

 $\mathcal{C}(\gamma)^* = Cone(W_{12}, W_{124}, W_{24}, W_{234}, W_{34}).$

The dual associahedron $\mathcal{A}(\gamma)^*$, canonically embedded in \mathcal{K}_5^* , shares two faces with $\mathcal{A}(\alpha)^*$. This can be read off from figure 8 or computed by hand. The shared faces are

 $Int = \mathcal{A}(\alpha)^* \cap \mathcal{A}(\gamma)^* = Conv(W_{234}, W_{34}) \cup Conv(W_{34}, W_{12}).$

The valuation I of this intersection is

$$I_{Cone(Int)}(Z,1) = \frac{1}{W_{34} \cdot ZW_{234} \cdot Z} + \frac{1}{W_{34} \cdot ZW_{12} \cdot Z}$$

The amplitude is

$$m(\alpha|\gamma) = Val(Int) = \frac{1}{s_{23}s_{234}} + \frac{1}{s_{12}s_{34}}$$

Example 4.7. As a final example, the cone $C(\delta)^*$ for $\delta = 13524$ is given by

$$\mathcal{C}(\delta)^* = Cone(W_{13}, W_{24}, W_{124}, W_{14}, W_{134})$$

One can readily verify that $\mathcal{C}(\delta)^* \cap \mathcal{C}(\alpha)^* = \emptyset$.²⁰ This means that the dual associahedra $\mathcal{A}(\alpha)^*$ and $\mathcal{A}(\delta)^*$ share no common faces in the canonical embedding into \mathcal{K}_5^* . (This can also be read off from figure 8.) This means that the amplitude vanishes:

$$m(\alpha|\delta) = Val(\mathcal{A}(\alpha)^* \cap \mathcal{A}(\delta)^*) = 0.$$

²⁰This is an exercise in linear reduction. One should write the generators of $\mathcal{C}(\delta)^*$ as linear combinations of the generators of $\mathcal{C}(\alpha)^*$ and observe that each generator of $\mathcal{C}(\delta)^*$ lies outside of $\mathcal{C}(\alpha)^*$.



Figure 11. The lattice of points with integer Mandelstam variables, L, and its dual lattice, L^* .

5 A connection with the KLT kernel

In this section we state a connection between the dual associahedra in \mathcal{K}_n^* and the inverse KLT kernel. We denote the inverse KLT kernel by $m_{\alpha'}(\alpha|\beta)$, following Mizera [12]. In section 5.3 we present a new formula for the diagonal elements $m_{\alpha'}(\alpha|\alpha)$. This formula involves a discrete sum over a lattice which we introduce in section 5.1. In kinematic space, \mathcal{K}_n , the lattice is a standard lattice of points for which the Mandelstam variables s_{ij} take integer values (or, if you prefer, are integer multiples of $1/\alpha'$). Before we prove the formula in section 5.3, we review some technical preliminaries in sections 5.1 and 5.2. But to put this discussion in context, let us begin by recalling that Kawai-Lewellen-Tye derived a relation of the form

$$A_{\text{closed}} = \sum_{\alpha,\beta} A_{\text{open}}(\alpha) \text{KLT}(\alpha|\beta) A_{\text{open}}(\beta)$$

between closed string tree amplitudes and open string amplitudes. This relation was derived in [11] using an analytic continuation argument to deform the integration contours of the amplitudes in such a way that A_{closed} factorises as shown. The kernel, $\text{KLT}(\alpha|\beta)$, can be inferred given knowledge of the string amplitudes. However, in [12], Mizera conjectured an algorithmic description of the inverse kernel, $m_{\alpha'}(\alpha|\beta)$, that makes no reference to the string amplitudes (and their associated hypergeometric functions). In subsequent work, Mizera showed that his formulas for $m_{\alpha'}(\alpha|\beta)$ can be regarded as (twisted) intersection pairings of the associahedra that tile $\mathcal{M}_{0,n}(\mathbb{R})$ [14]. For this reason, it is interesting that $m_{\alpha'}(\alpha|\alpha)$ appears naturally in our present context, where we are concerned with intersecting dual associahedra. For further speculations about whether the two presentations are related, see section 6.

5.1 The lattice

Recall from section 2 that $\mathcal{V}_n = \mathbb{R}^{n(n-3)/2+1}$ is the vector space with coordinates s_{ij} for all pairs i, j strictly less than n. Let \tilde{L} be the standard lattice $\mathbb{Z}^{n(n-3)/2+1} \subset \mathcal{V}_n$ of points for which the s_{ij} are all integers. Given this, we have a lattice $L = \tilde{L} \cap \mathcal{K}_n$ in kinematic space.²¹ It turns out that one can choose generators for the cones $\mathcal{C}(\alpha)$ from among vectors in the lattice L.²² For this reason, the cones $\mathcal{C}(\alpha)$ are called 'rational' cones, with respect to the lattice L. Dual to L, we have the lattice L^* in \mathcal{K}_n^* . The cones $\mathcal{C}(\alpha)^*$ are generated by vectors in L^* . Indeed, the vectors W_S that we have been using as generators for $\mathcal{C}(\alpha)^*$ are L^* -vectors. This is because

$$Z \cdot W_S \in \mathbb{Z}$$
 for all $Z \in L$, (5.1)

which is the condition for W_S to be a vector in L^* . To illustrate all of this, consider again the example of four points. A basis for the lattice L in $\mathcal{K}_3 \subset \mathbb{R}^3$ is

$$L = \mathbb{Z} \left(\begin{bmatrix} 1\\0\\-1 \end{bmatrix}, \begin{bmatrix} 0\\1\\-1 \end{bmatrix} \right).$$

Alternatively, we can employ the coordinates on \mathcal{K}_3 that we introduced in equation (2.2) (section 2.1). In these coordinates,

$$L = \mathbb{Z}\left(\begin{bmatrix} 1/2\\\sqrt{3}/2 \end{bmatrix}, \begin{bmatrix} 1\\0 \end{bmatrix} \right).$$

This is just the triangular lattice, see figure 11. On the other hand, the dual lattice is

$$L^* = \mathbb{Z}\left(W_{12}, W_{23}\right),$$

which is the triangular lattice rotated by $\pi/6$. Notice that $W_{13} = -W_{12} - W_{13}$, and so W_{13} is also a vector in L^* .

5.2 Preliminaries

Before we can present our formula for $m_{\alpha'}(\alpha|\alpha)$ in theorem 5.2, we need to introduce a sum over lattice points which is the direct analogy of the integral I_C that we introduced in equation (4.1) (section 4.1). Let P be a polyhedron in \mathcal{K}_n^* . We are interested in the sum

$$S_P(V, \alpha') = \sum_{W \in P \cap L^*} e^{-\alpha' W \cdot V}$$

Just as the integral I_P is the related to the volume of P, the sum S_P is related to the number of L^* -points contained in P. If P is a bounded polytope, the number of L^* -points would be given by

$$#(P \cap L^*) = \lim_{\alpha' \to 0} S_P(V, \alpha').$$

$$\{(1, -1, 0, 0, \ldots), (0, 1, -1, 0, \ldots), (0, 0, 1, -1, \ldots), \ldots\},\$$

which are all vectors in L and span the hyperplane \mathcal{K}_n .

²¹It is clear that \mathcal{K}_n is a rational subspace with respect to L. That is, we can choose a basis that spans \mathcal{K}_n from among vectors lying in L. Indeed, for some ordering of the s_{ij} we could adopt a basis of the form

²²The proof is as follows. The vectors W_S defined in equation (2.4) are L^* -vectors, as shown in equation (5.1) of the main text. This means that the hyperplanes $W_S \cdot Z = 0$ in \mathcal{K}_n are *L*-rational. The rays of the cone $\mathcal{C}(\alpha)$ are given by intersecting these hyperplanes. It follows that the rays are *L*-rational and, therefore, they are generated by some vector in *L*.

For now, let α' be positive and real. When P is a cone, the sum S_P is well defined if V is in the interior of the dual cone P^* .²³ Now fix some cone

$$C = \operatorname{Cone}(W_1, \ldots, W_k),$$

where $k \leq \dim \mathcal{K}_n^*$. The sum is particularly easy to evaluate if the vectors W_i are all generators of the lattice L^* . One finds

$$S_C(V, \alpha') = \prod_{i=1}^k \frac{1}{1 - e^{-\alpha' W_i \cdot V}},$$
(5.2)

for $V \in C^*$.²⁴ In section 4 we studied the integral $I_{\partial C}$ over the boundary, ∂C , of a cone C. We conclude this subsection by evaluating the discrete sum $S_{\partial C}$. We find a large number of new terms, when compared with $I_{\partial C}$, because of the valuation property of the sum S_P . In effect, the sum S_P 'sees' not just the faces $\operatorname{Cone}(F)$ of the cone C, but also their intersections $\operatorname{Cone}(F_i) \cap \operatorname{Cone}(F_j)$. This turns out to be crucial for the realtionship with the inverse KLT kernel, which we explain in section 5.3.

Lemma 5.1. Let C be a cone in the vector space \mathcal{K}^* whose boundary is generated by some set of polytopes $\{F\}$ of the same dimension that intersect in polytopes of strictly lower dimension. Then

$$S_{\partial C}(Z,\alpha') = -\sum_{k=1}^{k} (-1)^k \sum_{F_1,\dots,F_k} S_{Cone(F_1\cap\dots\cap F_k)}(Z,\alpha'),$$

where the summations are over all k-tuples, $\{F_1, \ldots, F_k\}$, such that the intersection $F_1 \cap \ldots \cap F_k$ is not empty.

Proof. The sum S_P is a 'valuation' in the sense that it behaves like a volume:

$$S_{P\cup Q} = S_P + S_Q - S_{P\cap Q},\tag{5.3}$$

for two polyhedra P and Q. The lemma is a consequence of this property, writing

$$\partial C = \bigcup_{i} \operatorname{Cone}(F_i).$$

To arrive at the formula explicitly is an exercise in induction, making use of the observation that

$$\operatorname{Cone}(F_i) \cap \operatorname{Cone}(F_j) = \operatorname{Cone}(F_i \cap F_j)$$

and using standard set theory identities such as $(A \cup B) \cap X = (A \cap X) \cup (B \cap X)$.

²³Recall that, for such a $V, W \cdot V \ge 0$ for all $W \in P$. The convergence of the sum $S_P(V, \alpha')$ then follows from the usual convergence statement:

$$\sum_{n=0}^{\infty} x^n$$

converges to 1/(1-x) if |x| < 1.

²⁴See appendix A.4, for the analogous formulas when W_i are not lattice generators and for when $k > \dim \mathcal{K}_n^*$. The general idea is to repeatedly apply

$$\sum_{n=0}^{\infty} e^{-n\alpha' W_i \cdot Z} = \frac{1}{1 - e^{-\alpha' W_i \cdot Z}}.$$

5.3 The formula

We will now observe that the sum $S_{\partial C(\alpha)}$, defined above, is related to the diagonal components $m_{\alpha'}(\alpha|\alpha)$ of the inverse KLT kernel.²⁵

Theorem 5.2. For a dual associahedron $\mathcal{A}(\alpha)^*$ canonically embedded in \mathcal{K}_n^* (as described in section 2),

$$S_{Cone(\partial \mathcal{A}(\alpha)^*)}(Z, 2\pi i\alpha') = 1 - m_{\alpha'}(\alpha | \alpha)$$

The left hand side can be found by, for instance, evaluating $S_{Cone\partial \mathcal{A}(\alpha)^*}(Z, \alpha')$ and then analytically continuing the result by replacing α' with $2\pi i \alpha'$.

Proof. This follows directly from lemma 5.1. Indeed, combining the lemma with equation (5.2),

$$1 - S_{\text{Cone}(\partial \mathcal{A}(\alpha)^*)}(Z, \alpha') = 1 + \sum_{k=1}^{n-2-k} \sum_{F_1, \dots, F_k} \prod_{i=1}^{n-2-k} \frac{1}{e^{-\alpha' W_{F(i)} \cdot Z} - 1}$$

Since the sum only includes k-tuples (F_1, \ldots, F_k) such that $F_1 \cap \ldots \cap F_k \neq \emptyset$, any $F_1 \cap \ldots \cap F_k$ appearing in the sum has n-2-k vertices W_{I_a} . Then $W_{I_a} \cdot V = -S_{I_a}$, where S_{I_a} are the n-2-k propagators associated to the codimension-k face $F_1 \cap \ldots \cap F_k$. (Dually, they are associated to a codimension-(n-2-k) face in $\mathcal{A}(\alpha)$.) Then

$$1 - S_{\text{Cone}(\partial \mathcal{A}(\alpha)^*)}(Z, 2\pi i \alpha') = 1 + \sum_{k=1}^{n} \sum_{F_1, \dots, F_k} \prod_{a=1}^{n-2-k} \frac{1}{e^{2\pi i \alpha' S_{I_a} \cdot Z} - 1}$$

This is the expression for $m_{\alpha'}(\alpha|\alpha)$ discussed by Mizera in [14] (equation (4.19) of Mizera's paper). The formula first appears in [15].

To illustrate theorem 5.2, consider the four point example. For instance,

$$\partial \mathcal{A}(1234)^* = \{W_{12}\} \cup \{W_{23}\}.$$

It follows that

$$1 - S_{\text{Cone}(W_{12})} - S_{\text{Cone}(W_{23})} = 1 + \frac{1}{e^{-\alpha' s_{12}} - 1} + \frac{1}{e^{-\alpha' s_{23}} - 1}$$

After an analytic continuation, $\alpha' \mapsto 2\pi i \alpha'$, this expression becomes

$$m_{\alpha'}(1234|1234) = -\frac{1}{2i\tan(\pi\alpha' s_{12})} - \frac{1}{2i\tan(\pi\alpha' s_{23})}$$

which is the formula also given by Mizera. We recover the amplitude m(1234|1234) from the pole in α' ,

$$m(1234) = \oint \mathrm{d}\alpha' m_{\alpha'}(1234) = \frac{1}{-s_{12}} + \frac{1}{-s_{23}},$$

which follows since $1/\tan(x) \simeq 1/x + \mathcal{O}(x^0)$.

²⁵Note added in proof: at the time of writing, we did not recognise that $1 - S_{\mathcal{C}(\alpha) \cap \mathcal{C}(\beta)}$ can be regarded as giving an off-diagonal element, $m_{\alpha'}(\alpha|\beta)$. The result of this sum differs from the standard function given for $m_{\alpha'}(\alpha|\beta)$ by an overall phase which can be absorbed by a redefinition of the open string amplitudes appearing on either side of the KLT relation. We thank Sebastian Mizera for discussions on this point.

6 Further comments

The main result in this paper is the formula for $m(\alpha|\beta)$ presented in section 4, theorem 4.1. The formula is based on an embedding of dual associahedra $\mathcal{A}(\alpha)^*$ into dual kinematic space \mathcal{K}_n^* . The amplitude $m(\alpha|\beta)$ can then be expressed in terms of the shared faces of $\mathcal{A}(\alpha)^*$ and $\mathcal{A}(\beta)^*$. Arguably, what we have done is "trivial" in the sense that the dual associahedra do no more than express Feynman diagrammatics in a geometric setting. The intersecting faces of the dual associahedra are just a fancy way to describe a sum over trivalent graphs which are (α, β) -planar. There is, however, something interesting about the new presentation. The embedded associahedra $\mathcal{A}(\alpha)^*$ tile a larger object in dual kinematic space. This is the object pictured in figure 5 for four points, where it is a triangle, and in figure 8 for five points, where it resembles a degenerate or 'halved' dodecahedron. What is this object in general? And, given the role that $m(\alpha|\beta)$ plays in the double copy relation, how can this object be related to Yang-Mills and gravity amplitudes? We offer some speculations below, and describe many unresolved loose ends.

The open string moduli space. In section 3 we introduced (n-1)!/2 embeddings of the boundary of the dual associahedron into \mathcal{K}_n^* . We defined the union of these embeddings,

$$FN_n = \bigcup_{\alpha} \partial \mathcal{A}(\alpha)^* \subset \mathcal{K}_n^*.$$
(6.1)

 FN_n does not bound a convex polytope in \mathcal{K}_n^* , except when n = 4 (in which case it bounds an equilateral triangle). The first loose end is to study FN_n in its own right and explain how it is related to $\mathcal{M}_{0,n}(\mathbb{R})$. The open string moduli space, $\mathcal{M}_{0,n}(\mathbb{R})$, after compactification, is tiled by (n-1)!/2 copies of the associahedron. (See, for instance, theorem 3.1.3 of [13].) There is a duality map from each associahedron $\mathcal{A}(\alpha)$ in $\mathcal{M}_{0,n}(\mathbb{R})$ to the corresponding dual associahedron $\mathcal{A}(\alpha)^*$. In this sense, FN_n is something like a 'dual' of the open string moduli space. With some effort, this could probably be made into a precise statement about dual polytopes.²⁶ A duality statement might help us to present a relationship between the intersecting dual associahedra in this paper and the (twisted) intersections of associahedra that have been explored recently by Mizera in [14].

The fan of cones $\mathcal{C}(\alpha)$. A second loose end concerns the cones $\mathcal{C}(\alpha)$ in kinematic space \mathcal{K}_n introduced by AHBHY. There are (n-1)!/2 cones cut out by

$$2^{n-1} - n - 1$$

hyperplanes. At four points, the three cones, $C(\alpha)$, are the 'tangent cones' of a triangle.²⁷ Indeed, consider the triangle

$$T_4 = \operatorname{Conv}(e_{12}, e_{13}, e_{23})$$

²⁶A possible route would be to exploit the relationship of $\mathcal{M}_{0,n}(\mathbb{R})$ to the permutoassociahedron (introduced by Kapranov in [16]) which can be realised as a convex polytope, as shown in [17]. As a convex polytope, we can take the polytope dual of the permutoassociahedron, from which we might be able to recover FN_n .

 $^{^{27}}$ In general, the tangent space at the vertex of a polytope is a cone. See section A.1 and especially equation (A.2) for the definitions.



Figure 12. At four points, the cones $C(\alpha)$ defined by AHBHY arise as the tangent cones to a triangle.

in \mathcal{K}_4 . The vectors e_{12} and e_{23} are as in equation (2.2), while e_{13} is defined as $-e_{12} - e_{23}$. The three vertices of T_4 have tangent cones such as, at the vertex e_{12} ,

$$\operatorname{Cone}(e_{13} - e_{12}, e_{23} - e_{12}).$$

But this cone is just the cone C(2314) as defined by AHBHY. Indeed, we see that C(1234)is the tangent cone of T_4 at e_{23} and C(3124) is the tangent cone of T_4 at e_{13} . We illustrate this in figure 12. It is no longer true at higher points that the cones $C(\alpha)$ arise as the tangent cones to a polytope. This can be seen at five points using the following fact. If a set of cones C_i are the tangent cones of a convex polytope, then the union of their duals, $\cup C_i^*$, is the entire vector space. At five points, one can see that the union $\cup C(\alpha)^*$ is not the whole vector space. (The union doesn't contain, for instance, $-W_{123}$. This is because $-W_{123} = W_{14} + W_{24} + W_{34}$, and there is no cone for which W_{14} , W_{24} , and W_{34} are all generators.) However, it may be possible to give a polytopal interpretation to the fan obtained by including also the cones $-C(\alpha)$. At four points, this fan has six cones and they are the duals of the tangent cones to the hexagon in \mathcal{K}_4^* . The hexagon is a regular permutohedron. At higher points, the fan $\bigcup_{\alpha} C(\alpha) \cup -C(\alpha)$ has (n-1)! cones and it is natural to speculate that it might be dual to the tangent fan of a permutohedron in \mathcal{K}_n^* which has (n-1)! vertices. Though we have not yet verified this.²⁸

The double copy. The amplitudes $m(\alpha|\beta)$ form the inverse kernel for the field theory double copy relation

$$M_n = \sum_{\alpha,\beta} A(\alpha) m^{-1}(\alpha|\beta) A(\beta)$$
(6.2)

between gravity amplitudes M_n and Yang-Mills partial amplitudes $A(\alpha)$. Whilst there are many presentations of the amplitudes $m(\alpha|\beta)$, the emphasis given by AHBHY, and

 $^{^{28}}$ A result of reference [19] says that a polytope is a regular permutohedron if its face lattice is the Weyl fan. For work on expanding cones as a sum of Weyl fans (or their duals, called 'plates'), see N. Early's recent paper, which might help in these computations [20].

pursued in this paper, is combinatorial. This prompts a question: can the Yang-Mills partial amplitudes $A(\alpha)$ be related to a combinatorial object? And can the field theory KLT relation, equation (6.2), be given a combinatorial interpretation? Speculations of this kind have been widespread ever since Bern-Cachazo-Johansson suggested that Yang-Mills partial amplitudes can be written in the form

$$A_{YM} = \sum_{\alpha,\beta} n(\alpha)c(\beta)m(\alpha|\beta),$$

where $c(\beta)$ are colour factors and $n(\alpha)$ are numerators that obey Jacobi-type relations [8]. The field theory KLT relation, equation (6.2), then reads

$$M_n = \sum_{\alpha,\beta} n(\alpha)n(\beta)m(\alpha|\beta)$$

A combinatorial presentation of these identities would perhaps be interesting, especially if it relied on FN_n (discussed following equation (6.1)) which is, in some sense, dual to $\mathcal{M}_{0,n}(\mathbb{R})$.

Kawai-Lewellen-Tye. In section 5, we observed that the inverse KLT kernel is related to point counting in the cones generated by the dual associahedra $\mathcal{A}(\alpha)^*$. Is this a fluke? Or is it related to something more interesting? It is noteworthy that the introduction of the lattice L breaks the projectivity of the original construction by introducing a fixed scale. The importance of the lattice L also suggests that we take advantage of the links with toric geometry. As emphasised below, in 'the toric dictionary,' our formulas for $m(\alpha|\beta)$ have a natural interpretation in terms of integrals over cycles. What is the toric interpretation of our formula for $m_{\alpha'}(\alpha|\beta)$? One might hope that intersection theory on the appropriate toric variety could be used to encode the full inverse KLT kernel. It would then be necessary to understand how our intersecting dual associahedra (or toric cycles) are related to Mizera's presentation of $m_{\alpha'}(\alpha|\beta)$ as the intersection of associahedra in $\mathcal{M}_{0,n}(\mathbb{R})$.

The toric dictionary. We conclude by describing the natural correspondence between cones and toric varieties. Under this correspondence, theorem 4.1 takes on a new significance: the amplitudes $m(\alpha|\beta)$ can be understood as coming from an integral over subvarieties in some toric variety. These integrals can be explicitly evaluated using the Duistermaat-Heckman ('localisation') formula. In this paragraph, we will briefly describe the toric dictionary as it bears on our results in section 4. Some more details are included in appendix A.3, but the relevant facts are as follows. Given a cone $C^* \subset \mathcal{K}_n^*$ and a lattice $L^* \subset \mathcal{K}_n^*$, a standard construction produces an associated variety X_{C^*} with a toric action. The torus T_{C^*} that acts on X_{C^*} may be identified with the unit cell of the lattice L in \mathcal{K}_n . Given these identifications, the moment map for the T_{C^*} action is a map

$$\mu: X_{C^*} \to \operatorname{Lie}(T_{C^*})^* = \mathcal{K}_m^*$$

and the image of this map turns out to be C^* . This means that we can write

$$I_{C^*}(Z,1) = \int_{C^*} e^{-Z \cdot W} \mathrm{d}W|_{C^*} = \int_{X_{C^*}} e^{-Z \cdot \mu(x)} \mathrm{d}x,$$

where dx is the pushforward of dW. The torus action T_{C^*} has a fixed point $x_o \in X_{C^*}$. If the cone C^* is based at the origin, $\mu(x_o) = 0$. The Duistermaat-Heckmann formula then evaluates the integral as

$$\int_{X_{C^*}} e^{-Z \cdot \mu(x)} \mathrm{d}x = \frac{1}{\det_{x_o}(-Z)},$$

where Z is regarded as a generator for a 1-parameter subgroup of T_{C^*} and acts on X_{C^*} at x_o . It turns out that this is just a toric version of the formula

$$\int_{C^*} e^{-Z \cdot W} \mathrm{d}W|_{C^*} = \frac{1}{\prod_i (Z \cdot W_i)}$$

where $C^* = \operatorname{Cone}(W_1, W_2, \ldots)$. Things get more interesting when more than one cone is involved. A union of cones is called a fan (provided that the intersections of the cones are cones). Consider a fan formed from multiple cones C^*, D^*, \ldots in \mathcal{K}_n^* . A standard construction associates a toric variety to this fan. This is constructed by glueing together the toric varieties X_{C^*}, X_{D^*}, \ldots associated to each cone. If the cones C^* and D^* intersect in a cone $C^* \cap D^*$, then $C^* \cap D^*$ defines a subvariety $X_{C^* \cap D^*}$ in both X_{C^*} and X_{D^*} . We can then glue X_{C^*} and X_{D^*} together by identifying this subvariety. Returning to the context of section 4, let Y be the toric variety associated to the fan $\operatorname{Cone}(FL_n)$. (See section 3 for FL_n .) Then all cones $C^* \subset FN_n$ give rise to subvarieties X_{C^*} in Y. If C^* is the intersection of two other cones, then X_{C^*} is also an intersection of toric subvarieties. Since all the n-point amplitudes $m(\alpha|\beta)$ can be written in terms of the integrals I_{C^*} for various $C^* \subset FN_n$, we can rewrite theorem 4.1 as

$$m(\alpha|\beta) = \int_{X_{\alpha\beta}} e^{-Z \cdot \mu(x)} \mathrm{d}x, \qquad (6.3)$$

where $X_{\alpha\beta}$ is some cycle (or toric subvariety) in Y.²⁹ This rewriting of theorem 4.1 is strictly tautological, but it may suggest new ways forward. In particular, notice that the exponent in the integrand, $-Z \cdot \mu(x)$, is a (perfect) Morse function on the toric variety Y. (See [21], for instance.) In fact, we get a whole family of Morse functions by varying Z. Given these observations, we might consider arriving at the amplitudes $m(\alpha|\beta)$ using a supersymmetric-quantum-mechanics model with Y as its target space (in the manner of [22] or [23]). It is not clear whether such a model would be useful for understanding the double copy relation. Moreover, it remains to work out whether or not intersections in the toric variety Y are related to the components $m_{\alpha'}(\alpha|\beta)$ of the inverse KLT kernel.

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²⁹ $X_{\alpha\beta}$ is homologous to what we might write as $X_{\operatorname{Cone}(\partial \mathcal{A}(\alpha)^*)} \cap X_{\operatorname{Cone}(\partial \mathcal{A}(\beta)^*)}$.

A Review of polyhedral cones and convex polytopes

This appendix reviews some facts about polyhedral cones in vector spaces and in discrete lattices. This topic is closely related to toric geometry and a fast-paced review of polyhedral cones appears in appendix A of Oda's textbook on toric geometry [24]. Ref. [25] begins with a substantial review of polyhedra in lattices. I have also relied heavily on [26]. Section A.5 on Euler-Maclaurin formulas presents a recent result that first appears in [27] following earlier work in [28] and [29]. Useful lecture notes on these subjects appear on A. Barvinok's university webpage, some of which have been published in book form, [30].

A.1 Cones

Let V be a vector space of dimension dim V. A polyhedron is a subset $P \subset V$ defined by finitely many linear inequalities and linear equalities. In other words, a polyhedron is the intersection of some planes and half-spaces. In particular, a polytope is a bounded polyhedron. A polytope P is a convex polytope if for all $Z, Z' \in P$, the midpoint (Z+Z')/2is in P, too. One way to produce convex polytopes, is from their vertices. Take some vectors $Z_1, \ldots, Z_k \in V$ as the vertices of a convex polytope P. Then one way to express P is as the convex hull of these points

$$P = \text{Conv}(Z_1, \dots, Z_k) = \left\{ \sum_{i=1}^k c_i Z_i \mid \sum_{i=1}^k c_i = 1 \text{ and } 0 \le c_i \le 1 \,\forall i \right\}.$$

For any polyhedron P, we can define the affine space containing P, Aff(P), as the smallest hyperplane in V containing all of P. By the *interior* of P, Int(P), we mean the interior of P regarded as a subset of Aff(P). The *boundary* of P is then $\partial P = P \setminus Int(P)$. The study of polyhedra can often be reduced to a study of 'polyhedral cones' using theorem A.1, which we will come to shortly. In general, a *cone* is a subset $S \subset V$ which is (i) convex, and (ii) for $Z \in P$, λZ is also in P for all positive real numbers λ . A cone is a *polyhedral cone* if it has flat sides. (i.e. if the cone is also a polyhedron in the sense described above.) One way to produce polyhedral cones is by taking the *conic hull* of some vectors,

$$\operatorname{Cone}(Z_1, \dots, Z_k) = \left\{ \sum_{i=1}^k c_i Z_i \, | \, 0 \le c_i \, \forall \, i \right\}.$$
(A.1)

A cone is a *pointed cone* if its apex is a point. This means that a pointed cone contains no fully extended line (i.e. it contains only half-lines). A polyhedron P can be decomposed into cones. Let Vert(P) be the vertices of P. For any vertex $Z \in Vert(P)$ define the *tangent cone*

$$\operatorname{Tan}_P(Z) = \{ Z + X \mid Z + \lambda X \in P \text{ for some } \lambda > 0 \}$$

and the cone of directions

$$\operatorname{Dir}_{P}(Z) = \{ X \mid Z + \lambda X \in P \text{ for some } \lambda > 0 \}.$$
(A.2)

Clearly these two cones are related by a translation:

$$\operatorname{Tan}_P(Z) = Z + \operatorname{Dir}_P(Z).$$

A useful observation is that P can be decomposed into the sum of its tangent cones.



Figure 13. An example of the polyhedral decomposition into cones, theorem A.1.

Theorem A.1. For any polyhedron P,

$$P \simeq \sum_{Z \in Vert(P)} Tan_P(Z),$$

where \simeq denotes equality modulo the addition or subtraction of polyhedra containing fully extended lines.³⁰

Example A.2. Consider the polyhedron P in $V = \mathbb{R}^2$ with vertices $Z_1 = (0,0)$ and $Z_2 = (1,0)$ and tangent cones

$$Tan_P(Z_1) = Cone\left(\begin{bmatrix} 0\\1\end{bmatrix}, \begin{bmatrix} 1\\0\end{bmatrix}\right)$$

and

$$Tan_P(Z_2) = Z_2 + Cone\left(\begin{bmatrix} 0\\1\end{bmatrix}, \begin{bmatrix} -1\\0\end{bmatrix}\right).$$

Then the sum of these two cones is

$$Tan_P(Z_1) + Tan_P(Z_2) = P + H,$$

where H is the upper-half-plane $H = \{(x, y) \in V | y \ge 0\}$. Since H contains a fully extended line, we conclude that

$$Tan_P(Z_1) + Tan_P(Z_2) \simeq P_2$$

as in the theorem. See figure 13.

For any set S we may define its *dual*

$$S^* = \{ W \in V^* \, | \, W \cdot Z \ge -1 \, \forall \, Z \in S \}.$$

If C is a cone, this definition implies that

$$C^* = \{ W \in V^* \,|\, W \cdot Z \ge 0 \,\forall \, Z \in C \}.$$
(A.3)

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We now mention an interesting result, related to the decomposition in theorem A.1.

 $^{^{30}}$ This is theorem 3.5 in [25].



Figure 14. Another example of theorem A.3.

Theorem A.3. Let P be a polytope in V with vertices $Vert_P$. Then

$$\bigcup_{Z \in Vert_P} Dir_P(Z)^* \simeq V^*,$$

where \simeq denotes equality as sets modulo excision by polyhedra contained in proper subspaces of V^* .

Example A.2 is a non-example of this theorem, since the infinite strip is not a polytope (it is not bounded). Indeed, the union of the dual cones $\text{Dir}_P(Z_1)^*$ and $\text{Dir}_P(Z_2)^*$ in example A.2 only cover a half-space in $V^* = \mathbb{R}^2$. For a bonafide example of theorem A.3 consider the equalitateral triangle.

Example A.4. The equilateral triangle is the convex hull of three points

$$e_1 = (0, 1),$$
 $e_2 = (\sqrt{3}/2, -1/2)$ $e_3 = (-\sqrt{3}/2, -1/2).$

Let $T = Conv(e_1, e_2, e_3)$ be the triangle. Then the direction cones are, for instance,

$$Dir_T(e_1) = Cone(e_2 - e_1, e_3 - e_1).$$

Notice that e_3 is orthogonal to $e_2 - e_1$ and e_2 is orthogonal to $e_3 - e_1$. Based on this, one can show that the dual cone is

$$Dir_T(e_1)^* = Cone(e_2, e_3).$$

Likewise,

$$Dir_{T}(e_{2})^{*} = Cone(e_{3}, e_{1})$$
 and $Dir_{T}(e_{3})^{*} = Cone(e_{1}, e_{2}).$

It is clear that these three cones fill the vector space.

This example appears in the text in connection with the dual associahedra for the four point amplitude. See section 2.3. For another example of the theorem, consider the hexagon shown in figure 14. The cones defined by the vertices of the hexagon overlap with each other. However, the dual cones tile the dual space in agreement with the theorem.

A.2 Continuous valuation

We now define a valuation (or 'volume') on polyhedra, and on cones in particular, that plays a significant role in section 4 of the main text. Let $P \subset V$ be a polyhedron containing no fully extended line. For $W \in V^*$ we consider the integral

$$I_P(W, \alpha') = \int_P e^{-\alpha' W \cdot Z} \mathrm{d}Z_P.$$

When this integral does not converge, we set $I_P(W, \alpha') = 0$. If the interior of P has dimension d (possibly smaller than dim V) the measure dZ_P appearing in the integral is the d-dimensional Euclidean volume on the interior of P. Notice the following shift property,

$$I_{Z_o+P}(W) = e^{-\alpha' W \cdot Z_o} I_P(W).$$

We now compute an example.

Example A.5. Consider the cone C = Cone((1,0), (0,1)), which is the first quadrant of the plane. If $W = (w_1, w_2)$ are coordinates on V^* ,

$$I_C(W,\alpha') = \int_0^\infty da \int_0^\infty db \, e^{-\alpha' a w_1 - \alpha' b w_2} = (-1)^2 \frac{1}{(\alpha')^2 w_1 w_2}.$$

The result of the calculation in example A.5 is a function which is divergent as $w_1 \to 0$ or $w_2 \to 0$. Indeed, $w_1 = 0$ and $w_2 = 0$ are the boundaries of the dual cone C^* . In general, if C is a cone, then the integral

$$I_C(W, \alpha')$$

is well defined for all $W \in C^*$ and diverges to positive infinity as W approaches the boundary of C^* . (See [24] proposition A.10.) The calculation in example A.5 generalises to an arbitrary polyhedral cone.

Theorem A.6. If $C = Cone(Z_1, \ldots, Z_k)$ be a polyhedral cone with $k \leq \dim V$ we compute that

$$I_C(W) = (-1)^k \frac{Vol(Box(Z_1, \dots, Z_k))}{\prod_{i=1}^k W \cdot Z_i},$$

where $Box(Z_1, \ldots, Z_k)$ is the unit (open-closed) box $\{\sum_{i=1}^k c_i Z_i \mid 0 \le c_i < 1 \forall i\}$.

This theorem follows by computing I for the standard cone $(\mathbb{R}^+)^k$ in \mathbb{R}^k and then using the map from $(\mathbb{R}^+)^k$ to C. What happens when k is larger than dim V? In this case, we can evaluate I_C for $C = \text{Cone}(Z_1, \ldots, Z_k)$ by giving a triangulation of the polytope $P = \text{Conv}(Z_1, \ldots, Z_k)$. Suppose that I_a are a collection of subsets $I_a \subset \{1, \ldots, k\}$ with length dim V. Associated to each such subset is a cone

$$C_a = \operatorname{Cone}\{Z_i \,|\, i \in I_a\}.$$



Figure 15. Decomposing a polyhedral cone using a triangulation.

If the subsets I_a define a triangulation of P, then

$$C = \sum_{a} C_{a}.$$

We sketch an example in figure 15. By the usual linearity of integration, we arrive at the following result.

Theorem A.7. Let $C = Cone(Z_1, \ldots, Z_k) \subset V$ for $k > \dim V$ and suppose that some subsets $I_a \subset \{1, \ldots, k\}$ define a triangulation of $P = Conv(Z_1, \ldots, Z_k)$ in (d-1)-simplices. Then

$$I_C(W) = \sum_a I_{C_a}(W) = \sum_a (-1)^{\dim V} \frac{Vol(Box_a)}{\prod_{i \in I_a} W \cdot Z_i},$$

where Box_a is the unit box generated by the Z_i with $i \in I_a$.

We can use the decomposition theorem, Theorem A.1, to compute I_P for any polyhedron P as

$$I_P(W,\alpha') = \sum_{Z \in \operatorname{Vert}(P)} I_{\operatorname{Tan}_P(Z)}(W,\alpha') = \sum_{Z \in \operatorname{Vert}(P)} e^{-\alpha' W \cdot Z} I_{\operatorname{Dir}_P(Z)}(W,\alpha').$$
(A.4)

The volume of P is ostensibly given by the limit $\alpha' \to 0$. We can extract the volume of P from $I_P(W, \alpha')$ by considering its Laurent expansion in α' and taking the $(\alpha')^0$ term. This is tractable because theorems A.6 and A.7 show us that

$$I_{\text{Dir}_P(Z)}(W, \alpha')$$

is always homogeneous in α' . If the interior of P has dimension d, then $I_{\text{Dir}_P(Z)}(W, \alpha')$ is homogeneous with weight -d.

Theorem A.8. Let P be a polyhedron as above. The $(\alpha')^0$ term in the Laurent expansion of $I_P(W, \alpha')$ is

$$I_P^{(0)}(W) = \sum_{Z \in Vert(P)} \frac{1}{d!} (W \cdot Z)^d I_{Dir_P(Z)}(W).$$

When P is a polytope, this expression is the volume of P. However, it is also well defined when P is not bounded. We give two examples.

Example A.9. Let P be the strip defined in example A.2. If $W = (w_1, w_2)$ are the dual coordinates, we find

$$I_P^{(0)}(W) = \frac{1}{2}\frac{w_1}{w_2}.$$

So, even though P has infinite Euclidean volume, $I_P^{(0)}(W)$ is still an interesting function. Indeed, the function takes values in [0,0.5] with extreme values on the boundary of P.

Example A.10. Let T be the triangle introduced in example A.4. The integrals $I_{Dir_P(Z)}(W)$ evaluate to give, for example,

$$I_{Dir_{T}(e_{1})}(W) = \frac{3\sqrt{3}}{2} \frac{4}{3} \frac{1}{w_{1}^{2} - 3w_{2}^{2}}$$

Summing the contributions in theorem A.8 gives

$$I_T^{(0)}(W) = \frac{3\sqrt{3}}{2}.$$

In this case we see that $I_T^{(0)}(w)$ is the Euclidean area of the triangle, which is what we expect.

A.3 Lattices and toric varieties

Now let L be a lattice in V. A vector Z is L-rational if an integer multiple of it, nZ, is a lattice point. More generally, a subspace $W \subset V$ is L-rational if it is the affine space generated by points in L. (Equivalently, W is rational if $W \cap L$ is a lattice.) A polyhedral Cis rational if it can be generated by a collection of lattice points. In sections A.4 and A.5 we discuss some concrete constructions based on rational cones in a lattice — these are used in section 5.3 of the main text to discuss a possible relation between the cones in kinematic space and the inverse KLT kernel. In this section, we will briefly recall the construction of toric varieties from cones and polyhedra. Our reason for doing this is to point out that formulas for the integral $I_P(W)$ like theorem A.6 and equation (A.4) are, in fact, disguised forms of the Duistermaat-Heckman formula. Readers not interested in this connection may skip to section A.4.

Consider the rational cone, $C = \text{Cone}(Z_1, \ldots, Z_k)$, where the Z_i are all lattice points. Let Aff(C) be the affine space containing C, as above. Clearly

$$\operatorname{Aff}(C) = \operatorname{Span}_{\mathbb{R}}(Z_1, \dots, Z_k).$$

We can consider the lattice

$$M = \mathbb{Z}[Z_1, \ldots, Z_k] \subset \operatorname{Aff}(C),$$

which is a sublattice of L. This lattice defines a torus

$$T = \operatorname{Aff}(C)/M$$

which we can identify with the open-closed unit box

Box
$$(Z_1, \ldots, Z_k) = \left\{ \sum_{i=1}^k c_i Z_i \, | \, 0 \le c_i < 1 \right\}.$$

And the image of L in T, which is $G = L \cap \operatorname{Aff}(C)/M$, is a finite subgroup of T. Associated to the lattice M we can consider the algebra of characters. For any $Z \in M$,

$$\chi_Z(W) = e^{2\pi i W \cdot Z}$$

and we define the algebra $\mathbb{C}[M]$ to be generated by these characters with the multiplication $\chi_Z \cdot \chi_{Z'} = \chi_{Z+Z'}$. Let $L \cap \operatorname{Aff}(C)$ have generators $\mathbb{Z}[\hat{Z}_1, \ldots, \hat{Z}_k]$. Then the algebra associated to $L \cap \operatorname{Aff}(C)$ is generated by the characters $T_i^{\pm 1} = \chi_{+\hat{Z}_i}$. So,

$$\mathbb{C}[L \cap \operatorname{Aff}(C)] = \mathbb{C}[T_1^{\pm 1}, \dots, T_k^{\pm 1}].$$

The characters of M, χ_Z with $Z \in M$, are positive integer powers of the $T_i^{\pm 1}$. So we have an inclusion

$$\mathbb{C}[M] \hookrightarrow \mathbb{C}[T_1^{\pm 1}, \dots, T_k^{\pm 1}]$$

The variety $\operatorname{Spec}\mathbb{C}[M]$ is the toric variety associated to the cone C^* . (For an introduction to these ideas, see Fulton [31], chapter one.)

Example A.11. If $C^* = Cone(e_1, e_2)$ in $V^* = \mathbb{R}^3$ with lattice $L = \mathbb{Z}[e_1, e_2, e_3]$, the dual cone is $C = Cone(e_1, e_2, e_3, -e_3)$. The algebra of characters of $M = \mathbb{Z}_{\geq 0}[e_1, e_2, \pm e_3]$ is

$$\mathbb{C}[M] = \mathbb{C}(T_1, T_2, T_3, T_3^{-1}).$$

The associated toric variety is

$$X = Spec \mathbb{C}[M] = \mathbb{C} \times \mathbb{C} \times \mathbb{C}^*.$$

The torus T = V/M acts on X in the obvious way. There are no fixed points. Each subcone $F \subset C$ is associated with a torus embedded in X. And, in particular, the points $\{0\}$ is associated with the torus $\mathbb{C}^* \times \mathbb{C}^* \times \mathbb{C}^* \subset X$.

Example A.12. If instead we consider $C^* = Cone(e_1, e_2, e_3)$ in $V^* = \mathbb{R}^3$, the associated toric variety is

$$X = Spec \mathbb{C}[M] = \mathbb{C} \times \mathbb{C} \times \mathbb{C}.$$

This has a single fixed point, namely $x_o = (0,0,0) \in X$. In general, if the generators of C^* span V^* , then X is a non-singular variety of the form $\mathbb{C}^{\times n}$ and there is a single fixed point, the origin.

The reason we have recalled the construction of toric varieties is that, in this context, the formula we obtained for $I_C(W)$, theorem A.6, is an example of the Duistermaat-Heckman formula. We will give a terse explanation of this. The key observation is made in [26]. For more on Duistermaat-Heckman, the standard reference is [32]. Let C be a cone with dim V generators in V and let X be the associated toric variety. We can regard the torus T which acts on X as an open box $Box \subset V$. So the Lie algebra of T is V and the moment map is a map

$$\mu: X \to \operatorname{Lie}(T)^* = V^*$$

Then, for $Z \in V$, Duistermaat-Heckman evaluates the integral

$$\int_{X} e^{Z \cdot \mu(x)} \mathrm{d}x = e^{Z \cdot \mu(x_o)} \frac{1}{\det_{x_o} Z},\tag{A.5}$$

where $\det_{x_o} Z$ is the determinant of the action of $Z \in V$ on X at the fixed point x_o . By using the moment map, the integral on the left-hand-side can be identified with

$$\int_{X} e^{Z \cdot \mu(x)} \mathrm{d}x = \int_{C^*} e^{Z \cdot W} \mathrm{d}W|_{C^*}.$$
 (A.6)

On the other hand, the moment map sends x_o to the vertex of the cone, W_o . We then identify the right-hand-side as

$$e^{Z \cdot \mu(x_o)} \frac{1}{\det_{x_o} Z} = e^{Z \cdot W_o} \frac{1}{\prod (-W \cdot Z_i)}$$

In general, a toric variety can be associated to any 'fan' of cones. (A fan is a union of cones whose faces meet each other to form sub-cones, and so on.) This is a standard construction and involves glueing together the toric varieties associated to each cone. See [31], chapter 1. In particular, we can consider the fan of cones F_P associated to a polytope P. The fan F_P is the union of all dual cones $\text{Dir}_P(Z)^*$ for vertices $Z \in \text{Vert}_P$. Let X be the associated toric variety. Then the Duistermaat-Heckman formula reads

$$\int_{X} e^{W \cdot \mu(x)} \mathrm{d}x = \sum_{\text{fixed points}} e^{W \cdot \mu(x)} \frac{1}{\det_{x} W}.$$
 (A.7)

Remarkably, this is just the formula, equation (A.4), that we obtained earlier for $I_P(W, \alpha')$ with α' set to 1. Brion makes this observation in [26], though he may not have been the first. The correspondence between the two formulas is roughly as follows. For each of the dual cones $\text{Dir}_P(Z)^*$, there is a single fixed point $x_Z \in X$ of the torus action on X (just as, for a single cone, there was a single fixed point). The moment map μ maps x_Z to the vertex $Z \in V$. This identifies the right-hand-side of equation (A.7) with equation (A.4). We will not elaborate these ideas in any detail here, — but notice that equation (A.6), for instance, gives another interpretation to the results in section 4 described in the main text.

A.4 Discrete valuation

Given a rational cone C in V with a lattice L, we will define a function $S_C(W)$ on V^* . By analogy with the definition of $I_C(W)$ in section A.2, consider the sum

$$S_C(W, \alpha') = \sum_{Z \in C \cap L} e^{-\alpha' W \cdot Z}.$$

This is well defined when W is in the dual cone C^* . Otherwise, when W is not in the dual cone, the summation is not well defined and we set $S_C(W)$ to zero. Notice, as for $I_C(W)$, the translation property

$$S_{Z_o+C}(W,\alpha') = e^{-\alpha'W \cdot Z_o} S_C(W,\alpha').$$

We can evaluate S_C explicitly in the following simple example.

Example A.13. Let C = Cone((1,0), (0,1)) be the first quadrant of the plane, as in example A.5. Take L to be the standard lattice $L = \mathbb{Z}^2$. Then writing $W = (w_1, w_2)$ for dual coordinates,

$$S_C(W) = \sum_{a=0}^{\infty} \sum_{b=0}^{\infty} e^{-aw_1 - bw_2} = \frac{1}{(1 - e^{-w_1})(1 - e^{-w_2})},$$

provided that $e^{w_1}, e^{w_2} < 1$.

Notice that, in example A.13, e^{w_1} and e^{w_2} are less than 1 for $w_1, w_2 > 0$. That is, the sum $S_C(W)$ is defined precisely for W in the dual cone C^* which is the first quadrant of the plane. At the boundary, $w_1 = 0$ or $w_2 = 0$, the sum diverges. The calculation in example A.13 is easy to generalise. The following is the lattice analog of Theorem A.6.

Theorem A.14. Let $C = Cone(Z_1, \ldots, Z_k)$ for $k \leq \dim V$. Then

$$S_C(W, \alpha') = \left(\sum_{Box \cap L} e^{-\alpha' W \cdot Z}\right) \prod_{i=1}^k \frac{1}{1 - e^{-\alpha' W \cdot Z_i}},$$

where Box is the unit open-closed box generated by the Z_i .

Just as I_P can be used to compute polytope volumes, S_P can be used to count lattice points on the interior of a polytope. Consider the Laurent expansion of $S_P(W, \alpha')$ around $\alpha' = 0$. Then the zero-order term $S_P^{(0)}(W)$ in the Laurent series is ostensibly the number of interior lattice points,

$$S_P^{(0)}(W) = \#(P \cap L).$$

This is strictly true when P is a polytope. When P is not bounded, the right-hand-side is no longer defined, but the left hand side may still be defined.

Example A.15. Let's count the number of lattice points in the triangle T, introduced in example A.4. The decomposition into cones gives a sum of terms of the form

$$S_{e_1+Dir_T(e_1)}(W,\alpha') = \frac{e^{-\alpha'w_1}}{(1-e^{-\alpha'(w_2-w_1)})(1-e^{-\alpha'(w_3-w_1)})},$$

where $w_i = W \cdot e_i$. Since $e_1 + e_2 + e_3 = 0$, we likewise have $w_1 + w_2 + w_3 = 0$. Summing these gives

$$S_T^{(0)}(W) = 1.$$

Indeed, T contains one point of the lattice $L = \mathbb{Z}[e_1, e_2, e_3]$ in its interior.

Example A.16. Let's return to the infinite strip P as defined in example A.2. We have

$$S_P(W, \alpha') = S_{Dir_P(Z_1)} + e^{-\alpha' W \cdot Z_2} S_{Dir_P(Z_2)}$$

= $\frac{1}{(1 - e^{-\alpha' W_1})(1 - e^{-\alpha' W_2})} + \frac{e^{-\alpha' W_2}}{(1 - e^{\alpha' W_1})(1 - e^{-\alpha' W_2})}$

Using the Laurent series

$$\frac{1}{1-e^{\tau}} = -\frac{1}{\tau} + \frac{1}{2} + \dots,$$

we find that the zero-order term in the Laurent series is

$$S_P^{(0)}(w) = 0.$$

A.5 Generalised Euler-Maclaurin formulas

How are the valuations S_P and I_P related to each other? In one dimension, this question is answered by the Euler-Maclaurin formula. Recall the result that, up to a remainder term,

$$\sum_{n=a}^{b} f(n) = \int_{a}^{b} dx f(x) - \sum_{n=1}^{b} \frac{b_n}{n!} f^{(n-1)}(a) + \sum_{n=1}^{b} (-1)^n \frac{b_n}{n!} f^{(n-1)}(b).$$

In our case, we replace f(x) with the exponential $\exp(-\alpha' wz)$. We can present the sums explicitly by making use of the generating function

$$\sum_{m=0} \frac{b_n}{n!} t^n = \frac{t}{1 - e^{-t}}.$$

Then

$$\sum_{n=a}^{b} e^{-\alpha'wz} - \int_{a}^{b} \mathrm{d}x e^{-\alpha'wz} = -\left(\frac{1}{e^{\alpha'w} - 1} - \frac{1}{\alpha'w}\right) e^{\alpha'wa} - \left(\frac{1}{e^{-\alpha'w} - 1} - \frac{1}{-\alpha'w}\right) e^{\alpha'wb}.$$
(A.8)

This curious looking formula has a generalisation to polytopes in higher dimensions. The generalisation is given by Brion, Vergne, and Berline, and their result shows that the valuation S_C , for C a cone, can be expressed in terms of integrals I_F , where the F are subcones of C. We present their result as the following theorem. For brevity, we temporarily omit α' from our formulas, which is equivalent to setting $\alpha' = 1$.

Theorem A.17. There exists some measure μ on rational polyhedra such that

$$S_P(W) = \sum_Q \mu_{P/Q}(W) I(P \cap Q)(W), \tag{A.9}$$

where the sum is over all rational subspaces Q which are tangent to one of the faces (of any dimension) of P.

Notice that a quotient vector space P/Q inherits a lattice \hat{L} from the lattice $L \subset V$. This is given by the projection of L.³¹ It is in this sense that P/Q is a rational polyhedron. Berline-Vergne define $\mu_C(W)$ on cones C inductively, beginning with

$$\mu_{\{0\}}(W) \equiv 1.$$

Putting $C = \text{Cone}(Z_1)$ we impose equation (A.9) to find

$$S_C(W) = I_C(W) + \mu_C(W).$$

We thus infer that

$$\mu_C(W) = \frac{1}{1 - e^{-W \cdot Z_1}} + \frac{1}{W \cdot Z_1}$$

Indeed, this is precisely what we already discovered from the Euler-Maclaurin formula — see equation (A.8). Putting $C = \text{Cone}(Z_1, Z_2)$ in (A.9) we find

$$\mu_C(W) = S_C(W) + \mu_{\text{Cone}(\hat{Z}_1)}(W) \frac{1}{W \cdot Z_1} + \mu_{\text{Cone}(\hat{Z}_2)}(W) \frac{1}{W \cdot Z_2} - \frac{1}{W \cdot Z_1 W \cdot Z_2}$$

where

$$\hat{Z}_1 = Z_1 - \frac{Z_1 \cdot Z_2}{Z_2 \cdot Z_2} Z_2$$
 and $\hat{Z}_2 = Z_2 - \frac{Z_1 \cdot Z_2}{Z_1 \cdot Z_1} Z_1.$

Clearly, this construction could be continued and be used to define $\mu_C(W)$ for all cones C.

Relevance for amplitudes. In the context of our results in sections 4 and 5, these generalised Euler-Maclaurin formulas imply that the diagonal entries $m_{\alpha'}(\alpha|\alpha)$ of the inverse KLT kernel have an exact expansion in terms of the amplitudes $m(\alpha|\beta)$ at finite values of α' . The coefficients in this expansion, given by μ , are complicated expressions, but they can be determined algorithmically and satisfy nice properties (μ is a valuation on polytopes). It is not clear whether this expansion is of any interest physically.

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³¹The lattice \hat{L} is not to be identified with $W^{\perp} \cap L$, which is usually a strict sublattice of \hat{L} .

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