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Second order transport coefficients of nonconformal fluids from compactified Dp-branes

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ABSTRACT: All the 7 dynamical second order transport coefficients of the nonconformal fluids that correspond to Dp-branes with one or more world-volume directions compactified are derived via fluid/gravity correspondence. The conditions considered in this paper include D4-brane with 1, 2 or 3 compact directions, D3-brane with 1 or 2 compact directions, as well as D2-brane with 1 direction compactified. The derived second order transport coefficients satisfy the Haack-Yarom, Romatschke and Kleinert-Probst relations.

KEYWORDS: Holography and quark-gluon plasmas, Gauge-gravity correspondence, D-branes

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1 Introduction

Fluid/gravity correspondence [1, 2] is an effective method of calculating the dynamical second order transport coefficients for various kinds of gravity backgrounds. After it was proposed, this framework has been used in asymptotically Anti de Sitter (AdS) black holes in various dimensions [3–5], with background scalar field [6] and with background vector charge [7, 8].

Branes in string/M theory are the vacuum solutions of 10/11 dimensional supergravity [9]. The methods used in previous studies on the transport properties of branes are the Minkowski AdS/CFT correspondence [10–16], the stretched horizon formalism [17–19] and the boundary derivative expansion method in Fefferman-Graham coordinate [20]. The most remarkable observation that one can make from these studies is that the relativistic fluids

dual to brane systems are nonconformal since background dilaton exists. Following [21], reference [22] proposes the most general formulation of the second order constitutive relations in nonconformal regime, which sets the stage for later systematic investigations on nonconformal relativistic fluids.

Based on these previous works, [23–26] calculate the first and second order transport coefficients for D-branes via the fluid/gravity correspondence and have got some new results. Of those references, [23, 24] study the D4-brane with 1 world-volume direction compact and [26] calculates all the dynamical second order transport coefficients for Dp-brane. This paper can be seen as a follow-up work of [26] and we are aiming to calculate all the dynamical second order transport coefficients for Dp-brane with q of its world-volume direction(s) compactified, where $1 \leq q \leq p - 1$. This can be seen as our first motivation. The maximum of q can not be p since we need to leave at least one dimension on the world-volume for the fluid to live in. We denote a Dp-brane with q direction(s) compact as D(p-q)-brane, such that the cases considered here are D(4-1), D(4-2), D(4-3), D(3-1), D(3-2) and D(2-1)-brane. Actually, the D(4-1) case has been studied in [23, 24], thus we will not offer the details for this case but will include its results in the final expression for the second order constitutive relation.

Our second motivation is to give a final answer to the question proposed in [13] which calculates the first order transport coefficients for Dp-brane. In the conclusion of that paper, in order to make their result compatible with that of [12], the authors proposed an expression for the ratio ζ/η . We quote it here in our convention as

$$\frac{\zeta}{\eta} = \frac{2(p-3)^2 + 2q(5-p)}{(p-q)(9-p)}. \tag{1.1}$$

After discovering the above can cover both their results and that of [12], the authors of [13] say that they do not know whether this compatibility is just a coincidence or due to some reason behind the formula. Now we can answer this question: the above expression of ζ/η is for the compactified Dp-brane, or in our notation, the D(p-q)-brane. When setting $q = 0$, it will be back to the cases in [13] and when setting $p = 4, q = 1$, it reproduces the result of [12].

The structure of the paper is as follows. In section 2 we will reduce the 10 dimensional background of near-extremal black Dp-brane to a $p - q + 2$ dimensional Einstein-dilaton theory. In section 3 we will answer the question proposed in [13] by offering the result of the first order constitutive relation for D(p-q)-brane. Section 4 is a preparation for the second order and then in section 5 we will solve the differential equations for all the second order perturbations. The final result of second order stress energy tensor and transport coefficients will be given in section 6. Section 7 is a summary about this paper.

2 The reduction from 10 to $p - q + 2$ dimensions

This section shows how to reduce the 10 dimensional background of Dp-brane to a $p - q + 2$ dimensional spacetime. Since this work is subsequent to [26], the conventions are the same

as that reference. We will then be brief in offering the details of the definitions which are the same as before and will also stress the differences.

The action for Dp-brane in 10 dimension contains the bulk term, the Gibbons-Hawking surface term and the counter term. It reads as

$$S = \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-G} \left[\mathcal{R} - \frac{1}{2} (\nabla_{\hat{M}} \phi)^2 - \frac{g_s^2}{2(8-p)!} e^{\frac{p-3}{2}\phi} F_{8-p}^2 \right] - \frac{1}{\kappa_{10}^2} \int d^9x \sqrt{-HK} + \frac{1}{\kappa_{10}^2} \int d^9x \sqrt{-H} \frac{9-p}{2L} e^{\frac{3-p}{4(7-p)}\phi}. \quad (2.1)$$

The equations of motion (EOM) derived from this action admit the following background solution:

$$ds^2 = \left[\left(\frac{r}{L_p} \right)^{\frac{(7-p)^2}{8}} \left(-f(r) dt^2 + \delta_{ij} dx^i dx^j \right) + \left(\frac{L_p}{r} \right)^{\frac{(p+1)(7-p)}{8}} \frac{dr^2}{f(r)} \right] + \left(\frac{r}{L_p} \right)^{\frac{(7-p)^2}{8}} \delta_{mn} dy^m dy^n + \left(\frac{L_p}{r} \right)^{\frac{(p+1)(7-p)}{8}} r^2 d\Omega_{8-p}^2, \quad (2.2)$$

$$e^\phi = \left(\frac{r}{L_p} \right)^{\frac{(p-3)(7-p)}{4}}, \quad F_{\theta_1 \dots \theta_{8-p}} = g_s^{-1} Q_p \sqrt{\gamma_{8-p}}, \quad (2.3)$$

where $f(r) = 1 - \frac{r_H^{7-p}}{r^{7-p}}$ and $Q_p = (7-p)L_p^{7-p}$. Here $L_p^{7-p} = \frac{(2\pi l_s)^{7-p} g_s N}{(7-p)\Omega_{8-p}}$, which relates the charge of the Dp-brane and will become the characteristic length of the reduced spacetime.

The 10 dimensional line element eq. (2.2) has three parts and represents a spacetime with structure of $\mathcal{M}_{p-q+2} \times \mathbf{T}^q \times \mathbf{S}^{8-p}$, of which the coordinate is $x^{\hat{M}} = (x^M, y^m, \theta^a)$. Therein, $x^M = (x^\mu, r)$ are the coordinates of \mathcal{M}_{p-q+2} with $\mu = 0, 1, \dots, p-q$; y^m are the coordinates of the compactified directions on D-branes that forms a q -dimensional torus \mathbf{T}^q with $m = 1, \dots, q$; and θ^a are the coordinates of $8-p$ dimensional sphere \mathbf{S}^{8-p} with $a = 1, \dots, 8-p$.

In the following we will implement the dimensional reduction on the compact dimensions of the gravity side \mathbf{T}^q and \mathbf{S}^{8-p} by integrating the coordinates y^m and θ^a , respectively. The dimensional reduction on these two compact subparts of the bulk metric has different significance. The reduction on \mathbf{S}^{8-p} is just to reconcile the Dp-brane metric with the framework of fluid/gravity correspondence. The reduction on \mathbf{T}^q can modify the metric of the reduced bulk gravity, hence causing changes on the results of the second order transport coefficients.

We use the ansatz

$$ds^2 = e^{2\alpha_1 A} g_{MN} dx^M dx^N + e^{2\alpha_2 A} \left(e^{2\beta_1 B} \delta_{mn} dy^m dy^n + e^{2\beta_2 B} L_p^2 d\Omega_{8-p}^2 \right) \quad (2.4)$$

to perform the dimensional reduction. If we denote the metric for the above ansatz as $G_{\hat{M}\hat{N}} = \{e^{2\alpha_1 A} g_{MN}, e^{2\alpha_2 A + 2\beta_1 B} \delta_{mn}, e^{2\alpha_2 A + 2\beta_2 B} L_p^2 \gamma_{ab}\}$ then the 9D induced boundary metric is $H_{\hat{M}\hat{N}} = G_{\hat{M}\hat{N}} - \mathbf{n}_{\hat{M}} \mathbf{n}_{\hat{N}}$ and the corresponding line element is

$$ds^2 = e^{2\alpha_1 A} h_{MN} dx^M dx^N + e^{2\alpha_2 A} \left(e^{2\beta_1 B} \delta_{mn} dy^m dy^n + e^{2\beta_2 B} L_p^2 d\Omega_{8-p}^2 \right), \quad (2.5)$$

where h_{MN} is the boundary metric of \mathcal{M}_{p-q+2} . The unit norm on the 9D boundary is

$$\mathbf{n}_{\hat{M}} = \frac{\nabla_{\hat{M}} r}{\sqrt{G^{\hat{N}\hat{P}} \nabla_{\hat{N}} r \nabla_{\hat{P}} r}} = (\mathbf{n}_M, \mathbf{n}_m, \mathbf{n}_a) = (e^{\alpha_1 A} n_M, 0, 0), \quad (2.6)$$

with n_M the unit norm on the boundary of \mathcal{M}_{p-q+2} .

As said in [23], the requirement that the reduced action is also in Einstein frame gives two relations among those 4 parameters in eq. (2.4). So one may set them as

$$\alpha_1 = -\frac{8-p+q}{p-q}, \quad \alpha_2 = 1, \quad \beta_1 = -\frac{8-p}{q}, \quad \beta_2 = 1. \quad (2.7)$$

The reduced action for the $p-q+2$ dimensional theory turns out to be

$$\begin{aligned} S = & \frac{1}{2\kappa_{p-q+2}^2} \int d^{p-q+2} x \sqrt{-g} \left[R - \frac{1}{2} (\partial\phi)^2 - \frac{8(8-p+q)}{p-q} (\partial A)^2 \right. \\ & \left. - \frac{(8-p)(8-p+q)}{q} (\partial B)^2 + V(\phi, A, B) \right] - \frac{1}{\kappa_{p-q+2}^2} \int d^{p-q+1} x \sqrt{-h} K \\ & + \frac{1}{\kappa_{p-q+2}^2} \int d^{p-q+1} x \sqrt{-h} \frac{9-p}{2L_p} \exp \left[-\frac{8-p+q}{p-q} A + \frac{3-p}{4(7-p)} \phi \right], \quad (2.8) \end{aligned}$$

$$\begin{aligned} V(\phi, A, B) = & \frac{(7-p)(8-p)}{L_p^2} \exp \left(-\frac{16}{p-q} A - 2B \right) \\ & - \frac{Q_p^2}{2L_p^{2(8-p)}} \exp \left[\frac{p-3}{2} \phi - \frac{2(p-q)(7-p) + 16}{p-q} A - 2(8-p)B \right]. \quad (2.9) \end{aligned}$$

In the above,

$$\frac{1}{2\kappa_{p-q+2}^2} = \frac{L_p^{8-p} \Omega_{8-p} V_q}{2\kappa_{10}^2}, \quad (2.10)$$

where $V_q = \int d^q y = \beta_1 \cdots \beta_q$ is the volume of the compact directions on the Dp-brane world-volume. $2\kappa_{p-q+2}^2$ is the surface gravity of the $p-q+2$ dimensional reduced gravity theory and it relates with that of Dp-brane [26] by

$$\frac{1}{2\kappa_{p-q+2}^2} = \frac{V_q}{2\kappa_{p+2}^2}. \quad (2.11)$$

The reduced background fields that solve the reduced action (2.8) are

$$ds^2 = \left(\frac{r}{L_p} \right)^{\frac{9-p}{p-q}} \left(-f(r) dt^2 + d\vec{x}^2 \right) + \left(\frac{r}{L_p} \right)^{\frac{(p^2-8p+9)+q(7-p)}{p-q}} \frac{dr^2}{f(r)}, \quad (2.12)$$

$$e^\phi = \left(\frac{r}{L_p} \right)^{\frac{(p-3)(7-p)}{4}}, \quad (2.13)$$

$$e^A = \left(\frac{r}{L_p} \right)^{\frac{(p-3)^2}{16} + \frac{q(5-p)}{2(8-p+q)}}, \quad e^B = \left(\frac{r}{L_p} \right)^{-\frac{q(5-p)}{2(8-p+q)}}. \quad (2.14)$$

For the cases of $p = 3$, the scalar fields A , B still not vanish. This suggests that the D3-brane will correspond to nonconformal relativistic fluids after compactification. To put it more generalized, dimensional reduction can make the dual fluid of a conformal background nonconformal. This tells us that we can perform a dimensional reduction on AdS black hole backgrounds and calculate the second order transport coefficients of the dual fluids, which may give a direct check on the proposal in [20].

The reduced theory is actually a $p - q + 2$ dimensional Einstein gravity coupled with one background scalar — for which one can choose the dilaton. This is because the three scalars in eqs. (2.13) and (2.14) are not independent as can be seen from the following:

$$\begin{aligned}
 A &= \left[\frac{p-3}{4(7-p)} + \frac{2q(5-p)}{(8-p+q)(p-3)(7-p)} \right] \phi, \\
 B &= \frac{-2q(5-p)}{(8-p+q)(p-3)(7-p)} \phi.
 \end{aligned}
 \tag{2.15}$$

The above are not appropriate for $p = 3$ because of the terms $p - 3$ in the denominator. But from eq. (2.14) one can figure out that $A = -B$ when $p = 3$ and dilaton vanishes. Anyway, there is always just one independent scalar field in the background of the $p - q + 2$ dimensional reduced gravity theory.

In Eddington-Finkelstein coordinate, the boosted form of eq. (2.12) is

$$\begin{aligned}
 ds^2 &= -r^{\frac{9-p}{p-q}} f(r) u_\mu u_\nu dx^\mu dx^\nu + r^{\frac{9-p}{p-q}} P_{\mu\nu} dx^\mu dx^\nu - 2r^{\frac{(p-3)(p-6)+q(7-p)}{2(p-q)}} u_\mu dx^\mu dr, \\
 u^\mu &= \frac{(1, \vec{\beta})}{\sqrt{1 - \vec{\beta}^2}}, \quad P_{\mu\nu} = \eta_{\mu\nu} + u_\mu u_\nu
 \end{aligned}
 \tag{2.16}$$

where now $x^0 = v$ is the in-going Eddington time. The above represents ideal relativistic fluid living on the boundary of \mathcal{M}_{p-q+2} . The fluid/gravity correspondence [1, 2] states that to mimic the real fluid flow, one should first promote the boost parameter u^μ and r_H to be x^μ dependent, then add the (also x^μ dependent) metric perturbations, and finally solve them from EOM in the order of partial derivatives of x^μ in some local patch on the boundary. The global solution of the metric that is dual to arbitrary boundary flow can then be got by stitching together all local solutions. The boundary coordinate dependent global metric ansatz can be set as

$$\begin{aligned}
 ds^2 &= -r^{\frac{9-p}{p-q}} [f(r_H(x), r) - k(r_H(x), u^\alpha(x), r)] u_\mu(x) u_\nu(x) dx^\mu dx^\nu \\
 &+ 2r^{\frac{9-p}{p-q}} P_\mu^\rho(u^\alpha(x)) w_\rho(u^\alpha(x), r) u_\nu(x) dx^\mu dx^\nu \\
 &+ r^{\frac{9-p}{p-q}} [P_{\mu\nu}(u^\alpha(x)) + \alpha_{\mu\nu}(r_H(x), u^\alpha(x), r) + h(r_H(x), u^\alpha(x), r) P_{\mu\nu}(u^\sigma(x))] dx^\mu dx^\nu \\
 &- 2r^{\frac{(p-3)(p-6)+q(7-p)}{2(p-q)}} [1 + j(r_H(x), u^\alpha(x), r)] u_\mu(x) dx^\mu dr.
 \end{aligned}
 \tag{2.17}$$

This metric contains not only the boosted black brane metric but also the artificially set perturbations, which are the tensor perturbation $\alpha_{\mu\nu}(r_H(x), u^\alpha(x), r)$, the vector perturbation $w_\rho(u^\alpha(x), r)$ and the scalar perturbations $k(r_H(x), u^\alpha(x), r)$, $h(r_H(x), u^\alpha(x), r)$ and $j(r_H(x), u^\alpha(x), r)$. They depend on the boundary coordinates x^μ through the proper velocity u^μ and the horizon parameter r_H . All the perturbations will be solved tubewisely [1]

in the bulk of the reduced $p - q + 2$ dimensional nonconformal spacetime by choosing some special point in x^μ directions — which is just $x^\mu = 0$, following [1].

The tubewise solving procedure of eq. (2.17) is performed by expanding it in terms of derivatives of boundary coordinates of u^μ and r_H . The EOM derived from eq. (2.8) that eq. (2.17) satisfies order by order reads

$$E_{MN} - T_{MN} = 0, \tag{2.18}$$

$$\nabla^2 \phi - \frac{(p-3)(7-p)^2}{4L_p^2} \exp \left[\frac{p-3}{2} \phi - \frac{2(p-q)(7-p)+16}{p-q} A - 2(8-p)B \right] = 0, \tag{2.19}$$

$$\begin{aligned} \nabla^2 A - \frac{(7-p)(8-p)}{(8-p+q)L_p^2} e^{-\frac{16}{p-q}A-2B} + \frac{(7-p)^2[(p-q)(7-p)+8]}{16(8-p+q)L_p^2} \\ \times \exp \left[\frac{p-3}{2} \phi - \frac{2(p-q)(7-p)+16}{p-q} A - 2(8-p)B \right] = 0, \end{aligned} \tag{2.20}$$

$$\begin{aligned} \nabla^2 B - \frac{q(7-p)}{(8-p+q)L_p^2} e^{-\frac{16}{p-q}A-2B} + \frac{q(7-p)^2}{2(8-p+q)L_p^2} \\ \times \exp \left[\frac{p-3}{2} \phi - \frac{2(p-q)(7-p)+16}{p-q} A - 2(8-p)B \right] = 0, \end{aligned} \tag{2.21}$$

where E_{MN} is the Einstein tensor and the energy momentum tensor on the gravity side is defined as

$$\begin{aligned} T_{MN} = \frac{1}{2} \left(\partial_M \phi \partial_N \phi - \frac{1}{2} g_{MN} (\partial \phi)^2 \right) + \frac{8(8-p+q)}{p-q} \left(\partial_M A \partial_N A - \frac{1}{2} g_{MN} (\partial A)^2 \right) \\ + \frac{(8-p)(8-p+q)}{q} \left(\partial_M B \partial_N B - \frac{1}{2} g_{MN} (\partial B)^2 \right) + \frac{1}{2} g_{MN} V. \end{aligned} \tag{2.22}$$

The EOM of ϕ , A and B are actually not independent from each other. This is because both A and B are proportional to ϕ , as can be seen from eq. (2.15). The consequence of this fact is that the differential equations that derived from eq. (2.19) to eq. (2.21) are the same, as long as we do not open the perturbations for ϕ , A and B . We will take advantage of this in solving the scalar perturbations for the compactified D3-brane: since there is no dilaton's EOM at hand for D3-brane, we will use the EOM of A instead.

3 The first order calculation

3.1 Solving the perturbations

We need only to solve the compactified Dp-brane with $2 \leq p \leq 4$. Because D5 and D6-brane do not have dual fluids, as has been pointed out in [26] and references therein. D1-brane can not be compactified since it has only one spatial dimension. We need at least 1 spatial dimension for the relativistic fluids to live in and support their viscous terms.

We expand eq. (2.17) to the first order of the derivative of x^μ at $x^\mu = 0$ and get

$$\begin{aligned}
 ds^2 = & r^{\frac{9-p}{p-q}} \left[- \left(f(r) - \frac{(7-p)r_H^{6-p}}{r^{7-p}} \delta r_H - k^{(1)}(r) \right) dv^2 \right. \\
 & \left. + 2((f-1)\delta\beta_i - w_i^{(1)}(r)) dv dx^i + (\delta_{ij} + \alpha_{ij}^{(1)}(r) + h^{(1)}(r)\delta_{ij}) dx^i dx^j \right] \\
 & + 2r^{\frac{(p-3)(p-6)+q(7-p)}{2(p-q)}} (1 + j^{(1)}(r)) dv dr - 2r^{\frac{(p-3)(p-6)+q(7-p)}{2(p-q)}} \delta\beta_i dx^i dr \quad (3.1)
 \end{aligned}$$

All the perturbations depend only on r after derivative expansion, which coincides with the fact that the fluid within the patch associated with that point is in local equilibrium. The above first order expanded metric depends on x^μ linearly through $g_{vv}^{(1)}$, $g_{vi}^{(1)}$ and $g_{ir}^{(1)}$, which contains $\delta r_H = x^\mu \partial_\mu r_H$ and $\delta\beta_i = x^\mu \partial_\mu \beta_i$. The dependence on r is much more complicated than on x^μ since the perturbations will turn out to be complicate functions of r . The perturbations will be solved in groups of how they transform under $SO(p-q)$. All the components of the traceless tensor and vector perturbations are on the same footing because of the $SO(p-q)$ invariance.

The tensor perturbation satisfies the traceless tensor part of the Einstein equation

$$E_{ij} - \frac{1}{p-q} \delta_{ij} \delta^{kl} E_{kl} - \left(T_{ij} - \frac{1}{p-q} \delta_{ij} \delta^{kl} T_{kl} \right) = 0, \quad (3.2)$$

After we put eq. (3.1) in, it gives

$$\partial_r (r^{8-p} f(r) \partial_r \alpha_{ij}^{(1)}(r)) + (9-p)r^{\frac{7-p}{2}} \sigma_{ij} = 0 \quad (3.3)$$

which does not depend on q . The reason is that the compactification on world-volume directions will decrease the spatial dimensions of Dp-brane. This will cause the tensor and vector perturbations losing some of their components. But compactification will not change the differential equations of the tensor and vector perturbations since every component of them satisfies the same differential equation because of the isotropy in all spatial directions. The existing components of the tensor and vector perturbations satisfy the same differential equations as the reduced ones. Compactification will only affects the spatial trace of the bulk metric since taking the spatial trace is summing over all spatial dimensions. Considering α_{ij} is the traceless part of the metric perturbations and does not relate to the sum of spatial components, it should be no wonder that eq. (3.3) does not depend on q .

The $SO(p-q)$ invariance of the bulk metric allows us to set $\alpha_{ij}^{(1)} = F(r)\sigma_{ij}$, then

$$(r^{8-p} f F')' + (9-p)r^{\frac{7-p}{2}} = 0. \quad (3.4)$$

The solution of the above has been solved in [26] as

$$\begin{aligned}
 F(r) = & \frac{1}{3r_H^{1/2}} \left[2\sqrt{3} \arctan \frac{\sqrt{3rr_H}}{r-r_H} \right. \\
 & \left. + \ln \frac{(\sqrt{r} + \sqrt{r_H})^4 (r + \sqrt{rr_H} + r_H)^2 (r^2 + rr_H + r_H^2)}{r^6} \right], \quad (\text{D4-brane})
 \end{aligned}$$

$$\begin{aligned}
 F(r) &= \frac{1}{2r_H} \left[2 \arctan \frac{r_H}{r} + \ln \frac{(r+r_H)^2(r^2+r_H^2)}{r^4} \right], & \text{(D3-brane)} \\
 F(r) &= \frac{2}{5r_H^{3/2}} \left[2 \sin \frac{2\pi}{5} \arctan \frac{2 \sin \frac{\pi}{5} \sqrt{rr_H}}{r-r_H} - 2 \sin \frac{\pi}{5} \arctan \frac{2 \sin \frac{2\pi}{5} \sqrt{rr_H}}{r-r_H} \right. \\
 &\quad - 2 \cos \frac{2\pi}{5} \operatorname{artanh} \frac{2 \cos \frac{\pi}{5} \sqrt{rr_H}}{r+r_H} - 2 \cos \frac{\pi}{5} \operatorname{artanh} \frac{2 \cos \frac{2\pi}{5} \sqrt{rr_H}}{r+r_H} \\
 &\quad \left. + \ln \frac{(\sqrt{r} + \sqrt{r_H})^2(r^4 + r^3 r_H + r^2 r_H^2 + r r_H^3 + r_H^4)}{r^5} \right]. & \text{(D2-brane)} \quad (3.5)
 \end{aligned}$$

There are two components of the Einstein equation can be used to solve the vector perturbation: the $(0i)$ and (ri) components. In the original framework, (ri) component of Einstein equation

$$E_{ri} - T_{ri} = 0 \quad (3.6)$$

is used to solve $w_i^{(1)}$ as

$$w_i^{(1)}(r) = a(r) \partial_0 \beta_i, \quad a(r) = -\frac{2}{(5-p)r^{\frac{5-p}{2}}} \quad (3.7)$$

The linear combination of the $(0i)$ and (ri) components of Einstein equation

$$g^{r0}(E_{0i} - T_{0i}) + g^{rr}(E_{ri} - T_{ri}) = 0, \quad (3.8)$$

is used to give a constraint relation between the two first order vector viscous terms as

$$\frac{1}{r_H} \partial_i r_H = -\frac{2}{5-p} \partial_0 \beta_i, \quad (3.9)$$

Both the first order vector perturbation eq. (3.7) and constraint equation eq. (3.9) are not q -dependent. The reason for the vector perturbation has been explained below eq. (3.3). The vector constraint equation eq. (3.9) is a linear combination of the vector components of Einstein equation, which should also not be q -dependent following similar argument as what is below eq. (3.3).

To solve the scalar perturbations, we have (00) , $(0r)$, (rr) and (ii) (with i summed) components of Einstein equation, as well as the EOM of ϕ . But there are only 3 unknown scalar functions, thus two of the five equations are redundant. The (ii) (with i summed) component turns out to be more complex than the others thus can be omitted. A linear combination of the (00) and $(0r)$ components of Einstein equation

$$g^{r0}(E_{00} - T_{00}) + g^{rr}(E_{r0} - T_{r0}) = 0 \quad (3.10)$$

does not contain any scalar perturbations and it is called the first scalar constraint equation. To put eq. (3.1) into the above one has

$$\frac{1}{r_H} \partial_0 r_H = -\frac{2}{9-p} \partial \beta. \quad (3.11)$$

Since the above comes from the linear combination of (00) and (0r) components of Einstein equation. Both of them are independent of the spatial trace part of the metric. It is reasonable that eq. (3.11) contains no q .

A linear combination of (0r) and (rr) components of Einstein equation (which is also called the second scalar constraint equation)

$$g^{rr}(E_{rr} - T_{rr}) + g^{r0}(E_{r0} - T_{r0}) = 0, \quad (3.12)$$

the (rr) component itself

$$E_{rr} - T_{rr} = 0, \quad (3.13)$$

and the EOM of ϕ eq. (2.19) are chosen to solve the 3 scalar perturbations, whose first order differential equations are

$$(r^{7-p}k_{(1)})' + 2(7-p)r^{6-p}j_{(1)} - \left[(p-q)r^{7-p} - \frac{2(p-q)}{9-p}r_H^{7-p} \right] h'_{(1)} - 2r^{\frac{7-p}{2}}\partial\beta = 0, \quad (3.14)$$

$$2(p-q)rh''_{(1)} + (7-p)(p-q)h'_{(1)} - 2(9-p)j'_{(1)} = 0, \quad (3.15)$$

$$2(r^{7-p}k_{(1)})' + 2r^{7-p}fj'_{(1)} + 4(7-p)r^{6-p}j_{(1)} - (p-q)r^{7-p}fh'_{(1)} - 2r^{\frac{7-p}{2}}\partial\beta = 0. \quad (3.16)$$

Again, SO($p-q$) invariance allows us to set $\chi^{(1)} = F_\chi\partial\beta$ with $\chi = \{h, j, k\}$. The viscous term $\partial\beta$ can be removed:

$$(r^{7-p}F_k)' + 2(7-p)r^{6-p}F_j - \left[(p-q)r^{7-p} - \frac{2(p-q)}{9-p}r_H^{7-p} \right] F'_h - 2r^{\frac{7-p}{2}} = 0,$$

$$2(p-q)rF''_h + (7-p)(p-q)F'_h - 2(9-p)F'_j = 0,$$

$$2(r^{7-p}F_k)' + 2r^{7-p}fF'_j + 4(7-p)r^{6-p}F_j - (p-q)r^{7-p}fF'_h - 2r^{\frac{7-p}{2}} = 0. \quad (3.17)$$

The first order solutions for the 3 scalar perturbations are

$$F_h = \frac{1}{p-q}F, \quad F_j = -\frac{2}{9-p} \frac{r^{\frac{9-p}{2}} - r_H^{\frac{9-p}{2}}}{r^{7-p} - r_H^{7-p}} + \frac{5-p}{2(9-p)}F,$$

$$F_k = \frac{4}{(9-p)r^{\frac{5-p}{2}}} - \frac{1}{9-p} \left(5-p + \frac{2r_H^{7-p}}{r^{7-p}} \right) F. \quad (3.18)$$

From the solution one can see that among the 3 first order scalar perturbations, only h is affected by compactification. This is because h is the spatial trace part of the metric perturbations. It is sensitive to the total number of the spatial dimensions and should be under affect of the compactification. From the view of boundary fluid, $h^{(1)}$ directly relates to the spatial trace part of $\partial_\mu u_\nu$, i.e. $\partial \cdot u$. Thus it is reasonable for $h^{(1)}$ to be dependent on the total number of the spatial dimensions. This is the reason that only h is affected by the compactification among the 3 scalar perturbations. The factor $p-q$ in the denominator

of the expression of F_h is exactly the sign that the number of spatial directions of the compactified Dp-brane is now $p - q$ but not p as in [26].

The differential equations of F_j and F_k do not depend on q . This can be seen by substituting $(p - q)F_h$ for F in eq. (3.17), then $(p - q)$ disappears from all the differential equations in eq. (3.17). Then one can get the expressions for F_k and F_j in terms of F . In such way we do not need to solve F_j and F_k out explicitly, which is much more difficult than to find the relations directly as in eq. (3.18). The q -independence of k and j may also be figured out in the way that they are the (00) and (0r) components of the perturbative metric, respectively. Both of them have nothing to do with the number of the spatial dimensions. Compactification on spatial directions should not affect the scalar perturbations k and j .

3.2 Constitutive relation on the boundary

The constitutive relation of the fluid living on the boundary can be got from taking large r limit for the Brown-York tensor of the bulk theory:

$$T_{\mu\nu} = \frac{1}{\kappa_{p-q+2}^2} \lim_{r \rightarrow \infty} \left(\frac{r}{L_p} \right)^{\frac{(9-p)(p-q-1)}{2(p-q)}} \left(K_{\mu\nu} - h_{\mu\nu} K - \frac{9-p}{2L_p} \left(\frac{r}{L_p} \right)^{-\frac{(p-3)^2+q(5-p)}{2(p-q)}} h_{\mu\nu} \right), \quad (3.19)$$

The third term in the parenthesis of the above is from the counter term of the action. Note that the difference of the two powers of the factor $\frac{r}{L_p}$ is always 2 for all Dp-brane with or without compactification. Here

$$K_{\mu\nu} = -\frac{1}{2} (n^\rho \partial_\rho h_{\mu\nu} + \partial_\mu n^\rho h_{\rho\nu} + \partial_\nu n^\rho h_{\rho\mu}), \quad (3.20)$$

and $K = h^{\mu\nu} K_{\mu\nu}$. From the factor $\left(\frac{r}{L_p} \right)^{\frac{(9-p)(p-q-1)}{2(p-q)}}$ in eq. (3.19) we can see that compactification of the bulk gravity will change the asymptotic behavior of the Brown-York tensor in the framework of fluid/gravity correspondence.

Putting eq. (3.1) into eq. (3.19) with all solved first order perturbations we then have the first order constitutive relation of the boundary fluid as

$$T_{\mu\nu} = \frac{1}{2\kappa_{p-q+2}^2} \left[\frac{r_H^{7-p}}{L_p^{8-p}} \left(\frac{9-p}{2} u_\mu u_\nu + \frac{5-p}{2} P_{\mu\nu} \right) - \left(\frac{r_H}{L_p} \right)^{\frac{9-p}{2}} \left(2\sigma_{\mu\nu} + \frac{2(p-3)^2 + 2q(5-p)}{(p-q)(9-p)} P_{\mu\nu} \partial u \right) \right]. \quad (3.21)$$

Then thermal and transport coefficients of first order can be got as

$$\begin{aligned} \varepsilon &= \frac{1}{2\kappa_{p-q+2}^2} \frac{9-p}{2} \frac{r_H^{7-p}}{L_p^{8-p}}, & \mathbf{p} &= \frac{1}{2\kappa_{p-q+2}^2} \frac{5-p}{2} \frac{r_H^{7-p}}{L_p^{8-p}}, \\ \eta &= \frac{1}{2\kappa_{p-q+2}^2} \left(\frac{r_H}{L_p} \right)^{\frac{9-p}{2}}, & \zeta &= \frac{1}{2\kappa_{p-q+2}^2} \frac{2(p-3)^2 + 2q(5-p)}{(p-q)(9-p)} \left(\frac{r_H}{L_p} \right)^{\frac{9-p}{2}}. \end{aligned} \quad (3.22)$$

Scalars of $SO(p-q)$	Vectors of $SO(p-q)$	Tensors of $SO(p-q)$
$\mathbf{s}_1 = \frac{1}{r_H} \partial_0^2 r_H$	$\mathbf{v}_{1i} = \frac{1}{r_H} \partial_0 \partial_i r_H$	$\mathbf{t}_{1ij} = \frac{1}{r_H} \partial_i \partial_j r_H - \frac{1}{p-q} \delta_{ij} \mathbf{s}_3$
$\mathbf{s}_2 = \partial_0 \partial_i \beta_i$	$\mathbf{v}_{2i} = \partial_0^2 \beta_i$	$\mathbf{t}_{2ij} = \partial_0 \Omega_{ij}$
$\mathbf{s}_3 = \frac{1}{r_H} \partial_i^2 r_H$	$\mathbf{v}_{3i} = \partial_j^2 \beta_i$	$\mathbf{t}_{3ij} = \partial_0 \sigma_{ij}$
$\mathfrak{S}_1 = \partial_0 \beta_i \partial_0 \beta_i$	$\mathbf{v}_{4i} = \partial_j \Omega_{ij}$	$\mathfrak{T}_{1ij} = \partial_0 \beta_i \partial_0 \beta_j - \frac{1}{p-q} \delta_{ij} \mathfrak{S}_1$
$\mathfrak{S}_3 = (\partial_i \beta_i)^2$	$\mathbf{v}_{5i} = \partial_j \sigma_{ij}$	$\mathfrak{T}_{2ij} = \sigma_{[i}{}^k \Omega_{j]k}$
$\mathfrak{S}_4 = \Omega_{ij} \Omega_{ij}$	$\mathfrak{V}_{1i} = \partial_0 \beta_i \partial \beta$	$\mathfrak{T}_{3ij} = \Omega_{ij} \partial \beta$
$\mathfrak{S}_5 = \sigma_{ij} \sigma_{ij}$	$\mathfrak{V}_{2i} = \partial_0 \beta_j \Omega_{ij}$	$\mathfrak{T}_{4ij} = \sigma_{ij} \partial \beta$
	$\mathfrak{V}_{3i} = \partial_0 \beta_j \sigma_{ij}$	$\mathfrak{T}_{5ij} = \Omega_i{}^k \Omega_{jk} - \frac{1}{p-q} \delta_{ij} \mathfrak{S}_4$
		$\mathfrak{T}_{6ij} = \sigma_i{}^k \sigma_{jk} - \frac{1}{p-q} \delta_{ij} \mathfrak{S}_5$
		$\mathfrak{T}_{7ij} = \sigma_{(i}{}^k \Omega_{j)k}$

Table 1. The list of $SO(p-q)$ invariant second order viscous terms for $1 \leq p-q \leq 3$. $\mathfrak{T}_{2,5,6}$ do not exist when $p-q=2$. At $p-q=1$, only $\mathbf{s}_{1,2,3}$, $\mathfrak{S}_{1,3}$, $\mathbf{v}_{1,2,3}$ and \mathfrak{V}_1 exist.

In the above, only the bulk viscosity depends on the number of compact directions q since it relates with the spatial trace of the metric perturbation for which only $h^{(1)} = F_h \partial \beta$ contributes.

From these results, one can see that the ratio ζ/η that we get via fluid/gravity correspondence is exactly what is proposed in [13] as quoted in eq. (1.1). Thus the expression for the ratio ζ/η given in [13] is indeed for compactified Dp-brane. The Hawking temperature T for compactified Dp-brane is the same as Dp-brane. So the entropy density $s = (\varepsilon+p)/T$ as well as the ratio η/s are also the same as that of Dp-brane case.

4 Preliminaries on second order calculation

4.1 The second order constraints and Navier-Stokes equations

In the first order calculation we have two equations that without any perturbation, which are actually the Navier-Stokes equations of the boundary fluid at the first order. But from the gravity viewpoint they are the first order scalar eq. (3.11) and vector eq. (3.9) constraint. They relate the following 5 first order partial derivative terms of the collective modes of different type:

$$\partial_0 r_H, \quad \partial \beta, \quad \partial_i r_H, \quad \partial_0 \beta_i, \quad \sigma_{ij} = \partial_{(i} \beta_{j)} - \frac{1}{p-q} \delta_{ij} \partial \beta. \quad (4.1)$$

Thus only three of the above are actually independent. We choose $\partial \beta$, $\partial_0 \beta_i$ and σ_{ij} as the first order scalar, vector and tensor viscous term, respectively. That's why the first order perturbations can be written as $\chi^{(1)} = F_\chi \partial \beta$, $w_i^{(1)} = a(r) \partial_0 \beta_i$ and $\alpha_{ij}^{(1)} = F(r) \sigma_{ij}$.

The second order solving procedure follows the same rule. But now we have much more viscous terms since its complexity grows nonlinearly. These second order viscous

terms satisfy 6 identities which can be extract from $\partial_\mu \partial^\rho T_{\rho\nu}^{(0)} = 0$. $T_{\mu\nu}^{(0)}$ here is the thermodynamical part of eq. (3.21) with r_H and u^μ depend on x^μ . After we expand $T_{\mu\nu}^{(0)}$ and take different components of $\partial_\mu \partial^\rho T_{\rho\nu}^{(0)} = 0$, it gives

$$\frac{9-p}{2} \frac{1}{r_H} \partial_0^2 r_H + \partial_0 \partial \beta - \frac{2}{9-p} (\partial \beta)^2 - \frac{4}{5-p} \partial_0 \beta_i \partial_0 \beta_i = 0, \quad (4.2)$$

$$\frac{5-p}{2} \frac{1}{r_H} \partial_i^2 r_H + \partial_0 \partial \beta - \frac{2}{5-p} \partial_0 \beta_i \partial_0 \beta_i - \frac{5-p}{9-p} (\partial \beta)^2 + \partial_i \beta_j \partial_j \beta_i = 0, \quad (4.3)$$

$$\frac{5-p}{2} \frac{1}{r_H} \partial_0 \partial_i r_H + \partial_0^2 \beta_i - \frac{7-p}{9-p} \partial_0 \beta_i \partial \beta + \partial_0 \beta_j \partial_j \beta_i = 0, \quad (4.4)$$

$$\frac{9-p}{2} \frac{1}{r_H} \partial_0 \partial_i r_H + \partial_i \partial \beta - \frac{2}{5-p} \partial_0 \beta_i \partial \beta - \frac{4}{5-p} \partial_0 \beta_j \partial_i \beta_j = 0, \quad (4.5)$$

$$\partial_0 \Omega_{ij} - \frac{5-p}{9-p} \Omega_{ij} \partial \beta - \partial_k \beta_{[i} \partial_j] \beta_k = 0, \quad (4.6)$$

$$\frac{5-p}{2} \frac{1}{r_H} \partial_i \partial_j r_H + \partial_0 \partial_{(i} \beta_{j)} - \frac{2}{5-p} \partial_0 \beta_i \partial_0 \beta_j - \frac{5-p}{9-p} \partial_{(i} \beta_{j)} \partial \beta + \partial_k \beta_{(i} \partial_j) \beta_k = 0. \quad (4.7)$$

These are the same as the cases of Dp-brane without compact direction. But after reexpressed in terms of the spatial viscous tensors in table 1, they actually depend on q :

$$\mathbf{s}_1 + \frac{2}{9-p} \mathbf{s}_2 - \frac{8}{(9-p)(5-p)} \mathfrak{S}_1 - \frac{4}{(9-p)^2} \mathfrak{S}_3 = 0, \quad (4.8)$$

$$\mathbf{s}_2 + \frac{5-p}{2} \mathbf{s}_3 - \frac{2}{5-p} \mathfrak{S}_1 + \frac{(p-3)^2 + q(5-p)}{(p-q)(9-p)} \mathfrak{S}_3 - \mathfrak{S}_4 + \mathfrak{S}_5 = 0, \quad (4.9)$$

$$\mathbf{v}_1 + \frac{2}{5-p} \mathbf{v}_2 + \frac{2(p^2 - 8p + 9 + q(7-p))}{(p-q)(9-p)(5-p)} \mathfrak{V}_1 - \frac{2}{5-p} \mathfrak{V}_2 + \frac{2}{5-p} \mathfrak{V}_3 = 0, \quad (4.10)$$

$$\mathbf{v}_1 + \frac{2(p-q)(\mathbf{v}_4 + \mathbf{v}_5)}{(p-q-1)(9-p)} - \frac{4(p-q+2)}{(p-q)(9-p)(5-p)} \mathfrak{V}_1 - \frac{8(\mathfrak{V}_2 + \mathfrak{V}_3)}{(9-p)(5-p)} = 0, \quad (4.11)$$

$$\mathbf{t}_2 - 2\mathfrak{T}_2 + \frac{p^2 - 7p + 18 + q(5-p)}{(p-q)(9-p)} \mathfrak{T}_3 = 0, \quad (4.12)$$

$$\mathbf{t}_1 + \frac{2}{5-p} \mathbf{t}_3 - \frac{4}{(5-p)^2} \mathfrak{T}_1 + \frac{2(p^2 - 7p + 18) + 2q(5-p)}{(p-q)(9-p)(5-p)} \mathfrak{T}_4 - \frac{2}{5-p} \mathfrak{T}_5 + \frac{2}{5-p} \mathfrak{T}_6 = 0. \quad (4.13)$$

The above 6 identities are used in deriving the differential equations for the second order perturbations as well as the second order Navier-Stokes equations. We will derive the latter in this section from a hydrodynamical viewpoint. They can be gained by expanding eq. (3.21) to the second order of partial derivatives and putting it into $\partial^\mu T_{\mu\nu}^{(0+1)} = 0$, where we add the superscript $(0+1)$ to emphasize the difference with $T_{\mu\nu}^{(0)}$. The second order Navier-Stokes equations are

$$\begin{aligned} \frac{1}{r_H^{(p-3)/2}} \partial_0 r_H^{(1)} &= \frac{4(p-3)^2 + 4q(5-p)}{(p-q)(9-p)^2(7-p)} \mathfrak{S}_3 + \frac{4}{(9-p)(7-p)} \mathfrak{S}_5, \\ \frac{1}{r_H^{(p-3)/2}} \partial_i r_H^{(1)} &= \frac{[4(p-3)^2 + 4q(5-p)] \mathbf{v}_4 + 16(p-q) \mathbf{v}_5}{(p-q-1)(9-p)(7-p)(5-p)} \end{aligned} \quad (4.14)$$

$$\begin{aligned}
& + \frac{2(p-1)(p^2 - 22p + 77) - 2q(p^2 - 22p + 85)}{(p-q)(9-p)(7-p)(5-p)^2} \mathfrak{A}_1 \\
& - \frac{2(19-3p)}{(9-p)(7-p)(5-p)} \mathfrak{A}_2 - \frac{2(p^2 - 14p + 77)}{(9-p)(7-p)(5-p)^2} \mathfrak{A}_3. \tag{4.15}
\end{aligned}$$

They are more complex than the Dp-brane case, when set $q = 0$, they are back to the form of Dp-brane. The above second order Navier-Stokes equations can also be got through the constraints of Einstein equation eqs. (3.8) and (3.10), which will be specified later when we solve the second order perturbations.

4.2 The second order expanded metric

The calculation of the second order needs first to expand eq. (2.17) to the second order of partial derivatives of r_H and β_i , and the result is

$$\begin{aligned}
ds^2 = & -r^{\frac{9-p}{p-q}} \left[f - (1-f)\delta\beta_i\delta\beta_i - \frac{(7-p)r_H^{6-p}}{r^{7-p}}\delta r_H - \frac{(7-p)r_H^{6-p}}{2r^{7-p}}\delta^2 r_H \right. \\
& - \frac{(7-p)r_H^{6-p}}{r^{7-p}}\delta r_H^{(1)} - \frac{(7-p)(6-p)r_H^{5-p}}{2r^{7-p}}(\delta r_H)^2 - (F_k + \delta F_k)\partial\beta \\
& - F_k(\delta\partial\beta + \delta\beta_i\partial_0\beta_i) - 2a(r)\delta\beta_i\partial_0\beta_i - k^{(2)}(r) \left. \right] dv^2 \\
& - 2r^{\frac{9-p}{p-q}} \left[(1-f) \left(\delta\beta_i + \frac{1}{2}\delta^2\beta_i \right) + a(\partial_0\beta_i + \delta\partial_0\beta_i + \delta\beta_j\partial_j\beta_i) \right. \\
& + \frac{(7-p)r_H^{6-p}}{r^{7-p}}\delta r_H\delta\beta_i + F_k\partial\beta\delta\beta_i + F\delta\beta_j\partial_{(i}\beta_{j)} + w_i^{(2)}(r) \left. \right] dv dx^i \\
& + 2r^{\frac{(p-3)(p-6)+q(7-p)}{2(p-q)}} \left[1 + (F_j + \delta F_j)\partial\beta + F_j(\delta\partial\beta + \delta\beta_i\partial_0\beta_i) \right. \\
& + \frac{1}{2}\delta\beta_i\delta\beta_i + j^{(2)}(r) \left. \right] dv dr \\
& + r^{\frac{9-p}{p-q}} \left[\delta_{ij} + (1-f)\delta\beta_i\delta\beta_j + 2a\delta\beta_{(i}\partial_{|0|}\beta_{j)} + (F + \delta F)\partial_{(i}\beta_{j)} \right. \\
& + F \left(\delta\partial_{(i}\beta_{j)} + \delta\beta_{(i}\partial_{|0|}\beta_{j)} \right) + \alpha_{ij}^{(2)}(r) + h^{(2)}(r)\delta_{ij} \left. \right] dx^i dx^j \\
& - 2r^{\frac{(p-3)(p-6)+q(7-p)}{2(p-q)}} \left(\delta\beta_i + \frac{1}{2}\delta^2\beta_i + F_j\partial\beta\delta\beta_i \right) dx^i dr. \tag{4.16}
\end{aligned}$$

Here we have $\delta^2\psi = x^\mu x^\nu \partial_\mu \partial_\nu \psi$ with ψ either r_H or β_i , while

$$\delta\mathcal{F}(r_H(x), r) = -\frac{(5-p)\mathcal{F}(r) + 2r\mathcal{F}'(r)}{2r_H} \delta r_H \tag{4.17}$$

with \mathcal{F} standing for any of the F , F_j and F_k . The above does not change compared with the Dp-brane case since F , F_j and F_k do not change.

Each component of the above second order expanded metric has the structure of $g_{MN} = r^{(\dots)}[\dots]$, the expressions inside the brackets $[\dots]$ do not change while terms inside the parenthesis (\dots) have changed compared with the Dp-brane case. The components of the second order expanded metric are now the second order polynomials of the boundary coordinates x^μ through terms like $\delta^2\psi$ or $(\delta\psi)^2$ with ψ being r_H or β_i .

5 Solving the second order perturbations

In the following subsections we will solve the second order perturbations for all the compactified Dp-branes with $p = 2, 3, 4$ and $1 \leq q \leq p-1$. The cases can be divided into three classes, in terms of $p - q$, they are $p - q = 1, 2, 3$. Since $p - q = 3$ is just the D4-brane with 1 direction compactified. We refer the reader to [24] for details and will omit the solving procedure in this work despite the conventions used here is a little different from that of [24].

5.1 The tensor part

We will solve the second order tensor perturbations for all compactified Dp-brane in this subsection. Since the D(4-1)-brane case has been solved in [24], the cases contained here are only $p - q = 2$ which include D(4-2) and D(3-1)-brane. The situation of $p - q = 1$ should not be considered because the spatial viscous tensors can only be constructed in spacetime with more than 1 spatial dimensions.

The differential equation of the tensor perturbation is

$$\frac{d}{dr} \left(r^{8-p} f(r) \frac{d\alpha_{ij}^{(2)}}{dr} \right) = S_{ij}(r), \quad (5.1)$$

which is an exact second order linear inhomogeneous differential equation. The inhomogeneous term $S_{ij}(r)$ is called the source term in [1]. This equation is singular at $r = 0$, $r = r_H$ and $r \rightarrow \infty$. Its solution can be got by integration twice as

$$\alpha_{ij}^{(2)}(r) = \int_r^\infty \frac{-1}{x^{8-p} f} dx \int_1^x S_{ij}(y) dy, \quad (5.2)$$

We will solve the above equation by specifying p, q with explicit values. When putting the second order expanded metric eq. (4.16) into eq. (3.2) for D4-brane with 2 compact directions, one gets

$$\begin{aligned} \partial_r(r^4 f \partial_r \alpha_{ij}^{(2)}) = & \left(6r - \frac{5}{2} r^{\frac{3}{2}} F - 2r^{\frac{5}{2}} F' \right) (t_3 + \mathfrak{T}_1) + \left[6r - \frac{1}{2} r^{\frac{3}{2}} F - \frac{21}{5} r^{\frac{5}{2}} F' \right. \\ & - \frac{4}{5} r^{\frac{7}{2}} F'' + \frac{1}{2} r^4 f F'^2 + (4r^3 - 1) F F' + r^4 f F F'' + 5r^{\frac{3}{2}} F_j - 15r^2 F F_j \\ & + 2(4r^3 - 1) F' F_j - (5r^3 - 2) F F'_j + r^4 f F' F'_j + 2r^4 f F'' F_j \\ & \left. - \frac{15}{2} r^2 F F_k + 4r^3 F' F_k - \frac{13}{2} r^3 F F'_k + r^4 F' F'_k + r^4 F'' F_k - r^4 F F''_k \right] \mathfrak{T}_4 \\ & - \left(4r + 5r^{\frac{3}{2}} F + 4r^{\frac{5}{2}} F' \right) \mathfrak{T}_7, \end{aligned} \quad (5.3)$$

The source term on the right hand side of the above equation has the same structure as D2-brane. The differences lie only in the coefficient functions of each spatial viscous tensors. Using eq. (5.2), one has

$$\begin{aligned} \alpha_{ij}^{(2)} = & \left(\frac{4}{3} - \frac{\pi}{9\sqrt{3}} - \frac{\ln 3}{3} \right) \frac{1}{r^3} (t_3 + \mathfrak{T}_1) + \left[\frac{4}{r} + \left(\frac{16}{15} - \frac{\pi}{45\sqrt{3}} - \frac{\ln 3}{15} \right) \frac{1}{r^3} \right] \mathfrak{T}_4 \\ & + \left[\frac{8}{r} - \left(\frac{2\pi}{9\sqrt{3}} + \frac{2\ln 3}{3} \right) \frac{1}{r^3} \right] \mathfrak{T}_7 \end{aligned} \quad (5.4)$$

This solution is in a form of an asymptotic series of r . It is got by making a series expansion in between the two sequential integrations for y and x , as has been specified in [24]. Expanding the source term at the beginning will lead to wrong results for $\alpha_{ij}^{(2)}$ and $h^{(2)}$. But this will not affect the integration for $w_i^{(2)}$, $j^{(2)}$ and $k^{(2)}$. We think the reason is that both the integrations for $\alpha_{ij}^{(2)}$ and $h^{(2)}$ depend on the emblackening factor $f(r)$ while $w_i^{(2)}$ do not. As for $j^{(2)}$ and $k^{(2)}$, they are calculated by integration only once. This exchangeability between the integral of y and the asymptotic expansion will give us great help for the calculations.

Though near extremal D3-brane can be reduced trivially to AdS₅ black hole whose dual fluid is conformal. After compactification, it will correspond to nonconformal fluid. The tensor perturbation for D(3-1)-brane satisfies the differential equation

$$\begin{aligned} \partial_r(r^5 f \partial_r \alpha_{ij}^{(2)}) = & (2r - 3r^2 F - 2r^3 F') (\mathfrak{t}_3 + \mathfrak{T}_1) + \left[2r - r^2 F - \frac{13}{3} r^3 F' - \frac{2}{3} r^4 F'' \right. \\ & + (5r^4 - 1) F F' + r^5 f F F'' + \frac{1}{2} r^5 f F'^2 + 6r^2 F_j - 24r^3 F F_j \\ & + 2(5r^4 - 1) F' F_j - 2(3r^4 - 1) F F'_j + r^5 f F' F'_j + 2r^5 f F'' F_j \\ & \left. - 12r^3 F F_k + 5r^4 F' F_k - 8r^4 F F'_k + r^5 F' F'_k + r^5 F'' F_k - r^5 F F''_k \right] \mathfrak{T}_4 \\ & - (4r + 6r^2 F + 4r^3 F') \mathfrak{T}_7, \end{aligned} \tag{5.5}$$

using eq. (5.2) one gets the solution as

$$\alpha_{ij}^{(2)} = \left(\frac{1}{2} - \frac{\ln 2}{4} \right) \frac{1}{r^4} (\mathfrak{t}_3 + \mathfrak{T}_1) + \left[\frac{1}{r^2} + \left(\frac{1}{3} - \frac{\ln 2}{12} \right) \frac{1}{r^4} \right] \mathfrak{T}_4 + \left(\frac{2}{r^2} - \frac{\ln 2}{2r^4} \right) \mathfrak{T}_7. \tag{5.6}$$

Compared with D3-brane case where the coefficient function of $\mathfrak{t}_3 + \mathfrak{T}_1$ is proportional to that of \mathfrak{T}_4 , these two coefficient functions here are not proportional to each other, suggesting the background is not conformal now.

From the calculation of this subsection, the $1/r$ dependent behavior of the tensor perturbations of D(p-q)-brane does not change compared with that of Dp-brane. Perturbations of other types share the same feature as the tensor perturbation, as can be seen from the solutions of vector and scalar perturbations. This shows that the compactification will not affect the asymptotic behavior of the metric perturbations. Actually, the asymptotic behavior of the metric perturbations are mainly determined by the power of r in the emblackening factor $f(r)$ which is independent of the compact directions q .

5.2 The vector part

There are a constraint plus a dynamical equation for the vector part. Both the cases of $p - q = 2$ and 1 need to be analyzed here.

5.2.1 The constraint equation

The vector constraint equation can be got by feeding eq. (4.16) into eq. (3.8). The result is just the spatial component of the Navier-Stokes equation of second order, as one can check by substituting specific values for p and q into eq. (4.15).

Let's start with $p - q = 2$ first. For D(4-2)-brane, the vector constraint equation is

$$\begin{aligned}
\partial_i r_H^{(1)} = & \left[-\frac{4}{3}r^{\frac{5}{2}} - \frac{2}{3}r^4 f F' + \frac{4}{3}(5r^3 - 2)F_j + \frac{10}{3}r^3 F_k + \frac{4}{3}r^4 F'_k \right] \mathbf{v}_4 \\
& + \left[\frac{4}{3}(5r^3 - 2)F_j + \frac{10}{3}r^3 F_k + \frac{4}{3}r^4 F'_k \right] \mathbf{v}_5 \\
& + \left[\frac{13}{3}r^{\frac{5}{2}} + \frac{1}{6}r(13r^3 - 1)F' - (15r^3 + 2)F_j - \frac{15}{2}r^3 F_k - \frac{7}{3}r^4 F'_k \right] \mathfrak{Y}_1 \\
& + \left[-2r^{\frac{5}{2}} - r^4 f F' + \frac{2}{3}(5r^3 - 2)F_j + \frac{5}{3}r^3 F_k + \frac{2}{3}r^4 F'_k \right] \mathfrak{Y}_2 \\
& + \left[-6r^{\frac{5}{2}} - 3r^4 f F' + \frac{2}{3}(5r^3 - 2)F_j + \frac{5}{3}r^3 F_k + \frac{2}{3}r^4 F'_k \right] \mathfrak{Y}_3.
\end{aligned} \tag{5.7}$$

After expansion with respect to $1/r$, we have

$$\partial_i r_H^{(1)} = \frac{4}{5}\mathbf{v}_4 + \frac{32}{15}\mathbf{v}_5 - \frac{11}{15}\mathfrak{Y}_1 - \frac{14}{15}\mathfrak{Y}_2 - \frac{74}{15}\mathfrak{Y}_3 \tag{5.8}$$

By the same token, one has the vector constraint for D(3-1)-brane as

$$\begin{aligned}
\partial_i r_H^{(1)} = & \left[-\frac{1}{2}r^3 - \frac{1}{4}r^5 f F' + (3r^4 - 1)F_j + \frac{3}{2}r^4 F_k + \frac{1}{2}r^5 F'_k \right] \mathbf{v}_4 \\
& + \left[(3r^4 - 1)F_j + \frac{3}{2}r^4 F_k + \frac{1}{2}r^5 F'_k \right] \mathbf{v}_5 \\
& + \left[\frac{9}{8}r^3 + \frac{1}{16}r(9r^4 - 1)F' - \frac{3}{4}(5r^4 + 1)F_j - \frac{15}{8}r^4 F_k - \frac{3}{8}r^5 F'_k \right] \mathfrak{Y}_1 \\
& + \left[-\frac{3}{4}r^3 - \frac{3}{8}r^5 f F' + \frac{1}{2}(3r^4 - 1)F_j + \frac{3}{4}r^4 F_k + \frac{1}{4}r^5 F'_k \right] \mathfrak{Y}_2 \\
& + \left[-\frac{5}{4}r^3 - \frac{5}{8}r^5 f F' + \frac{1}{2}(3r^4 - 1)F_j + \frac{3}{4}r^4 F_k + \frac{1}{4}r^5 F'_k \right] \mathfrak{Y}_3.
\end{aligned} \tag{5.9}$$

After expansion, the above becomes

$$\partial_i r_H^{(1)} = \frac{1}{6}\mathbf{v}_4 + \frac{2}{3}\mathbf{v}_5 + \frac{1}{8}\mathfrak{Y}_1 - \frac{5}{12}\mathfrak{Y}_2 - \frac{11}{12}\mathfrak{Y}_3. \tag{5.10}$$

The above two cases have two spatial directions for the dual fluid. Next we consider situations for $p - q = 1$, the vector viscous term in this case are only $\mathbf{v}_{1,2,3}$ and \mathfrak{Y}_1 . Thus the equations will turn out to be much simple now. In the order of D(4-3), D(3-2) and D(2-1)-brane, the vector constraint equations can be read as

$$\begin{aligned}
\partial_i r_H^{(1)} = & \left[\frac{2}{3}(5r^3 - 2)F_j + \frac{5}{3}r^3 F_k + \frac{2}{3}r^4 F'_k \right] \mathbf{v}_3 + \left[\frac{4}{3}r^{\frac{5}{2}} + \frac{2}{3}r(r^3 + 2)F' \right. \\
& \left. - \frac{8}{3}(5r^3 + 1)F_j - \frac{20}{3}r^3 F_k - 2r^4 F'_k \right] \mathfrak{Y}_1,
\end{aligned} \tag{5.11}$$

$$\begin{aligned}
\partial_i r_H^{(1)} = & \left[\frac{1}{2}(3r^4 - 1)F_j + \frac{3}{4}r^4 F_k + \frac{1}{4}r^5 F'_k \right] \mathbf{v}_3 + \left[\frac{1}{2}r^3 + \frac{1}{4}r(r^4 + 1)F' \right. \\
& \left. - (3r^4 + 1)F_j - \frac{3}{2}r^4 F_k - \frac{1}{4}r^5 F'_k \right] \mathfrak{Y}_1,
\end{aligned} \tag{5.12}$$

$$\begin{aligned} \partial_i r_H^{(1)} = & \left[\frac{2}{15}(7r^5 - 2)F_j + \frac{7}{15}r^5 F_k + \frac{2}{15}r^6 F'_k \right] \mathbf{v}_3 + \left[\frac{4}{15}r^{\frac{7}{2}} + \frac{2}{45}r(3r^5 + 2)F' \right. \\ & \left. - \frac{8}{45}(7r^5 + 3)F_j - \frac{28}{45}r^5 F_k - \frac{2}{45}r^6 F'_k \right] \mathfrak{A}_1, \end{aligned} \quad (5.13)$$

which, after expansion with respect to $1/r$, lead to

$$\partial_i r_H^{(1)} = \frac{16}{15} \mathbf{v}_3 - \frac{16}{5} \mathfrak{A}_1, \quad \text{D(4-3)-brane} \quad (5.14)$$

$$\partial_i r_H^{(1)} = \frac{1}{3} \mathbf{v}_3 - \frac{1}{3} \mathfrak{A}_1, \quad \text{D(3-2)-brane} \quad (5.15)$$

$$\partial_i r_H^{(1)} = \frac{16}{105} \mathbf{v}_3 - \frac{16}{315} \mathfrak{A}_1. \quad \text{D(2-1)-brane} \quad (5.16)$$

5.2.2 The dynamical equation

The differential equation of $w_i^{(2)}$ can be got by plugging the second order expanded metric eq. (4.16) into eq. (3.6):

$$\frac{d}{dr} \left(r^{8-p} \frac{dw_i^{(2)}}{dr} \right) = S_i(r), \quad (5.17)$$

where S_i in the right hand side is the source term. This equation is similar like eq. (5.1), the only difference is that the above does not have the emblackening factor in the left hand side, thus no singularity at $r = r_H$. The above equation can be solved by direct integration as

$$w_i^{(2)} = \int_r^\infty \frac{1}{x^{8-p}} dx \int_x^\infty S_i(y) dy. \quad (5.18)$$

Let's solve the $p - q = 2$ cases first. For D(4-2)-brane, the differential equation of $w_i^{(2)}$ is

$$\begin{aligned} \partial_r(r^4 \partial_r w_i^{(2)}) = & \left(-2r - r^{\frac{5}{2}} F' + 5r^{\frac{3}{2}} F_j - 2r^{\frac{5}{2}} F'_j \right) \mathbf{v}_4 + \left(5r^{\frac{3}{2}} F_j - 2r^{\frac{5}{2}} F'_j \right) \mathbf{v}_5 \\ & + \left(\frac{1}{2}r - \frac{19}{4}r^{\frac{5}{2}} F' - 2r^{\frac{7}{2}} F'' + \frac{25}{4}r^{\frac{3}{2}} F_j + \frac{5}{2}r^{\frac{5}{2}} F'_j - 2r^{\frac{7}{2}} F''_j \right) \mathfrak{A}_1 \\ & + \left(5r - \frac{3}{2}r^{\frac{5}{2}} F' + \frac{5}{2}r^{\frac{3}{2}} F_j - r^{\frac{5}{2}} F'_j \right) \mathfrak{A}_2 \\ & + \left(-r - \frac{9}{2}r^{\frac{5}{2}} F' + \frac{5}{2}r^{\frac{3}{2}} F_j - r^{\frac{5}{2}} F'_j \right) \mathfrak{A}_3, \end{aligned} \quad (5.19)$$

of which the solution is very simple

$$w_i^{(2)} = -\frac{2}{r} \mathfrak{A}_1 - \frac{4}{r} \mathfrak{A}_2 - \frac{4}{r} \mathfrak{A}_3. \quad (5.20)$$

The vector dynamical equation for the second order perturbation of D(3-1)-brane is

$$\begin{aligned} \partial_r(r^5 \partial_r w_i^{(2)}) = & \left(-2r - r^3 F' + 6r^2 F_j - 2r^3 F'_j \right) \mathbf{v}_4 + \left(6r^2 F_j - 2r^3 F'_j \right) \mathbf{v}_5 \\ & + \left(\frac{1}{2}r - \frac{11}{4}r^3 F' - r^4 F'' + \frac{15}{2}r^2 F_j + \frac{3}{2}r^3 F'_j - r^4 F''_j \right) \mathfrak{A}_1 \end{aligned}$$

$$\begin{aligned}
& + \left(r - \frac{3}{2}r^3 F' + 3r^2 F_j - r^3 F_j' \right) \mathfrak{Y}_2 \\
& + \left(-r - \frac{5}{2}r^3 F' + 3r^2 F_j - r^3 F_j' \right) \mathfrak{Y}_3,
\end{aligned} \tag{5.21}$$

with the solution is

$$w_i^{(2)} = -\frac{1}{2r^2} \mathfrak{Y}_1 - \frac{1}{r^2} \mathfrak{Y}_2 - \frac{1}{r^2} \mathfrak{Y}_3. \tag{5.22}$$

For the cases of $p - q = 1$, the differential equation for $w_i^{(2)}$ becomes simple again. In the order of D(4-3), D(3-2) and D(2-1)-brane, they are

$$\begin{aligned}
\partial_r(r^4 \partial_r w_i^{(2)}) & = \left(\frac{5}{2}r^{\frac{3}{2}} F_j - r^{\frac{5}{2}} F_j' \right) \mathfrak{v}_3 \\
& + \left(-7r^{\frac{5}{2}} F' - 2r^{\frac{7}{2}} F'' + \frac{15}{2}r^{\frac{3}{2}} F_j + 2r^{\frac{5}{2}} F_j' - 2r^{\frac{7}{2}} F_j'' \right) \mathfrak{Y}_1,
\end{aligned} \tag{5.23}$$

$$\begin{aligned}
\partial_r(r^5 \partial_r w_i^{(2)}) & = \left(3r^2 F_j - r^3 F_j' \right) \mathfrak{v}_3 \\
& + \left(-4r^3 F' - r^4 F'' + 9r^2 F_j + r^3 F_j' - r^4 F_j'' \right) \mathfrak{Y}_1,
\end{aligned} \tag{5.24}$$

$$\begin{aligned}
\partial_r(r^6 \partial_r w_i^{(2)}) & = \left(\frac{7}{2}r^{\frac{5}{2}} F_j - r^{\frac{7}{2}} F_j' \right) \mathfrak{v}_3 \\
& + \left(-3r^{\frac{7}{2}} F' - \frac{2}{3}r^{\frac{9}{2}} F'' + \frac{21}{2}r^{\frac{5}{2}} F_j + \frac{2}{3}r^{\frac{7}{2}} F_j' - \frac{2}{3}r^{\frac{9}{2}} F_j'' \right) \mathfrak{Y}_1.
\end{aligned} \tag{5.25}$$

The above 3 equations can also be identified by the power of r in the left hand side: this power is $8 - p$ for D(p-q)-brane. The solutions of the above are

$$w_i^{(2)} = -\frac{4}{r} \mathfrak{Y}_1, \quad \text{D(4-3)-brane} \tag{5.26}$$

$$w_i^{(2)} = -\frac{1}{r^2} \mathfrak{Y}_1, \quad \text{D(3-2)-brane} \tag{5.27}$$

$$w_i^{(2)} = -\frac{4}{9r^3} \mathfrak{Y}_1. \quad \text{D(2-1)-brane} \tag{5.28}$$

5.3 The scalar part

The scalar perturbations play an important role in the nonconformal background. They are also more difficult to solve than the tensor and vector part because the scalar perturbations are mixed together in the differential equations derived from scalar components of Einstein equation and the EOM of dilaton (or the scalar A).

The differential equation of $h^{(2)}$ for D(p-q)-brane is

$$\frac{d}{dr} \left(r^{8-p} f(r) \frac{dh^{(2)}}{dr} \right) = S_h(r), \tag{5.29}$$

whose solution can be written as

$$h^{(2)}(r) = \int_r^\infty \frac{-1}{x^{8-p} f} dx \int_1^x S_h(y) dy. \tag{5.30}$$

This is very like $\alpha_{ij}^{(2)}$, which can be seen from eqs. (5.1) and (5.2).

The differential equations for j and k are actually first order ones thus they can be solved by integrating once as

$$j^{(2)}(r) = - \int_r^\infty S_j(x) dx, \quad (5.31)$$

$$k^{(2)}(r) = - \frac{1}{r^{7-p}} \int_r^\infty S_k(x) dx. \quad (5.32)$$

$\frac{1}{r^{7-p}}$ is present in front of the integration. This is because k always appear as $(r^{7-p}k)$ in the differential equations.

5.3.1 The first scalar constraint

We still first deal with the cases of $p - q = 2$ in the scalar sector. The first scalar constraint does not contain any unknown perturbation and it can be calculated by plugging the second order expanded metric into eq. (3.10), which gives for the D(4-2)-brane case as

$$\begin{aligned} \partial_0 r_H^{(1)} = & \left(\frac{4}{5r^{\frac{1}{2}}} + \frac{4}{15}r^{\frac{5}{2}} - \frac{1}{5}F + \frac{2}{15}r^4 fF' - \frac{2}{3}r^3 fF_j - \frac{1}{3}r^3 F_k \right) \mathfrak{s}_2 + \frac{2}{5r^{\frac{1}{2}}} \mathfrak{s}_3 \\ & + \left(-\frac{8}{5r^{\frac{1}{2}}} + \frac{4}{15}r^{\frac{5}{2}} - \frac{1}{5}F + \frac{2}{15}r^4 fF' - \frac{2}{3}r^3 fF_j - \frac{1}{3}r^3 F_k \right) \mathfrak{G}_1 \\ & + \left(\frac{6}{25r^{\frac{1}{2}}} + \frac{2}{15}r^{\frac{5}{2}} - \frac{1}{25}F + \frac{11}{75}r^4 fF' + \frac{4}{75}r^5 fF'' - \frac{2}{15}r^3 fF_j - \frac{4}{15}r^4 fF'_j \right. \\ & \left. - \frac{1}{15}r^3 F_k \right) \mathfrak{G}_3 - \frac{4}{5r^{\frac{1}{2}}} \mathfrak{G}_4 + \left(\frac{4}{5r^{\frac{1}{2}}} + \frac{4}{15}r^{\frac{5}{2}} + \frac{2}{15}r^4 fF' \right) \mathfrak{G}_5 \end{aligned} \quad (5.33)$$

After expanding in powers of $1/r$, one has

$$\partial_0 r_H^{(1)} = \frac{2}{25} \mathfrak{G}_3 + \frac{4}{15} \mathfrak{G}_5 + \frac{4}{5r^{1/2}} \left(\mathfrak{s}_2 + \frac{1}{2} \mathfrak{s}_3 - 2\mathfrak{G}_1 + \frac{3}{10} \mathfrak{G}_3 - \mathfrak{G}_4 + \mathfrak{G}_5 \right). \quad (5.34)$$

The expression inside the parenthesis is just the identity eq. (4.9) that satisfied by the second order spatial viscous terms thus gives a zero. So the above finally becomes

$$\partial_0 r_H^{(1)} = \frac{2}{25} \mathfrak{G}_3 + \frac{4}{15} \mathfrak{G}_5 \quad (5.35)$$

which is actually eq. (4.14) in the case of $p = 4, q = 2$.

Eq. (3.10) gives us for D(3-1)-brane the first scalar constraint as

$$\begin{aligned} \partial_0 r_H^{(1)} = & \left(\frac{1}{3r} + \frac{1}{6}r^3 - \frac{1}{6}F + \frac{1}{12}r^5 fF' - \frac{1}{2}r^4 fF_j - \frac{1}{4}r^4 F_k \right) \mathfrak{s}_2 + \frac{1}{3r} \mathfrak{s}_3 \\ & + \left(-\frac{1}{3r} + \frac{1}{6}r^3 - \frac{1}{6}F + \frac{1}{12}r^5 fF' - \frac{1}{2}r^4 fF_j - \frac{1}{4}r^4 F_k \right) \mathfrak{G}_1 \\ & + \left(\frac{1}{18r} + \frac{1}{12}r^3 - \frac{1}{18}F + \frac{7}{72}r^5 fF' + \frac{1}{36}r^6 fF'' - \frac{1}{6}r^4 fF_j - \frac{1}{6}r^5 fF'_j \right. \\ & \left. - \frac{1}{12}r^4 F_k \right) \mathfrak{G}_3 - \frac{1}{3r} \mathfrak{G}_4 + \left(\frac{1}{3r} + \frac{1}{6}r^3 + \frac{1}{12}r^5 fF' \right) \mathfrak{G}_5, \end{aligned} \quad (5.36)$$

which gives

$$\partial_0 r_H^{(1)} = \frac{1}{36} \mathfrak{G}_3 + \frac{1}{6} \mathfrak{G}_5 + \frac{1}{3r} \left(\mathfrak{s}_2 + \mathfrak{s}_3 - \mathfrak{G}_1 + \frac{1}{6} \mathfrak{G}_3 - \mathfrak{G}_4 + \mathfrak{G}_5 \right). \quad (5.37)$$

Again, terms in the parenthesis sum to zero by eq. (4.9) and one finally has the scalar component of Navier-Stokes equation as

$$\partial_0 r_H^{(1)} = \frac{1}{36} \mathfrak{S}_3 + \frac{1}{6} \mathfrak{S}_5. \quad (5.38)$$

Then we solve the situations of $p - q = 1$. The first scalar constraint for D(4-3), D(3-2) and D(2-1)-brane can be listed separately as

$$\begin{aligned} \partial_0 r_H^{(1)} = & \left(\frac{4}{5r^{\frac{1}{2}}} + \frac{4}{15} r^{\frac{5}{2}} - \frac{1}{5} F + \frac{2}{15} r^4 f F' - \frac{2}{3} r^3 f F_j - \frac{1}{3} r^3 F_k \right) \mathbf{s}_2 + \frac{2}{5r^{\frac{1}{2}}} \mathbf{s}_3 \\ & + \left(-\frac{8}{5r^{\frac{1}{2}}} + \frac{4}{15} r^{\frac{5}{2}} - \frac{1}{5} F + \frac{2}{15} r^4 f F' - \frac{2}{3} r^3 f F_j - \frac{1}{3} r^3 F_k \right) \mathfrak{S}_1 \\ & + \left(\frac{16}{25r^{\frac{1}{2}}} + \frac{4}{15} r^{\frac{5}{2}} - \frac{1}{25} F + \frac{16}{75} r^4 f F' + \frac{4}{75} r^5 f F'' - \frac{2}{15} r^3 f F_j \right. \\ & \left. - \frac{4}{15} r^4 f F'_j - \frac{1}{15} r^3 F_k \right) \mathfrak{S}_3, \end{aligned} \quad (5.39)$$

$$\begin{aligned} \partial_0 r_H^{(1)} = & \left(\frac{1}{3r} + \frac{1}{6} r^3 - \frac{1}{6} F + \frac{1}{12} r^5 f F' - \frac{1}{2} r^4 f F_j - \frac{1}{4} r^4 F_k \right) \mathbf{s}_2 + \frac{1}{3r} \mathbf{s}_3 \\ & + \left(-\frac{1}{3r} + \frac{1}{6} r^3 - \frac{1}{6} F + \frac{1}{12} r^5 f F' - \frac{1}{2} r^4 f F_j - \frac{1}{4} r^4 F_k \right) \mathfrak{S}_1 \\ & + \left(\frac{2}{9r} + \frac{1}{6} r^3 - \frac{1}{18} F + \frac{5}{36} r^5 f F' + \frac{1}{36} r^6 f F'' - \frac{1}{6} r^4 f F_j \right. \\ & \left. - \frac{1}{6} r^5 f F'_j - \frac{1}{12} r^4 F_k \right) \mathfrak{S}_3, \end{aligned} \quad (5.40)$$

$$\begin{aligned} \partial_0 r_H^{(1)} = & \left(\frac{4}{21r^{\frac{3}{2}}} + \frac{4}{35} r^{\frac{7}{2}} - \frac{1}{7} F + \frac{2}{35} r^6 f F' - \frac{2}{5} r^5 f F_j - \frac{1}{5} r^5 F_k \right) \mathbf{s}_2 + \frac{2}{7r^{\frac{3}{2}}} \mathbf{s}_3 \\ & + \left(-\frac{8}{63r^{\frac{3}{2}}} + \frac{4}{35} r^{\frac{7}{2}} - \frac{1}{7} F + \frac{2}{35} r^6 f F' - \frac{2}{5} r^5 f F_j - \frac{1}{5} r^5 F_k \right) \mathfrak{S}_1 \\ & + \left(\frac{16}{147r^{\frac{3}{2}}} + \frac{4}{35} r^{\frac{7}{2}} - \frac{3}{49} F + \frac{24}{245} r^6 f F' + \frac{4}{245} r^7 f F'' - \frac{6}{35} r^5 f F_j \right. \\ & \left. - \frac{4}{35} r^6 f F'_j - \frac{3}{35} r^5 F_k \right) \mathfrak{S}_3. \end{aligned} \quad (5.41)$$

After expansion in terms of powers of inverse r , the above become

$$\partial_0 r_H^{(1)} = \frac{16}{75} \mathfrak{S}_3 + \frac{4}{5r^{1/2}} \left(\mathbf{s}_2 + \frac{1}{2} \mathbf{s}_3 - 2\mathfrak{S}_1 + \frac{4}{5} \mathfrak{S}_3 \right), \quad \text{D(4-3)-brane} \quad (5.42)$$

$$\partial_0 r_H^{(1)} = \frac{1}{9} \mathfrak{S}_3 + \frac{1}{3r} \left(\mathbf{s}_2 + \mathbf{s}_3 - \mathfrak{S}_1 + \frac{2}{3} \mathfrak{S}_3 \right), \quad \text{D(3-2)-brane} \quad (5.43)$$

$$\partial_0 r_H^{(1)} = \frac{16}{245} \mathfrak{S}_3 + \frac{4}{21r^{3/2}} \left(\mathbf{s}_2 + \frac{3}{2} \mathbf{s}_3 - \frac{2}{3} \mathfrak{S}_1 + \frac{4}{7} \mathfrak{S}_3 \right). \quad \text{D(2-1)-brane} \quad (5.44)$$

Once more, terms in the parenthesis of the above are all zero in the name of eq. (4.9). Thus one has finally

$$\partial_0 r_H^{(1)} = \frac{16}{75} \mathfrak{S}_3, \quad \text{D(4-3)-brane} \quad (5.45)$$

$$\partial_0 r_H^{(1)} = \frac{1}{9} \mathfrak{S}_3, \quad \text{D(3-2)-brane} \quad (5.46)$$

$$\partial_0 r_H^{(1)} = \frac{16}{245} \mathfrak{S}_3. \quad \text{D(2-1)-brane} \quad (5.47)$$

5.3.2 The scalar dynamical equations

To solve the 3 scalar perturbations, we need the second scalar constraint eq. (3.12), the (rr) component of Einstein equation eq. (3.13) and the EOM of dilaton eq. (2.19) or the scalar A eq. (2.20). Note the compactified D3-brane does not have dilaton field in the background, so we need EOM of A to solve the scalar perturbations. For compactified D4 and D2-brane, one can just use EOM of dilaton.

We still tackle the cases of $p - q = 2$ first. In the case of D(4-2)-brane, after feeding the equations mentioned above with eq. (4.16), we get

$$\begin{aligned}
 & 5(r^3 k_{(2)})' + 30r^2 j_{(2)} - 2(5r^3 - 2)h'_{(2)} \\
 &= \left(-16r + 5r^{\frac{3}{2}}F + 2r^{\frac{5}{2}}F'\right) \mathfrak{s}_2 + \left(24r + 5r^{\frac{3}{2}}F + 2r^{\frac{5}{2}}F'\right) \mathfrak{S}_1 \\
 &+ \left[-7r + r^{\frac{3}{2}}F + \frac{26}{5}r^{\frac{5}{2}}F' + \frac{4}{5}r^{\frac{7}{2}}F'' - \frac{1}{2}(5r^3 - 2)FF' + \frac{1}{4}r^4 fF'^2\right. \\
 &- 10r^{\frac{3}{2}}F_j + 45r^2 F_j^2 - 2(5r^3 - 2)F'F_j - 5r^3 F'F_k - r^4 F'F'_k + 30r^2 F_j F_k \\
 &\left. + 10r^3 F_j F'_k\right] \mathfrak{S}_3 + \left(\frac{2}{r^2} + 20r\right) \mathfrak{S}_4 - \left[18r + (5r^3 - 2)FF' + \frac{1}{2}r^4 fF'^2\right] \mathfrak{S}_5, \quad (5.48)
 \end{aligned}$$

$$\begin{aligned}
 & 2rh''_{(2)} + 3h'_{(2)} - 5j'_{(2)} \\
 &= \left(\frac{3}{4}FF' + \frac{1}{2}rFF'' + \frac{1}{4}rF'^2 + rF'F'_j - 5F_j F'_j\right) \mathfrak{S}_3 \\
 &+ \frac{2}{r^2} \mathfrak{S}_4 + \left(\frac{3}{2}FF' + rFF'' + \frac{1}{2}rF'^2\right) \mathfrak{S}_5, \quad (5.49)
 \end{aligned}$$

and

$$\begin{aligned}
 & (r^3 k_{(2)})' + r^3 f j'_{(2)} + 6r^2 j_{(2)} - r^3 f h'_{(2)} \\
 &= \left(-2r + \frac{1}{2}r^{\frac{3}{2}}F\right) \mathfrak{s}_2 + \left(6r + \frac{1}{2}r^{\frac{3}{2}}F\right) \mathfrak{S}_1 \quad (5.50) \\
 &+ \left[-r + \frac{1}{10}r^{\frac{3}{2}}F + \frac{1}{5}r^{\frac{5}{2}}F' - \frac{1}{4}r^3 fFF' - r^{\frac{3}{2}}F_j + 9r^2 F_j^2 - r^3 fF'F_j + 3r^3 fF_j F'_j\right. \\
 &\left.- \frac{1}{2}r^3 F'F_k + 6r^2 F_j F_k + r^3 F'_j F_k + 2r^3 F_j F'_k\right] \mathfrak{S}_3 + 2r \mathfrak{S}_4 - \left(2r + \frac{1}{2}r^3 fFF'\right) \mathfrak{S}_5,
 \end{aligned}$$

respectively. Then the perturbations can be solved as

$$\begin{aligned}
 h^{(2)} &= \left(\frac{2}{3} - \frac{\pi}{18\sqrt{3}} - \frac{\ln 3}{6}\right) \frac{1}{r^3} (\mathfrak{s}_2 + \mathfrak{S}_1) + \left[\frac{1}{r} + \left(\frac{1}{5} - \frac{\pi}{90\sqrt{3}} - \frac{\ln 3}{30}\right) \frac{1}{r^3}\right] \mathfrak{S}_3 \\
 &+ \frac{2}{r} \mathfrak{S}_4 + \left(\frac{2}{r} + \frac{2}{3r^3}\right) \mathfrak{S}_5, \quad (5.51)
 \end{aligned}$$

$$j^{(2)} = -\left(\frac{2}{3} - \frac{\pi}{18\sqrt{3}} - \frac{\ln 3}{6}\right) \frac{1}{r^3} (\mathfrak{s}_2 + \mathfrak{S}_1) - \left(\frac{1}{5} - \frac{\pi}{90\sqrt{3}} - \frac{\ln 3}{30}\right) \frac{1}{r^3} \mathfrak{S}_3 - \frac{2}{3r^3} \mathfrak{S}_5, \quad (5.52)$$

$$\begin{aligned}
 k^{(2)} &= \left[-\left(\frac{4}{15} - \frac{\pi}{45\sqrt{3}} - \frac{\ln 3}{15}\right) \frac{1}{r^3} + \frac{1}{10r^4} - \left(\frac{2}{3} - \frac{\pi}{18\sqrt{3}} - \frac{\ln 3}{6}\right) \frac{1}{r^6}\right] \mathfrak{s}_2 \\
 &+ \left[\frac{4}{r} - \left(\frac{4}{15} - \frac{\pi}{45\sqrt{3}} - \frac{\ln 3}{15}\right) \frac{1}{r^3} + \frac{1}{10r^4} - \left(\frac{2}{3} - \frac{\pi}{18\sqrt{3}} - \frac{\ln 3}{6}\right) \frac{1}{r^6}\right] \mathfrak{S}_1
 \end{aligned}$$

$$\begin{aligned}
& + \left[- \left(\frac{2}{25} - \frac{\pi}{225\sqrt{3}} - \frac{\ln 3}{75} \right) \frac{1}{r^3} + \frac{12}{35r^{7/2}} + \frac{47}{25r^4} - \left(\frac{7}{45} - \frac{\pi}{90\sqrt{3}} - \frac{\ln 3}{30} \right) \frac{1}{r^6} \right] \mathfrak{S}_3 \\
& - \frac{2}{r^4} \mathfrak{S}_4 - \left(\frac{4}{15r^3} - \frac{8}{7r^{7/2}} - \frac{1}{r^4} + \frac{2}{3r^6} \right) \mathfrak{S}_5.
\end{aligned} \tag{5.53}$$

By the same token as the previous case, one has for D(3-1)-brane the differential equations derived from the second scalar constraint, the (rr) component of Einstein equation and the EOM of the scalar A as

$$\begin{aligned}
& 3(r^4 k_{(2)})' + 24r^3 j_{(2)} - 2(3r^4 - 1)h'_{(2)} \\
& = (-4r + 3r^2 F + r^3 F') \mathfrak{s}_2 \\
& + (2r + 3r^2 F + r^3 F') \mathfrak{S}_1 + \left[-\frac{3}{2}r + r^2 F + \frac{8}{3}r^3 F' + \frac{1}{3}r^4 F'' - \frac{1}{2}(3r^4 - 1)FF' \right. \\
& + \frac{1}{8}r^5 fF'^2 - 6r^2 F_j + 36r^3 F_j^2 - 2(3r^4 - 1)F'F_j - 3r^4 F'F_k - \frac{1}{2}r^5 F'F'_k + 24r^3 F_j F_k \\
& \left. + 6r^4 F_j F'_k \right] \mathfrak{S}_3 + \left(\frac{1}{r^3} + 6r \right) \mathfrak{S}_4 - \left[5r + (3r^4 - 1)FF' + \frac{1}{4}r^5 fF'^2 \right] \mathfrak{S}_5,
\end{aligned} \tag{5.54}$$

$$\begin{aligned}
& rh''_{(2)} + 2h'_{(2)} - 3j'_{(2)} \\
& = \left(\frac{1}{2}FF' + \frac{1}{4}rFF'' + \frac{1}{8}rF'^2 + \frac{1}{2}rF'F'_j - 3F_j F'_j \right) \mathfrak{S}_3 \\
& + \frac{1}{r^3} \mathfrak{S}_4 + \left(FF' + \frac{1}{2}rFF'' + \frac{1}{4}rF'^2 \right) \mathfrak{S}_5,
\end{aligned} \tag{5.55}$$

$$\begin{aligned}
& (r^4 k_{(2)})' + r^4 f j'_{(2)} + 8r^3 j_{(2)} - r^4 f h'_{(2)} \\
& = \left(-r + \frac{1}{2}r^2 F \right) \mathfrak{s}_2 + \left(r + \frac{1}{2}r^2 F \right) \mathfrak{S}_1 \\
& + \left[-\frac{1}{2}r + \frac{1}{6}r^2 F + \frac{1}{6}r^3 F' - \frac{1}{4}r^4 fFF' - r^2 F_j + 12r^3 F_j^2 + 3r^4 fF_j F'_j - r^4 fF'F_j \right. \\
& \left. - \frac{1}{2}r^4 F'F_k + 8r^3 F_j F_k + r^4 F'_j F_k + 2r^4 F_j F'_k \right] \mathfrak{S}_3 + r \mathfrak{S}_4 - \left(r + \frac{1}{2}r^4 fFF' \right) \mathfrak{S}_5.
\end{aligned} \tag{5.56}$$

The solutions of the above are

$$\begin{aligned}
h^{(2)} & = \left(\frac{1}{4} - \frac{\ln 2}{8} \right) \frac{1}{r^4} (\mathfrak{s}_2 + \mathfrak{S}_1) + \left[\frac{1}{4r^2} + \left(\frac{1}{24} - \frac{\ln 2}{24} \right) \frac{1}{r^4} \right] \mathfrak{S}_3 \\
& + \frac{1}{2r^2} \mathfrak{S}_4 + \left(\frac{1}{2r^2} + \frac{1}{4r^4} \right) \mathfrak{S}_5,
\end{aligned} \tag{5.57}$$

$$j^{(2)} = - \left(\frac{1}{4} - \frac{\ln 2}{8} \right) \frac{1}{r^4} (\mathfrak{s}_2 + \mathfrak{S}_1) - \left(\frac{1}{24} - \frac{\ln 2}{24} \right) \frac{1}{r^4} \mathfrak{S}_3 - \frac{1}{4r^4} \mathfrak{S}_5, \tag{5.58}$$

$$\begin{aligned}
k^{(2)} & = \left[- \left(\frac{1}{6} - \frac{\ln 2}{12} \right) \frac{1}{r^4} + \frac{1}{18r^6} - \left(\frac{1}{4} - \frac{\ln 2}{8} \right) \frac{1}{r^8} \right] \mathfrak{s}_2 \\
& + \left[\frac{1}{r^2} - \left(\frac{1}{6} - \frac{\ln 2}{12} \right) \frac{1}{r^4} + \frac{1}{18r^6} - \left(\frac{1}{4} - \frac{\ln 2}{8} \right) \frac{1}{r^8} \right] \mathfrak{S}_1 \\
& + \left[- \left(\frac{1}{36} - \frac{\ln 2}{36} \right) \frac{1}{r^4} + \frac{1}{15r^5} + \frac{37}{108r^6} + \frac{\ln 2}{24r^8} \right] \mathfrak{S}_3 \\
& - \frac{1}{2r^6} \mathfrak{S}_4 - \left(\frac{1}{6r^4} - \frac{2}{5r^5} - \frac{1}{6r^6} + \frac{1}{4r^8} \right) \mathfrak{S}_5.
\end{aligned} \tag{5.59}$$

For the situations of one spatial dimension after compactification, i.e. $p - q = 1$, the D(4-3)-brane has the differential equations as

$$\begin{aligned}
 & 5(r^3 k_{(2)})' + 30r^2 j_{(2)} - (5r^3 - 2)h'_{(2)} \\
 &= \left(-16r + 5r^{\frac{3}{2}}F + 2r^{\frac{5}{2}}F'\right) \mathbf{s}_2 \\
 &+ \left(24r + 5r^{\frac{3}{2}}F + 2r^{\frac{5}{2}}F'\right) \mathfrak{S}_1 + \left[-16r + r^{\frac{3}{2}}F + \frac{26}{5}r^{\frac{5}{2}}F' + \frac{4}{5}r^{\frac{7}{2}}F''\right. \\
 &- (5r^3 - 2)FF' - 10r^{\frac{3}{2}}F_j + 45r^2F_j^2 - 2(5r^3 - 2)F'F_j - 5r^3F'F_k - r^4F'F'_k \\
 &\left. + 30r^2F_jF_k + 10r^3F_jF'_k\right] \mathfrak{S}_3, \tag{5.60}
 \end{aligned}$$

$$\begin{aligned}
 & 2rh''_{(2)} + 3h'_{(2)} - 10j'_{(2)} \\
 &= \left(3FF' + 2rFF'' + rF'^2 + 2rF'F'_j - 10F_jF'_j\right) \mathfrak{S}_3, \tag{5.61}
 \end{aligned}$$

$$\begin{aligned}
 & (r^3 k_{(2)})' + r^3 f j'_{(2)} + 6r^2 j_{(2)} - \frac{1}{2}r^3 f h'_{(2)} \\
 &= \left(-2r + \frac{1}{2}r^{\frac{3}{2}}F\right) \mathbf{s}_2 + \left(6r + \frac{1}{2}r^{\frac{3}{2}}F\right) \mathfrak{S}_1 \\
 &+ \left[-2r + \frac{1}{10}r^{\frac{3}{2}}F + \frac{1}{5}r^{\frac{5}{2}}F' - \frac{1}{2}r^3 f FF' - r^{\frac{3}{2}}F_j + 9r^2F_j^2 + 3r^3 f F_jF'_j\right. \\
 &\left.- r^3 f F'F_j - \frac{1}{2}r^3 F'F_k + 6r^2F_jF_k + r^3F'_jF_k + 2r^3F_jF'_k\right] \mathfrak{S}_3. \tag{5.62}
 \end{aligned}$$

Which can be solved as

$$h^{(2)} = \left(\frac{4}{3} - \frac{\pi}{9\sqrt{3}} - \frac{\ln 3}{3}\right) \frac{1}{r^3} (\mathbf{s}_2 + \mathfrak{S}_1) + \left[\frac{4}{r} + \left(\frac{16}{15} - \frac{\pi}{45\sqrt{3}} - \frac{\ln 3}{15}\right) \frac{1}{r^3}\right] \mathfrak{S}_3, \tag{5.63}$$

$$j^{(2)} = -\left(\frac{2}{3} - \frac{\pi}{18\sqrt{3}} - \frac{\ln 3}{6}\right) \frac{1}{r^3} (\mathbf{s}_2 + \mathfrak{S}_1) - \left(\frac{5}{18} - \frac{\pi}{90\sqrt{3}} - \frac{\ln 3}{30}\right) \frac{1}{r^3} \mathfrak{S}_3, \tag{5.64}$$

$$\begin{aligned}
 k^{(2)} &= \left[-\left(\frac{4}{15} - \frac{\pi}{45\sqrt{3}} - \frac{\ln 3}{15}\right) \frac{1}{r^3} + \frac{1}{10r^4} - \left(\frac{2}{3} - \frac{\pi}{18\sqrt{3}} - \frac{\ln 3}{6}\right) \frac{1}{r^6}\right] \mathbf{s}_2 \\
 &+ \left[\frac{4}{r} - \left(\frac{4}{15} - \frac{\pi}{45\sqrt{3}} - \frac{\ln 3}{15}\right) \frac{1}{r^3} + \frac{1}{10r^4} - \left(\frac{2}{3} - \frac{\pi}{18\sqrt{3}} - \frac{\ln 3}{6}\right) \frac{1}{r^6}\right] \mathfrak{S}_1 \\
 &+ \left[-\left(\frac{16}{75} - \frac{\pi}{225\sqrt{3}} - \frac{\ln 3}{75}\right) \frac{1}{r^3} + \frac{32}{35r^{7/2}} + \frac{119}{50r^4}\right. \\
 &\left.- \left(\frac{22}{45} - \frac{\pi}{90\sqrt{3}} - \frac{\ln 3}{30}\right) \frac{1}{r^6}\right] \mathfrak{S}_3. \tag{5.65}
 \end{aligned}$$

For the D(3-2)-brane, we have

$$\begin{aligned}
 & 3(r^4 k_{(2)})' + 24r^3 j_{(2)} - (3r^4 - 1)h'_{(2)} \\
 &= \left(-4r + 3r^2F + r^3F'\right) \mathbf{s}_2 + \left(2r + 3r^2F + r^3F'\right) \mathfrak{S}_1 \\
 &+ \left[-4r + r^2F + \frac{8}{3}r^3F' + \frac{1}{3}r^4F'' - (3r^4 - 1)FF' - 6r^2F_j + 36r^3F_j^2\right. \\
 &\left.- 2(3r^4 - 1)F'F_j - 3r^4F'F_k - \frac{1}{2}r^5F'F'_k + 24r^3F_jF_k + 6r^4F_jF'_k\right] \mathfrak{S}_3, \tag{5.66}
 \end{aligned}$$

$$\begin{aligned}
 & rh''_{(2)} + 2h'_{(2)} - 6j'_{(2)} \\
 &= \left(2FF' + rFF'' + \frac{1}{2}rF'^2 - 6F_jF'_j + rF'F'_j \right) \mathfrak{S}_3,
 \end{aligned} \tag{5.67}$$

$$\begin{aligned}
 & (r^4k_{(2)})' + r^4fj'_{(2)} + 8r^3j_{(2)} - \frac{1}{2}r^4fh'_{(2)} \\
 &= \left(-r + \frac{1}{2}r^2F \right) \mathfrak{s}_2 + \left(r + \frac{1}{2}r^2F \right) \mathfrak{S}_1 \\
 &+ \left[-r + \frac{1}{6}r^2F + \frac{1}{6}r^3F' - \frac{1}{2}r^4fFF' - r^2F_j + 12r^3F_j^2 + 3r^4fF_jF'_j - r^4fF'F_j \right. \\
 &\left. - \frac{1}{2}r^4F'F_k + 8r^3F_jF_k + r^4F'_jF_k + 2r^4F_jF'_k \right] \mathfrak{S}_3.
 \end{aligned} \tag{5.68}$$

Note the last differential equation is from the EOM of A . The solutions of the above are

$$h^{(2)} = \left(\frac{1}{2} - \frac{\ln 2}{4} \right) \frac{1}{r^4} (\mathfrak{s}_2 + \mathfrak{S}_1) + \left[\frac{1}{r^2} + \left(\frac{1}{3} - \frac{\ln 2}{12} \right) \frac{1}{r^4} \right] \mathfrak{S}_3, \tag{5.69}$$

$$j^{(2)} = - \left(\frac{1}{4} - \frac{\ln 2}{8} \right) \frac{1}{r^4} (\mathfrak{s}_2 + \mathfrak{S}_1) - \left(\frac{1}{6} - \frac{\ln 2}{24} \right) \frac{1}{r^4} \mathfrak{S}_3, \tag{5.70}$$

$$\begin{aligned}
 k^{(2)} &= \left[- \left(\frac{1}{6} - \frac{\ln 2}{12} \right) \frac{1}{r^4} + \frac{1}{18r^6} - \left(\frac{1}{4} - \frac{\ln 2}{8} \right) \frac{1}{r^8} \right] \mathfrak{s}_2 \\
 &+ \left[\frac{1}{r^2} - \left(\frac{1}{6} - \frac{\ln 2}{12} \right) \frac{1}{r^4} + \frac{1}{18r^6} - \left(\frac{1}{4} - \frac{\ln 2}{8} \right) \frac{1}{r^8} \right] \mathfrak{S}_1 \\
 &+ \left[- \left(\frac{1}{9} - \frac{\ln 2}{36} \right) \frac{1}{r^4} + \frac{4}{15r^5} + \frac{23}{54r^6} - \left(\frac{1}{8} - \frac{\ln 2}{24} \right) \frac{1}{r^8} \right] \mathfrak{S}_3.
 \end{aligned} \tag{5.71}$$

The differential equations for the second order scalar perturbations of D(2-1)-brane can be got as

$$\begin{aligned}
 & 7(r^5k_{(2)})' + 70r^4j_{(2)} - (7r^5 - 2)h'_{(2)} \\
 &= \left(-\frac{16}{3}r + 7r^{\frac{5}{2}}F + 2r^{\frac{7}{2}}F' \right) \mathfrak{s}_2 + \left(\frac{8}{9}r + 7r^{\frac{5}{2}}F + 2r^{\frac{7}{2}}F' \right) \mathfrak{S}_1 \\
 &+ \left[-\frac{16}{3}r + 3r^{\frac{5}{2}}F + \frac{38}{7}r^{\frac{7}{2}}F' + \frac{4}{7}r^{\frac{9}{2}}F'' - (7r^5 - 2)FF' - 14r^{\frac{5}{2}}F_j + 105r^4F_j^2 \right. \\
 &\left. - 2(7r^5 - 2)F'F_j - 7r^5F'F_k - r^6F'F'_k + 70r^4F_jF_k + 14r^5F_jF'_k \right] \mathfrak{S}_3,
 \end{aligned} \tag{5.72}$$

$$\begin{aligned}
 & 2rh''_{(2)} + 5h'_{(2)} - 14j'_{(2)} \\
 &= \left(5FF' + 2rFF'' + rF'^2 - 14F_jF'_j + 2rF'F'_j \right) \mathfrak{S}_3,
 \end{aligned} \tag{5.73}$$

$$\begin{aligned}
 & (r^5k_{(2)})' + r^5fj'_{(2)} + 10r^4j_{(2)} - \frac{1}{2}r^5fh'_{(2)} \\
 &= \left(-\frac{2}{3}r + \frac{1}{2}r^{\frac{5}{2}}F \right) \mathfrak{s}_2 + \left(\frac{2}{9}r + \frac{1}{2}r^{\frac{5}{2}}F \right) \mathfrak{S}_1 \\
 &+ \left[-\frac{2}{3}r + \frac{3}{14}r^{\frac{5}{2}}F + \frac{1}{7}r^{\frac{7}{2}}F' - \frac{1}{2}r^5fFF' - r^{\frac{5}{2}}F_j + 15r^4F_j^2 + 3r^5fF_jF'_j \right. \\
 &\left. - r^5fF'F_j - \frac{1}{2}r^5F'F_k + 10r^4F_jF_k + r^5F'_jF_k + 2r^5F_jF'_k \right] \mathfrak{S}_3.
 \end{aligned} \tag{5.74}$$

Solving these perturbations is tough because $F(r)$ of D2-brane is more complex than the other Dp-brane, as can be seen from eq. (3.5). One can simplify the calculation by first making the series expansion with respect to $1/r$ and then do the integral when solving the differential equations of $j^{(2)}$ and $k^{(2)}$. The solutions are

$$h^{(2)} = \left(\frac{4}{15} + \frac{\pi}{25} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{1}{5\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{\ln 5}{10} \right) \frac{1}{r^5} (\mathbf{s}_2 + \mathfrak{S}_1) + \left[\frac{4}{9r^3} + \left(\frac{16}{105} + \frac{3\pi}{175} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{3}{35\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{3}{70} \ln 5 \right) \frac{1}{r^5} \right] \mathfrak{S}_3, \quad (5.75)$$

$$j^{(2)} = - \left(\frac{2}{15} + \frac{\pi}{50} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{1}{10\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{\ln 5}{20} \right) \frac{1}{r^5} (\mathbf{s}_2 + \mathfrak{S}_1) - \left(\frac{8}{105} + \frac{3\pi}{350} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{3}{70\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{3 \ln 5}{140} \right) \frac{1}{r^5} \mathfrak{S}_3, \quad (5.76)$$

$$k^{(2)} = \left[- \left(\frac{4}{35} + \frac{3\pi}{175} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{3}{35\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{3 \ln 5}{70} \right) \frac{1}{r^5} + \frac{1}{28r^8} - \left(\frac{2}{15} + \frac{\pi}{50} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{1}{10\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{\ln 5}{20} \right) \frac{1}{r^{10}} \right] \mathbf{s}_2 + \left[\frac{4}{9r^3} - \left(\frac{4}{35} + \frac{3\pi}{175} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{3}{35\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{3 \ln 5}{70} \right) \frac{1}{r^5} + \frac{1}{28r^8} - \left(\frac{2}{15} + \frac{\pi}{50} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{1}{10\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{\ln 5}{20} \right) \frac{1}{r^{10}} \right] \mathfrak{S}_1 + \left[- \left(\frac{16}{245} + \frac{9\pi}{1225} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{9}{245\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{9 \ln 5}{490} \right) \frac{1}{r^5} + \frac{32}{273r^{\frac{13}{2}}} + \frac{253}{1764r^8} - \left(\frac{22}{525} + \frac{3\pi}{350} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{3}{70\sqrt{5}} \operatorname{arccoth} \sqrt{5} - \frac{3 \ln 5}{140} \right) \frac{1}{r^{10}} \right] \mathfrak{S}_3. \quad (5.77)$$

6 The second order constitutive relation

In this section, we will offer all the results of the second order constitutive relations for the D(p-q)-brane with $2 \leq p \leq 4$ and $1 \leq q \leq p - 1$. We will omit D(4-1)-brane case here as in the previous sections, for its results can be found in [24]. One will find the following substitutions very helpful to gain the final form of the constitutive relations:

$$\begin{aligned} t_{3ij} &= \partial_0 \sigma_{ij} \rightarrow \langle D \partial_\mu u_\nu \rangle, & \mathfrak{T}_{1ij} &= \partial_0 \beta_i \partial_0 \beta_j - \frac{1}{p-q} \delta_{ij} \mathbf{s}_1 \rightarrow Du_{\langle \mu} Du_{\nu \rangle}, \\ \mathfrak{T}_{4ij} &= \sigma_{ij} \partial \beta \rightarrow \sigma_{\mu\nu} \partial_\rho u^\rho, & \mathfrak{T}_{7ij} &= \sigma_{(i}^k \Omega_{j)k} \rightarrow \sigma_{\langle \mu}^\rho \Omega_{\nu \rangle \rho} \end{aligned} \quad (6.1)$$

are for the spatial viscous tensors and

$$\mathbf{s}_2 + \mathfrak{S}_1 = \partial_0 \partial \beta + \partial_0 \beta_i \partial_0 \beta_i \rightarrow D \partial u, \quad \mathfrak{S}_3 = (\partial \beta)^2 \rightarrow (\partial u)^2, \quad \mathfrak{S}_5 = \sigma_{ij}^2 \rightarrow \sigma_{\mu\nu}^2 \quad (6.2)$$

are for the spatial viscous scalar terms. Here the angle brackets stand for the transverse traceless symmetrization of rank 2 tensors, i.e.

$$A_{\langle\mu\nu\rangle} = P_{\mu}^{\rho}P_{\nu}^{\lambda}A_{(\rho\lambda)} - \frac{1}{p-q}P_{\mu\nu}P^{\rho\lambda}A_{\rho\lambda}. \quad (6.3)$$

The situation of $p - q = 2$ includes the D(4-2) and D(3-1)-brane cases. For the D(4-2)-brane, one has

$$\begin{aligned} T_{\mu\nu} = & \frac{1}{2\kappa_4^2} \left\{ \frac{r_H^3}{L_4^4} \left(\frac{5}{2}u_{\mu}u_{\nu} + \frac{1}{2}P_{\mu\nu} \right) - \left(\frac{r_H}{L_4} \right)^{\frac{5}{2}} \left(2\sigma_{\mu\nu} + \frac{3}{5}\partial_{\rho}u^{\rho}P_{\mu\nu} \right) \right. \\ & + \frac{r_H^2}{L_4} \left[\left(2 - \frac{\pi}{6\sqrt{3}} - \frac{\ln 3}{2} \right) \cdot 2 \left(\langle D\sigma_{\mu\nu} \rangle + \frac{1}{2}\sigma_{\mu\nu}\partial u \right) \right. \\ & + \left. \left(\frac{6}{5} + \frac{\pi}{10\sqrt{3}} + \frac{3\ln 3}{10} \right) \frac{2\sigma_{\mu\nu}\partial u}{2} - \left(\frac{\pi}{3\sqrt{3}} + \ln 3 \right) \cdot 2\sigma_{\langle\mu}{}^{\rho}\Omega_{\nu\rangle\rho} \right] \\ & + \frac{r_H^2}{L_4}P_{\mu\nu} \left[\left(\frac{6}{5} - \frac{\pi}{10\sqrt{3}} - \frac{3\ln 3}{10} \right) D(\partial u) \right. \\ & \left. \left. + \left(\frac{9}{25} - \frac{\pi}{50\sqrt{3}} - \frac{3\ln 3}{50} \right) (\partial u)^2 + \frac{3}{10} \cdot 4\sigma_{\rho\lambda}^2 \right] \right\}. \quad (6.4) \end{aligned}$$

With the second order transport coefficients can be read from the above as

$$\begin{aligned} \eta\tau_{\pi} &= \frac{1}{2\kappa_4^2} \left(2 - \frac{\pi}{6\sqrt{3}} - \frac{\ln 3}{2} \right) \frac{r_H^2}{L}, & \eta\tau_{\pi}^* &= \frac{1}{2\kappa_4^2} \left(\frac{6}{5} + \frac{\pi}{10\sqrt{3}} + \frac{3\ln 3}{10} \right) \frac{r_H^2}{L}, \\ \lambda_2 &= -\frac{1}{2\kappa_4^2} \left(\frac{\pi}{3\sqrt{3}} + \ln 3 \right) \frac{r_H^2}{L}, & \zeta\tau_{\Pi} &= \frac{1}{2\kappa_4^2} \left(\frac{6}{5} - \frac{\pi}{10\sqrt{3}} - \frac{3\ln 3}{10} \right) \frac{r_H^2}{L}, \\ \xi_1 &= \frac{1}{2\kappa_4^2} \frac{3}{10} \frac{r_H^2}{L}, & \xi_2 &= \frac{1}{2\kappa_4^2} \left(\frac{9}{25} - \frac{\pi}{50\sqrt{3}} - \frac{3\ln 3}{50} \right) \frac{r_H^2}{L}. \quad (6.5) \end{aligned}$$

The D(3-1)-brane has the second order stress energy tensor as

$$\begin{aligned} T_{\mu\nu} = & \frac{1}{2\kappa_4^2} \left\{ \frac{r_H^4}{L_3^5} \left(3u_{\mu}u_{\nu} + P_{\mu\nu} \right) - \left(\frac{r_H}{L_3} \right)^3 \left(2\sigma_{\mu\nu} + \frac{1}{3}P_{\mu\nu}\partial_{\rho}u^{\rho} \right) \right. \\ & + \frac{r_H^2}{L_3} \left[\left(1 - \frac{\ln 2}{2} \right) \cdot 2 \left(\langle D\sigma_{\mu\nu} \rangle + \frac{1}{2}\sigma_{\mu\nu}\partial u \right) + \left(\frac{1}{3} + \frac{\ln 2}{6} \right) \frac{2\sigma_{\mu\nu}\partial u}{2} - \ln 2 \cdot 2\sigma_{\langle\mu}{}^{\rho}\Omega_{\nu\rangle\rho} \right] \\ & \left. + P_{\mu\nu} \frac{r_H^2}{L_3} \left[\left(\frac{1}{3} - \frac{\ln 2}{6} \right) D(\partial u) + \left(\frac{1}{18} - \frac{\ln 2}{18} \right) (\partial u)^2 + \frac{1}{12} \cdot 4\sigma_{\rho\lambda}^2 \right] \right\}. \quad (6.6) \end{aligned}$$

The corresponding second order transport coefficients are

$$\begin{aligned} \eta\tau_{\pi} &= \frac{1}{2\kappa_4^2} \left(1 - \frac{\ln 2}{2} \right) \frac{r_H^2}{L_3}, & \eta\tau_{\pi}^* &= \frac{1}{2\kappa_4^2} \left(\frac{1}{3} + \frac{\ln 2}{6} \right) \frac{r_H^2}{L_3}, & \lambda_2 &= -\frac{1}{2\kappa_4^2} \ln 2 \frac{r_H^2}{L_3}, \\ \zeta\tau_{\Pi} &= \frac{1}{2\kappa_4^2} \left(\frac{1}{3} - \frac{\ln 2}{6} \right) \frac{r_H^2}{L_3}, & \xi_1 &= \frac{1}{2\kappa_4^2} \frac{1}{12} \frac{r_H^2}{L_3}, & \xi_2 &= \frac{1}{2\kappa_4^2} \left(\frac{1}{18} - \frac{\ln 2}{18} \right) \frac{r_H^2}{L_3}. \quad (6.7) \end{aligned}$$

When $p - q = 1$, one has the D(4-3), D(3-2) and D(2-1)-brane. These cases are like the D1-brane in [26]. The spatial viscous terms do not contain σ_{ij} or Ω_{ij} since there is only one spatial direction. Thus the results will be much simpler than the cases of $p - q = 2$. The second order constitutive relation for D(4-3)-brane is

$$T_{\mu\nu} = \frac{1}{2\kappa_3^2} \left\{ \frac{r_H^3}{L_4^4} \left(\frac{5}{2} u_\mu u_\nu + \frac{1}{2} P_{\mu\nu} \right) - \left(\frac{r_H}{L_4} \right)^{\frac{5}{2}} \frac{8}{5} P_{\mu\nu} \partial_\rho u^\rho \right. \\ \left. + \frac{r_H^2}{L_4} P_{\mu\nu} \left[\left(\frac{16}{5} - \frac{4\pi}{15\sqrt{3}} - \frac{4\ln 3}{5} \right) D(\partial u) + \left(\frac{64}{25} - \frac{4\pi}{75\sqrt{3}} - \frac{4\ln 3}{25} \right) (\partial u)^2 \right] \right\}. \quad (6.8)$$

There are only two transport coefficients at the second order:

$$\zeta\tau_{\text{II}} = \frac{1}{2\kappa_3^2} \left(\frac{16}{5} - \frac{4\pi}{15\sqrt{3}} - \frac{4\ln 3}{5} \right) \frac{r_H^2}{L_4}, \quad \xi_2 = \frac{1}{2\kappa_3^2} \left(\frac{64}{25} - \frac{4\pi}{75\sqrt{3}} - \frac{4\ln 3}{25} \right) \frac{r_H^2}{L_4} \quad (6.9)$$

In the D(3-2)-brane case one has

$$T_{\mu\nu} = \frac{1}{2\kappa_3^2} \left\{ \frac{r_H^4}{L_3^5} \left(3u_\mu u_\nu + P_{\mu\nu} \right) - \left(\frac{r_H}{L_3} \right)^3 \frac{4}{3} P_{\mu\nu} \partial_\rho u^\rho \right. \\ \left. + P_{\mu\nu} \frac{r_H^2}{L_3} \left[\left(\frac{4}{3} - \frac{2\ln 2}{3} \right) D(\partial u) + \left(\frac{8}{9} - \frac{2\ln 2}{9} \right) (\partial u)^2 \right] \right\}, \quad (6.10)$$

whose second order transport coefficients can be read as

$$\zeta\tau_{\text{II}} = \frac{1}{2\kappa_3^2} \left(\frac{4}{3} - \frac{2\ln 2}{3} \right) \frac{r_H^2}{L_3}, \quad \xi_2 = \frac{1}{2\kappa_3^2} \left(\frac{8}{9} - \frac{2\ln 2}{9} \right) \frac{r_H^2}{L_3}. \quad (6.11)$$

The situation of D(2-1)-brane gives us

$$T_{\mu\nu} = \frac{1}{2\kappa_3^2} \left\{ \frac{r_H^5}{L_2^6} \left(\frac{7}{2} u_\mu u_\nu + \frac{3}{2} P_{\mu\nu} \right) - \left(\frac{r_H}{L_2} \right)^{\frac{7}{2}} \frac{8}{7} P_{\mu\nu} \partial u \right. \\ \left. + P_{\mu\nu} \frac{r_H^2}{L_2} \left[\left(\frac{16}{21} + \frac{4\pi}{35} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{4}{7\sqrt{5}} \operatorname{arcoth} \sqrt{5} - \frac{2\ln 5}{7} \right) D\partial u \right. \right. \\ \left. \left. + \left(\frac{64}{147} + \frac{12\pi}{245} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{12}{49\sqrt{5}} \operatorname{arcoth} \sqrt{5} - \frac{6\ln 5}{49} \right) (\partial u)^2 \right] \right\}, \quad (6.12)$$

with the second order coefficients are

$$\zeta\tau_{\text{II}} = \frac{1}{2\kappa_3^2} \left(\frac{16}{21} + \frac{4\pi}{35} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{4}{7\sqrt{5}} \operatorname{arcoth} \sqrt{5} - \frac{2\ln 5}{7} \right) \frac{r_H^2}{L_2}, \\ \xi_2 = \frac{1}{2\kappa_3^2} \left(\frac{64}{147} + \frac{12\pi}{245} \sqrt{1 - \frac{2}{\sqrt{5}}} + \frac{12}{49\sqrt{5}} \operatorname{arcoth} \sqrt{5} - \frac{6\ln 5}{49} \right) \frac{r_H^2}{L_2}. \quad (6.13)$$

All the results listed in this section can be rewritten in a universal form by introducing the Harmonic number and the second order constitutive relations for the Dp-brane with q

directions of their world-volume compactified can be recast in the form as

$$\begin{aligned}
T_{\mu\nu} = & \frac{1}{2\kappa_{p-q+2}^2} \left\{ \frac{r_H^{7-p}}{L_p^{8-p}} \left(\frac{9-p}{2} u_\mu u_\nu + \frac{5-p}{2} P_{\mu\nu} \right) \right. \\
& - \left(\frac{r_H}{L_p} \right)^{\frac{9-p}{2}} \left(2\sigma_{\mu\nu} + \frac{2(p-3)^2 + 2q(5-p)}{(p-q)(9-p)} P_{\mu\nu} \partial u \right) \\
& + \frac{r_H^2}{L_p} \left[\left(\frac{1}{5-p} + \frac{1}{7-p} H_{\frac{5-p}{7-p}} \right) \cdot 2 \left(\langle D \sigma_{\mu\nu} \rangle + \frac{1}{p-q} \sigma_{\mu\nu} \partial u \right) \right. \\
& + \left(\frac{3(p-3)^2 + 3q(5-p)}{(5-p)(9-p)} - \frac{(p-3)^2 + q(5-p)}{(7-p)(9-p)} H_{\frac{5-p}{7-p}} \right) \frac{2\sigma_{\mu\nu} \partial u}{p-q} \\
& + \left. \frac{1}{5-p} \cdot 4\sigma_{\langle \mu}{}^\rho \sigma_{\nu \rangle \rho} + \left(-\frac{2}{5-p} + \frac{2}{7-p} H_{\frac{5-p}{7-p}} \right) \cdot 2\sigma_{\langle \mu}{}^\rho \Omega_{\nu \rangle \rho} \right] \\
& + P_{\mu\nu} \frac{r_H^2}{L_p} \left[\left(\frac{2(p-3)^2 + 2q(5-p)}{(p-q)(5-p)(9-p)} + \frac{2(p-3)^2 + 2q(5-p)}{(p-q)(7-p)(9-p)} H_{\frac{5-p}{7-p}} \right) D(\partial u) \right. \\
& + \left(\frac{[2(p-3)^2 + 2q(5-p)][(3p^2 - 17p + 18) + 3q(5-p)]}{(p-q)^2(5-p)(9-p)^2} \right. \\
& + \left. \frac{(5-p)[2(p-3)^2 + 2q(5-p)]}{(p-q)(7-p)(9-p)^2} H_{\frac{5-p}{7-p}} \right) (\partial u)^2 \\
& \left. + \frac{(p-3)^2 + q(5-p)}{(p-q)(5-p)(9-p)} \cdot 4\sigma_{\alpha\beta}^2 \right\}. \tag{6.14}
\end{aligned}$$

from which one can read all the second order transport coefficients as

$$\begin{aligned}
\eta\tau_\pi &= \frac{1}{2\kappa_{p-q+2}^2} \left(\frac{1}{5-p} + \frac{1}{7-p} H_{\frac{5-p}{7-p}} \right) \frac{r_H^2}{L_p}, \\
\eta\tau_\pi^* &= \frac{1}{2\kappa_{p-q+2}^2} \left[\frac{3(p-3)^2 + 3q(5-p)}{(5-p)(9-p)} - \frac{(p-3)^2 + q(5-p)}{(7-p)(9-p)} H_{\frac{5-p}{7-p}} \right] \frac{r_H^2}{L_p}, \\
\lambda_1 &= \frac{1}{2\kappa_{p-q+2}^2} \frac{1}{5-p} \frac{r_H^2}{L_p}, \quad \lambda_2 = \frac{1}{2\kappa_{p-q+2}^2} \left(-\frac{2}{5-p} + \frac{2}{7-p} H_{\frac{5-p}{7-p}} \right) \frac{r_H^2}{L_p}, \\
\zeta\tau_\Pi &= \frac{1}{2\kappa_{p-q+2}^2} \left[\frac{2(p-3)^2 + 2q(5-p)}{(p-q)(5-p)(9-p)} + \frac{2(p-3)^2 + 2q(5-p)}{(p-q)(7-p)(9-p)} H_{\frac{5-p}{7-p}} \right] \frac{r_H^2}{L_p}, \\
\xi_1 &= \frac{1}{2\kappa_{p-q+2}^2} \frac{(p-3)^2 + q(5-p)}{(p-q)(5-p)(9-p)} \frac{r_H^2}{L_p}, \\
\xi_2 &= \frac{1}{2\kappa_{p-q+2}^2} \left[\frac{[2(p-3)^2 + 2q(5-p)][(3p^2 - 17p + 18) + 3q(5-p)]}{(p-q)^2(5-p)(9-p)^2} \right. \\
& \left. + \frac{(5-p)[2(p-3)^2 + 2q(5-p)]}{(p-q)(7-p)(9-p)^2} H_{\frac{5-p}{7-p}} \right] \frac{r_H^2}{L_p}. \tag{6.15}
\end{aligned}$$

We have added the result of D(4-1)-brane into the above formulas, that's why the viscous tensor relates with λ_1 appears. The definition of the Harmonic number $H_{\frac{5-p}{7-p}}$ with its

special value for the cases of Dp-brane with $1 \leq p \leq 4$ can be found in [26]. Eq. (6.14) and (6.15) will reproduce the results in [26] and [24] by separately setting $q = 0$ and $p = 4, q = 1$.

One can easily see that of the 7 second order transport coefficients that we derived, only $\eta\tau_\pi^*, \zeta\tau_\Pi, \xi_1, \xi_2$ depend on the number of the compact directions of the brane, i.e., their expressions contain q while $\eta\tau_\pi, \lambda_1, \lambda_2$ do not. If we take a closer look at how q enters the formulae, we will find that q appears in the form of

$$\frac{2(p-3)^2 + 2q(5-p)}{(p-q)(9-p)}, \tag{6.16}$$

which is just the numerical factor of bulk viscosity ζ . Only $\eta\tau_\pi^*$ is a little special. The $p-q$ in the denominator is with the viscous tensor $\frac{2\sigma_{\mu\nu}\partial u}{p-q}$ by definition.

Since the number of spatial dimensions of the compactified Dp-brane is also the number of the spatial dimensions of the boundary fluid. Compactifying the Dp-brane will reduce the number of the spatial dimensions of the boundary fluid, too. For the boundary fluid, the bulk flow viscous terms $\partial u, D(\partial u), (\partial u)^2$ and $\sigma_{\mu\nu}\partial u$ contain the sum of diagonal spatial components and the bulk flow viscous term $\sigma_{\alpha\beta}^2$ sums the nondiagonal spatial components. There will always be some spatial components lost if one or more spatial directions of Dp-branes are compactified. So the bulk flow viscous terms are affected by compactification and their transport coefficients $\eta\tau_\pi^*, \zeta\tau_\Pi, \xi_1$ and ξ_2 should include q . On the contrary, the shear flow viscous terms $\langle D\sigma_{\mu\nu} \rangle, \sigma_{\langle\mu}{}^\rho\sigma_{\nu\rangle\rho}, \sigma_{\langle\mu}{}^\rho\Omega_{\nu\rangle\rho}$ are tensorial in nature. By definition they do not sum spatial components. Thus their transport coefficients $\eta\tau_\pi, \lambda_1, \lambda_2$ should not be under affected by compactification and do not contain q .

As one can check, the results of eq. (6.15) satisfy the Haack-Yarom relation¹ $4\lambda_1 + \lambda_2 = 2\eta\tau_\pi$, the Romatschke relations [22]

$$\tau_\pi = \tau_\Pi, \quad \xi_1 = \frac{1}{p-q}[1 - (p-q)c_s^2]\lambda_1 \tag{6.17}$$

and also the Kleinert-Probst relations [28]

$$\begin{aligned} \eta\tau_\pi^* &= (1 - (p-q)c_s^2)(4\lambda_1 - \eta\tau_\pi), \\ \xi_2 &= \frac{2}{(p-q)^2} (1 - (p-q)c_s^2) [(1 - 2(p-q)c_s^2)2\lambda_1 + (p-q)c_s^2\eta\tau_\pi]. \end{aligned} \tag{6.18}$$

Since $p-q$ is the spatial dimension of the dual relativistic fluid, if we denote d as the dimension of the relativistic fluid, then we have $p-q = d-1$. So the above Romatschke and Kleinert-Probst relations still have the same form as in [26]. The above 5 relations are among the 7 dynamical transport coefficients, one may draw a conclusion that only 2 of the 7 dynamical transport coefficients are actually independent. For which one can choose $\eta\tau_\pi$ and λ_1 , since they appear in every relation. Ref. [29] derives 5 relations among the 8 thermodynamical second order transport coefficients,² so only 3 of them are actually

¹It is firstly found in 5 dimensional charged AdS black hole [7] and is firstly confirmed in nonconformal regime holographically in [27].

²Following the nomenclature of [30].

independent. For which one can choose λ_3 , λ_4 and κ . Thus one may make a bold guess that at second order, nonconformal relativistic fluids have at most 5 independent flow modes:

$$\langle D\sigma_{\mu\nu} \rangle, \quad \sigma_{\langle\mu}{}^\rho\sigma_{\nu\rangle\rho}, \quad \Omega_{\langle\mu}{}^\rho\Omega_{\nu\rangle\rho}, \quad \nabla_{\langle\mu}^\perp \ln s \nabla_{\nu\rangle}^\perp \ln s, \quad \left(R_{\langle\mu\nu\rangle} - 2u^\rho u^\sigma R_{\rho\langle\mu\nu\rangle\sigma} \right), \quad (6.19)$$

which relate to the 5 independent second order transport coefficients. Here $\nabla_\mu^\perp = P_\mu^\rho \nabla_\rho$ is the covariant derivative perpendicular to the direction of u^μ .

One can also reformulate the results of this work in terms of field theory quantities, which can be found in table 2. Compared with the field theory reformulation of the transport coefficients of Dp-brane [26], here we use $\Gamma\left(\frac{7-p}{2}\right)$ but not $\Gamma\left(\frac{9-p}{2}\right)$. Thus the dimensionless part of the results here is different from that of Dp-brane. We have also set the characteristic energy scale of the Dp-brane the same as the Hawking temperature, i.e. $\Lambda = T$ to make the expressions neat and clear. Another obvious difference from the Dp-brane case is that V_q appears in all the thermal quantities as well as the transport coefficients which comes from the integration on the compact directions of Dp-branes.

The compactification on the gravity side in holographic relativistic hydrodynamics introduces the volume of the compact directions V_q to balance the dimension of the physical results. V_q in the results shown in table 2 will neutralize the extra mass dimension from T to make the dimension of thermal quantities and transport coefficients to be right.

7 Summary and discussions

In this paper, we derive all the second order dynamical transport coefficients for Dp-brane with q directions of their world-volume compactified. To be more specific, the situations considered in this paper include the D4-brane with 2 and 3 directions of its world-volume compactified, D3-brane with 1 and 2 directions compactified, as well as D2-brane with 1 direction compactified. This work can be seen as a generalization of [26] which considers only the Dp-brane, thus by setting $q = 0$ one can reproduce all the results of [26]. Through [23, 24, 26] and this work, we have finished the calculation for all the dynamical second order transport coefficients for Dp-brane with or without compactified dimensions.

Through the calculation we can see that the compactification on the world-volume of Dp-brane only affects the scalar perturbation h , which can be seen as the spatial trace part of the perturbative metric and relates to the number of the spatial dimensions of the metric. At the first order, $h^{(1)}$ contributes to the spatial diagonal components of the Brown-York tensor on the gravity side and determines the value of ζ . From the boundary fluid viewpoint, $h^{(1)}$ is dual to the viscous term ∂u of the stress tensor. Thus the bulk viscosity ζ is dependent on the number of compact dimensions q . At the second order, the bulk flow viscous terms of $\eta\tau_\pi^*$, $\zeta\tau_\Pi$, ξ_1 , ξ_2 relate with ∂u . Their results all contain the numerical factor of ζ and depend on the number of compact directions q . Whereas results of $\eta\tau_\pi$, λ_1 , λ_2 do not contain q since their viscous tensors do not need to sum the spatial components and are not affected by compactification.

The second order transport coefficients for the compactified Dp-brane also satisfy the Haack-Yarom, Romatschke and Kleinert-Probst relations as in [26]. Calculation shows the dispersion relations of the D(p-q)-brane do not change compared with Dp-brane case.

ε	$\frac{9-p}{2} \left(\frac{4\pi}{7-p} \right)^2 \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p+1}$
\mathfrak{p}	$\frac{5-p}{2} \left(\frac{4\pi}{7-p} \right)^2 \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p+1}$
η	$\left(\frac{4\pi}{7-p} \right)^2 \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^p$
ζ	$\frac{2(p-3)^2+2q(5-p)}{(p-q)(9-p)} \left(\frac{4\pi}{7-p} \right)^2 \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^p$
$\eta\tau_\pi$	$\left(\frac{1}{5-p} + \frac{1}{7-p} H_{\frac{5-p}{7-p}} \right) \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p-1}$
$\eta\tau_\pi^*$	$\left[\frac{3(p-3)^2+3q(5-p)}{(5-p)(9-p)} - \frac{(p-3)^2+q(5-p)}{(7-p)(9-p)} H_{\frac{5-p}{7-p}} \right] \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p-1}$
λ_1	$\frac{1}{5-p} \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p-1}$
λ_2	$\left(-\frac{2}{5-p} + \frac{2}{7-p} H_{\frac{5-p}{7-p}} \right) \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p-1}$
$\zeta\tau_\Pi$	$\left[\frac{2(p-3)^2+2q(5-p)}{(p-q)(5-p)(9-p)} + \frac{2(p-3)^2+2q(5-p)}{(p-q)(7-p)(9-p)} H_{\frac{5-p}{7-p}} \right] \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p-1}$
ξ_1	$\frac{(p-3)^2+q(5-p)}{(p-q)(5-p)(9-p)} \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p-1}$
ξ_2	$\left[\frac{[2(p-3)^2+2q(5-p)][(3p^2-17p+18)+3q(5-p)]}{(p-q)^2(5-p)(9-p)^2} + \frac{(5-p)[2(p-3)^2+2q(5-p)]}{(p-q)(7-p)(9-p)^2} H_{\frac{5-p}{7-p}} \right] \\ \times \frac{2^{\frac{2(6-p)}{5-p}} \pi^{\frac{3-p}{5-p}}}{(7-p)^{\frac{9-p}{5-p}}} \Gamma \left(\frac{7-p}{2} \right)^{\frac{2}{5-p}} \lambda_{p+1}^{\frac{p-3}{5-p}} N^2 V_q T^{p-1}$

Table 2. The dual field theory reformulation of the results of compactified Dp-brane. Here V_q is the volume of the compact dimensions of Dp-brane, $\lambda_{p+1} = g_{p+1}^2 N T^{p-3}$ with $g_{p+1}^2 = (2\pi)^{p-2} g_s l_s^{p-3}$. λ_3 and ξ_3 for the compactified Dp-brane are both zero.

Although the calculation of the dispersion relations connect with the linear viscous terms ∂u and $D(\partial u)$. From the discussions above, the dispersion relations should depend on the number of compact directions q . But the q -dependent factors cancel each other in the calculation process making the final results of the dispersion relations contain no q .

We know from [26] that near-extremal black D3-brane will lead to the asymptotically AdS₅ black hole after integrating out the unit 5-sphere. Keep on compactifying one or more world-volume directions of the near-extremal black D3-brane equals to compactify the AdS₅ black hole. From the results of the compactified D3-brane in this work, one can see that compactification on AdS₅ black hole can lead to nonconformal results for the dual

fluid's transport coefficients. This reminds us that we can also get nonconformal transport coefficients from compactified AdS black holes of other dimensions. This calculation may give a direct check on the method of obtaining nonconformal hydrodynamical stress-energy tensor from a conformal one that is proposed in [20].

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A Christoffel symbol and Ricci quantities of the reduction ansatz

The reduction ansatz used in this work is

$$ds^2 = e^{2\alpha_1 A} g_{MN} dx^M dx^N + e^{2\alpha_2 A} \left(e^{2\beta_1 B} \delta_{mn} dy^m dy^n + e^{2\beta_2 B} L_p^2 d\Omega_{8-p}^2 \right). \quad (\text{A.1})$$

We separately denote $\tilde{\Gamma}_{\hat{N}\hat{P}}^{\hat{M}}$ and Γ_{NP}^M as the Christoffel symbols in 10 dimension and $p-q+2$ dimensional reduced theory. ${}^\Omega\Gamma_{bc}^a$ is the Christoffel symbol on $8-p$ dimensional unit sphere. The Christoffel symbols of the reduction ansatz can be listed as

$$\begin{aligned} \tilde{\Gamma}_{NP}^M &= \Gamma_{NP}^M + \alpha_1 (\delta_N^M \partial_P A + \delta_P^M \partial_N A - g_{NP} \nabla^M A), \\ \tilde{\Gamma}_{mn}^M &= -(\alpha_2 \nabla^M A + \beta_1 \nabla^M B) e^{(-2\alpha_1 + 2\alpha_2)A + 2\beta_1 B} \delta_{mn}, \\ \tilde{\Gamma}_{Mm}^n &= (\alpha_2 \partial_M A + \beta_1 \partial_M B) \delta_m^n, \\ \tilde{\Gamma}_{ab}^M &= -(\alpha_2 \nabla^M A + \beta_2 \nabla^M B) e^{(-2\alpha_1 + 2\alpha_2)A + 2\beta_2 B} \gamma_{ab} L_p^2, \\ \tilde{\Gamma}_{Mb}^a &= (\alpha_2 \partial_M A + \beta_2 \partial_M B) \delta_b^a, \\ \tilde{\Gamma}_{bc}^a &= {}^\Omega\Gamma_{bc}^a. \end{aligned} \quad (\text{A.2})$$

The following relations are useful in the calculation:

$$\begin{aligned} \tilde{\Gamma}_{MN}^N &= \Gamma_{MN}^N + (p-q+2)\alpha_1 \partial_M A, \\ \tilde{\Gamma}_{M\hat{N}}^{\hat{N}} &= \Gamma_{MN}^N + [(p-q+2)\alpha_1 + (8-p+q)\alpha_2] \partial_M A + [q\beta_1 + (8-p)\beta_2] \partial_M B. \end{aligned} \quad (\text{A.3})$$

Here $\tilde{\Gamma}_{M\hat{N}}^{\hat{N}} = \tilde{\Gamma}_{MN}^N + \tilde{\Gamma}_{Mn}^n + \tilde{\Gamma}_{Ma}^a$. The 10 dimensional Ricci tensors are

$$\begin{aligned} \mathcal{R}_{MN} &= R_{MN} - [(p-q)\alpha_1 + (8-p+q)\alpha_2] \nabla_M \nabla_N A - \alpha_1 g_{MN} \nabla_P \nabla^P A \\ &\quad - [q\beta_1 + (8-p)\beta_2] \nabla_M \nabla_N B \\ &\quad + \left[(p-q)\alpha_1^2 + 2(8-p+q)\alpha_1\alpha_2 - (8-p+q)\alpha_2^2 \right] \partial_M A \partial_N A \\ &\quad - \left[(p-q)\alpha_1^2 + (8-p+q)\alpha_1\alpha_2 \right] g_{MN} (\partial A)^2 \\ &\quad + (\alpha_1 - \alpha_2) [q\beta_1 + (8-p)\beta_2] (\partial_M A \partial_N B + \partial_N A \partial_M B) \\ &\quad - \alpha_1 [q\beta_1 + (8-p)\beta_2] g_{MN} \partial_P A \partial^P B - \left[q\beta_1^2 + (8-p)\beta_2^2 \right] \partial_M B \partial_N B; \end{aligned} \quad (\text{A.4})$$

$$\begin{aligned} \mathcal{R}_{mn} = & - \left[\alpha_2 \nabla^2 A + \beta_1 \nabla^2 B + \left((p-q)\alpha_1\alpha_2 + (8-p+q)\alpha_2^2 \right) (\partial A)^2 \right. \\ & + \left. \left((p-q)\alpha_1\beta_1 + (8-p+2q)\alpha_2\beta_1 + (8-p)\alpha_2\beta_2 \right) \partial A \partial B \right. \\ & + \left. \left(q\beta_1^2 + (8-p)\beta_1\beta_2 \right) (\partial B)^2 \right] e^{(-2\alpha_1+2\alpha_2)A+2\beta_1 B} \delta_{mn}; \end{aligned} \quad (\text{A.5})$$

$$\begin{aligned} \mathcal{R}_{ab} = & (7-p)\gamma_{ab} - \left[\alpha_2 \nabla^2 A + \beta_2 \nabla^2 B + \left((p-q)\alpha_1\alpha_2 + (8-p+q)\alpha_2^2 \right) (\partial A)^2 \right. \\ & + \left. \left((p-q)\alpha_1\beta_2 + q\alpha_2\beta_1 + (16-2p+q)\alpha_2\beta_2 \right) \partial A \partial B \right. \\ & + \left. \left(q\beta_1\beta_2 + (8-p)\beta_2^2 \right) (\partial B)^2 \right] e^{(-2\alpha_1+2\alpha_2)A+2\beta_2 B} \gamma_{ab} L_p^2. \end{aligned} \quad (\text{A.6})$$

From the above we can get the Ricci scalar as

$$\begin{aligned} \mathcal{R} = & \frac{(7-p)(8-p)}{L_p^2} e^{-2\alpha_2 A - 2\beta_2 B} + e^{-2\alpha_1 A} \left[R - 2((p-q+1)\alpha_1 + (8-p+q)\alpha_2) \nabla^2 A \right. \\ & - 2(q\beta_1 + (8-p)\beta_2) \nabla^2 B - \left((p-q)(p-q+1)\alpha_1^2 + 2(p-q)(8-p+q)\alpha_1\alpha_2 \right. \\ & + \left. (8-p+q)(9-p+q)\alpha_2^2 \right) (\partial A)^2 - 2(q(q-p)\alpha_1\beta_1 + (8-p)(p-q)\alpha_1\beta_2 \\ & + q(9-p+q)\alpha_2\beta_1 + (8-p)(9-p+q)\alpha_2\beta_2) \partial A \partial B \\ & \left. - \left(q(q+1)\beta_1^2 + 2q(8-p)\beta_1\beta_2 + (8-p)(9-p)\beta_2^2 \right) (\partial B)^2 \right]. \end{aligned} \quad (\text{A.7})$$

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