

Determining the CP nature of spin-0 mediators in associated production of dark matter and $t\bar{t}$ pairs

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ABSTRACT: In the framework of spin-0 s -channel simplified models, we explore the possibility of assessing the structure of dark matter interactions through the associated production of dark matter and $t\bar{t}$ pairs. To this purpose, final states with two leptons are considered and the kinematic properties of the dilepton system is studied. We develop a realistic analysis strategy and provide a detailed evaluation of the achievable sensitivity for the dark matter signal assuming integrated luminosities of 300 fb^{-1} and 3 ab^{-1} at the 14 TeV LHC. Furthermore, upper limits on the mediator masses for which the two different CP hypotheses can be distinguished are derived. The obtained limits on the signal strengths are finally translated into constraints on the parameter space of two spin-0 simplified models including a scenario with an extended Higgs sector.

KEYWORDS: Dark matter, Hadron-Hadron scattering (experiments)

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1 Introduction

Searches for dark matter (DM) particles constitute a key part of the physics programme at the LHC. The minimal experimental signature of DM production at a hadron collider consists of an excess of events with a final-state object X recoiling against large amounts of missing transverse energy ($E_{\text{T}}^{\text{miss}}$). In LHC Run I and II, the ATLAS and CMS collaborations have examined a variety of such mono- X signatures involving jets of hadrons, photons, electroweak (EW) gauge bosons, top and bottom quarks as well as the Higgs boson in the final state. The cross section limits obtained from these $E_{\text{T}}^{\text{miss}}$ searches have been interpreted in the context of three different classes of theories: ultraviolet complete models, effective field theories [1–6] and to an increasing degree also simplified models [3, 5, 7–20]. In this article, we focus on the third class of theories which has the advantage of introducing only a minimal number of additional parameters compared to the SM, while evading the limitations that DM effective field theories face at the LHC. Specifically we will study simplified scalar and pseudoscalar mediator models which allow for s -channel DM pair production.

New spin-0 DM mediators coupling to the SM fermions necessarily carry SM flavour quantum numbers or conversely break the SM flavour symmetry. Extra sources of flavour breaking are however severely restricted since most flavour measurements agree well with the corresponding SM predictions. The simplest explanation of this empirical fact is to assume that the minimal flavour breaking consistent with the observed fermion hierarchy is also realised beyond the SM. This assumption is known as the minimal flavour violation (MFV) hypothesis [21] and implies that the couplings between any new neutral spin-0 state and SM matter are proportional to the fermion masses. One is thus naturally led to consider DM spin-0 mediators that couple most strongly to the third generation. Similar

to the SM Higgs, such resonances can be produced at the LHC through loop-induced gluon fusion or in association with top (or bottom) quarks before they decay to the heaviest kinematically allowed SM final state or DM.

The two main channels that have been used up to now at the LHC to search for scalars and pseudoscalars with large invisible decay widths are $pp \rightarrow \text{jets} + E_{\text{T}}^{\text{miss}}$ [22–29] and $pp \rightarrow t\bar{t} + E_{\text{T}}^{\text{miss}}$ [26, 28–34]. While for realistic cuts the signal rate is larger in the first case, one needs to study two-jet correlations [24, 35, 36] to gain information on the structure of the portal interactions. Associated production of DM and $t\bar{t}$ pairs is also sensitive to the nature of the mediator couplings through the kinematics of the top quarks, and through their polarisation which is accessible by studying the top-decay products. The fully leptonic channel with both W from $t \rightarrow bW$ decaying to $\ell\nu_{\ell}$ has compared to $pp \rightarrow jj + E_{\text{T}}^{\text{miss}}$ the advantage of a clean final state, and the kinematic correlations between the two charged leptons can be used as a CP analyser of the underlying $pp \rightarrow t\bar{t} + E_{\text{T}}^{\text{miss}}$ process. An interesting variable that has been proposed for separating $t\bar{t}$ production in the SM from new physics is the distance of the two leptons in the plane transverse to the beam direction (azimuthal plane) $\Delta\phi_{\ell\ell}$. This observable has been studied in some detail in the recent paper [34] and found to have promising discriminating power to disentangle scalar from pseudoscalar mediators.¹

In this work an alternative search strategy is devised which exploits the properties of $\cos\theta_{\ell\ell} \equiv \tanh(\Delta\eta_{\ell\ell}/2)$ where $\Delta\eta_{\ell\ell}$ is the difference in pseudorapidity of the two charged leptons. This variable has been introduced in [38] as a proxy of the polar angle for the production of a pair of hypothetical particles both of which decay into a visible and an invisible particle. In the case of the $pp \rightarrow t\bar{t} + E_{\text{T}}^{\text{miss}}$ signal, we find that in contrast to the $\Delta\phi_{\ell\ell}$ distribution used in [34], the $\cos\theta_{\ell\ell}$ distribution is affected only in a minor way by the selections in the transverse plane that are needed to suppress the dominant SM background from top-quark final states to manageable levels. The variable $\cos\theta_{\ell\ell}$ is thus a promising observable when analysing $t\bar{t} + E_{\text{T}}^{\text{miss}}$ events since it provides good CP-discriminating power also when experimental selection cuts are employed. Based on it we develop a realistic strategy for the detection of the DM signal at future LHC runs and determine the exclusion limits on the signal models using both a simple cut-and-count experiment as well as a shape-fit analysis. For the selection criteria adopted in this study the discovery reach in the $t\bar{t} + E_{\text{T}}^{\text{miss}}$ channel depends sensitively on the systematic uncertainty assigned to the estimate of SM backgrounds, which are dominated by $t\bar{t}Z$. A good experimental understanding of this process is hence a prerequisite to exploit the large data sets that the LHC is expected to provide at high luminosity (HL-LHC). We also derive upper limits on the mediator masses for which the two different CP hypotheses can be distinguished and translate the obtained model-independent limits on the signal strengths into constraints on the parameter space of two specific simplified DM models.

This article is structured as follows. In section 2 we discuss the relevant spin-0 interactions, while the basic properties of the $t\bar{t} + E_{\text{T}}^{\text{miss}}$ signal are studied in section 3. A description of our Monte Carlo (MC) simulations is presented in section 4. This part

¹Other observables such as $\Delta\eta_{\ell\ell}$ and $\Delta\eta_{b\bar{b}}$ [28, 37] as well as $m_{t\bar{t}}$ [29] have been proposed as probes of the CP nature of the DM $t\bar{t}$ interactions, but a detailed evaluation of the achievable sensitivity for the DM signal has not been performed in these cases.

contains a discussion of the signal and background generation and explains how physics objects are built in our analysis. The actual analysis strategy is presented in section 5 spelling out all selection criteria and illustrating their impact on the DM signal and the SM backgrounds. In section 6 we present the numerical results of our analysis providing a detailed evaluation of the achievable sensitivity for the $t\bar{t} + E_T^{\text{miss}}$ signature at future LHC runs. The constraints on the parameter space of two different spin-0 simplified DM models are also derived in this section. Our conclusions and a brief outlook are given in section 7.

2 Simplified models

The simplified DM models we are considering in our article can be described by the following interactions

$$\begin{aligned}\mathcal{L}_\phi &\supset -g_\chi \phi \bar{\chi} \chi - \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} g_q y_q \bar{q} q, \\ \mathcal{L}_a &\supset -i g_\chi a \bar{\chi} \gamma_5 \chi - i \frac{a}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} g_q y_q \bar{q} \gamma_5 q.\end{aligned}\tag{2.1}$$

Here ϕ is a scalar while a is a pseudoscalar, χ represents the DM particle which we assume to be a Dirac fermion, g_χ is a dark sector Yukawa coupling and $y_q = \sqrt{2} m_q / v$ are the SM quark Yukawa couplings with $v \simeq 246$ GeV the Higgs vacuum expectation value (VEV).

Notice that the simplified models (2.1) are valid descriptions of the physics below the EW scale as long as the new scalar ϕ does not mix strongly with the SM Higgs boson. In such a case the model dependence associated to the full scalar sector is captured by the portal couplings g_χ and g_q . The simplest choice of quark couplings compatible with the MFV hypothesis is universal $g_q = g_v$ and realised for instance in the spin-0 simplified models recommended by the ATLAS/CMS DM Forum (DMF) [20]. If the SM Higgs sector is extended to a two Higgs doublet model (2HDM) other MFV coupling patterns are however possible. For instance within the alignment/decoupling limit of 2HDM of type II (2HDMII) plus singlet extensions the relevant couplings are given by (see for example [39])

$$g_\chi = y_\chi \cos \theta, \quad g_{u,c,t} = \pm \sin \theta \cot \beta, \quad g_{d,s,b} = -\sin \theta \tan \beta.\tag{2.2}$$

Here y_χ denotes the coupling of the singlet to DM and θ is the mixing angle between the singlet and the heavy CP-even or CP-odd Higgs state and the $+$ ($-$) sign in $g_{u,c,t}$ holds for a scalar (pseudoscalar). Furthermore $\tan \beta$ represents the ratio of VEVs of the two Higgs doublets and for small $\tan \beta$ one obtains the hierarchy $g_t y_t \gg g_q y_q$ with $q = u, d, s, c, b$.

3 Anatomy of the signal

In figure 1 we show the production cross section for $pp \rightarrow t\bar{t} + E_T^{\text{miss}}$ (left panel) and the corresponding fraction of events arising from gluon fusion (right panel) for a centre-of-mass energy (\sqrt{s}) of 14 TeV. The displayed results have been obtained at next-to-leading order (NLO) with the help of `MadGraph5_aMC@NLO` [40] employing the `DMsimp` implementation [29] of the simplified models (2.1) and `NNPDF3.0` parton distribution functions (PDFs) [41]. From the left panel one observes that for very low mediator masses

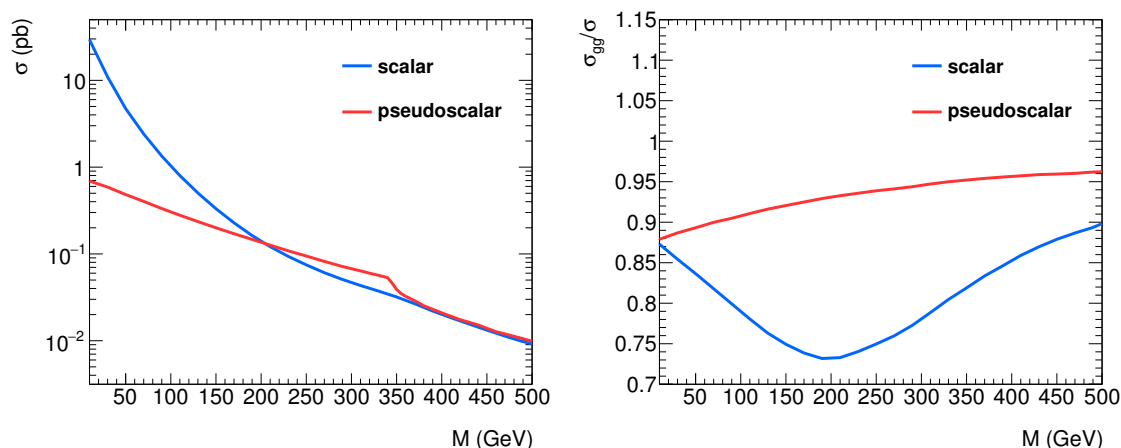


Figure 1. Left: total production cross section for $pp \rightarrow t\bar{t} + E_T^{\text{miss}}$ as a function of the mediator mass. Right: mediator mass dependence of the ratio of gluon-fusion production rate to the total production cross section. Both panels correspond to $\sqrt{s} = 14$ TeV, employ $m_\chi = 1$ GeV and $g_\chi = g_t = 1$ and assume a minimal decay width for the mediator. The predictions for a scalar (pseudoscalar) mediator are shown in blue (red).

$M = M_\phi$ or M_a the cross section associated to scalar exchange (blue curve) is larger than that for a pseudoscalar (red curve) by more than an order of magnitude. At around $M \simeq 200$ GeV the two predictions then become alike, while at higher masses the rate for pseudoscalar production is always slightly larger than that for a scalar. In the right plot, one sees that at the LHC the gluon-fusion channel is the dominant production mode independently of the CP nature of the mediator and amounts to roughly 85% of the total cross section for $M \simeq 10$ GeV. The functional dependence of σ_{gg}/σ is however different in the two cases. While in the CP-even case the fraction of gluon-fusion initiated events first decreases until about $M \simeq 200$ GeV and then starts rising, in the case of the CP-odd mediator the ratio σ_{gg}/σ is a steadily increasing function of M .

The features observed in figure 1 can be understood qualitatively in terms of two physical effects [29]. The first effect is related to the fact that a spin-0 state which has a mass much lighter than all of the relevant energy scales in a process $pp \rightarrow X$ can be treated as a parton which is radiated off the individual particles in the final state X . The process $pp \rightarrow t\bar{t} + \phi/a$ ($\phi/a \rightarrow \chi\bar{\chi}$) can thus be thought as $pp \rightarrow t\bar{t}$ followed by the radiation of ϕ/a from the final-state heavy quark lines with a subsequent decay of the spin-0 mediator to DM. This procedure is guaranteed to correctly reproduce the collinear divergencies associated with the emission of a massless ϕ/a state. The observed radiation pattern is determined by the leading (universal) fragmentation function $f_{t \rightarrow \phi/a}(x)$ which take the form [42, 43]

$$\begin{aligned}
 f_{t \rightarrow \phi}(x) &= \frac{g_t^2}{(4\pi)^2} \left[\frac{4(1-x)}{x} + x \ln \left(\frac{s}{m_t^2} \right) \right], \\
 f_{t \rightarrow a}(x) &= \frac{g_t^2}{(4\pi)^2} \left[x \ln \left(\frac{s}{m_t^2} \right) \right],
 \end{aligned}
 \tag{3.1}$$

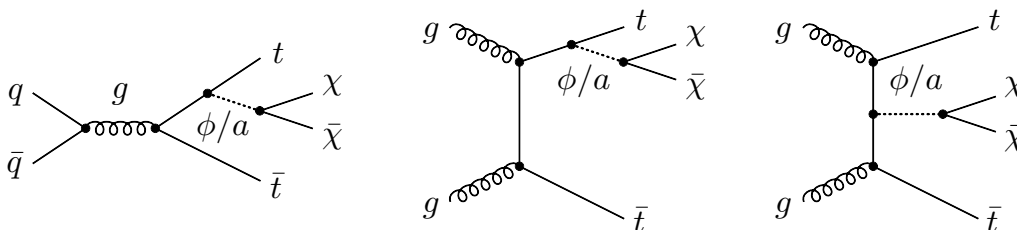


Figure 2. Examples of LO diagrams that give rise to a $t\bar{t} + E_{\text{T}}^{\text{miss}}$ signature through the exchange of a colourless spin-0 mediator. In the quark-fusion channel (left) only contributions from mediator fragmentation appear, while in the case of the gluon-fusion channel both mediator-fragmentation (center) and top-fusion (right) diagrams are present.

in the simplified models described by (2.1). These results are valid for $s \gg 4m_t^2 \gg M^2$ and $\ln(s/m_t^2) \ll 1$ where $\sqrt{s} = 2E/x$ with E the energy of the emitted spin-0 particle. From (3.1) one sees that due to the soft singularity proportional to $1/x$ a light scalar is radiated off top quarks preferentially with small energy (or equivalent small momentum fraction x). The soft-enhanced term is instead absent in the case of the pseudoscalar mediator. These features explain the order of magnitude difference between the total rates of the scalar and pseudoscalar mediators for masses $M \ll 2m_t$.

The second important difference between $\sigma(pp \rightarrow t\bar{t} + \phi)$ and $\sigma(pp \rightarrow t\bar{t} + a)$ with ϕ and a subsequently decaying to DM can be understood by considering the spin-averaged and colour-averaged squared matrix elements for the production of an on-shell spin-0 state with mass $M = \sqrt{s}$ from a top-quark pair. The corresponding squared matrix elements are given by

$$\overline{\sum} |\mathcal{M}(t\bar{t} \rightarrow \phi)|^2 = \frac{g_t^2 s}{12} \beta^2, \quad \overline{\sum} |\mathcal{M}(t\bar{t} \rightarrow a)|^2 = \frac{g_t^2 s}{12}, \quad (3.2)$$

with $\beta = \sqrt{1 - 4m_t^2/s}$ the velocity of the top quarks in the top-pair rest frame. From the above formulas one observes that close to the $t\bar{t}$ threshold located at $4m_t^2$ the production of a scalar in top-fusion is compared to that of a pseudoscalar suppressed by two powers of β . It follows that in cases where either the DM pair or the mediator is produced close to threshold, the production cross section of the pseudoscalar mediator is expected to be larger than that of a scalar. This is precisely what one observes in the left panel of figure 1. As it leads to a pronounced kink in the pseudoscalar case, the opening of the $t\bar{t}$ threshold is also clearly visible in this plot. The threshold suppression of $t\bar{t} \rightarrow \phi$ production finally explains the M dependence of the ratio σ_{gg}/σ with a dip at $M \simeq 200$ GeV as shown on the right in the latter figure.

The above observations can also be used to identify which leading order (LO) diagrams give the dominant contribution to the $t\bar{t} + E_{\text{T}}^{\text{miss}}$ signature in the case of a scalar or pseudoscalar. Representative examples of the three possible tree-level topologies are shown in figure 2. From the previous discussion it should be clear that at the LHC the $\sigma(pp \rightarrow t\bar{t} + \phi (\phi \rightarrow \chi\bar{\chi}))$ cross section is dominated by the gluon-fusion graph with a mediator fragmentation shown in the centre of the latter figure. In the case of the pseudoscalar

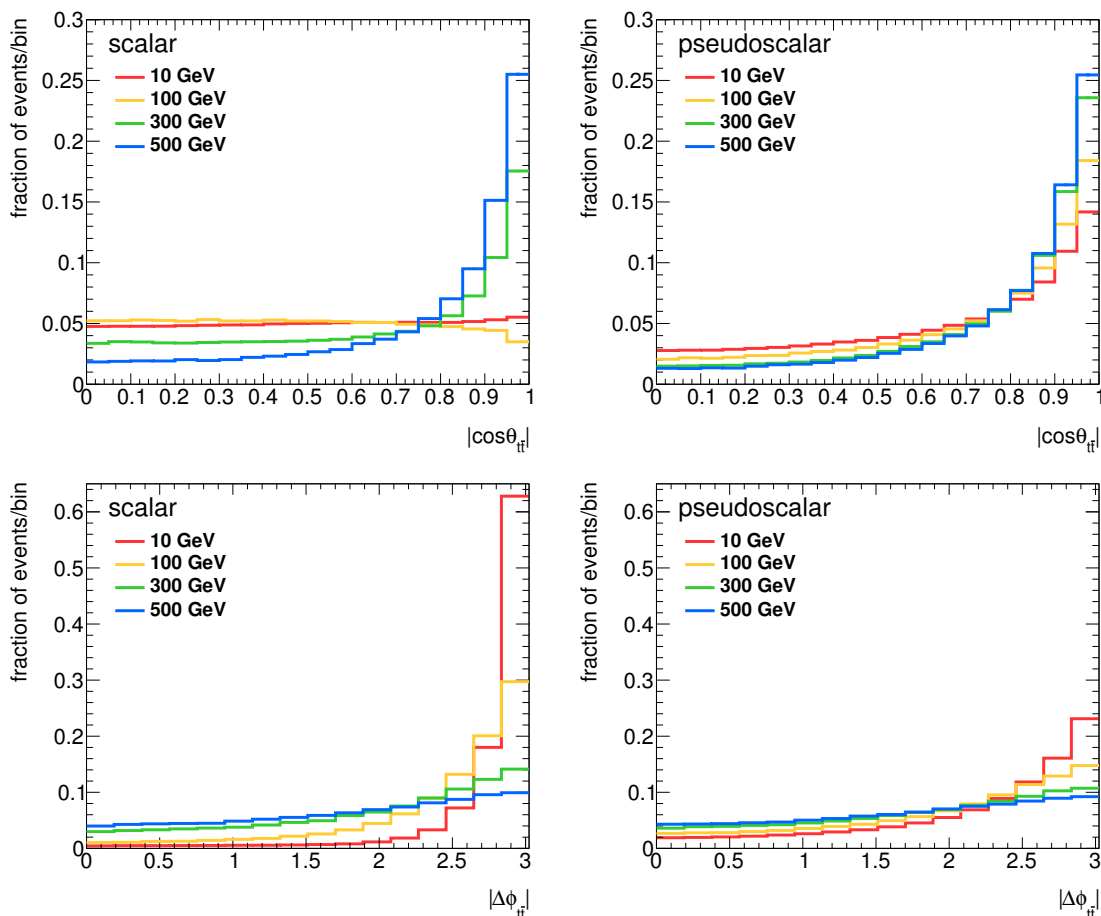


Figure 3. Normalised distributions of the $|\cos\theta_{t\bar{t}}|$ (upper row) and $|\Delta\phi_{t\bar{t}}|$ (lower row) variables for four different scalar (left) and pseudoscalar (right) benchmark models. The red, yellow, green and blue curves correspond to mediator masses of 10 GeV, 100 GeV, 300 GeV and 500 GeV, respectively.

cross section $\sigma(pp \rightarrow t\bar{t} + a (a \rightarrow \chi\bar{\chi}))$, on the other hand, both mediator-fragmentation and top-fusion diagrams in gluon-fusion are relevant. The latter contribution is displayed on the right in figure 2.

In the case of $t\bar{t} + E_T^{\text{miss}}$ production information on the CP nature of the coupling between the mediator and top quarks is encoded in the correlations between the final-state top quarks and their decay products. The two variables that we will study in this section are the $\cos\theta_{t\bar{t}} \equiv \tanh(\Delta\eta_{t\bar{t}}/2)$ variable and the azimuthal angle difference $\Delta\phi_{t\bar{t}}$ of the $t\bar{t}$ system. In figure 3 we present NLO predictions for the $|\cos\theta_{t\bar{t}}|$ and $|\Delta\phi_{t\bar{t}}|$ distributions for four different realisations of the scalar and pseudoscalar models introduced in (2.1). The shown spectra are all normalised to unity. Comparing the two panels in the upper row of the figure, one observes that for a light mediator of 10 GeV (red curves) and 100 GeV (yellow curves) the normalised $|\cos\theta_{t\bar{t}}|$ distributions are almost flat in the scalar case, while they are enhanced toward larger values of $|\cos\theta_{t\bar{t}}|$ if the mediator is a pseudoscalar. This feature can be understood by recalling that as a result of the soft singularity of the frag-

mentation function $f_{t \rightarrow \phi}(x)$ in (3.1) emissions of soft ϕ fields are favoured. Compared to a pseudoscalar, a scalar of the same mass hence tends to be produced more forward, which in turn means that the accompanying top quarks are produced more central. In consequence the difference in pseudorapidity of the top-antitop pair $\Delta\eta_{t\bar{t}}$ and likewise $\cos\theta_{t\bar{t}}$ are expected to be smaller on average for a light scalar than a pseudoscalar. For a heavier mediator of mass 300 GeV (green curves) and 500 GeV (blue curves) the soft enhancement of scalar fragmentation is no longer active so that the $|\cos\theta_{t\bar{t}}|$ spectra are peaked at 1 irrespectively of the CP nature of the mediator.

Similar arguments can be used to qualitatively explain the results for the $|\Delta\phi_{t\bar{t}}|$ distributions as shown in the lower panels of figure 3. One first notices that all spectra peak at the maximum allowed angle of π . However, for pseudoscalar mediators the distributions at large $|\Delta\phi_{t\bar{t}}|$ values are flatter in comparison to the scalar cases, in particular for the low mass benchmarks. This can be understood as follows. In the CP-even case soft emissions are favoured, and consequentially the top and antitop quarks prefer to fly in opposite directions in the transverse plane, which corresponds to large $|\Delta\phi_{t\bar{t}}|$ values. Because relative to the scalar case the emission of a pseudoscalar tends to be harder, for a CP-odd state the top quark and the mediator instead end up in a back-to-back position in the transverse plane and as a result the $|\Delta\phi_{t\bar{t}}|$ spectrum flattens. The impact of soft effects is less pronounced for heavier mediators for which the $|\Delta\phi_{t\bar{t}}|$ distributions become all very similar.

The above discussion should have made clear that the $t\bar{t}$ system can be used as an analyser of the CP properties of the mediation mechanism in $t\bar{t} + E_{\text{T}}^{\text{miss}}$ production. Given the presence of four invisible particles in the dilepton final state, the directions of the two top quarks are however experimentally not directly accessible. Indirection information on the relative orientation of the top quarks has hence to be obtained from measurements of the angular distributions of the charged lepton pairs resulting from top decays. The question is whether the discriminating power of angular variables such as $\cos\theta_{\ell\ell}$ and $\Delta\phi_{\ell\ell}$ survive the experimental cuts necessary to extract the DM signal from the SM backgrounds, and for which range of mediator masses and couplings it will be possible to perform the needed measurements with the foreseen integrated luminosity of the LHC. These issues will be addressed in the following two sections through detailed MC analyses.

4 MC simulations

In this section we provide a brief description of the MC simulations used to generate both the DM signal and the SM backgrounds and explain how electrons, muons, photons, jets and $E_{\text{T}}^{\text{miss}}$ are built in our detector simulation. Throughout our analysis we will consider pp collisions at $\sqrt{s} = 14$ TeV.

4.1 Signal generation

As in section 3 our signal samples are generated at NLO using the `DMsimp` simplified model implementation together with `MadGraph5_aMC@NLO` and `NNPDF3.0` PDFs. The final-state top quarks and W bosons are decayed with `MadSpin` [44] and the events are showered with `PYTHIA 8.2` [45] using the `FxFx` NLO jet matching prescription [46]. We consider five

different values of the mediator mass M , varying from 10 GeV to 500 GeV for both the case of a scalar and a pseudoscalar mediator. The mass of the DM particles is set to $m_\chi = 1$ GeV and we employ $g_\chi = g_t = 1$ for the mediator couplings to tops and to DM. The width of the mediator is assumed to be minimal and calculated at tree level using `MadGraph5_aMC@NLO`. Since in the narrow width approximation the signal prediction factorises into the cross section for $pp \rightarrow t\bar{t} + \phi/a$ production times the $\phi/a \rightarrow \chi\bar{\chi}$ branching ratio, changing the mediator width only leads to a rescaling of the signal strength. The experimental acceptance is instead insensitive to the mediator width, and therefore it is sufficient to generate samples for a single coupling choice. The predictions for other couplings values can then be simply obtained by scaling with the corresponding invisible branching ratio of the mediator.

4.2 Background generation

In order to describe the $t\bar{t} + E_{\text{T}}^{\text{miss}}$ backgrounds accurately, SM processes involving at least two leptons coming from the decay of vector bosons are generated. Backgrounds either with fake electrons from jet misidentification or with real non-isolated leptons from the decay of heavy flavours are not considered in our analysis, as a reliable estimate of these backgrounds would require a simulation of detector effects beyond the scope of this work.² The backgrounds from $t\bar{t}$ [48], tW [49], WW , WZ and ZZ production [50, 51] were all generate at NLO with `POWHEG BOX` [52]. The jets + Z samples are generated at LO with `MadGraph5_aMC@NLO` and contain up to four jets. `MadGraph5_aMC@NLO` is also used to simulate the $t\bar{t}V$ backgrounds with $V = W, Z$ at LO with a multiplicity of up to two jets. All partonic events are showered with `PYTHIA 8.2`. The samples produced with `POWHEG BOX` are normalised to the NLO cross section given by the generator, except $t\bar{t}$ which is normalised to the cross section obtained at NNLO plus next-to-next-to-leading logarithmic accuracy [53, 54]. The jets + W/Z samples are normalised to the known NNLO cross sections [55, 56], while in the case of the $t\bar{t}V$ samples the NLO cross sections calculated with `MadGraph5_aMC@NLO` are used as normalisations.

4.3 Detector smearing

Starting from the stable particles in the generator output, the following physics objects are built: electrons, muons, photons, jets and $E_{\text{T}}^{\text{miss}}$. Jets are constructed from the true momenta of particles interacting in the calorimeters except muons using an anti- k_t algorithm [57] with a parameter $R = 0.4$, as implemented in `FastJet` [58]. The variable $\vec{p}_{\text{T}}^{\text{miss}}$ with magnitude $E_{\text{T}}^{\text{miss}}$ is defined at truth level, i.e. before any detector effects are applied, as the negative of the vector sum of the transverse momenta of all the invisible particles (neutrinos and DM particles in our case). The effect of the detector on the kinematic quantities of interest is simulated by applying a Gaussian smearing to the momenta of the different reconstructed objects and reconstruction efficiency factors. The parametrisation

²The ATLAS analysis [47] employs kinematic variables similar to the ones used in our study and finds that in the relevant signal region the background from non-prompt leptons is negligible. We take this as a strong indication that the background from non-prompt leptons is also not an issue in our case.

of the smearing and reconstruction efficiency as a function of momentum and pseudorapidity of the objects is tuned to mimic the performance reported by ATLAS for Run I at the LHC [59]. The discrimination of signal from background is strongly affected by the assumed experimental smearing of E_T^{miss} , which is the main handle to tame the enormous $t\bar{t}$ background. To this purpose, the transverse momenta of unsmeared jets, electrons and muons are subtracted from the truth E_T^{miss} and replaced by the corresponding smeared quantities. The residual truth imbalance is then smeared as a function of the scalar sum of the transverse momenta of the particles not assigned to jets or electrons.

5 Analysis strategy

The final state targeted in our study contains two leptons, two jets with b -hadrons (b -jets) from the decay of the top quarks as well as a significant amount of E_T^{miss} associated to both the DM particles and the neutrinos from the $W \rightarrow \ell\nu_\ell$ decays.

In order to understand which discriminators and kinematic variables are useful to separate signal and background, one first has to recall the different types of SM backgrounds that can resemble the feature of a $t\bar{t} + E_T^{\text{miss}}$ signal. In fact, the SM backgrounds can be divided into three distinct classes:

- 1) Top backgrounds: the final states $t\bar{t}$ and tW with one or two top quarks fall into this class. These backgrounds are characterised by significant jet activity from the top decays and the presence of b -jets. The two neutrinos from $W \rightarrow \ell\nu_\ell$ represent the dominant E_T^{miss} contribution. The $t\bar{t}$ and tW backgrounds can mimic the signal if extra E_T^{miss} arises either from neutrinos produced in heavy-quark decays or from jet-momenta mismeasurement in the detector.
- 2) Reducible backgrounds: the second class of backgrounds comprises events from WW , WZ , ZZ and jets + Z production where jets arise from QCD radiation. Such configurations have compared to the $t\bar{t} + E_T^{\text{miss}}$ signal on average less jet activity and often no b -tagged jets. In addition the amount of E_T^{miss} for these backgrounds is typically small compared to that of the DM signal.
- 3) Irreducible backgrounds: the final states $t\bar{t}V$ form the third class of backgrounds. Here the E_T^{miss} signature arises from the decays $W \rightarrow \ell\nu_\ell$ or $Z \rightarrow \nu_\ell\bar{\nu}_\ell$. These backgrounds are irreducible in the sense that the resulting final-state configurations of $t\bar{t}V$ and the $t\bar{t} + E_T^{\text{miss}}$ signal are very similar.

The natural variable for separating DM signals from SM backgrounds is E_T^{miss} , which in our case turns out to be very powerful in reducing both the top and the reducible backgrounds. Unfortunately, a selection in E_T^{miss} alone is not enough to fully tame the overwhelming top backgrounds, which exhibit long E_T^{miss} tails. A variable that is effective in suppressing all backgrounds where the two leptons are produced in the decay of two W bosons is the m_{T2} variable [60, 61]

$$m_{T2}^2 \left(\vec{p}_T^{\ell_i}, \vec{p}_T^{\ell_j}, \vec{p}_T^{\text{miss}} \right) \equiv \min_{\vec{q}_T^1 + \vec{q}_T^2 = \vec{p}_T^{\text{miss}}} \left\{ \max \left[m_T^2 \left(\vec{p}_T^{\ell_i}, \vec{q}_T^1 \right), m_T^2 \left(\vec{p}_T^{\ell_j}, \vec{q}_T^2 \right) \right] \right\}, \quad (5.1)$$

which can be calculated using the momenta $\vec{p}_T^{\ell_i}$ and $\vec{p}_T^{\ell_j}$ of the two leptons and \vec{p}_T^{miss} . In (5.1) the parameter m_T denotes the transverse mass and \vec{q}_T^1 and \vec{q}_T^2 are auxiliary vectors. The minimum is taken over all the possible choices of \vec{q}_T^1 and \vec{q}_T^2 which satisfy the equality $\vec{q}_T^1 + \vec{q}_T^2 = \vec{p}_T^{\text{miss}}$. The m_{T2} variable has a sharp upper limit at m_W for W -boson induced backgrounds, while for the $t\bar{t} + E_T^{\text{miss}}$ signal the presence of extra E_T^{miss} leads to a tail in the m_{T2} distribution. All the other backgrounds, in particular those including the leptonic decays of Z bosons have rapidly falling E_T^{miss} and m_{T2} distributions, and can therefore be strongly suppressed by appropriate cuts on these variables.

The first step in the analysis is the preselection. Events pass the preselection only if they have exactly two isolated oppositely charged leptons (electrons, muons or one of each flavour) with $p_T^{\ell_1} > 25$ GeV, $p_T^{\ell_2} > 20$ GeV, $|\eta_\ell| < 2.5$ and an invariant mass that satisfies $m_{\ell\ell} > 20$ GeV. The η_ℓ and p_T^ℓ requirements ensure that leptons are reconstructed with high efficiency. If the charged signal leptons are of the same flavour the additional requirement $m_{\ell\ell} \in [71, 111]$ GeV is imposed to veto events where the charged lepton pair arises from a $Z \rightarrow \ell^+\ell^-$ decay.

A further set of cuts aims at a basic reduction of the top and reducible backgrounds. From an inspection of the topology of the top events with both high E_T^{miss} and m_{T2} , we find that the $t\bar{t}$ system often recoils against a high- p_T jet. For events in which the leading jet (j_1) comes from the decay of a top, the minimum invariant mass of j_1 with the two leptons $m_{j_1\ell\ell}^{\text{min}}$ in the event has to be lower than about 150 GeV. As an additional requirement we thus impose $m_{j_1\ell\ell}^{\text{min}} < 150$ GeV which rejects events with a high- p_T jet produced by QCD radiation, thereby suppressing both the top and reducible backgrounds. The $m_{j_1\ell\ell}^{\text{min}}$ cut has an efficiency of around 90% on the signal. In order to further suppress the reducible backgrounds to well below the level of the top backgrounds, the following requirements are employed. Events are required to have at least one b -jet with $p_T^j > 30$ GeV and all reconstructed jets with $p_T^j > 25$ GeV within $|\eta_\ell| < 2.5$ have to satisfy $\Delta\phi_{\text{min}} > 0.2$, where $\Delta\phi_{\text{min}}$ is defined to be the angle between \vec{p}_T^j and \vec{p}_T^{miss} for the jet closest to E_T^{miss} in the azimuthal plane. The latter requirement suppresses events where the E_T^{miss} is in part an artefact of jet mismeasurement.

The distribution of events in the $E_T^{\text{miss}}-m_{T2}$ plane after applying the first two sets of cuts is shown in figure 4 for the three classes of SM backgrounds and for a benchmark signal point. The signal prediction corresponds to a scalar mediator with mass $M_\phi = 100$ GeV and assumes $m_\chi = 1$ GeV and $g_\chi = g_t = 1$. The edge structure of the m_{T2} variable for the W -boson induced backgrounds naturally imposes the requirement $m_{T2} > 90$ GeV, whereas a selection $E_T^{\text{miss}} > 150$ GeV ensures the robustness of the analysis for backgrounds with instrumental E_T^{miss} . From the distributions of events in the upper right panel, one observes that imposing these selection criteria strongly suppresses the reducible backgrounds. To further reduce the top backgrounds, we construct the following linear combination from E_T^{miss} and m_{T2} :

$$C_{\text{em}} \equiv m_{T2} + 0.2 \cdot (200 \text{ GeV} - E_T^{\text{miss}}). \tag{5.2}$$

The C_{em} distribution after all other selections requirements have been applied is shown in figure 5 for the various backgrounds and our benchmark signal. The optimal cut on C_{em}

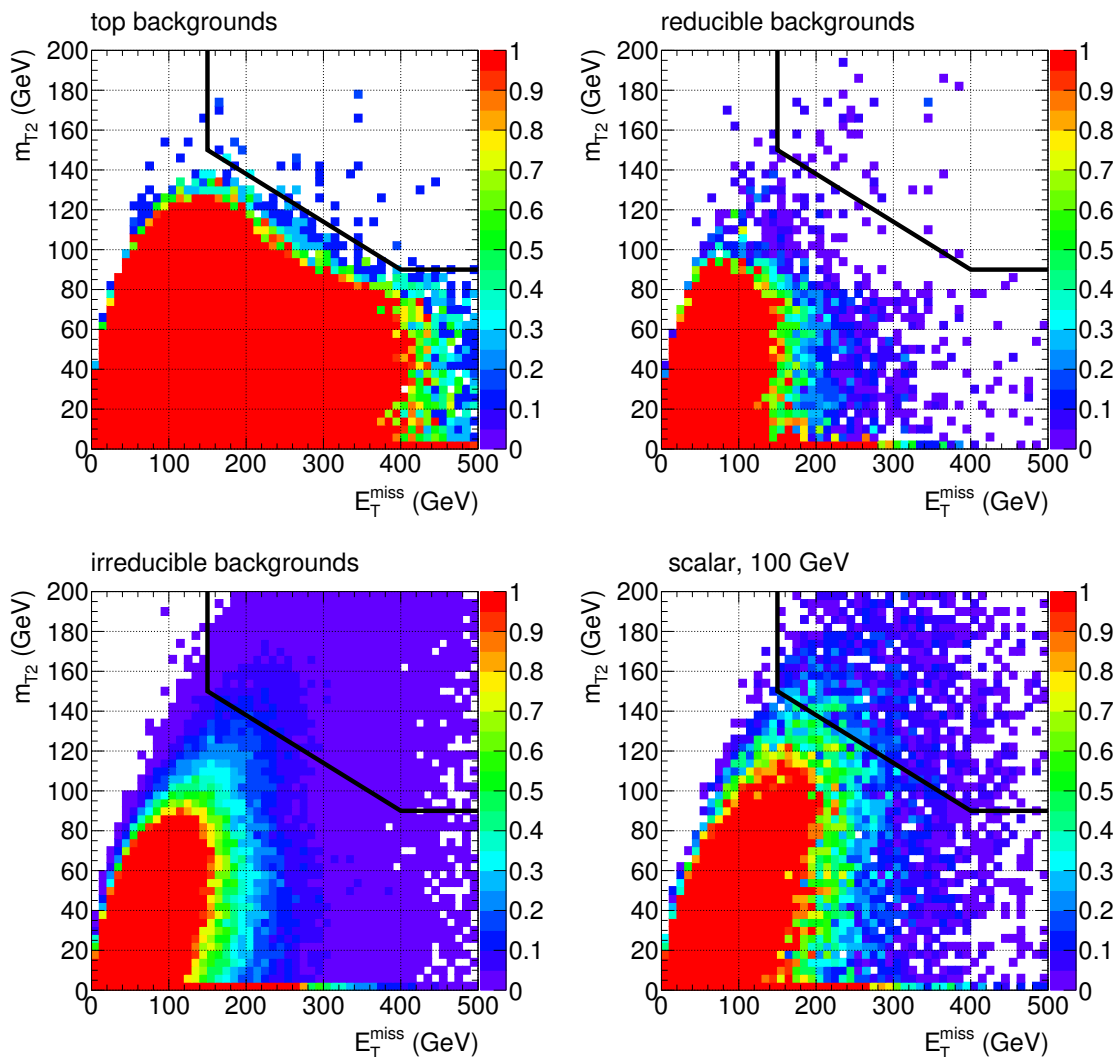


Figure 4. Distributions of events in the $E_T^{\text{miss}}-m_{T2}$ plane for the three different types of backgrounds and a signal point. The signal prediction corresponds to a scalar mediator with mass $M_\phi = 100$ GeV and assumes $m_\chi = 1$ GeV and $g_\chi = g_t = 1$. The selection cuts as defined in the text except E_T^{miss} , m_{T2} and C_{em} are imposed. The scale on the z axis is saturated at 1 event/bin/(100 fb $^{-1}$). The area in the upper right corner above the black line represents the signal region used in our analysis.

for the benchmark signal was established by minimising the value of the coupling $g = g_\chi = g_t$ which can be excluded at 95% CL for an integrated luminosity of 300 fb $^{-1}$, resulting in a requirement $C_{\text{em}} > 130$ GeV. It was checked explicitly that this requirement provides an adequate sensitivity over the whole considered range of mediator masses.

Figure 5 shows that the chosen criteria allows for an adequate reduction of the top backgrounds, while keeping an acceptable signal statistics for the considered model point. After applying the C_{em} cut the residual background is dominated by $t\bar{t}Z$ with subsequent $Z \rightarrow \nu_\ell \bar{\nu}_\ell$ decays. The area defined by the combined E_T^{miss} , m_{T2} and C_{em} re-

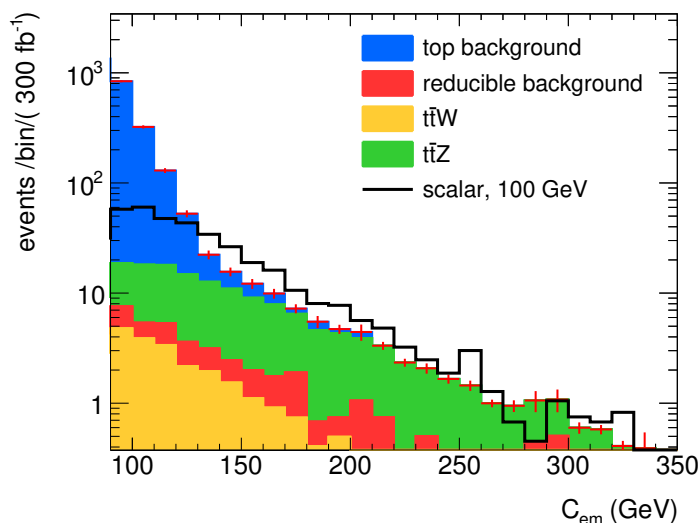


Figure 5. Distribution of the C_{em} variable after preselection, basic background suppression and imposing the E_T^{miss} and m_{T2} cuts as detailed in the text. The coloured histograms are stacked and represent the different SM backgrounds, while the black line corresponds to the C_{em} distribution as predicted by a scalar mediator with $M_\phi = 100$ GeV, $m_\chi = 1$ GeV and $g_\chi = g_t = 1$. All predictions have been obtained for 300 fb^{-1} of 14 TeV LHC data and the red error bars are the statistical errors of our background simulations.

quirements is indicated by the black lines in the panels of figure 4. One sees that after imposing all three selection criteria the huge majority of events arising from $t\bar{t}$ and tW are rejected. The total background for an integrated luminosity of 300 fb^{-1} is approximately 100 events of which 60% are from $t\bar{t}Z$, 20% are associated to the top backgrounds, 10% are from reducible backgrounds and 10% are due to $t\bar{t}W$. In the mass range from 10 GeV to 500 GeV the signal efficiency varies from 2‰ to 2% (6‰ to 2%) for scalar (pseudoscalar) mediators.

Having developed a realistic strategy for the detection of the DM signal, we can now come back to the angular variables and study to which extent the $\cos\theta_{\ell\ell}$ and $|\Delta\phi_{\ell\ell}|$ distributions are distorted by the selection requirements. To illustrate the impact of the selections, we show in figures 6 and 7 predictions for the angular distributions for eight different benchmark models before and after imposing the full set of selection cuts. From the panels in figure 6 one observes that if no cuts are imposed both $|\cos\theta_{\ell\ell}|$ and $|\Delta\phi_{\ell\ell}|$ provide sensitivity to the mass and the CP nature of the mediator particle. Even before applying any event selection the shape differences observed in $|\cos\theta_{\ell\ell}|$ are however more pronounced than those in $|\Delta\phi_{\ell\ell}|$.

The results presented in figure 7 demonstrate that applying the signal selections has a notable impact on the distributions of both angular observables. However, while for the $|\cos\theta_{\ell\ell}|$ variable the shape of the distributions still differ visibly for the chosen mediator masses and types, in the case of $|\Delta\phi_{\ell\ell}|$ almost all differences are washed out once experimental selections cuts are employed. Notice that the distortion of the $|\Delta\phi_{\ell\ell}|$ spectra is

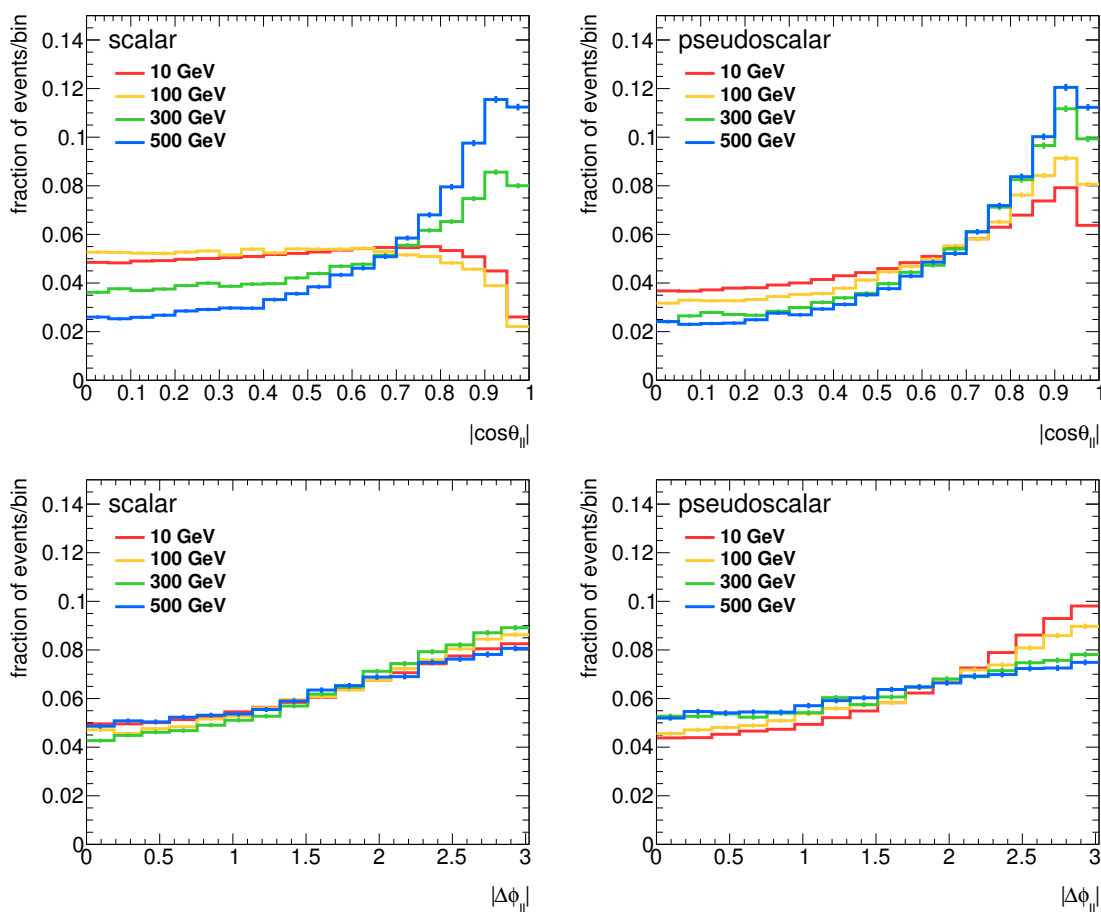


Figure 6. Normalised distributions of the $|\cos\theta_{\ell\ell}|$ (upper row) and $|\Delta\phi_{\ell\ell}|$ (lower row) variables for four different scalar (left) and pseudoscalar (right) benchmark models before imposing any selection. The style and colour coding of the curves follows the one of figure 3. The shown error bars are the statistical errors associated to our MC simulations.

most pronounced for values of $|\Delta\phi_{\ell\ell}|$ close to π , which corresponds to kinematic configurations where the two leptons are back-to-back in the azimuthal plane. The imposed m_{T2} cut typically removes such events thereby strongly reducing the discriminating power of the azimuthal angle difference $\Delta\phi_{\ell\ell}$.³ The observable $|\cos\theta_{\ell\ell}|$ instead still shows promising discriminating properties in a realistic experimental analysis with E_T^{miss} and m_{T2} cuts. We will therefore use this observables in the next section as our CP analyser.

6 Results

On the basis of the selection criteria defined in the previous section, we can study the LHC sensitivity to the $t\bar{t} + E_T^{\text{miss}}$ signature for an integrated luminosity of 300 fb^{-1} and

³The recent analysis [34] that uses $\Delta\phi_{\ell\ell}$ as a CP analyser claims that only a cut of $E_T^{\text{miss}} > 200 \text{ GeV}$ but no m_{T2} selection is needed to reduce the SM backgrounds to manageable levels. We are unable to reproduce this finding.

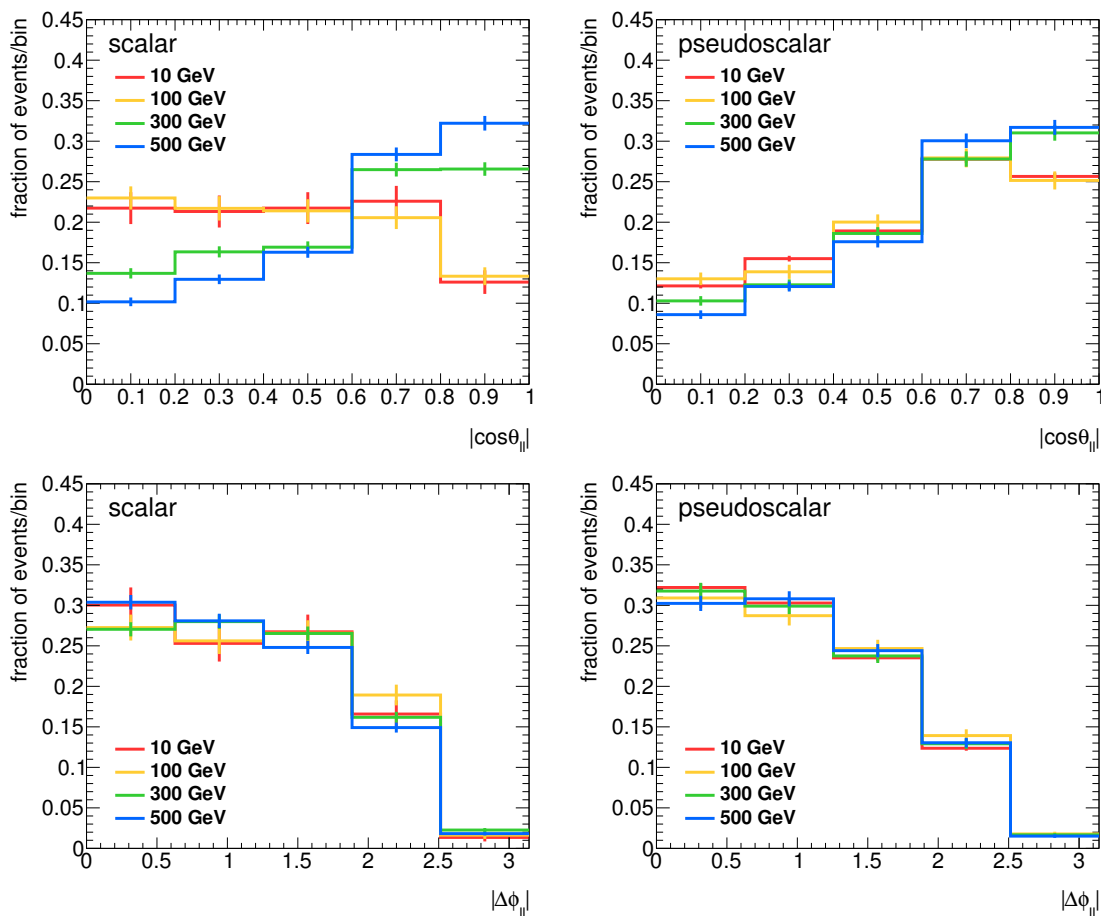


Figure 7. Normalised distributions of the $|\cos\theta_{\ell\ell}|$ (upper row) and $|\Delta\phi_{\ell\ell}|$ (lower row) variables for four different scalar (left) and pseudoscalar (right) benchmark models after imposing all selection requirements. The style and colour coding of the curves resembles the one of figure 3. The shown error bars are the statistical errors associated to our MC simulations.

3 ab^{-1} at $\sqrt{s} = 14 \text{ TeV}$. A profiled likelihood ratio test statistic is used to evaluate the upper limit on the ratio of the signal yield to that predicted in the DMF spin-0 simplified models with $m_\chi = 1 \text{ GeV}$ and $g_\chi = g_t = 1$. This ratio will be called signal strength in the following and denoted by μ . The CL_s method [62] is used to derive the confidence level (CL) of the exclusion limits and signal models with CL_s values below 0.05 are said to be excluded at 95% CL. The statistical analysis has been performed employing the `HistFitter` toolkit [63].

Our sensitivity study is performed in two ways. First by performing a simple counting experiment and second by including shape information in the form of a 5-bin likelihood fit to the $|\cos\theta_{\ell\ell}|$ distributions. The inclusion of shape information is motivated by the observation that the distributions of events as a function of the pseudorapidity difference of the dilepton pair is different for signal and background. This feature is illustrated in figure 8 which compares the predictions for a scalar (blue curve) and pseudoscalar (red curve) assuming $M = 100 \text{ GeV}$, $m_\chi = 1 \text{ GeV}$ and $g_\chi = g_t = 1$ with the SM background (black curve).

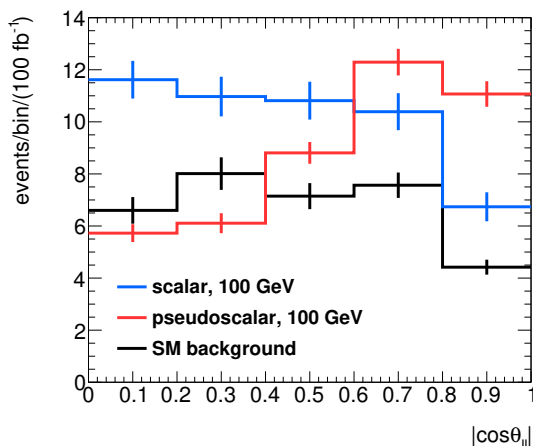


Figure 8. Distribution of the $|\cos\theta_{\ell\ell}|$ variable after employing the full selection requirements as specified in section 5. The normalisation corresponds to the numbers of events expected for 100 fb^{-1} at $\sqrt{s} = 14\text{ TeV}$. The error bars indicate the errors on the generated MC statistics.

Given the presence of a sizeable irreducible background surviving all the selections, the experimental sensitivity will be largely determined by the systematic uncertainty on the estimate of the SM backgrounds. Such an error has two main sources: on the one hand, uncertainties on the parameters of the detector performance such as the energy scale for hadronic jets and the identification efficiency for leptons, and on the other hand, uncertainties plaguing the MC modelling of SM processes. Depending on the process and on the kinematic selection, the total uncertainty can vary between a few percent and a few tens of percent. The present analysis does not select extreme kinematic configurations for the dominant $t\bar{t}Z$ background, and it therefore should be possible to control the experimental systematics at the 10% to 30% level. In the following, we will assume a systematic error of 20% on the backgrounds in the case of the counting experiment. In the case of the 5-bin shape fits we will consider background uncertainties of both 30% and 20%, fully correlated across the bins. We have checked that in the absence of an external measurement (e.g. a background control region) which profiles uncertainties, the use of correlated uncertainties in the shape fit provides the most conservative results.

The results of our sensitivity study are displayed in figure 9. Notice that the results shown for 3 ab^{-1} rely on the assumption that the $E_{\text{T}}^{\text{miss}}$ measurement performance in the very harsh experimental conditions of the HL-LHC will be equivalent to the one achieved during LHC Run I. As expected from the shapes of the distributions in figures 7 and 8, the 5-bin likelihood fit provides a significant improvement over the counting experiment for high-mass mediators irrespectively of their CP nature. The gain in sensitivity at lower mass depends strongly on the assumption on the systematic uncertainty of the SM background. For instance assuming a 20% systematics on the counting experiment and a 30% background error on the shape fit, we find that the shape analysis will have larger discriminating power than the simple cut-and-count strategy for $M_{\phi} \gtrsim 300\text{ GeV}$ and $M_a \gtrsim 100\text{ GeV}$ with 300 fb^{-1} of integrated luminosity. If the background for the shape fit can instead be

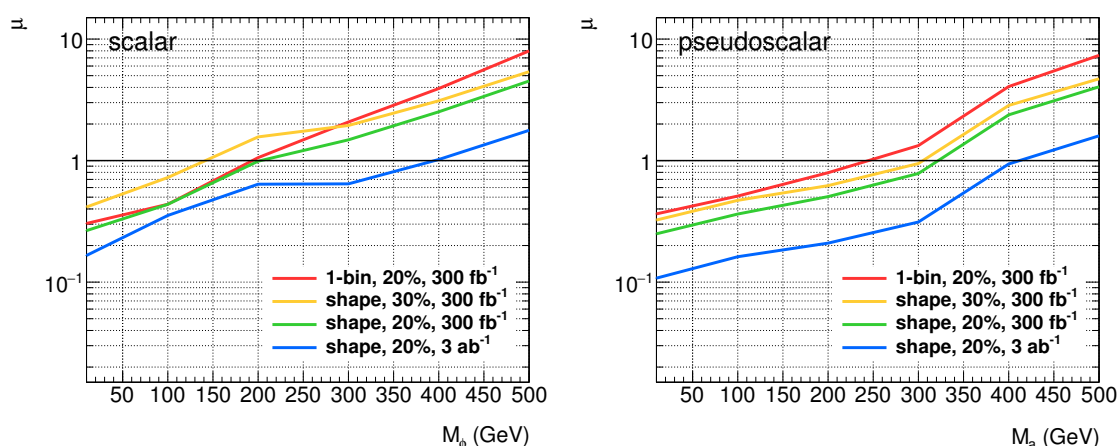


Figure 9. Value of the signal strength that can be excluded at 95% CL as a function of the mass for scalar (left) and pseudoscalar (right) mediators. The reach with 300 fb^{-1} of $\sqrt{s} = 14 \text{ TeV}$ data is given for a simple counting experiment assuming a 20% systematic background uncertainty (red curves) and for 5-bin shape fits with both 30% (yellow curves) and 20% (green curves) errors. A hypothetical shape-fit scenario based on 3 ab^{-1} and 20% systematics is also shown (blue curves).

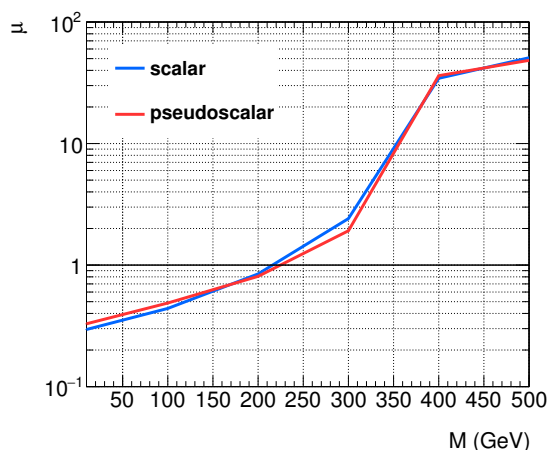


Figure 10. Value of the signal strength for which the scalar (pseudoscalar) hypothesis can be excluded at 95% CL in favour of the pseudoscalar (scalar) hypothesis. The reach is given for a shape fit using the $|\cos\theta_{\ell\ell}|$ variable, assuming 300 fb^{-1} of 14 TeV LHC data and a systematic uncertainty of 20% on the SM background.

estimated with an error of 20%, including shape information is expected to be the superior strategy over almost the entire range of considered masses. In fact, at the LHC with 3 ab^{-1} of data it should be possible to exclude spin-0 models that predict $\mu = 1$ for mediator masses up to around 400 GeV using the 5-bin likelihood fit employed in our study. The observed strong dependence of the reach on the assumption on the systematic background uncertainty shows that a good experimental understanding of $t\bar{t}Z$ production within the SM will be a key ingredient to a possible discovery of DM in the $t\bar{t} + E_T^{\text{miss}}$ channel.

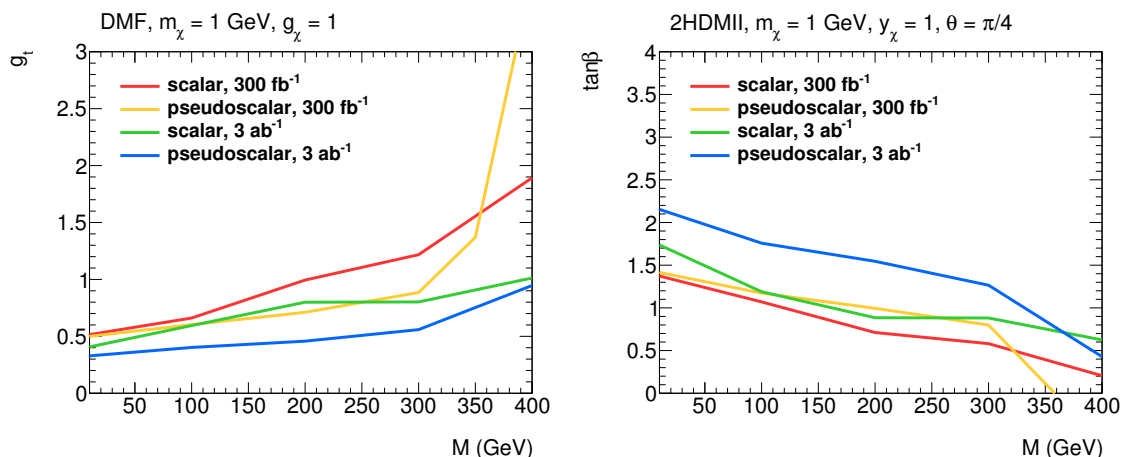


Figure 11. Left: value of the coupling g_t that can be excluded at 95% CL in the DMF spin-0 models. The shown limits correspond to $m_\chi = 1$ GeV and $g_\chi = 1$ and the parameter space above the coloured curves is ruled out. Right: value of $\tan\beta$ that can be excluded at 95% CL in the alignment/decoupling limit of the 2HDMII plus singlet models. The relevant model parameters are $m_\chi = 1$ GeV, $y_\chi = 1$ and $\theta = \pi/4$ and the exclusion holds for the parameter space to the bottom-left of the coloured curves. All limits have been obtained from our 5-bin shape-fit analysis assuming a systematic error of 20% on the SM background.

We also perform a hypothesis test between the scalar and pseudoscalar mediator hypotheses as a function of the mediator mass. Figure 10 shows the value of μ for which the scalar hypothesis can be excluded at 95% CL in favour of the pseudoscalar one (blue curve) and vice versa (red curve). Our statistical analysis is based on a 5-bin shape fit of the $|\cos\theta_{\ell\ell}|$ distributions and employs standard maximum likelihood estimator techniques (see for instance [64]) that are implemented in the RooFit/RooStat package [65]. From the figure it is evident that based on 300 fb^{-1} of $\sqrt{s} = 14 \text{ TeV}$ data and under the assumption that the SM backgrounds can be determined with an uncertainty of 20%, it should be possible to distinguish between the two CP hypotheses for masses $M \lesssim 200 \text{ GeV}$ in scenarios that lead to a signal strength of 1 or lower.

While we have assumed in our hypothesis test that M is already known, constraining the mediator mass, its couplings and CP nature simultaneously using our proposal seems feasible. For instance constraining both M and g_t should be possible by combining the information on the C_{em} variable (or m_{T2} observable) with that on the fiducial $t\bar{t} + E_T^{\text{miss}}$ cross section. The obtained constraints may be further strengthened by including besides the $t\bar{t} + E_T^{\text{miss}}$ channel other E_T^{miss} measurements such as mono-jets into a global fit. Such a fit would depend on the detailed consideration of additional search channels and of their experimental uncertainties, and is outside the the scope of the present study.

The model-independent limits on the signal strengths derived in figure 9 can be used to constrain the parameter space of the two classes of simplified DM models that have been introduced in section 2. In the left (right) panel of figure 11 we present the 95% CL exclusion bounds in the $M-g_t$ ($M-\tan\beta$) plane that apply in the case of the DMF (2HDMII plus singlet) models. The constraints are obtained from a shape fit with 20% background

uncertainties and the red, yellow, green and blue lines illustrate the scalar and pseudoscalar case assuming integrated luminosities of 300 fb^{-1} and 3 ab^{-1} . From the left panel we see that for the parameter choices $m_\chi = 1 \text{ GeV}$ and $g_\chi = 1$ values of $g_t \lesssim 1$ can be ruled out in the mass ranges $M_\phi \lesssim 200 \text{ GeV}$ and $M_a \lesssim 320 \text{ GeV}$ after 300 fb^{-1} of data have been collected. With ten times more luminosity these exclusions then extend up to mediator masses close to 400 GeV in the DMF models.

In the case of the 2HDMII plus singlet models in the alignment/decoupling limit, one sees from the right panel in figure 11 that with 300 fb^{-1} data model realisations with $\tan \beta \gtrsim 1$ can be excluded for $M_\phi \lesssim 120 \text{ GeV}$ and $M_a \lesssim 200 \text{ GeV}$. The corresponding choice of parameters is $m_\chi = 1 \text{ GeV}$, $y_\chi = 1$ and $\theta = \pi/4$. With 3 ab^{-1} of integrated luminosity the quoted mass limits are expected to improve to $M_\phi \lesssim 160 \text{ GeV}$ and $M_a \lesssim 330 \text{ GeV}$.

It is also apparent from both panels that below the $t\bar{t}$ threshold the limits on pseudoscalar models are always stronger than those on scalar scenarios. This feature can be understood by realising that our $|\cos \theta_{\ell\ell}|$ shape analysis has larger discriminating power for a CP-odd than for a CP-even spin-0 portal state, as one would expect from figure 8. Notice finally that above the $t\bar{t}$ threshold the constraints in the $M-g_t$ and $M-\tan \beta$ planes as depicted in figure 11 start to weaken because the branching ratios of $\phi/a \rightarrow \chi\bar{\chi}$ are no longer 100%. This feature is most pronounced in the case of our pseudoscalar scenario with 300 fb^{-1} of data. In this parameter space region ditop resonance searches can provide relevant constraints [66, 67] on both the DMF as well as the 2HDMII plus singlet models.

7 Conclusions

In this article, we have studied the prospects of future LHC runs to probe spin-0 interactions between DM and top quarks via the $t\bar{t} + E_T^{\text{miss}}$ signature. This final state is particularly interesting, since it is expected to have an appreciable rate in simplified s -channel scalar and pseudoscalar models that satisfy both the constraints from quark-flavour and Higgs physics. Examples of such models are provided by the spin-0 scenarios recommended by the ATLAS/CMS DMF and 2HDMII plus singlet extensions in the alignment/decoupling limit and low values of $\tan \beta$.

In order to understand which kinematic variables are useful to separate signal and SM backgrounds, we have first analysed the basic properties of the $t\bar{t} + E_T^{\text{miss}}$ signal. By identifying the dominant production topologies, qualitative explanations of the mediator mass and type dependence of the total production cross sections σ and their gluon-fusion fractions σ_{gg}/σ have been provided. It has also been shown that by dividing the signal into mediator-fragmentation and top-fusion diagrams allows one to understand the basic features of the $\cos \theta_{t\bar{t}}$ distributions and the azimuthal angle difference $\Delta\phi_{t\bar{t}}$ of the $t\bar{t}$ system. In all cases we have confirmed our general expectations with explicit MC simulations, finding that the angular correlations of the top-antitop pair provide in principle powerful probes of the CP properties of the mediation mechanism in $t\bar{t} + E_T^{\text{miss}}$ production.

Given the presence of four invisible particles in the dilepton final state, the relative orientation of the top quarks is however experimentally not directly accessible. Any infor-

mation on the $t\bar{t}$ system must thus be indirectly obtained from the kinematical distributions of the dilepton pair resulting from $t \rightarrow bW (W \rightarrow \ell\nu_\ell)$. In our work we have developed a realistic analysis strategy to enhance the small signal-to-background ratio of the $t\bar{t} + E_{\text{T}}^{\text{miss}}$ dilepton channel. We found that only a combination of selections based on the $E_{\text{T}}^{\text{miss}}$ and the $m_{\text{T}2}$ variables is able to suppress the overwhelming $t\bar{t}$ and tW backgrounds sufficiently. The impact of these cuts on the $\cos\theta_{\ell\ell}$ and $\Delta\phi_{\ell\ell}$ distributions has then been studied in detail. While the pseudorapidity difference of the dilepton pair was found to be affected by the selections in the transverse plane only in a minor way, the azimuthal angle difference was shown to lose most of its discriminating power due to the $m_{\text{T}2}$ requirement. In a realistic experimental search for a $t\bar{t} + E_{\text{T}}^{\text{miss}}$ signal, the $\Delta\phi_{\ell\ell}$ variable thus seems to be of limited use in probing the CP nature of the DM $t\bar{t}$ interactions.

Focusing on the $\cos\theta_{\ell\ell}$ observable we have then presented a comprehensive sensitivity study for the considered DM signal. We have derived the 95% CL exclusion limits on the signal strengths μ as a function of the mediator mass M that follow from both a counting experiment and a 5-bin shape fit to the $|\cos\theta_{\ell\ell}|$ distributions. Our analysis shows that including shape information generically allows to improve the reach of the $t\bar{t} + E_{\text{T}}^{\text{miss}}$ searches. The actual improvement however turns out to depend sensitively on the assumption about the systematic uncertainty on the irreducible SM backgrounds. In order to exploit the full potential of $t\bar{t} + E_{\text{T}}^{\text{miss}}$ searches it is thus crucial to refine the modelling of the $E_{\text{T}}^{\text{miss}}$ and $m_{\text{T}2}$ distributions in $t\bar{t}Z$ production. We have furthermore found that with of 300 fb^{-1} data and under the assumption that the SM backgrounds are known to 20%, it should be possible to distinguish between the two CP hypotheses for mediator masses up to around 200 GeV in all scenarios that lead to $\mu \lesssim 1$. Finally, we have presented 95% CL exclusion bounds in the $M-g_t$ and $M-\tan\beta$ planes which allow to illustrate the LHC reach in the context of the DMF spin-0 simplified models and 2HDMII plus singlet extensions, respectively. The upshot of our recast is that spin-0 simplified models that lead to an effective coupling strength of order 1 of the mediators to top-quark pairs can typically be tested for masses up to (or even above) the $t\bar{t}$ threshold at the HL-LHC.

The DM top-quark couplings of spin-0 s -channel simplified models can be probed at the LHC also in other channels such as $pp \rightarrow \text{jets} + E_{\text{T}}^{\text{miss}}$ and $pp \rightarrow t\bar{t}$. In both cases it should be possible to use angular correlations to either enhance the sensitivity of these searches or to determine the CP properties of the involved mediators. Since the $pp \rightarrow t\bar{t} + E_{\text{T}}^{\text{miss}}$, $pp \rightarrow \text{jets} + E_{\text{T}}^{\text{miss}}$ and $pp \rightarrow t\bar{t}$ probe different aspects of the underlying theory, we believe that it is crucial to combine all available information on the DM top-quark couplings to fully exploit the LHC potential. We look forward to further investigations in this direction.

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