# Homogenization of Attractors for Semilinear Parabolic Equations on Manifolds with Complicated Microstructure (*). 

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#### Abstract

An approach to a homogenized description of solutions of the Cauchy problem for parabolic equations on Riemannian manifolds with complicated microstructure is presented. This approach covers both linear and non-linear cases and makes it possible to establish a connection between global attractors of the initial problem of the homogenized one.


## 1. - Introduction.

We consider on an $n$-dimensional ( $n \geqslant 2$ ) Riemannian manifold $M_{\varepsilon}$ of complicated microstructure depending on $\varepsilon>0$ the following initial-boundary problem

$$
\begin{equation*}
\frac{\partial u^{\varepsilon}}{\partial t}-\Delta_{\varepsilon} u^{\varepsilon}+f\left(u^{\varepsilon}\right)=h^{\varepsilon}(x), \quad x \in M_{\varepsilon}, \quad t>0 \tag{1}
\end{equation*}
$$

$$
\frac{\partial u^{\varepsilon}}{\partial v_{\varepsilon}}=0, \quad x \in \partial M_{\varepsilon}, \quad t>0
$$

$$
\begin{equation*}
u^{\varepsilon}(x, 0)=u_{0}^{\varepsilon}(x) \tag{3}
\end{equation*}
$$

Here $\Delta_{\varepsilon}$ is the Laplace operator on $M_{\varepsilon}, \partial / \partial \nu_{\varepsilon}$ is the outer normal derivative on the boundary $\partial M_{\varepsilon}$ of $M_{\varepsilon}, f(u)$ is a smooth real function on $\boldsymbol{R}^{1}$ and $h^{\varepsilon}(x), u_{0}^{\varepsilon}: M_{\varepsilon} \rightarrow \boldsymbol{R}^{1}$ are given functions. We suppose that the local structure of the manifold $M_{\varepsilon}$ becomes more and more complicated, when $\varepsilon$ tends to zero.

This paper deal with the study of the asymptotic behaviour of the solution $u^{\varepsilon}(x, t)$

[^0]and of the global attractor $\mathfrak{G}_{\varepsilon}$ of problem (1)-(3) when $\varepsilon \rightarrow 0$. One of the main goals here is to learn how the transition to homogenized ( $\varepsilon \rightarrow 0$ ) description reflects on the long-time ( $t \rightarrow+\infty$ ) dynamics.

Under certain conditions on the manifold $M_{\varepsilon}$ and non-linear term $f(u)$ we first prove that for any finite time interval the limit behaviour of $u^{\varepsilon}(x, t)$ is described by a solution of the Cauchy problem for a system of two coupled equations. After that we study the long-time dynamics of this homogenized system and show that it possesses a finite-dimensional global attractor $\mathcal{G}$ (for definitions and basic facts on attractors see, e.g. $[1,4,7,17])$. We investigate the structure of $\mathfrak{G}$ and prove that global attractors $\mathcal{C}_{\varepsilon}$ tend to A in a suitable sense.

In the linear case $(f(u) \equiv 0)$ a similar homogenization problem has been studied in [2]. It has been proved that the asymptotic of $u^{\varepsilon}(x, t)$ is described by a linear diffusion equation with a term non-local in time. This term can be interpretated as memory of the medium (on the memory phenomena for linear homogenized models see also [11-14]). The method developed in [2] essentially relies on the linearity of the problem. The main ingredients there are the Laplace transformation in time and the study of the corresponding stationary problem by variational methods. Unlike [2] the approach presented here can be applied both to linear and non-linear cases. For the linear case the homogenized coupled system can be reduced to a single diffusion equation with memory term of the same form as in [2].

We also note that the dependence of attractors on parameters for various singularly perturbated systems has been studied by many authors (see, e.g. [1, 3, 5, 7, 8, 10, $16]$ and the references therein). In this paper we rely on some ideas presented in [3, 5 , 7, 8].

The paper is organized as follows. In Section 2 we describe the structure of the manifold $M_{\varepsilon}$ introduce some notations and give preliminary results concerning the properties of solutions of the problem (1)-(3), when $\varepsilon>0$ is fixed. In Section 3 we formulate our main results. The rest of the paper is devoted to the proofs of the Theorems of Section 3. Section 4 contains the proof of the estimates which guarantee the compactness of the family $\left\{u_{\varepsilon}: \varepsilon \rightarrow 0\right\}$. In Section 5 we make the limit transition in the weak form of problem (1)-(3). The main point here is to choose the testing function. In Section 6 we study properties of the homogenized and prove the upper semicontinuity of global attractor $\mathfrak{a}_{\varepsilon}$ of the problem (1)-(3), when $\varepsilon \rightarrow 0$.

## 2. - Preliminary consideration.

Now we describe the structure of the manifold $M_{\varepsilon}$. Let $\Omega$ be a smooth bounded domain in $\boldsymbol{R}^{n}(n \geqslant 2)$ and let

$$
F_{\varepsilon}=\bigcup_{j \in N_{\varepsilon}} F\left(x^{i}, a_{\varepsilon}\right)
$$

## L. Boutet de Monvel - I. D. Chueshov - E. Ya. Khruslov: Homogenization, etc.

be a union of balls $F\left(x^{i}, a_{\varepsilon}\right)$ of radius $a_{\varepsilon} \ll \varepsilon\left(\lim _{\varepsilon \rightarrow 0} a_{\varepsilon} \varepsilon^{-1}=0\right)$ with centers in $x^{j}=j \varepsilon$ $\left(j \in \boldsymbol{Z}^{n}\right)$ such that $F\left(x^{i}, a_{\varepsilon}\right) \in \Omega$. Here $N_{\varepsilon}$ stands for the corresponding set of multiindexes $j \in \boldsymbol{Z}^{n}$. In $\boldsymbol{R}^{n+1}$ we consider the surfaces (below $x=\left(x_{1}, \ldots, x_{n}\right) \in \boldsymbol{R}^{n}, y \in \boldsymbol{R}^{1}$, $\left.(x, y) \in \boldsymbol{R}^{n+1}\right):$

$$
\Omega_{\varepsilon}=\left\{(x ; 0) \in \boldsymbol{R}^{n+1}: x \in \Omega \backslash F_{\varepsilon}\right\}
$$

and

$$
B_{\varepsilon}^{j}=(j \varepsilon ; 0)+B_{\varepsilon}, \quad j \in N_{i} \subset \boldsymbol{Z}^{n}
$$

where

$$
B_{\varepsilon}=\left\{(x, y) \in \boldsymbol{R}^{n+1}:|x|^{2}+\left(y-\sqrt{b^{2} \varepsilon^{2}-a_{\varepsilon}^{2}}\right)^{2}=b^{2} \varepsilon^{2}, y \leqslant 0\right\} .
$$

Here $b$ is a parameter such that $a_{\varepsilon} \varepsilon^{-1}<b<1$. We assume that

$$
M_{\varepsilon}=\Omega_{\varepsilon} \cup\left(\bigcup_{j \in N_{\varepsilon}} B_{\varepsilon}^{j}\right),
$$

i.e. $M_{\varepsilon}$ consists of a piece of flat submanifold in $\boldsymbol{R}^{n+1}$ with bubbles $B_{\varepsilon}^{j}$. We define a Riemannian structure on $M_{\varepsilon}$ by a $C^{\infty}$ metric tensor

$$
g^{\varepsilon}(x)=\left\{g_{\alpha \beta}^{\varepsilon}(x) ; \alpha, \beta=1,2, \ldots, n\right\}, \quad x \in M_{\varepsilon},
$$

and assume the following:
(i) the metric coincides with the euclidean metric of $\boldsymbol{R}^{n+1}$ on $\Omega_{\varepsilon}$;
(ii) the metric is the same for all bubbles $B_{\varepsilon}^{j}, j \in N_{\varepsilon}$;
(iii) there exist positive constants $C_{1}$ and $C_{2}$ such that

$$
\begin{equation*}
C_{1} \varepsilon^{n}|\xi|^{2} \leqslant \sum_{\alpha \beta} g_{\alpha \beta}^{\varepsilon}(x) \xi_{\alpha} \xi_{\beta} \leqslant C_{2} \varepsilon^{n}|\xi|^{2}, \quad \varepsilon>0 \tag{2.1}
\end{equation*}
$$

for all $x \in B_{\varepsilon}^{j}, j \in N_{\varepsilon}$ andfor all $\xi \in \boldsymbol{R}^{n}$.
The main object of this paper is the problem (1)-(3) on the Riemannian manifold ( $M_{\varepsilon}, g^{\varepsilon}$ ), which can be treated as a model of diffusion in a medium with traps. The corresponding Laplace operator $\Delta_{\varepsilon}$ is of the form

$$
\Delta_{\varepsilon}=\frac{1}{\sqrt{\left|g^{\varepsilon}\right|}} \sum_{a, \beta} \frac{\partial}{\partial x_{a}}\left(\sqrt{\left|g^{\varepsilon}\right|} g_{\varepsilon}^{\alpha \beta} \frac{\partial}{\partial x_{\beta}}\right)
$$

where $\left|g^{\varepsilon}\right|=\operatorname{det} g^{\varepsilon}$ and $g_{\varepsilon}^{\alpha \beta}$ are the components of the inverse of the tensor $g^{\varepsilon}$. We also assume that the function $f(u) \in C^{2}\left(\boldsymbol{R}^{1}\right)$ possesses the property:

$$
\begin{equation*}
\sup \left\{\left|f^{\prime}(u)\right|: u \in \boldsymbol{R}^{1}\right\}<\infty \tag{2.2}
\end{equation*}
$$

and there exists a constant $\eta>0$ such that

$$
\begin{gather*}
u f(u) \geqslant \eta u^{2}-C_{1},  \tag{2.3}\\
\mathscr{F}(u) \equiv \int_{0}^{u} f(\xi) d \xi \geqslant \eta u^{2}-C_{2} . \tag{2.4}
\end{gather*}
$$

Below $d x$ represents the surface measure on $M_{\varepsilon}$. In local coordinates $\left\{x_{1}, \ldots, x_{n}\right\}$ we have $d x=\sqrt{\left|g^{\varepsilon}\right|} d x_{1} \ldots d x_{n}$. We also denote $H^{l}\left(V_{\varepsilon}\right)$ the Sobolev space of order $l$ on a submanifold $V_{\varepsilon} \subseteq M_{\varepsilon}$ and $H_{0}^{l}\left(V_{\varepsilon}\right)$ for closure of $C_{0}^{\infty}\left(V_{\varepsilon}\right)$ in $H^{l}\left(V_{\varepsilon}\right)$. We denote by $\|\cdot\|_{l, \varepsilon}$ the norm $H^{l}\left(M_{\varepsilon}\right)$ and by $\|\cdot\|_{\varepsilon}$ and $(\cdot, \cdot)_{\varepsilon}$ the norm and inner product in $L^{2}\left(M_{\varepsilon}\right)$. In certain obvious cases the index $\varepsilon$ in norms and inner products will be omitted.

By standard way (see, e.g. [9,15]) we can prove the following existence and uniqueness theorem.

Theorem 2.1. - Let $u_{0}^{\varepsilon}$ and $h^{\varepsilon}$ belong to $L^{2}\left(M_{\varepsilon}\right)$. Then for any interval $[0, T]$ problem (1)-(3) has a unique solution $u^{\varepsilon}(t)=u^{\varepsilon}(x, t)$ such that

$$
\begin{gather*}
u^{\varepsilon}(t) \in C\left(0 . T ; L^{2}\left(M_{\varepsilon}\right)\right) \cap L^{2}\left(0, T ; H^{1}\left(M_{\varepsilon}\right)\right)  \tag{2.5}\\
\left\|u^{\varepsilon}(t)\right\|_{\varepsilon}^{2} \int_{0}^{t}\left(\left\|\nabla_{\varepsilon} u^{\varepsilon}\right\|_{\varepsilon}^{2}+\eta\left\|u^{\varepsilon}\right\|_{\varepsilon}^{2}\right) d \tau \leqslant\left\|u_{0}^{\varepsilon}\right\|_{\varepsilon}^{2}+C_{1}\left(1+\left\|h^{\varepsilon}\right\|_{\varepsilon}^{2}\right) \tag{2.6}
\end{gather*}
$$

and

$$
\begin{equation*}
\left\|u^{\varepsilon}(t)\right\|_{\varepsilon}^{2} \leqslant\left\|u_{0}^{\varepsilon}\right\|_{\varepsilon}^{2} e^{-\eta t}+C_{2}\left(1+\left\|h^{\varepsilon}\right\|_{\varepsilon}^{2}\right)\left(1-e^{-\eta t}\right), \tag{2.7}
\end{equation*}
$$

where $C_{1}$ and $C_{2}$ are independent of $\varepsilon$. The solution $u^{\varepsilon}(t)$ has the following properties:
i) if $u_{0}^{\varepsilon} \in H^{1}\left(M_{\varepsilon}\right)$, then

$$
u^{\varepsilon}(t) \in C\left(0, T ; H^{1}\left(M_{\varepsilon}\right)\right) \cap L^{2}\left(0, T ; H^{2}\left(M_{\varepsilon}\right)\right)
$$

and

$$
\frac{\partial u^{\varepsilon}}{\partial t} \in L^{2}\left(0, T ; L^{2}\left(M_{\varepsilon}\right)\right) ;
$$

ii) if

$$
u_{0}^{\varepsilon} \in\left\{v \in H^{2}\left(M_{\varepsilon}\right): \frac{\partial v}{\partial n}=0 \text { on } \partial M_{\varepsilon}\right\} \equiv H_{N}^{2}\left(M_{\varepsilon}\right)
$$

then

$$
u^{\varepsilon}(t) \in C\left(0, T ; H_{N}^{2}\left(M_{\varepsilon}\right)\right)
$$

and

$$
\frac{\partial u^{\varepsilon}}{\partial t} \in C\left(0, T ; L^{2}\left(M_{\varepsilon}\right)\right) \cap L^{2}\left(0, T ; H^{1}\left(M_{\varepsilon}\right)\right) .
$$

To obtain additional estimates for the solutions $u^{\varepsilon}(t)$ we introduce on $H^{1}\left(M_{\varepsilon}\right)$ the Lyapunov function

$$
\begin{equation*}
V_{\varepsilon}(u)=\frac{1}{2}\left\|\nabla_{\varepsilon} u\right\|_{\varepsilon}^{2}+\int_{M_{\varepsilon}} \mathscr{F}(u(x)) d x-\left(h^{\varepsilon}, u\right)_{\varepsilon} \tag{2.8}
\end{equation*}
$$

It is clear that $V_{\varepsilon}$ is continuous on $H^{1}\left(M_{\varepsilon}\right)$ and there exist positive constants $\alpha_{j}$ and $\beta_{j}$ independent of $\varepsilon$ such that

$$
\begin{equation*}
\alpha_{1}\|u\|_{1, \varepsilon}^{2}-\beta_{1} \leqslant V_{\varepsilon}(u) \leqslant \alpha_{2}\|u\|_{1, \varepsilon}^{2}+\beta_{2} \tag{2.9}
\end{equation*}
$$

Here we assume that $\left\|h^{\varepsilon}\right\|_{\varepsilon} \leqslant C$ for all $0<\varepsilon \leqslant \varepsilon_{0}$.
One can easily prove (see, e.g. [1,7,17]) that the solution $u^{\varepsilon}(t)$ of problem (1)-(3) with $u_{0}^{\varepsilon} \in H^{1}\left(M_{\varepsilon}\right)$ satisfies

$$
\begin{equation*}
V_{\varepsilon}\left(u^{\varepsilon}(t)\right)+\int_{0}^{t}\left\|\partial_{t} u^{\varepsilon}(\tau)\right\|_{\varepsilon}^{2} d \tau=V_{\varepsilon}\left(u_{0}^{\varepsilon}\right) \tag{2.10}
\end{equation*}
$$

Lemma 2.1. - Let $u_{0}^{\varepsilon} \in H_{N}^{2}\left(M_{\varepsilon}\right)$. Then

$$
\begin{equation*}
\left\|\frac{\partial u^{\varepsilon}}{\partial t}\right\|_{\varepsilon}^{2}+2 \int_{0}^{t}\left\|\nabla_{\varepsilon} \frac{\partial u^{\varepsilon}}{\partial t}\right\|_{\varepsilon}^{2} d \tau \leqslant C_{1}+C_{2} V\left(u_{0}^{\varepsilon}\right)+\left\|u_{1}^{\varepsilon}\right\|_{\varepsilon}^{2} \tag{2.11}
\end{equation*}
$$

where $u_{1}^{\varepsilon}=\Delta_{\varepsilon} u_{0}^{\varepsilon}-f\left(u_{0}^{\varepsilon}\right)+h^{\varepsilon}$ and $C_{1,2}$ are independent of $\varepsilon$.
Proof. - Theorem 2.1 implies that $w^{\varepsilon}(t)=\partial u^{\varepsilon} / \partial t$ is a solution of the following problem:

$$
\begin{equation*}
\frac{\partial \omega^{\varepsilon}}{\partial t}-\Delta_{\varepsilon} w^{\varepsilon}+f^{\prime}\left(u^{\varepsilon}(t)\right) w^{\varepsilon}=0, \quad \frac{\partial \omega^{\varepsilon}}{\partial n}=0 \quad \text { on } \quad \partial M_{\tilde{\varepsilon}}, w^{\varepsilon}(x, 0)=u_{1}^{\varepsilon}(x) \tag{2.12}
\end{equation*}
$$

Since $\left|f^{\prime}(u)\right| \leqslant C$ it is clear that

$$
\begin{equation*}
\frac{1}{2} \frac{d}{d t}\left\|w^{\varepsilon}(t)\right\|_{\varepsilon}^{2}+\left\|\nabla_{\varepsilon} w^{\varepsilon}(t)\right\|_{\varepsilon}^{2} \leqslant C\left\|w^{\varepsilon}(t)\right\|_{\varepsilon}^{2} \tag{2.13}
\end{equation*}
$$

Therefore (2.11) follows from (2.10) and (2.13).
Remark 2.1. - From (2.10) and (2.13) it also follows that

$$
\begin{equation*}
t\left\|\frac{\partial u^{\varepsilon}}{\partial t}\right\|_{\varepsilon}^{2} \leqslant e^{C_{1} t}\left\{V\left(u_{0}^{\varepsilon}\right)+C_{2}\right\} . \tag{2.14}
\end{equation*}
$$

Therefore using (1), (2.7), (2.9), (2.10) we have

$$
\begin{equation*}
t\left\|\Delta_{\varepsilon} u^{\varepsilon}(t)\right\|_{\varepsilon}^{2} \leqslant C_{1} e^{C t}\left(1+\left\|u^{\varepsilon}\right\|_{1, \varepsilon}^{2}\right) \tag{2.15}
\end{equation*}
$$

if we assume that $\left\|h_{\varepsilon}\right\|_{\varepsilon} \leqslant C$ for all $0<\varepsilon \leqslant \varepsilon_{0}$.
Theorem 2.1 makes it possible to define an evolution operator $S_{t}^{\epsilon}$ on the space $H^{1}\left(M_{\varepsilon}\right)$ by the formula $S_{t}^{\varepsilon} u_{0}^{\varepsilon}=u^{\varepsilon}(t)$, where $u^{\varepsilon}(t)$ is the solution of the problem (1)(3). It is not difficult to show that $S_{i}^{\varepsilon}$ is a $C^{1}$-smooth nonlinear semigroup in the space $H^{1}\left(M_{\varepsilon}\right)$ and to prove (see, e.g. [1,17]) the following

Theorem 2.2. - The dynamical system $\left(S_{t}^{\varepsilon}, H^{1}\left(M_{\varepsilon}\right)\right)$ for every $\varepsilon>0$ has compact global attractor, i.e. there is a compact set $\mathfrak{Q}_{\varepsilon}$ in $H^{1}\left(M_{\varepsilon}\right)$ such that $S_{t}^{\ell} \mathfrak{O}_{\varepsilon}=\mathfrak{Q}_{\varepsilon}$ for $t \geqslant 0$ and

$$
\lim _{t a+\infty} \sup \left\{\operatorname{dist}_{H^{1}\left(M_{\varepsilon}\right)}\left(S_{t}^{\varepsilon} v, \mathfrak{a}_{\varepsilon}\right): v \in B\right\}=0
$$

for any bounded set $B$ in $H^{1}\left(M_{\varepsilon}\right)$. This attractor $\mathfrak{a}_{\varepsilon}$ has finite Hausdorff dimension.

Remark 2.2. - Using (2.7), (2.11), (2.15) and the formula

$$
u^{\varepsilon}(t)=e^{-L_{\varepsilon, \gamma} t} u_{0}^{\varepsilon}+\int_{0}^{t} e^{-L_{\varepsilon, y}(t-\tau)}\left(\gamma u^{\varepsilon}(\tau)-f\left(u^{\varepsilon}(\tau)\right)+h^{\varepsilon}\right) d \tau
$$

where $L_{\varepsilon, \gamma}=-\Delta_{\varepsilon}+\gamma$ with the Neumann boundary condition on $\partial M_{\varepsilon}, \gamma>0$, it is easy to show that for any trajectory $u^{\varepsilon}(t)$ lying in the attractor $\mathfrak{G}_{\varepsilon}$ we have the estimates:

$$
\begin{equation*}
\left\|\frac{\partial u^{\varepsilon}}{\partial t}(t)\right\|_{\varepsilon}^{2}+\left\|\Delta_{\varepsilon} u^{\varepsilon}(t)\right\|_{\varepsilon}^{2}+C\left\|\nabla_{\varepsilon} u^{\varepsilon}(t)\right\|_{\varepsilon}^{2}+\left\|u^{\varepsilon}(t)\right\|_{\varepsilon}^{2}<C_{1} \tag{2.16}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{-\infty}^{\infty}\left(\left\|\frac{\partial u^{\varepsilon}}{\partial t}(t)\right\|_{\varepsilon}^{2}+\left\|\nabla_{\varepsilon} \frac{\partial u^{\varepsilon}}{\partial t}(t)\right\|_{\varepsilon}^{2}\right) d t \leqslant C_{2}, \tag{2.17}
\end{equation*}
$$

where $C_{1}$ and $C_{2}$ are independent of $\varepsilon, 0<\varepsilon \leqslant \varepsilon_{0}$.

## 3. - Formulation of main results.

We introduce a parameter to describe the asymptotic behaviour of manifolds. For simplicity we will suppose $0 \in \Omega$, and denote

$$
G_{\varepsilon}=\left\{(x ; 0) \in \boldsymbol{R}^{n+1}: a_{\varepsilon} \leqslant|x|<\frac{\varepsilon}{2}\right\}, \quad D_{\varepsilon}=B_{\varepsilon} \cup G_{\varepsilon},
$$

We set

$$
\begin{equation*}
\lambda_{\varepsilon}=\inf \left\{\frac{\left\|\nabla_{\varepsilon} v\right\|_{L^{2}\left(D_{\varepsilon}\right)}^{2}}{\|v\|_{L^{2}\left(D_{\varepsilon}\right)}^{2}}: v \in H_{0}^{1}\left(D_{\varepsilon}\right)\right\} . \tag{3.1}
\end{equation*}
$$

$\lambda_{\varepsilon}$ is the first eigenvalue of the Dirichlet problem

$$
\begin{equation*}
\Delta_{\varepsilon}+\lambda_{\varepsilon} v=0, \quad x \in D_{\varepsilon} ; \quad v=0, \quad x \in \partial D_{\varepsilon} \tag{3.2}
\end{equation*}
$$

Our main assumption concerning to behaviour of the bubbles $B_{\varepsilon}^{j}$ (and manifold $M_{\varepsilon}$ ) is the existence of the limits

$$
\begin{equation*}
\lambda=\lim _{\varepsilon \rightarrow 0} \lambda_{\varepsilon} \quad \text { and } \quad \mu=\lim _{\varepsilon \rightarrow 0} \varepsilon^{-n} m_{\varepsilon}>0, \tag{3.3}
\end{equation*}
$$

where

$$
\begin{equation*}
m_{\varepsilon} \doteq \operatorname{Vol}\left(B_{\varepsilon}\right)=\int_{B_{\varepsilon}} \sqrt{\left|g^{\varepsilon}\right|} d x_{1} \ldots d x_{n} \tag{3.3}
\end{equation*}
$$

Remark 3.1. - It is easy to see that

$$
0<\lambda_{\varepsilon} \leqslant C \begin{cases}a_{\varepsilon}^{n-2} \varepsilon^{-n}, & n>2, \\ \left|\operatorname{In} \mathbf{a}_{\varepsilon}\right|^{-1} \varepsilon^{-2}, & n=2 .\end{cases}
$$

Moreover, if the metric on $M_{\varepsilon}$ coincides with the metric induced from $\boldsymbol{R}^{n+1}$ outside of small neighbourhoods of the boundaries $\partial B_{\varepsilon}^{j}$, one can prove that the condition

$$
a_{\varepsilon}= \begin{cases}a \varepsilon^{n /(n-2)}, & n>2, \\ \exp \left(-1 / \varepsilon^{2}\right), & n=2,\end{cases}
$$

implies that limits (3.3) exist and $\lambda=(1 / 2) a^{n-2} b^{-n}$ and $\mu=\omega_{n}$, where $\omega_{n}$ is the volume of the unit sphere in $\boldsymbol{R}^{n+1}$ (see [2] for a closely related assertion). From this observation and (2.1) it also follows that for existence of limits (3.3) it is necessary that

$$
C_{1} \varepsilon^{n /(n-2)} \leqslant a_{\varepsilon} \leqslant C_{2} \varepsilon^{n /(n-2)} \quad \text { for } n \geqslant 3
$$

and

$$
C_{1} \exp \left(-1 / \varepsilon^{2}\right) \leqslant a_{\varepsilon} \leqslant C_{2} \exp \left(-1 / \varepsilon^{2}\right) \text { for } n=2 .
$$

Let $P_{\varepsilon}$ be a bounded operator from $L^{2}\left(M_{\varepsilon}\right)$ into $L^{2}(\Omega)$ defined by the formula

$$
\left(P_{\varepsilon} u\right)(x)= \begin{cases}u(x), & x \in \Omega_{\varepsilon}, \\ 0 & x \in \Omega \backslash \Omega_{\varepsilon},\end{cases}
$$

and let $Q_{\varepsilon}$ be the operator which maps a function $u \in L^{2}\left(M_{\varepsilon}\right)$ into poly-linear spline
$Q_{\varepsilon} u$ associated with a net $\left\{x^{j}=j \varepsilon, j \in N_{\varepsilon}\right\}$ such that

$$
\left(Q_{\varepsilon} u\right)\left(x^{j}\right)=\frac{1}{m_{\varepsilon}} \int_{B_{\varepsilon}^{i}} u(x) d x, \quad j \in N_{\varepsilon} .
$$

It is clear that $Q_{\varepsilon}$ is a linear bounded operator from $L^{2}\left(M_{\varepsilon}\right)$ into $H^{1}\left(\Omega_{N_{\varepsilon}}\right)$, where $\Omega_{N_{\varepsilon}}$ is the union of elementary cubes corresponding to the net $\left\{j \varepsilon: j \in N_{\varepsilon}\right\}$. If we set $Q_{\varepsilon} u(x)=0$ for $x \in \Omega \backslash \Omega_{N_{\varepsilon}}$, we can also consider $Q_{\varepsilon}$ as a bounded operator from $L^{2}\left(M_{\varepsilon}\right)$ into $L^{2}(\Omega)$.

The first main result of the paper is the following
Theorem 3.1. - Let $u^{\varepsilon}(t)$ be the solution of the problem (1)-(3). Assume that i) for any $\varepsilon \in\left(0, \varepsilon_{0}\right)$, we have

$$
\left\|u_{0}^{\varepsilon}\right\|_{1, \varepsilon}+\left\|\nabla Q_{\varepsilon} u_{0}^{\varepsilon}\right\|_{L^{2}\left(\Omega_{\left.N_{\varepsilon}\right)}\right)} \leqslant C
$$

and

$$
\left\|h^{\varepsilon}\right\|_{1, \varepsilon}+\left\|\nabla Q_{\varepsilon} u^{\varepsilon}\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)} \leqslant C,
$$

where the constant $C$ is independent of $\varepsilon$;
ii) there exist functions $u_{0}, v_{0}, h_{1}, h_{2}$ from $L^{2}(\Omega)$ such that $P_{\varepsilon} u_{0}^{\varepsilon} \rightarrow u_{0}, Q_{\varepsilon} u_{0}^{\varepsilon} \rightarrow$ $\rightarrow v_{0}, P_{\varepsilon} h^{\varepsilon} \rightarrow h_{1}, Q_{\varepsilon} h^{\varepsilon} \rightarrow h_{2}$ strongly in $L^{2}(\Omega) ;$
iii) there exist limits (3.3).

Then for any interval $[0, T]$ we have that

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}\left\{\max _{[0, T]}\left\|P_{\varepsilon} u^{\varepsilon}(t)-u(t)\right\|_{L^{2}(\Omega)}^{2}+\max _{[0, T]}\left\|Q_{\varepsilon} u^{\varepsilon}(t)-v(t)\right\|_{L^{2}(\Omega)}^{2}\right\}=0, \tag{3.4}
\end{equation*}
$$

where the pair of functions $u(t)=u(x, t)$ and $v(t)=v(x, t)$ is the solution of the problem:

$$
\begin{equation*}
\left.\frac{\partial u}{\partial n}\right|_{\partial \Omega}=0,\left.\quad u\right|_{t=0}=u_{0}(x) \tag{3.6}
\end{equation*}
$$

$$
\begin{equation*}
\frac{\partial u}{\partial t}+\lambda(v-u)+f(v)=h_{2}(x), \quad x \in \Omega, \quad t>0 \tag{3.7}
\end{equation*}
$$

$$
\begin{equation*}
\frac{\partial u}{\partial t}-\Delta u+\lambda \mu(u-v)+f(u)=h_{1}(x), \quad x \in \Omega, t>0, \tag{3.5}
\end{equation*}
$$

$$
\begin{equation*}
\left.v\right|_{t=0}=v_{0}(x), \tag{3.8}
\end{equation*}
$$

The proof of this theorem consists of two parts. The main point of the first one is to obtain a uniform estimate

$$
\begin{equation*}
\int_{0}^{T}\left\|\nabla Q_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2} d t<C \tag{3.9}
\end{equation*}
$$

In the second part we make a limit transition in the equation (1) on testing functions of special structure. In order to prove the uniqueness of limits we also use the following

TheOrem 3.2. - Assume that (2.2)-(2.4) are satisfied and $U_{0}=\left(u_{0}, v_{0}\right) \in \mathscr{F}_{0}=$ $=L^{2}(\Omega) \times L^{2}(\Omega)$. Then the problem (3.5)-(3.8) has a unique generalized solution $U(t)=(u(t), v(t))$ belonging to the space $C\left(\boldsymbol{R}_{+}, \mathscr{F}_{0}\right)$. Moreover, if $U_{0} \in \mathscr{F}_{1}=H^{1}(\Omega) \times$ $\times L^{2}(\Omega)$ then

$$
\begin{equation*}
U(t) \in C\left(\boldsymbol{R}_{+}, \mathscr{F}_{1}\right) \quad \text { and } \quad \frac{d}{d t} U(t) \in L^{2}\left(\boldsymbol{R}_{+}, \mathscr{F}_{0}\right) \tag{3.10}
\end{equation*}
$$

if $U_{0} \in \mathscr{J}_{2}=H^{1}(\Omega) \times L^{2}(\Omega)$ and $h_{2} \in H^{1}(\Omega)$ then

$$
\begin{equation*}
U(t) \in C\left(\boldsymbol{R}_{+}, \mathscr{F}_{2}\right) \quad \text { and } \quad \frac{d}{d t} U(T) \in L^{2}\left(\boldsymbol{R}_{+}, L^{2}(\Omega) \times H^{1}(\Omega)\right) . \tag{3.11}
\end{equation*}
$$

The proof of this theorem is of standard character and relies on the methods presented in [9].

Theorem 3.2 allows us to define the evolutionary semigroup $S_{t}$ in each of the spaces $\mathscr{F}_{i}$ by the formula $S_{t} U_{0}=U(t)$, where $U(t)$ is the solution of the problem (3.5)(3.8). If we consider this semigroup in $\mathscr{F}_{2}$, then we can prove the following assertion on the existence of a global attractor.

Theorem 3.3. - Assume that (2.2)-(2.4) are satisfied and

$$
\begin{equation*}
\lambda+\inf \left\{f^{\prime}(u): u \in \boldsymbol{R}^{1}\right\}>0, \quad h_{2}(x) \in H^{1}(\Omega) \tag{3.12}
\end{equation*}
$$

Then the dynamical system $\left(S_{t}, \mathfrak{F}_{2}\right)$ has a weak global attractor $\mathfrak{F}$. This attractor has finite Hausdorff dimension as a compact set in $\mathscr{F}_{0}$.

In order to prove this theorem we rely on certain results from [6,17]. Recall (see $[1,4,17]$ ) that a weak global attractor $\mathfrak{G}$ is a bounded weakly closed set in $\mathscr{F}_{2}$ such that (i) $S_{t} \mathfrak{a}=\mathfrak{A}$ for any $t>0$ and (ii) for any weak neighbourhood $\mathcal{O}$ of $A$ and for any bounded set $B \subset \mathscr{F}_{2}$ we have $S_{t} B \subset \mathcal{O}$, when $t \geqslant t_{0}(B, \mathcal{O})$.

At last using Theorem 3.1 and estimates (2.16) and (2.17) we prove the second main result of the paper.

Theorem 3.4. - Assume that (2.2)-(2.4), (3.12) and the assumptions of Theorem 3.1 are satisfied. Then we have

$$
\left.\lim _{\varepsilon \rightarrow 0} \sup _{u_{\varepsilon} \in \mathbb{Q}_{\varepsilon}}\left\{\inf _{(u, v) \in \mathbb{Q}}\left\|P_{\varepsilon} u^{\varepsilon}-u\right\|_{L^{2}(\Omega)}^{2}+\left\|Q_{\varepsilon} u^{\varepsilon}-v\right\|_{L^{2}(\Omega)}^{2}\right)\right\}=0 .
$$

This theorem means that the global attractor $\mathfrak{G}_{\varepsilon}$ of problem (1)-(3) tends to a weak global attractor $\mathfrak{C}$ of the homogenized system (3.5)-(3.8).

## 4. - Uniform estimates.

Now we begin the proof of Theorem 3.1. In this section we establish our main Lemma 4.1 on uniform boundness of the norms $\left\|Q_{\varepsilon} u^{\varepsilon}\right\|_{H^{1}\left(\Omega_{N_{\varepsilon}} \times(0, T)\right)}$. This lemma and estimates for $P_{\varepsilon} u^{\varepsilon}$ which directly follow from (2.6) and (2.10) make it possible to extract from $\left\{P_{\varepsilon} u^{\varepsilon}\right\}$ and $\left\{Q_{\varepsilon} u^{\varepsilon}\right\}$ subsequences strongly convergent in $L^{2}(\Omega \times(0, T))$. Below we consider the case $n \geqslant 3$ only. For the case $n=2$ the consideration should be repeated word by word with slight modifications in the estimates. We assume that the conditions (i)-(iii) of Theorem 3.1 are satisfied.

At first we note that (2.7) and (2.10) imply that the solution $u^{\varepsilon}(x, t)$ satisfies the estimate

$$
\begin{equation*}
\left\|u^{\varepsilon}(t)\right\|_{\varepsilon}^{2}+\left\|\nabla_{\varepsilon} u^{\varepsilon}(t)\right\|_{\varepsilon}^{2}+\int_{0}^{t}\left\|\frac{\partial u^{\varepsilon}}{\partial t}\right\|_{\varepsilon}^{2} d \tau \leqslant C_{T} \tag{4.0}
\end{equation*}
$$

for any $t \in[0, T]$. Since the metric $g^{\varepsilon}$ coincides with the euclidean one on $\Omega_{\varepsilon}$, we have

$$
\begin{equation*}
\left\|P_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}(\Omega)}^{2}+\left\|\nabla_{\varepsilon} P_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}(\Omega)}^{2}+\int_{0}^{t}\left\|P_{\varepsilon} \frac{\partial u^{\varepsilon}}{\partial t}\right\|_{L^{2}(\Omega)}^{2} d \tau \leqslant C_{T} . \tag{4.1}
\end{equation*}
$$

The remaining part of this section is devoted to the proof of a similar estimate for $Q_{\varepsilon} u^{\varepsilon}(t)$.

Let us introduce the following notation:

$$
\begin{gathered}
u_{k}^{\varepsilon}(x, t)=u^{\varepsilon}\left(x^{k}+x, t\right), \quad x^{k}=l \varepsilon, \quad k \in N_{\varepsilon}, \quad x \in D_{\varepsilon} ; \\
u_{k}^{\varepsilon, i n}(t)=\frac{1}{m_{\varepsilon}} \int_{B_{\varepsilon}} u_{k}^{\varepsilon}(x, t) d x ; \\
u_{k}^{\varepsilon, e x}(t)=\frac{1}{m_{\varepsilon}^{\prime}} \int_{G_{\varepsilon}} u_{k}^{\varepsilon}(x, t) d x ;
\end{gathered}
$$

where $u^{\varepsilon}(x, t)$ is the solution of problem (1)-(3), the sets $B_{\varepsilon}, G_{\varepsilon}$ and $D_{\varepsilon}$ are defined in Sections 2 and $3, m_{\varepsilon}=\operatorname{Vol}\left(B_{\varepsilon}\right)$ and $m_{\varepsilon}^{\prime}=\operatorname{Vol}\left(G_{\varepsilon}\right)$. We also use the notation

$$
w^{\varepsilon} \equiv w_{k l}^{\varepsilon}(x, t)=u_{k}^{\varepsilon}(x, t)-u_{l}^{\varepsilon}(x, t), \quad x \in D_{\varepsilon}, \quad k, l \in N_{\varepsilon}
$$

and

$$
w^{\#}=w_{k l}^{\# \#}(t)=u_{k}^{\varepsilon_{k}^{,}, \#}(t)-u_{l}^{\varepsilon_{1}, \#}(t), \quad k, l \in N_{\varepsilon}
$$

where \# is either «in» or «ex».

It is clear from (2.7) and (2.10) that for any $t \geqslant 0$

$$
\begin{equation*}
\left\|Q_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}(\Omega)}^{2}+\int_{0}^{t}\left\|Q_{\varepsilon} \frac{\partial u^{\varepsilon}}{\partial t}\right\|_{L^{2}(\Omega)}^{2} d \tau \leqslant C_{T} \tag{4.2}
\end{equation*}
$$

The main result of this section is
Lemma 4.1. - For any $T>0$ we have

$$
\left\|\nabla Q_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2} \leqslant C_{T}, \quad t \in[0, T]
$$

where $C_{T}$ is a constant independent of $\varepsilon$.
In order to prove this Lemma it is sufficient to obtain appropriate estimates for $w_{k l}^{i n}(t)$. We will use the following preliminary assertions.

Lemma 4.2. - Let $a_{\varepsilon} \leqslant a \varepsilon^{n /(n-2)}(n>2)$ and let $v^{\varepsilon}(x) \in H_{0}^{1}\left(D_{\varepsilon}\right)$ be the solution of the problem (3.2) such that

$$
\begin{equation*}
\int_{B_{\varepsilon}} v^{\varepsilon}(x) d x=m_{\varepsilon} . \tag{4.3}
\end{equation*}
$$

Then we have the following estimates:
(i) $\left|D^{\alpha} v^{\varepsilon}(x)\right| \leqslant \frac{C \varepsilon^{n}}{|x|^{n-2+|\alpha|}}$ for $x \in \widetilde{G}_{\varepsilon}$ and $|x| \geqslant \varepsilon / 4$;
(ii) $\int_{G_{\varepsilon}}\left|v^{\varepsilon}(x)\right|^{2} d x \leqslant C \varepsilon^{n+2}$;
(iii) $\int_{D_{\varepsilon}}\left|v^{\varepsilon}(x)\right|^{2} d x=\int_{B_{\varepsilon}}\left|v^{\varepsilon}(x)\right|^{2} d x+O\left(\varepsilon^{n+2}\right)=m_{\varepsilon}+O\left(\varepsilon^{n+2}\right)$;
(iv) $\int_{D_{\varepsilon}}\left|\nabla_{\varepsilon} v^{\varepsilon}(x)\right|^{2} d x=\lambda_{\varepsilon} m_{\varepsilon}+O\left(\varepsilon^{n+2}\right)$;
(v) $\int_{\Gamma\left(a_{\varepsilon}\right)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma=\lambda_{\varepsilon} m_{\varepsilon}$ and $\int_{\Gamma(\varepsilon / 2)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma=\lambda_{\varepsilon} m_{\varepsilon}+O\left(\varepsilon^{n+1}\right) ;$
where $\Gamma\left(a_{\varepsilon}\right)$ and $\Gamma(\varepsilon / 2)$ are the inner and outer boundaries of the ring $G_{\varepsilon}$, and the normal vector $n$ is directed towards the center of the ring $G_{\varepsilon}$.

Proof. - It is easy to see that for $v^{\varepsilon}(x)$ we have the following inequalities of

Poincaré and Friedrichs type:

$$
\begin{equation*}
\int_{B_{\varepsilon}}\left(v^{\varepsilon}-1\right)^{2} d x \leqslant C \varepsilon^{2} \int_{B_{\varepsilon}}\left|\nabla_{\varepsilon} v^{\varepsilon}\right|^{2} d x ; \tag{4.4}
\end{equation*}
$$

$$
\begin{equation*}
\int_{G_{\varepsilon}}\left|v^{\varepsilon}\right|^{2} d x \leqslant C \varepsilon^{2} \int_{G_{\varepsilon}}\left|\nabla v_{\varepsilon} v^{\varepsilon}\right|^{2} d x \tag{4.5}
\end{equation*}
$$

Since

$$
\int_{D_{\varepsilon}}\left|\nabla_{\varepsilon} v^{\varepsilon}\right|^{2} d x=\lambda_{\varepsilon}\left\{m_{\varepsilon}+; \int_{B_{\varepsilon}}\left(v^{\varepsilon}-1\right)^{2} d x+\int_{G_{\varepsilon}}\left|v^{\varepsilon}\right|^{2} d x\right\}
$$

from (3.3), (4.4) and (4.5) we have

$$
\begin{equation*}
\int_{D_{\varepsilon}}\left|\nabla_{\varepsilon} v^{\varepsilon}\right|^{2} d x \leqslant C \varepsilon^{n} \tag{4.6}
\end{equation*}
$$

and the property (iv) follows. Now (ii) and (iii) follow from (4.5) and (4.6). Using . Green's formula we get

$$
\int_{\Gamma \varepsilon / 2)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma=\int_{D_{\varepsilon}}\left|\nabla_{\varepsilon} v^{\varepsilon}\right|^{2} d x-\int_{D_{\varepsilon}} \Delta v^{\varepsilon}\left(1-v^{\varepsilon}\right) d x
$$

Therefore using (3.2)and (4.3) we obtain

$$
\int_{\Gamma(\varepsilon / 2)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma=\lambda_{\varepsilon} m_{\varepsilon}+\lambda_{\varepsilon} \int_{G_{\varepsilon}} v^{\varepsilon} d x .
$$

Hence (v) follows from (iii) and from the obvious formula:

$$
\int_{\Gamma\left(a_{\varepsilon}\right)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma=\int_{\Gamma(\varepsilon / 2)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma-\lambda_{\varepsilon} \int_{G_{\varepsilon}} v^{\varepsilon} d x .
$$

We now prove (i). Let $\Gamma(x, y)$ be the generalized solution of the problem:

$$
\Delta \Gamma(x, y)+\lambda_{\varepsilon} \Gamma(x, y)=-\delta(x-y) \text { for } x, y \in K_{\varepsilon},\left.\quad \Gamma_{\varepsilon}(x, y)\right|_{x \in \partial K_{\varepsilon}}=0,
$$

where $K_{\varepsilon}=\left\{x \in \boldsymbol{R}^{n}:|x|<\varepsilon / 2\right\}$. It is well known that

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{y}^{\beta} \Gamma(x, y)\right| \leqslant C|x-y|^{-n+2-|\alpha|-|\beta|} . \tag{4.7}
\end{equation*}
$$

For $y \in G_{\varepsilon},|y| \geqslant 2 a_{\varepsilon}$, and $x \in D_{\varepsilon}$ we define the function

$$
R_{\varepsilon}(x, y)= \begin{cases}\Gamma(0, y), & x \in B_{\varepsilon} \\ \Gamma(x, y)+(\Gamma(0, y)-\Gamma(x, y)) \varphi\left(\frac{x}{a_{\varepsilon}}\right), & x \in G_{\varepsilon}\end{cases}
$$

where $\varphi(x) \in C_{0}^{\infty}\left(\boldsymbol{R}^{n}\right)$ satisfies:

$$
\varphi(x)=1 \quad \text { for }|x| \leqslant 1 \quad \varphi(x)=0 \quad \text { for }|x| \geqslant \frac{3}{2} .
$$

It is clear that $R_{\varepsilon}(x, y)=0$, when $x \in \partial D_{\varepsilon}$ and

$$
\left(\Delta_{\varepsilon, x}+\lambda_{\varepsilon}\right) R_{\varepsilon}(x, y)= \begin{cases}\lambda_{\varepsilon} \Gamma(0, y), & x \in B_{\varepsilon} \\ -\delta(x-y)+\theta_{\varepsilon}(x, y), & x \in G_{\varepsilon}\end{cases}
$$

where

$$
\theta_{\varepsilon}(x, y)=a_{\varepsilon}^{-2} \Delta \varphi(\Gamma(0, y)-\Gamma(x, y))-2 a_{\varepsilon}^{-1} \sum_{i=1}^{n} \frac{\partial \varphi}{\partial x_{i}} \frac{\partial}{\partial x_{i}} \Gamma(x, y)+\lambda_{\varepsilon} \Gamma(0, y) \varphi .
$$

Using the Green formula we have

$$
\int_{D_{\varepsilon}} v^{\varepsilon}(x)\left(\Delta_{\varepsilon, x}+\lambda_{\varepsilon}\right) R_{\varepsilon}(x, y) d x=0
$$

Consequently,

$$
\begin{equation*}
v^{\varepsilon}(y)=\lambda_{\varepsilon} m_{\varepsilon} \Gamma(0, y)+\int_{G+_{\varepsilon}} \theta_{\varepsilon}(x, y) v^{\varepsilon}(x) d x \tag{4.8}
\end{equation*}
$$

for $y \in G_{\varepsilon}$ and $|y| \geqslant 2 a_{\varepsilon}$, where $G_{\varepsilon}^{*}=\left\{x \in G_{\varepsilon} ;|x| \leqslant(3 / 2) a_{\varepsilon}\right\}$. The property (i) then follows from (ii), (4.7) and (4.8).

Lemma 4.3. - Let $v^{\varepsilon}(x)$ be as in Lemma 4.2. Then we have

$$
\begin{equation*}
\frac{d}{d t} \int_{D_{\varepsilon}} w^{\varepsilon}(t) v^{\varepsilon} d x+\left(\lambda_{\varepsilon}+I_{k l}^{\varepsilon}(t)\right) \int_{D_{\varepsilon}} w^{\varepsilon}(t) v^{\varepsilon} d x=R_{k l}^{\varepsilon}(t) \tag{4.9}
\end{equation*}
$$

where

$$
I_{k l}^{\varepsilon}(t)=\frac{1}{m_{\varepsilon}} \int_{0}^{1} d \tau \int_{B_{\varepsilon}} f^{\prime}\left(u_{k}^{\varepsilon}(t)+\tau\left(u_{l}^{\varepsilon}(t)-u_{k}^{\varepsilon}(t)\right)\right) v^{\varepsilon} d x
$$

and the quantity $R_{k l}^{\varepsilon}(t)$ admits the estimate

$$
\begin{aligned}
& \left|R_{k l}^{\varepsilon}(t)\right| \leqslant \\
& \quad C_{1} m_{\varepsilon}\left(\left|w^{e x}(t)\right|+\left|h^{i n}\right|\right)+ \\
& \quad+C_{2} \varepsilon^{n / 2+1}\left\{\left\|w^{\varepsilon}(t)\right\|_{L^{2}\left(G_{\varepsilon}\right)}+\left\|\nabla_{\varepsilon} w^{\varepsilon}(t)\right\|_{L^{2}\left(D_{\varepsilon}\right)}+\left\|h_{k}\right\|_{L^{2}\left(G_{\varepsilon}\right)}+\left\|\nabla_{\varepsilon} h_{k l}\right\|_{L^{2}\left(G_{\varepsilon}\right)}\right\} .
\end{aligned}
$$

Here $h_{k l}(x)=h\left(x_{k}+x\right)-h\left(x_{l}+x\right)$ for $x \in D_{\varepsilon}, k, l \in N_{\varepsilon}$ and $h^{i n}$ is defined in the same manner as $u^{\varepsilon, \text { in }}$.

Proof. - We use the equation

$$
\begin{equation*}
\frac{d}{d t} \int_{D_{\varepsilon}} w^{\varepsilon} v^{\varepsilon} d x-\int_{D_{\varepsilon}} \Delta w^{\varepsilon} v^{\varepsilon} d x+\int_{D_{\varepsilon}}\left(f\left(u_{k}^{\varepsilon}\right)-f\left(u_{l}^{\varepsilon}\right)\right) v^{\varepsilon} d x=\int_{D_{\varepsilon}} h_{k l} v^{\varepsilon} d x, \tag{4.10}
\end{equation*}
$$

which follows from (1) and we use the following Lemmas.
Lemma 4.4

$$
\begin{equation*}
\left|\int_{D_{\varepsilon}} \Delta w^{\varepsilon} v^{\varepsilon} d x+\lambda_{\varepsilon} \int_{D_{\varepsilon}} w^{\varepsilon} v^{\varepsilon} d x\right| \leqslant C\left(\varepsilon^{n}\left|w^{e x}\right|+\varepsilon^{n / 2+1}\left\|\nabla w^{\varepsilon}(t)\right\|_{L^{2}\left(G_{\varepsilon}\right)}\right) . \tag{4.11}
\end{equation*}
$$

Proof. - Using Green formula we get

$$
\int_{D_{\varepsilon}} \Delta w^{\varepsilon} v^{\varepsilon} d x+\lambda_{\varepsilon} \int_{D_{c}} w^{\varepsilon} v^{\varepsilon} d x=-w^{e x} \int_{\Gamma(\varepsilon / 2)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma+\int_{\Gamma(\varepsilon / 2)}\left(w^{e x}-w^{\varepsilon}\right) \frac{\partial v^{\varepsilon}}{\partial n} d \sigma
$$

Lemma 4.2 (i) implies that

$$
\left|\int_{\Gamma(\varepsilon / 2)}\left(w^{e x}-w^{\varepsilon}\right) \frac{\partial v^{\varepsilon}}{\partial n} d \sigma\right| \leqslant C \varepsilon \int_{\Gamma(\varepsilon / 2)}\left|w^{e x}-w^{\varepsilon}\right| d \sigma
$$

Therefore, using the trace theorem and the Poincaré inequality in $G_{\varepsilon}$ we obtain (4.11).

Lemma 4.5.

$$
\begin{align*}
\left|\int_{D_{\varepsilon}}\left(f\left(u_{k}^{\varepsilon}\right)-f\left(u_{l}^{\varepsilon}\right)\right) v^{\varepsilon} d x-I_{k l}^{\varepsilon}(t) \int_{D_{\varepsilon}} w^{\varepsilon} v^{\varepsilon} d x\right| & \leqslant  \tag{4.12}\\
& \leqslant C \varepsilon^{n / 2+1}\left\{\left\|\nabla_{\varepsilon} w^{\varepsilon}\right\|_{L^{2}\left(B_{\varepsilon}\right)}+\left\|w^{\varepsilon}\right\|_{L^{2}\left(G_{\varepsilon}\right)}\right\}
\end{align*}
$$

Proof. - It is clear from (2.2) and Lemma 4.2 (ii) that

$$
\left|\int_{D_{\varepsilon}}\left(f\left(u_{k}^{\varepsilon}\right)-f\left(u_{l}^{\varepsilon}\right)\right) v^{\varepsilon} d x\right| \leqslant C\|w\|_{L^{2}\left(G_{\varepsilon}\right)}\|v\|_{L^{2}\left(G_{\varepsilon}\right)} \leqslant C \varepsilon^{n / 2+1}\|w\|_{L^{2}\left(G_{\varepsilon}\right)} .
$$

Using Lemma 4.2 and the Hölder and Poincaré inequalities we also get

$$
\begin{align*}
\left|\int_{D_{\varepsilon}} w^{\varepsilon} v^{\varepsilon} d x-w^{i n} m_{\varepsilon}\right| \leqslant\left|\int_{D_{\varepsilon}} w^{\varepsilon} v^{\varepsilon} d x\right| & +\left|\int_{B_{\varepsilon}}\left(w^{\varepsilon}-w^{i n}\right) v^{\varepsilon} d x\right| \leqslant  \tag{4.13}\\
& \leqslant C \varepsilon^{n / 2+1}\left(\left\|\nabla_{\varepsilon} w^{\varepsilon}\right\|_{L^{\varepsilon}\left(B_{\varepsilon}\right)}+\left\|w^{\varepsilon}\right\|_{L^{2}\left(G_{\varepsilon}\right)}\right)
\end{align*}
$$

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Obviously

$$
\begin{aligned}
& \left|\int_{B_{\varepsilon}}\left(f\left(u_{k}^{\varepsilon}\right)-f\left(u_{l}^{\varepsilon}\right)\right) v^{\varepsilon} d x-w^{i n} m_{\varepsilon} I_{k l}^{\varepsilon}(t)\right| \leqslant C \int_{B_{\varepsilon}}\left|w^{\varepsilon}-w^{i n} \| v^{\varepsilon}\right| d x \leqslant \\
& \leqslant C \varepsilon^{n / 2+1}\left\|\nabla_{\varepsilon} w^{\varepsilon}\right\|_{L^{2}\left(B_{\epsilon}\right)} .
\end{aligned}
$$

Now in order to obtain (4.12) it is sufficient to note that

$$
\begin{equation*}
\inf _{u \in R} f^{\prime}(u) \leqslant I_{k l}^{e}(t) \leqslant \sup _{u \in R} f^{\prime}(u) . \tag{4.14}
\end{equation*}
$$

This follows from the fact that $v^{\varepsilon}$, first eigenfunction of the Dirichlet problem is positive in $D_{\varepsilon}$.

As in (4.13) we have

$$
\left|\int_{D_{\varepsilon}} h_{k l} v^{\varepsilon} d x\right| \leqslant\left|h^{i n}\right| m_{\varepsilon}+C \varepsilon^{n / 2+1}\left(\left\|\nabla_{\varepsilon} h_{k l}\right\|_{L^{2}\left(B_{\varepsilon}\right)}+\left\|h_{k l}\right\|_{L^{2}\left(G_{\varepsilon}\right)}\right) .
$$

Therefore Lemma 4.4 and 4.5 give equality (4.9), and this proves Lemma 4.3.
Using (4.9) and (4.14) we get

$$
\begin{equation*}
\frac{1}{2} \frac{d}{d t}\left(\int_{D_{\varepsilon}} w^{\varepsilon} v^{\varepsilon} d x\right)^{2}+\left(\lambda_{\varepsilon}+L-\delta\right)\left(\int_{D_{\varepsilon}} w^{\varepsilon} v^{\varepsilon} d x\right)^{2} \leqslant \frac{1}{4 \delta}\left[R_{k l}^{\varepsilon}\right]^{2} \tag{4.15}
\end{equation*}
$$

for any positive $\delta$, where $L=\inf _{u \in R} f^{\prime}(u)$. Then it follows from Gronwall's lemma that we have

$$
\begin{equation*}
\left(\int_{D_{\varepsilon}} w^{\varepsilon}(t) v^{\varepsilon} d x\right)^{2} \leqslant e^{-\alpha_{\varepsilon} t}\left(\int_{D_{\varepsilon}} w^{\varepsilon}(0) v^{\varepsilon} d x\right)^{2}+\frac{1}{2 \delta} \int_{0}^{t} e^{-\alpha_{\varepsilon}(t-\tau)}\left[R_{k l}^{\varepsilon}(\tau)\right]^{2} d \tau \tag{4.16}
\end{equation*}
$$

for any $\delta>0$, with $\alpha_{\varepsilon}=2\left(\lambda_{\varepsilon}+L-\delta\right)$.
It is clear that

$$
\left[R_{k l}^{\varepsilon}(t)\right]^{2} \leqslant C_{1} m_{\varepsilon}^{2}\left(\left|w^{e x}\right|^{2}+\left|h^{i n}\right|^{2}\right)+C_{2} \varepsilon^{n+2}\left(Y_{k}^{\varepsilon}(t)+Y_{l}^{\varepsilon}(t)\right)
$$

where

$$
Y_{l}^{\varepsilon}(t)=\left\|u^{\varepsilon}(t)\right\|_{H^{1}\left(D_{\varepsilon}^{i}\right)}^{2}+\|h\|_{H^{1}\left(D_{k}^{i}\right)}^{2}
$$

Here $D_{\varepsilon}^{j}=(j \varepsilon ; 0)+D_{\varepsilon}$ for $j \in N_{\varepsilon}$. Therefore

$$
\begin{aligned}
& \sum_{k \in N_{\varepsilon}} \sum_{l \in \sigma(k)} \frac{1}{m_{\varepsilon}}\left[R_{k l}^{\varepsilon}(t)\right]^{2} \leqslant C_{1} \varepsilon^{2}\left(\left\|u^{\varepsilon}(t)\right\|_{H^{1}\left(M_{\varepsilon}\right)}^{2}+\|h\|_{H^{1}\left(M_{\varepsilon}\right)}^{2}\right)+ \\
&+C_{2} \sum_{k \in N_{\varepsilon} l \in \sigma(k)} \sum_{l} \varepsilon^{n}\left(\left|u_{k}^{e x}-u_{l}^{e x}\right|^{2}+\left|h_{k}^{i n}-h_{l}^{i n}\right|^{2}\right)
\end{aligned}
$$

where $\sigma(k)=\bar{\sigma}(k) \cap N_{\varepsilon}$ and $\bar{\sigma}(k)$ is the set of the nearest neighbours of $k$ in $\boldsymbol{Z}^{n}$. Since any function $\varphi \in H^{1}\left(\Omega_{\varepsilon}\right)$ can be extended to $\widetilde{\varphi} \in H^{1}(\Omega)$ such that

$$
\|\widetilde{\varphi}\|_{H_{1}(\Omega)} \leqslant C\|\varphi\|_{H_{1}\left(\Omega_{\varepsilon}\right)}
$$

with constant $C$ independent of $\varepsilon$, one can easily verify that

$$
\sum_{k \in N_{\varepsilon}} \sum_{l \in \sigma(k)} \varepsilon^{n}\left|u_{k}^{e x}-u_{l}^{e x}\right|^{2} \leqslant C \varepsilon^{2} \int_{\Omega_{e}}|\nabla u(t)|^{2} d x .
$$

Therefore using inequality

$$
\begin{equation*}
C_{1} \varepsilon^{2}\left\|\nabla Q_{\varepsilon} h\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2} \leqslant \sum_{k \in N_{\varepsilon} l} \sum_{l \in \sigma(k)} \varepsilon^{n}\left|h_{k}^{i n}-h_{l}^{i n}\right|^{2} \leqslant C_{2} \varepsilon^{2}\left\|\nabla Q_{\varepsilon} h\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2}, \tag{4.17}
\end{equation*}
$$

we obtain

$$
\begin{equation*}
\sum_{k \in N_{\varepsilon}} \sum_{l \in \sigma(k)} \frac{1}{m_{\varepsilon}}\left[R_{k l}^{\varepsilon}(t)\right]^{2} \leqslant C \varepsilon^{2}\left(\left\|\nabla Q_{\varepsilon} h\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2}+\left\|u^{\varepsilon}(t)\right\|_{1, \varepsilon}^{2}+\|h\|_{1, \varepsilon}^{2}\right) \tag{4.18}
\end{equation*}
$$

Now using (4.13), (4.16) and (4.17) with $u^{\varepsilon}(t)$ instead of $h$ we conclude

$$
\begin{align*}
\left\|\nabla Q_{\varepsilon} u(t)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2} & \leqslant C_{1}\left\|u^{\varepsilon}(t)\right\|_{1, \varepsilon}^{2}+C_{2}\left(\left\|\nabla Q_{\varepsilon} u_{0}\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2}+\left\|u_{0}\right\|_{1, \varepsilon}^{2}\right) e^{-\alpha_{\varepsilon} t}+  \tag{4.19}\\
& +C_{3} \delta^{-1} \int_{0}^{t} e^{-\alpha_{\varepsilon}(t-\tau)}\left\{\left\|u^{\varepsilon}(\tau)\right\|_{1 . \varepsilon}^{2}+\|h\|_{1, \varepsilon}^{2}+\left\|\nabla Q_{\varepsilon} h\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right.}^{2}\right\} d \tau .
\end{align*}
$$

Therefore the assertion of Lemma 4.1 follows from (4.0) and condition (i) of Theorem 3.1.

Remark 4.1. - Assume that

$$
\lambda+\inf _{u \in R} f^{\prime}(u)>0 .
$$

Then we have $\alpha_{\varepsilon}>0$ for $\varepsilon$ and $\delta$ small enough. Therefore for any trajectory $u^{\varepsilon}(t)$ lying in the attractor $\mathcal{G}_{\varepsilon}$ from (4.19) and Remark 2.2 it is easy to see that

$$
\begin{equation*}
\left\|\nabla Q_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2} \leqslant C_{1}+C_{2}\left(1+\left\|\nabla Q_{\varepsilon} u^{\varepsilon}(s)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2}\right) e^{-a_{\varepsilon}(t-s)} \tag{4.20}
\end{equation*}
$$

for all $t \geqslant s$. Since

$$
\left\|\nabla Q_{\varepsilon} u^{\varepsilon}(s)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2}<C_{\varepsilon}, \quad-\infty<s<\infty,
$$

then letting $s \rightarrow-\infty$ in (4.20) we conclude that

$$
\begin{equation*}
\left\|\nabla Q_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2} \leqslant C_{1} \tag{4.21}
\end{equation*}
$$

for any $u^{\varepsilon}(t) \in \mathfrak{Q}_{\varepsilon}$, where $\varepsilon$ is small enough.

## 5. - Limit transition.

Let $R_{1}^{\varepsilon}$ be a linear continuation operator from $\Omega_{\varepsilon}$ to $\Omega$ possesses the properties:
i) $R_{1}^{\varepsilon}: H^{l}\left(\Omega_{\varepsilon}\right) \rightarrow H^{l}(\Omega)$ for $l=0,1$ such that

$$
\left\|R_{1}^{\varepsilon} \phi\right\|_{H^{1}(\Omega)} \leqslant C\|\phi\|_{H^{1}\left(\Omega_{\varepsilon}\right)}, \quad l=0,1
$$

where $C$ is a constant, independent of $\varepsilon$;
ii) $R_{1}^{\varepsilon} \phi=\phi$ on $\Omega_{\varepsilon}$ for all $\phi \in L^{2}\left(\Omega_{\varepsilon}\right)$.

We also denote by $R_{2}^{\varepsilon}$ a continuation operator from $\Omega_{N_{\varepsilon}}$ to $\Omega$ with similar properties. The existence of such operators is easily proved, in view of the structure of the domains $\Omega_{\varepsilon}$ and $\Omega_{N_{\varepsilon}}$.

Let $\bar{u}^{\varepsilon}(t)=R_{1}^{\varepsilon} P_{\varepsilon}^{\varepsilon} u^{\varepsilon}(t)$ and $\bar{v}^{\varepsilon}(t)=R_{2}^{\varepsilon} Q_{\varepsilon} u^{\varepsilon}(t)$, where $u^{\varepsilon}(t)$ is the solution of problem (1)-(3). Then it follows from (4.1)-(4.3) that the family $\left\{\left(\bar{u}^{\varepsilon}(t) ; \bar{v}^{\varepsilon}(t)\right)\right\}$ is a precompact set in the space $C\left(0, T ; L^{2}(\Omega) \times L^{2}(\Omega)\right)$, when $\varepsilon$ goes to zero.

In this section we prove that any limiting point $\left(u(t) ; v(t) \in C\left(0, T ; L^{2}(\Omega) \times\right.\right.$ $\left.\times L^{2}(\Omega)\right)$ of the family $\left\{\left(\bar{u}_{\varepsilon}(t) ; \bar{v}_{\varepsilon}(t)\right): \varepsilon \rightarrow 0\right\}$ is a weak solution of problem (3.5)-(3.8). Below we assume that the conditions (i)-(iii) of Theorem 3.1 are satisfied.

We first rewrite problem (1)-(3) in a weak form. Let $u^{\varepsilon}(t)$ be the solution of (1)-(3). We denote

$$
\begin{align*}
& J_{\varepsilon}\left(u^{\varepsilon} ; \psi ; V_{\varepsilon}\right)=-\left(u_{0}^{\varepsilon}, \psi(0)\right)_{L^{2}\left(V_{\varepsilon}\right)}-\int_{0}^{T}\left(u^{\varepsilon}(t), \partial_{t} \psi(t)\right)_{L^{2}\left(V_{\varepsilon}\right)} d t-  \tag{5.1}\\
& \quad-\int_{0}^{T}\left(u^{\varepsilon}(t), \Delta_{\varepsilon} \psi(t)\right)_{L^{2}\left(V_{\varepsilon}\right)} d t+\int_{0}^{T}\left(f\left(u^{\varepsilon}(t)\right), \psi(t)\right)_{L^{2}\left(V_{\varepsilon}\right)} d t-\int_{0}^{T}(h, \psi(t))_{L^{2}\left(V_{\varepsilon}\right)} d t,
\end{align*}
$$

where $\psi(x, t)$ belongs to the class

$$
\mathfrak{L}_{T}=\left\{\psi(x, t) \in L^{2}\left(0, T ; H^{2}\left(M_{\varepsilon}\right)\right): \partial_{t} \psi(x, t) \in L^{2}\left(0, T ; L^{2}\left(M_{\varepsilon}\right)\right), \psi(T)=0\right\}
$$

and $V_{\varepsilon}$ is a submanifold of $M_{\varepsilon}$. It is clear that

$$
\begin{equation*}
J_{\varepsilon}\left(u^{\varepsilon} ; \psi ; M_{\varepsilon}\right)=0 \quad \text { for all } \psi \in \mathfrak{L}_{T} \tag{5.2}
\end{equation*}
$$

Now we suppose $\psi(t, x)=\beta(t) \xi_{\varepsilon}(x)$, where $\beta(t) \in C^{1}(0, T), \beta(T) \neq 0$, and the function $\xi_{\varepsilon}(x)$ is constructed as follows. If $x \in \Omega_{\varepsilon}$, we set

$$
\xi_{\varepsilon}(x)=\zeta(x)+\sum_{i \in N_{\varepsilon}}(\zeta(i \varepsilon)-\zeta(x)) \varphi\left(\frac{x-i \varepsilon}{4 a_{\varepsilon}}\right)+\sum_{i \in N_{\varepsilon}}(\eta(i \varepsilon)-\zeta(i \varepsilon)) v^{\varepsilon}(x-i \varepsilon) \varphi\left(\frac{x-i \varepsilon}{\varepsilon}\right)
$$

and for $x \in B_{\varepsilon}^{i}$ we suppose

$$
\xi_{\varepsilon}(x)=\eta(i \varepsilon)+(\zeta(i \varepsilon)-\eta(i \varepsilon))\left(1-v^{\varepsilon}(x-i \varepsilon)\right),
$$

where $\zeta(x)$ and $\eta(x)$ are smooth functions on $\Omega$ and $\varphi(x)=\bar{\varphi}(|x|) \in C_{0}^{\infty}\left(\boldsymbol{R}^{n}\right)$ possesses the properties $0 \leqslant \varphi \leqslant 1 ; \bar{\varphi}(r)=1$ for $r \leqslant 1 / 4 ; \bar{\varphi}(r)=0$, if $r \geqslant 1 / 3$. It is clear that $\xi_{\varepsilon}(x)$ is a smooth function on $M_{\varepsilon}$.

Lemma 5.1. - The function $\xi_{\varepsilon}(x)$ has the following properties:

$$
\begin{equation*}
P_{\varepsilon} \xi_{\varepsilon} \rightarrow \zeta(x) \text { strongly in } L^{2}(\Omega), \tag{5.3}
\end{equation*}
$$

$$
\begin{equation*}
P_{\varepsilon} \Delta \xi_{\varepsilon} \rightarrow \Delta \zeta+\lambda \mu(\eta(x)-\zeta(x)) \quad \text { weakly in } L^{2}(\Omega) \tag{5.4}
\end{equation*}
$$

when $\varepsilon \rightarrow 0$.
Proof. - Using Lemma 4.2 (ii) we have

$$
\left\|P_{\varepsilon} \xi_{\varepsilon}-\xi\right\|_{L_{2}\left(\Omega_{\varepsilon}\right)} \leqslant C\left(a_{\varepsilon}^{n+1}+\varepsilon^{n+1}\right)\left|N_{\varepsilon}\right| \rightarrow 0
$$

when $\varepsilon \rightarrow 0$. Here $\left|N_{\varepsilon}\right|$ is a number of elements of $N_{\varepsilon}$. Since $\operatorname{Vol}\left(\Omega \backslash \Omega_{\varepsilon}\right) \rightarrow 0$, we ob$\operatorname{tain}$ (5.3).

In order to prove (5.4) we first note that Remark 3.1 implies that

$$
\begin{equation*}
P_{\varepsilon} \Delta\left(\sum_{i \in N_{\varepsilon}}(\zeta(i \varepsilon)-\zeta(x)) \varphi\left(\frac{x-i \varepsilon}{4 a_{\varepsilon}}\right)\right) \rightarrow 0, \quad \varepsilon \rightarrow 0 \tag{5.5}
\end{equation*}
$$

strongly in $L^{2}(\Omega)$. Therefore it is sufficient to consider the term

$$
\chi_{\varepsilon}(x)=\sum_{i \in N_{\varepsilon}}(\eta(i \varepsilon)-\zeta(i \varepsilon)) v^{\varepsilon}(x-i \varepsilon) \varphi\left(\frac{x-i \varepsilon}{\varepsilon}\right)
$$

It follows from Lemma 4.2 (i), (ii) that

$$
\int_{F_{\varepsilon}}\left|\Delta\left(v^{\varepsilon}(x) \varphi\left(\frac{x}{\varepsilon}\right)\right)\right|^{2} d x \leqslant C \varepsilon^{n}
$$

so the family $\left\{P_{\varepsilon}\left(\Delta \chi_{\varepsilon}\right)\right\}$ is bounded in $L^{2}(\Omega)$. Consequently on this family, weak convergence is the same as weak convergence on smooth functions $\theta(x) \in C_{0}^{\infty}(\Omega)$.

It is clear from Lemma 4.2 (i), (ii) that we have

$$
\begin{equation*}
\int_{G_{\varepsilon}} \Delta\left[v^{\varepsilon}(x) \varphi\left(\frac{x}{\varepsilon}\right)\right] \theta(x) d x=\theta(0) \int_{G_{\varepsilon}} \Delta\left[v^{\varepsilon}(x) \varphi\left(\frac{x}{\varepsilon}\right)\right] d x+O\left(\varepsilon^{n+1}\right) \tag{5.6}
\end{equation*}
$$

Using Green's formula and Lemma 4.2 (v) we get

$$
\int_{G_{\varepsilon}} \Delta\left(v^{\varepsilon}(x) \varphi\left(\frac{x}{\varepsilon}\right)\right) d x=\int_{\Gamma\left(\alpha_{\varepsilon}\right)} \frac{\partial v^{\varepsilon}}{\partial n} d \sigma=\lambda_{\varepsilon} m_{\varepsilon}
$$

Therefore (5.6) gives

$$
\int_{\Omega_{\varepsilon}} P_{\varepsilon} \chi_{\varepsilon} \theta(x) d x=\sum_{i \in N_{\varepsilon}} \lambda_{\varepsilon} m_{\varepsilon}(\eta(i \varepsilon)-\zeta(i \varepsilon)) \theta(i \varepsilon)+O(\varepsilon) .
$$

From this and (5.5) follows (5.4). This proves Lemma 5.1.
If (u;v) is a limit point in $C\left(0, T ; L^{2}(\Omega) \times L^{2}(\Omega)\right)$ of the family $\left\{\left(\bar{u}_{\varepsilon}(t) ; \bar{v}_{\varepsilon}(t)\right)\right.$ : $\varepsilon \rightarrow 0\}$, there exists a sequence $\left\{\varepsilon_{k}\right\}, \varepsilon_{k} \rightarrow 0$, such that

$$
\begin{equation*}
\max _{[0, T]}\left\|u(t)-\widetilde{u}_{\varepsilon_{k}}(t)\right\|_{L^{2}(\Omega)}+\max _{[0, T]}\left\|v(t)-\widetilde{v}_{\varepsilon_{k}}(t)\right\|_{L^{2}(\Omega)} \rightarrow 0, \tag{5.7}
\end{equation*}
$$

when $k \rightarrow \infty$. Therefore Lemma 5.1 implies that

$$
\begin{equation*}
\lim _{k \rightarrow \infty} J_{\varepsilon_{k}}\left(u^{\varepsilon_{k}} ; \beta \xi_{\varepsilon_{k}}, \Omega_{\varepsilon_{k}}\right)=J_{1}(u ; \eta, \zeta) \tag{5.8}
\end{equation*}
$$

where

$$
\begin{align*}
J_{1}(u ; \eta, \zeta) & =-\left(w_{0}, \zeta\right)_{L^{2}(\Omega)} \beta(0)-\int_{0}^{T}(u(t), \zeta)_{L^{2}(\Omega)} \beta^{\prime}(t) d t-  \tag{5.9}\\
& -\int_{0}^{T}(u(t), \Delta \zeta+\lambda \mu(\eta-\zeta))_{L^{2}(\Omega)} \beta(t) d t+\int_{0}^{T}\left(f(u(t))-h_{1}, \zeta\right)_{L^{2}(\Omega)} \beta d t .
\end{align*}
$$

We now study the asymptotic behaviour of $J_{\varepsilon_{k}}\left(u^{\varepsilon_{k}} ; \beta \xi_{\varepsilon_{k}}, M_{\varepsilon_{k}} \backslash \Omega_{\varepsilon_{k}}\right)$.
Lemma 5.2. - Let

$$
\gamma_{\varepsilon} \equiv \sum_{j \in N_{\varepsilon} B_{\varepsilon}^{\prime}} \int_{B_{i}} f\left(u^{\varepsilon}(x, t)\right) \xi_{\varepsilon}(x) d x
$$

Then for any interval $[0, T]$ we have

$$
\begin{equation*}
\left|\gamma_{\varepsilon}-\sum_{j \in N_{\varepsilon}} f\left(u_{j}^{\varepsilon, i n}(t)\right) \eta(j \varepsilon) m_{\varepsilon}\right| \leqslant C_{T} \varepsilon, \quad t \in[0, T] . \tag{5.10}
\end{equation*}
$$

Proof. - The Poincaré inequality and the structure of $\xi_{\varepsilon}$ on $B_{\varepsilon}^{j}$ give $\left|\int_{B_{\varepsilon}^{l}} f\left(u^{\varepsilon}\right) \xi_{\varepsilon}(x) d x-u_{j}^{\varepsilon, i n}(t) \int_{B_{\varepsilon}^{j}} \xi_{\varepsilon}(x) d x\right| \leqslant \varepsilon\left\|\nabla_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(B_{\varepsilon}^{\prime}\right)}\left(\left\|v^{\varepsilon}\right\|_{L^{2}\left(B_{\varepsilon}^{j}\right)}+\varepsilon\left\|\nabla_{\varepsilon} v^{\varepsilon}\right\|_{L^{2}\left(B_{\varepsilon}^{j}\right)}\right)$.

Therefore it follows from Lemma 4.24(iii), (iv) and (4.10) that we have

$$
\begin{equation*}
\left|\gamma_{\varepsilon}-\sum_{j \in N_{\varepsilon}} f\left(u_{j}^{\varepsilon, i n}(t)\right) \int_{B_{\varepsilon}^{j}} \xi_{\varepsilon}(x) d x\right| \leqslant C_{T} \varepsilon, \quad t \in[0, T] . \tag{5.11}
\end{equation*}
$$

It is also clear from (2.2), the Poincaré inequality and Lemma 4.2 (iv) that

$$
\begin{aligned}
\mid \int_{B_{\varepsilon}^{j}} f\left(u_{j}^{\varepsilon, i n}\right)(\zeta(j \varepsilon)-\eta(j \varepsilon))(1 & \left.-v^{\varepsilon}(x-j \varepsilon)\right) d x \mid \leqslant \\
& \leqslant C \varepsilon^{n+1}\left(1+\mid u_{j}^{\varepsilon,}, i n\right. \\
& t) \mid) \leqslant C \varepsilon^{n / 2+1}\left(\varepsilon^{n / 2}+\left\|u^{\varepsilon}(t)\right\|_{L^{2}\left(B_{\varepsilon}^{j}\right)}\right)
\end{aligned}
$$

So (5.10) follows from (5.11).
As above we can conclude that

$$
\begin{equation*}
\left|\sum_{j \in N_{\varepsilon}}\left\{\int_{B_{\varepsilon}^{j}} u^{\varepsilon}(t) \xi_{\varepsilon} d x-u_{j}^{\varepsilon, i n}(t) \eta(j \varepsilon) m_{\varepsilon}\right\}\right| \leqslant C_{1} \varepsilon \tag{5.12}
\end{equation*}
$$

for all $t \in[0, T]$ and

$$
\begin{equation*}
\left|\sum_{j \in N_{\varepsilon}}\left\{\int_{B_{\varepsilon}^{j}} h \xi_{\varepsilon} d x-h(j \varepsilon) \eta(j \varepsilon) m_{\varepsilon}\right\}\right| \leqslant C_{2} \varepsilon . \tag{5.13}
\end{equation*}
$$

Using the equality

$$
\Delta_{\varepsilon} \xi_{\varepsilon}(x)=\lambda_{\varepsilon}(\zeta(j \varepsilon)-\eta(j \varepsilon)) v^{\varepsilon}(x-j \varepsilon), \quad x \in B_{\varepsilon}^{j}
$$

and Lemma 4.2 it is easy to see that

$$
\begin{equation*}
\sum_{j \in N_{\varepsilon}} \int_{B_{\varepsilon}^{i}} u^{\varepsilon}(t, x) \Delta_{\varepsilon} \xi_{\varepsilon}(x) d x=\lambda_{\varepsilon} m_{\varepsilon} \sum_{j \in N_{\varepsilon}} u_{j}^{\varepsilon, \text { in }}(t)(\zeta(j \varepsilon)-\eta(j \varepsilon))+O(\varepsilon) . \tag{5.14}
\end{equation*}
$$

Thus, it follows from (5.7), (5.10), (5.12)-(5.14) that

$$
\begin{equation*}
\lim _{k \rightarrow \infty} J_{\varepsilon_{k}}\left(u^{\varepsilon_{k}} ; \beta \xi_{\varepsilon_{k}} M_{\varepsilon_{k}} \backslash \Omega_{\varepsilon_{k}}\right)=\mu J_{2}(v ; \eta, \xi), \tag{5.15}
\end{equation*}
$$

where

$$
\begin{align*}
J_{2}(v ; \eta, \zeta)= & -\left(v_{0}, \eta\right)_{L^{2}(\Omega)} \beta(0)-\int_{0}^{T}(v(t), \eta)_{L^{2}(\Omega)} \beta^{\prime}(t) d t-  \tag{5.16}\\
& -\lambda \int_{0}^{T}(v(t), \xi-\eta)_{L^{2}(\Omega)} \beta^{\prime}(t) d t+\int_{0}^{T}\left(f(v(t))-h_{2}, \eta\right)_{L^{2}(\Omega)} \beta(t) d t
\end{align*}
$$

Equations (5.2), (5.8) and (5.15) imply that

$$
\begin{equation*}
J_{1}(u ; \eta, \zeta)+\mu J_{2}(v ; \eta, \zeta)=0 \tag{5.17}
\end{equation*}
$$

for any limiting point $(u(t) ; v(t)) \in C\left(0, T ; L^{2}(\Omega) \times L^{2}(\Omega)\right)$ of the family $\left\{\left(\bar{u}_{\varepsilon}(t) ; \bar{v}_{\varepsilon}(t)\right): \varepsilon \rightarrow 0\right\}$. Here $J_{1}(u ; \eta, \zeta)$ and $J_{2}(v ; \eta, \zeta)$ are defined by (5.9) and (5.16), where $\beta(t) \in C^{1}(0, T), \beta(T)=0$ and $\eta(x), \zeta(x)$ are any smooth functions on $\Omega$. There-
fore it follows easily from (5.17) that $(u(t) ; v(t))$ is a weak solution of problem (3.5)-(3.8).

Remark 5.1. - Existence theorem for solutions from $C\left(0, T ; L^{2}(\Omega) \times L^{2}(\Omega)\right)$ of the problem (3.5)-(3.8) also follow from the considerations above, under certain conditions concerning the functions $u_{0}, v_{0}, h_{1}$ and $h_{2}$ (see the assumption (ii) of Theorem 3.1).

In order to complete the proof of Theorem 3.1 we only need to prove the uniqueness theorem for the system (3.5)-(3.8). We will do this in the following section.

## 6. - Properties of the homogenized system.

In this section we prove Theorems 2.2-2.4. We rewrite equations (3.5)-(3.8) as the following system of first order evolution equations in the space $\mathscr{F}_{0}=L^{2}(\Omega) \times$ $\times L^{2}(\Omega)$ :

$$
\begin{equation*}
\frac{d}{d t} U+A U=B(U),\left.\quad U\right|_{t=0}=U_{0} \tag{6.1}
\end{equation*}
$$

where

$$
U=\binom{u}{v}, \quad A=\left(\begin{array}{cc}
-\Delta+\lambda \mu & 0 \\
0 & \lambda
\end{array}\right), \quad B(U)=\binom{\lambda \mu v-f(u)+h_{1}}{\lambda u-f(v)+h_{2}} .
$$

It is easy to see that $A$ is a positive self-adjoint operator in $\mathscr{F}_{0}$ such that

$$
\begin{equation*}
(A U, U)_{\mathcal{J}_{0}} \geqslant\|\nabla u\|_{L^{2}(\Omega)}+\gamma \lambda\|U\|_{\mathcal{J}_{0}}^{2}, \quad U \in \mathcal{O}\left(A^{1 / 2}\right), \tag{6.2}
\end{equation*}
$$

where $\gamma=\min (1, \mu)$, and

$$
\begin{equation*}
\|B(U)\|_{\mathscr{J}_{0}} \leqslant M_{1}(1+\|U\|)_{J_{0}}, \quad\left\|B\left(U_{1}\right)-B\left(U_{2}\right)\right\| \leqslant M_{2}\left\|U_{1}-U_{2}\right\|_{\widetilde{J}_{0}} \tag{6.3}
\end{equation*}
$$

If we consider de equation (6.1) in the integral form

$$
\begin{equation*}
U(t)=e^{-A t} U_{0}+\int_{0}^{t} e^{-A(t-\tau)} B(U(\tau)) d \tau \tag{6.4}
\end{equation*}
$$

then using the fixed point method in the space $C\left(0, T ; \mathscr{F}_{0}\right)$ we can easily prove the existence and uniqueness of solutions for $T<T_{0}$, when $T_{0}$ is small enough. It is clear that the function $U(t)$ gives a generalized solution of the system (3.5)-(3.8) on the interval [ $0, T], T<T_{0}$. Using standard methods (see, e.g. [9,15]) and the properties (2.2)-(2.4) of the function $f(u)$ we see that

$$
\begin{equation*}
\|U(t)\|_{\mathcal{J}_{0}} \leqslant C_{1}\left\|U_{0}\right\|_{\mathscr{H}_{0}} e^{-\omega t}+C_{2}\left(1-e^{-\omega t}\right), \tag{6.5}
\end{equation*}
$$

where $\omega, C_{1}$ and $C_{2}$ are positive constants. This estimate allows us to extend the sol-
ution $U(t)$ on the whole of $\boldsymbol{R}_{+}$. The proofs of the properties (3.10) and (3.11) are also of standard character (for similar consideration see, e.g. [9,15]). This proves Theorem 3.2.

Let $S_{t}$ be the evolutionary semigroup defined by the formula $S_{t} U_{0}=U(t)$, where $U(t)$ is the solution of problem (6.1). Since

$$
\left\|A^{\beta} e^{-t A}\right\| \leqslant C t^{-\beta} e^{-\lambda \gamma t}, \quad t>0, \quad 0<\beta<1,
$$

and $\mathcal{O}\left(A^{1 / 2}\right)=\mathscr{F}_{1}=H^{1}(\Omega) \times L^{2}(\Omega)$, (6.4) and (6.5) imply that $S_{t}$ has the following the following dissipativity property: there exists a constant $R>0$ such that for any bounded set $B$ in $\mathscr{F}_{0}$ we have

$$
\begin{equation*}
\left\|S_{t} U_{0}\right\|_{\widetilde{\sigma}_{1}} \leqslant R \quad \text { for all } U_{0} \in B \text { and } t \geqslant t_{0}(B) . \tag{6.6}
\end{equation*}
$$

Lemma 6.1. - Assume that (3.12) is satisfied. Then $S_{t}$ is a dissipative semigroup in the space $\mathscr{F}_{2}=H^{1}(\Omega) \times H^{1}(\Omega)$, i.e. there exists a constant $R^{*}>0$ such that for any bounded set $B$ in $\mathscr{F}_{2}$ we have

$$
\begin{equation*}
\left\|S_{t} U_{0}\right\|_{\sqrt{2}^{2}} \leqslant R^{*} \quad \text { for all } U_{0} \in B \text { and } t \geqslant t_{0}(B) . \tag{6.7}
\end{equation*}
$$

Proof. - Using (3.7) and (3.8) we see that the function $w_{k}(x, t)=\partial_{x_{k}} v(x, t)$ satisfies the equation

$$
\frac{d}{d t} w_{k}(t)+\left(\lambda+f^{\prime}(v(t))\right) w_{k}=\lambda \partial_{x_{k}} u+\partial_{x_{k}} h_{2}
$$

Therefore from (3.12) we get

$$
\frac{1}{2} \frac{d}{d t}\left\|w_{k}(t)\right\|_{L^{2}(\Omega)}^{2}+\delta\left\|w_{k}(t)\right\|_{L^{2}(\Omega)}^{2} \leqslant C\left(\|u\|_{H^{1}(\Omega)}^{2}+\left\|h_{2}\right\|_{H^{2}(\Omega)}^{2}\right)
$$

where $\delta>0$. Using (6.6) we obtain

$$
\left\|\partial_{x_{k}} v(t)\right\|_{L^{2}(\Omega)}^{2} \leqslant\|v(s)\|_{H^{1}(\Omega)}^{2} e^{-2 \delta(t-s)}+C_{R}, \quad t \geqslant s \geqslant t_{0}(B) .
$$

This estimate and (6.6) imply (6.7).
Lemma 6.2. - The semigroup $S_{t}$ is weakly closed in $\mathfrak{F}_{2}$, i.e. for any $t>0$, the conditions: $U_{n} \rightarrow U$ and $S_{t} U_{n} \rightarrow V$ weakly in $\mathscr{F}_{2}$ for $n \rightarrow \infty$ impliy $V=S_{t} U$.

Proof. - The lemma follows from equation (6.4) and from the compactness of the imbedding $\mathscr{F}_{2} \rightarrow \mathscr{F}_{0}$.

Lemma 6.1 and 6.2 make it possible to use the results from [1] and to guarantee the existence of weak global attractor $\mathfrak{a}$ for the dynamical system $\left(S_{t}, \mathscr{F}_{2}\right)$. This attractor is a bounded weakly closed set in $\mathscr{F}_{2}=H^{1}(\Omega) \times H^{1}(\Omega)$. It is also easy to see that the initial data from $\mathscr{F}_{2}$ is $C^{1}$ with respect to the semigroup $S_{t}$. Therefore in order to prove the finiteness of the Hausdorff dimension of $\mathfrak{a}$ we can use the approach presented in [17].

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Let us consider the first variation equation corresponding to (6.1):

$$
\frac{d}{d t} W+A W=B^{\prime}(U(t)) W
$$

for a trajectory $U(t)$ lying on the attractor $\mathfrak{a}$. As in [17] (see also [4]) it is necessary to estimate the quantity

$$
\sigma_{N}(t)=\operatorname{tr}\left\{(A-B(U(t))) Q_{N}\right\}
$$

for any $N$ dimensional orthoprojector $Q_{N}$ in the space $\mathscr{F}_{0}$ such that $Q_{N} \mathscr{F}_{0} \subset \mathscr{F}_{1}$. It is clear that for $W=\left(w_{1} ; w_{2}\right) \in \mathscr{F}_{1}$ we have

$$
\left(\left[A-B^{\prime}(U(t))\right] W, W\right)_{J_{0}} \geqslant\left\|\nabla w_{1}\right\|^{2}+\alpha_{\delta}\left\|w_{1}\right\|^{2}+\beta_{\delta}\left\|\nabla w_{2}\right\|^{2},
$$

where

$$
\alpha_{\delta}=\lambda_{\mu}-\frac{\lambda^{2}\left(1+\mu^{2}\right)}{4 \delta}+\inf _{u} f^{\prime}(u)
$$

and

$$
\beta_{\delta}=\lambda+\inf _{u} f^{\prime}(u)-\delta
$$

for any $\delta>0$ such that $\beta_{\delta}>0$. Let $\left\{W^{k}=\left(w_{1}^{k} ; w_{2}^{k}\right)\right\}_{k=1}^{N}$ be an orthonormal basis of $Q_{N} \mathscr{F}_{0}$. Using the equality

$$
\sum_{k=1}^{N}\left\|w_{2}^{k}\right\|^{2}=N-\sum_{k=1}^{N}\left\|w_{1}^{k}\right\|^{2}
$$

we get

$$
\sigma_{N}(t)=\beta_{\delta} N+\sum_{k=1}^{N}\left\|w_{1}^{k}\right\|^{2}+\left(a_{\delta}-\beta_{\delta}\right) \sum_{k=1}^{N}\left\|w_{2}^{k}\right\|^{2}
$$

Now we use the following version of the Sobolev-Lieb-Thirring inequality

$$
k_{1} \sum_{k=1}^{N}\left\|w_{1}^{k}\right\|^{2}+\frac{k_{2}}{[d(\Omega)]^{2}} \int_{\Omega} \varrho(x) d x \geqslant \int_{\Omega} \varrho(x)^{1+2 / n} d x
$$

which follows from [6, Theorem 2.1]. Here $\varrho(x)=\sum_{k=1}^{N}\left[w_{1}^{k}\right]^{2}, k_{1}$ and $k_{2}$ are constants depending on $n$ and on the shape of $\Omega, d(\Omega)$ is the diameter of $\Omega$. We obtain

$$
\sigma_{N}(t)=\beta_{\delta} N+\int_{\Omega}\left\{\frac{1}{k_{1}} \varrho(x)^{1+2 / n}-\gamma \varrho(x)\right\} d x
$$

where

$$
\gamma=\frac{k_{2}}{k_{1}}[d(\Omega)]^{-2}+\lambda+\frac{\lambda^{2}\left(1+\mu^{2}\right)}{4 \delta} .
$$

Since

$$
z^{1+2 / n}-\gamma k_{1} z \geqslant-\frac{2}{n}\left(\frac{\gamma k_{1} n}{n+2}\right)^{1+n / 2}
$$

for any $z>0$, we have

$$
\sigma_{N}(t) \geqslant \beta_{\delta} N-\frac{|\Omega|}{k_{1}} \frac{2}{n}\left(\frac{\gamma k_{1} n}{n+2}\right)^{1+n / 2}
$$

Therefore (see, e.g. [17]) an estimate for the Hausdorff dimension of $\mathfrak{G}$ as a compact set in $\mathscr{F}_{0}$ can be found from the condition

$$
N>\frac{2}{n} \frac{|\Omega|}{k_{1} \beta_{\delta}}\left(\frac{\gamma k_{1} n}{n+2}\right)^{1+n / 2}
$$

This proves Theorem 3.3.
Remark 6.1. - It is easy to see that the function

$$
\begin{aligned}
& V(u, v)=\frac{1}{2}\left(\|\nabla u\|_{L^{2}(\Omega)}^{2}+\lambda \mu\|u-v\|_{L^{2}(\Omega)}^{2}\right)+ \\
&+\int_{\Omega}(F(u)+\mu F(v)) d x-\left(h_{1}, u\right)_{L^{2}(\Omega)}^{2}-\mu\left(h_{2}, v\right)_{L^{2}(\Omega)}^{2}
\end{aligned}
$$

is continuous on $\mathscr{F}_{2}$ and has the following properties

$$
V(u(t), v(t))+\int_{0}^{t}\left(\left\|\frac{\partial u}{\partial t}(\tau)\right\|_{L^{2}(\Omega)}^{2}+\mu\left\|\frac{\partial u}{\partial t}(\tau)\right\|_{L^{2}(\Omega)}^{2}\right) d t=V\left(u_{0}, v_{0}\right)
$$

where $(u(t), v(t))$ is the solution of problem (3.5)-(3.8). This property implies (see, e.g. $[1,4,7,17]) \mathfrak{a}=\Re_{+}(\mathcal{N})$, where $\mathcal{N}$ is the set of stationary solutions of the system (3.5)-(3.8) and $\pi_{+}(\mathcal{N})$ is the unstable manifold of $\mathcal{N}$. In particular this means that any trajectory of $S_{t} U_{0}$ goes to $\mathcal{N}$, when $t \rightarrow+\infty$.

Remark 6.2. - The assumption (3.12) is of prime importance in the proof of Theorem 3.3. Indeed, for any $\delta>0$ it is easy to find a function $f(u)$ satisfying (2.2)-(2.4) such that

$$
f\left(v_{0}\right)=0, \quad \lambda+f^{\prime}\left(v_{0}\right)=-\delta
$$

for some $v_{0} \in \boldsymbol{R}$. In this case the pair ( $v_{0} ; v_{0}$ ) gives a stationary solution of (3.5)-(3.8). The linearization of (3.5) and (3.7) near ( $v_{0} ; v_{0}$ ) has the form

$$
\frac{d}{d t}\binom{u}{v}+A_{v_{0}}\binom{u}{v}=0, \quad \text { where } A_{v_{0}}=\left(\begin{array}{cc}
-\Delta+\lambda \mu+f^{\prime}\left(v_{0}\right) & -\lambda \mu \\
-\lambda & \lambda+f^{\prime}\left(v_{0}\right)
\end{array}\right)
$$

Simple considerations show that for $\mu=1$ the operator $A_{v_{0}}$ has an infinite dimensional spectral subspace corresponding to eigenvalues belonging to $\{\lambda: \operatorname{Re} \lambda<0\}$. Therefore the instable manifold $\mathscr{N}_{+}\left(v_{0} ; v_{0}\right)$ has infinite dimension. Thus the asymptotic behaviour of the system (3.5)-(3.8) without assumption (3.12) cannot be described by a finite-dimensional global attractor.

Now we prove Theorem 3.4 on the upper semicontinuity of the family $\left\{\mathfrak{C}_{\varepsilon}: \varepsilon>0\right\}$ of attractors for the problem (1)-(3), when $\varepsilon \rightarrow 0$.

It follows from Remarks 2.2 and 4.1 that for any trajectory $\left\{u^{\varepsilon}(t)-\infty<t<\infty\right\}$, belonging to the attractor $\mathfrak{a}_{\varepsilon}$ we have the uniform estimates:

$$
\begin{equation*}
\left\|P_{\varepsilon} \frac{\partial u^{\varepsilon}}{\partial t}\right\|_{L^{2}\left(\Omega_{\varepsilon}\right)}^{2}+\left\|\nabla P_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{\varepsilon}\right)}^{2}+\left\|P_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{\varepsilon}\right)}^{2} \leqslant C_{1}, \tag{6.8}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|Q_{\varepsilon} \frac{\partial u^{\varepsilon}}{\partial t}\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2}+\left\|\nabla Q_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2}+\left\|Q_{\varepsilon} u^{\varepsilon}(t)\right\|_{L^{2}\left(\Omega_{N_{\varepsilon}}\right)}^{2} \leqslant C_{2} \tag{6.9}
\end{equation*}
$$

for all $t \in(-\infty, \infty)$ and $\varepsilon$ small enough.
Let $u_{0}^{\varepsilon} \in \mathfrak{C}_{\varepsilon}$. Then there exist a trajectory $\left\{u^{\varepsilon}(t):-\infty<t<\infty\right\} \subset \mathfrak{C}_{\varepsilon}$ such that $u^{\varepsilon}(0)=u_{0}^{\varepsilon}$ and (6.8) and (6.9) are satisfied. Therefore, as in Section 5 we can find a solution (u(t); v(t)) of (3.5)-(3.8) belonging to $C\left(a, b ; \mathscr{F}_{0}\right)$ for any $a, b$ such that $-\infty<a<a<+\infty$,

$$
\begin{equation*}
\max _{[a, b]}\left\|P_{\varepsilon_{k}} u^{\varepsilon_{k}}(t)-u(t)\right\|_{L^{2}(\Omega)}+\max _{[a, b]}\left\|Q_{\varepsilon_{k}} u^{\varepsilon_{k}}(t)-v(t)\right\|_{L^{2}(\Omega)} \rightarrow 0 \tag{6.10}
\end{equation*}
$$

for some subsequence $\left\{\varepsilon_{k}\right\}, \varepsilon_{k} \rightarrow 0$. From (6.8) and (6.9) we also get

$$
\left\|\frac{\partial u^{\varepsilon}}{\partial t}\right\|_{L^{2}\left(\Omega_{\varepsilon}\right)}^{2}+\|\nabla u\|_{L^{2}\left(\Omega_{\varepsilon}\right)}^{2}+\|u\|_{L^{2}\left(\Omega_{e}\right)}^{2} \leqslant C_{1},
$$

and

$$
\left\|\frac{\partial u^{\varepsilon}}{\partial t}\right\|_{L^{2}(\Omega)}^{2}+\|\nabla v\|_{L^{2}(\Omega)}^{2}+\|v\|_{L^{2}(\Omega)}^{2^{2}} \leqslant C_{2},
$$

for all $t$. Consequently $U(t)=(u(t) ; v(t))$ belongs to a weak global attractor $\mathfrak{a}$. Therefore from (6.10) it is easy to extract the assertion of Theorem 3.4.

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