# Non-variational weakly coupled elliptic systems 

Mónica Clapp ${ }^{1}$ • Andrzej Szulkin ${ }^{2}$

Received: 26 July 2021 / Accepted: 23 February 2022 / Published online: 19 March 2022
© The Author(s) 2022

To the memory of Harold S. Shapiro.
A.S. was Harold's student.

He is forever grateful for all inspiration and encouragement.


#### Abstract

We establish the existence of a nonnegative fully nontrivial solution to a non-variational weakly coupled competitive elliptic system. We show that this kind of solutions belong to a topological manifold of Nehari-type, and apply a degree-theoretical argument on this manifold to derive existence.


Keywords Weakly coupled elliptic system • Positive solution • Uniform bound • Nehari manifold • Brouwer degree • Synchronized solutions

Mathematics Subject Classification 35J57 - 35J61 - 35B09 • 47H11

## 1 Introduction and statement of results

In this paper we consider the existence of solutions to the elliptic system

$$
\left\{\begin{array}{l}
-\Delta u_{i}=\mu_{i} u_{i}^{p}+\sum_{j \neq i} \lambda_{i j} u_{i}^{\alpha_{i j}} u_{j}^{\beta_{i j}},  \tag{1.1}\\
u_{i} \geq 0, u_{i} \not \equiv 0 \text { in } \Omega, \\
u_{i} \in H_{0}^{1}(\Omega), \quad i, j=1, \ldots, \ell,
\end{array}\right.
$$

[^0]where $\Omega$ is a smooth bounded domain in $\mathbb{R}^{N}, N \geq 2,1<p<\frac{N+2}{N-2}$ if $N \geq 3$, $1<p<\infty$ if $N=2, \mu_{i}>0, \lambda_{i j}<0, \alpha_{i j}, \beta_{i j}>0$ and $\alpha_{i j}+\beta_{i j}<p$ for $i, j=1, \ldots, \ell, j \neq i$. This system arises as a model for the steady state distribution of $\ell$ competing species coexisting in $\Omega$. Here $u_{i}$ represents the density of the $i$-th population, $\mu_{i}$ corresponds to the attraction between the species of the same kind, or more generally, $\mu_{i} u_{i}^{p}$ can be replaced by $f_{i}\left(u_{i}\right)$ and represent internal forces. The parameters $\lambda_{i j}, \lambda_{j i}$ (which may not be equal) correspond to the interaction (repulsion) between different species. In particular, if $\alpha_{i j}=\beta_{i j}=1$, then the interaction is of the Lotka-Volterra type while $\alpha_{i j}=1, \beta_{i j}=2$ corresponds to the interaction which appears in the Bose-Einstein condensates. In the latter case one also has $\lambda_{i j}=\lambda_{j i}$ and the system is variational.

In what follows we do not assume $\lambda_{i j}=\lambda_{j i}$ or $\beta_{i j}=\alpha_{j i}$. The system (1.1) is non-variational except for some very special choices of $\lambda_{i j}, \alpha_{i j}$ and $\beta_{i j}$. While there is an extensive literature concerning the existence (and multiplicity) of solutions for variational systems like (1.1), there are not so many results in the non-variational case. Here we could mention [1, 6-9] where, however, the right-hand sides are quite different from ours. In particular, in [6-9] the interaction term is of the Lotka-Volterra type (or is a variant of it) while the terms $f_{i}\left(u_{i}\right)$ are different from $\mu_{i} u_{i}^{p}$. For these $f_{i}$ one obtains uniform bounds on the solutions when $\lambda_{i j} \rightarrow-\infty$. Existence of such bounds allows to study the limiting behaviour of solutions. To be more precise, if $\lambda_{i j, n} \rightarrow-\infty$ and ( $u_{1, n}, \ldots, u_{l, n}$ ) is a corresponding solution with uniform bound on each component, then one expects that $u_{i, n} \rightarrow u_{i}$ (in an appropriate space) and $u_{i}(x) \cdot u_{j}(x)=0$ a.e. in $\Omega$ for all $i \neq j$, i.e. different components separate spatially. This has been studied in the above mentioned papers. In [1,6] the emphasis is in fact on the properties of limiting configurations, including regularity of free boundaries between the components.

The main result of this paper is the following
Theorem 1.1 The system (1.1) has a solution.
Existence proofs in the above-mentioned papers do not seem to be applicable here. Our problem can be reformulated as an operator equation in the space $\mathcal{H}:=H_{0}^{1}(\Omega)^{\ell}$ and one can use degree theory to obtain a nontrivial solution. However, this could give a semitrivial solution (i.e. $u_{i}=0$ for some but not all $i$ ). To rule out such solutions we introduce a Nehari-type manifold on which all $u$ are fully nontrivial in the sense that no $u_{i}$ is identically zero, and then we apply a degree-theoretical argument on this manifold.

We do not know if there always exist solutions for (1.1) which are uniformly bounded, see Problem 5.5. Moreover, as we shall see in Sect. 5, under a suitable choice of exponents and parameters and for $\ell=2$ there exists a sequence of solutions which are synchronized in the sense that $u_{i, n}=t_{i, n} v_{n}(i=1,2)$ and such that $\left\|u_{i, n}\right\| \rightarrow \infty$ as $\lambda_{12, n}, \lambda_{21, n} \rightarrow-\infty$. So the components neither separate spatially nor are bounded.

Let $u_{i}^{+}:=\max \left\{u_{i}, 0\right\}, u_{i}^{-}:=\min \left\{u_{i}, 0\right\}$, and consider the system

$$
\left\{\begin{array}{l}
-\Delta u_{i}=\mu_{i}\left(u_{i}^{+}\right)^{p}+\sum_{j \neq i} \lambda_{i j}\left(u_{i}^{+}\right)^{\alpha_{i j}}\left(u_{j}^{+}\right)^{\beta_{i j}},  \tag{1.2}\\
u_{i} \in H_{0}^{1}(\Omega), \quad i, j=1, \ldots, \ell .
\end{array}\right.
$$

In Proposition 3.4(v) we shall show that any fully nontrivial solution to this system also solves (1.1).

In what follows we shall work with (1.2) and we shall also need the parametrized system

$$
\left\{\begin{array}{l}
-\Delta u_{i}=\mu_{i}\left(u_{i}^{+}\right)^{p}+t \sum_{j \neq i} \lambda_{i j}\left(u_{i}^{+}\right)^{\alpha_{i j}}\left(u_{j}^{+}\right)^{\beta_{i j}}  \tag{1.3}\\
u_{i} \in H_{0}^{1}(\Omega), \quad i, j=1, \ldots, \ell, \quad 0 \leq t \leq 1
\end{array}\right.
$$

Note that (1.3) homotopies (1.2) to an uncoupled system. Since

$$
t \sum_{j \neq i}\left|\lambda_{i j}\right|\left(u_{i}^{+}\right)^{\alpha_{i j}}\left(u_{j}^{+}\right)^{\beta_{i j}} \leq C\left(1+\left(u_{1}^{+}\right)^{q}+\cdots+\left(u_{\ell}^{+}\right)^{q}\right),
$$

where $\alpha_{i j}+\beta_{i j} \leq q<p$ for all $i, j$, the following statement holds true.
Lemma 1.2 All solutions $u=\left(u_{1}, \ldots, u_{\ell}\right)$ of (1.3) are uniformly bounded in $L^{\infty}(\Omega)$ and hence in $H_{0}^{1}(\Omega)$. This bound is independent of $t \in[0,1]$.

This has been shown, in a much more general setting, in [13] for a single equation and in [11] for two equations. It is easy to see that the argument in [11] extends to an arbitrary number of equations. In both papers a blow-up procedure is used in order to reduce the problem to a Liouville-type result. For the reader's convenience, in Appendix A we shall provide a simple proof of such reduction, adapted to our special case. The assumption $q<p$ is crucial for the validity of this lemma. Indeed, in [10] it has been shown that the conclusion may fail if $q=p$.

The paper is organized as follows. In Sect. 2 we state and prove a lemma for functions in $\mathbb{R}^{\ell}$. In Sect. 3 we define a Nehari-type manifold $\mathcal{N}$ similar to the one introduced in [5]. We also show that solutions to (1.2) correspond to solutions for an operator equation in an open subset of the product of the unit spheres $\mathscr{S}_{i} \subset H_{0}^{1}(\Omega)$, $1 \leq i \leq \ell$. The idea comes from [4]. To our knowledge, this is the first time a Neharitype manifold appears in a non-variational setting. Theorem 1.1 is proved in Sect. 4 and synchronized solutions are discussed in Sect. 5. As we have already mentioned, Lemma 1.2 is proved in Appendix A.

In the proof of Theorem 1.1 we shall employ a topological degree argument. Since our operator is not admissible for common infinite-dimensional degree theories, we introduce a sequence of finite-dimensional ("Galerkin-like") approximations and use the Brouwer degree, see (4.6) and (4.9-4.11) below.

## 2 A lemma on functions in $\mathbb{R}^{\ell}$

Let $a_{i}, \alpha_{i j}, \beta_{i j}>0, b_{i}, d_{i j} \geq 0, \alpha_{i j}+\beta_{i j}<p$ for all $i, j=1, \ldots, \ell, j \neq i$. Define $M:(0, \infty)^{\ell} \rightarrow \mathbb{R}^{\ell}$ as

$$
M(s):=\left(M_{1}(s), \ldots, M_{\ell}(s)\right),
$$

where

$$
M_{i}(s):=a_{i} s_{i}-b_{i} s_{i}^{p}+\sum_{j \neq i} d_{i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}}, \quad i, j=1, \ldots, \ell .
$$

Lemma 2.1 (i) If $b_{i}=0$ for some $i$, then $M(s) \neq 0$ for any $s \in(0, \infty)^{\ell}$.
(ii) If $b_{i}>0$ for all $i$, then there exists $s \in(0, \infty)^{\ell}$ such that $M(s)=0$.

Moreover, if $0<a \leq a_{i} \leq \bar{a}, 0<b \leq b_{i} \leq \bar{b}$ and $d_{i j} \leq \bar{d}$ for all $i, j$, then there exist $0<r<R$, depending only on $a, \bar{a}, b, \bar{b}, \bar{d}$, such that $s \in(r, R)^{\ell}$.
(iii) The solution $s$ in (ii) is unique.
(iv) The solution $s$ in (ii) depends continuously on $a_{i}, b_{i}>0, d_{i j} \geq 0$.

Proof (i): If $b_{i}=0$ then

$$
M_{i}(s)=a_{i} s_{i}+\sum_{j \neq i} d_{i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}}>0 \quad \text { for all } s \in(0, \infty)^{\ell} .
$$

(ii) : Let $0<r<R$ be such that, for every $i, j=1, \ldots, \ell$,

$$
\begin{aligned}
a_{i} t-b_{i} t^{p}>0 & \text { if } t \in(0, r] \\
a_{i} t-b_{i} t^{p}+\sum_{j \neq i} d_{i j} t^{\alpha_{i j}+\beta_{i j}}<0 & \text { if } t \in[R, \infty)
\end{aligned}
$$

(such $R$ exists because $\alpha_{i j}+\beta_{i j}<p$ ). If $s=\left(s_{1}, \ldots, s_{\ell}\right) \in(0, \infty)^{\ell}$ and $s_{i} \geq s_{j}$ for all $j$, then

$$
M_{i}(s)=a_{i} s_{i}-b_{i} s_{i}^{p}+\sum_{j \neq i} d_{i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}} \leq a_{i} s_{i}-b_{i} s_{i}^{p}+\sum_{j \neq i} d_{i j} s_{i}^{\alpha_{i j}+\beta_{i j}} .
$$

Therefore, $M_{i}(s)<0$ whenever $s_{i}=\max \left\{s_{1}, \ldots, s_{\ell}\right\} \geq R$, and $M_{i}(s)>0$ if $0<s_{i} \leq r$. If $a \leq a_{i} \leq \bar{a}, b \leq b_{i} \leq \bar{b}, d_{i j} \leq \bar{d}$, then
$a_{i} t-b_{i} t^{p} \geq a t-\bar{b} t^{p}, \quad a_{i} t-b_{i} t^{p}+\sum_{j \neq i} d_{i j} t^{\alpha_{i j}+\beta_{i j}} \leq \bar{a} t-b t^{p}+\sum_{j \neq i} \bar{d} t^{\alpha_{i j}+\beta_{i j}}$,
so $r, R$ may be chosen as claimed.
Let

$$
G(s):=\rho-s \quad \text { where } \quad \rho:=\frac{r+R}{2}(1, \ldots, 1) \text {. }
$$

Then $H(s, \tau):=\tau M(s)+(1-\tau) G(s) \neq 0$ on the boundary of $[r, R]^{\ell}$ for every $\tau \in[0,1]$. Hence this is an admissible homotopy for the Brouwer degree (see e.g. [18,Appendix D] for the definition and properties of this degree). So

$$
\operatorname{deg}\left(M,(r, R)^{\ell}, \rho\right)=\operatorname{deg}\left(G,(r, R)^{\ell}, \rho\right)=(-1)^{\ell}
$$

and $M(s)=0$ must have a solution.
(iii) : If $M\left(s_{1}^{0}, \ldots, s_{\ell}^{0}\right)=0$, then $\widetilde{M}(1, \ldots, 1)=0$ where

$$
\tilde{M}_{i}(s)=\widetilde{a}_{i} s_{i}-\widetilde{b}_{1} s_{i}^{p}+\sum_{j \neq i} \widetilde{d}_{i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}}
$$

with $\widetilde{a}_{i}:=a_{i} s_{i}^{0}, \widetilde{b}_{i}:=b_{i}\left(s_{i}^{0}\right)^{p}, \widetilde{d}_{i j}:=d_{i j}\left(s_{i}^{0}\right)^{\alpha_{i j}}\left(s_{j}^{0}\right)^{\beta_{i j}}$. So we may assume without loss of generality that $M(1, \ldots, 1)=0$. Then,

$$
a_{i}-b_{i}+\sum_{j \neq i} d_{i j}=0
$$

Suppose there is another solution $s=\left(s_{1}, \ldots, s_{\ell}\right)$. Then, using the previous identity, we get

$$
0=a_{i} s_{i}-b_{i} s_{i}^{p}+\sum_{j \neq i} d_{i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}}=a_{i} s_{i}-\left(a_{i}+\sum_{j \neq i} d_{i j}\right) s_{i}^{p}+\sum_{j \neq i} d_{i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}},
$$

and after rearranging the terms,

$$
a_{i}\left(s_{i}-s_{i}^{p}\right)=\sum_{j \neq i} d_{i j}\left(s_{i}^{p}-s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}}\right)
$$

There are two possible cases: If $s_{i}>1$ for some $i$, we may assume without loss of generality that $s_{i} \geq s_{j}$ for all $j$. Then the left-hand side above is negative while the right-hand side is $\geq 0$, a contradiction. If, on the other hand, $0<s_{i}<1$ for some $i$, we may assume $s_{i} \leq s_{j}$ for all $j$. Now the left-hand side is positive and the right-hand side is $\leq 0$, a contradiction again.
(iv) : If $a_{n, i}, a_{i}, b_{n, i}, b_{i}>0, d_{n, i}, d_{i} \geq 0, a_{n, i} \rightarrow a_{i}, b_{n, i} \rightarrow b_{i}, d_{n, i j} \rightarrow d_{i j}$ then, as in (ii), there exist $0<r<R$ such that the unique solution $s_{n}$ to

$$
M_{n, i}(s):=a_{n, i} s_{i}-b_{n, i} s_{i}^{p}+\sum_{j \neq i} d_{n, i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}}=0, \quad i, j=1, \ldots, \ell
$$

belongs to $[r, R]^{\ell}$ for every $n$. Passing to a subsequence, we have that $s_{n} \rightarrow s \in[r, R]^{\ell}$ and $M(s)=0$.

## 3 A Nehari-type manifold

Let $\mathcal{H}:=H_{0}^{1}(\Omega)^{\ell}, u=\left(u_{1}, \ldots, u_{\ell}\right) \in \mathcal{H}$. As convenient norms in $H_{0}^{1}(\Omega)$ and $\mathcal{H}$ we choose

$$
\left\|u_{i}\right\|:=\left(\int_{\Omega}\left|\nabla u_{i}\right|^{2}\right)^{\frac{1}{2}} \quad \text { and } \quad\|u\|:=\left(\left\|u_{1}\right\|^{2}+\cdots+\left\|u_{\ell}\right\|^{2}\right)^{\frac{1}{2}}
$$

and we denote by $\langle\cdot, \cdot\rangle$ the inner product in $H_{0}^{1}(\Omega)$. Let

$$
I(u):=\left(I_{1}(u), \ldots, I_{\ell}(u)\right)
$$

where $I_{i}: H_{0}^{1}(\Omega) \rightarrow H_{0}^{1}(\Omega)$ are given by

$$
\begin{equation*}
I_{i}(u):=u_{i}-K_{i}(u) \tag{3.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\langle K_{i}(u), v\right\rangle:=\int_{\Omega} \mu_{i}\left(u_{i}^{+}\right)^{p} v+\sum_{j \neq i} \lambda_{i j} \int_{\Omega}\left(u_{i}^{+}\right)^{\alpha_{i j}}\left(u_{j}^{+}\right)^{\beta_{i j}} v \quad \forall v \in H_{0}^{1}(\Omega) . \tag{3.2}
\end{equation*}
$$

Lemma 3.1 If $u_{n} \rightharpoonup u$ weakly in $\mathcal{H}$, then $K_{i}\left(u_{n}\right) \rightarrow K_{i}(u)$ strongly in $H_{0}^{1}(\Omega)$ for each $i=1, \ldots, \ell$.

Proof Since $p, \alpha_{i j}+\beta_{i j}<\frac{N+2}{N-2}$ for $N \geq 3$, after passing to a subsequence $u_{n, i}^{+} \rightarrow u_{i}^{+}$ strongly in $L^{p+1}(\Omega)$ and in $L^{\alpha_{i j}+\beta_{i j}+1}(\Omega)$ for every $j \neq i$. Using Hölder's and Sobolev's inequalities we obtain

$$
\begin{aligned}
& \left|\left\langle K_{i}\left(u_{n}\right)-K_{i}(u), v\right\rangle\right| \leq C\left(\left|\left(u_{n, i}^{+}\right)^{p}-\left(u_{i}^{+}\right)^{p}\right|_{\frac{p+1}{p}}\right. \\
& \quad+\sum_{j \neq i}\left|\left(u_{n, i}^{+}\right)^{\alpha_{i j}}-\left(u_{i}^{+}\right)^{\alpha_{i j}}\right| \frac{\alpha_{i j}+\beta_{i j}+1}{\alpha_{i j}} \\
& \left.\quad \sum_{j \neq i}\left|\left(u_{n, j}^{+}\right)^{\beta_{i j}}-\left(u_{j}^{+}\right)^{\beta_{i j}}\right|_{\frac{\alpha_{i j}+\beta_{i j}+1}{\beta_{i j}}}\right)\|v\|,
\end{aligned}
$$

where $|\cdot|_{r}$ denotes the norm in $L^{r}(\Omega)$. From [18,Theorem A.2] we derive

$$
\sup _{v \neq 0} \frac{\left|\left\langle K_{i}\left(u_{n}\right)-K_{i}(u), v\right\rangle\right|}{\|v\|} \longrightarrow 0
$$

Hence, $K_{i}\left(u_{n}\right) \rightarrow K_{i}(u)$ strongly in $H_{0}^{1}(\Omega)$, as claimed.
We define a Nehari-type set $\mathcal{N}$ by putting

$$
\mathcal{N}:=\left\{u \in \mathcal{H}: u_{i} \neq 0 \text { and }\left\langle I_{i}(u), u_{i}\right\rangle=0 \text { for all } i=1, \ldots, \ell\right\} .
$$

Lemma $3.2 \mathcal{N}$ is closed in $\mathcal{H}$.
Proof Since $\lambda_{i j}<0$, it follows from the Sobolev inequality that

$$
\left\|u_{i}\right\|^{2} \leq \mu_{i} \int_{\Omega}\left(u_{i}^{+}\right)^{p+1} \leq C_{i}\left\|u_{i}\right\|^{p+1}
$$

for some $C_{i}>0$. Hence there exists $d_{0}>0$ such that, if $\left(u_{1}, \ldots, u_{\ell}\right) \in \mathcal{N}$, then $\left\|u_{i}\right\| \geq d_{0}$ for all $i$. This shows that $\mathcal{N}$ is closed in $\mathcal{H}$.

For $u:=\left(u_{1}, \ldots, u_{\ell}\right) \in \mathcal{H}, s:=\left(s_{1}, \ldots, s_{\ell}\right) \in(0, \infty)^{\ell}$ and $s u:=$ $\left(s_{1} u_{1}, \ldots, s_{\ell} u_{\ell}\right)$, we define

$$
M_{u}(s):=\left(M_{u, 1}(s), \ldots, M_{u, \ell}(s)\right),
$$

where

$$
M_{u, i}(s):=\left\langle I_{i}(s u), u_{i}\right\rangle=a_{u, i} s_{i}-b_{u, i} s_{i}^{p}+\sum_{j \neq i} d_{u, i j} s_{i}^{\alpha_{i j}} s_{j}^{\beta_{i j}}
$$

and

$$
a_{u, i}:=\left\|u_{i}\right\|^{2}, \quad b_{u, i}:=\int_{\Omega} \mu_{i}\left(u_{i}^{+}\right)^{p+1}, \quad d_{u, i j}:=\int_{\Omega}\left(-\lambda_{i j}\right)\left(u_{i}^{+}\right)^{\alpha_{i j}+1}\left(u_{j}^{+}\right)^{\beta_{i j}} .
$$

Lemma 3.3 (i) If $a_{u, i} \neq 0$ and $b_{u, i}=0$ for some $i$, then $M_{u}(s) \neq 0$ for any $s \in(0, \infty)^{\ell}$.
(ii) If $a_{u, i}, b_{u, i}>0$ for all $i$, then there exists a unique $s_{u} \in(0, \infty)^{\ell}$ such that $M_{u}\left(s_{u}\right)=0$. Moreover, if $0<a \leq a_{u, i} \leq \bar{a}, 0<b \leq b_{u, i} \leq \bar{b}$ and $d_{u, i j} \leq \bar{d}$ for all $i, j$, then there exist $0<r<R$, depending only on $a, \bar{a}, b, \bar{b}, \bar{d}$, such that $s_{u} \in(r, R)^{\ell}$.

Proof This is an immediate consequence of Lemma 2.1.
Let

$$
\mathscr{S}:=\left\{v \in H_{0}^{1}(\Omega):\|v\|=1\right\}, \quad \mathcal{T}:=\mathscr{S}^{\ell}
$$

and

$$
\begin{align*}
\mathcal{U}: & =\left\{u \in \mathcal{T}: s_{u} \in(0, \infty)^{\ell} \text { exists with } M_{u}\left(s_{u}\right)=0\right\} \\
& =\left\{u \in \mathcal{T}: u_{i}^{+} \neq 0 \text { for all } i=1, \ldots, \ell\right\} . \tag{3.3}
\end{align*}
$$

The tangent space of $\mathcal{T}$ at $u$ is

$$
\begin{equation*}
T_{u}(\mathcal{T}):=\left\{\left(v_{1}, \ldots, v_{\ell}\right) \in \mathcal{H}:\left\langle u_{i}, v_{i}\right\rangle=0 \text { for all } i=1, \ldots, \ell\right\} \tag{3.4}
\end{equation*}
$$

Proposition 3.4 (i) $\mathcal{U}$ is a nonempty open subset of $\mathcal{T}$ and $\mathcal{U} \neq \mathcal{T}$.
(ii) The mapping $m: \mathcal{U} \rightarrow \mathcal{N}$ given by $m(u):=s_{u} u$ is a homeomorphism. In particular, $\mathcal{N}$ is a topological manifold.
(iii) If $\left(u_{n}\right)$ is a sequence in $\mathcal{U}$ such that $u_{n} \rightarrow u \in \partial \mathcal{U}$, then $s_{u_{n}} \rightarrow \infty$ (and hence $\left.\left\|m\left(u_{n}\right)\right\| \rightarrow \infty\right)$.
(iv) Let $S: \mathcal{U} \rightarrow \mathcal{H}$ be given by

$$
S(u):=I\left(s_{u} u\right)=s_{u} u-K\left(s_{u} u\right) .
$$

Then $S(u) \in T_{u}(\mathcal{U})$ for every $u \in \mathcal{U}$.
(v) $S(u)=0$ if and only if $m(u)=s_{u} u$ is a solution for (1.1).

Proof $(i)$ : That $\mathcal{U}$ is neither empty nor the whole $\mathcal{T}$ is obvious and, since $u \mapsto u_{i}^{+}$ is continuous [2,Lemma 2.3], it is easily seen from the second line of (3.3) that $\mathcal{U}$ is open in $\mathcal{T}$.
(ii) : If $u \in \mathcal{U}$, then $s_{u} u \in \mathcal{N}$ because $\left\langle I_{i}\left(s_{u} u\right), s_{u, i} u_{i}\right\rangle=s_{u, i} M_{u, i}\left(s_{u}\right)=0$ for all $i$. So $m$ is well defined. If $\left(u_{n}\right)$ is a sequence in $\mathcal{U}$ and $u_{n} \rightarrow u \in \mathcal{U}$, then $a_{u_{n}, i} \rightarrow a_{u, i}$, $b_{u_{n}, i} \rightarrow b_{u, i}$ and $d_{u_{n}, i j} \rightarrow d_{u, i j}$ for all $i, j$. By Lemma 2.1(iv), $s_{u_{n}} \rightarrow s_{u}$. Hence, $m$ is continuous.

If $u \in \mathcal{N}$, then $u_{i}^{+} \neq 0$ for all $i$. Otherwise, $0=\left\langle I_{i}(u), u_{i}\right\rangle=\left\|u_{i}\right\|^{2}$, a contradiction. Hence, the inverse of $m$ satisfies

$$
m^{-1}(u):=\left(\frac{u_{1}}{\left\|u_{1}\right\|}, \ldots, \frac{u_{\ell}}{\left\|u_{\ell}\right\|}\right) \in \mathcal{U}
$$

and it is obviously continuous.
(iii) : Let $\left(u_{n}\right)$ be a sequence in $\mathcal{U}$ such that $u_{n} \rightarrow u \in \partial \mathcal{U}$. If $\left(s_{u_{n}}\right)$ is bounded, then, after passing to a subsequence, $s_{u_{n}} \rightarrow s_{*}$. Since $\mathcal{N}$ is closed, $s_{*} u \in \mathcal{N}$ and hence $u \in \mathcal{U}$. This is impossible because $\mathcal{U}$ is open.
(iv) : Since $\left\langle I_{i}\left(s_{u} u\right), u_{i}\right\rangle=M_{u, i}\left(s_{u}\right)=0$ for all $i$, we have that $S(u) \in T_{u}(\mathcal{T})$ according to (3.4).
(v) : If $u \in \mathcal{U}$ satisfies $S(u)=0$, then $\bar{u}:=s_{u} u \in \mathcal{N}$ and $\bar{u}$ is a weak solution to the system (1.2) (see (3.1) and (3.2)). Multiplying the $i$-th equation in (1.2) by $u_{i}^{-}:=\min \left\{\bar{u}_{i}, 0\right\}$ and integrating gives $\int_{\Omega}\left|\nabla u_{i}^{-}\right|^{2}=0$. Hence $u_{i}^{-}=0$, i.e., $\bar{u}_{i} \geq 0$ for all $i$. As $\bar{u} \in \mathcal{N}$, we have that $\bar{u}_{i} \neq 0$. This proves that $\bar{u}$ solves (1.1). The converse is obvious.

Remark 3.5 If $\alpha_{i j} \geq 1$ for all $i$ and all $j \neq i$, then, as $\bar{u}_{i}$ above satisfies the $i$-th equation in (1.1), we have

$$
-\Delta \bar{u}_{i}+c(x) \bar{u}_{i} \geq 0 \quad \text { where } c(x):=-\sum_{j \neq i} \lambda_{i j} \bar{u}_{i}^{\alpha_{i j}-1} \bar{u}_{j}^{\beta_{i j}} .
$$

Since all $u_{i}$ are continuous in $\bar{\Omega}$ and $c \geq 0$, it follows from the strong maximum principle (see e.g. [14,Theorem 3.5]) that our solution is strictly positive in $\Omega$ in this case.

## 4 Proof of Theorem 1.1

In this section the sub- or superscript $t$ will be used in order to emphasize that we are concerned with the system (1.3). So, e.g.,

$$
\begin{equation*}
I_{t}(u):=u-K_{t}(u), \quad S_{t}(u):=I_{t}\left(s_{u}^{t} u\right)=s_{u}^{t} u-K_{t}\left(s_{u}^{t} u\right), \tag{4.1}
\end{equation*}
$$

with

$$
\left\langle K_{t, i}(u), v\right\rangle:=\int_{\Omega} \mu_{i}\left(u_{i}^{+}\right)^{p} v+t \sum_{j \neq i} \lambda_{i j} \int_{\Omega}\left(u_{i}^{+}\right)^{\alpha_{i j}}\left(u_{j}^{+}\right)^{\beta_{i j}} v,
$$

and

$$
\mathcal{N}_{t}:=\left\{u \in \mathcal{H}: u_{i} \neq 0,\left\langle I_{t, i}(u), u_{i}\right\rangle=0 \text { for all } i=1, \ldots, \ell\right\}
$$

According to this notation, $\mathcal{N}_{1}=\mathcal{N}$. When $t=1$, we shall sometimes omit the subor superscript $t$.

Consider first the system (1.3) with $t=0$. In this case the equations are uncoupled, the set

$$
\mathcal{N}_{0}=\left\{u \in \mathcal{H}: u_{i} \neq 0,\left\|u_{i}\right\|^{2}=\int_{\Omega} \mu_{i}\left(u_{i}^{+}\right)^{p+1} \text { for all } i=1, \ldots, \ell\right\}
$$

is the product of the usual Nehari manifolds associated to these equations, and the components of $s_{u}^{0}=\left(s_{u, 1}^{0} \ldots, s_{u, \ell}^{0}\right)$ are

$$
s_{u, i}^{0}=\left(\int_{\Omega} \mu_{i}\left(u_{i}^{+}\right)^{p+1}\right)^{-\frac{1}{p-1}}, \quad u \in \mathcal{U}
$$

The function $I_{0}$ (cf. (4.1)) is the gradient vector field of the functional $\mathcal{J}: \mathcal{H} \rightarrow \mathbb{R}$ defined by

$$
\mathcal{J}(u):=\frac{1}{2} \sum_{i=1}^{\ell}\left\|u_{i}\right\|^{2}-\frac{1}{p+1} \sum_{i=1}^{\ell} \int_{\Omega} \mu_{i}\left(u_{i}^{+}\right)^{p+1} .
$$

Note that

$$
\mathcal{J}(u)=\left(\frac{1}{2}-\frac{1}{p+1}\right)\|u\|^{2} \quad \text { if } u \in \mathcal{N}_{0} .
$$

$\mathcal{J}$ has a minimizer $\tilde{u}_{0}=\left(\widetilde{u}_{0,1}, \ldots, \widetilde{u}_{0, \ell}\right)$ on $\mathcal{N}_{0}$ with $\widetilde{u}_{0, i}>0$ and $\widetilde{u}_{0}$ is a solution to the system (1.3) with $t=0$. Each component $\widetilde{u}_{0, i}$ is a positive least energy solution to the $i$-th equation of this system. Let $\Psi: \mathcal{U} \rightarrow \mathbb{R}$ be given by

$$
\begin{align*}
\Psi(u): & =\mathcal{J}\left(s_{u}^{0} u\right)=\left(\frac{1}{2}-\frac{1}{p+1}\right)\left|s_{u}^{0}\right|^{2} \\
& =\left(\frac{1}{2}-\frac{1}{p+1}\right) \sum_{i=1}^{\ell}\left(\int_{\Omega} \mu_{i}\left(u_{i}^{+}\right)^{p+1}\right)^{-\frac{2}{p-1}} . \tag{4.2}
\end{align*}
$$

By [18,Proposition 1.12] one has that $\Psi \in \mathcal{C}^{2}(\mathcal{U}, \mathbb{R})$. It is easily seen that

$$
\begin{equation*}
\Psi^{\prime}(u) v=\mathcal{J}^{\prime}\left(s_{u}^{0} u\right)\left[s_{u}^{0} v\right]=\left\langle S_{0}(u), s_{u}^{0} v\right\rangle \text { for all } u \in \mathcal{U}, v \in T_{u}(\mathcal{U}) \tag{4.3}
\end{equation*}
$$

and that $u$ is a critical point of $\Psi$ if and only if $u \in \mathcal{U}$ and $m_{0}(u)=s_{u}^{0} u$ is a critical point of $\mathcal{J}$, see [4,Theorem 3.3]. Let $u_{0}:=m_{0}^{-1}\left(\widetilde{u}_{0}\right)$. Then $u_{0}$ is a minimizer for $\Psi$.

Invoking Lemma 1.2 we may choose $R>0$ such that all solutions to the systems (1.3) are contained in the open ball $B_{R}(0) \subset \mathcal{H}$, where $R$ is independent of $t \in[0,1]$. Then, by Proposition 3.4,

$$
\begin{equation*}
\left\{u \in \mathcal{U}: S_{t}(u)=0\right\} \subset \mathcal{V}_{t}:=m_{t}^{-1}\left(B_{R}(0) \cap \mathcal{N}_{t}\right) \tag{4.4}
\end{equation*}
$$

For $a \leq d$ let

$$
\Psi^{d}:=\{u \in \mathcal{U}: \Psi(u) \leq d\}, \quad \Psi_{a}^{d}:=\{u \in \mathcal{U}: a \leq \Psi(u) \leq d\} .
$$

It follows from Proposition 3.4(iii) that the set $\Psi^{d}$ is closed in $\mathcal{T}$ for any $d \in \mathbb{R}$. Note that $\lambda_{i j}<0$ implies $s_{u, i}^{t} \geq s_{u, i}^{0}$ for every $u \in \mathcal{U}, t \in[0,1], i=1, \ldots \ell$. So if $\left|s_{u}^{t}\right|<R$, then $\left|s_{u}^{0}\right|<R$; hence $\mathcal{V}_{t} \subset \mathcal{V}_{0}$ and, setting $c:=\left(\frac{1}{2}-\frac{1}{p+1}\right) R^{2}$, we have that the closure of $\mathcal{V}_{t}$ in $\mathcal{T}$ satisfies

$$
\begin{equation*}
\overline{\mathcal{V}}_{t} \subset \overline{\mathcal{V}}_{0} \subset \Psi^{c} \quad \forall t \in[0,1] \tag{4.5}
\end{equation*}
$$

For each $i=1, \ldots, \ell$ and $k \geq 2$ we choose an ascending sequence ( $E_{k, i}$ ) of linear subspaces of $H_{0}^{1}(\Omega)$ such that $\operatorname{dim} E_{k, i}=k, u_{0, i} \in E_{2, i}$ ( $u_{0}$ is the minimizer chosen above) and $\overline{\bigcup_{k \geq 1} E_{k, i}}=H_{0}^{1}(\Omega)$. We define

$$
\begin{equation*}
E_{k}:=E_{k, 1} \times \cdots \times E_{k, \ell} \subset \mathcal{H} \tag{4.6}
\end{equation*}
$$

and denote by $P_{k}$ the orthogonal projector of $\mathcal{H}$ onto $E_{k}$.
Lemma 4.1 Given $d>0$ there exists $k_{d} \in \mathbb{N}$ such that

$$
P_{k}\left(S_{t}(u)\right) \neq 0 \quad \text { for all } u \in\left(\Psi^{d} \backslash \mathcal{V}_{t}\right) \cap E_{k}, k \geq k_{d}, t \in[0,1]
$$

Proof Arguing by contradiction, assume that there exist $k_{n} \rightarrow \infty, t_{n} \in[0,1]$ and $u_{n} \in\left(\Psi^{d} \backslash \mathcal{V}_{t_{n}}\right) \cap E_{k_{n}}$ such that

$$
\begin{equation*}
P_{k_{n}}\left(S_{t_{n}}\left(u_{n}\right)\right)=s_{u_{n}}^{t_{n}} u_{n}-P_{k_{n}} K_{t_{n}}\left(s_{u_{n}}^{t_{n}} u_{n}\right)=0 \quad \forall n \in \mathbb{N} . \tag{4.7}
\end{equation*}
$$

As $u_{n} \in \Psi^{d}$, we derive from (4.2) that $\int_{\Omega} \mu_{i}\left(u_{n, i}^{+}\right)^{p+1} \geq b$ for some $b>0$ and all $n, i$. In the notation of Lemma 3.3, we have $a_{u_{n}, i}=1$ and, using the Hölder and the Sobolev inequalities, $b \leq b_{u_{n}, i} \leq \bar{b}$ and

$$
d_{u_{n}, i j}=t_{n} \int_{\Omega}\left(-\lambda_{i j}\right)\left(u_{n, i}^{+}\right)^{\alpha_{i j}+1}\left(u_{n, j}^{+}\right)^{\beta_{i j}} \leq \bar{d}
$$

for some $\bar{b}, \bar{d}>0$. So Lemma 3.3 asserts that $\left(s_{u_{n}, i}^{t_{n}}\right)$ is bounded and bounded away from 0 for each $i$. Therefore, after passing to a subsequence, $s_{u_{n}, i}^{t_{n}} \rightarrow s_{i}>0, t_{n} \rightarrow t$
and $u_{n} \rightharpoonup u$ weakly in $\mathcal{H}$. By Lemma 3.1, $K_{t_{n}}\left(s_{u_{n}}^{t_{n}} u_{n}\right) \rightarrow K_{t}(s u)$ strongly in $\mathcal{H}$, and we easily deduce that $P_{k_{n}} K_{t_{n}}\left(s_{u_{n}}^{t_{n}} u_{n}\right) \rightarrow K_{t}(s u)$ strongly in $\mathcal{H}$. Now we derive from (4.7) that $s_{u_{n}}^{t_{n}} u_{n} \rightarrow s u$ strongly in $\mathcal{H}$ and $s u-K_{t}(s u)=0$. Therefore, $s u \in \mathcal{N}_{t}, s=s_{u}^{t}$ and $S_{t}(u)=0$. On the other hand, as $u_{n} \notin \mathcal{V}_{t_{n}}$, we have that $\left\|s_{u_{n}}^{t_{n}} u_{n}\right\| \geq R$. Hence, $\left\|s_{u}^{t} u\right\| \geq R$. This is a contradiction.

Lemma 4.2 Let $c$ be as in (4.5). Then $\Psi^{c} \cap E_{k}$ is contractible in itself for each large enough $k$.

Proof Let $\eta:[0,1] \times \mathcal{U} \rightarrow \mathcal{U}$ be given by

$$
\eta(\tau, u):=\left(\frac{(1-\tau) u_{1}+\tau u_{0,1}}{\left\|(1-\tau) u_{1}+\tau u_{0,1}\right\|}, \ldots, \frac{(1-\tau) u_{\ell}+\tau u_{0, \ell}}{\left\|(1-\tau) u_{\ell}+\tau u_{0, \ell}\right\|}\right)
$$

where $u_{0}$ is the previously chosen minimizer for $\Psi$ on $\mathcal{U}$. Note that $\eta$ is well defined and maps into $\mathcal{U}$ because $u_{0, i}>0$ in $\Omega$ and $u_{i}^{+} \neq 0$ for all $i$. Moreover, if $u \in E_{k}$, then $\eta(\tau, u) \in E_{k}$ for each $k \geq 2$. So $\eta$ is a deformation of $\mathcal{U} \cap E_{k}$ into $u_{0}$ and, in particular, of $\Psi^{c} \cap E_{k}$ into $u_{0}$ in $\mathcal{U} \cap E_{k}$.

We claim that there exists $\delta_{0}>0$ such that

$$
\int_{\Omega}\left[\left((1-\tau) u_{i}+\tau u_{0, i}\right)^{+}\right]^{p+1} \geq \delta_{0} \text { for all } \tau \in[0,1], u \in \Psi^{c}, i=1, \ldots, \ell
$$

Otherwise, there would exist $\tau_{n} \in[0,1]$ and $u_{n} \in \Psi^{c}$ such that

$$
\begin{equation*}
\left(1-\tau_{n}\right) \int_{\Omega}\left(u_{n, i}^{+}\right)^{p+1} \leq \int_{\Omega}\left[\left(\left(1-\tau_{n}\right) u_{n, i}+\tau_{n} u_{0, i}\right)^{+}\right]^{p+1} \rightarrow 0 \tag{4.8}
\end{equation*}
$$

(the inequality is satisfied because $u_{0, i}>0$ ). From (4.2) we see that there exists $\delta>0$ such that $\int_{\Omega}\left(u_{i}^{+}\right)^{p+1} \geq \delta$ for all $u \in \Psi^{c}$ and all $i$. Hence, $\tau_{n} \rightarrow 1$. Since $\left(u_{n}\right)$ is bounded in $\mathcal{H}$, a subsequence of ( $u_{n, i}$ ) converges in $L^{p+1}(\Omega)$. Therefore,

$$
\int_{\Omega}\left[\left(\left(1-\tau_{n}\right) u_{n, i}+\tau_{n} u_{0, i}\right)^{+}\right]^{p+1} \rightarrow \int_{\Omega} u_{0, i}^{p+1} \geq \delta,
$$

a contradiction to (4.8).
So, for every $\tau \in[0,1], u \in \Psi^{c}, i=1, \ldots, \ell$, we have

$$
\begin{aligned}
\int_{\Omega}\left(\eta_{i}(\tau, u)^{+}\right)^{p+1} & =\int_{\Omega} \frac{\left[\left((1-\tau) u_{i}+\tau u_{0, i}\right)^{+}\right]^{p+1}}{\left\|(1-\tau) u_{i}+\tau u_{0, i}\right\|^{p+1}} \\
& \geq \int_{\Omega}\left[\left((1-\tau) u_{i}+\tau u_{0, i}\right)^{+}\right]^{p+1} \geq \delta_{0}
\end{aligned}
$$

and we deduce from (4.2) that there exists $d>c$ such that

$$
\eta(\tau, u) \in \Psi^{d} \cap E_{k} \quad \text { for all } \tau \in[0,1], u \in \Psi^{c} \cap E_{k}, k \geq 2
$$

Next we show that $\left.\Psi\right|_{\mathcal{U} \cap E_{k}}$ does not have a critical value in $[c, d]$ for any large enough $k$. Indeed, if $u_{k} \in \Psi_{c}^{d}$ is a critical point of $\left.\Psi\right|_{\mathcal{U} \cap E_{k}}$, then, according to (4.3),

$$
\left\langle S_{0}\left(u_{k}\right), s_{u_{k}}^{0} v\right\rangle=0 \quad \text { for all } v \in T_{u_{k}}\left(\mathcal{U} \cap E_{k}\right)
$$

i.e., $P_{k} S_{0}\left(u_{k}\right)=0$. Since $u_{k} \in \Psi_{c}^{d} \subset \Psi^{d} \backslash \mathcal{V}_{t}$ (see (4.5)), $k<k_{d}$ according to Lemma 4.1.

Now Proposition 3.4(iii) allows us to use the negative gradient flow of $\left.\Psi\right|_{\mathcal{U} \cap E_{k}}$ in the standard way to obtain a retraction $\varrho: \Psi^{d} \cap E_{k} \rightarrow \Psi^{c} \cap E_{k}$; see, e.g., [3,Theorem I.3.2]. Then, $\varrho \circ \eta:[0,1] \times\left(\Psi^{c} \cap E_{k}\right) \rightarrow \Psi^{c} \cap E_{k}$ is a deformation of $\Psi^{c} \cap E_{k}$ into a point.

The following statement is an immediate consequence of Lemma 4.2 and basic properties of homology (see e.g. [12,Sections III. 4 and III.5]).

Corollary 4.3 Denote the $q$-th singular homology with coefficients in a field $\mathbb{F}$ by $\mathrm{H}_{q}(\cdot)$. Then $\mathrm{H}_{0}\left(\Psi^{c} \cap E_{k}\right)=\mathbb{F}$ and $\mathrm{H}_{q}\left(\Psi^{c} \cap E_{k}\right)=0$ for $q \neq 0$. In particular, the Euler characteristic

$$
\chi\left(\Psi^{c} \cap E_{k}\right):=\sum_{q \geq 0}(-1)^{q} \operatorname{dim}_{\mathbb{F}} \mathrm{H}_{q}\left(\Psi^{c} \cap E_{k}\right)=1
$$

for every large enough $k$.
For $u_{0}$ as above, let

$$
\sigma_{i}: \mathscr{S} \backslash\left\{-u_{0, i}\right\} \rightarrow\left(\mathbb{R} u_{0, i}\right)^{\perp}=: F_{i}
$$

be the stereographic projection. The product $\sigma=\left(\sigma_{1}, \ldots, \sigma_{\ell}\right)$ of the stereographic projections is a diffeomorphism. So its derivative at $u$

$$
\sigma^{\prime}(u): T_{u}(\mathcal{U}) \rightarrow F:=F_{1} \times \ldots \times F_{\ell}
$$

is an isomorphism for every $u \in \mathcal{U}$. Note that, as $u_{0, i} \in E_{2, i}$, we have that $\sigma_{i}((\mathscr{S} \cap$ $\left.\left.E_{k}\right) \backslash\left\{-u_{0, i}\right\}\right) \subset F_{i} \cap E_{k}$ for all $k \geq 2$.

Proof of Theorem 1.1 Let $\mathcal{O}:=\sigma\left(\mathcal{V}_{0}\right)$ with $\mathcal{V}_{0}$ as in (4.4). As $u_{0} \in \mathcal{V}_{0}$ we have that $0 \in \mathcal{O}$, and as $\overline{\mathcal{V}}_{0} \subset \mathcal{U}$ and $-u_{0} \notin \mathcal{U}, \mathcal{O}$ is bounded in $F$. Set $\mathcal{O}_{k}:=\mathcal{O} \cap E_{k}$ and $F_{k}:=$ $F \cap E_{k}$. Then $\mathcal{O}_{k}$ is a bounded open neighborhood of 0 in $F_{k}$, and $\overline{\mathcal{O}}_{k} \subset \sigma\left(\Psi^{c} \cap E_{k}\right)$ for $c$ as in (4.5).

Fix $k_{c} \in \mathbb{N}$ as in Lemma 4.1. Recall that

$$
S_{t}(u)=s_{u}^{t} u-K_{t}\left(s_{u}^{t} u\right) \in T_{u}(\mathcal{U}) \quad \forall u \in \mathcal{U}
$$

(see Proposition 3.4(iv)). Define $G_{t, k}: \sigma\left(\mathcal{U} \cap E_{k}\right) \rightarrow F_{k}$ by

$$
\begin{equation*}
G_{t, k}(w):=\left(\sigma^{\prime}\left(\sigma^{-1}(w)\right) \circ P_{k} \circ S_{t} \circ \sigma^{-1}\right)(w) \tag{4.9}
\end{equation*}
$$

Note that

$$
G_{t, k}(w)=0 \Longleftrightarrow P_{k}\left(S_{t}\left(\sigma^{-1}(w)\right)\right)=0
$$

So, if $k \geq k_{c}, w \in \sigma\left(\mathcal{U} \cap E_{k}\right)$ and $G_{t, k}(w)=0$, Lemma 4.1 asserts that $w \in \mathcal{O}_{k}$. In particular, $G_{t, k}(w) \neq 0$ for every $w \in \partial \mathcal{O}_{k}$. From the homotopy and the excision properties of the Brouwer degree we get that

$$
\begin{equation*}
\operatorname{deg}\left(G_{1, k}, \mathcal{O}_{k}, 0\right)=\operatorname{deg}\left(G_{0, k}, \mathcal{O}_{k}, 0\right)=\operatorname{deg}\left(G_{0, k}, \sigma\left(\Psi^{c} \cap E_{k}\right), 0\right) \tag{4.10}
\end{equation*}
$$

On the other hand, using (4.3) and (4.9) we get

$$
\begin{aligned}
\left(\Psi \circ \sigma^{-1}\right)^{\prime}(w) z & =\Psi^{\prime}\left(\sigma^{-1}(w)\right)\left[\left(\sigma^{-1}\right)^{\prime}(w) z\right] \\
& =s_{\sigma^{-1}(w)}^{0}\left\langle P_{k}\left(S_{0}\left(\sigma^{-1}(w)\right)\right),\left(\sigma^{-1}\right)^{\prime}(w) z\right\rangle \\
& =s_{\sigma^{-1}(w)}^{0}\left\langle\left(\sigma^{-1}\right)^{\prime}(w)\left(G_{0, k}(w)\right),\left(\sigma^{-1}\right)^{\prime}(w) z\right\rangle \\
& =\frac{4 s_{\sigma^{-1}(w)}^{0}}{\left(\|w\|^{2}+1\right)^{2}}\left\langle G_{0, k}(w), z\right\rangle \quad \forall w \in \sigma\left(\mathcal{U} \cap E_{k}\right), z \in F_{k} .
\end{aligned}
$$

The last identity is obtained by a simple calculation, see e.g. [15,Lemma 3.4]. Since $\left.\left(\Psi \circ \sigma^{-1}\right)\right|_{\sigma\left(\mathcal{U} \cap E_{k}\right)}$ is of class $\mathcal{C}^{2}($ see $(4.2))$ and $-\left(\Psi \circ \sigma^{-1}\right)^{\prime}(w)$ points into $\left(\Psi \circ \sigma^{-1}\right)^{c}$ for all $w \in\left(\Psi \circ \sigma^{-1}\right)_{c}^{c}$, from [3,Theorem II.3.3] and Corollary 4.3 we obtain

$$
\begin{align*}
\operatorname{deg}\left(G_{0, k}, \sigma\left(\Psi^{c} \cap E_{k}\right), 0\right) & =\operatorname{deg}\left(\left(\left.\left(\Psi \circ \sigma^{-1}\right)\right|_{\sigma\left(\mathcal{U} \cap E_{k}\right)}\right)^{\prime}, \sigma\left(\Psi^{c} \cap E_{k}\right), 0\right) \\
& =\chi\left(\sigma\left(\Psi^{c} \cap E_{k}\right)\right)=\chi\left(\Psi^{c} \cap E_{k}\right)=1 . \tag{4.11}
\end{align*}
$$

Combining (4.10) and (4.11) gives

$$
\operatorname{deg}\left(G_{1, k}, \mathcal{O}_{k}, 0\right)=1
$$

Hence, for each $k \geq k_{c}$ there exists $w_{k} \in \mathcal{O}_{k}$ such that $G_{k, 1}\left(w_{k}\right)=0$. Then $u_{k}:=$ $\sigma^{-1}\left(w_{k}\right) \in \mathcal{V}_{0} \cap E_{k} \subset \Psi^{c} \cap E_{k}$ satisfies $P_{k}\left(S\left(u_{k}\right)\right)=0$, i.e.,

$$
\begin{equation*}
s_{u_{k}} u_{k}=P_{k} K\left(s_{u_{k}} u_{k}\right) \tag{4.12}
\end{equation*}
$$

As in the proof of Lemma 4.1 (with $t_{n}$ replaced by 1 and $s_{u_{n}}^{t_{n}} u_{n}$ by $s_{u_{k}} u_{k}$ ) one shows that $\left(s_{u_{k}, i}\right)$ is bounded and bounded away from 0 for each $i$. So passing to a subsequence, $s_{u_{k}} \rightarrow s$ and $u_{k} \rightharpoonup u$ weakly in $\mathcal{H}$. Taking limits in (4.12) and using Lemma 3.1, we obtain that $s_{u_{k}} u_{k} \rightarrow s u$ strongly in $\mathcal{H}$ and $s u=K(s u)$. Hence, $s u \in \mathcal{N}, s=s_{u}$ and $S(u)=s_{u} u-K\left(s_{u} u\right)=0$. So, according to Proposition 3.4(v), $s_{u} u$ is a solution to (1.1).

## 5 Synchronized solutions

A solution $u=\left(u_{1}, \ldots, u_{\ell}\right)$ to (1.1) is called synchronized if $u_{i}=t_{i} v$ and $u_{j}=t_{j} v$ for some $i \neq j, v \in H_{0}^{1}(\Omega) \backslash\{0\}$ and $t_{1}, t_{2}>0$. In this section we consider a system of 2 equations:

$$
\left\{\begin{array}{l}
-\Delta u_{1}=\mu_{1} u_{1}^{p}+\lambda_{12} u_{12}^{\alpha_{12}} u_{2}^{\beta_{12}}  \tag{5.1}\\
-\Delta u_{2}=\mu_{2} u_{2}^{p}+\lambda_{21} u_{2}^{\alpha_{21}} u_{1}^{\beta_{21}} \\
u_{1}, u_{2} \geq 0 \text { in } \Omega, \quad u_{1}, u_{2} \in H_{0}^{1}(\Omega) \backslash\{0\} .
\end{array}\right.
$$

Recall that according to our assumptions $\alpha_{12}+\beta_{12}<p$ and $\alpha_{21}+\beta_{21}<p$.
Theorem 5.1 The system (5.1) has a synchronized solution if and only if $\alpha_{12}+\beta_{12}=$ $\alpha_{21}+\beta_{21}=: q$ and

$$
\begin{equation*}
\frac{\lambda_{12}}{\lambda_{21}}=\left(\frac{\mu_{1}}{\mu_{2}}\right)^{\left(\alpha_{21}-\beta_{12}-1\right) /(p-1)} \tag{5.2}
\end{equation*}
$$

Proof Inserting $u_{1}=t_{1} v, u_{2}=t_{2} v$ into (5.1) we obtain

$$
\left\{\begin{array}{l}
-t_{1} \Delta v=\mu_{1} t_{1}^{p} v^{p}+\lambda_{12} t_{1}^{\alpha_{12}} t_{2}^{\beta_{12}} v^{\alpha_{12}+\beta_{12}} \\
-t_{2} \Delta v=\mu_{2} t_{2}^{p} v^{p}+\lambda_{21} t_{2}^{\alpha_{21}} t_{1}^{\beta_{21}} v^{\alpha_{21}+\beta_{21}}
\end{array}\right.
$$

Dividing the first equation by $t_{1}$, the second one by $t_{2}$ and subtracting gives

$$
\left(\mu_{1} t_{1}^{p-1}-\mu_{2} t_{2}^{p-1}\right) v^{p}+\left(\lambda_{12} t_{1}^{\alpha_{12}-1} t_{2}^{\beta_{12}} v^{\alpha_{12}+\beta_{12}}-\lambda_{21} t_{2}^{\alpha_{21}-1} t_{1}^{\beta_{21}} v^{\alpha_{21}+\beta_{21}}\right)=0
$$

So $\alpha_{12}+\beta_{12}=\alpha_{21}+\beta_{21}=q$,

$$
\mu_{1} t_{1}^{p-1}-\mu_{2} t_{2}^{p-1}=0 \quad \text { and } \quad \lambda_{12} t_{1}^{\alpha_{12}-1} t_{2}^{\beta_{12}}=\lambda_{21} t_{2}^{\alpha_{21}-1} t_{1}^{\beta_{21}}
$$

Inserting the solution

$$
t_{2}=\left(\frac{\mu_{1}}{\mu_{2}}\right)^{1 /(p-1)} t_{1}
$$

of the first equation into the second one gives (5.2).
We have shown that the conditions in Theorem 5.1 are necessary. It remains to show that they are also sufficient. To this aim observe that, if $\alpha_{12}+\beta_{12}=\alpha_{21}+\beta_{21}=: q$, (5.2) holds true, and $w$ satisfies

$$
\begin{equation*}
-\Delta w=\mu_{1} w^{p}-a w^{q}, \quad w \geq 0, w \in H_{0}^{1}(\Omega) \backslash\{0\} \tag{5.3}
\end{equation*}
$$

with

$$
a:=-\lambda_{12}\left(\frac{\mu_{1}}{\mu_{2}}\right)^{\beta_{12} /(p-1)},
$$

then $\left(w,\left(\frac{\mu_{1}}{\mu_{2}}\right)^{1 /(p-1)} w\right)$ solves the system (5.1). Consider the functional

$$
\Phi(w):=\frac{1}{2} \int_{\Omega}|\nabla w|^{2}+\frac{a}{q+1} \int_{\Omega}\left(w^{+}\right)^{q+1}-\frac{\mu_{1}}{p+1} \int_{\Omega}\left(w^{+}\right)^{p+1} .
$$

By standard arguments (see e.g. [17] or [18]), $\Phi$ is of class $\mathcal{C}^{1}$ and critical points of $\Phi$ are solutions to the equation

$$
\begin{equation*}
-\Delta w+a\left(w^{+}\right)^{q}=\mu_{1}\left(w^{+}\right)^{p} . \tag{5.4}
\end{equation*}
$$

We shall complete the proof by showing that $\Phi$ has a nontrivial critical point $w \geq 0$. We use the mountain pass theorem (see e.g. [17] or [18]). By easy calculations (as e.g. in [18,Proof of Theorem 1.19]), $\Phi$ has the mountain pass geometry. Here it is important that $p>2$ and $p>q$. Next we show that $\Phi$ satisfies the Palais-Smale condition. Let $\left(w_{n}\right)$ be such that $\Phi\left(w_{n}\right) \rightarrow c$ and $\Phi^{\prime}\left(w_{n}\right) \rightarrow 0$. Then

$$
\begin{aligned}
c+1+\left\|w_{n}\right\| & \geq \Phi\left(w_{n}\right)-\frac{1}{p+1} \Phi^{\prime}\left(w_{n}\right) w_{n} \\
& =\left(\frac{1}{2}-\frac{1}{p+1}\right) \int_{\Omega}\left|\nabla w_{n}\right|^{2}+a\left(\frac{1}{q+1}-\frac{1}{p+1}\right) \int_{\Omega}\left(w_{n}^{+}\right)^{q+1}
\end{aligned}
$$

for all $n$ large enough. Hence $\left(w_{n}\right)$ is bounded, so passing to a subsequence, $w_{n} \rightarrow w$ weakly in $H_{0}^{1}(\Omega)$, and strongly in $L^{p}(\Omega)$ and $L^{q}(\Omega)$. It follows by a standard argument (see e.g. [18,Proof of Lemma 1.20]) that $w_{n} \rightarrow w$ strongly also in $H_{0}^{1}(\Omega)$. Finally, multiplying (5.4) by $w^{-}$gives $\int_{\Omega}\left|\nabla w^{-}\right|^{2}=0$, so $w^{-}=0$. The proof is complete.

Remark 5.2 It is easy to show that if $q=p$, then there are no synchronized solutions for $-\lambda_{i j}$ sufficiently large, as is well known in the variational case, see e.g. [4,Proposition 3.2].

Remark 5.3 Let $\lambda_{i j, n}<0, i \neq j$, and let $u_{n}=\left(u_{n, 1}, \ldots, u_{n, \ell}\right)$ be a solution to (1.1) with $\lambda_{i j}$ replaced by $\lambda_{i j, n}$. It is easy to see that, if the sequence $\left(u_{n}\right)$ is bounded in $\mathcal{H}$, the components $u_{n, i}$ separate spatially as $\lambda_{i j, n} \rightarrow-\infty$. More precisely, after passing to a subsequence, $u_{n, i} \rightarrow u_{i} \neq 0$ weakly in $H_{0}^{1}(\Omega)$ and strongly in $L^{p}(\Omega)$ for each $i$, and $u_{i}(x) \cdot u_{j}(x)=0$ a.e. in $\Omega$ for $i \neq j$. There is an extensive literature on spatial separation of solutions and limiting profiles, under the assumption that the sequence $\left(u_{n}\right)$ is bounded and under different assumptions on the nonlinearities. See e.g. [6, 7, $16]$ and the references therein.

Obviously, synchronized solutions to (1.1) do not separate spatially. So we cannot expect the sequence $\left(w_{n}\right)$ given by (5.3) to be bounded. Indeed, we have the following

Proposition 5.4 Let $\left(w_{n}\right)$ be a sequence of solutions to (5.3) with $a=a_{n}$. If $a_{n} \rightarrow \infty$, then $\left(w_{n}\right)$ is unbounded in $H_{0}^{1}(\Omega)$.

Proof Suppose ( $w_{n}$ ) is bounded. Then, passing to a subsequence, $w_{n} \rightarrow w$ weakly in $H_{0}^{1}(\Omega)$, strongly in $L^{p}(\Omega)$ and in $L^{q}(\Omega)$. Since

$$
\int_{\Omega}\left|\nabla w_{n}\right|^{2}+a_{n} \int_{\Omega} w_{n}^{q}=\mu_{1} \int_{\Omega} w_{n}^{p}
$$

we have that $w_{n} \rightarrow 0$ in $L^{q}(\Omega)$. So $w=0$ and therefore $w_{n} \rightarrow 0$ strongly in $H_{0}^{1}(\Omega)$. This is a contradiction because by the Sobolev inequality,

$$
\int_{\Omega}\left|\nabla w_{n}\right|^{2} \leq \mu_{1} \int_{\Omega} w_{n}^{p} \leq C\left(\int_{\Omega}\left|\nabla w_{n}\right|^{2}\right)^{p / 2}
$$

for some constant $C$, so $\left\|w_{n}\right\|$ is bounded away from 0 .
It is well known that, when the system (1.1) is variational, least energy solutions are bounded in $\mathcal{H}$, independently of $\lambda_{i j}$. We close this section with the following open question.

Problem 5.5 Given $\lambda_{i j, n} \rightarrow-\infty$ for $i \neq j$, does the system (1.1) with $\lambda_{i j}$ replaced by $\lambda_{i j, n}$ have a solution $u_{n}$ such that the sequence $\left(u_{n, i}\right)$ is bounded in $H_{0}^{1}(\Omega)$ for all $i$ ?

Acknowledgements The authors thank the referee for useful suggestions.
Funding Open access funding provided by Stockholm University.
Data Availability Not applicable as no data were generated or analysed.
Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http://creativecommons.org/licenses/by/4.0/.

## A Appendix

In this appendix we prove Lemma 1.2. We employ some arguments which may be found in [11, 13]. First we note that by standard regularity results the solutions $u_{i}$ of (1.3) are in $\mathcal{C}^{2}(\Omega) \cap \mathcal{C}(\bar{\Omega})$.

Suppose there exists a sequence of solutions $\left(u_{n}\right)$ with $\left|u_{n}\right|_{\infty} \rightarrow \infty$. Passing to a subsequence, we may assume $\left|u_{n, i}\right|_{\infty} \rightarrow \infty$ and $\left|u_{n, i}\right|_{\infty} \geq\left|u_{n, j}\right|_{\infty}$ for some $i$ and
all $j$. There exists $x_{n} \in \Omega$ such that

$$
\max _{x \in \Omega} u_{n, i}(x)=u_{n, i}\left(x_{n}\right)
$$

Let $\beta:=\frac{2}{p-1}$ and choose $\varrho_{n}$ so that

$$
\varrho_{n}^{\beta}\left|u_{n, i}\right|_{\infty}=1
$$

Then $\varrho_{n} \rightarrow 0$ and passing to a subsequence, $x_{n} \rightarrow x_{0} \in \bar{\Omega}$. Let

$$
\Omega_{n}:=\left\{y \in \mathbb{R}^{N}: \varrho_{n} y+x_{n} \in \Omega\right\}
$$

and

$$
\begin{equation*}
v_{n, j}(y):=\varrho_{n}^{\beta} u_{n, j}\left(\varrho_{n} y+x_{n}\right), \quad j=1, \ldots, \ell \tag{A.1}
\end{equation*}
$$

Then,

$$
\begin{equation*}
0 \leq v_{n, i} \leq 1, \quad v_{n, i}(0)=1 \quad \text { and }\left.v_{n, j}\right|_{\partial \Omega_{n}}=0 \text { for all } j \tag{A.2}
\end{equation*}
$$

Passing to a subsequence, there are two possible cases and we shall complete the proof by ruling out both of them. Denote the distance from $x$ to a set $A$ by $d(x, A)$.
Case 1. $\frac{d\left(x_{n}, \partial \Omega\right)}{\varrho_{n}} \rightarrow \infty$.
Since $\varrho_{n} y+x_{n} \in \Omega$ if $|y|<\frac{d\left(x_{n}, \partial \Omega\right)}{\varrho_{n}}$, for each $R>0$ there exists $n_{0}$ such that $B_{R}(0) \subset \Omega_{n}$ whenever $n \geq n_{0}$. For $y \in B_{R}(0)$ and $n \geq n_{0}$ we have

$$
\begin{aligned}
-\Delta_{y} v_{n, i} & =\varrho_{n}^{\beta+2} \Delta_{x} u_{n, i}=\varrho_{n}^{\beta+2}\left(\mu_{i} u_{n, i}^{p}+\sum_{j \neq i} \lambda_{i j} u_{n, i}^{\alpha_{i j}} u_{n, j}^{\beta_{i j}}\right) \\
& =\varrho_{n}^{\beta+2-\beta p} \mu_{i} v_{n, i}^{p}+\sum_{j \neq i} \varrho_{n}^{\beta+2-\beta\left(\alpha_{i j}+\beta_{i j}\right)} \lambda_{i j} v_{n, i}^{\alpha_{i j}} v_{n, j}^{\beta_{i j}}
\end{aligned}
$$

Since $\beta+2-\beta p=0$ and $\gamma_{i j}:=\beta+2-\beta\left(\alpha_{i j}+\beta_{i j}\right)>0$, we can re-write this identity as

$$
-\Delta v_{n, i}=\mu_{i} v_{n, i}^{p}+\sum_{j \neq i} \varrho_{n}^{\gamma_{i j}} \lambda_{i j} v_{n, i}^{\alpha_{i j}} v_{n, j}^{\beta_{i j}}
$$

By elliptic estimates, $\left(v_{n, i}\right)$ is bounded in $W^{2, q}\left(B_{R}(0)\right)$ for some $q>N$. So passing to a subsequence, $v_{n, i} \rightarrow v_{i}$ weakly in $W^{2, q}\left(B_{R}(0)\right)$ and strongly in $\mathcal{C}^{1}\left(B_{R}(0)\right)$. Since $\varrho_{n}^{\gamma_{i j}} \rightarrow 0, v_{i}$ is a nonnegative solution to the equation

$$
-\Delta v=\mu_{i} v^{p}
$$

in $B_{R}(0)$. Let now $R_{m} \rightarrow \infty$. Then for each $m$ we get a solution $v_{i, m}$ of the above equation in $B_{R_{m}}(0)$. Passing to subsequences and applying the diagonal procedure, we see that $v_{i, m} \rightarrow w$, weakly in $W_{l o c}^{2, q}\left(\mathbb{R}^{N}\right)$ and strongly in $\mathcal{C}_{l o c}^{1}\left(\mathbb{R}^{N}\right)$. So $-\Delta w=\mu_{i} w^{p}$ in $\mathbb{R}^{N}, w \geq 0, w(0)=1$ according to (A.2), and $w \in \mathcal{C}^{2}\left(\mathbb{R}^{N}\right)$ by Schauder estimates. Replacing $w$ with $c w$ for a suitable $c>0$ we may assume $\mu_{i}=1$. Hence it follows from [13, Theorem 1.2] that $w=0$ which rules out Case 1.
Case 2. $\frac{d\left(x_{n}, \partial \Omega\right)}{\varrho_{n}} \rightarrow d \in[0, \infty)$.
It is clear that $x_{0} \in \partial \Omega$ and we may assume without loss of generality that $x_{0}=0$ and $v=(0, \ldots, 0,1)$ is the unit outer normal to $\partial \Omega$ at $x_{0}$. Let

$$
\mathbb{H}^{N}:=\left\{y \in \mathbb{R}^{N}: y_{N}<d\right\} \quad \text { where } y=\left(y_{1}, \ldots, y_{N}\right)
$$

We shall need the following result.
Lemma A. 1 (i) Let $A \subset \mathbb{H}^{N}$ be compact. Then there exists $n_{0}$ such that $\varrho_{n} y+x_{n} \in \Omega$ for all $n \geq n_{0}$ and $y \in A$.
(ii) Let $A \subset \mathbb{R}^{N} \backslash \overline{\mathbb{H}^{N}}$ be compact. Then there exists $n_{0}$ such that $\varrho_{n} y+x_{n} \notin \Omega$ for all $n \geq n_{0}$ and $y \in A$.

Proof (i) : Since $A$ is compact, there exists $\varepsilon>0$ such that $y_{N}<d-2 \varepsilon$ for all $y \in A$. For each $n$ there exists $\widehat{x}_{n} \in \partial \Omega$ which is closest to $x_{n}$, i.e., $d\left(x_{n}, \partial \Omega\right)=\left|x_{n}-\widehat{x}_{n}\right|$. As $\partial \Omega$ is tangent to the hyperplane $x_{N}=0$ at 0 ,

$$
\frac{\left|x_{n}-\widehat{x}_{n}\right|}{\varrho_{n}}=\frac{\widehat{x}_{n, N}-x_{n, N}}{\varrho_{n}}+o(1) .
$$

Therefore,

$$
\frac{x_{n, N}-\widehat{x}_{n, N}}{\varrho_{n}}<-d+\varepsilon \quad \text { and } \quad y_{N}+\frac{x_{n, N}-\widehat{x}_{n, N}}{\varrho_{n}}<-\varepsilon
$$

for all $y \in A$ if $n$ is large enough. There exists $C>0$ such that

$$
\left|y+\frac{x_{n}-\widehat{x}_{n}}{\varrho_{n}}\right| \leq C \quad \text { for all } y \in A
$$

Using this, we see that there is $n_{0}$ such that, if $n \geq n_{0}$ and $y \in A$, then $\varrho_{n} y+x_{n}-\widehat{x}_{n} \in \Omega$ and, as $\widehat{x}_{n} \in \partial \Omega$ and $\partial \Omega$ is tangent to the hyperplane $x_{N}=0$ at $0, \varrho_{n} y+x_{n}=$ $\varrho_{n} y+\left(x_{n}-\widehat{x}_{n}\right)+\widehat{x}_{n} \in \Omega$.
(ii) : This time $y_{N}>d+2 \varepsilon$ for $y \in A$,

$$
\frac{x_{n, N}-\widehat{x}_{n, N}}{\varrho_{n}}>-d-\varepsilon \quad \text { and } \quad y_{n, N}+\frac{x_{n, N}-\widehat{x}_{n, N}}{\varrho_{n}}>\varepsilon
$$

if $n$ is sufficiently large, and the conclusion follows by a similar argument as above.

Now we can continue with Case 2 . Let $\omega_{R}:=B_{R}(0) \cap\left\{y \in \mathbb{R}^{N}: y_{N}<d-1 / R\right\}$. Then $\bar{\omega}_{R} \subset \Omega_{n}$ for all $n \geq n_{0}$ by Lemma A.1(i). Let $v_{n, i}$ be given by (A.1) and using (A.2) extend it by 0 outside $\Omega_{n}$. According to Lemma A.1(ii), if $A \subset \mathbb{R}^{N} \backslash \overline{\mathbb{H}^{N}}$ is compact, then $\varrho_{n} y+x_{n} \notin \Omega$ for all $y \in A$ and $n$ large enough. So

$$
\begin{equation*}
\lim _{n \rightarrow \infty} v_{n, i}(x)=0 \text { for all } x \notin \mathbb{H}^{N} \tag{A.3}
\end{equation*}
$$

We can repeat the argument of Case 1 which now gives a nonegative solution to the equation $-\Delta w=\mu_{i} w^{p}$ in $\mathbb{H}^{N}$ such that $w(0)=1$. By (A.3), $w=0$ on $\partial \Omega$. As before, $w \in \mathcal{C}^{2}\left(\mathbb{H}^{N}\right)$, and since the extended functions $v_{n, i}$ are continuous in $\mathbb{R}^{N}$, $w \in \mathcal{C}^{0}\left(\overline{\mathbb{H}^{N}}\right)$. So $w=0$ according to [13,Theorem 1.3], a contradiction. Hence also Case 2 is ruled out.

## References

1. Caffarelli, L., Patrizi, S., Quitalo, V.: On a long range segregation model. J. Eur. Math. Soc. 19, 3575-3628 (2017)
2. Castro, A., Cossio, J., Neuberger, J.M.: A sign-changing solution for a superlinear Dirichlet problem. Rocky Mountain J. Math. 27, 1041-1053 (1997)
3. Chang, K.C.: Infinite Dimensional Morse Theory and Multiple Solution Problems. Birkhäuser, Boston (1993)
4. Clapp, M., Szulkin, A.: A simple variational approach to weakly coupled competitive elliptic systems. Nonl. Diff. Eq. Appl. 26:26, 21 pp (2019)
5. Conti, M., Terracini, S., Verzini, G.: Nehari's problem and competing species systems. Ann. Inst. H. Poincaré Anal. Non Linéaire 19, 871-888 (2002)
6. Conti, M., Terracini, S., Verzini, G.: Asymptotic estimates for the spatial segregation of competitive systems. Adv. in Math. 195, 524-560 (2005)
7. Crooks, E.C.M., Dancer, E.N.: Highly nonlinear large-competition limits of elliptic systems. Nonl. Anal. 73, 1447-1457 (2010)
8. Dancer, E.N., Du, Y.: Competing species equations with diffusion, large interactions, and jumping nonlinearities. J. Diff. Eq. 114, 434-475 (1994)
9. Dancer, E.N., Du, Y.: Positive solutions for a three-species competition with diffusion-I. General existence results, II. The case of equal birth rates. Nonl. Anal. 24, 337-357 and 359-373 (1995)
10. Dancer, E.N., Wei, J., Weth, T.: A priori bounds versus multiple existence of positive solutions for a nonlinear Schrödinger system. Ann. Inst. H. Poincaré - Anal. Non Linéaire 27, 953-969 (2010)
11. de Figueiredo, D.G., Yang, J.F.: A priori bounds for positive solutions of a non-variational elliptic system. Comm. PDE 26, 2305-2321 (2001)
12. Dold, A.: Lectures on Algebraic Topology, 2nd edn. Springer-Verlag, Berlin (1980)
13. Gidas, B., Spruck, J.: A priori bounds for positive solutions of nonlinear elliptic equations. Comm. PDE 6, 883-901 (1981)
14. Gilbarg, D., Trudinger, N.S.: Elliptic Partial Differential Equations of Second Order. Springer-Verlag, Berlin (2001)
15. Lee, J.M.: Riemannian manifolds. An introduction to curvature. Graduate Texts in Mathematics 176, Springer-Verlag, New York (1997)
16. Soave, N., Tavares, H., Terracini, S., Zilio, A.: Hölder bounds and regularity of emerging free boundaries for strongly competing Schrödinger equations with nontrivial grouping. Nonl. Anal. 138, 388-427 (2016)
17. Struwe, M.: Variational Methods, 4th edn. Springer-Verlag, Berlin (2008)
18. Willem, M.: Minimax Theorems. Birkhäuser, Boston (1996)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.


[^0]:    M. Clapp was partially supported by CONACYT grant A1-S-10457 (Mexico).
    $\boxtimes$ Andrzej Szulkin
    andrzejs@math.su.se
    Mónica Clapp
    monica.clapp@im.unam.mx
    1 Instituto de Matemáticas, Universidad Nacional Autónoma de México, Circuito Exterior, Ciudad Universitaria, 04510 Coyoacán, Ciudad de México, Mexico
    2 Department of Mathematics, Stockholm University, 10691 Stockholm, Sweden

