Erratum to: Gravity currents in two-layer stratified media

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This brief communication corrects a minor error in the presentation of the model results described in §2 of Tan et al. [1] (referred to hereafter as T. *al.*). Previously, we predicted that the solution of Eqs. 2.10 and 2.11 becomes unphysical for $g'_{12}/g'_{02} \geq 3/4$ irrespective of h'_1 . Upon closer inspection, we observe that model breakdown occurs at $g'_{12}/g'_{02} = 3/4$ when $h'_1 \lesssim 0.50$ but extends beyond this value when $h'_1 \gtrsim 0.50$. This is illustrated in Fig. 2a which shows the surface plot of Fr vs. g'_{12}/g'_{02} and h'_1 and is identical to Fig. 3a of T. *al.* with the exception of the aforementioned correction. In like fashion, Figs. 1, 2, 3 and 4 are, respectively, the corrected versions of Figs. 2, 3, 4 and 5 from T. *al.* (Note, that the geometric variables h'_1 , h_1 , h'_2 and h_2 used here have been non-dimensionalized by H , the channel depth.) As these new figures make clear, the region of parameter space where the model is valid and, by extension, where the gravity current is predicted to be supercritical, is larger than described in the original manuscript. Notwithstanding this correction, the principal conclusions of T. *al.* remain unchanged.

In addition to calculating head loss along streamlines as we do with Eqs. 2.14 and 2.15 of T. *al.*, one can also evaluate the global dissipation by determining the change of $D = \int u(p + \frac{1}{2}\rho u^2 + \rho g z) dz$ from far upstream to far downstream. The expression for ΔD , written in non-dimensional form reads

$$\Delta D = \text{Fr} \left[\frac{g'_{12}}{g'_{02}} (h'_2 h_1 - h'_1 h'_2) + \frac{g'_{01}}{g'_{02}} h'_1 h_0 + h'_2 h_0 - \frac{1}{2} \text{Fr}^2 \left(\frac{h_1^3}{h_1^2} + \frac{h_2^3}{h_2^2} \right) \right] \quad (1)$$

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Fig. 1 As in Fig. 2 of T. *al.* except that the region of model validity is extended past $g'_{12}/g'_{02} = 3/4$ when $h'_1 \gtrsim 0.5$. Broken lines for large h'_1 and g'_{12}/g'_{02} indicate model breakdown

[2]. Figure 5 shows ΔD vs. g'_{12}/g'_{02} and h'_1 using results of Eqs. 2.10 (panel a) and 2.11 (panel b). From Fig. 5a, we note that ΔD attains its maximum value when $h'_1 \simeq 0.49$ and $g'_{12}/g'_{02} \simeq 0.50$ while from Fig. 5b, ΔD attains its maximum value when $h'_1 \simeq 0.58$ and $g'_{12}/g'_{02} \simeq 0.43$. Although the location of this maximum varies slightly depending on the model assumptions, the surface plots of Fig. 5a, b are otherwise very similar. In particular, $\Delta D \rightarrow 0$ when $h'_1 \rightarrow 0$, $h'_1 \rightarrow 1$ or $g'_{12}/g'_{02} \rightarrow 0$. Also, and consistent with Figs. 2b and 4b, $\Delta D \rightarrow 0$ when $g'_{12}/g'_{02} = 3/4$ with $h'_1 \lesssim 0.5$.

Fig. 2 As in Fig.3 of T. *al.* except that the region of model validity is extended past $g'_{12}/g'_{02} = 3/4$ when $h'_1 \gtrsim 0.5$

Fig. 3 As in Fig. 4 of T. *al.* except that the region of model validity is extended past $g'_{12}/g'_{02} = 3/4$ when $h'_1 \gtrsim 0.5$. Broken lines for large h'_1 and g'_{12}/g'_{02} indicate model breakdown

Fig. 4 As in Fig. 5 of T. *al.* except that the region of model validity is extended past $g'_{12}/g'_{02} = 3/4$ when $h'_1 \gtrsim 0.5$

Fig. 5 The non-dimensional dissipation function, ΔD , defined by (1), plotted against h'_1 and g'_{12}/g'_{02} . Panels **a** and **b** corresponds to the dissipation associated with Eqs. 2.10 and 2.11 of T. *al.*, respectively

References

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