# Branches of positive solutions for some semilinear Schrödinger equations 

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## 1 Introduction

Semilinear elliptic problems on all of $\mathbb{R}^{N}$ of the form

$$
\begin{equation*}
-\Delta u+q(x) u=\lambda u+g(x) h(u) u, \quad x \in \mathbb{R}^{N} \tag{P}
\end{equation*}
$$

have been widely investigated under various assumptions on $q, g$ and $h$, see, for example, $[6,7,13]$ and references therein.

In particular, the results of [7] deal with the case in which $q \in L^{\infty}, g>0$, $h(0)=0, h(s)$ is strictly decreasing and $h(s) \rightarrow-\infty$ as $s \rightarrow+\infty$ and yield the existence of a bifurcation branch of positive solutions of $(P)$ that, roughly, blows up (in a suitable Lebesgue norm) as $\lambda$ tends to $a$, the infimum of the essential spectrum of the linear Schrödinger operator $-\Delta+q$.

The main purpose of the present paper is to consider that case in which $h$ has, roughly, the same asymptotic behaviour but is not necessarily decreasing and $g(x)$ can possibly vanish in a bounded domain $\Omega_{0} \subset \mathbb{R}^{N}$.

These specific features of $h$ and $g$ are motivated also by some problems arising in Nonlinear Optics. ${ }^{1}$ Actually, the study of nonlinear modes in a layered structure leads to a Schrödinger equation of the form (see [2,8,9 and 13])

$$
\begin{equation*}
u^{\prime \prime}+\varepsilon\left(x, u^{2}\right) u=k^{2} u, \quad x \in \mathbb{R}, \tag{1.1}
\end{equation*}
$$

[^0]where, whenever the material out of the layer is assumed to be defocusing,
\[

\varepsilon\left(x, u^{2}\right)= $$
\begin{cases}\alpha_{1} & \text { for }|x| \leqq d \\ \alpha_{2}-u^{2} & \text { for }|x|>d\end{cases}
$$
\]

Setting $\lambda=-k^{2}$,

$$
q(x)=\left\{\begin{array}{ll}
-\alpha_{1} & \text { for }|x| \leqq d \\
-\alpha_{2} & \text { for }|x|>d
\end{array} \quad g(x)= \begin{cases}0 & \text { for }|x| \leqq d \\
1 & \text { for }|x|>d\end{cases}\right.
$$

and $h(u)=-u^{2},(1.1)$ becomes of the form $(P)$ with a nonlinear term like that discussed in the present paper.

In order to find bifurcation branches of positive solutions of $(P)$ we approximate $(P)$ with Dirichlet boundary value problems on balls. Actually, we deal in Sect. 2 with an elliptic eigenvalue problem such as

$$
\begin{cases}-\Delta u+q(x) u=\lambda u+g(x) h(u) u, & x \in \Omega \\ u(x)=0, & x \in \partial \Omega\end{cases}
$$

where $\Omega$ is a general bounded domain in $\mathbb{R}^{N}$. When $g$ vanishes on a subset $\Omega_{0}$ of $\Omega$, problem $\left(P_{\Omega}\right)$ has been studied by Alama and Tarantello in [1] by variational methods, see also [10]. Unlike [1] we use here bifurcation theory and improve those results by showing that the branch of positive solutions of $\left(P_{\Omega}\right)$ bifurcating from the trivial solution at $\lambda=\lambda_{\Omega}$ (the first eigenvalue of $-\Delta+q$ on $W_{0}^{1,2}(\Omega)$, blows up in $L^{p}, p \geqq 1$, as $\lambda \rightarrow \lambda_{\Omega_{0}}$; moreover $\left(P_{\Omega}\right)$ has no positive solutions for $\lambda \geqq \lambda_{\Omega_{0}}$, see Theorem 2.6.

In Sect. 3 we turn to problem $(P)$ and prove a general global bifurcation result, see Theorem 3.4, under an uniform a-priori estimate on the branches of the approximated problems. Taking $\Omega=B_{R}$ and letting $R \rightarrow \infty$, this a-priori estimate allows us to show that the branches of the approximated problems converge, in an appropriate sense, to a branch of positive solutions of ( $P$ ). As a first application, we handle a problem studied by Brézis and Kamin in [5], see Theorem 3.6.

In Sect. 4 we still deal with $(P)$ and show that the bound above can be actually found provided that, roughly, the principal eigenvalue $\Lambda$ of $-\Delta+q$ on $W^{1,2}\left(\mathbb{R}^{N}\right)$ is smaller than $q$, see Theorem 4.4 for the precise statement. In particular this result applies when $g(x)>0$ on $\mathbb{R}^{N}$ yielding an improvement of the existence results of Edelson and Stuart [7].

Finally, in Sect. 5 we assume that $\Omega_{0} \neq \emptyset$ and, roughly, $\lambda_{\Omega_{0}}<a$, and show, by an appropriate comparison with problems with decreasing nonlinearities, that the branch of positive solutions of $(P)$ blows up in $L^{p}, p \geqq 1$, iff $\lambda \uparrow \lambda_{\Omega_{0}}$, see Theorem 5.2. Such a result is in striking contrast with that in [7], see Remark 5.3.

Notation. In the sequel $\Omega$ denotes a bounded domain of $\mathbb{R}^{N}$ with (smooth) boundary $\partial \Omega$.
$W_{0}^{1,2}(\Omega)$ or $W_{0}^{1,2}\left(\mathbb{R}^{N}\right)$ denote Sobolev spaces and $L^{p}=L^{p}(\Omega)$ or $L^{p}=$ $L^{p}\left(\mathbb{R}^{N}\right)$ denote Lebesgue spaces. For brevity and whenever unambiguous, the
indication of $\Omega$ or $\mathbb{R}^{N}$ will be omitted. The standard norm in $L^{p}$ will be denoted by $|u|_{p}$.

In the rest of the paper we will often extend a function $z$ outside its support by setting $z(x) \equiv 0$ therein. To keep the notation as light as possible, this extended function will still be denoted by $z$.

## 2 Problems on bounded domains

In this section we deal with problem $\left(P_{\Omega}\right)$. We set $E=W_{0}^{1,2}(\Omega)$, endowed with norm $\|u\|^{2}:=\int|\nabla u|^{2}$. The first eigenvalue of the linear problem

$$
\begin{cases}-\Delta u+q(x) u=\lambda u, & x \in \Omega \\ u(x)=0, & x \in \partial \Omega\end{cases}
$$

will be denoted by $\lambda_{\Omega}[q]$, or simply by $\lambda_{\Omega}$. We also denote by $\phi_{\Omega}$, the eigenfunction corresponding to $\lambda_{\Omega}$, with $\phi_{\Omega}(x)>0$ and $\left|\phi_{\Omega}\right|_{2}=1$. We shall use the variational characterization of $\lambda_{\Omega}$, namely:

$$
\lambda_{\Omega}=\inf _{u \in E,|u|_{2}=1} \int_{\Omega}\left[|\nabla u|^{2}+q u^{2}\right]
$$

We consider problem $\left(P_{\Omega}\right)$ and assume
(Q) $\quad q \in L^{\infty}(\Omega)$;
$\left(A_{1}\right) \quad g \in L^{\infty}(\Omega), \quad g \geqq 0$, and there exists a (bounded) domain $\Omega_{0}$ with smooth $\left(\right.$ say $\left.C^{1, v}\right)$ boundary $\partial \Omega_{0}$ such that $\Omega_{0} \subset \Omega$ and $g(x)=0$ iff $x \in \Omega_{0} ;$
$\left(A_{2}\right) \quad h \in C\left(\mathbb{R}^{+}\right), h(0)=0, \exists \kappa>0, c_{0}>0$ such that $h(s) \leqq c_{0} s^{\kappa}, \forall s>0$ and $h(s) \rightarrow-\infty$ as $s \rightarrow+\infty$.

In particular, we explicitly point out that $\left(A_{2}\right)$ implies there exists $C_{0} \geqq 0$ such that

$$
\begin{equation*}
h(s) \leqq C_{0} \tag{2.1}
\end{equation*}
$$

In the sequel it is understood that assumptions $(Q)$ and $\left(A_{1}-A_{2}\right)$ hold true.
a) Some preliminary Lemmas. Although the following lemma is perhaps well known, we give an outline of the proof for the reader's convenience.

Hereafter, by a positive solution of $\left(P_{\Omega}\right)$ we mean an $u \in E, u>0$, which solves $\left(P_{\Omega}\right)$ weakly. Actually, in our case, weak solutions belong to $C^{1, v}$. Of course, if $q$ and $g$ are Hölder continuous, $u$ will become a classical solution.

Lemma 2.1 From $\left(\lambda_{\Omega}, 0\right)$ bifurcates an unbounded branch of positive solutions of problem $\left(P_{\Omega}\right)$.

Proof. Let $M \geqq 0$ be such that $q(x)+M \geqq 0$ in $\Omega$, and consider the problem:

$$
\begin{cases}-\Delta u+[q(x)+M] u=\mu u+g(x) h(u) u & \text { in } \Omega  \tag{2.2}\\ u(x)=0 & \text { on } \partial \Omega\end{cases}
$$

Trivially, $(\lambda, u)$ is a solution of problem $\left(P_{\Omega}\right)$ iff $(\mu, u)=(\lambda+M, u)$ is a solution of problem (2.2).

Letting $K=(-\Delta+[q+M])^{-1}$, equation (2.2) can be written as

$$
\begin{equation*}
u=\mu K u+K(g h(u) u) . \tag{2.3}
\end{equation*}
$$

It is immediate to check that problem (2.3) satisfies the hypotheses of the Global Bifurcation Theorem of Rabinowitz (see [11]). Then a branch of positive solutions $(\mu, u)$ of (2.2) bifurcates from $\lambda_{\Omega}[q+M]=\lambda_{\Omega}[q]+M$ and yields a branch of solutions of $\left(P_{\Omega}\right)$ bifurcating from $\lambda_{\Omega}$. Moreover, by standard arguments it can be proved that this branch remains in the interior of the cone of positive functions of $C_{0}^{1}(\Omega)$. Since $\lambda_{\Omega}$ is the only eigenvalue with corresponding positive eigenfunction, the branch cannot meet another eigenvalue different from $\lambda_{\Omega}$, and thus it is unbounded.

Let $\lambda_{0}=\lambda_{\Omega_{0}} \quad$ and $\quad \phi_{0}=\phi_{\Omega_{0}}$.
Lemma 2.2 There exists $\lambda \in \mathbb{R}$ such that for every positive solutions $(\lambda, u)$ of ( $P_{\Omega}$ ) one has

$$
\lambda<\lambda<\lambda_{0}
$$

Proof. We set $q=\inf _{\Omega} q(x)$ and $g=\sup _{\Omega} g(x)$. Let $(\lambda, u)$ be a positive solution of $\left(P_{\Omega}\right)$. Then

$$
\|u\|^{2}+\int_{\Omega} q(x) u^{2}=\lambda|u|_{2}^{2}+\int_{\Omega} g(x) h(u) u^{2}
$$

and one has:

$$
\lambda=\frac{\|u\|^{2}+\int_{\Omega} q(x) u^{2}-\int_{\Omega} g(x) h(u) u^{2}}{|u|_{2}^{2}} \geqq \lambda:=q-g C_{0}
$$

where $C_{0}$ is given in (2.1).
Next, from $\left(P_{\Omega}\right)$ it follows that

$$
-\int_{\Omega} \Delta u \phi_{0}+\int_{\Omega} q(x) u \phi_{0}=\lambda \int_{\Omega} u \phi_{0}+\int_{\Omega} g(x) h(u) u \phi_{0}
$$

Since $g(x) \equiv 0$ on $\Omega_{0}$ one infers that

$$
\begin{equation*}
-\int_{\Omega_{0}} \Delta u \phi_{0}-\int_{\Omega_{0}} q(x) u \phi_{0}=\lambda \int_{\Omega_{0}} u \phi_{0} . \tag{2.4}
\end{equation*}
$$

Since $u>0$ and $\partial \phi_{0} / \partial \mathbf{n}<0$ on $\partial \Omega_{0}\left(\mathbf{n}\right.$ denotes the outer unit normal at $\left.\Omega_{0}\right)$, an integration by parts yields

$$
-\int_{\Omega_{0}} \Delta u \phi_{0}<-\int_{\Omega_{0}} u \Delta \phi_{0}=\int_{\Omega_{0}} u\left[\lambda_{0}-q(x)\right] \phi_{0} .
$$

This and (2.4) imply that $\lambda<\lambda_{0}$.

Remark 2.3 Let us explicitely point out that the lower bound $\lambda$ does not depend upon $\Omega$.
b) Blow up as $\lambda \uparrow \lambda_{0}$. We begin with the following Lemma.

Lemma 2.4 Let $\left(\lambda_{n}, u_{n}\right)$ be a sequence of positive solutions of $\left(P_{\Omega}\right)$ such that $\left|u_{n}\right|_{2} \rightarrow \infty$. Then $\lambda_{n} \rightarrow \lambda_{0}, \lambda_{n}<\lambda_{0}$.

Proof. By Lemma 2.2 we can assume that $\lambda_{n} \rightarrow \widehat{\lambda} \in\left[\lambda, \lambda_{0}\right]$. Setting $z_{n}=u_{n} /\left|u_{n}\right|_{2}$ one has

$$
\begin{cases}-\Delta z_{n}+q(x) z_{n}=\lambda_{n} z_{n}+g(x) h\left(u_{n}\right) z_{n}, & x \in \Omega  \tag{2.5}\\ z_{n}(x)=0 & x \in \partial \Omega\end{cases}
$$

Hence

$$
\int_{\Omega}\left|\nabla z_{n}\right|^{2}+\int_{\Omega} q(x) z_{n}^{2}=\lambda_{n} \int_{\Omega} z_{n}^{2}+\int_{\Omega} g(x) h\left(u_{n}\right) z_{n}^{2}
$$

Using (2.1) and since $\left|z_{n}\right|_{2}=1$ we infer from (2.5)

$$
\left\|z_{n}\right\|^{2}=\int_{\Omega}\left|\nabla z_{n}\right|^{2} \leqq \int_{\Omega}\left(\lambda_{n}-q(x)+g(x) C_{0}\right) z_{n}^{2} \leqq c_{1}
$$

Thus, up to a subsequence, $z_{n} \rightarrow z$ strongly in $L^{2}(\Omega)$, and $|z|_{2}=1$.
Let $D$ be any domain such that $D \subset \Omega \backslash \Omega_{0}$ and let $\phi \in C_{0}^{\infty}(D)$. From (2.5) it follows

$$
-\int_{D} z_{n} \Delta \phi+\int_{D} q(x) z_{n} \phi-\lambda_{n} \int_{D} z_{n} \phi=\int_{D} g(x) h\left(u_{n}\right) z_{n} \phi .
$$

If $z(x)>0$ for a.e. $x \in D$, one has $u_{n}(x)=z_{n}(x)\left|u_{n}\right|_{2} \rightarrow \infty$ a.e. in $D$. Since $g>0$ in $D,\left(A_{2}\right)$ implies

$$
\int_{D} g(x) h\left(u_{n}\right) z_{n} \phi \rightarrow-\infty
$$

On the other hand

$$
\int_{D}\left[-z_{n} \Delta \phi+q(x) z_{n} \phi-\lambda_{n} z_{n} \phi\right] \rightarrow \int_{D}[-z \Delta \phi+q(x) z \phi-\widehat{\lambda} z \phi]>-\infty
$$

a contradiction. This shows that $z(x)=0$ a.e. in $\Omega \backslash \Omega_{0}$. Recalling that (see [4], Proposition IX.18)

$$
W_{0}^{1,2}\left(\Omega_{0}\right)=\left\{u \in W_{0}^{1,2}(\Omega): u=0 \text { a.e. in } \Omega \backslash \Omega_{0}\right\}
$$

it follows that $z \in W_{0}^{1,2}\left(\Omega_{0}\right)$.
Now, let $\varphi \in C_{0}^{\infty}\left(\Omega_{0}\right)$. Since $g(x) \equiv 0$ in $\Omega_{0}$, one has

$$
\int_{\Omega} g(x) h\left(u_{n}\right) z_{n} \varphi=0
$$

for all $n$. Then, multiplying (2.5) by $\varphi$ and integrating by parts, one finds

$$
-\int_{\Omega_{0}} z_{n} \Delta \varphi+\int_{\Omega_{0}} q(x) z_{n} \varphi=\lambda_{n} \int_{\Omega_{0}} z_{n} \varphi
$$

Passing to the limit one has

$$
-\int_{\Omega_{0}} z \Delta \varphi+\int_{\Omega_{0}} q(x) z \varphi=\widehat{\lambda} \int_{\Omega_{0}} z \varphi, \quad \forall \varphi \in C_{0}^{\infty}\left(\Omega_{0}\right)
$$

This shows that $z$ satisfies

$$
-\Delta z+q(x) z=\widehat{\lambda} z \quad \text { in } \Omega_{0}
$$

Since $z \in W_{0}^{1,2}\left(\Omega_{0}\right), z \geqq 0$ in $\Omega_{0}$ and $|z|_{2}=1$, it follows that $\widehat{\lambda}=\lambda_{0}$ and $z=\phi_{0}$.

Remarks 2.5 (i) The preceding proof shows that $u_{n} \simeq\left|u_{n}\right|_{2} \phi_{0}$, for $n$ large.
(ii) When $\Omega_{0}=\emptyset$, or else when $\Omega_{0}$ has zero Lebesgue measure, the preceding arguments prove that $\lambda_{n} \rightarrow+\infty$. Actually, if not, the first part of the proof shows that $z(x)=0$ a.e. in $\Omega$, which is in contradiction with the fact that $|z|_{2}=1$.
c) Branches of positive solutions. We are in position to prove the main result of this section. In the sequel we shall denote by $\Pi$ the canonical projection of $\mathbb{R} \times E$ onto $\mathbb{R}$.
Theorem 2.6 Assume $(Q)$ and $\left(A_{1}-A_{2}\right)$ hold. Then from $\left(\lambda_{\Omega}, 0\right)$ bifurcates an unbounded branch $S_{\Omega}$ of positive solutions of $\left(P_{\Omega}\right)$ such that:
(i) $\Pi\left(S_{\Omega}\right)=\left[\lambda^{*}, \lambda_{0}\left[\right.\right.$, for some $\lambda^{*} \leqq \lambda_{\Omega}$;
(ii) On $S_{\Omega}$ one has that, for all $p \geqq 1,|u|_{p} \rightarrow+\infty$ iff $\lambda \uparrow \lambda_{0}$.

Proof. It suffices to apply Lemmas 2.1, 2.2, 2.4 and Remark 2.5(i).
Remarks 2.7 (i) The preceding Theorem improves some results of [1].
(ii) If $h$ is a decreasing function, then $\left(P_{\Omega}\right)$ has a unique positive solution for all $\lambda \in] \lambda_{\Omega}, \lambda_{0}[$. In such a case the branch of positive solutions is a graph and could have been found by using sub- and super-solutions.
(iii) Completing Remark 2.5(ii), when $\Omega_{0}$ has zero Lebesgue measure or it is possibly empty, there is an unbounded branch $S_{\Omega}$ of positive solutions of $\left(P_{\Omega}\right)$ such that $\Pi\left(S_{\Omega}\right)=\left[\lambda^{*},+\infty[\right.$.
(iv) According to the classical Theorem of Bifurcation from the simple eigenvalue, the behavior of the branch $S_{\Omega}$ near $\lambda_{\Omega}$ depends on the sign of $h$ near 0 . In particular, if $h(s)>0$ in a right neighbourhood of $s=0$, then $\lambda^{*}<\lambda_{\Omega}$.

## 3 A general global bifurcation result

We turn to the Schrödinger equation on all of $\mathbb{R}^{N}$

$$
\begin{equation*}
-\Delta u+q(x) u=\lambda u+g(x) h(u) u, \quad x \in \mathbb{R}^{N} \tag{P}
\end{equation*}
$$

where, hereafter, $N \geqq$. In the sequel $q, g$ are supposed to satisfy $(Q)$ and $\left(A_{1}\right)$ with $\Omega=\mathbb{R}^{N}$, as well as $h$ is assumed to verify $\left(A_{2}\right)$. In any case, it is worth recalling that $\Omega_{0}$ is still assumed to be a bounded domain. We also set

$$
a=\liminf _{|x| \rightarrow \infty} q(x)(>-\infty)
$$

We shall work in the Sobolev space $H=W^{1,2}\left(\mathbb{R}^{N}\right)$ equipped with the usual norm

$$
\|u\|^{2}=\int_{\mathbb{R}^{N}}\left[|\nabla u|^{2}+u^{2}\right] d x .
$$

A positive solution of $(P)$ will be an $u \in H, u>0$, which satisfies $(P)$ weakly.
Let us recall (see, for example, [3] or [12]) that if ( $Q$ ) holds then the spectrum of the linear eigenvalue problem

$$
\begin{equation*}
-\Delta u+q(x) u=\lambda u, \quad x \in \mathbb{R}^{N} \tag{3.1}
\end{equation*}
$$

contains eigenvalues provided

$$
\begin{equation*}
\Lambda=\Lambda[q]:=\inf _{u \in H,|u|_{2}=1} \int_{\mathbb{R}^{N}}\left[|\nabla u|^{2}+q u^{2}\right]<a . \tag{3.2}
\end{equation*}
$$

Moreover, if (3.2) holds then $\Lambda$ is the principal eigenvalue of (3.1).
The existence of a branch of positive solutions for $(P)$ will be proved by an approximating procedure, carried out by means of the following topological lemma (see [14], Theorem 9.1).

Lemma 3.1 Let $X_{n}$ be a sequence of connected subsets of a complete metric space X. If
(i) $\lim \inf \left(X_{n}\right) \neq \emptyset$;
(ii) $\bigcup X_{n}$ is precompact;

Then $\lim \sup \left(X_{n}\right)$ is not empty, compact and connected.
Above, $\lim \inf \left(X_{n}\right)$ and $\lim \sup \left(X_{n}\right)$ denote the set of all $x \in X$ such that any neighbourhood of $x$ intersects all but finitely many of $X_{n}$, infinitely many of $X_{n}$ respectively. In order to use the preceding Lemma, let $B_{R}$ be the ball in $\mathbb{R}^{N}$ centered at the origin, of radius $R>0$ and let $\left(P_{R}\right)$ denote problem $\left(P_{\Omega}\right)$ with $\Omega=B_{R}$. We will always take $R$ sufficiently large in such a way that $\Omega_{0} \subset B_{R}$. We also set $\lambda_{R}:=\lambda_{B_{R}}$. According to Theorem 2.6, there exists a branch $S_{R}$, of positive solutions of ( $P_{R}$ ), that bifurcates from ( $\lambda_{R}, 0$ ), and blows up as $\lambda \uparrow \lambda_{0}$.

It is well known that $\lambda_{R}$ (is decreasing with respect to $R>0$ and) converges to $\Lambda$ as $R \rightarrow+\infty$. For future reference we add that, in particular, since $\lambda<$ $\lambda_{R}<\lambda_{0}$ for $R$ large, one has that $\lambda \leqq \Lambda<\lambda_{0}$.

To carry over the limiting procedure an uniform a-priori bound is in order. We suppose
(B) There exist $b \in] \Lambda, \lambda_{0}\left[\right.$ and $\Psi \in L^{\infty} \cap L^{2}, \Psi>0$, such that $u<\Psi$, for all $(\lambda, u) \in S_{R}$, uniformly in $R>0$ and $\lambda \in[\lambda, b]$.

We set $T=[\lambda, b]$ and $X=T \times H$. Let $R_{n} \rightarrow+\infty$ and denote by $X_{n}$ the connected component of the set $\left\{(\lambda, u) \in S_{R_{n}}: \lambda \in T\right\}$, such that $\left(\lambda_{R}, 0\right) \in X_{n}$. In view of the properties of $S_{R_{n}}$ discussed in the preceding section, $X_{n} \neq \emptyset$ and $b$ belongs to $\Pi\left(X_{n}\right)$, for all $n$ large.

Lemma 3.2 $\bigcup X_{n}$ is precompact in $X$.
Proof. Let $\left(\lambda_{k}, u_{k}\right) \in \bigcup X_{n}$ (we assume that the diameters of $\operatorname{supp}\left(u_{k}\right) \rightarrow \infty$, otherwise, the result is trivial). By ( $B$ ), it follows that

$$
\begin{equation*}
\left|u_{k}\right|_{2} \leqq c_{1} \tag{3.3}
\end{equation*}
$$

Moreover, one has

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left|\nabla u_{k}\right|^{2}+\int_{\mathbb{R}^{N}} q(x) u_{k}^{2}=\lambda_{k} \int_{\mathbb{R}^{N}} u_{k}^{2}+\int_{\mathbb{R}^{N}} g(x) h\left(u_{k}\right) u_{k}^{2} . \tag{3.4}
\end{equation*}
$$

Since $h\left(u_{k}\right) \leqq C_{0}$ then (3.3) and (3.4) imply $\left|\nabla u_{k}\right|_{2} \leqq c_{2}$ and hence

$$
\left\|u_{k}\right\| \leqq c_{3}
$$

Therefore, up to a subsequence, $u_{k} \rightarrow u$ in $H$ and in $L^{2}$. It is easy to see that $u$ satisfies

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} \nabla u \nabla \phi+\int_{\mathbb{R}^{N}} q(x) u \phi=\lambda \int_{\mathbb{R}^{N}} u \phi+\int_{\mathbb{R}^{N}} g(x) h(u) u \phi, \quad \forall \phi \in C_{0}^{\infty} . \tag{3.5}
\end{equation*}
$$

By density, (3.5) holds for all $\phi \in H$ and, in particular, for $\phi=u_{k}$. To prove that $u_{k}$ converges strongly to $u$ we consider

$$
\left\|u_{k}-u\right\|^{2}=\left\|u_{k}\right\|^{2}+\|u\|^{2}-2 \int_{\mathbb{R}^{N}} \nabla u_{k} \nabla u-2 \int_{\mathbb{R}^{N}} u_{k} u
$$

From (3.4) we find

$$
\begin{equation*}
\left\|u_{k}\right\|^{2}=\int_{\mathbb{R}^{N}}\left|\nabla u_{k}\right|^{2}+\int_{\mathbb{R}^{N}} u_{k}^{2}=\int_{\mathbb{R}^{N}} F_{\lambda_{k}}\left(x, u_{k}\right) u_{k}+\int_{\mathbb{R}^{N}} u_{k}^{2} \tag{3.6}
\end{equation*}
$$

where

$$
\begin{equation*}
F_{\lambda}(x, u)=\lambda u-q(x) u+g(x) h(u) u \tag{3.7}
\end{equation*}
$$

From (3.5) with $\phi=u_{k}$ and $\phi=u$ respectively, we infer

$$
\begin{align*}
\int_{\mathbb{R}^{N}} \nabla u_{k} \nabla u & =\int_{\mathbb{R}^{N}} F_{\lambda}(x, u) u_{k}  \tag{3.8}\\
\int_{\mathbb{R}^{N}}|\nabla u|^{2} & =\int_{\mathbb{R}^{N}} F_{\lambda}(x, u) u \tag{3.9}
\end{align*}
$$

Putting together (3.6), (3.8) and (3.9) we find

$$
\begin{array}{r}
\left\|u_{k}-u\right\|^{2}=\int_{\mathbb{R}^{N}} F_{\lambda_{k}}\left(x, u_{k}\right) u_{k}+\int_{\mathbb{R}^{N}} u_{k}^{2}+\int_{\mathbb{R}^{N}} F_{\lambda}(x, u) u+\int_{\mathbb{R}^{N}} u^{2}-2 \int_{\mathbb{R}^{N}} F_{\lambda}(x, u) u_{k}-2 \int_{\mathbb{R}^{N}} u_{k} u \\
=\int_{\mathbb{R}^{N}}\left[F_{\lambda_{k}}\left(x, u_{k}\right)-F_{\lambda}(x, u)\right] u_{k}+\int_{\mathbb{R}^{N}} F_{\lambda}(x, u)\left[u-u_{k}\right]+\int_{\mathbb{R}^{N}} u_{k}\left[u_{k}-u\right]+\int_{\mathbb{R}^{N}} u\left[u-u_{k}\right] .
\end{array}
$$

Since $u_{k}<\Psi$ we deduce

$$
\begin{aligned}
\left\|u_{k}-u\right\|^{2} \leqq & \int_{\mathbb{R}^{N}}\left|F_{\lambda_{k}}\left(x, u_{k}\right)-F_{\lambda}(x, u)\right| \Psi+\int_{\mathbb{R}^{N}}\left|F_{\lambda}(x, u)\right|\left|u-u_{k}\right| \\
& +\int_{\mathbb{R}^{N}}\left|u_{k}-u\right| \Psi+\int_{\mathbb{R}^{N}}|u|\left|u-u_{k}\right|
\end{aligned}
$$

The last three integrals converge to zero. As for the first one, since $F_{\lambda}$ is locally Lipschitzian, one has

$$
\int_{\mathbb{R}^{N}}\left|F_{\lambda_{k}}\left(x, u_{k}\right)-F_{\lambda}(x, u)\right| \Psi \leqq\left|\lambda_{k}-\lambda\right| \int_{\mathbb{R}^{N}} \Psi^{2}+c_{4} \int_{\mathbb{R}^{N}}\left|u_{k}-u\right| \Psi \rightarrow 0
$$

In conclusion, it follows that $u_{k} \rightarrow u$ strongly in $H$.
Let $S:=\lim \sup \left(X_{n}\right) \backslash\{(\Lambda, 0)\}$.
Lemma 3.3 If $(\lambda, u) \in S$ then $u$ is a (nontrivial) positive solution of $(P)$.
Proof. It is easy to check that any $(\lambda, u) \in S$ is a non-negative solution of $(P)$. We need to prove that $u \neq 0$. We shall consider separately the cases $\lambda>\Lambda$ and $\lambda<\Lambda$ (obviously, when $\lambda=\Lambda$ the result is trivial, according to the definition of $S$ ).

Step 1. Let $\lambda>\Lambda$ and suppose there exists a sequence $\left(\lambda_{k}, u_{k}\right) \in X_{n_{k}}, n_{k} \rightarrow \infty$, such that $\left(\lambda_{k}, u_{k}\right) \rightarrow(\lambda, 0)$. We can assume that all $\lambda_{k}>\Lambda+\delta$ for some $\delta>0$. Since $\lambda_{R} \downarrow \Lambda$, there exists $R>0$ such that $\lambda_{R}<\Lambda+\delta<\lambda$. Let $m$ be such that $R_{n_{k}}>R$ for all $k>m$. It is easy to check that, for such $k,\left(\lambda_{k}, u_{k}\right)$ are super-solutions of the problem

$$
\begin{cases}-\Delta u+q(x) u=(\Lambda+\delta) u+g(x) h(u) u, & |x|<R  \tag{3.10}\\ u(x)=0 & |x|=R\end{cases}
$$

Small subsolutions of type $\varepsilon_{k} \phi_{R}$ can also be obtained, and thus we can find a sequence of positive solutions of (3.10) that converges to 0 . In other words $(\Lambda+\delta, 0)$ is bifurcation of positive solutions for $\left(P_{R}\right)$. Since $\left(\lambda_{R}, 0\right)$ is the only possible bifurcation for positive solutions for $\left(P_{R}\right)$, this is a contradiction.
Step 2. Suppose that $\lambda<\Lambda$ and let $M>0$ be such that $q(x)+M \geqq 1$ in $\mathbb{R}^{N}$. It is convenient to endow the space $H$ with the norm

$$
\left\|\left.u\left|\|^{2}:=\int_{\mathbb{R}^{N}}\right| \nabla u\right|^{2}+\int_{\mathbb{R}^{N}}[q+M] u^{2},\right.
$$

which is equivalent to $\|\cdot\|$. By the variational expression of $\Lambda$, we deduce that

$$
0<(\Lambda+M) \int_{\mathbb{R}^{N}} u^{2} \leqq\| \| u\| \|^{2}, \quad \forall u \in H
$$

Let $0<\sigma<\kappa$ be such that $h(s)<c_{1} s^{\sigma}$ and $H \subset L^{\sigma+2}\left(\mathbb{R}^{N}\right)$.
Let $(\lambda, w)$ be a positive solution of some $\left(P_{R}\right)$. Then

$$
\int_{\mathbb{R}^{N}}\left(|\nabla w|^{2}+q(x) w^{2}\right)=\lambda \int_{\mathbb{R}^{N}} w^{2}+\int_{\mathbb{R}^{N}} g(x) h(w) w^{2}
$$

and hence

$$
\begin{aligned}
\|\|w\|\|^{2} & \left.=(\lambda+M) \int_{\mathbb{R}^{N}} w^{2}+\int_{\mathbb{R}^{N}} g(x) h(w) w^{2} \leqq \frac{\lambda+M}{\Lambda+M}\| \| w \right\rvert\, \|^{2}+c_{2} \int_{\mathbb{R}^{N}} w^{\sigma+2} \\
& \leqq \frac{\lambda+M}{\Lambda+M}\left|\left\|w\left|\left\|^{2}+c_{3}\right\|\right|\right\| w\| \|^{\sigma+2}\right.
\end{aligned}
$$

In particular, for $\lambda<\Lambda$,

$$
c_{3} \mid\|w\| \|^{\sigma} \geqq 1-\frac{\lambda+M}{\Lambda+M}>0
$$

Then the approximating branchs are uniformly bounded away from zero.
Theorem 3.4 Suppose that $(Q),\left(A_{1}-A_{2}\right)$ and $(B)$ hold. Then there exists a branch, $S$, of positive solutions of $(P)$ bifurcating from $(\Lambda, 0)$, such that $\Pi(S)=\left[\Lambda^{*}, b\right]$, for some $\Lambda^{*} \leqq \Lambda$.

Proof. We apply Lemma 3.1 with $X$ and $X_{n}$ defined before. Since $\lambda_{R_{n}} \downarrow \Lambda$, it immediately follows that $(\Lambda, 0) \in \lim \inf X_{n}$. Moreover, Lemma 3.2 shows that the compactness property (ii) of Lemma 3.1 is satisfied. As a consequence $S=$ $\left.\lim \sup \left(X_{n}\right) \backslash\{\Lambda, 0)\right\}$ is (not empty), connected and $(\Lambda, 0) \in S$. By Lemma 3.3 $(\lambda, u) \in S$ is a positive solution of $(P)$. Moreover, since $\Lambda<b<\lambda_{0}$ then $b \in \Pi\left(X_{n}\right)$ for all $n$ large, and the preceding compactness arguments imply that $b \in \Pi(S)$.

Remarks 3.5 (i) Similarly as in Remark 2.5(ii), if $g(x)>0$ a.e. in $\mathbb{R}^{N}$ then it is understood that $b$ could be any number greater than $\Lambda$.
(ii) We have found solutions in $W^{1,2}$. In fact in the specific applications discussed in the sequel the a-priori bound $\Psi$ will have an exponential decay at zero as $|x| \rightarrow \infty$ and thus the solutions on $S$ will have the same asymptotic behaviour.
(iii) Assumption ( $B$ ) can be slightly weakened. Actually it suffices to assume
( $\widetilde{B}) \quad$ There exist $b \in] \Lambda, \lambda_{0}\left[\right.$ and $\Psi \in L^{\infty} \cap L^{2}, \Psi>0$, such that $u<\Psi$, for all $(\lambda, u) \in X_{n}$, uniformly in $n$ and $\lambda \in[\lambda, b]$.

If for each $\lambda \in[\lambda, b]$ problem $\left(P_{R}\right)$ has a unique positive solution, $(B)$ and $(\widetilde{B})$ obviously coincide. However, in the applications discussed later on, where multiple solutions could arise, we will be able to prove $(\widetilde{B})$, only.
(iv) According to Remark 2.3 one obviously has that $\Lambda^{*} \geqq \lambda$.

Theorem 3.4 deals with the specific problem $(P)$. But a similar result holds true for a more general equation like

$$
-\Delta u=F_{\lambda}(x, u), \quad x \in \mathbb{R}^{N}
$$

where $F$ is locally Lipschitzian. It suffices that each approximated problem on $B_{R}$ possesses a branch of positive solutions bifurcating from some $\left(\lambda_{R}, 0\right)$, that $\lambda_{R}$ converge to some $\Lambda$ and that assumption $(B)$, or $(\widetilde{B})$, holds for some $\lambda, b$ such that $\lambda<\Lambda<b$. The preceding arguments can be still carried out and yield the existence of a global bifurcating branch.

This is, for example, the case of problem

$$
\begin{equation*}
-\Delta u=\lambda \rho(x) u^{\alpha}, \quad x \in \mathbb{R}^{N} \tag{3.11}
\end{equation*}
$$

where $\rho \in L^{\infty}$ and $0<\alpha<1$, studied in [5]. Here we can show:

Theorem 3.6 Assume that there exists $U \in L^{\infty} \cap L^{2}$ satisfying

$$
\begin{equation*}
-\Delta U=\rho(x), \quad x \in \mathbb{R}^{N} \tag{3.12}
\end{equation*}
$$

Then there exists a branch $S$ of positive, $W^{1,2}$ solutions of (3.11) bifurcating from $(0,0)$ and such that $\Pi(S)=[0, \infty)$.

Proof. We only give an outline of the proof and leave the details to the reader. For every $R>0$ the boundary value problem

$$
-\Delta u=\lambda \rho(x) u^{\alpha}, \quad x \in B_{R}, \quad u=0, x \in \partial B_{R}
$$

has a bifurcating branch $S_{R}$ emanating from $(0,0)$ such that $\Pi\left(S_{R}\right)=[0, \infty)$. Moreover, one easily shows that assumption $(B)$ holds true with $\Lambda=\lambda=0$, any $b>0$ and $\Psi(x)=C U(x)$, where $U$ satisfies (3.12) and $C>0$ is sufficiently large. Then the result follows by the arguments discussed before.

Remark 3.7 In [5] it is only assumed that $U \in L^{\infty}$. On the other hand, we find here (a branch of) solutions in $W^{1,2}$, not merely bounded solutions.

## 4 Existence of a global branch of positive solutions of (P)

A global branch of positive solutions of problem $(P)$ will be found through an application of Theorem 3.4. Here we assume in addition to $(Q),\left(A_{1}-A_{2}\right)$, that
$\left(A_{3}\right) \quad$ for all compact $\mathscr{K} \subset \mathbb{R}^{N} \backslash \Omega_{0}$ there exists $g_{0}=g_{0}(\mathscr{K})$, such that $\inf \{g(x): x \in \mathscr{K}\} \geqq g_{0}>0 ;$
$\left(A_{4}\right) \quad \Lambda<a-g C_{0}$,
where according to the notation introduced in Sect. $2, g=\sup _{\mathbb{R}^{N}} g(x)$. We set

$$
l=\min \left\{\lambda_{0}, a-g C_{0}\right\}
$$

Let us point out that $\left(A_{4}\right)$ implies $\Lambda<l$, because $\Lambda<\lambda_{0}$.
Our goal is to show that assumption $(\widetilde{B})$ holds for any $b<l$.
a) Construction of $\Psi$. We assume that $\Omega_{0}$ is not emptry: if not, the first part of this subsection can be avoided.

Let $b<l \leqq a-g C_{0}$, and let $\Omega_{\delta}$ denote the $\delta$-neighborhood of $\Omega_{0}$. Taking $\delta>0$ sufficiently small, there results

$$
\begin{equation*}
b<\lambda_{\Omega_{\delta}}<\lambda_{0} . \tag{4.1}
\end{equation*}
$$

Fix $a^{\prime}$, with $b+g C_{0}<a^{\prime}<a$, and let $\beta(x)$ denote the function defined by setting

$$
\beta(x):=\left(q(x)-a^{\prime}\right)^{-} .
$$

Since $a^{\prime}<a$, it follows that the support $K=\operatorname{supp}(\beta)$ of $\beta$ is compact. Letting $\rho>0$ be such that

$$
K \cup \Omega_{\delta} \subset B_{\rho}
$$

we set $\gamma_{\alpha}(x)=\gamma_{\alpha}(|x|), \alpha>0$, where

$$
\gamma_{\alpha}(t)= \begin{cases}-\alpha & \text { if } t \leqq \rho \\ \alpha(t-\rho-1) & \text { if } \rho<t<\rho+1, \\ 0 & \text { if } t \geqq \rho+1\end{cases}
$$

Consider the linear Schrödinger equation

$$
\begin{equation*}
-\Delta u+\gamma_{\alpha}(x) u=\mu u, \quad x \in \mathbb{R}^{N} \tag{4.2}
\end{equation*}
$$

and let

$$
\mu_{\alpha}=\inf _{u \in H,|u|_{2}=1} \int_{\mathbb{R}^{N}}\left[|\nabla u|^{2}+\gamma_{\alpha}(x) u^{2}\right] .
$$

The properties of $\mu_{\alpha}$ are collected in the following Lemma.
Lemma 4.1 (i) $\mu_{\alpha} \leqq 0$ and there exists $\alpha^{*} \geqq 0$ such that $\mu_{\alpha}<\mu_{\alpha^{*}}=0$ provided $\alpha>\alpha^{*}$;
(ii) if $\alpha>\alpha^{*}$ then $\mu_{\alpha}$ is the principal eigenvalue of (4.2) with corresponding positive eigenfunction $\varphi_{\alpha} \in W^{2, p}\left(\mathbb{R}^{N}\right)$ for all $p>1$.
(iii) $\mu_{\alpha}$ depends continuously on $\alpha$;

Proof. Since $\gamma_{\alpha}(x) \leqq 0$ it follows that

$$
\mu_{\alpha} \leqq \inf _{|u|_{2}=1}|\nabla u|_{2}^{2}=0
$$

Moreover, $\mu_{\alpha}<0$ provided $\alpha$ is large enough (for example, $\alpha>\lambda_{B_{\rho}}[0]$ suffices). Property (ii) is well known (see, for example, [3] or [12]) and (iii) is standard.

From Lemma 4.1 it follows that there exists $\alpha_{0}>0$ such that for $\mu_{0}:=\mu_{\alpha_{0}}$ one has

$$
\begin{equation*}
b-a^{\prime}+g C_{0}<\mu_{0}<0 \tag{4.3}
\end{equation*}
$$

Let $\psi=\psi_{b}: \mathbb{R}^{N} \rightarrow \mathbb{R}$ be any strictly positive $C^{2}$-function such that

$$
\psi(x)= \begin{cases}\phi_{\delta}(x) & x \in \Omega_{\delta / 2} \\ \varphi_{0}(x) & x \in \mathbb{R}^{N} \backslash B_{\rho}\end{cases}
$$

where $\varphi_{0}=\varphi_{\alpha_{0}}$ and $\phi_{\delta} \in W_{0}^{1,2}\left(\Omega_{\delta}\right)$ satisfies

$$
-\Delta \phi_{\delta}+q(x) \phi_{\delta}=\lambda_{\Omega_{\delta}} \phi_{\delta}, \quad x \in \Omega_{\delta}
$$

Finally we define $\Psi(x)=\Psi_{b}(x):=C \psi(x)$.
b) A priori bounds. We first show

Lemma 4.2 There exists $C>0$ such that $\Psi=C \psi$ is a supersolution for $\left(P_{R}\right)$ for all $R>0$ and all $\lambda \leqq b$.

Proof. Taking $C$ large enough, one has that $h(\Psi(x))<0$ for all $x \in \Omega_{\delta / 2}$. Then, the definition of $\psi,(4.1)$ and $\lambda \leqq b$ imply

$$
-\Delta \Psi+q(x) \Psi=\lambda_{\Omega_{\delta}} \Psi>b \Psi \geqq \lambda \Psi+g(x) h(\Psi) \Psi, \quad x \in \Omega_{\delta / 2}
$$

Let

$$
M=\sup \left\{\Delta \psi-q(x) \psi+b \psi: x \in B_{\rho+1} \backslash \Omega_{\delta / 2}\right\}
$$

Since $\psi(x) \geqq c_{1}>0$ on $B_{\rho+1} \backslash \Omega_{\delta / 2},\left(A_{2}-A_{3}\right)$ allow us to take $C$ in such a way that

$$
c_{1} g(x) h(C \psi(x)) \leqq-M, \quad \forall x \in B_{\rho+1} \backslash \Omega_{\delta / 2}
$$

whence

$$
-\Delta \Psi+q(x) \Psi \geqq b \Psi-M C \geqq \lambda \Psi+g(x) h(\Psi) \Psi, \quad \forall x \in B_{\rho+1} \backslash \Omega_{\delta / 2}
$$

Finally, for $|x| \geqq \rho+1, \gamma_{\alpha} \equiv 0$ and $\psi=\varphi_{0}$ satisfies $-\Delta \psi=\mu_{0} \psi$. Using (4.3) and since $\operatorname{supp}(\beta) \subset B_{\rho}$, it follows that

$$
\begin{aligned}
-\Delta \Psi+q(x) \Psi & =\left(\mu_{0}+q(x)\right) \Psi \geqq\left(\mu_{0}+a^{\prime}\right) \Psi \\
& \geqq\left(b+g C_{0}\right) \Psi \geqq \lambda \Psi+g(x) h(\Psi) \Psi
\end{aligned}
$$

Let us remark that if $\Omega_{0}=\emptyset$ one can simply define $\Psi(x)=C \varphi_{0}(x)$ and the Lemma follows from the preceding inequality.

Lemma 4.3 For all $b<l$ condition $(\widetilde{B})$ holds with the function $\Psi$ defined above.

Proof. We use the notation introduced in Sect. 3. For a fixed $n$, we set $B_{n}=B_{R_{n}}$. Let $K>0$ be such that $F_{\lambda}(x, s)+K s$ (see (3.7) for the definition) is strictly increasing with respect to $s \in\left[0, \max _{B_{n}} \Psi\right]$, for all $\lambda \in[\lambda, b]$. Let $v_{n}$ denote the solution of the b.v.p.

$$
\begin{cases}-\Delta v_{n}+K v_{n}=F_{b}(x, \Psi)+K \Psi, & |x|<R_{n} \\ v_{n}(x)=0, & |x|=R_{n}\end{cases}
$$

Since $F_{b}(x, \Psi)+K \Psi \geqq 0$, the maximum principle implies that $v_{n}$ is strictly positive in $B_{n}$, and has negative outer normal derivative at every point on $\partial B_{n}$. This means that $v_{n}$ lies in the interior of the cone, $\mathscr{P}$, of positive functions of $C_{0}^{1}\left(B_{n}\right)$. Moreover, by Lemma 4.2 it follows that $-\Delta \Psi \geqq F_{b}(x, \Psi)$, whence

$$
-\Delta\left(\Psi-v_{n}\right)+K\left(\Psi-v_{n}\right) \geqq F_{b}(x, \Psi)+K \Psi-F_{b}(x, \Psi)-K \Psi=0, \quad|x|<R_{n}
$$

Since, for $|x|=R_{n}$ one has $\Psi(x)-v_{n}(x)=\Psi(x)>0$, the maximum principle again implies

$$
0<v_{n}(x)<\Psi(x), \quad \forall|x|<R .
$$

Furthermore, $F_{b}(x, \cdot)+K$ increasing and $F_{b}(x, s) \geqq F_{\lambda}(x, s)$ for all $b \geqq \lambda$, yield

$$
-\Delta v_{n}>F_{b}\left(x, v_{n}\right) \geqq F_{\lambda}\left(x, v_{n}\right), \quad \forall|x|<R .
$$

In particular, $v_{n} \in C_{0}^{1}\left(B_{n}\right)$ is a super-solution (but not a solution) of $\left(P_{R_{n}}\right)$ for every $\lambda \leqq b$. To conclude the proof it remains to show that $v_{n}>u$ whenever $(\lambda, u) \in X_{n}$. Consider the set $\Sigma_{n}=\left\{\left(\lambda, v_{n}-u\right):(\lambda, u) \in X_{n}\right\}$. Let us explicitely point out that $\Sigma_{n}$ is connected because $X_{n}$ is. Moreover observe that $\left(\lambda_{R_{n}}, v_{n}\right) \in \Sigma_{n}$ and therefore $\Sigma_{n} \cap(T \times \mathscr{P}) \neq \emptyset$. We claim that $\Sigma_{n} \subset T \times \mathscr{P}$. If not, the connection of $\Sigma_{n}$ allows us to find $(\lambda, u) \in X_{n}$ such that $v_{n}-u \in \partial \mathscr{P}$. In particular, since $v_{n}$ is not a solution of $\left(P_{R_{n}}\right)$, one infers that $v_{n} \geqq(\not \equiv) u$ in $B_{n}$. Then

$$
-\Delta\left(v_{n}-u\right)+K\left(v_{n}-u\right) \geqq F_{\lambda}\left(x, v_{n}\right)+K v_{n}-F_{\lambda}(x, u)-K u \geqq 0, \quad \forall|x|<R_{n}
$$

By the strong maximum principle, we deduce that $v_{n}-u \in \mathscr{\mathscr { P }}$, which is a contradiction. Then $\forall(\lambda, u) \in X_{n}$ one has

$$
0<u(x)<v_{n}(x)<\Psi(x), \quad \forall|x|<R_{n} .
$$

This completes the proof.
c) Global bifurcation. We are now in position to prove:

Theorem 4.4 Suppose that $(Q),\left(A_{1}-A_{2}-A_{3}-A_{4}\right)$ hold. Then there exist a branch, $S$, of positive solutions of $(P)$ bifurcating from $(\Lambda, 0)$, such that $\Pi(S)=\left[\Lambda^{*}, l\left[\right.\right.$, for some $\Lambda^{*} \leqq \Lambda$.

Moreover, if $\Omega_{0} \neq \emptyset$ and
$\left(A_{5}\right) \quad \lambda_{0}<a-g C_{0}$,
then $\Pi\left(S=\left[\Lambda^{*}, \lambda_{0}[\right.\right.$.
Proof. By Lemma $4.3(B)$ holds for all $b<l$. Then an application of Theorem 3.4 yields the existence of a bifurcation branch $S_{b}$ of positive solutions of $(P)$ such that $\Pi\left(S_{b}\right)=\left[\Lambda^{*}, b\right]$. Taking $S=\bigcup_{b} S_{b}$ it follows that $\Pi(S)=\left[\Lambda^{*}, l[\right.$.

Finally, to prove the last statement, it suffices to remark that $(P)$ has no positive solution for $\lambda \geqq \lambda_{0}$. Otherwise, if $u_{0}$ is a positive solution of $(P)$ for some $\lambda \geqq \lambda_{0}$, such a $u_{0}$ is a supersolution of $\left(P_{\Omega}\right)$, for such $\lambda$, for any bounded domain $\Omega$ with $\Omega_{0} \subset \Omega$. Since $\varepsilon \phi_{0}$ (see notation introduced in Sect. 2) is a subsolution for any $\varepsilon>0$ small, it follows that $\left(P_{\Omega}\right)$ has a positive solution for that $\lambda \geqq \lambda_{0}$, in contradiction with Lemma 2.2.

Remarks 4.5 (i) Since $\Lambda<\lambda_{0}$ (see Sect. 3), condition $\left(A_{5}\right)$ implies $\left(A_{4}\right)$.
(ii) When $\Omega_{0}=\emptyset$ Theorem 4.4 improves the existence results of [7], where, in addition to $\left(A_{1}-A_{2}-A_{4}\right)$ and to some regularity and growth restriction on $q$ and $g$, it is assumed that $g(x)>0$ on $\mathbb{R}^{N}$ and that $h(s)$ is decreasing. Of course, in this case $\left(A_{4}\right)$ is nothing but $\Lambda<a$.
(iii) Let us point out that $\left(A_{4}\right)$ is, in general, a necessary condition for the existence of a positive solution of $(P)$. For example, see [7], if $g(x)>0, h(s)$ is decreasing and $q(x) \rightarrow a$ as $|x| \rightarrow \infty$, then $(P)$ has no positive solution if $\lambda \geqq a$.

A case in which $\left(A_{4}\right)$ holds is, for example, when there exists $R>0$ such that $q(x) \leqq-\lambda_{B_{R}}[0]+a-g C_{0}$ in $B_{R}$.
(iv) Theorem 4.4 applies to the specific problem (1.1) discussed in the Introduction. In such a case one immediately finds that $\left(A_{5}\right)$ holds provided $-\alpha_{1}+\pi^{2} / 4 d^{2}<-\alpha_{2}$.

## 5 Blow up

In this section we will study the behaviour of $S$ when $\lambda \uparrow \lambda_{0}$. Here we deal with the case $\Omega_{0} \neq \emptyset$, and assume that ( $A_{5}$ ) holds.

Consider $\widehat{h}: \mathbb{R}^{+} \rightarrow \mathbb{R}$, a strictly decreasing function satisfying $\left(A_{2}\right)$ and such that

$$
\widehat{h}(s)<h(s), \quad \forall s>0
$$

Let $k: \mathbb{R}^{+} \rightarrow \mathbb{R}$ be a strictly decreasing function such that

$$
k(s)>h(s), \quad \forall s \geqq 0
$$

Setting $\widetilde{h}(s)=k(s)-k(0)$, suppose that $k$ has been chosen in such a way that $\widetilde{h}$ satisfies $\left(A_{2}\right)$.

Denote by $\left(\widehat{P}_{R}\right)$, resp. $(P)$, the problem $\left(P_{R}\right)$, resp. $(P)$, with $\widehat{h}$ instead of $h$, and moreover, setting

$$
\widetilde{q}(x)=q(x)-g(x) k(0), \quad \forall x \in \mathbb{R}^{N},
$$

we denote by $\left(\widetilde{P}_{R}\right)$, resp. $(\widetilde{P})$, the problem $\left(P_{R}\right)$, resp. $(P)$, with $\widetilde{h}$ and $\widetilde{q}$ instead of $h$ and $q$. We also set $\widetilde{a}=\lim \inf \widetilde{q}(x)$ and $\widetilde{\Lambda}=\Lambda[\widetilde{q}]$.

Remark that $\lambda_{\Omega_{0}}[\widetilde{q}]=\lambda_{0}$, because $\widetilde{q}(x)=q(x)$, for all $x \in \Omega_{0}$. Moreover, since $k(0)$ can be taken sufficiently close to $C_{0}$, it follows that

$$
\begin{equation*}
\tilde{\Lambda}<\lambda_{0}<\widetilde{a} \tag{5.1}
\end{equation*}
$$

We denote by $\widehat{S}_{R}, \widetilde{S}_{R}$, the branches of positive solutions of $\left(\widehat{P}_{R}\right)$, and $\left(\widetilde{P}_{R}\right)$, found in Theorem 2.6. By Remark 2.7(ii), for all $\lambda \in \Pi\left(\widehat{S}_{R}\right) \cap \Pi\left(\widetilde{S}_{R}\right)$, there exists a unique $\widehat{u}_{\lambda, R}$ such that $\left(\lambda, \widehat{u}_{\lambda, R}\right) \in \widehat{S}_{R}$, and a unique $\widetilde{u}_{\lambda, R}$ such that $\left(\lambda, \widetilde{u}_{\lambda, R}\right) \in \widetilde{S}_{R}$.

According to (5.1) Theorem 4.4 applies to problem $(\widetilde{P})$ and yields a branch of positive solutions $\widetilde{S}$.

Roughly, we will show that $S$ 'lies between' the branches $\widehat{S}_{R_{0}}$ and $\widetilde{S}$ in such a way that the blow up of the latters will imply that of $S$.
Lemma 5.1 Fixed $R_{0}>0$, with $B_{R_{0}} \supset \Omega_{0}$, then $\forall R \geqq R_{0}$, there results

$$
\begin{gather*}
\widehat{u}_{\lambda, R_{0}} \leqq u_{R}, \quad \forall\left(\lambda, u_{R}\right) \in S_{R}, \quad \lambda>\lambda_{R_{0}}  \tag{5.2}\\
u_{R} \leqq \widetilde{u}_{\lambda, R} \quad \forall\left(\lambda, u_{R}\right) \in S_{R} \tag{5.3}
\end{gather*}
$$

Proof. Since $\widehat{h}<h$ and $R_{0}<R$, it follows that $u_{R}$ is a super-solution for problem $\left(\widehat{P}_{R_{0}}\right)$. As usual, a small sub-solution of $\left(\widehat{P}_{R_{0}}\right)$ can be found as before, and thus the unique solution $\widehat{u}_{\lambda, R_{0}}$ of ( $\widehat{P}_{R_{0}}$ ) satisfies $\widehat{u}_{\lambda, R_{0}} \leqq u_{R}$.

Similarly, it is easy to check that $u_{R}$ is a sub-solution of $\left(\widetilde{P}_{R}\right)$, and $\Psi=C \psi$ is a super-solution of ( $\widetilde{P}_{R}$ ) (provided that $C$ is sufficiently large). Hence $u_{R} \leqq \widetilde{u}_{\lambda, R}$.

Theorem 5.2 Suppose $(Q)$ and $\left(A_{1}-A_{2}-A_{3}-A_{5}\right)$ hold. Then, on $S$ one has that for all $p \geqq 1,|u|_{p} \rightarrow \infty$, iff $\lambda \uparrow \lambda_{0}$.

Proof. (a) Let $\lambda \uparrow \lambda_{0}$. By (h) of Theorem 2.6, $\left|\widehat{u}_{\lambda, R_{0}}\right|_{p} \rightarrow \infty$. Using (5.2) it follows that $|u|_{p} \rightarrow \infty$ as $\lambda \uparrow \lambda_{0}$.
(b) Since $\left(\widetilde{P}_{R}\right)$ has a unique solution for all $\lambda \in \Pi\left(\widetilde{S}_{R}\right)$, the $\Psi$ found in Subsect. 4 b , satisfies not merely $(\widetilde{B})$ but also $(B)$ and therefore it becomes an a-priori bound for all the solutions of $(\widetilde{P})$, with $\lambda \leqq b$. This implies that $\widetilde{S}$ can only blow up in $L^{p}$ as $\lambda \rightarrow \lambda_{0}$. Taking limits in (5.3) one finds that $u \leqq \widetilde{u}$ for all $(\lambda, u) \in S$ and $(\lambda, \widetilde{u}) \in \widetilde{S}$. Then $|u|_{p} \rightarrow \infty$ on $S$ implies that $\lambda \rightarrow \lambda_{0}$.

Remark 5.3 When $\Omega_{0}=\emptyset$ the blow up of $S$ has been studied in Sects. 6 and 7 of [7]. It is worth pointing out that Theorem 5.2 is quite different from the results of [7]. Actually, in the latter the branch blows up at $\lambda=a$ in $L^{p}$, for a certain range of $p$, related to the dimension $N$ and to the asymptotic properties of $q, g$ and $h$. We also note that such a behaviour of $S$ remains valid also if $\Omega_{0} \neq \emptyset, \Lambda<a<\lambda_{0}$ and $q, h$ satisfy the same assumptions of [7]: it suffices to use, with minor changes, the arguments in [7].

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