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Development of ultra-pure NaI(Tl) detectors for the COSINE-200 experiment

B. J. Park^{1,2}, J. J. Choi³, J. S. Choe², O. Gileva², C. Ha^{4,a}, A. Iltis⁵, E. J. Jeon², D. Y. Kim², K. W. Kim², S. K. Kim³, Y. D. Kim^{1,2}, Y. J. Ko², C. H. Lee², H. S. Lee^{2,1,b}, I. S. Lee², M. H. Lee^{1,2}, S. H. Lee^{1,2}, S. J. Ra², J. K. Son², K. A. Shin²

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Abstract The annual modulation signal observed by the DAMA experiment is a long-standing question in the community of dark matter direct detection. This necessitates an independent verification of its existence using the same detection technique. The COSINE-100 experiment has been operating with 106 kg of low-background NaI(Tl) detectors providing interesting checks on the DAMA signal. However, due to higher backgrounds in the NaI(Tl) crystals used in COSINE-100 relative to those used for DAMA, it was difficult to reach final conclusions. Since the start of COSINE-100 data taking in 2016, we also have initiated a program to develop ultra-pure NaI(Tl) crystals for COSINE-200, the next phase of the experiment. The program includes efforts of raw powder purification, ultra-pure NaI(Tl) crystal growth, and detector assembly techniques. After extensive research and development of NaI(Tl) crystal growth, we have successfully grown a few small-size (0.61-0.78 kg) thalliumdoped crystals with high radio-purity. A high light yield has been achieved by improvements of our detector assembly technique. Here we report the ultra-pure NaI(Tl) detector developments at the Institute for Basic Science, Korea. The technique developed here will be applied to the production of NaI(Tl) detectors for the COSINE-200 experiment.

1 Introduction

Although numerous astronomical observations support the conclusion that most of the matter in the universe is invisible dark matter, an understanding of its nature and interactions remains elusive [1–3]. Dark matter phenomenon can be explained by new particles, such as weakly interacting massive particles (WIMPs) [4,5]. Even though tremendous efforts to search for evidence of WIMP dark matter by directly detecting nuclei recoiling from WIMP-nucleus interactions are being pursued, no definitive signal has been observed [6,7]. One exception is the DAMA experiment that, using an array of NaI(Tl) detectors [8,9], observes an annual modulation of event rates that can be interpreted as resulting from WIMP-nuclei interactions [10,11]. However, this result has been the subject of a continuing debate because the WIMP-nucleon cross sections inferred from the DAMA modulation are in conflict with limits from other experiments [12–23].

An unambiguous verification of the DAMA signal by independent experiments using the same NaI(Tl) crystal target material is mandatory. Experimental efforts by several groups using NaI(Tl) as the target medium are currently underway [24–30]. Even though interesting physics results from COSINE-100 [31,32] and ANAIS-112 [33] have been reported, none of these efforts have yet achieved a background level similar or lower than that of the DAMA experiment. This large background is a strong impediment to the establishment of any unambiguous conclusion of the DAMA's observation.

From our COSINE-100 experience, we have identified the crystal growth as the key component for reducing the overall background, of which radioactive ⁴⁰K and ²¹⁰Pb contaminants are the most troublesome [29,34,35]. The COSINE-100 detector is immersed in an active liquid scintillator veto system that reduces the impact of ⁴⁰K impurities [26,36,37] but ²¹⁰Pb has to be eliminated at the crystal production stage.



¹ IBS School, University of Science and Technology (UST), Daejeon 34113, Republic of Korea

² Center for Underground Physics, Institute for Basic Science (IBS), Daejeon 34126, Republic of Korea

³ Department of Physics and Astronomy, Seoul National University, Seoul 08826, Republic of Korea

⁴ Department of Physics, Chung-Ang University, Seoul 06973, Republic of Korea

⁵ Damavan Imaging, 10430 Troyes, France

^a e-mail: changhyon.ha@gmail.com (corresponding author)

^b e-mail: hyunsulee@ibs.re.kr (corresponding author)

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Fig. 1 Two dedicated Kyropoulos growers for the NaI(Tl) crystals. **a** A small-sized grower for low-background NaI(Tl) R&D and **b** a full-sized grower for 120 kg size NaI(Tl) crystal production. Both growers have an accompanying setup to facilitate hot crystal annealing

The PICO-LON experiment made a crystal with a ^{210}Pb contamination as low as $29\pm7~\mu Bq/kg$ by using a cation resin-based NaI powder purification system, but had difficulties reproducing additional crystals with similar quality [27,38]. The SABRE experiment recently produced a NaI(Tl) crystal with ^{210}Pb and ^{40}K contaminations of 0.34 \pm 0.04 mBq/kg and 4.3 \pm 0.2 ppb [30], respectively.

As part of an effort to upgrade the on-going COSINE-100 experiment to the next-phase COSINE-200 experiment, we have conducted a R&D aimed at low-background NaI(Tl) crystal growth. An initial chemical purification stage based on a recrystallization method demonstrated sufficient reduction of K and Pb in the powder [39]. Based on this, a mass production facility for the powder purification was assembled in our institute [40]. Two custom-made Kyropoulos machines were implemented at the same site for both the demonstration of low-background crystal growth and the timely production of 200 kg of NaI(Tl) crystals [41]. In the meantime, we developed detector assembly techniques for low-background NaI(Tl) detectors. In this article, our progress in the development of low-background NaI(Tl) detectors and the prospects of the COSINE-200 experiment is described.

2 Crystal growth

We have built two dedicated Kyropoulos growers for NaI(Tl) crystals. A small R&D grower that is used for proof of principle measurements is shown in Fig. 1a. This grower is specifically designed to grow pilot NaI(Tl) crystals with relatively large temperature gradient and fast ingot rotation. The typical growth time is approximately 24 h for a 70 mm diameter and 80 mm high ingot. Because of this short crystal growth time, we can minimize radon contamination from emanations from the grower's materials and the ambient air.

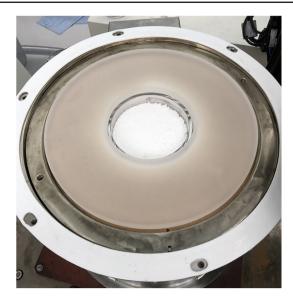


Fig. 2 A top view of the grower with NaI powder loaded in the middle crucible, with the chamber cupola removed. A hollow-disk-shape quartz lid covers the top of the refractories to prevent fumes from coming out when heated

For the small grower, we use a 12 cm diameter, 10 cm high quartz crucible that is surrounded by heaters and refractories inside the chamber. The chamber is cylindrical with dimensions of 34 cm diameter and 50 cm height. During crystal growth, N₂ gas was continuously flushed via a thallium trap with a flow rate of 10 l/min. A same-size annealing machine is included as shown in Fig. 1a. Figure 2 shows the top view of the growing setup before the start the growth process. The quartz crucible can contain up to 1.7 kg of the NaI powder and can produce a 1.1 kg crystal ingot with typical dimensions of: diameter 60–80 mm height 50–80 mm as shown in Fig. 3.

With this system, we successfully produced several thallium-doped NaI(Tl) crystals with low-background NaI powder. Based on this success, we constructed a full-size grower that accommodates a 61 cm diameter and 44 cm height quartz crucible for the growth of approximately 120 kg crystal ingots (see Fig. 1b). The conceptual design of the full-size grower is similar to that of the smaller one, in order to minimize differences in growth conditions in terms of internal materials and thermal gradient. We expect approximately 3–4 days crystal growth times for 120 kg ingots (40 cm diameter and 25 cm height).

All thallium-doped low-background crystals shown in Fig. 3 start with low-background NaI powder from the Sigma-Aldrich (SA) company named "Astro-Grade". Previous studies [26] show that this grade powder contains low potassium (less than 10 ppb level) and lead (a few ppb level). Typically, the NaI powder from this vendor does not contain significant amounts of uranium and thorium.



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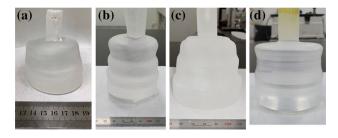


Fig. 3 Four NaI(Tl) crystal ingots produced by the small grower: **a** NaI-025, **b** NaI-034, **c** NaI-035, and **d** NaI-036

When the crystal growth is finished, sample pieces from the crystal ingot are measured with inductively coupled plasma mass spectrometry (ICP-MS) to study impurity contaminations. The first grown crystal, NaI-025, was highly contaminated with both K and Pb, which was found to originate from fumes produced by the refractories as shown in Table 1. The refractories were replaced with ones fabricated from relatively radiopure materials, and this resulted in significantly reduced impurities in the fumes. In addition, the implementation of a quartz plate that covers the top of the refractories prevented further diffusion of the fumes into the NaI crucible (see Fig. 2). Subsequently grown crystals, NaI-034, NaI-035, and NaI-036, had significantly reduced K and Pb contaminations, as can be seen in Table 2.

We tested different types and amounts of thallium iodide (TII) from different vendors as the dopant in the NaI(Tl) crystals. Initially, we used 99.999% purity TII powder from SA. However, this powder created black dust in the melt that was traced to carbon impurities in the powder. This dust made it difficult to optimize crystal growth conditions and the produced ingots contained black speckles throughout their entire bulk. We speculate that the carbon contamination in the TII powder may have been caused by plastic or polytetrafluoroethylene (PTFE) material used in the chemical purification process. After a variety of tests for different types of TII, we found that the bead type TII does not produce any significant black dust either in the melt or the ingot. To confirm the reproducibility of this, we also tested TII from Alpha-Aesar (AA) in both powder and bead forms and concluded that the bead type contains less carbon. Contamination levels

Table 1 ICP-MS results for: initial materials; fumes produced during the growing process

Sample	K (ppb)	Pb (ppb)	U (ppt)	Th (ppt)
NaI (Astro-grade)	5	1.3	< 10	< 10
TlI (SA powder)	3400	500	< 10,000	< 10,000
TlI (SA bead)	150	100	< 10,000	< 10,000
TlI (AA powder)	2100	600	< 10,000	< 10,000
TlI (AA bead)	4400	600	< 20,000	< 10,000
Fumes with old refractories	1,021,235	10,407	< 10	< 10
Fumes with new refractories	320	25	<10	< 10

of radioisotopes for various TII samples were measured with ICP-MS (see Table 1).

Additionally, since the NaI(Tl) melt directly touches the crucibles, we checked the impact of the crucible type by exchanging the quartz crucible with a platinum one. The platinum crucible provided a lower contamination of K and Pb although the crystals produced with quartz crucibles also had ⁴⁰K and ²¹⁰Pb levels that are sufficiently low for the COSINE-200 experiment. Specifications of the crystal growth conditions and ICP-MS measurement results for various NaI(Tl) crystals grown in the laboratory are summarized in Table 2.

3 Experimental setup

3.1 NaI(Tl) crystal detectors

Once the crystal ingots are prepared, samples from the top and bottom sections were cut out with a bandsaw for ICP-MS measurements. The remaining bulk of the crystal was machined to a cylindrical shape using a lathe in a low-humidity, N₂ gas flushed glovebox shown in Fig. 4. While shaping the crystal, mineral oil was continuously poured on it to prevent cracks from developing and suppress NaI(Tl) dust in the environment. Four separate ingots were assembled as detectors after the crystal surfaces were polished, also in the low-humidity, N₂ gas flushed glovebox. For NaI-025 and NaI-036, we did not machine the barrel surfaces, while NaI-034 and NaI-035 were machined into cylindrical forms. Studies show that the machined cylindrical surface improves light collection while not affecting the background contamination level.

The crystal was dry-polished using aluminum oxide lapping films ranging from 400 to 8000 grits. After the polishing, several layers of soft PTFE sheets were wrapped on the barrel surface as a diffusive reflector and an optical interface (silicon rubber EJ-560) was coupled to the top and bottom surfaces. A 5 mm-thick cylindrical copper casing with quartz windows at each end is used to hermetically encapsulate the crystal as shown in Fig. 5. Hamamatsu 3-inch PMTs (R12669SEL)



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Table 2 The top six entries list growth conditions of each crystal ingot, where SA and AA indicate Sigma-Aldrich and Alpha-Aesar, respectively. The middle two entries list the masses of the crystal ingots and

polished detectors. The lower five entries list the ICP-MS results of the grown crystals, where top and bottom measurements are averaged and uncertainties are half the differences of the two measurements

Crystal	NaI-025	NaI-034	NaI-035	NaI-036
Growth date	May/24/2018	Sep/18/2019	Nov/12/2019	Feb/4/2020
Quartz cover	No	Yes	Yes	Yes
Refractory	Old	New	New	New
Crucible	Quartz	Platinum	Quartz	Quartz
TlI loaded (mol%)	0.1	0.1	0.12	1.0
TlI powder	SA powder	SA powder	AA bead	SA bead
Mass (ingot) (kg)	0.79	1.18	1.16	1.08
Mass (detector) (kg)	0.65	0.67	0.61	0.78
K (ppb)	740 ± 68	8±3	13±6	23±9
Pb (ppb)	170±15	0.5 ± 0.1	2±1	10±1
U (ppt)	< 10	< 10	< 10	< 10
Th (ppt)	< 10	< 10	< 10	< 10
Tl (mol %)	0.008 ± 0.001	0.015 ± 0.001	0.023 ± 0.001	0.072 ± 0.003

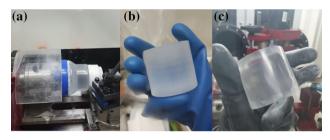


Fig. 4 NaI-035 crystal machining and polishing. **a** The crystal is machined into a cylindrical form using a lathe. The finished product before (**b**) and after (**c**) the surface polishing is shown

are coupled via an optical interface to each quartz window. The PMTs are fixed with PTFE guard rings by means of stainless-steel bolts.

All assembly procedures were performed inside the glovebox in which the humidity level was maintained to be less than a few tens of ppm (H_2O) using N_2 gas and a molecular sieve trap. All assembly parts were cleaned using diluted Citranox liquid with sonication, and baked in an oven before they were moved into the glovebox. Several weeks after the assembly was finished, a visual inspection showed no degradation of the crystal in terms of its transparency. For evaluation, we used an underground low-background setup for further detailed measurements.

3.2 Shielding structure

To evaluate the background contamination levels in the NaI(Tl) crystals, we used the same experimental apparatus that was used for previous NaI(Tl) crystal R&D at the Yangyang Underground Laboratory (Y2L) [24,34]. This includes a 12-module array of CsI(Tl) crystals inside a shield that is comprised of 10 cm of copper, 5 cm of polyethylene,

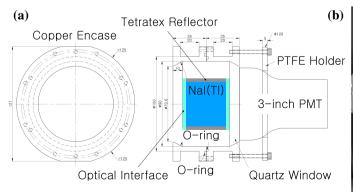




Fig. 5 NaI-035 detector. The diagram a shows the encapsulation design and the photo b shows the final detector



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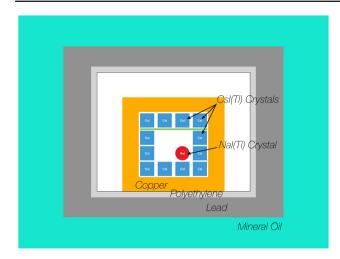


Fig. 6 A schematic view of the Y2L R&D setup. The NaI(Tl) crystal (circle) is placed within the CsI(Tl) crystal array (squares) to tag external and coincident components of the testing crystal. An acrylic table (green solid line) supports top-side CsI(Tl) crystals. These were installed inside a shield. Each NaI(Tl) crystal was tested one by one in this setup

15 cm of lead, and 30 cm of liquid-scintillator-loaded mineral oil [42,43] as shown in Fig. 6. The shielding stops external neutrons and gamma rays, and vetoes cosmic-ray muons. The NaI(Tl) crystals were tested inside the CsI(Tl) detector array with an arrangement indicated schematically in Fig. 6.

3.3 Electronics

Signals from PMTs attached at each end of the NaI(Tl) crystal are amplified by custom made preamplifiers: the high-gain anode signal by a factor of 30 and the low-gain signal from the 5th stage dynode by a factor of 100. The amplified signals are digitized by a 500 MHz, 12-bit flash analog-to-digital converters (FADC). The FADC module has four channels that accommodates one NaI(Tl) crystal's data. Triggers from individual channels are generated by field programmable gate arrays (FPGAs) embedded in the FADC. The final trigger decision for an event is made by a trigger and clock board (TCB) that also synchronizes the timing of the different modules. A trigger is generated when an anode signal corresponding to one or more photoelectrons (PEs) occur in each PMT within a 200 ns time window. Both the low-gain dynode and high-gain anode signals are recorded whenever the anode signals produce a trigger. The dynode signals do not produce triggers. The dynamic range of the anode signals is up to 100 keV while that for the dynode signals is up to 5000 keV.

Signals from the CsI(Tl) crystals are amplified by a gain of 10 and digitized in a charge-sensitive 62.5 MHz ADC (SADC). The SADC returns the integrated charge and the time of the signal. We use a 2048 ns integration time to contain CsI(Tl) signals considering their decay times. The SADC channels do not generate triggers.

If the NaI(Tl) crystal satisfies the trigger condition, the TCB transmits trigger signals to the FADC and the SADC. Data of the NaI(Tl) crystal and the CsI(Tl) crystals are stored for the FADC trigger. For the FADC channel, this corresponds to an 8 μ s-long waveform staring approximately 2.4 μ s before the time of the trigger position. For each SADC channel this corresponds to the maximum integrated charge within an 8 μ s search window and to the associated time of that maximum. This system has been used in COSINE-100 experiment and descriptions are provided elsewhere [26,44].

3.4 Energy calibration and light yields

The low energy anode signals of the NaI(Tl) crystals are calibrated with 59.54 keV γ -rays emitted from an 241 Am source. The high energy dynode signals are calibrated with 609 keV (214 Bi) and 1460 keV (40 K) lines that are produced by contaminants in the crystal.

The charge distribution of single PE signals is obtained by identifying isolated clusters in the decay tail of the 59.54 keV γ -ray signal (2–4 μ s after the signal start). From this, the light yield is determined from the ratio of the total deposited charge to the single PE's mean charge, scaled to 59.54 keV.

A high NaI(Tl) crystal light yield is crucial for enabling the COSINE-200 experiment to operate with a 1 keV energy threshold [45] that matches the DAMA experiment's conditions [9]. The doping concentration of thallium is one of the most important factors that determine the light yield, as described in Refs. [46,47]. We originally targeted a doping level of 0.1 mol% for an optimal light yield based on the above-noted references. To accomplish this, we applied approximately 0.1 mol% TII powder into the initial mix for the NaI-025, NaI-034, and NaI-035 crystals. However, the measured doping levels in the final produced crystals were significantly reduced to 0.01–0.02 mol% (Table 1). This was mainly attributable to the evaporation of a large fraction of the thallium during the high-temperature growth period. We measured relatively low light yields from these crystals; between 9 and 12 number of PEs (NPEs) per keV. Because of light absorption by the black impurities inside the crystals, slightly lower light yields in the crystals grown with TII powder (NaI-025 and NaI-034) were observed (see Table 3).

For NaI-036, 1 mol% of bead-type TII from SA was added to the initial mix and a 0.07 mol% Tl concentration was measured in the ingot after growth. In this crystal, we obtain the highest light yield of 17.1±0.5 NPE/keV. Further increases in the light yield are possible with the improved encapsulation scheme that is described in Ref. [48]. Note that the COSINE-100 crystals and DAMA crystals have light yields of approximately 15 NPE/keV and 5–10 NPE/keV, respectively, and are able to operate with a 1 keV threshold [9,45]. If the optical quality in crystals from the full-sized grower can match that of NaI-036, the light yield will be sufficient



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Table 3 Measured results of detector masses, light yield (LY), and the internal contaminations from this study compared with C6 of COSINE-100 [35], DAMA crystals [49,50], and SABRE result [30]

Crystal		NaI-025	NaI-034	NaI-035	NaI-036	COSINE-100	DAMA	SABRE
Mass (k	g)	0.65	0.67	0.61	0.78	12.5	9.7	3.4
LY (NPI	E/keV)	10.4 ± 1.6	9.5±1.1	11.8 ± 1.8	17.1 ± 0.5	14.6 ± 1.5	5-10	16.4 ± 0.3
nat K (pp	ob)	684 ± 100	< 62	< 42	< 53	17±3	<20	4.3 ± 0.2
²¹⁰ Pb (n	nBq/kg)	3.8 ± 0.3	0.05 ± 0.09	0.01 ± 0.02	0.42 ± 0.27	1.87 ± 0.09	0.01-0.03	0.34 ± 0.04
^{238}U	²¹⁴ Po (μBq/kg)	26±7	46±7	13±6	451±48	< 0.25	_	_
	²¹⁸ Po (μBq/kg)	_	64±11	8±6	_	_	_	_
	Average (µBq/kg)	26±7	51±7	11±4	451±48	< 0.25	8.7-124	_
²³² Th	216 Po (μ Bq/kg)	< 6	35±5	7±2	< 20	2.5±0.8	2–31	_

for operation of the COSINE-200 experiment with a 1 keV energy threshold.

3.5 PMT noise background and low-energy event selection

Photomultiplier tubes generate low-energy noise signals that are produced by a variety of mechanisms [24,49]. The DAMA group reported a signal selection criterion for efficiently removing PMT noise that exploits the fast decay time characteristics of PMT noise pulses. The ratio of "fast" charge (0-50 ns) and "slow" charge (100-600 ns) is used as a single parameter to discriminate fast PMT-induced noise pulse from genuine scintillation-induced signals [49,50]. In addition, our COSINE-100 experience showed that further selections are required based on the asymmetry between the two PMT signals and the average charge of the isolated clusters that comprise each signal [24,26]. A further optimized selection using a multivariable machine-learning algorithm, the boosted decision trees [51], was developed for the COSINE-100 experiment [31,45]. However, we only applied three parameter-based selection criteria discussed in Refs. [24,26] because the short-term measurements of the R&D data were not sufficient to produce adequate training samples.

4 Background understanding

4.1 Internal natural backgrounds

To produce ultra-low-background NaI(Tl) crystals, one should understand the internal contamination of natural radioisotopes. Contamination levels for various radioisotopes are described for each of the four crystals below. Table 3 shows a summary of the internal background measurements that are discussed in this section. The results are compared with the DAMA [49,50], and the SABRE [30] crystals, and crystal-6 (C6) of the COSINE-100 experiment [35]. (C6 has

the lowest background level among all the COSINE-100 crystals.)

4.1.1 40K background

One of the most troublesome sources of background for WIMP searches with NaI(Tl) crystals is contamination from 40 K. Its natural abundance is roughly 0.012 % of the total amount of potassium (nat K). About 10 % of 40 K decays produce a 1460 keV γ -ray in coincidence with a 3.2 keV X-ray. If the 1460 keV γ -ray escapes the crystal and only the 3.2 keV X-ray is detected, an event is produced that is similar in energy to that expected for a WIMP-nuclei interaction [35,52,53]. The nat K contents in the DAMA crystals are in the 10–20 ppb range [49], for the SABRE crystal it is 4.3 \pm 0.2 ppb [30], and for the COSINE-100 C6 it is 16.8 \pm 2.5 ppb [35].

In our apparatus (see Fig. 6), 40 K decays can be identified by coincidences between 1460 keV γ -rays in the CsI(Tl) detectors and 3.2 keV X-rays in the tested NaI(Tl) crystal. The 40 K background level in each crystal is determined by comparing the measured coincidence rate with a GEANT4-simulated rate using the method described in Refs. [24,34].

Figure 7 shows low energy spectra of the NaI(Tl) crystals associated with a 1460 keV γ tagged in the surrounding CsI(Tl) crystals. A clear 3.2 keV peak in NaI-025 (Fig. 7a) corresponds to a contamination of 684 ± 100 ppb nat K. In contrast, we found no evidence for potassium in NaI-035 and set a limit of 42 ppb (90% confidence level). Similarly, < 62 ppb and < 53 ppb are obtained for NaI-034 and NaI-036, respectively, which are consistent with the ICP-MS measurements in Table 2. Due to the small size of the tested crystals and the brief, 1-month long measurement period, there are large uncertainties in the nat K measurements. Because the ICP-MS results provided consistent and much precise results on nat K, we use the results from the ICP-MS measurements for the analysis discussed below.



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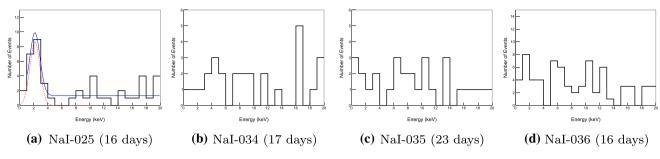


Fig. 7 The 40 K 3.2 keV spectra in the NaI(Tl) crystals tagged by a 1460 keV γ in the surrounding CsI(Tl) crystals. **a** A distinct 3.2 keV X-ray peak in NaI-025 reflects its high contamination of K. We model the energy spectrum (solid line) assuming a Gaussian 40 K signal (red-

dotted line) with a flat background. However, there is no evidence for the 3.2 keV line in NaI-034 ($\bf b$) and NaI-035 ($\bf c$) and NaI036 ($\bf d$). The exposure time for each crystal is denoted in parentheses

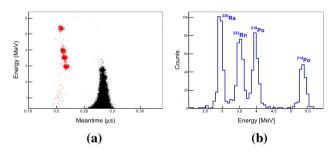


Fig. 8 A scatter plot of the mean time versus the measured energy for the NaI-036 crystal with 16 days exposure is shown (a). Here α -induced events (red dots) are well separated from β or γ -induced events (black dots). **b** The energy spectrum of the selected α -induced events is shown; each peak represents a specific daughter alpha decaying component from ^{226}Ra

4.1.2 α analysis

Alpha-induced events inside the crystals can be identified by the mean time of their signals, defined as,

$$\langle t \rangle = \frac{\sum_{i} A_{i} t_{i}}{\sum_{i} A_{i}}.$$

Here A_i and t_i are the charge and time in each digitized bin of a recorded event waveform. Figure 8a shows a scatter plot of the energy versus the mean time for the NaI-036 crystal. Alpha-induced events are clearly separated from β or γ -induced events because of their faster decay times. Selected α events (red points) are used to understand the internal contamination of heavy elements such as 238 U, 232 Th, and 210 Pb. We also analyzed the energy spectra of these α events (see Fig. 8b) in order to understand the contamination. Results of our analysis of the α contaminations are summarized in Table 3.

4.1.3 ²¹⁰Pb background

In the COSINE-100 experiment, the main source of the alpha contamination was due to decays of 210 Po ($E_{\alpha} = 5.4$ MeV)

nuclei that originated from a 222 Rn contamination [26,35]. This was confirmed by the observation of a 46.5 keV γ peak that is a characteristic of 210 Pb. In the 222 Rn decay chain, 210 Po (half-life $t_{1/2}=138$ days) is produced by the beta decay of 210 Pb ($t_{1/2}=22$ years) and 210 Bi ($t_{1/2}=5$ days), and subsequently decays via α emission to 206 Pb. Thus, if the 210 Po contamination is zero at the time the 222 Rn exposure occurs, it will grow with a characteristic time of τ_{210} Po = 200 days until an equilibrium level is reached. Thus, the total alpha rate over time can provide information about the amount of the 210 Pb contamination. We modeled the measured alpha rate (N) as following formula,

$$N = N_{\text{Pb210}}(1 - e^{-(t-t_0)/\tau_{210}}_{\text{Po}}) + C.$$

where $N_{\rm Pb210}$ is asymptotic $^{210}{\rm Pb}$ amount at infinite time, t_0 is contamination date fixed at the crystal was grown as summarized in Table 2, and C is the total alpha rate without $^{210}{\rm Po}$ decay. We can determine the N from the mean time distribution (Fig. 8). Figure 9 shows the measured alpha rates in the crystals as a function of data-taking time. The time-dependent fit with the $^{210}{\rm Po}$ time-dependent model overlaid. From these fits, the $^{210}{\rm Pb}$ level in NaI-025 is $3.80\pm0.30~{\rm mBq/kg}$ and, no significant $^{210}{\rm Pb}$ is observed for NaI-034 & NaI-035 as one can see in Fig. 9. The measured level of $^{210}{\rm Pb}$ in NaI-036 (0.42 $\pm0.27~{\rm mBq/kg}$) is consistent with the $^{226}{\rm Ra}$ contamination (0.45 $\pm0.05~{\rm mBq/kg}$).

4.1.4 ²³²Th background

Contaminants from the 232 Th decay chain can be identified by studying three-alpha delayed time coincidences of 224 Ra, 220 Rn and 216 Po, where the half-lives are 3.66 day, 55.6 s, and 0.145 s, respectively. Due to the short half-life of 216 Po ($t_{1/2}$ =0.145 s), identification of two α particles from 220 Rn and 216 Po is relatively straight-forward and background free, as can be seen in Fig. 10a. With the selection of two consecutive α signals, 224 Ra can also be identified. No significant 232 Th α -decay events in NaI-025 and NaI-036 are measured.



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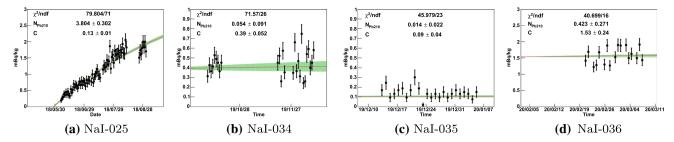


Fig. 9 The total alpha rates in the crystals as a function of time are modeled as a contamination of 222 Rn (and/or 210 Pb) at the crystal growth dates described in the text. Here black points are data, red-solid lines are the best fit models, and green bands are 1σ uncertainty regions of the fit. **a** An increase in the alpha rates corresponding to the 210 Po lifetime is evident in NaI-025. No significant increases in alpha rates indicate

that there are no significant 210 Pb contaminations present in NaI-034 (**b**) or NaI-035 (**c**) during the crystal growth. **d** 226 Ra contamination in NaI-036 resulted in a relatively large number of measured alpha events. The total alpha rate is consistent with a 0.45 mBq/kg contamination of 210 Pb that originated from 226 Ra decays

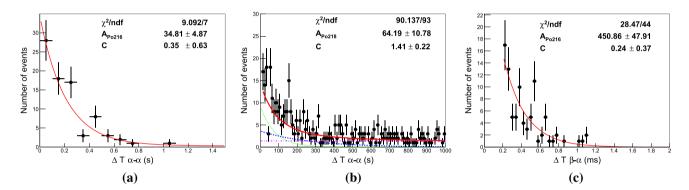


Fig. 10 The time difference (ΔT) distribution of data (points) and exponential fits (red-solid line) between two successive α -induced events are presented. Here decay time used in the fit is fixed to the known life-time of each radioisotope and measured rates are converted into activities (A) in unit of $\mu Bq/kg$. ΔT between two α decays of a $^{220}Rn \rightarrow ^{216}Po$ (half-life 0.145 s) and **b** $^{224}Ra \rightarrow ^{220}Rn$ (green-dotted line, half-life 55.6 s), where the activity of ^{220}Rn is constraint to

²¹⁶Po measurement in (**a**), and ²²²Rn→²¹⁸Po (blue-dashed line, half-life 186 s) together with random coincidence events (purple-dashed-dotted line), in the NaI-034 crystal; (**c**) Δ T between successive β and α events from ²¹⁴Bi→²¹⁴Po (half-life 0.164 ms) in the NaI-036 crystal are shown. Due to dead time of the data acquisition system, events with Δ T less than 0.2 ms were not recorded

However, the presence of α - α delayed coincidence events in NaI-034 and NaI-035 is evident in Fig. 10 with measured activity levels of 35±5 μ Bq/kg and 7±2 μ Bq/kg for NaI-034 and NaI-035, respectively. Assuming equilibrium of the 232 Th chain, this is similar to the 2–31 μ Bq/kg level in the DAMA crystals.

4.1.5 ²³⁸U background

²³⁸U is one of the most common radioisotopes in nature primarily because of its long decay time. A study of its decay products allows us to understand its contamination level. Delayed coincidence α – α events with $t_{1/2}=3.10$ min time difference from ²²²Rn \rightarrow ²¹⁸Po can be used to infer the ²³⁸U contamination levels, as shown in Fig. 10b. The measured rate for ²²⁴Ra \rightarrow ²²⁰Rn ($t_{1/2}=55.6$ s) is extracted from the ²³²Th level, which is determined from ²²⁰Rn \rightarrow ²¹⁶Po decays. In this way, we infer ²¹⁸Po levels of 64±11 μBq/kg and 8±6

 μ Bq/kg for NaI-034 and NaI-035, respectively. Because of the high total α rates in NaI-025 and NaI-036, we cannot measure 3 min half-life α - α delayed coincidence events.

Another method for studying the 238 U chain exploits the 164.3 μ s half-life 214 Po α -decay, which follows its production via the β -decay of 214 Bi. Due to the 200 μ s dead time of the trigger system, only events in the higher end of the exponentially falling distribution are acquired but the number of these events is sufficient to measure the 238 U contamination (see in Fig. 10c). The measured activity calculation includes a correction for the effect of the inefficiency caused by the 200 μ s dead time. NaI-036 shows a relatively large 214 Po contamination of 451 \pm 48 μ Bq/kg, while NaI-025, NaI-034, and NaI-036 show contaminations of 26 \pm 7 μ Bq/kg, 46 \pm 7 μ Bq/kg, and 13 \pm 6 μ Bq/kg, respectively. The contamination in NaI-036 is considerably larger than those of the DAMA crystals (8.7–124 μ Bq/kg) but the other crystal contamination levels are similar to those of DAMA.



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The energy spectrum of α events in Fig. 8b can be used to determine whether or not the ²³⁸U chain is in equilibrium, especially for NaI-036. The α energy spectrum of NaI-036 can be well explained by ²²⁶Ra decays with a 0.45±0.05 mBq/kg activity with no higher components from the ²³⁸U chain. At the start of the measurement time period, we did not see evidence for 210 Po decay events in the α spectrum but the evolution plot in Fig. 9d can be explained by an approximately 0.45 mBq/kg of the ²¹⁰Pb contamination coming from ²²⁶Ra. The ²²⁶Ra in NaI-036 can be attributed to impurities contained in the TII powder. Specially, the SA powder (initial loading 0.1 mol% of NaI-034) and SA beads (initial loading 1 mol% of NaI-036) contribute consistent contaminations of ²²⁶Ra to the crystals. If this is the case, we expect that the use of TII beads from Alpha Aesar (NaI-035) could reduce the ²²⁶Ra contamination significantly.

4.2 External background

In the COSINE-100 experiment, the external background contributions in the region of interest (ROI) are negligible [35] because of the active veto provided by the 2200 L liquid scintillator [37]. However, the R&D setup contains a significant level of external background that is more than 10 times that for a full-size COSINE crystal [52]. Because the PMTs are the main contributor to the external background, a small crystal will have relatively larger background contribution per kilogram from external sources. Since we use the same type of PMTs in the COSINE-100 experiment, we can estimate the external background contribution using the GEANT4 simulation based on previous studies [35,52].

4.3 Cosmogenic radionuclides

The cosmogenic production of radioactive isotopes in the NaI(Tl) crystal contributes a non-negligible amount of background in ROI of the COSINE-100 experiment [35,54]. This is mainly due to long-lived nuclides such as ³H and ²²Na. All crystals of the COSINE-100 experiment were produced at the Alpha Spectra company located at Grand Junction, Colorado, USA. Due to its high altitude of about 1400 m above sea level as well as the long delivery time to Korea, the crystals were exposed to a significant flux of cosmic-ray muons and their spallation products and, as a result, we measure a relatively high level of cosmogenic radioactive isotopes. However, the crystals that are used for this study were all grown in Daejeon, Korea (70 m in altitude) and delivered underground in a short time period (typically within a few weeks). Based on our previous study, described in Ref. [54], 1-month exposure time near sea level can produce 0.004 mBq/kg of ³H and 0.05 mBq/kg of ²²Na. These activities are an order of magnitude lower than measured activities in COSINE-100 C6, which are 0.11 mBq/kg and 0.73 mBq/kg [54] of ³H and ²²Na, respectively. Thus, we expect that the background contributions from these long-lived cosmogenic radioisotopes will be sufficiently low for the COSINE-200 experiment.

5 Background modeling and prospects of COSINE-200

For a quantitative understanding of background in these recently grown crystals, we use the GEANT4-based simulation that was developed for the background modeling of the COSINE-100 NaI(Tl) crystals [35,52]. We take values of the contamination levels discussed in Sect. 4 as input. Because of the low statistics of the underground ⁴⁰K measurements, we use the ICP-MS results for ^{nat}K from Table 2.

We use a log-likelihood method to fit the data. The fitting range is 3 keV-3 MeV and we perform a simultaneous fit of four-channels: single-hit low-energy (3-60 keV) and high energy (60 keV-3 MeV), multiple-hit low-energy and high energy. Here, multiple hit corresponds to events in which there is one or more coincident hits recorded in any of the surrounding CsI(Tl) crystals. The internal backgrounds are constrained to the measured uncertainties. We do not consider surface ²¹⁰Pb components in the crystals because the crystal surface is polished in a nitrogen gas environment and we do not see any indication of increases in energy spectra in the low-energy region that are characteristic of surface ²¹⁰Pb contamination [55]. The backgrounds from the PMTs that are attached to the NaI(Tl) and CsI(Tl) crystals are measured with a high purity Germanium detector [35,52]. These values are constrained to be within 50% of the measured results because exact locations of these radioisotopes are uncertain. The long-lived cosmogenic radioisotopes are constrained to be within 50% of its calculated production values that are estimated in Section 4.3 while other shortlived cosmogenic components are floated. Additional external backgrounds from the CsI(Tl) crystals and copper surface are also included as free parameters but the final results are cross-checked with previous measurements made with the same setup [52].

Figure 11 shows the fitted NaI-035 results for all of the simulated background components plotted as lines with different colors. The overall energy spectra are well matched to the data for both single-hit and multi-hit events. The small sizes of the crystals as well as the non-existence of an active veto detector in the R&D setup results in the external background being the dominant contributor to these spectra. However, this contribution will be significantly reduced with the large, COSINE-100-sized, crystals and the liquid scintillator veto system that will be used in the upcoming COSINE-200 setup. Table 4 shows a summary of input values and fitted results from this background modeling.

To estimate the expected background in the COSINE-200 crystal, we generate simulated background spectra assuming



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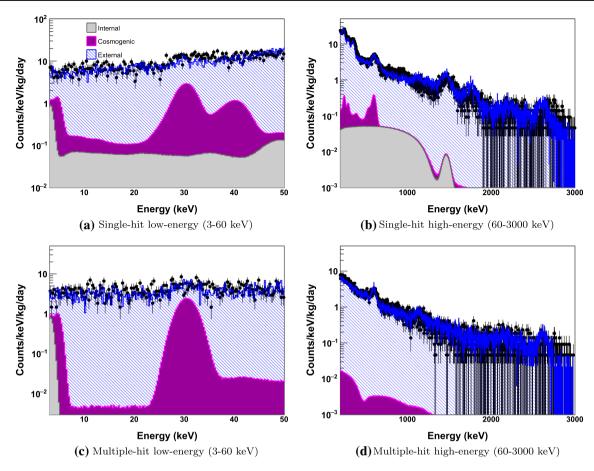


Fig. 11 Measured single-hit and multiple-hit background spectra of NaI-035 (black points) are determined for all simulated backgrounds using a simultaneous fit to the four-channels. The external component

(blue-dashed filled) is the dominant contributor; this will be reduced with full-size detectors situated in the liquid scintillator veto system

the same internal and cosmogenic radioisotopes in the crystals. Because of the relatively small contribution, the background modeling fit does not provide more accurate results for measurements of a few internal and cosmogenic ³H components, as shown in Table 4. The more accurate results between the measured activities shown in Table 3 and the fitted activities shown in Table 4 are used. We consider the crystal size enlarged to that of the C6 (12.5 kg) and installed inside a COSINE-100-like shield. The latter consists of a 2,200 L liquid scintillator tank, a 3 cm thick copper, a 20 cm thick lead, and an array of plastic counters for muon tagging [56] as described in Refs. [26,57]. We assume levels of external background that are similar to those seen in the COSINE-100 detector [35]. Because of significantly reduced ²¹⁰Pb contaminations, the internal background contributions are reduced. We also expect that the lower exposure to cosmic muons will significantly reduce the level of long-lived cosmogenic such as ³H and ²²Na. Figure 12a shows the expected background spectrum for the COSINE-200 detector assuming the internal background as NaI-035 background measure-

ments in Table 4. By taking advantage of the 2200 L active veto detector as well as the enlarged detector size, a significantly reduced background in low-energy single-hit events is expected. In the 1–6 keV ROI, the expected background level is approximately 0.33 counts/kg/day/keV, which is below that of the DAMA experiment. After time spent underground, the short-lived cosmogenic radioisotopes will decay away as shown in Fig. 12b. After 2 years, the background is expected to decrease to about 0.13 counts/kg/day/keV. If we consider different background levels obtained from NaI-034 and NaI-036, the expected background from the COSINE-200 detectors would be slightly increased as can be seen in Fig. 12c. However, it is still less than 1 counts/kg/day/keV that is target background level for the COSINE-200 experiment [34].

6 Conclusion

In the quest for dark matter direct detection, the origin of the DAMA modulation signal remains the subject of a long-



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Table 4 Summary of the measured and fitted radioactive contaminants in the modeling of the NaI-035 crystal

Background sources	Isotopes	Activities(mBq/l	Activities(mBq/kg)			
		Measured	Input constraint	Fitted		
Internal	²³⁸ U	0.011±0.004	±0.004	0.01±0.01		
	²²⁸ Th	0.007 ± 0.002	± 0.002	0.006 ± 0.004		
	40 K	0.39 ± 0.18	± 0.18	0.52 ± 0.18		
	²¹⁰ Pb	0.014 ± 0.022	± 0.022	0.009 ± 0.009		
Cosmogenic	^{125}I	_	_	0.39 ± 0.01		
	¹²¹ Te	_	_	0.80 ± 0.03		
	121m Te	_	_	0.06 ± 0.01		
	$^{123m}{ m Te}$	_	_	0.10 ± 0.09		
	$^{125m}{ m Te}$	_	_	0.14 ± 0.01		
	$^{127m}{ m Te}$	_	_	0.19 ± 0.02		
	¹⁰⁹ Cd	_	_	0.007 ± 0.001		
	¹¹³ Sn	_	_	0.015 ± 0.01		
	²² Na	0.05	± 0.03	0.05 ± 0.01		
	^{3}H	0.004	± 0.002	0.005 ± 0.005		
NaI PMTs	^{238}U	25±5	±13	27.2 ± 11.2		
	²³² Th	12±5	±6	16.8 ± 1.2		
	40 K	58±5	±29	58.6 ± 18.0		
CsI PMTs	^{238}U	78±4	±39.1	81.0 ± 14.6		
	²³² Th	26±4	±13	22.5 ± 2.3		
	$^{40}\mathrm{K}$	504±72	±252	395.7±7.2		

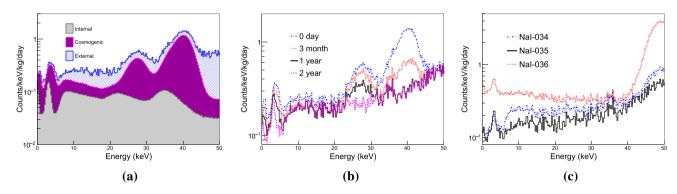


Fig. 12 a The expected background spectrum of a COSINE-200 crystal for low-energy single-hit events assuming the same internal and cosmogenic radioisotopes as NaI-035 but, installed in the COSINE-100 shield with an enlarged crystal size (12.5 kg). **b** The time evolution of the background level due to decays of short-lived cosmogenic radioisotopes after 0 days (blue-dashed-dotted line), 3 months

(red-dotted line), 1 year (purple-dashed line), and 2 years (black-solid line). c The expected background spectra of the COSINE-200 detectors assuming background levels obtained from NaI-034 (blue-dashed-dotted line), NaI-035 (black-solid line), and NaI-036 (red-dotted line) after being underground for 2 years

standing debate. To expand the currently ongoing COSINE-100 experiment and reach a model-independent conclusion, we propose the COSINE-200 experiment with 200 kg ultrapure NaI(Tl) crystals. As a part of a program to reduce background levels, we have grown four small-size NaI(Tl) crystals in our laboratory. In this article, we present results from performance studies of these R&D devices that demonstrate the feasibility of the production of low-background NaI(Tl) crys-

tals suitable for the COSINE-200 experiment. It is necessary to have further R&D on full-size crystals, and this is currently underway. Assuming the full-size crystals of similar quality as the NaI-035 detector that is installed in the COSINE-100-like shield, our study shows that an expected background level will be less than 0.5 counts/kg/day/keV and lower than those of the DAMA crystals. In this case COSINE-200 will



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be able to produce a definitive and unambiguous test of the DAMA experiment's annual modulation signal.

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References

- D. Clowe et al., A direct empirical proof of the existence of dark matter. Astrophys. J. 648, L109 (2006)
- N. Aghanim et al., (Planck Collaboration), Planck 2018 results. VI. Cosmological parameters. arXiv:1807.06209
- G. Bertone, D. Hooper, History of dark matter. Rev. Mod. Phys. 90, 045002 (2018)
- B.W. Lee, S. Weinberg, Cosmological lower bound on heavyneutrino masses. Phys. Rev. Lett. 39, 165 (1977)
- M.W. Goodman, E. Witten, Detectability of certain dark matter candidates. Phys. Rev. D 31, 3059 (1985)
- T.M. Undagoitia, L. Rauch, Dark matter direct-detection experiments. J. Phys. G 43, 013001 (2016)
- M. Schumann, Direct detection of WIMP dark matter: concepts and status. J. Phys. G 46, 103003 (2019)
- R. Bernabei et al., (DAMA/LIBRA Collaboration), Final model independent result of DAMA/LIBRA-phase1. Eur. Phys. J. C 73, 2648 (2013)
- R. Bernabei et al., (DAMA/LIBRA Collaboration), First Model Independent Results from DAMA/LIBRA-Phase2. Nucl. Phys. At. Energy 19, 307 (2018)
- C. Savage, G. Gelmini, P. Gondolo, K. Freese, Compatibility of DAMA/LIBRA dark matter detection with other searches. JCAP 0904, 010 (2009)
- Y. J. Ko et al., (COSINE-100 Collaboration), Comparison between DAMA/LIBRA and COSINE-100 in the light of quenching factors. JCAP 1911, 008 (2019)
- M. Tanabashi et al., (Particle Data Group Collaboration), Review of particle physics. Phys. Rev. D 98, 030001 (2018)
- S.C. Kim et al., New limits on interactions between weakly interacting massive particles and nucleons obtained with CsI(Tl) crystal detectors. Phys. Rev. Lett. 108, 181301 (2012)

- G. Angloher et al., (CRESST-II Collaboration), Results on low mass WIMPs using an upgraded CRESST-II detector. Eur. Phys. J. C 74, 3184 (2014)
- R. Agnese et al., (SuperCDMS Collaboration), Search for low-mass weakly interacting massive particles with SuperCDMS. Phys. Rev. Lett. 112, 241302 (2014)
- K. Abe et al., (XMASS Collaboration), Direct dark matter search by annual modulation in XMASS-I. Phys. Lett. B 759, 272 (2016)
- D. S. Akerib et al., (LUX Collaboration), Results from a search for dark matter in the complete LUX exposure. Phys. Rev. Lett. 118, 021303 (2017)
- E. Aprile et al., (XENON Collaboration), Search for electronic recoil event rate modulation with 4 years of XENON100 data.' Phys. Rev. Lett. 118, 101101 (2017)
- R. Agnese et al., (SuperCDMS Collaboration), Results from the super cryogenic dark matter search experiment at Soudan. Phys. Rev. Lett. 120, 061802 (2018)
- E. Aprile et al., (XENON Collaboration), Dark matter search results from a one tonne × year exposure of XENON1T. Phys. Rev. Lett. 121, 111302 (2018)
- P. Agnes et al., (DarkSide Collaboration), Low-mass dark matter search with the DarkSide-50 experiment. Phys. Rev. Lett. 121, 081307 (2018)
- D.S. Akerib et al., (LUX Collaboration), Search for annual and diurnal rate modulations in the LUX experiment. Phys. Rev. D 98, 062005 (2018)
- A.H. Abdelhameed et al., (CRESST Collaboration), First results from the CRESST-III low-mass dark matter program. Phys. Rev. D 100, 102002 (2019)
- K.W. Kim et al., Tests on NaI(Tl) crystals for WIMP search at the Yangyang Underground Laboratory. Astropart. Phys. 62, 249 (2015)
- M. Antonello et al., (SABRE Collaboration), The SABRE project and the SABRE Proof-of-Principle. Eur. Phys. J. C 79, 363 (2019)
- G. Adhikari et al., (COSINE-100 Collaboration), Initial performance of the COSINE-100 experiment. Eur. Phys. J. C 78, 107 (2018)
- K.-I. Fushimi, Low background measurement by means of NaI(Tl) scintillator: improvement of sensitivity for cosmic dark matter. RADIOISOTOPES 67, 101 (2018)
- I. Coarasa et al., ANAIS-112 sensitivity in the search for dark matter annual modulation. Eur. Phys. J. C 79, 233 (2019)
- J. Amare et al., Performance of ANAIS-112 experiment after the first year of data taking. Eur. Phys. J. C 79, 228 (2019)
- B. Suerfu, M. Wada, W. Peloso, M. Souza, F. Calaprice, J. Tower,
 G. Ciampi, Growth of ultra-high purity NaI(Tl) crystal for dark matter searches. Phys. Rev. Res. 2, 013223 (2020)
- G. Adhikari et al., (COSINE-100 Collaboration), An experiment to search for dark-matter interactions using sodium iodide detectors. Nature 564, 83 (2018)
- G. Adhikari et al., (COSINE-100 Collaboration), Search for a dark matter-induced annual modulation signal in NaI(Tl) with the COSINE-100 experiment. Phys. Rev. Lett. 123, 031302 (2019)
- J. Amare et al., First results on dark matter annual modulation from ANAIS-112 experiment. Phys. Rev. Lett. 123, 031301 (2019)
- P. Adhikari et al., (KIMS Collaboration), Understanding internal backgrounds in NaI(Tl) crystals toward a 200 kg array for the KIMS-NaI experiment. Eur. Phys. J. C 76, 185 (2016)
- P. Adhikari et al., (COSINE-100 Collaboration), Background model for the NaI(Tl) crystals in COSINE-100. Eur. Phys. J. C 78, 490 (2018)
- J. S. Park et al., (KIMS Collaboration), Performance of a prototype active veto system using liquid scintillator for a dark matter search experiment. Nucl. Instrum. Methods A 851, 103 (2017)
- G. Adhikari et al., The COSINE-100 liquid scintillator veto system. arXiv:2004.03463



Eur. Phys. J. C (2020) 80:814 Page 13 of 13 814

- Y. Kanemitsu, D. Chernyak, H. Ejiri, K. Fushimi, K. Hata, R. Hazama, H. Ikeda, K. Imagawa, K. Inoue, A. Kozlov, R. Orito, T. Shima, Y. Takemoto, S. Umehara, K. Yasuda, S. Yoshida, Purification of the NaI(tl) crystal for dark matter search project PICOLON. J. Phys. Conf. Ser. 1468, 012054 (2020)
- K. Shin, O. Gileva, Y. Kim, H.S. Lee, H. Park, Reduction of the radioactivity in sodium iodide (NaI) powder by recrystallization method. J. Radioanal. Nucl. Chem. 317, 1329 (2018)
- K. Shin, J. Choe, O. Gileva, A. Iltis, Y. Kim, C. Lee, H.S. Lee, M.H. Lee, H.K. Park, A facility for mass production of ultra-pure NaI powder for the COSINE-200 experiment. JINST 15, C07031 (2020)
- 41. S. Ra et al., Scintillation crystal growth at the CUP. PoS ICHEP2018, 668 (2019)
- 42. H. S. Lee et al., (KIMS Collaboration), First limit on WIMP cross section with low background CsI(Tl) crystal detector. Phys. Lett. B 633, 201 (2006)
- 43. H.S. Lee et al., Development of low-background CsI(Tl) crystals for WIMP search. Nucl. Instrum. Methods A **571**, 644 (2007)
- G. Adhikari et al., (COSINE-100 Collaboration), The COSINE-100 data acquisition system. JINST 13, P09006 (2018)
- G. Adhikari et al., Lowering the energy threshold in COSINE-100 dark matter searches. arXiv:2005.13784
- S. Kubota, F. Shiraishi, Y. Takami, Scintillation process in NaI(Tl): comparison with scintillation models. J. Phys. Soc. Jpn. 69, 3435 (2000)
- L. Trefilova, A. Kudin, L. Kovaleva, B. Zaslavsky, D. Zosim, S. Bondarenko, Concentration dependence of the light yield and energy resolution of NaI: Tl and CsI: Tl crystals excited by gamma, soft X-rays and alpha particles. Nucl. Instrum. Methods A 486, 474 (2002)

- J. Choi, B. Park, C. Ha, K. Kim, S. Kim, Y. Kim, Y. Ko, H. Lee, S. Lee, and S. Olsen, Improving the light yield of NaI(Tl) crystal detectors. Nucl. Instrum. Methods A 981, 164556 (2020). https:// doi.org/10.1016/j.nima.2020.164556
- R. Bernabei et al., (DAMA/LIBRA Collaboration), The DAMA/LIBRA apparatus. Nucl. Instrum. Meth. A 592, 297 (2008)
- R. Bernabei et al., (DAMA/LIBRA Collaboration), Performances of the new high quantum efficiency PMTs in DAMA/LIBRA. JINST 7, P03009 (2012)
- J.H. Friedman, Greedy function approximation: a gradient boosting machine. Ann. Stat. 29, 1189 (2001)
- G. Adhikari et al., (KIMS Collaboration), Understanding NaI(Tl) crystal background for dark matter searches. Eur. Phys. J. C 77 437 (2017)
- J. Amare et al., Analysis of backgrounds for the ANAIS-112 dark matter experiment. Eur. Phys. J. C 79, 412 (2019)
- E. Barbosa de Souza et al., (COSINE-100 Collaboration), Study of cosmogenic radionuclides in the COSINE-100 NaI(Tl) detectors. Astropart. Phys. 115, 102390 (2020)
- G.H. Yu, C. Ha, E.J. Jeon, K.W. Kim, N.Y. Kim, Y.D. Kim, H.S. Lee, H.K. Park, and C. Rott, Depth-profile study of 210Pb in the surface of an NaI(T) crystal. arXiv:2001.06132
- H. Prihtiadi et al., (COSINE-100 Collaboration), Muon detector for the COSINE-100 experiment. JINST 13, T02007 (2018)
- C. Ha et al., (COSINE-100 Collaboration), First direct search for inelastic boosted dark matter with COSINE-100. Phys. Rev. Lett. 122, 131802 (2019)

