



Cyclicity of all anti-NMHV and N^2 MHV tree amplitudes in $\mathcal{N}=4$ SYM

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Abstract This article proves the cyclicity of anti-NMHV and N^2 MHV tree amplitudes in planar $\mathcal{N}=4$ SYM up to any number of external particles as an interesting application of positive Grassmannian geometry. In this proof the two-fold simplex-like structures of tree amplitudes introduced in 1609.08627 play a key role, as the cyclicity of amplitudes will induce similar simplex-like structures for the boundary generators of homological identities. For this purpose, we only need a part of all distinct boundary generators, and the relevant identities only involve BCFW-like cells. The manifest cyclic invariance in this geometric representation reflects one of the invariant characteristics of amplitudes, though they are obtained by the scheme-dependent BCFW recursion relation.

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1 Introduction

$\mathcal{N}=4$ super Yang–Mills theory has been the most understood quantum field theory so far. In recent years, tremendous progress on the scattering amplitudes of planar $\mathcal{N}=4$ SYM was made through its connection to positive Grassmannian, momentum twistors and on-shell diagrams [1, 2]. In this new context, the well known BCFW recursion relation [3] using massless spinors has an elegant generalization [4] in momentum twistor space. This efficient machinery is powerful for generating tree amplitudes and loop integrands, and it can manifest dual superconformal invariance of planar $\mathcal{N}=4$

SYM. There is a MATHEMATICA package “positroids” to implement these results, with investigations on various mathematical aspects of positive Grassmannian [5]. More relevant background on amplitudes can be found in [6, 7].

In particular, tree amplitudes in planar $\mathcal{N}=4$ SYM have an impressive simplicity in the language of positive Grassmannian in momentum twistor space, namely the so-called two-fold simplex-like structures [8], a concise review can be referred in [9]. In terms of Grassmannian geometry representatives specifying linear dependencies of different ranks and empty slots for null columns, information of amplitudes can be compactly captured by finite numbers of fully-spanning cells and their growing parameters. Given a fixed k , as $(k+2)$ is the number of negative helicities, there is no new full cell beyond $n=4k+1$, then after we identify all full cells with their growing parameters at this critical n , N^k MHV amplitudes are known once for all up to any number of external particles. This is an extension following the logic similar to [10, 11].

With the aid of this purely geometric description, homological identities can be understood in a much more intuitive way, and most of them turn out to be the secret incarnation of the simple NMHV identity. A part of these identities are crucial for interconnecting different BCFW cells, and hence different BCFW recursion schemes [1]. Explicitly in this work, we would like to manifest the cyclicity of amplitudes of two specific classes: the anti-NMHV and the N^2 MHV families, by applying the simplex-like structures of both the amplitudes and boundary generators of identities. From [8] we have fully understood the structures of anti-NMHV, NMHV, N^2 MHV and N^3 MHV families, while only the cyclicity of NMHV family and $n=7, 8$ anti-NMHV amplitudes has been shown. It is then desirable to see more nontrivial examples and attempt to extract the general pattern from them. And the manifest cyclic invariance reflects one of the invariant characteristics of amplitudes, which is obscured by the scheme-dependent BCFW recursion relation. This cyclicity is not

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manifest in the amplituhedron setting [12–15] as well, as the triangulation process usually chooses some fixed labels of particles for simplicity. Moreover, though the NMHV identity has an obvious geometric interpretation as different triangulations give the same invariant sum of “volumes”, for more general $N^k\text{MHV}$ identities with $k \geq 2$ the corresponding geometric pictures are unclear yet. From the perspective of cyclicity, we may find more intuition of these identities which interconnect BCFW cells of different Grassmannian geometric configurations. We may even go further to find their counterparts in the context of amplituhedron, and in particular, explore their relation to the sign-flip triangulation [13, 14].

As a helpful warmup exercise, we now reconsider the cyclicity of NMHV family in a more formal way before we derive its generalization for $k \geq 2$ in this work.

Recall the NMHV $n=6$ amplitude in terms of empty slots in the default recursion scheme is given by

$$Y_6^1 = [6] + [4] + [2], \quad (1.1)$$

where Y_n^k is the Yangian invariant related to the n -particle amplitude with $(k+2)$ negative helicities, via $A_n^k = A_n^{\text{MHV}} Y_n^k$ (the MHV sector means $k=0$), and an empty slot $[i]$ of $k=1, n=6$ denotes the commonly used 5-bracket with entry i removed [1, 8]. Then the difference between Y_6^1 and its cyclically shifted (by $+1$) counterpart is

$$Y_6^1 - Y_{6,+1}^1 = -[1] + [2] - [3] + [4] - [5] + [6] \equiv I_{123456}, \quad (1.2)$$

here the 6-term NMHV identity of labels 1, 2, 3, 4, 5, 6 is defined as I_{123456} . According to the simplex-like structures of amplitudes, namely the quadratic mode with growing parameters $(6, 4, 2)$, given by (in this triangle-shape sum each entry is multiplied by its corresponding vertical and horizontal factors)

we have the following relation for $n=7$ as an example:

$$\begin{aligned} Y_7^1 - Y_{7,+1}^1 &= \begin{pmatrix} [23] \\ [27] [25] \\ [67] [47] [45] \end{pmatrix} - \begin{pmatrix} [34] \\ [31] [36] \\ [71] [51] [56] \end{pmatrix} \\ &= \begin{pmatrix} [3] I_{124567} \\ [7] I_{123456} [5] I_{123467} \end{pmatrix}. \end{aligned} \quad (1.4)$$

This already completes the proof of cyclicity for the general Y_n^1 , since the growing parameters $(7, 5, 3)$ of I_{123456} have been identified, and as n increases I_{123456} also follows the simplex-like growing pattern.

We see that $(7, 5, 3)$ are closely related to $(6, 4, 2)$ of Y_n^1 , which shows how the cyclicity of amplitudes induces similar simplex-like structures for the relevant homological identities. This intriguing feature will appear in a much more nontrivial form for $N^2\text{MHV}$ amplitudes.

2 Cyclicity of anti-NMHV amplitudes

Before moving to the $N^2\text{MHV}$ family, let's first consider the cyclicity of all anti-NMHV amplitudes, since this is in fact the nontrivial starting point for all $N^k\text{MHV}$ cases. More explicitly, recall that for a given k non-vanishing amplitudes start with the anti-MHV sector $n=k+4$, which contains just one top cell, then the first interesting case is the anti-NMHV sector $n=k+5$. It can be rearranged in the similar form of a triangle-shape sum as its parity conjugate, namely the NMHV sector.

The anti-NMHV triangle-like pattern can be clearly observed in the series of examples below:

$$Y_6^1 = \begin{pmatrix} [2] \left\{ \begin{matrix} [6] \\ [4] \end{matrix} \right\} \end{pmatrix}, \quad (2.1)$$

$$Y_7^2 = \begin{pmatrix} [2] \quad (23) \left\{ \begin{matrix} (67) \\ (45) \end{matrix} \right\} \left\{ \begin{matrix} [7] \\ (45)(71) \\ [5] \end{matrix} \right\} \end{pmatrix}, \quad (2.2)$$

$$Y_n^1 = \begin{pmatrix} [234 \dots n-4] & & & & & 1 \\ \vdots & & & & \ddots & \vdots \\ [234] & & & 1 & \dots & [\dots n-2] \\ [23] & & & [6] & \dots & [6 \dots n-2] \\ [2] & & 1 & [5] & [56] & \dots & [56 \dots n-2] \\ 1 & 1 & [4] & [45] & [456] & \dots & [456 \dots n-2] \\ \hline [678 \dots n] & [78 \dots n] & [8 \dots n] & [\dots n] & \dots & 1 \end{pmatrix}, \quad (1.3)$$

$$Y_8^3 = \left([2] \quad (23) \left\{ \begin{smallmatrix} (678) \\ (456) \end{smallmatrix} \right\} (234) \left\{ \begin{smallmatrix} (78) \\ (456)(781) \\ (56) \end{smallmatrix} \right\} \left\{ \begin{smallmatrix} [8] \\ (456)(81) \\ (56)(812) \\ [6] \end{smallmatrix} \right\} \right), \quad \text{for } k=2 \text{ (we have shifted boundary generator (12) in [8] to (23), and similar below), as well as} \quad (2.3)$$

$$\partial(23) = +[2] - [3] + (23)(456) - (23)(567) + (23)(678) - (23)(781), \quad (2.11)$$

$$Y_9^4 = \left([2] \quad (23) \left\{ \begin{smallmatrix} (6789) \\ (4567) \end{smallmatrix} \right\} (234) \left\{ \begin{smallmatrix} (789) \\ (4567)(7891) \\ (567) \end{smallmatrix} \right\} (2345) \left\{ \begin{smallmatrix} (89) \\ (4567)(891) \\ (567)(8912) \\ (67) \end{smallmatrix} \right\} \left\{ \begin{smallmatrix} [9] \\ (4567)(91) \\ (567)(912) \\ (67)(9123) \\ [7] \end{smallmatrix} \right\} \right), \quad (2.4)$$

and its general form can be proved by induction. Here we remind that, for $k \geq 2$, the empty slot $[i]$ again denotes the Grassmannian cell with column i removed, (ij) denotes that columns i, j are proportional and (ijk) , $(ijkl)$ similarly denote linear dependencies of various ranks (in particular, $(i_1 \dots i_k)$ denotes a $k \times k$ vanishing minor) as defined in [8].

From [8], it is already known that

$$Y_7^2 - Y_{7,+1}^2 = -\partial(23) - \partial(56) - \partial(71), \quad (2.5)$$

$$Y_8^3 - Y_{8,+1}^3 = \partial(23) + \partial(67) + \partial(81) + \partial(234)(567) + \partial(567)(812) + \partial(781)(234), \quad (2.6)$$

which manifest the cyclicity of Y_7^2 and Y_8^3 . As boundary generators, the $(4k+1)$ -dimensional cells above lead to the relevant identities after the formal ‘ ∂ ’ operation. These results can be rearranged in a more suggestive form as

$$Y_7^2 - Y_{7,+1}^2 = -\partial \left((23) \left\{ \begin{smallmatrix} (71) \\ (56) \end{smallmatrix} \right\} \right), \quad (2.7)$$

$$Y_8^3 - Y_{8,+1}^3 = \partial \left((23) \quad (234) \left\{ \begin{smallmatrix} (781) \\ (567) \end{smallmatrix} \right\} \left\{ \begin{smallmatrix} (81) \\ (567)(812) \\ (67) \end{smallmatrix} \right\} \right), \quad (2.8)$$

as well as a further $k=4$ extension of this pattern:

$$Y_9^4 - Y_{9,+1}^4 = -\partial \left((23) \quad (234) \left\{ \begin{smallmatrix} (7891) \\ (5678) \end{smallmatrix} \right\} (2345) \left\{ \begin{smallmatrix} (891) \\ (5678)(8912) \\ (678) \end{smallmatrix} \right\} \left\{ \begin{smallmatrix} (91) \\ (5678)(912) \\ (678)(9123) \\ (78) \end{smallmatrix} \right\} \right), \quad (2.9)$$

where the sign factor $(-)^{k+1}$ for each of these relations follows the convention of [1, 5]. And the types of homological identities used in (2.7) and (2.8), as already proved in [8], include

$$\partial(23) = -[2] + [3] - (23)(45) + (23)(56) - (23)(67) + (23)(71) \quad (2.10)$$

$$\partial(234)(567) = +(23)(567) - (34)(567) + (234)(56) - (234)(67) + (234)(567)(781) - (234)(567)(812) \quad (2.12)$$

for $k=3$, while those for $k=4$ used in (2.9) are new, as given by

$$\partial(23) = -[2] + [3] - (23)(4567) + (23)(5678) - (23)(6789) + (23)(7891), \quad (2.13)$$

$$\partial(234)(5678) = -(23)(5678) + (34)(5678) - (234)(567) + (234)(678) - (234)(5678)(7891) + (234)(5678)(8912), \quad (2.14)$$

$$\partial(2345)(678) = -(234)(678) + (345)(678) - (2345)(67) + (2345)(78) - (2345)(678)(8912) + (2345)(678)(9123), \quad (2.15)$$

$$\partial(2345)(5678)(8912) = -(234)(5678)(8912)$$

$$+ (345)(5678)(8912) - (2345)(567)(8912) + (2345)(678)(8912) - (2345)(5678)(891) + (2345)(5678)(912), \quad (2.16)$$

and they can be proved by using the similar matrix approach as done in [8]. The examples of anti-NMHV family again show how the cyclicity of amplitudes induces similar structures for the relevant identities, and such 6-term identities for any k can be easily guessed from the boundary generators then proved.

3 Cyclicity of N^2 MHV amplitudes

Now we will start with the cyclicity of N^2 MHV $n = 7$ amplitude, namely (2.5), to explore its generalization towards $n \geq 8$. Recall the N^2 MHV full cells along with their growing parameters are given by

$$G_{7,0} = \left\{ \begin{array}{c} (45)(71) \\ [5] \end{array} \right\} \quad (5) \quad (3.1)$$

$$G_{7,1} = (23) \left\{ \begin{array}{c} (67) \\ (45) \end{array} \right\} \quad (6, 4) \quad (3.2)$$

$$G_{8,1} = \left\{ \begin{array}{c} (234)_2(678)_2 \\ (456)_2(781)_2 \\ (23)(456)_2(81) \end{array} \right\} \quad \begin{array}{c} (7, 4) \\ (7, 5) \\ (6, 4) \end{array} \quad (3.3)$$

$$G_{9,2} = \left\{ \begin{array}{c} (2345)_2(6789)_2 \\ (23)(4567)_2(891)_2 \end{array} \right\} \quad \begin{array}{c} (8, 6, 4) \end{array} \quad (3.4)$$

and for notational convenience, below we will suppress subscript '2' for consecutive vanishing 2×2 minors, such as $(234)_2 \equiv (234) = (23)(34)$, which is unambiguous as we restrict the discussion to the N^2 MHV sector from now on. To understand how the full cells capture the information of amplitudes up to any number of external particles [8], let's give a brief review first.

According to the simplex-like structures of amplitudes similar to (1.3), for the general Y_n^k we have

and only $I_{i,1}$ in the bottom row needs to be identified, while $I_{i,1+j}$ can be obtained by performing a partial cyclic shift $i \rightarrow i+j$ except that label 1 is fixed, for all BCFW cells within $I_{i,1}$.

In the case of $k=2$, $I_{i,1}$ can be expressed in terms of $G_{i,m}$ above as

$$\begin{aligned} I_{7,1} &= G_{7,0} + G_{7,1}, \\ I_{8,1} &= G_{8,1} + (G_{7,0,2} + G_{7,1,2}), \\ I_{9,1} &= G_{9,2} + G_{8,1,2} + (G_{7,0,3} + G_{7,1,3}), \\ I_{10,1} &= G_{9,2,2} + G_{8,1,3} + (G_{7,0,4} + G_{7,1,4}), \\ I_{11,1} &= G_{9,2,3} + G_{8,1,4} + (G_{7,0,5} + G_{7,1,5}), \end{aligned} \quad (3.6)$$

and so on, where $G_{i,m}$ is purely made of full cells with growing mode m , namely these full cells have $(m+1)$ growing parameters, and the additional label l in $G_{i,m,l}$ represents its level during the simplex-like growth. Explicitly, up to level 3, $G_{7,0} \rightarrow G_{7,0,2} \rightarrow G_{7,0,3}$ of 0-mode is given by

$$\left\{ \begin{array}{c} (45)(71) \\ [5] \end{array} \right\} \rightarrow [5] \left\{ \begin{array}{c} (46)(81) \\ [6] \end{array} \right\} \rightarrow [56] \left\{ \begin{array}{c} (47)(91) \\ [7] \end{array} \right\} \quad (3.7)$$

and $G_{7,1} \rightarrow G_{7,1,2} \rightarrow G_{7,1,3}$ of 1-mode is given by

$$\begin{aligned} (23) \left\{ \begin{array}{c} (67) \\ (45) \end{array} \right\} &\rightarrow [6](23) \left\{ \begin{array}{c} (78) \\ (45) \end{array} \right\} \rightarrow [67](23) \left\{ \begin{array}{c} (89) \\ (45) \end{array} \right\} \\ &\rightarrow [4](23) \left\{ \begin{array}{c} (78) \\ (56) \end{array} \right\} \rightarrow [47](23) \left\{ \begin{array}{c} (89) \\ (56) \end{array} \right\} \\ &\rightarrow [45](23) \left\{ \begin{array}{c} (89) \\ (67) \end{array} \right\} \end{aligned} \quad (3.8)$$

and it is similar for $G_{8,1}$, finally $G_{9,2} \rightarrow G_{9,2,2} \rightarrow G_{9,2,3}$ of 2-mode is given by

$$Y_n^k = \left(\begin{array}{c|cccccc} [234 \dots n-k-3] & & & & & 1 \\ \vdots & & & & & \vdots \\ [234] & & & & 1 & \dots & I_{n-3,4} \\ [23] & & & 1 & I_{k+5,3} & \dots & I_{n-2,3} \\ [2] & & 1 & I_{k+5,2} & I_{k+6,2} & \dots & I_{n-1,2} \\ 1 & 1 & I_{k+5,1} & I_{k+6,1} & I_{k+7,1} & \dots & I_{n,1} \\ \hline & [k+5 \ k+6 \ k+7 \dots n] & [k+6 \ k+7 \dots n] & [k+7 \dots n] & [\dots n] & \dots & 1 \end{array} \right), \quad (3.5)$$

$$\begin{aligned}
 & \left\{ \begin{array}{l} (2345)(6789) \\ (23)(4567)(891) \end{array} \right\} \rightarrow [8] \left\{ \begin{array}{l} (2345)(679\ 10) \\ (23)(4567)(9\ 10\ 1) \end{array} \right\} \rightarrow [89] \left\{ \begin{array}{l} (2345)(67\ 10\ 11) \\ (23)(4567)(10\ 11\ 1) \end{array} \right\} \\
 & \rightarrow [6] \left\{ \begin{array}{l} (2345)(789\ 10) \\ (23)(4578)(9\ 10\ 1) \end{array} \right\} \rightarrow [69] \left\{ \begin{array}{l} (2345)(78\ 10\ 11) \\ (23)(4578)(10\ 11\ 1) \end{array} \right\} \\
 & \rightarrow [4] \left\{ \begin{array}{l} (2356)(789\ 10) \\ (23)(5678)(9\ 10\ 1) \end{array} \right\} \rightarrow [67] \left\{ \begin{array}{l} (2345)(89\ 10\ 11) \\ (23)(4589)(10\ 11\ 1) \end{array} \right\} \\
 & \rightarrow [49] \left\{ \begin{array}{l} (2356)(78\ 10\ 11) \\ (23)(5678)(10\ 11\ 1) \end{array} \right\} \\
 & \rightarrow [47] \left\{ \begin{array}{l} (2356)(89\ 10\ 11) \\ (23)(5689)(10\ 11\ 1) \end{array} \right\} \\
 & \rightarrow [45] \left\{ \begin{array}{l} (2367)(89\ 10\ 11) \\ (23)(6789)(10\ 11\ 1) \end{array} \right\}
 \end{aligned} \quad (3.9)$$

from which we can see that, knowing all full cells with their growing parameters at $n = 10$ is sufficient for generating Y_n^2 up to any n . Note the full cells (or fully-spanning cells) are named such that none of their i columns are removed when they first show up in $I_{i,1}$, while [5] in $G_{7,0}$ is an exception, as it is a descendent of the N^2 MHV top cell at $n = 6$ but put together with (45)(71) for convenience.

Given the summary above, the N^2 MHV $n = 8$ amplitude is

$$Y_8^2 = S_{8,2} + S_{8,1} + S_{8,0}, \quad (3.10)$$

where we have separated the terms containing 2, 1, 0 empty slots respectively by defining

$$\begin{aligned}
 S_{8,2} &= \begin{pmatrix} [23] \\ [28] [26] \\ [78] [58] [56] \end{pmatrix}, \\
 S_{8,1} &= \begin{pmatrix} [2](56)(81) \\ [2](34) \left\{ \begin{array}{l} (78) \\ (56) \end{array} \right\} \\ [8](45)(71) \quad [5](46)(81) \\ [8](23) \left\{ \begin{array}{l} (67) \\ (45) \end{array} \right\} \quad [6](23) \left\{ \begin{array}{l} (78) \\ (45) \end{array} \right\} \\ [4](23) \left\{ \begin{array}{l} (78) \\ (56) \end{array} \right\} \end{pmatrix}, \\
 S_{8,0} &= \left\{ \begin{array}{l} (234)(678) \\ (456)(781) \\ (23)(456)(81) \end{array} \right\},
 \end{aligned} \quad (3.11)$$

as indicated by the second subscript i of $S_{n,i}$. Now the cyclicity of Y_8^2 is separated into three parts:

$$\begin{aligned}
 Y_8^2 - Y_{8,+1}^2 &= (S_{8,2} - S_{8,2,+1}) + (S_{8,1} - S_{8,1,+1}) \\
 &\quad + (S_{8,0} - S_{8,0,+1}),
 \end{aligned} \quad (3.12)$$

where the third subscript ‘+1’ similarly denotes the cyclic shift. Straightforwardly we find

$$S_{8,2} - S_{8,2,+1} = \left\{ \begin{array}{l} [8] (-\partial(23) - \partial(56) - \partial(71)) \\ [6] (-\partial(23) - \partial(57) - \partial(81)) \\ [3] (-\partial(24) - \partial(67) - \partial(81)) \end{array} \right\}_1, \quad (3.13)$$

and ‘ $|_1$ ’ denotes the truncation that only keeps terms containing one empty slot. For example, (2.10) can be separated as

$$\begin{aligned}
 \partial(23)|_1 &= -[2] + [3], \quad \partial(23)|_0 \\
 &= -(23)(45) + (23)(56) - (23)(67) + (23)(71).
 \end{aligned} \quad (3.14)$$

Knowing growing parameters (8, 6, 3) of $n = 7$ boundary generators, or identities $(-\partial(23) - \partial(56) - \partial(71))$, we can denote this result as (B_n stands for boundary generators first induced by the cyclicity of Y_n^2)

$$B_7 = -(23) - (56) - (71) \quad (8, 6, 3) \quad (3.15)$$

so that

$$\begin{aligned}
 S_{8,2} - S_{8,2,+1} &= (\partial B_7|_1)_1, \\
 S_{8,1} - S_{8,1,+1} &= (\partial B_7|_0)_1 + \partial B_8|_1, \\
 S_{8,0} - S_{8,0,+1} &= \partial B_8|_0,
 \end{aligned} \quad (3.16)$$

where $(\partial B_7|_1)_1$ denotes the counterpart of $\partial B_7|_1$ when n increases from 7 to 8, according to its simplex-like growing pattern. In this way, we can figure out $\partial B_8|_1$ (and hence B_8) via the second relation above,

namely (in the 5th and 8th lines below we have added $+ [4](56)(81)$ and $- [4](56)(81)$ respectively)

$$\begin{aligned}
 \partial B_8 |_1 &= (S_{8,1} - S_{8,1,+1}) - (\partial B_7 |_0)_1 \\
 &= -[7](81)(34) + [8](71)(34) - [1](78)(34) \\
 &\quad + [2](34)(56) - [3](24)(56) + [4](23)(56) \\
 &\quad + [2](34)(78) - [3](24)(78) + [4](23)(78) \\
 &\quad - [3](56)(81) + [4](56)(81) - [5](34)(81) \\
 &\quad + [6](34)(81) + [8](34)(56) - [1](34)(56) \\
 &\quad + [8](12)(56) - [1](82)(56) + [2](81)(56) \\
 &\quad - [5](67)(34) + [6](57)(34) - [7](56)(34) \\
 &\quad - [4](56)(81) + [5](46)(81) - [6](45)(81) \\
 &= \partial (+ (781)(34) - (234)(56) - (234)(78) \\
 &\quad + (34)(56)(81) + (567)(34) - (812)(56) \\
 &\quad + (456)(81)) |_1.
 \end{aligned} \tag{3.17}$$

After identifying B_8 (by trial and error), we can check the third relation above, as

$$\begin{aligned}
 \partial B_8 |_0 &= + (781)(234) - (781)(345) + (781)(34)(56) \\
 &\quad - (234)(567) + (234)(56)(78) - (234)(56)(81) \\
 &\quad - (234)(56)(78) + (234)(678) - (234)(781) \\
 &\quad - (34)(567)(81) - (34)(56)(781) + (34)(56)(812) \\
 &\quad + (234)(56)(81) \\
 &\quad + (567)(81)(34) - (567)(12)(34) + (567)(234) \\
 &\quad - (812)(34)(56) + (812)(456) - (812)(567) \\
 &\quad + (456)(781) - (456)(812) + (456)(81)(23) \\
 &= S_{8,0} - S_{8,0,+1}
 \end{aligned} \tag{3.18}$$

nicely obeys the required consistency. The $n = 8$ identities used above can be referred in appendix A and they can be classified into five distinct types.

Next, for the cyclicity of Y_9^2 we similarly have

$$\begin{aligned}
 S_{9,2} - S_{9,2,+1} &= (\partial B_7 |_0)_2 + (\partial B_8 |_1)_1, \\
 S_{9,1} - S_{9,1,+1} &= (\partial B_8 |_0)_1 + \partial B_9 |_1, \\
 S_{9,0} - S_{9,0,+1} &= \partial B_9 |_0,
 \end{aligned} \tag{3.19}$$

where the simplex-like growing patterns give

$$S_{9,1} = \begin{pmatrix} [2] \left\{ \begin{array}{l} (345)(789) \\ (567)(891) \\ (34)(567)(91) \end{array} \right. & \left\{ \begin{array}{l} [7](234)(689) \\ [4](235)(789) \\ [7](456)(891) \\ [5](467)(891) \end{array} \right. \\ [9] \left\{ \begin{array}{l} (234)(678) \\ (456)(781) \\ (23)(456)(81) \end{array} \right. & \left\{ \begin{array}{l} [6](23)(457)(91) \\ [4](23)(567)(91) \end{array} \right. \end{pmatrix},$$

$$S_{9,0} = \left\{ \begin{array}{l} (2345)(6789) \\ (23)(4567)(891) \end{array} \right. \tag{3.20}$$

And via the first relation above we can figure out $(\partial B_8 |_1)_1$, or the growing parameters of B_8 , as

$$\begin{aligned}
 B_8 &= 781(34) - (234)(56) - (234)(78) + (34)(56)(81) \quad (9, 7, 5, 3) \\
 &\quad + (567)(34) \quad (9, 8, 6, 3) \\
 &\quad - (812)(56) \quad (9, 6, 3) \\
 &\quad + (456)(81) \quad (9, 6, 4)
 \end{aligned} \tag{3.21}$$

so that via the second relation we can similarly figure out $\partial B_9 |_1$ (and hence B_9), and the third one again serves as a consistency check. Explicitly, we find

$$\begin{aligned}
 B_9 &= - (2345)(789) - (7891)(345) \\
 &\quad + (8912)(567) - (5678)(234) \\
 &\quad - (5678)(91)(34) + (5678)(12)(34) \\
 &\quad - (912)(34)(567) + (4567)(891).
 \end{aligned} \tag{3.22}$$

Following exactly the same logic, for the cyclicity of Y_{10}^2 we have

$$\begin{aligned}
 S_{10,2} - S_{10,2,+1} &= (\partial B_8 |_0)_2 + (\partial B_9 |_1)_1, \\
 S_{10,1} - S_{10,1,+1} &= (\partial B_9 |_0)_1 + \partial B_{10} |_1, \\
 S_{10,0} - S_{10,0,+1} &= \partial B_{10} |_0,
 \end{aligned} \tag{3.23}$$

where the simplex-like growing patterns give

$$S_{10,1} = \begin{pmatrix} [2] \left\{ \begin{array}{l} (3456)(789 \ 10) \\ (34)(5678)(9 \ 10 \ 1) \end{array} \right. & [8] \left\{ \begin{array}{l} (2345)(679 \ 10) \\ (23)(4567)(9 \ 10 \ 1) \\ (2345)(789 \ 10) \\ (23)(4578)(9 \ 10 \ 1) \end{array} \right. \\ [10] \left\{ \begin{array}{l} (2345)(6789) \\ (23)(4567)(891) \end{array} \right. & [4] \left\{ \begin{array}{l} (2356)(789 \ 10) \\ (23)(5678)(9 \ 10 \ 1) \end{array} \right. \end{pmatrix}, \tag{3.24}$$

note that $S_{10,0} = 0$ since there is no new full cell at $n \geq 10$ [8], and hence $\partial B_{10} |_0 = 0$. Explicitly, we find

$$\begin{aligned}
 B_9 &= - (2345)(789) - (7891)(345) \quad (10, 8, 5, 3) \\
 &\quad + (8912)(567) - (5678)(234) - (5678)(91)(34) \\
 &\quad + (5678)(12)(34) \quad (10, 8, 6, 3) \\
 &\quad - (912)(34)(567) \quad (10, 7, 5, 3) \\
 &\quad + (4567)(891) \quad (10, 7, 5)
 \end{aligned} \tag{3.25}$$

as well as

$$\begin{aligned}
 B_{10} &= + (789 \ 10 \ 1)(3456) - (23456)(789 \ 10) \\
 &\quad + (5678)(9 \ 10 \ 1 \ 2)(34) + (45678)(9 \ 10 \ 1 \ 2) \\
 &\quad - (45678)(9 \ 10 \ 1)(23).
 \end{aligned} \tag{3.26}$$

Finally, for the cyclicity of Y_{11}^2 we have

$$\begin{aligned} S_{11,2} - S_{11,2,+1} &= (\partial B_9 | 0)_2 + (\partial B_{10} | 1)_1, \\ S_{11,1} - S_{11,1,+1} &= (\partial B_{10} | 0)_1 + \partial B_{11} | 1, \\ S_{11,0} - S_{11,0,+1} &= \partial B_{11} | 0, \end{aligned} \quad (3.27)$$

and explicitly we find

$$\begin{aligned} B_{10} = & +(789\ 10\ 1)(3456) - (23456)(789\ 10) + (5678)(9\ 10\ 1\ 2)(34) \quad (11, 9, 7, 5, 3) \\ & +(45678)(9\ 10\ 1\ 2) - (45678)(9\ 10\ 1)(23) \quad (11, 9, 7, 5) \end{aligned} \quad (3.28)$$

which leads to $(\partial B_{10} | 0)_1 = 0$, and hence $\partial B_{11} | 1 = 0$. From $S_{11,0} = 0$ we also have $\partial B_{11} | 0 = 0$, therefore it is safe to conclude that $B_{11} = 0$. We can summarize all these intriguing results as

$$B_7 = -(23) - (56) - (71) \quad (8, 6, 3) \quad (3.29)$$

$$\begin{aligned} B_8 = & +(781)(34) - (234)(56) - (234)(78) + (34)(56)(81) \quad (9, 7, 5, 3) \\ & +(567)(34) \quad (9, 8, 6, 3) \\ & -(812)(56) \quad (9, 6, 3) \\ & +(456)(81) \quad (9, 6, 4) \end{aligned} \quad (3.30)$$

$$\begin{aligned} B_9 = & -(2345)(789) - (7891)(345) \quad (10, 8, 5, 3) \\ & +(8912)(567) - (5678)(234) - (5678)(91)(34) + (5678)(12)(34) \quad (10, 8, 6, 3) \\ & -(912)(34)(567) \quad (10, 7, 5, 3) \\ & +(4567)(891) \quad (10, 7, 5) \end{aligned} \quad (3.31)$$

$$\begin{aligned} B_{10} = & +(789\ 10\ 1)(3456) - (23456)(789\ 10) + (5678)(9\ 10\ 1\ 2)(34) \quad (11, 9, 7, 5, 3) \\ & +(45678)(9\ 10\ 1\ 2) - (45678)(9\ 10\ 1)(23) \quad (11, 9, 7, 5) \end{aligned} \quad (3.32)$$

which terminate at $n = 10$ like the full cells. With B_7, B_8, B_9, B_{10} and the growing parameters of relevant boundary generators identified, the cyclicity of Y_n^2 for any n is proved. These identities are classified into 1, 5, 6, 4 distinct types with respect to $n = 7, 8, 9, 10$ in Appendix 1.

A final remark is, not all N^2 MHV homological identities are required for this proof. Especially, those involving the quadratic cell at $n = 8$, namely $(12)(34)(56)(78)$, or the composite-linear cell at $n = 9$, namely $(123)(456)(789)$, are irrelevant. These two non-BCFW-like cells will lead to extra non-unity factors along with the 5-brackets [8], which cannot be generated by recursion. Therefore it is desirable to find that they do not appear at all in the proof of cyclicity for N^2 MHV amplitudes, not even appear as canceling pairs in the intermediate steps.

Since the cyclicity of tree amplitudes can be divided into many sub-equalities in terms of homological identities, the latter in fact have some kind of invariant meaning if we reshuffle an identity as “terms with plus signs = terms with minus signs”. It is definitely an interesting and geometrically profound direction to explore this more refined invariance in

the context of amplituhedron and sign flips [13, 14]. Especially, the $k \geq 2$ identities should give nontrivial insights on the connection between positive Grassmannian and the sign-flip triangulation. Of course, the cyclicity at higher k will demand us to work out identities up to $n = 4k + 2$ at least, for example $n = 14$ for $k = 3$, which is a straightforward but very lengthy task, since a more transparent pattern for all k 's still awaits to be found.

Data Availability Statement This manuscript has no associated data or the data will not be deposited. [Authors' comment: This is a theoretical work that does not require data.]

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A Relevant N^2 MHV homological identities

Below we list all distinct N^2 MHV homological identities that are relevant in this work. Note that we have discarded boundary cells that fail to have kinematical supports in terms of momentum twistors, but still we abuse the term “homological” here while the actual kinematics also matters [8, 9].

$n=7$

$$\partial(12) = -[1] + [2] - (12)(34) + (12)(45) - (12)(56) + (12)(67). \quad (\text{A.1})$$

$n=8$

$$\begin{aligned} \partial(123)(45) = & -[1](23)(45) + [2](13)(45) \\ & -[3](12)(45) + (123)(456) \\ & - (123)(45)(67) + (123)(45)(78). \end{aligned} \quad (\text{A.2})$$

$$\begin{aligned} \partial(123)(56) = & -[1](23)(56) + [2](13)(56) \\ & -[3](12)(56) + (123)(456) \\ & - (123)(567) + (123)(56)(78). \end{aligned} \quad (\text{A.3})$$

$$\begin{aligned} \partial(123)(67) = & -[1](23)(67) + [2](13)(67) \\ & -[3](12)(67) + (123)(45)(67) \\ & - (123)(567) + (123)(678). \end{aligned} \quad (\text{A.4})$$

$$\begin{aligned} \partial(123)(78) = & -[1](23)(78) + [2](13)(78) \\ & -[3](12)(78) + (123)(45)(78) \\ & - (123)(56)(78) + (123)(678). \end{aligned} \quad (\text{A.5})$$

$$\begin{aligned} \partial(12)(34)(67) = & -[1](34)(67) + [2](34)(67) \\ & -[3](12)(67) + [4](12)(67) \\ & + [6](12)(34) - [7](12)(34) \\ & - (12)(345)(67) - (12)(34)(567) \\ & + (12)(34)(678) + (812)(34)(67). \end{aligned} \quad (\text{A.6})$$

$n=9$

$$\begin{aligned} \partial(1234)(567) = & -[1](234)(567) + [2](134)(567) \\ & -[3](124)(567) + [4](123)(567) \\ & - (1234)(5678) + (1234)(567)(89). \end{aligned} \quad (\text{A.7})$$

$$\begin{aligned} \partial(1234)(678) = & -[1](234)(678) + [2](134)(678) \\ & -[3](124)(678) + [4](123)(678) \\ & - (1234)(5678) + (1234)(6789). \end{aligned} \quad (\text{A.8})$$

$$\begin{aligned} \partial(1234)(789) = & -[1](234)(789) + [2](134)(789) \\ & -[3](124)(789) + [4](123)(789) \\ & - (1234)(56)(789) \\ & + (1234)(6789). \end{aligned} \quad (\text{A.9})$$

$$\begin{aligned} \partial(1234)(56)(89) = & -[1](234)(56)(89) + [2](134)(56)(89) \\ & -[3](124)(56)(89) + [4](123)(56)(89) \end{aligned}$$

$$\begin{aligned} & - (1234)(567)(89) \\ & + (1234)(56)(789). \end{aligned} \quad (\text{A.10})$$

$$\begin{aligned} \partial(1234)(67)(89) = & -[1](234)(67)(89) \\ & + [2](134)(67)(89) \\ & - [3](124)(67)(89) \\ & + [4](123)(67)(89) \\ & - (1234)(567)(89) \\ & + (1234)(6789). \end{aligned} \quad (\text{A.11})$$

$$\begin{aligned} \partial(123)(45)(678) = & -[1](23)(45)(678) \\ & + [2](13)(45)(678) \\ & - [3](12)(45)(678) \\ & + [4](123)(678) \\ & - [5](123)(678) \\ & + [6](123)(45)(78) \\ & - [7](123)(45)(68) \\ & + [8](123)(45)(67) \\ & + (123)(45)(6789) \\ & - (9123)(45)(678). \end{aligned} \quad (\text{A.12})$$

$n=10$

$$\begin{aligned} \partial(12345)(6789) = & -[1](2345)(6789) \\ & + [2](1345)(6789) \\ & - [3](1245)(6789) \\ & + [4](1235)(6789) \\ & - [5](1234)(6789) \\ & + (12345)(6789 \ 10). \end{aligned} \quad (\text{A.13})$$

$$\begin{aligned} \partial(12345)(789 \ 10) = & -[1](2345)(789 \ 10) \\ & + [2](1345)(789 \ 10) \\ & - [3](1245)(678 \ 10) \\ & + [4](1235)(789 \ 10) \\ & - [5](1234)(789 \ 10) \\ & + (12345)(6789 \ 10). \end{aligned} \quad (\text{A.14})$$

$$\begin{aligned} \partial(12345)(678)(9 \ 10) = & -[1](2345)(678)(9 \ 10) \\ & + [2](1345)(678)(9 \ 10) \\ & - [3](1245)(678)(9 \ 10) \\ & + [4](1235)(678)(9 \ 10) \\ & - [5](1234)(678)(9 \ 10) \\ & + (12345)(6789 \ 10). \end{aligned} \quad (\text{A.15})$$

$$\begin{aligned} \partial(1234)(5678)(9 \ 10) = & -[1](234)(5678)(9 \ 10) \\ & + [2](134)(5678)(9 \ 10) \\ & - [3](124)(5678)(9 \ 10) \\ & + [4](123)(5678)(9 \ 10) \\ & - [5](1234)(678)(9 \ 10) \end{aligned}$$

$$\begin{aligned}
&+[6](1234)(578)(9\ 10) \\
&-[7](1234)(568)(9\ 10) \\
&+[8](1234)(567)(9\ 10) \\
&-[9](1234)(5678) \\
&+[10](1234)(5678). \quad (A.16)
\end{aligned}$$

References

1. N. Arkani-Hamed, J.L. Bourjaily, F. Cachazo, A.B. Goncharov, A. Postnikov, J. Trnka, Scattering amplitudes and the positive Grassmannian, [arXiv:1212.5605](#) [hep-th]
2. N. Arkani-Hamed, F. Cachazo, C. Cheung, The Grassmannian origin of dual superconformal invariance. *JHEP* **1003**, 036 (2010). [arXiv:0909.0483](#) [hep-th]
3. R. Britto, F. Cachazo, B. Feng, E. Witten, Direct proof of tree-level recursion relation in Yang-Mills theory. *Phys. Rev. Lett.* **94**, 181602 (2005). [arXiv:hep-th/0501052](#)
4. N. Arkani-Hamed, J.L. Bourjaily, F. Cachazo, S. Caron-Huot, J. Trnka, The all-loop integrand for scattering amplitudes in planar $N=4$ SYM. *JHEP* **1101**, 041 (2011). [arXiv:1008.2958](#) [hep-th]
5. J.L. Bourjaily, positroids, plabic graphs, and scattering amplitudes in mathematica, [arXiv:1212.6974](#) [hep-th]
6. H. Elvang, Y.T. Huang, Scattering amplitudes, [arXiv:1308.1697](#) [hep-th]
7. J.M. Henn, J.C. Plefka, *Scattering Amplitudes in Gauge Theories* (Springer, Berlin, 2014)
8. J. Rao, Positivity, Grassmannian geometry and simplex-like structures of scattering amplitudes. *JHEP* **1712**, 147 (2017). [arXiv:1609.08627](#) [hep-th]
9. J. Rao, Simplex-like structures of maximally supersymmetric scattering amplitudes, [arXiv:1712.10000](#) [hep-th]
10. J.M. Drummond, J.M. Henn, All tree-level amplitudes in $N=4$ SYM. *JHEP* **0904**, 018 (2009). [arXiv:0808.2475](#) [hep-th]
11. N. Arkani-Hamed, J. Bourjaily, F. Cachazo, J. Trnka, Unification of residues and Grassmannian dualities. *JHEP* **1101**, 049 (2011). [arXiv:0912.4912](#) [hep-th]
12. N. Arkani-Hamed, J. Trnka, The Amplituhedron. *JHEP* **1410**, 030 (2014). [arXiv:1312.2007](#) [hep-th]
13. N. Arkani-Hamed, H. Thomas, J. Trnka, Unwinding the amplituhedron in binary. *JHEP* **1801**, 016 (2018). [arXiv:1704.05069](#) [hep-th]
14. R. Kojima, C. Langer, Sign flip triangulations of the Amplituhedron, [arXiv:2001.06473](#) [hep-th]
15. T. Lukowski, R. Moerman, Boundaries of the Amplituhedron with amplituhedron Boundaries, [arXiv:2002.07146](#) [hep-th]