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SPECIAL ISSUE: Optical Gain Materials towards Enhanced Light-Matter Interactions

Light-induced phase transition and photochromism in all-inorganic two-dimensional Cs₂PbI₂Cl₂ perovskite

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ABSTRACT Organic-inorganic hybrid two-dimensional (2D) ruddlesden-popper (RP) perovskites with fantastic optoelectronic properties and good stability have attracted tremendous attention for the potential applications in photovoltaics and electroluminescence. Recently, a new allinorganic Cs₂PbI₂Cl₂ 2D perovskite has been proposed with excellent excitonic absorption and improved ambient and thermal stability. Herein, an interesting light-induced phase transition and photochromism in the Cs₂PbI₂Cl₂ were reported. Under low fluence light pumping, the room temperature photoluminescence (PL) of Cs₂PbI₂Cl₂ is dominated by a weak violet excitonic emission peaked at 412 nm. Surprisingly, the emission color gradually changes from violet to bright red while increasing the laser pumping fluence. This photochromic effect is determined to be caused by forming CsPbI₃ phase within the Cs₂PbI₂Cl₂ crystals, which is efficiently driven by thermal energy. Due to protection by the Cs₂PbI₂Cl₂ matrix, the embedded CsPbI₃ nanocrystals show improved stability than standard pure CsPbI₃. Therefore, the Cs₂PbI₂Cl₂ perovskite with photochromic feature may find applications in optical encryption, as preliminarily shown in this work.

Keywords: $Cs_2PbI_2Cl_2$, light-induced phase transition, photochromism, optical encryption

INTRODUCTION

Metal-halide perovskites have demonstrated tremendous potentials in photovoltaics, light-emitting diodes, lasers and photodetectors for their excellent optoelectronic properties [1–9]. Among the lead-perovskites family, in recent years, cesium lead halide perovskites, due to their outstanding optoelectronic properties and improved sta-

bility, have attracted much more attention [10–14]. Based on the connection manner of the PbX₆ octahedra, in the aspect of dimensionality, cesium lead halide perovskites can be classified into three-dimensional (3D) CsPbX₃ (corner-shared octahedra, X is halogen), 2D Cs₂PbX₄ or CsPb₂X₅ (octahedra planes), 1D Cs₃PbX₅ (octahedra chains) and 0D Cs₄PbX₆ crystal structures. Among them, 2D ruddlesden-popper (RP) phase Cs₂PbX₄ owns an important and special position and gains more and more attention, due to its stable optoelectronic properties and substantially improved stability. Very recently, Li and coworkers [15] reported 2D RP phase Cs₂PbI₂Cl₂ bulk materials synthesized by a solid-state method and characterized their excellent in-plane UV-light photoresponse and markedly ambient and thermal stability. In addition, theoretical studies predicted a rather high, layer-dependent in-plane electron mobilities as high as $\sim 9.39 \times 10^{3} \text{ cm}^{2} \text{ V}^{-1} \text{ s}^{-1}$ in the 2D bilayer Cs₂PbI₂Cl₂ [16]. Then shortly after that, a strong excitonic absorption but with a weak photoluminescence (PL) quantum yield has been found in the photophysical research of Cs₂PbI₂Cl₂ nanoparticles [17,18]. Further studies showed that the strong PL quenching may be caused by the combination of iodide segregation and defective RP phase interfaces. However, noticeably, doping and post-synthetic chemical transformation of Cs₂PbI₂Cl₂ may provide application opportunities in the optoelectronic field [18,19].

In general, phase transition (or degeneration) occurs when the temperature reaches up to a specific value. This thermo-induced phase transition originates from the transformation or reconstruction of crystal structure. Theoretically, phase transition may occur even at room temperature if the raw material receives enough energy

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from incident light *via* the photothermal process [20]. Such a transition can be termed as a light-induced structural phase transition. For a semiconductor, light-induced structural phase transition often leads to modulation of morphological, electrical and optical properties [21–23]. Accordingly, such materials with light-induced phase transition may have great applications in signal transformation and optical encryption fields.

Herein, the light-induced phase transition features of all-inorganic perovskite 2D RP phase $Cs_2PbI_2Cl_2$ have been demonstrated for the first time. It shows that the room temperature PL spectra are dominated by weak violet (strong red) emission before (after) phase transition. According to X-ray diffraction (XRD) and transmission electron microscopy (TEM) measurements, the unexpected red emission band comes from reconstruction of CsPbI₃ nanocrystals inside the Cs₂PbI₂Cl₂. Due to the protection of Cs₂PbI₂Cl₂ composite shows an improved stability than pure CsPbI₃. Finally, optical encryption based on Cs₂PbI₂Cl₂ is preliminarily performed.

EXPERIMENTAL SECTION

Preparation of large sized Cs₂PbI₂Cl₂ microplates

The square Cs₂PbI₂Cl₂ microplates were synthesized by the antisolvent vapor-assisted slow crystallization method [24]. All the raw materials were purchased from Sigma-Aldrich (USA). According to stoichiometric ratio of Cs₂PbI₂Cl₂, 0.1 mmol CsI (grade 99.99%) and 0.05 mmol PbCl₂ (grade 99.99%) (2CsI+PbCl₂) were dissolved in 3 mL N,N-dimethylformamide (DMF, anhydrous, grade 99.99%) solvent, which is defined as precursor solution I. 0.1 mmol CsCl (grade 99.99%) and 0.05 mmol PbI₂ (grade 99.99%) (2CsCl+PbI₂) were dissolved in 1.5 mL DMF and 1.5 mL dimethyl sulfoxide (DMSO, anhydrous, grade 99.99%), which is defined as precursor solution II. A glass platform was immersed in the precursor solution I or II in a small bottle, which was placed in 3 mL antisolvent dichloromethane (DCM, anhydrous, grade 99.99%) contained in a beaker, and then the beaker was sealed. After 24 h at room temperature, all DCM had slowly diffused into the precursor solution, which drove the slow crystallization of Cs₂PbI₂Cl₂. Then, the glass platform was taken out and placed on the heating stage at 100°C for 5 s. The above method was used to prepare dense square Cs₂PbI₂Cl₂ microplates with sizes of 10–20 μ m. In addition, another method as follow was used to prepare sparse square Cs₂PbI₂Cl₂ microplates. The glass platform was dripped with 200 µL solution I or II and suspended above the 3 mL antisolvent DCM. After the growth was kept at room temperature for 24 h, the glass platform was taken out and placed on the heating stage at 100° C for 5 s.

Preparation of small sized $Cs_2PbI_2Cl_2$ microcrystalline thin film

The Cs₂PbI₂Cl₂ microcrystalline thin film was synthesized by the antisolvent dripping-assisted fast crystallization method [24]. 3 mL antisolvent DCM was dripped into precursor solution I or II in a bottle holding a glass platform. The obtained suspension was stirred for 10 s at 500 r s⁻¹, then stranded at room temperature. After 6 h, the glass platform was taken out and placed on the heating stage at 100°C for 5 s.

Characterizations

The morphologies and sizes of the obtained Cs₂PbI₂Cl₂ were characterized by scanning electron microscopy (SEM, Zeiss Zigma field emission scanning electron microscope, Germany). The elemental composition and distribution were measured by energy-dispersive X-ray spectroscopy (EDX, TEAM Apollo XL, USA). PL spectra of the samples were monitored by a micro Raman system (Horiba LABHRev-UV, USA) under 325 nm continuous wavelength (CW) laser excitation. Light-induced phase transition was realized by 325 nm CW laser excitation source. PL spectra of the heated Cs₂PbI₂Cl₂ at 550 K were collected by a PL3-211-P spectrometer (Horiba Jobin Yvon, USA) and a 450 W xenon lamp was used as the excitation source. Thermo-induced phase transition was realized by a homothermal heating stage. Absorption spectrum was measured by an ultraviolet visible nearinfrared spectrophotometer (UV3600) using BaSO₄ as a reference in the range of 200-800 nm. Phase compositions were tested by powder XRD (Rigaku Smartlab diffractometer, Japan) with Cu Ka radiation (λ =1.5405 Å) over the angular range of $10^{\circ} \le 2\theta \le 80^{\circ}$, operated at 45 kV and 200 mA. The scanning TEM image was obtained using an FEI Talos F200X (USA) at an accelerating voltage of 200 kV. The pulse laser-excited PL spectra were tested by a femtosecond amplified laser system (pulse width: 100 fs; repetition rate: 1000 Hz) with 400 nm excitation and emission light being coupled and collected through the SP-2300 spectrometer and CCD detector. All the above measurements were performed at room temperature.

RESULTS AND DISCUSSION

Herein, two types of Cs₂PbI₂Cl₂ microplates were syn-



Figure 1 (a) i. Synthesis schematic diagram of large square $Cs_2PbI_2Cl_2$ microplates. The precursor is $2CsI+PbCl_2$ or $2CsCI+PbI_2$. ii. Synthesis schematic diagram of small $Cs_2PbI_2Cl_2$ microplates. The DCM was injected at conditions of 500 r s⁻¹,10 s. The products were baked on the hot plate at 100°C for 5 s. (b) SEM image of large square $Cs_2PbI_2Cl_2$ microplates with size of 10–20 µm. (c) SEM image of small microplates with size of 1–2 µm, which is defined as a microcrystalline thin film in this work. (d) Linear absorption and PL spectra of $Cs_2PbI_2Cl_2$ microplates. The PL was monitored under low pump fluence of 325 nm CW laser.

thesized by the antisolvent-assisted crystallization method (Fig. 1a). Fig. 1b shows the top-view SEM image of the square Cs₂PbI₂Cl₂ microplates with typical edge lengths of 10-20 µm. These large sized microplates were obtained by antisolvent vapor-assisted slow crystallization method (Fig. 1a-i). Another type of Cs₂PbI₂Cl₂ microplates with sizes of as small as $1-2 \mu m$, looks like a thin film in the macroscopic scale (Fig. 1c). These small sized microplates were obtained by the antisolvent dripping-assisted fast crystallization method (Fig. 1a-ii). Small sized microcrystal thin film is denser than large sized one because the antisolvent dripping-assisted fast crystallization can induce many nucleation centers in a very short time. In addition, Cs₂PbI₂Cl₂ has the advantage of simplified preparation with low cost because it can be synthesized by different room-temperature solution-processed methods (Fig. S1). Fig. 1d shows the optical absorption and PL spectra of Cs₂PbI₂Cl₂ microplates. These microplates show a strong absorbance in the ultraviolet region with an absorption edge at 410 nm. Meanwhile, Cs₂PbI₂Cl₂ microplates show a violet emission peak at 412 nm under a low-level CW laser pumping. The optical bandgap of Cs₂PbI₂Cl₂ is estimated to be 3.01 eV from the linear extrapolation of Kubelka-Munk function (Fig. S2).

The XRD patterns of the as-grown microplates match well with that of the reported 2D RP phase Cs₂PbI₂Cl₂ single crystal (Fig. 2a) [15], which indicates that the obtained microplate is pure Cs₂PbI₂Cl₂ phase. EDX spectroscopy analysis shows that the molar ratio of Cl:I:Cs:Pb is in accordance with the stoichiometric ratio of Cs₂PbI₂Cl₂ (Fig. S3a), which further verifies the purity of the obtained Cs₂PbI₂Cl₂ microplates. Fig. 2b shows PL spectra of Cs₂PbI₂Cl₂ microplates under 325 nm CW laser excitation. Under low excitation intensity of 0.2 mW, only a violet emission peak at 412 nm can be observed. With the increase of excitation intensity, a new emission peak centered at about 637 nm unexpectedly emerges. The PL intensities of both two emission peaks increase with the increase of excitation intensity (Fig. 2c). However, the red emission band increases rapidly compared with the violet emission band. Contrastively, this red band emission was hardly observed under 150 fs pulsed 400 nm laser (Fig. S4). Compared with cold light source of pulse laser, CW laser is a hot light source that can induce thermal energy during the excitation/irradiation process. Thus, it can be speculated that the emergence of red emission band originates from light-induced thermal energy-driven phase transition.



Figure 2 (a) XRD patterns of the as-grown $Cs_2PbI_2Cl_2$ microplates, CW laser-irradiated $Cs_2PbI_2Cl_2$, 550 K-heated $Cs_2PbI_2Cl_2$ and reported $Cs_2PbI_2Cl_2$ [15], compared with standard cards of $CsPbI_3$, and Cs_4PbCl_6 . (b) PL spectra of a single $Cs_2PbI_2Cl_2$ microplate with increasing 325 nm CW laser excitation intensity. The top-left inset shows the amplifying PL spectra under low pumping levels. The bottom-left inset shows the PL image of the multiple microplates after high power CW laser excitation. The right insets show the optical images of a single $Cs_2PbI_2Cl_2$ microplate before and after high power CW laser excitation. The right insets show the optical images of 637 nm with the increase of CW laser excitation intensity. The inset shows the amplifying CW laser excitation power dependent PL intensity under low pumping levels. (d) PL spectra of the $Cs_2PbI_2Cl_2$ microplates before and after thermal treatments at 550 K under 325 nm excitation. The insets show the fluorescent images of unheated $Cs_2PbI_2Cl_2$ (right) under 325 nm xenon lamp excitation source.

To further verify the influence of thermal energy, the structural and optical properties of unheated and heated Cs₂PbI₂Cl₂ were investigated. As can be seen, it shows that heating at 550 K also can lead to the emergence of the red emission band (Fig. 2d). As shown in Fig. 2a, compared with original diffraction peaks, the XRD patterns of CW laser-irradiated and 550 K-heated Cs₂PbI₂Cl₂ exhibit many new similar diffraction peaks, which indicates the destruction of the crystal structure under thermal energy. Therefore, the light-induced phase transition in Fig. 2b originates from the CW laser source induced thermal energy-driven changes of crystal structure. This is not a complete phase transition because many diffraction peaks of Cs₂PbI₂Cl₂ can be clearly observed after light and thermal treatment. Partial new diffraction peaks match well with CsPbI₃ phase suggesting that thermal energy has resulted in the formation of CsPbI₃ phase. In spectroscopy characterizations, the characteristic emission peak of bulk CsPbI3 is around 700 nm according to previous publications [25-27]. To

accurately verify the origin of unexpected red emission around 640 nm, mixed CsCl and PbI₂ precursor solutionbased microcrystals were prepared. Experiments demonstrate that microcrystals obtained from CsCl and PbI₂ precursor solution are mixed phases of CsPbCl₃ and CsPbI₃ with two emission bands at about 502 and 700 nm (Fig. S5). This result suggests that CsPbI₂Cl phase is absent or instable. Thus, the new emerged red emission band should be ascribed to CsPbI₃ phase because only CsPbI₃ can exhibit red emission band among possibly new emerged phases. The PL intensity of new emerged red emission is much stronger than that of original violet emission, which may be due to that the PL quantum yield of the emerged CsPbI₃ phase is much higher than that of the original Cs₂PbI₂Cl₂ crystals [10,17,18,25]. These results indicate that Cs₂PbI₂Cl₂ has light-induced phase transition and photochromic features, as shown in the insets of Fig. 2b. After CW laser irradiation, very low power excitation intensity also can achieve its excitation and produce the red emission band (Fig. 2b), which

suggests the light/thermo-induced phase transition is irreversible. In addition, this CW laser excitation power dependent photochromic feature of Cs₂PbI₂Cl₂ enables its usage as a CW laser power detection material.

In order to further confirm the existence of CsPbI₃ phase in these microplates after light/thermo-induced phase transition in a figurative way, TEM analysis was used. Before CW laser irradiation, the high-resolution TEM image of the obtained Cs₂PbI₂Cl₂ shows a characteristic lattice spacing of 0.40 nm indexed to (110) crystal plane (Fig. 3a), which is in good agreement with previous studies [15,17]. After CW laser irradiation, the TEM result shows a characteristic Cs₂PbI₂Cl₂ lattice spacing of 0.40 nm [17], characteristic CsPbI₃ lattice spacings of 0.34 and 0.24 nm indexed to (015) and (200) crystal planes, respectively [28], and a much larger lattice spacing of 1.06 nm (Fig. 3b). The fringe with a significantly large lattice spacing of 1.06 nm is Moiré fringe, which indicates two overlapping lattice fringes between two crystals [24,29]. In general, a kind of small lattice fringes appear simultaneously with Moiré fringes, and other lattice lines with different orientations are not visible [30]. Herein, Moiré fringes in Fig. 3b result from the overlapping lattice fringes between Cs₂PbI₂Cl₂ and CsPbI₃ lattices. This is a direct evidence that demonstrates the formation of CsPbI₃ phase from the inside of Cs₂PbI₂Cl₂ crystals by heating or CW laser irradiation. According to the size limit of crystal lattices, the formed CsPbI₃ phase should be nanocrystals. The XRD patterns show that the new emerged CsPbI₃ phase preferentially aligned along the (023), (015), (016), (200) and (008) planes (Fig. 2a). Some main diffraction peaks typically observed in standard CsPbI₃ could not be clearly resolved here. Thus, the crystal orientation of the embedded CsPbI₃ phase should be limited by the original Cs₂PbI₂Cl₂ crystal. In addition, EDX analyses show that the content of the iodine element increases after high power CW laser excitation (Figs S3b and S6). Therefore, all the evidences include XRD patterns, red characteristic emission band, TEM Moiré fringes and EDX spectra have pointed to the formation of CsPbI₃ phase within Cs₂PbI₂Cl₂ crystals during the heating or ultraviolet CW laser irradiation process. Meanwhile, due to the quantum size effect, the red emission band is around 640 nm rather than the characteristic peak of 700 nm of bulk CsPbI₃ microcrystals. The broad full width at the half maximum of the red emission band can be ascribed to the inconsistent sizes of internal CsPbI₃ nanocrystals. Herein, we define the embedded CsPbI₃ nanocrystals in Cs₂PbI₂Cl₂ crystals as the CsPbI₃-in-Cs₂PbI₂Cl₂ composite.



Figure 3 (a) High-resolution TEM image of the synthetic $Cs_2PbI_2Cl_2$ microcrystal. (b) High-resolution TEM image of a $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ crystal shows the existent of Moiré fringe patterns with large lattice spacing of 1.06 nm. (c) Schematic diagram of the formation of $CsPbI_3$ nanocrystals. Bottom inserts show the fluorescent images of $Cs_2PbI_2Cl_2$ microcrystalline thin film before and after CW laser irradiation, respectively. (d) Normalized PL spectra of pure $CsPbI_3$ and $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ under 325 nm excitation. Left inset shows the optical image of $Cs_2PbI_2Cl_2$ microcrystalline thin film; right inset shows the optical image of $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystalline thin film; right inset shows the optical image of $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystalline thin film; right inset shows the optical image of $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystalline thin film; right inset shows the optical image of $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystalline thin film; right inset shows the optical image of $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystalline thin film; right inset shows the optical image of $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystalline thin film; right inset shows the optical image of $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystalline thin film with red emission along the "U" pattern.

Fig. 3c shows the schematic diagram of the formation of CsPbI₃ nanocrystals. The 325 nm CW laser irradiation will result in the formation of thermal energy that drives the chemical reaction and ionic migration at internal crystal lattice of Cs₂PbI₂Cl₂. Based on the comparisons of XRD patterns in Fig. 2a, many XRD diffraction peaks of CW laser-irradiated and 550 K-heated Cs₂PbI₂Cl₂ match well with the standard cards of CsPbI3 and Cs4PbCl6. Thus, in view of the emergence of red emission band, the dominant chemical reaction is Cs₂PbI₂Cl₂=2/3CsPbI₃+ $1/3Cs_4PbCl_6$ (I) during the phase transition process. The previous study has calculated that decomposition pathway I is thermodynamically favorable because of the positive decomposition energy [16]. After CW laser irradiation or 550 K heating, a few XRD diffraction peaks match well with other phases including PbCl₂, CsI, PbI₂, CsCl, CsPbCl₃ and Cs₄PbI₆, suggesting that some side reactions occurred during the light/thermo-induced degradation in Cs₂PbI₂Cl₂ (Fig. S7). Hence, the red emissive CsPbI₃ nanocrystal is one of the main products resulting from the light/thermo-induced phase transition. Contrastively, pulse laser excitation hardly leads to the phase transition/red emission (Fig. S4) since very less thermal energy has been produced during the excitation process.

The optical images in Fig. 3c show that very weak and almost invisible violet emission turned into bright and macroscopic red emission after CW laser processing in Cs₂PbI₂Cl₂ microcrystalline thin film. Significantly, the PL intensity of the new emerged red emission band is much larger than the intrinsic violet emission of Cs₂PbI₂Cl₂ (Figs 2b and 3d). This photochromic ability of Cs₂PbI₂Cl₂ enables it to be used as an optical image encryption material. For instance, ultraviolet CW laser processed patterns in Cs₂PbI₂Cl₂ thin films can be used as anti-counterfeit labels. As shown in the right inset of Fig. 3d, red emission along the "U" pattern can be clearly observed, while other regions do not show red emission. Thus, the marked red emission based on the designed patterns in CsPbI₃-in-Cs₂PbI₂Cl₂ thin films can be used as specific signals to be quickly identified.

The time-dependent PL intensity of phase transition and stability of $Cs_2PbI_2Cl_2$ is another issue for the evaluation of its practical application. Fig. 4a shows timedependent PL intensity of $Cs_2PbI_2Cl_2$ during the phase transition process. The red emission becomes bright with the increase of irradiation time due to the increase of $CsPbI_3$ phase (insets of Fig. 4a). The PL intensity of the red emission band increases in the first 50 s and after that keeps an approximate constant with the increase of irradiation time within 50-70 s. Hence, with the formation of CsPbI₃ within the crystals, the mixture involving CsPbI₃-in-Cs₂PbI₂Cl₂ composite will reach a stable phase state during the irreversible phase transition [31]. Only CsPbI₃ can exhibit red emission among all the mixture phases. Further, the photostability of CsPbI₃-in-Cs₂PbI₂Cl₂ and pure CsPbI₃ microcrystals were evaluated, as shown in Fig. 4b. Compared with the original intensity, the PL intensity of CsPbI3-in-Cs2PbI2Cl2 can remain about 60% while pure CsPbI3 only remain about 20% after 45 min of sustaining illumination with 325 nm CW laser (Fig. 4c, d). This result demonstrates that CsPbI₃-in-Cs₂PbI₂Cl₂ has better stability than pure CsPbI₃ since in the CsPbI₃-in-Cs₂PbI₂Cl₂ composite, CsPbI₃ nanocrystals are embedded in Cs₂PbI₂Cl₂ matrix microcrystals that offer a natural barrier to prevent the degeneration of internal CsPbI3 nanocrystals.

CONCLUSIONS

In summary, the light- and thermo-induced phase transition and photochromic feature of Cs₂PbI₂Cl₂ have been



Figure 4 (a) Time-dependent PL intensity of $Cs_2PbI_2Cl_2$ under 325 nm CW laser excitation during the phase transition process. Four insets show the fluorescence images of $Cs_2PbI_2Cl_2$ after different CW laser irradiation times of 0, 20, 40, and 60 s, respectively. (b) Integrated emission intensity changes of $Cs_2PbI_2Cl_2$ and pure $CsPbI_3$ under constant CW laser excitation of 325 nm in ambient atmosphere. Evolution of PL spectra after different illumination times under 325 nm CW laser excitation at about 24°C, 40% relative humidity atmosphere of (c) $CsPbI_3$ -in- $Cs_2PbI_2Cl_2$ microcrystals and (d) pure $CsPbI_3$ microcrystals.

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studied. Antisolvent vapor/dripping-assisted slow/fast crystallization methods are developed to synthesize Cs₂PbI₂Cl₂ microplates/thin films. The intrinsic violet emission band of Cs₂PbI₂Cl₂ is around 412 nm with a wide band gap. Both CW laser irradiation and thermal energy can lead to a new unexpected red emission band around 640 nm. The XRD analysis, characteristic red emission band, TEM Moiré fringes and EDX spectra indicate that this red emission band originates from the formation of CsPbI₃ nanocrystals within Cs₂PbI₂Cl₂ crystals. Compared with pure CsPbI₃, CsPbI₃-in-Cs₂PbI₂Cl₂ shows better stability. The PL intensity of the red emission increases dramatically with the increase of CW laser excitation power at a certain range, which enables Cs₂PbI₂Cl₂ to be used as power detection materials. The photochromic feature of Cs₂PbI₂Cl₂ contributes to the development of anti-counterfeit labels and specific identification applications based on the designed patterns. These results suggest Cs₂PbI₂Cl₂ can serve as promising power detection and optical image encryption materials with excellent photochromic ability toward practical functional devices.

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Conflict of interest The authors declare that they have no conflict of interest.

Supplementary information Supporting data are available in the online version of the paper.



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全无机二维钙钛矿Cs₂PbI₂Cl₂中的光致相变和光 致变色

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摘要 有机-无机杂化二维ruddlesden-popper (RP)钙钛矿因为其优 异的光电性能和良好的稳定性,在太阳能电池和发光器件领域应 用前景巨大,而受到广泛关注.最近,人们提出一类新型的二维RP 相全无机钙钛矿Cs2PbI2Cl2材料,这类材料具有优异的激子吸收性 能以及良好的环境和热稳定性.本文报道了Cs2PbI2Cl2中有趣的光 致相变和光致变色现象.在低功率连续激光激发下,Cs2PbI2Cl2的 室温光致发光以412 nm处的微弱的紫色发射为主.有趣的是,随着 激发功率的增加,发射颜色逐渐由紫色变为明亮的红色.这种光致 变色起源于Cs2PbI2Cl2晶体内由热能驱动的光致相变而形成的 CsPbI3纳米晶.由于Cs2PbI2Cl2基体微晶的保护作用,内嵌的CsPbI3 纳米晶的稳定性相比于纯的CsPbI3有明显提高.而具有光致变色特 征的Cs2PbI2Cl2可能会在光学加密中得到应用,这在本项工作中得 到了初步证实.