### ORIGINAL PAPER - PRODUCTION ENGINEERING



# Fractional flow in radial flow systems: a study for peripheral waterflood

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**Abstract** The Buckley–Leverett displacement mechanism has been used to predict the performance of waterflood. With Buckley-Leverett method, oil recovery from waterflood is calculated and required water injection volume to achieve that oil recovery is estimated. This method does provide a very useful tool in waterflood design. Our experience in oil industry and a thorough literature review indicates that Buckley-Leverett method was used to analyze waterflood project directly without any adjustment based on the real reservoir and production situations. By doing so, errors are introduced into the analysis. Buckley-Leverett method assumed that displacement occurs in a linear system. This is true for some waterflood scenarios while for others it is not. For some waterflood scenarios, a radial system is more appropriate than a linear system. In this study, we investigated the fractional flow in a radial system and derived the solutions to predict the performance of water displacing oil in radial system. With this radial displacement model, design and prediction of waterflood can be achieved by Buckley-Leverett method and our model, whichever fits the waterflood pattern. Considering the fact that many waterflood scenarios follow radial displacement, our model is an important supplement to Buckley-Leverett method.

### List of symbols

A Flow area

c<sub>w</sub> Water compressibility

 $f_{\rm w}$  Water fraction

h Reservoir thickness

k Reservoir permeability

 $k_{\rm ro}$  Relative permeability to oil

 $k_{\rm rw}$  Relative permeability to water

P<sub>c</sub> Capillary pressure

 $P_{\rm d}$  Threshold pressure in capillary pressure curve

p Pressure

 $p_{\rm o}$  Oil pressure

 $p_{\rm w}$  Water pressure

 $p_{\rm wf}$  Flowing bottomhole pressure

 $q_{\rm o}$  Oil rate

 $q_{\rm t}$  Total liquid rate

 $q_{\mathrm{w}}$  Water rate

r Radius from center of wellbore

r<sub>D</sub> Dimensionless radius

r<sub>e</sub> Reservoir outer boundary radius

 $r_{
m f}$  Displacement front position in radial system

 $r_{S_{w}}$  Position of any water saturation in radial system

 $r_{\rm w}$  Wellbore radius

 $\Delta r$  Radius incremental

S<sub>w</sub> Water saturation

 $S_{\rm wi}$  Irreducible water saturation

T Temperature

 $\Delta t$  Time period

t Time

V Volume

 $x_{\rm f}$  Displacement front position in linear system

φ Porosity

 $\rho_{\rm w}$  Water density

 $\mu_{\rm o}$  Oil viscosity

 $\mu_{\rm w}$  Water viscosity

λ Rock property parameter related to the distribution of pore sizes



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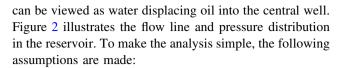
### Introduction

When reservoir engineers analyze the waterflood performance, they resort to the conventional frontal advance theory of Buckley and Leverett (1942). With Buckley-Leverett method, oil recovery from waterflood is calculated and required water injection volume to achieve that oil recovery is estimated. This method does provide a very useful tool in waterflood design. Welge (1952) proposed a tangent construction method to estimate the water saturation, water fraction at the water front and oil recovery factor. Several other investigators studied the multilayer reservoir waterflood performance. Stiles (1949) investigated the multilayer reservoir displacement by assuming the displacement velocity in a layer to be proportional to its absolute permeability. Dykstra and Parsons (1950) developed their famous multi-permeability model for noncommunicating layers without crossflow. Hearn (1971) derived expressions for communicating stratified reservoirs using the pseudo-relative permeability functions. El-Khatib (1999) advanced the closed form analytical solution for communicating stratified systems with log-normal permeability distributions. To the best of our knowledge, none of the study considered water displacing oil in a radial reservoir system.

It should be noted that Buckley-Leverett method and all the aforementioned studies assumed that displacement occurs in a linear system. This is true for some waterflood scenarios, while for others it is not. Our experience in oil industry and a thorough literature review indicates that petroleum engineers used Buckley-Leverett method to analyze the waterflood project directly without any adjustment based on the real reservoir and production situations such as production-injection patterns. By doing so, a lot of errors are introduced into the analysis. For some waterflood scenarios, a radial system is more appropriate than a linear system. Therefore, a radial displacement model is a necessary supplement to Buckley-Leverett linear displacement model. With both displacement models, design and prediction of waterflood can be achieved by selecting the appropriate model that fits the waterflood pattern. Considering the fact that many waterflood scenarios follow radial displacement, our model is very useful in field application.

## Derivation of fractional flow in a radial reservoir system

Figure 1 shows a circular reservoir with a well located in the center. For oil reservoir with strong peripheral water drive or surrounded by peripheral injectors, Fig. 1 can represent the displace procedure well. The fractional flow



- 1. A circular reservoir with constant height
- 2. Reservoir is homogenous in all rock properties
- 3. The dip angle of the formation is zero
- 4. Oil and water two-phase flow in reservoir, no gas presents in the reservoir
- 5. Compressibilities of oil and water are negligible
- 6. The variation in oil and water densities can be neglected
- 7. Constant reservoir temperature is applied
- 8. All rock properties do not change as pressure changes
- Constant oil and water viscosities during the displacement

Starting from Darcy's equation, we have oil and water flow rates to be calculated as:

$$q_{\rm o} = \frac{kk_{\rm ro}}{\mu_{\rm o}} \frac{\partial (Ap_{\rm o})}{\partial r} \tag{1}$$

$$q_{\rm w} = \frac{kk_{\rm rw}}{\mu_{\rm w}} \frac{\partial (Ap_{\rm w})}{\partial r} \tag{2}$$

where A is the flow area, k is the reservoir permeability,  $k_{\rm ro}$  is the relative permeability to oil,  $k_{\rm rw}$  is the relative permeability to water,  $p_{\rm o}$  is the oil pressure,  $p_{\rm w}$  is the water pressure,  $q_{\rm o}$  is the oil rate,  $q_{\rm w}$  is the water rate, r is the radius from wellbore,  $\mu_{\rm o}$  is the oil viscosity, and  $\mu_{\rm w}$  is the water viscosity.

Recalling the concept of capillary pressure we have

$$P_{\rm c} = p_{\rm o} - p_{\rm w} \tag{3}$$

where  $p_c$  is the capillary pressure.

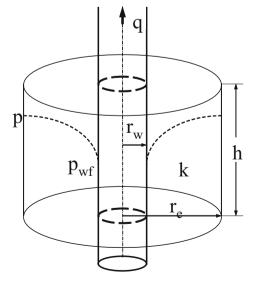


Fig. 1 A circular reservoir with a well located in the center



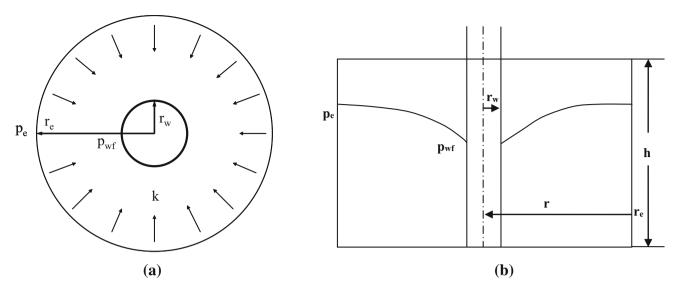


Fig. 2 Radial flow reservoir system: a plan view, b lateral view

Replacing water pressure by oil and capillary pressure Eq. (2) becomes

$$q_{\rm w} = \frac{kk_{\rm rw}}{\mu_{\rm w}} \frac{\partial [A(p_{\rm o} - p_{\rm c})]}{\partial r}. \tag{4}$$

Expressing in pressure gradient, Eqs. (1) and (4) are changed to

$$\frac{\partial (Ap_{\rm o})}{\partial r} = \frac{\mu_{\rm o}}{kk_{\rm ro}} q_{\rm o} \tag{5}$$

$$\frac{\partial (Ap_{\rm o})}{\partial r} - \frac{\partial (AP_{\rm c})}{\partial r} = \frac{\mu_{\rm w}}{kk_{\rm rw}} q_{\rm w}. \tag{6}$$

Subtracting Eq. (5) from (6), we obtain

$$-\frac{\partial (AP_{\rm c})}{\partial r} = \frac{\mu_{\rm w}}{kk_{\rm rw}}q_{\rm w} - \frac{\mu_{\rm o}}{kk_{\rm ro}}q_{\rm o}$$

or

$$\frac{\partial (AP_{\rm c})}{\partial r} = \frac{1}{k} \left( \frac{\mu_{\rm o}}{k_{\rm ro}} q_{\rm o} - \frac{\mu_{\rm w}}{k_{\rm rw}} q_{\rm w} \right).$$

Fig. 3 A control volume in a circular reservoir with a well located in the center

At this stage, we can introduce the concepts of total liquid rate and fractional flow, which are defined as:

$$q_{\rm t} = q_{\rm o} + q_{\rm w} \tag{8}$$

$$f_{\rm w} = \frac{q_{\rm w}}{q_{\rm t}} \tag{9}$$

where  $q_t$  is the total liquid rate and  $f_w$  is the water fraction. Substituting Eqs. (8) and (9) into (7) yields

$$f_{\rm W} = \frac{1 - \frac{\hat{o}(AP_{\rm c})}{\hat{o}r} \frac{kk_{\rm ro}}{q_{\rm t}\mu_{\rm o}}}{1 + \frac{k_{\rm ro}\mu_{\rm w}}{k_{\rm rw}\mu_{\rm o}}}.$$
 (10)

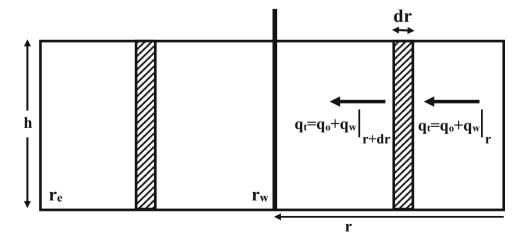
Flow area is defined as:

$$A = 2\pi rh \tag{11}$$

where h is the reservoir thickness.

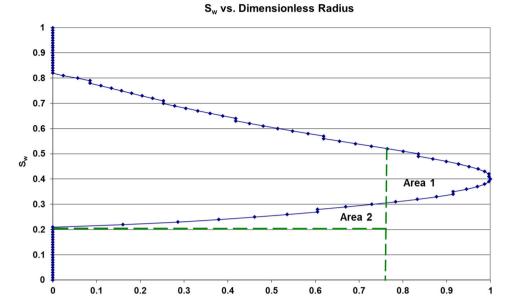
Substituting Eq. (11) into (10), we have

$$f_{\rm w} = \frac{1 - \frac{2\pi h k k_{\rm ro}}{q_{\rm t} \mu_{\rm o}} \left(\frac{r \hat{\rm o} P_{\rm c}}{\hat{\rm o} r} + P_{\rm c}\right)}{1 + \frac{k_{\rm ro} \mu_{\rm o}}{k_{\rm cr} \mu_{\rm o}}}.$$
(12)





**Fig. 4** The plot of water saturation versus dimensionless radius based on Eq. (31)



It should be noted that capillary pressure decreases as radius increases in this case. Therefore,  $\frac{\partial P_c}{\partial r}$  is negative. Comparing Eq. (12) with the fractional flow of linear displacement, we found that they are different. In linear displacement, the water saturation is calculated by (Buckley and Leverett 1942)

$$f_{\rm w} = \frac{1 - \frac{Akk_{\rm ro}}{q_1\mu_{\rm o}} \frac{\partial P_{\rm c}}{\partial x}}{1 + \frac{k_{\rm ro}\mu_{\rm w}}{k_{\rm rw}\mu_{\rm o}}}.$$
 (12a)

Therefore, linear displacement fractional flow equation cannot be used for radial displacement fractional flow. Equation (12) is the correct equation we should use in the radial system. If capillary pressure is negligible, Eq. (12) collapses into

$$f_{\rm w} = \frac{1}{1 + \frac{k_{\rm ro}\mu_{\rm w}}{k_{\rm rw}\mu_{\rm o}}}.$$
 (13)

Now, we derive the continuity equation of radial displacement. Considering the water displacing oil situation, material balance equation provides that the mass change in a control volume for a time period can be shown as Fig. 3. Since the flow direction is from reservoir outer boundary to the wellbore, for convenience, we define the reservoir boundary as the start point, where r=0, and the wellbore as the end point where  $r=r_{\rm e}$ . Therefore, in the dimensionless analysis, dimensionless radius can be defined as,  $r_{\rm D}=r/r_{\rm e}$ , the start point at the reservoir outer boundary will have  $r_{\rm D}=0$ , and the end point at the wellbore will have  $r_{\rm D}=1$ . Material balance gives us



$$[(q_{w}\rho_{w})_{r} - (q_{w}\rho_{w})_{r+\Delta r}]\Delta t = \pi h \{(r_{e} - r)^{2} - [r_{e} - (r + \Delta r)]^{2}\}\phi[(S_{w}\rho_{w})^{t+\Delta t} - (S_{w}\rho_{w})^{t}]$$
(14)

Wellbore

where  $S_{\rm w}$  is the water saturation,  $\rho_{\rm w}$  is the water density,  $\Delta t$  is the time period, t is the time,  $\Delta r$  is the radius incremental, r is the radius from reservoir outer boundary to wellbore,  $r_{\rm e}$  is the distance between reservoir outer boundary to wellbore, and  $\phi$  is the porosity.

Simplifying Eq. (14), we have

**Dimensionless Radius** 

$$[(q_{\mathbf{w}}\rho_{\mathbf{w}})_r - (q_{\mathbf{w}}\rho_{\mathbf{w}})_{r+\Delta r}]\Delta t = \pi h[2r_{\mathbf{e}}\Delta r - 2r\Delta r - (\Delta r)^2]\phi[(S_{\mathbf{w}}\rho_{\mathbf{w}})^{t+\Delta t} - (S_{\mathbf{w}}\rho_{\mathbf{w}})^t].$$
(15)

As  $\Delta r \rightarrow 0$  and  $\Delta t \rightarrow 0$ , we have

$$2r_{e}\Delta r - 2r\Delta r - (\Delta r)^{2} \approx 2r_{e}\Delta r - 2r\Delta r. \tag{16}$$

Equation (15) becomes partial differential equation

$$-\frac{\partial(q_{\rm w}\rho_{\rm w})}{\partial r} = (2r_{\rm e} - 2r)\pi h \phi \frac{\partial(S_{\rm w}\rho_{\rm w})}{\partial t}.$$
 (17)

Assuming constant density ("Appendix A" shows the derivation of governing equation including the change of fluid density), we have

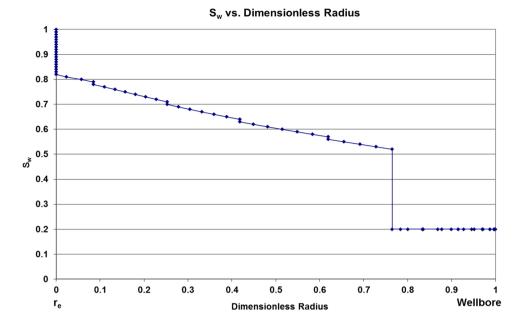
$$-\frac{\partial q_{\rm w}}{\partial r} = (2r_{\rm e} - 2r)\pi h \phi \frac{\partial S_{\rm w}}{\partial t}.$$
 (18)

Substituting Eq. (9) into (18) gives

$$-\frac{\partial (f_{\mathbf{w}}q_{\mathbf{t}})}{\partial r} = (2r_{\mathbf{e}} - 2r)\pi h \phi \frac{\partial S_{\mathbf{w}}}{\partial t}.$$
 (19)

If the water encroachment rate is constant, we have a constant total liquid rate. Equation (19) can be simplified to

**Fig. 5** The correct plot of water saturation versus dimensionless radius



$$-\frac{\partial f_{\rm w}}{\partial r} = \frac{(2r_{\rm e} - 2r)\pi h\phi}{q_{\rm t}} \frac{\partial S_{\rm w}}{\partial t}.$$
 (20)

Since water fraction is function of water saturation,  $f_{\rm w}(S_{\rm w})$ 

applying chain rule to partial differential equation results in

$$-\frac{\mathrm{d}f_{\mathrm{w}}}{\mathrm{d}S_{\mathrm{w}}}\frac{\partial S_{\mathrm{w}}}{\partial r} = \frac{(2r_{\mathrm{e}} - 2r)\pi h\phi}{q_{\mathrm{t}}}\frac{\partial S_{\mathrm{w}}}{\partial t}.$$
 (21)

At the first look Eq. (21) is similar to the Buckley–Leverett equation for linear displacement, we should notice that the term on the right-hand side before partial derivative is not constant.

Observing that water saturation is function of time, t, and position, r, we can express

$$dS_{w} = \frac{\partial S_{w}}{\partial t} dt + \frac{\partial S_{w}}{\partial r} dr.$$
 (22)

The fact that at the displacement front the water saturation is constant provides us a boundary condition.

 $\begin{tabular}{ll} \textbf{Table 1} & \textbf{The input data for water displacing oil in a radial system} \\ \end{tabular}$ 

Injection rate (BWPD)	50,000
Radius (ft)	5000
Reservoir thickness (ft)	50
Initial oil saturation (fraction)	0.8
Irreducible oil saturation (fraction)	0.2
Oil viscosity (cp)	0.93
Water viscosity (cp)	0.32
Porosity (fraction)	0.2

 Table 2
 The relative permeabilities and calculated parameters versus water saturation

S <sub>w</sub> (%)	S <sub>o</sub> (%)	$K_{\rm ro}$	$K_{\rm rw}$	f <sub>w</sub> (%)	df/dS <sub>w</sub>
0.0	100.0	0.4800	0.0000	0.0	
20.00	80.0	0.4800	0.0000	0.0	_
21.16	78.8	0.4594	0.0002	0.1	0.101
22.32	77.7	0.4393	0.0007	0.5	0.321
23.48	76.5	0.4196	0.0017	1.1	0.565
24.64	75.4	0.4005	0.0030	2.1	0.834
25.80	74.2	0.3819	0.0047	3.4	1.123
26.96	73.0	0.3638	0.0067	5.1	1.430
28.12	71.9	0.3461	0.0091	7.1	1.749
29.28	70.7	0.3290	0.0119	9.5	2.074
30.43	69.6	0.3123	0.0151	12.3	2.396
31.59	68.4	0.2961	0.0187	15.4	2.707
32.75	67.2	0.2804	0.0226	18.9	2.997
33.91	66.1	0.2651	0.0269	22.7	3.258
35.07	64.9	0.2504	0.0316	26.7	3.481
36.23	63.8	0.2361	0.0366	30.9	3.659
37.39	62.6	0.2222	0.0420	35.3	3.788
38.55	61.4	0.2088	0.0478	39.8	3.864
39.71	60.3	0.1959	0.0540	44.3	3.888
40.87	59.1	0.1835	0.0605	48.8	3.861
42.03	58.0	0.1715	0.0674	53.2	3.787
43.19	56.8	0.1599	0.0747	57.5	3.672
44.35	55.7	0.1488	0.0823	61.5	3.523
45.51	54.5	0.1381	0.0904	65.4	3.346
46.67	53.3	0.1279	0.0988	69.1	3.148
47.83	52.2	0.1181	0.1075	72.5	2.937
48.99	51.0	0.1087	0.1167	75.6	2.719
50.14	49.9	0.0998	0.1262	78.5	2.499



Table 2 continued

$S_{\rm w}~(\%)$	$S_{\rm o}~(\%)$	$K_{\rm ro}$	$K_{\mathrm{rw}}$	$f_{\rm w}~(\%)$	$\mathrm{d}f/\mathrm{d}S_{\mathrm{w}}$
51.30	48.7	0.0913	0.1361	81.2	2.282
52.46	47.5	0.0832	0.1464	83.6	2.071
53.62	46.4	0.0755	0.1570	85.7	1.868
54.78	45.2	0.0683	0.1680	87.7	1.677
55.94	44.1	0.0614	0.1794	89.4	1.497
57.10	42.9	0.0550	0.1912	91.0	1.331
58.26	41.7	0.0489	0.2033	92.3	1.177
59.42	40.6	0.0432	0.2158	93.5	1.036
60.58	39.4	0.0379	0.2287	94.6	0.907
61.74	38.3	0.0330	0.2420	95.5	0.791
62.90	37.1	0.0285	0.2556	96.3	0.686
64.06	35.9	0.0243	0.2696	97.0	0.591
65.22	34.8	0.0205	0.2840	97.6	0.507
66.38	33.6	0.0171	0.2987	98.1	0.431
67.54	32.5	0.0140	0.3138	98.5	0.364
68.70	31.3	0.0112	0.3293	98.8	0.304
69.86	30.1	0.0088	0.3452	99.1	0.251
71.01	29.0	0.0067	0.3615	99.4	0.205
72.17	27.8	0.0049	0.3781	99.6	0.164
73.33	26.7	0.0034	0.3951	99.7	0.128
74.49	25.5	0.0022	0.4124	99.8	0.097
75.65	24.3	0.0013	0.4302	99.9	0.070
76.81	23.2	0.0007	0.4483	99.9	0.047
77.97	22.0	0.0002	0.4668	100.0	0.028
79.13	20.9	0.0000	0.4856	100.0	0.013
80.29	19.7	0.0000	0.5000	100.0	0.002
81.45	18.6	0.0000	0.5000	100.0	_
82.61	17.4	0.0000	0.5000	100.0	_
83.77	16.2	0.0000	0.5000	100.0	_
84.93	15.1	0.0000	0.5000	100.0	_
86.09	13.9	0.0000	0.5000	100.0	_
87.25	12.8	0.0000	0.5000	100.0	_
88.41	11.6	0.0000	0.5000	100.0	_
89.57	10.4	0.0000	0.5000	100.0	_
90.72	9.3	0.0000	0.5000	100.0	_
91.88	8.1	0.0000	0.5000	100.0	_
93.04	7.0	0.0000	0.5000	100.0	_
94.20	5.8	0.0000	0.5000	100.0	_
95.36	4.6	0.0000	0.5000	100.0	_
96.52	3.5	0.0000	0.5000	100.0	_
97.68	2.3	0.0000	0.5000	100.0	_
98.84	1.2	0.0000	0.5000	100.0	_
100.00	0.0	0.0000	0.5000	100.0	_

$$\mathrm{d}S_{\mathrm{w}} = \frac{\partial S_{\mathrm{w}}}{\partial t} \mathrm{d}t + \frac{\partial S_{\mathrm{w}}}{\partial r} \mathrm{d}r = 0$$

or



**Table 3** Comparisons of locations of different water saturations at a waterflooding time for linear and radial systems

Water saturation	Location of water saturation in linear system (ft)	Location of water saturation in radial system (ft)
0.397	1389.7	751.3
0.409	1380.0	745.6
0.420	1353.7	730.2
0.432	1312.6	706.2
0.443	1259.2	675.2
0.455	1195.9	638.7
0.467	1125.4	598.5
0.478	1050.0	555.9
0.490	972.1	512.3
0.501	893.4	468.7
0.513	815.7	426.0
0.525	740.2	384.9
0.536	667.9	345.9
0.548	599.4	309.3
0.559	535.2	275.2
0.571	475.6	243.8
0.583	420.6	214.9
0.594	370.3	188.7
0.606	324.4	164.9
0.617	282.8	143.4
0.629	245.2	124.1
0.641	211.4	106.8
0.652	181.1	91.4
0.664	154.1	77.7
0.675	130.0	65.4
0.687	108.7	54.6
0.699	89.8	45.1
0.710	73.1	36.7
0.722	58.5	29.3
0.733	45.7	22.9
0.745	34.6	17.3
0.757	25.1	12.5
0.768	16.9	8.5
0.780	10.1	5.1
0.791	4.6	2.3
0.803	0.8	0.4
0.814	0.0	0.0

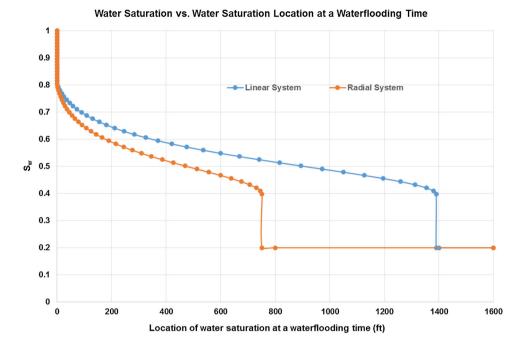
$$\frac{\partial S_{\mathbf{w}}}{\partial r} = -\frac{\partial S_{\mathbf{w}}}{\partial t} \frac{\mathrm{d}t}{\mathrm{d}r}.$$
 (23)

Substituting Eq. (23) into (21) yields

$$-\frac{\mathrm{d}f_{\mathrm{w}}}{\mathrm{d}S_{\mathrm{w}}}(-\frac{\partial S_{\mathrm{w}}}{\partial t}\frac{\mathrm{d}t}{\mathrm{d}r}) = \frac{(2r_{\mathrm{e}}-2r)\pi h\phi}{q_{\mathrm{t}}}\frac{\partial S_{\mathrm{w}}}{\partial t}$$

or

Fig. 6 Comparisons of locations of different water saturations at a waterflooding time for linear and radial systems



$$\frac{\mathrm{d}f_{\mathrm{w}}}{\mathrm{d}S_{\mathrm{w}}}\mathrm{d}t = \frac{(2r_{\mathrm{e}} - 2r)\pi h\phi}{q_{\mathrm{t}}}\mathrm{d}r. \tag{24}$$

Integrating Eq. (24) yields an equation for displacement front position,  $r_f$ .

$$\frac{\pi h \phi}{q_{\rm t}} (2r_{\rm e}r_{\rm f} - r_{\rm f}^2) = \left(\frac{\mathrm{d}f_{\rm w}}{\mathrm{d}S_{\rm w}}\right)_{\rm f} t$$

or

$$r_{\rm f}^2 - 2r_{\rm e}r_{\rm f} + \frac{tq_{\rm t}}{\pi h\phi} \left(\frac{\mathrm{d}f_{\rm w}}{\mathrm{d}S_{\rm w}}\right)_{\rm f} = 0 \tag{25}$$

where  $r_{\rm f}$  is the displacement front position in radial system. There are two solutions to Eq. (25), which are

$$r_{\rm f} = r_{\rm e} \pm \sqrt{r_{\rm e}^2 - \frac{tq_{\rm t}}{\pi h \phi} \left(\frac{\mathrm{d}f_{\rm w}}{\mathrm{d}S_{\rm w}}\right)_{\rm f}}.$$
 (26)

Obviously only one solution is correct to match with the physical phenomenon. Considering at the beginning of the displacement as  $t \to 0$ , we have  $r_f \to 0$ ; therefore, we can eliminate the solution

$$r_{\rm f} = r_{\rm e} + \sqrt{r_{\rm e}^2 - \frac{tq_{\rm t}}{\pi h \phi} \left(\frac{\mathrm{d}f_{\rm w}}{\mathrm{d}S_{\rm w}}\right)_{\rm f}}.$$
 (27)

Therefore, the correct solution is

$$r_{\rm f} = r_{\rm e} - \sqrt{r_{\rm e}^2 - \frac{tq_{\rm t}}{\pi h \phi} \left(\frac{\mathrm{d}f_{\rm w}}{\mathrm{d}S_{\rm w}}\right)_{\rm f}}.$$
 (28)

The distance between the wellbore and water front will be calculated by

**Table 4** Comparisons of water front location versus waterflooding time for linear and radial systems

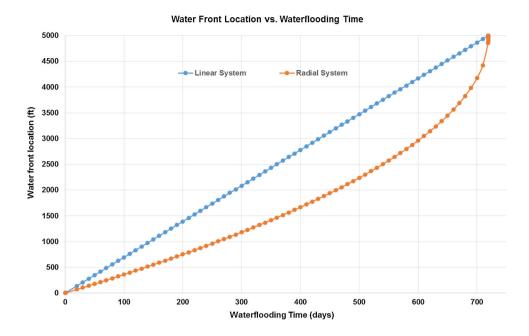
Waterflooding time (days)	Location of water front in linear system (ft)	Location of water front in radial system (ft)
0.1	0.7	0.3
1	6.9	3.5
20	139.0	70.0
30	208.5	105.3
50	347.4	176.8
100	694.8	360.4
150	1042.3	551.5
200	1389.7	751.3
250	1737.1	960.9
300	2084.5	1182.0
350	2431.9	1416.7
400	2779.3	1667.8
450	3126.8	1939.6
500	3474.2	2237.9
550	3821.6	2572.7
600	4169.0	2961.6
650	4516.4	3445.1
700	4863.8	4174.9
710	4933.3	4422.6
719.6	5000.0	5000.0

$$r_{\text{well-to-waterfront}} = r_{\text{e}} - r_{\text{f}} = \sqrt{r_{\text{e}}^2 - \frac{tq_{\text{t}}}{\pi h \phi} \left(\frac{\mathrm{d}f_{\text{w}}}{\mathrm{d}S_{\text{w}}}\right)_{\text{f}}}.$$
 (29)

Again, if we compare Eq. (28) with the linear displacement, we found that it is distinct from the



Fig. 7 Comparisons of water front location versus waterflooding time for linear and radial systems



displacement front of linear displacement system, which is

$$x_{\rm f} = \frac{tq_{\rm t}}{A\phi} \left(\frac{\mathrm{d}f_{\rm w}}{\mathrm{d}S_{\rm w}}\right)_{\rm f} \tag{30}$$

where  $x_f$  is the displacement front position in the linear system.

Therefore, linear displacement fractional flow equation cannot be used to locate the position of displacement front in a radial displacement system. Equation (28) is the correct equation we should use in the radial system.

For any water saturation,  $S_{\rm w}$ , the position can be calculated by

$$r_{S_{\rm w}} = r_{\rm e} - \sqrt{r_{\rm e}^2 - \frac{tq_{\rm t}}{\pi h \phi} \left(\frac{\mathrm{d}f_{\rm w}}{\mathrm{d}S_{\rm w}}\right)_{S_{\rm w}}} \tag{31}$$

where  $r_{S_{w}}$  is the position of any water saturation in radial system.

The plot of water saturation versus dimensionless radius based on Eq. (31) indicates that the plot needs to be modified to match with the physical model. Here, dimensionless radius is defined as:

$$r_{\rm D} = \frac{r}{r_{\rm e}}.\tag{32}$$

Figure 4 shows the plot of water saturation versus dimensionless radius based on Eq. (31).

Figure 4 gives two saturation values for the same position. Physically, it is impossible. The modification to Fig. 4 to get the correct water saturation distribution can be accomplished by determining displacement front position. To determine the displacement front location, one can

define a saturation discontinuity (or displacement front) at  $r_{\rm f}$  and balancing of the areas ahead of the front (Area 1) and below (Area 2) the saturation curve shown in Fig. 4. Then, the water saturation ahead of the displacement front should be the initial water saturation. The correct water saturation distribution is shown in Fig. 5.

### Case study

A case study was conducted to illustrate the analysis of radial water displacing oil procedure. The input data are shown in Table 1. Table 2 shows the relative permeabilities and calculated parameters versus water saturation.

The plot of water saturation versus dimensionless radius is shown in Fig. 5. If linear displacement equation is used, the position of the displacement front will be quite different. The processes of water displacing oil in linear and radial systems are compared to illustrate the difference. All inputs are the same in the comparison. The comparisons of locations of different water saturations at a waterflooding time for linear and radial systems are shown in Table 3 and Fig. 6. The comparisons of water front location at different waterflooding times for linear and radial systems are shown in Table 4 and Fig. 7. The differences in these tables and figures indicate that Buckley–Leverett linear displacement is not appropriate for peripheral waterflood reservoirs.

### **Conclusions**

The following conclusions can be drawn upon this study.



The performance of radial water displacing oil system is different from that of linear water displacing oil system.

If consider the effect of capillary pressure, equations used to calculate water fraction are different between linear and radial displacement systems. Linear displacement system equation cannot be used for the radial displacement system. Equation (12) should be used to estimate the water fraction.

Equations used to calculate the position of any water saturation are different between linear and radial displacement systems. Linear displacement system equation cannot be used for the radial displacement system. Equation (31) should be used to estimate the position of any water saturation.

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### Appendix A

Expanding Eq. (17) we have

$$-\rho_{\rm w} \frac{\partial q_{\rm w}}{\partial r} - q_{\rm w} \frac{\partial \rho_{\rm w}}{\partial r} = (2r_{\rm e} - 2r)\pi h \phi \left[ \rho_{\rm w} \frac{\partial S_{\rm w}}{\partial t} + S_{\rm w} \frac{\partial \rho_{\rm w}}{\partial t} \right]. \tag{33}$$

Applying chain rule we obtain

$$-\rho_{\rm w} \frac{\partial q_{\rm w}}{\partial r} - q_{\rm w} \frac{\partial \rho_{\rm w}}{\partial p_{\rm w}} \frac{\partial p_{\rm w}}{\partial r} = (2r_{\rm e} - 2r)$$

$$\pi h \phi \left[ \rho_{\rm w} \frac{\partial S_{\rm w}}{\partial t} + S_{\rm w} \frac{\partial \rho_{\rm w}}{\partial p_{\rm w}} \frac{\partial p_{\rm w}}{\partial t} \right].$$
(34)

Now introducing the concept of water compressibility, which is defined as

$$c_{\rm w} = -\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T = \frac{1}{\rho_{\rm w}} \left( \frac{\partial \rho_{\rm w}}{\partial p} \right). \tag{35}$$

Substituting Eq. (35) into (34) gives

$$-\rho_{\rm w} \frac{\partial q_{\rm w}}{\partial r} - q_{\rm w} c_{\rm w} \rho_{\rm w} \frac{\partial p_{\rm w}}{\partial r} = (2r_{\rm e} - 2r)$$

$$\pi h \phi \left[ \rho_{\rm w} \frac{\partial S_{\rm w}}{\partial t} + S_{\rm w} c_{\rm w} \rho_{\rm w} \frac{\partial p_{\rm w}}{\partial t} \right].$$
(36)

Substituting Eq. (9) into (36) yields

$$-\rho_{\rm w} \frac{\partial (f_{\rm w} q_{\rm t})}{\partial r} - f_{\rm w} q_{\rm t} c_{\rm w} \rho_{\rm w} \frac{\partial p_{\rm w}}{\partial r} = (2r_{\rm e} - 2r)$$

$$\pi h \phi \left[ \rho_{\rm w} \frac{\partial S_{\rm w}}{\partial t} + S_{\rm w} c_{\rm w} \rho_{\rm w} \frac{\partial p_{\rm w}}{\partial t} \right].$$
(37)

If the water encroachment rate is constant, we have a constant total liquid rate. Equation (37) can be simplified to

$$-\left[\frac{\partial f_{\mathbf{w}}}{\partial r} + f_{\mathbf{w}}c_{\mathbf{w}}\frac{\partial p_{\mathbf{w}}}{\partial r}\right] = \frac{(2r_{\mathbf{e}} - 2r)\pi h\phi}{q_{\mathbf{t}}}\left[\frac{\partial S_{\mathbf{w}}}{\partial t} + S_{\mathbf{w}}c_{\mathbf{w}}\frac{\partial p_{\mathbf{w}}}{\partial t}\right].$$
(38)

Expressing the pressure in terms of radius we have

$$p_{\rm w} = \frac{f_{\rm w} q_{\rm t} \mu_{\rm w}}{2\pi k k_{\rm rw} h} \ln \frac{r}{r_{\rm w}} + p_{\rm wf}.$$
 (39)

Taking the derivative of both sides of Eq. (39) with respect to r gives

$$\frac{\partial p_{\mathbf{w}}}{\partial r} = \frac{f_{\mathbf{w}} q_{\mathbf{t}} \mu_{\mathbf{w}}}{2\pi k k_{\mathbf{rw}} h} \frac{1}{r}.$$
(40)

For the same location (or radius in this case), water pressure change with respect to time can be approximated by capillary pressure change with respect to time,

$$\frac{\partial p_{\mathbf{w}}}{\partial t} = -\frac{\partial P_{\mathbf{c}}}{\partial t}.\tag{41}$$

According to Brooks and Corey (1964) capillary pressure model capillary pressure can be expressed as

$$P_{\rm c} = P_{\rm d} \left( \frac{S_{\rm w} - S_{\rm wi}}{1 - S_{\rm wi}} \right)^{\frac{-1}{\lambda}} \tag{42}$$

where,  $S_{\rm wi}$  = the irreducible water saturation,  $P_{\rm d}$  = the threshold pressure,  $\lambda$  = rock property parameter related to the distribution of pore sizes.

Brooks and Corey related the parameter  $\lambda$  to the distribution of pore sizes. For narrow distributions,  $\lambda$  is > 2; for wide distributions,  $\lambda$  is < 2.

Substituting Eq. (42) into (43) we have

$$\frac{\partial p_{\mathbf{w}}}{\partial t} = \frac{P_{\mathbf{d}}}{\lambda (1 - S_{\mathbf{w}i})} \left( \frac{S_{\mathbf{w}} - S_{\mathbf{w}i}}{1 - S_{\mathbf{w}i}} \right)^{\frac{-1}{\lambda} - 1} \frac{\partial S_{\mathbf{w}}}{\partial t}.$$
 (43)

Substituting Eqs. (40) and (43) into (38) yields

$$\frac{\partial f_{\mathbf{w}}}{\partial r} + f_{\mathbf{w}} c_{\mathbf{w}} \left( \frac{f_{\mathbf{w}} q_{\mathbf{t}} \mu_{\mathbf{w}}}{2\pi k k_{\mathbf{rw}} h} \frac{1}{r} \right) = \frac{(2r - 2r_{\mathbf{e}})\pi h \phi}{q_{\mathbf{t}}}$$

$$\left[ 1 + S_{\mathbf{w}} c_{\mathbf{w}} \frac{P_{\mathbf{d}}}{\lambda (1 - S_{\mathbf{w}i})} \left( \frac{S_{\mathbf{w}} - S_{\mathbf{w}i}}{1 - S_{\mathbf{w}i}} \right)^{\frac{-1}{\lambda} - 1} \right] \frac{\partial S_{\mathbf{w}}}{\partial t}.$$
(44)

Equation (44) can be solved numerically to obtain the location of any water saturation at any waterflooding time.



If the capillary pressure is small and can be ignored, Eq. (44) becomes

$$\frac{\partial f_{\rm w}}{\partial r} + f_{\rm w} c_{\rm w} \left( \frac{f_{\rm w} q_{\rm t} \mu_{\rm w}}{2\pi k k_{\rm rw} h} \frac{1}{r} \right) = \frac{(2r - 2r_{\rm e})\pi h \phi}{q_{\rm t}} \frac{\partial S_{\rm w}}{\partial t}. \tag{45}$$

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