

Meteoroid Layers in Planetary Atmospheres

J.G. Molina-Cuberos · J.J. López-Moreno · F. Arnold

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Abstract Metallic ions coming from the ablation of extraterrestrial dust, play a significant role in the distribution of ions in the Earth's ionosphere. Ions of magnesium and iron, and to a lesser extent, sodium, aluminium, calcium and nickel, are a permanent feature of the lower E-region. The presence of interplanetary dust at long distances from the Sun has been confirmed by the measurements obtained by several spacecrafts. As on Earth, the flux of interplanetary meteoroids can affect the ionospheric structure of other planets. The electron density of many planets show multiple narrow layers below the main ionospheric peak which are similar, in magnitude, to the upper ones. These layers could be due to long-lived metallic ions supplied by interplanetary dust and/or their satellites. In the case of Mars, the presence of a non-permanent ionospheric layer at altitudes ranging from 65 to 110 km has been confirmed and the ion $\text{Mg}^+\cdot\text{CO}_2$ identified. Here we present a review of the present status of observed low ionospheric layers in Venus, Mars, Jupiter, Saturn and Neptune together with meteoroid based models to explain the observations. Meteoroids could also affect the ionospheric structure of Titan, the largest Saturnian moon, and produce an ionospheric layer at around 700 km that could be investigated by Cassini.

Keywords Meteoroids · Ionosphere · Planetary atmosphere

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1 Introduction

The ablation of a continuous flux of extraterrestrial dust in the atmosphere gives rise to permanent layers of free neutral and ionized metal atoms in the 80–110 km altitude range. From the initial ground based observations of sodium at the end of the 1930s (Chapman 1938) to the in-situ measurements of metal ions by rocket-borne mass spectrometers (Grebowsky et al. 1998), the density profiles of metallic species as well as their latitude and temporal variability has been established. Magnesium and iron are the most abundant metal species in the atmosphere; others like sodium, aluminium, calcium and nickel are also present in a concentration, at least, one order of magnitude lower. The relative abundance of metal species in the atmosphere is roughly equal to the one exhibited in carbonaceous chondrites (Mason 1971).

In addition to general background of extraterrestrial dust, meteor showers, that are produced when the Earth crosses the dust stream left along a comet orbit, can increase the concentration of metals during short periods of time. The net mass influx to the Earth from each meteor shower is only a small fraction of the total yearly influx from the sporadic background (Hughes 1978). However, an increase by a factor of 2–3 in metallic concentration has been found during such events. The increase can be as much as one order of magnitude during a strong meteor shower (Kopp 1997; Grebowsky et al. 1998), which is often large enough to be manifested as a peak in the total ion and electron density profiles. Figure 1 shows two examples of the distribution of positive ions and electrons in the terrestrial atmosphere, where the increase in the metallic ion concentrations during the Perseid meteor shower is highlighted.

Dust detectors on board several space missions have observed that meteoroids are distributed through the whole interplanetary medium in the Solar System. Meteoroids can thus affect the atmospheric structure of other planets. In this paper we search for evidence of meteoroid layers in the atmospheres of extraterrestrial planets and we review the present knowledge of meteoroid modeling in the Solar System.

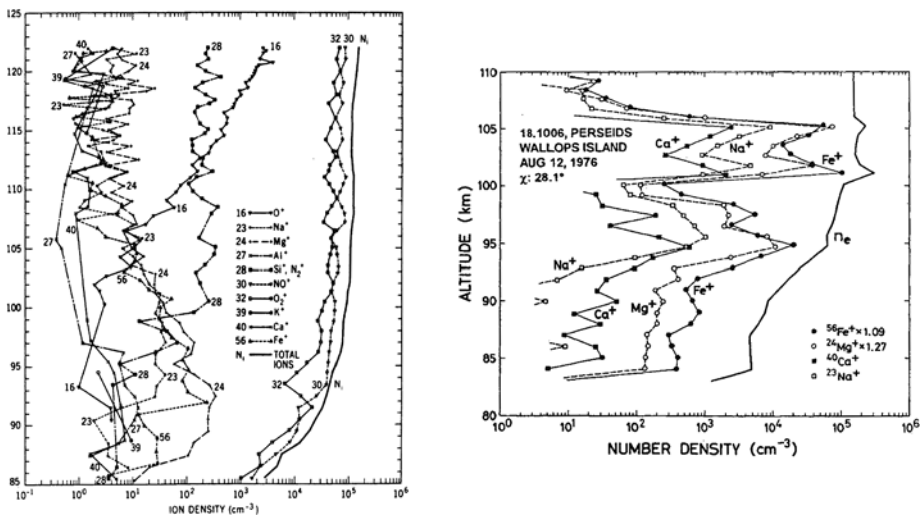


Fig. 1 Observed distribution of positive ion species in the terrestrial ionosphere during the following conditions: (left) daytime over Thumba (India) (Aikin and Goldbert 1973); (right) Perseid meteor shower on August 12, 1976, above Wallops Islands (Kopp 1997)

2 Evidence of Meteoroid Layers through the Solar System

Planetary ionospheres have been sounded by radio occultation techniques since the beginning of the 1960s; the measurements show a daytime ionosphere with a major peak mainly produced by solar radiation and photoelectrons. Below the main peak, one or more secondary ionospheric layers have been found. Examples of such layers have been found at Venus by Pioneer Venus (Kliore et al. 1979) and Mars by Mariner IV (Fjeldbo et al. 1966), Mars 4 and 5 (Savich et al. 1976) and Mars Express (Pätzold et al. 2005) among others, see Fig. 2. The number of missions to the external planets is not so comprehensive, Galileo (Hinson et al. 1997) and Voyager (Hinson et al. 1998) found evidence of such layers in the Jovian atmosphere, Cassini (Nagy et al. 2006) in the atmosphere of Saturn and Voyager at Uranus (Strobel et al. 1991) and Neptune (Lyons 1995), see Fig. 3.

2.1 Venus

The electron density profile measured by Pioneer Venus in the nightside ionosphere shows low altitude layers below the main ambient ionospheric layer, see Fig. 2. The altitude of the main peak is located at 142.2 ± 4.1 km, very close to the main peak of the dayside terminator ionosphere. The peak density is characterized by a great variability, with a magnitude ranging from 23×10^3 to 40×10^3 cm^{-3} (Kliore et al. 1979). A double-peak structure appears during two closely spaced orbits, 55 and 57, and on orbit 57 the structure appeared in both the entry and exit measurements. It seems that the appearance of such a double-peak structure is a relatively rare temporal phenomenon. The altitude of the layer, ~ 120 km, is in agreement with the maximum for meteoroid ablation. Other ionization sources, such as direct impact ionization by electron precipitation or protons into the nightside, could also explain the nature of the lower ionospheric layer.

2.2 Mars

The atmosphere of Mars has been sounded in more detail than the rest of the extraterrestrial solar system planets, and its ionosphere presents the strongest evidences of metallic layers. The magnesium ion $\text{Mg}^+ \cdot \text{CO}_2$ has even been identified as a constituent of the Martian ionosphere (Aikin and Maguire 2005). The daytime ionosphere is well characterized by a main layer produced by solar radiation at an altitude of 140 km with a number density of some teens of thousands electrons per cubic centimeter. Mariner IV found a secondary layer one order of magnitude lower at around 100 km, below the main photoionospheric peak (Fjeldbo et al. 1966). Some years later, the soviet Mars 4 and 5 (Savich et al. 1976) found a layer at around 80 km during nighttime similar, in magnitude, to the daytime one. Mars Express confirmed the existence of a sporadic layer between 65 and 110 km in altitude in 10 of 120 ionospheric electron concentration profiles (Pätzold et al. 2005). Figure 2 shows the measurements developed by Mars 4 and 5 (left) and Mars Express (right). The occurrence of the daytime layer was not limited to specific times of the day or locations, and part of it is hidden in the lower portion of the upper one (Pätzold et al. 2005). Mars Express did not find such layers in the 20 ionospheric observations at night, all of which were at high southern latitudes during winter. The observations are too limited to exclude the occurrence of a layer at night (Pätzold et al. 2005). Theoretical models considering meteoroids ablation indicate that the altitude and magnitude of the observed layers can be explained by long lived metallic ions deposited by meteoroids (Pesnell and Grebowsky 2000; Molina-Cuberos et al. 2003).

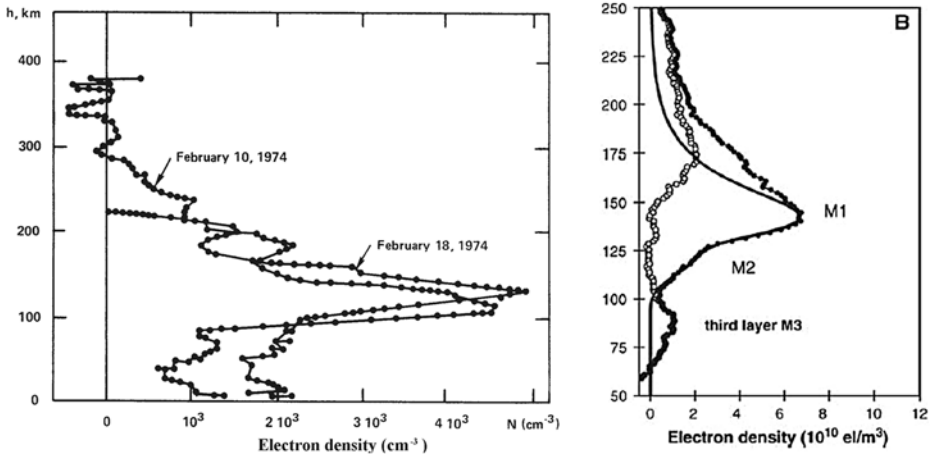
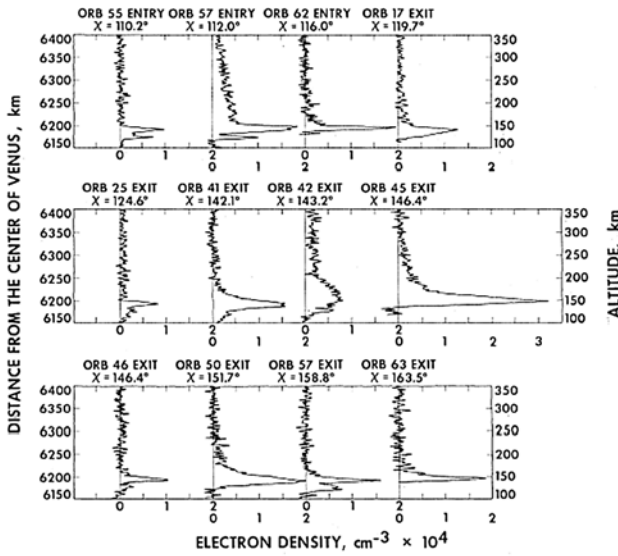


Fig. 2 Top electron concentration profiles in the nightside ionosphere of Venus measured by Pioneer Venus. From A.J. Kliore et al., *Science* 205:99–102 (July 6 1979). Reprinted with permission from AAAS. Lower-left distribution of electron concentration in the nighttime ionosphere of Mars measured by Mars 4 and 5 (Savich et al. 1976). Lower-right electron concentration in the Martian ionosphere observed by Mars Express (solid circles) and after subtracting a Chapman ionization model (open circles). From M. Pätzold et al., *Science* 310:837–839 (2005). Reprinted with permission from AAAS

2.3 Jupiter

The Voyager 2 fly-by provided most of the information about the lower ionospheric structure of Jupiter (Hinson et al. 1998). The electron concentration profile obtained during the egress contains two distinct layers: one is centred near 1000 km (relative to the 1 bar altitude) with a peak number density of $46 \times 10^4 \text{ cm}^{-3}$ and the structure of the other is more complex, see Fig. 3. It is formed by a group of fine layers situated between 300 and 500 km

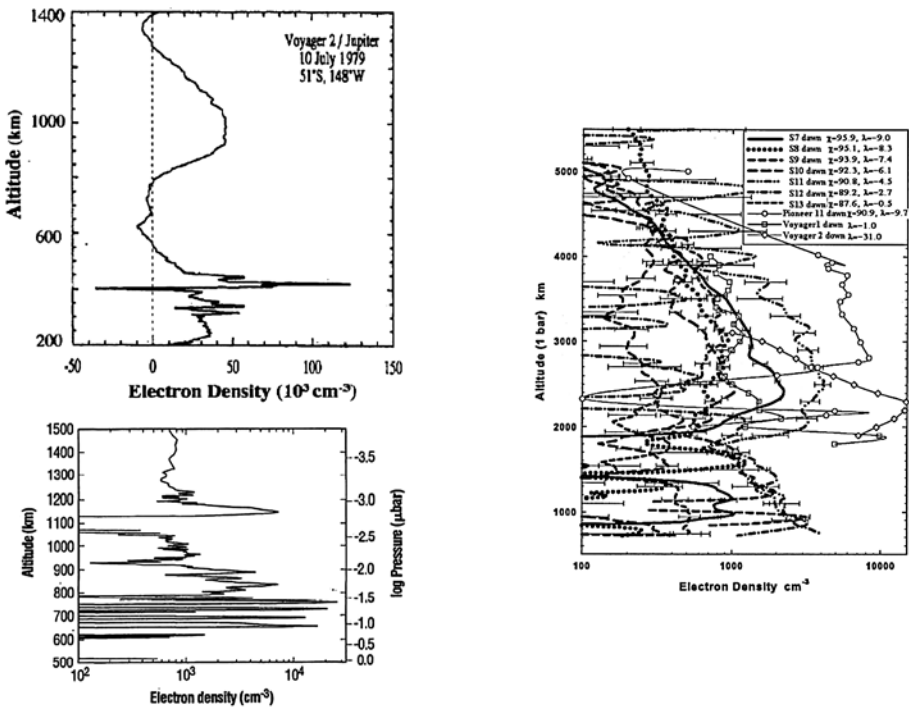


Fig. 3 Electron concentration profile at Jupiter (*top-left*) (Hinson et al. 1998), Neptune (*lower-left*) (From J.R. Lyons, *Science* 268:648 (1995). Reprinted with permission from AAAS) and Saturn (*right*) (Nagy et al. 2006)

with a concentration of $(20\text{--}120) \times 10^3 \text{ cm}^{-3}$ that might be formed in response to vertical shear in the zonal wind or plasma instabilities. Pre-Voyager theoretical models of the lower ionosphere predicted a layer of hydrocarbon ions in the 300–400 km altitude range (Kim and Fox 1994); although the calculated magnitude is around two order of magnitude smaller than the observed one. The difference between theory and observations seems too large to be reconciled by considering atmospheric processes. A plausible explanation is that the lowest layer is composed of long-lived metallic ions supplied by meteoroids or by the Galilean satellites (Hinson et al. 1998).

2.4 Saturn

Saturn presents quite a complex ionosphere where several ionization sources and physical processes take place. In addition to the solar and cosmic radiation, water inflow and particle impact have important influences on the ionospheric structure. The lower part also presents some layered structure, as it has been detected by Cassini (Nagy et al. 2006). Figure 3 shows the electron concentration profile for exit (dawn terminator) measured by Pioneer 11, Voyager 1 and 2, and Cassini. We can observe that only Cassini was able to determine the structure of the lower part of the ionosphere. The peak densities are, in general, larger for the dusk results than for the dawn ones. The profile corresponding to the dawn terminator presents a thick layer in the 900–1500 km range, well below the main peak placed at around 2500 km. For the dusk terminator, the layer is more sharp and could be partially included in

the upper main ionosphere. The magnitude of the lower ionospheric layer is, in both cases, of around 1000 cm^{-3} .

2.5 Neptune

Electron number density profile in Neptune observed during Voyager 2 occultation revealed sharp layers in the lower ionosphere with densities of around 10^4 cm^{-3} (Lyons 1995), see Fig. 3. The magnitude of the layers are even higher than the upper peak. It must be taken into account that the uncertainty in the electron abundance is high in the lower ionosphere, by as much as a factor of two, but the altitude of the layers is well determined from the phase of the received signal. A simple explanation for these layers is that the long-lived metallic ions are compressed by a horizontal wind with a vertical shear acting on the ions in the presence of a magnetic field (Lyons 1995).

3 The Interaction of Meteoroids with a Planetary Atmosphere

In parallel to the experimental observations, many theoretical studies have considered the effect of dust in the atmosphere of Venus (McAuliffe and Christou 2006), Mars (Adolfsson et al. 1996; Pesnell and Grebowsky 2000; Molina-Cuberos et al. 2003), Jupiter (Grebowsky 1981; Hinson et al. 1998; Kim et al. 2001), Saturn (Moses and Bass 2000), Titan (Ip 1990; English et al. 1996; Molina-Cuberos et al. 2001), Neptune (Moses 1992; Lyons 1995) and Triton (Pesnell et al. 2004).

The evaluation of meteoroid effects on planetary atmospheres requires the knowledge of the mass and velocity distributions of the meteoroid flux through the Solar System. Then the dynamical evolution of small particles through the atmosphere has to be calculated in order to determine the deposition profiles of neutrals and ions. The concentration of each metallic compound is calculated by solving the continuity and momentum equations.

3.1 Interplanetary Flux

Collisions between asteroidal parent bodies or grains released by comets are the major source of meteoroids in the Solar System (Liou et al. 1995; Gurnett et al. 1997). Exogenous sources also exist, particles coming from the local interstellar medium cross the Solar System on hyperbolic orbits (Grün et al. 1993). The experimental information of the dust distribution beyond the orbit of the Earth comes from the dust detectors on board of Pioneer 10 and 11, Ulysses, Galileo and Cassini (Humes 1980; Grün et al. 1993; Altobelli et al. 2007):

The measurements by the penetration detector of Pioneer 10 indicate that the spatial density of 10^{-9} g meteoroids is essentially constant between 1 and 18 AU (Humes 1980). The data obtained by the detector on Pioneer 11 show that meteoroids between 4 and 5 AU are not in circular orbits near the ecliptic plane, but they follow randomly inclined orbits of high eccentricity (Humes 1980), which implies a cometary origin. During the Saturn encounters, the on board detectors measured an increase in the flux of about three orders of magnitude, probably as a result of impacts from ring particles (Humes 1980). The data obtained with the penetration detector on board Pioneer 10 and 11 have similar shapes in spite of the differences on the threshold mass, 10^{-9} and 10^{-8} g, respectively, which means that the particle size distribution does not change strongly with the orbital radius (Cuzzi and Estrada 1998).

The high sensitivity Cassini Dust Analyzer (Altobelli et al. 2007) measured the dust particles between Jupiter and Saturn and found two main groups. The first group of impacts consists of particles on bound and prograde orbits coming from the dust ram direction, most probably on low eccentric and low inclined orbits, and they show a large spread in mass. The possible sources are short-period Jupiter family comets or circumsolar dust. Impactors of the second group were identified as interstellar dust particles, perhaps including a minority of beta-meteoroids. The upper limit value of the flux and the particles size are in very good agreement with what is expected from the Ulysses data (Grün et al. 1993) and model predictions.

The dust environment in the outer Solar System has been sounded by the Voyagers. The plasma wave instruments were able to detect a small but persistent level of dust impacts (up to 51 AU for Voyager 1 and up to 33 AU for voyager 2). The average number density obtained is estimated to be around $2 \times 10^{-8} \text{ m}^{-3}$, and the average mass of around 10^{-11} g (Gurnett et al. 1997). The ecliptic latitudes of the paths taken by Voyager 1 and 2 were quite different. After the flyby of Saturn at 9.5 AU, Voyager proceeded northward from the ecliptic plane at an asymptotic ecliptic latitude of about 35° . Voyager 2 remained very close to the ecliptic plane until the flyby of Neptune at 30 AU. Considering the differences in the spacecraft trajectories, the observed variations in number densities were small and (Gurnett et al. 1997) concluded that comets are the most likely source for interplanetary dust particles in the outer Solar System.

3.2 Entry Velocity

For particles on bound orbits, the velocity distribution of dust particles arriving a planet depends on the distance to the Sun and on the planetary gravitational field. The velocity decreases with the distance to the Sun. Cuzzi and Estrada (1998) found the relationship between the orbital velocity of the meteoroids and the distance to the Sun to be:

$$v(R) = \frac{v_1}{\sqrt{R_{AU}}}, \tag{1}$$

where v_1 is the velocity at 1 AU and R_{AU} the distance to the Sun.

Meteoroids penetrate the atmosphere at higher velocities than predicted by (1). The planetary gravitational field accelerates the particle and its orbit becomes parabolic. If v^* is the relative speed of a particle with respect to the planet, then the meteoroid speed $v(r)$ at a distance r from the planetary centre is:

$$v(r) = \sqrt{v_{esc}^2(r) + v^{*2}}, \tag{2}$$

where $v_{esc}(r)$ is the planetary escape velocity of an object at a distance r .

The gravitational field also produces an enhancement in the cross section of a planet and, therefore, an increase in the meteoroid flux by a factor (Bauer 1973):

$$G = 1 + \left(\frac{v_{esc}}{v^*}\right)^2. \tag{3}$$

Table 1 shows the characteristic velocities of meteoroids arriving at Solar System bodies with a noticeable atmosphere. V_{min} and V_{max} represent the minimum and maximum velocities reaching the top of the atmosphere, respectively, where meteoroids in heliocentric orbits are assumed. The minimum velocity corresponds to the planetary escape velocity at the top

Table 1 Characteristic velocities (km s⁻¹)

	V_{\min}	V_{\odot}	V_{orb}	V_{\max}
Venus	10	50	35	85
Earth	11	42	30	72
Mars	5	34	24	58
Jupiter	60	19	13	69
Saturn	35.5	13.7	9.7	42.5
Titan ^d	2.6	13.7	5.58	29.1
Uranus	21	9.6	6.8	26.6
Neptune	23.5	7.7	5.5	26.9

^dThe case of Titan includes the gravitational focus of Saturn, V_{orb} is the orbital velocity around Saturn

of the atmosphere. The maximum velocity corresponds to a particle with the solar system escape velocity V_{\odot} orbiting the Sun in a retrograde orbit. V_{orb} is the planetary orbital velocity.

It can be observed that the range of velocities ($V_{\max} - V_{\min}$) is very wide for terrestrial planets due to the combined effect of a low escape velocity and high orbital velocity. For the giant planets this velocity range is much smaller.

3.3 Meteoroid Ablation

The problem of determining the physical evolution of a particle penetrating the atmosphere was first treated in detail by Öpik (1958). Here we briefly describe the processes describing the loss of velocity and mass of a spherical small particle, based in the work of Lebedinets et al. (1973). We do not consider aspects like fragmentation, non spherical shape, mixed compositions or differential ablation which are usually important for the detailed modeling of the Terrestrial atmosphere or for high mass particles.

Meteoroids penetrating the atmosphere are accelerated by the planetary gravitational field and slowed down by collisions with atmospheric constituents. Collisions also remove part of the mass and heat the particle surface producing an additional loss of mass by evaporation. The increase of the particles' temperature by atmospheric collisions is balanced with thermal radiation and loss of heat through ablation.

The dynamical evolution of a small particle penetrating the atmosphere is calculated by solving the motion, ablation and energy equations. The motion equation relates the decrease in relative impact velocity v due to the drag of the atmosphere:

$$\cos \theta \frac{dv}{dz} = \frac{\Gamma A \rho v}{\delta^{2/3} m^{1/3}} \quad (4)$$

where θ is the entry angle, and v , m and δ are the meteoroid velocity, mass and density, respectively. The atmospheric drag depends on the drag coefficient, Γ , the atmospheric mass density, ρ , and on the meteoroid shape, through factor A .

The ablation equation relates the loss of mass m from a meteoroid due to evaporation and sputtering:

$$\cos \theta \frac{dm}{dz} = -\frac{4AK_1 m^{2/3}}{\delta^{2/3} v T^{1/2}} e^{-K_2/T} - \frac{\Lambda_S A \rho m^{2/3} v^2}{2Q^{2/3}} \quad (5)$$

where K_1 and K_2 are constants describing the dependence of the evaporation rate on temperature, T . Λ_S is the sputtering coefficient and Q the energy of evaporation of 1 g of the meteoroid.

Finally, the energy equation provides the thermal evolution of the particle as a function of the increase of temperature due to the heating by sputtering, thermal radiation from the body surface and the deposition of energy by evaporation:

$$\cos\theta \frac{dT}{dz} = \frac{4A\rho v^2}{8C\delta^{2/3}m^{1/3}}(\Lambda - \Lambda_s) - \frac{4A\sigma T^4}{C\delta^{2/3}vm^{1/3}} - \frac{4AK_1Q}{C\delta^{2/3}T^{1/2}m^{1/3}v}e^{-K_2/T} \tag{6}$$

where Λ the heat transfer coefficient, σ the Stefan–Boltzmann constant and C the heat capacity of the meteoroid substance.

The energy, mass and momentum are couple by a system of equations which usually demand a high precision numerical method and a small discretization grid in order to solve it. In principle, the effect of a velocity distribution has to be taken into account in the modeling, mainly if a high precision of the meteoroid mass deposition profile is required, as usually occurs to the Earth case. However, most of the numerical models of the extraterrestrial atmospheres simply consider a monochromatic distribution at the mean velocity, rather than a distribution of velocities, and analyse the effect of different incoming velocities. For low gravity planets a mean angle of 45 degrees for incoming particles can be used. The effect of non-vertical entry is to elevate the altitude at which ablation occurs. For massive planets the gravitational focus shifts the distribution of incident angle towards vertical.

Icy meteoroids coming from comets ablate more easily and at higher altitude than stony meteoroids, which are produced in the asteroid belt. Rocky material introduces a higher amount of metal constituents to the atmosphere than icy meteoroids.

Once the ablation of meteoroids is known, the linear concentration of the individual ion species can be calculated by

$$\alpha_i = -\frac{p_i\beta_i}{m_i v} \frac{dm}{dt} \tag{7}$$

where p_i is the ratio of atom i type to the total, m_i is the atomic mass, and β_i is the ionization probability, which depends on the ion produced and meteoroid velocity. For low velocity $\leq 35 \text{ km s}^{-1}$, for which no secondary ionization or recombination take place, Jones (1997) proposed an empirical expression:

$$\beta_i = k_i(v - v_i)^2 v^{0.8} \tag{8}$$

where k_i is an experimental value, which depends on the element and v_i is a cut off velocity, For high velocity particles, a more general expression can be used (Lebedinets et al. 1973):

$$\beta_i = Cv^{7/2}. \tag{9}$$

The production rate of i type ions (P_i) is calculated as:

$$P_i = \int \alpha_i f(m)dm, \tag{10}$$

where $f(m)$ is the flux per unit of micrometeoroidal mass.

4 Modelling Metallic Layers in Planetary Atmospheres

The meteoroid mass deposition profiles are calculated by adapting the flux of interplanetary dust and solving the dynamical evolution of the particles, as described in the above section.

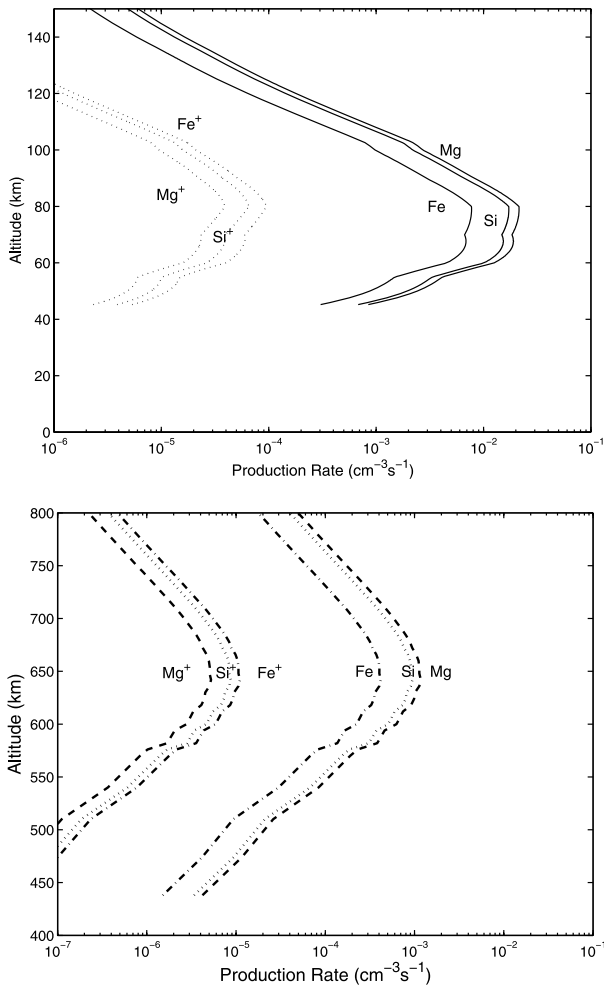


Fig. 4 Neutral and ionic deposition rates of Mg, Fe and Si due to the ablation of meteoroids with entry velocity of 18 km s^{-1} in the atmosphere of Mars (*top*) (Molina-Cuberos et al. 2003) and Titan (*bottom*) (adapted from Molina-Cuberos et al. 2001) to consider the same meteoroid composition than on Mars

We have used the model by Grün et al. (1985), which assumes an isotropic flux of meteoroids at Earth's orbit with effective density of 2.5 g cm^{-3} and mean velocity of v (1 AU) = 20 km s^{-1} . The model considers mass ranging from 10^{-18} to 100 g , although the main contribution to the total mass is due to particles ranging from 10^{-7} to 10^{-4} g .

Figure 4 shows the neutral and ion deposition rate of magnesium, iron and silicon in the atmosphere of Mars (*top*) and Titan (*bottom*), assuming the interplanetary dust is mainly composed of carbonaceous chondrites, that have a relative concentration of Mg = 6.1%, Si = 5.7% and Fe = 5.1% (Anders and Ebihara 1982).

Table 2 shows the altitude range where the maximum of the meteoroid ablation occurs and the altitude of the ionospheric layers that could consist of metallic ions. Please note that the ionospheric peak is located in all the cases quite close to the ablation altitude.

Table 2 Theoretical predictions of the altitude of the meteoroid deposition peak and altitude of ionospheric layer that could consist on metallic ions. ☉ represents ice meteoroids and ● silicate ones

Planet	Ablation (km)	Reference	Ionospheric peak (km)
Venus	110–120	<i>This work</i>	120–130 km
Earth	85–95	McNeil et al. 1998	90–100 km
Mars	75–85	Molina-Cuberos et al. 2003	65–100
Jupiter	300–400	Kim et al. 2001	300–550
Saturn	790–1290 ☉ 610–790 ●	Moses and Bass 2000 Moses and Bass 2000	900–1200
Titan	650–700	Molina-Cuberos et al. 2001	<i>No evidences</i>
Neptune	500–800 ● 250–500 ☉	Moses 1992 Moses 1992	600–1000 km

The altitude of the ablation depends on the physical characteristics of meteoroids (volatility, speed, size and composition). The radio-occultation measurements of ionospheric layers of long-lived metallic ions through the Solar System planets could, therefore, provide information about properties and composition of interplanetary dust. The distribution of metals is, however, affected by transport and layering.

Once the meteoroid deposition in the atmosphere is known, the concentration of each neutral and ion species is calculated from the continuity and momentum equations that, assuming steady state, may be expressed as:

$$P_i - n_i l_i = \frac{\partial}{\partial z} n_i v_i \tag{11}$$

$$v_i = -D_i \left(\frac{1}{n_i} \frac{\partial n_i}{\partial z} + \frac{1}{H_i} + \frac{1}{T} \frac{\partial T}{\partial z} \right) - K \left(\frac{1}{n_i} \frac{\partial n_i}{\partial z} + \frac{1}{H} + \frac{1}{T} \frac{\partial T}{\partial z} \right) \tag{12}$$

where i denotes the i th constituent, z the altitude, n_i the concentration, P_i the production, l_i the specific loss, T the temperature. v_i is the mean vertical velocity, D_i and K_i are molecular and eddy diffusion coefficients, H_i and H are the individual and atmospheric scale heights.

Vertical transport of metallic species is mainly produced by turbulent and molecular diffusion, the former being more effective at lower levels. Diffusion theory provide analytical expressions for the molecular diffusion coefficients D_i (Chapman and Cowling 1970). The turbulent diffusion is the least known factor in the modeling of an atmosphere. It is typically parameterized by means of an eddy diffusion coefficient K .

The production of metallic ions depends on the atmospheric and ionospheric characteristics. Other ionization sources, like solar radiation or electrons, provide atmospheric ions that transfer the charge to metallic atoms by charge exchange reactions, which is a very important source for metallic ions. In addition, the photoionization of metallic neutrals increases the production of metallic ions.

4.1 Earth

Meteoritic metals in the terrestrial atmosphere have been modeled in detail. The chemistry and temporal and spatial variations have been extensively treated in the literature (Swider 1969; Aikin and Goldbert 1973; Carter and Forbes 1999; Joiner and Aikin 1996). In particular, special attention has been paid to the most abundant metallic species, magnesium (McNeil et al. 1996; Plane and Helmer 1995), iron (Helmer et al. 1998; Carter and Forbes 1999), silicon (Kopp et al. 1995), potassium (Eska et al. 1999), sodium and calcium (McNeil et al. 1998; Plane et al. 1999) and the chemistry is well known.

Metals constitute a very small fraction of the total atmospheric constituents in the E-region; however the ionized fraction of metallic atoms is very high when compared with other atmospheric compounds. The reason for the relatively high concentration of metallic ions is their low electron recombination rate, several orders of magnitude slower than the recombination of the most abundant ambient ions, O_2^+ and NO^+ . The chemical lifetime of metallic ions is very long and their vertical distribution is strongly influenced by transport mechanisms such as eddy diffusion and layering due to wind shears and electric field. Metallic species are removed from the E-layer by downward transport. Three-body association reactions with atmospheric neutrals produce molecules in the gas phase that subsequently condensate and coagulate to aggregates and aerosols (Hungen et al. 1980). Figure 6 shows a schematic diagram of reactions involving magnesium species, as an example of the chemistry of metallic compounds in the terrestrial atmosphere. It was developed from the works by Plane and Helmer (1995), McNeil et al. (1996) and McNeil et al. (1998).

Magnesium ions are mainly produced by charge exchange with atmospheric ambient species (O_2^+ , NO^+ and O^+) (Grebowsky et al. 1998), and also by direct meteoric ionization and photoionization. The recombination of Mg^+ with electrons is not the main loss process of Mg ions. Three-body reactions of Mg^+ with O_2 and N_2 produce MgO_2^+ (Plane and Helmer 1995) and MgN_2^+ (McNeil et al. 1996) and two-body reaction with ozone leads to the formation of MgO^+ . These ions are recycled to neutral Mg through molecular dissociative recombination. The chemistry of neutral magnesium is determined by two/three-body reactions with oxygen species (O , O_2 and O_3), which produce MgO and MgO_2 . The final sink of Mg may be $Mg(OH)_2$ as obtained by Plane and Helmer (1995) or $MgCO_3$ as calculated McNeil et al. (1996).

Grebowsky et al. (1998) compiled all published studies describing rocket flights which measured meteoric ions between 1963 and 1991. They found that the observed Mg^+ concentrations are lower than those yielded by models and they also confirmed that meteor showers do have significant impact on the average ionospheric composition.

4.2 Mars

Terrestrial knowledge is the starting point to model the meteoroid effects on the atmosphere of Mars. The chemistry of metallic ions is quite similar to the terrestrial case, CO_2 playing the role of the third body in three-body reactions in Mars instead of N_2 on Earth.

Magnesium and iron ions are produced by direct meteoric ionization, photoionization and charge exchange with atmospheric ions, mainly O_2^+ . The last one is the main source for production of metallic ions. Magnesium and iron follow quite similar processes. The oxidation by ozone is the most efficient mode of converting atomic metals into neutral oxides. Once MgO and FeO are formed, three-body association of CO_2 provides carbonated metallic atoms, which are the more stable neutrals. Metallic ions can undergo electron recombination or be converted to oxygenated ions. At higher pressure, molecular association by three-body

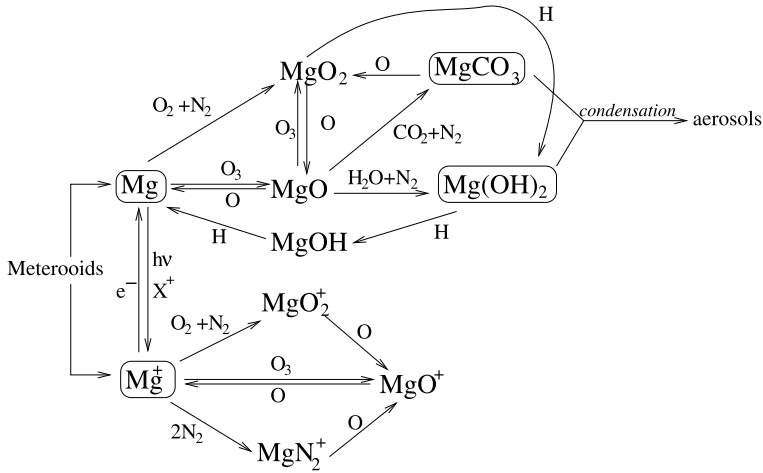
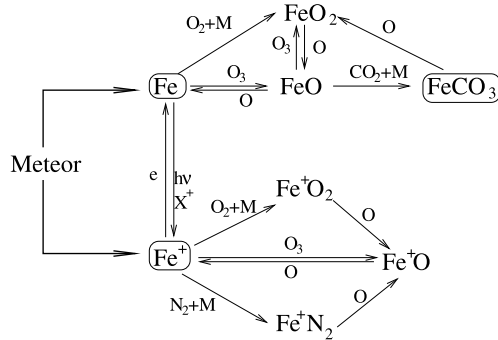


Fig. 5 Schematic diagram of the chemistry of magnesium species in the terrestrial atmosphere, where X^+ represents a non-metallic ion (mainly O^+ and NO^+) and $h\nu$ photoionization. Adapted from Plane and Helmer (1995), McNeil et al. (1996, 1998)

Fig. 6 Schematic diagram of the chemistry of iron species in the martian atmosphere, where X^+ represents a non-metallic ion and $h\nu$ photoionization. Adapted from Molina-Cuberos et al. (2003)



reactions produce molecular ions, which undergo molecular dissociative recombination to form Fe and Mg (Molina-Cuberos et al. 2003). Figure 6 shows a schematic diagram of the chemistry of iron species. For the case of Mg, it follows a general scheme similar to Fe, but with different rates.

Pesnell and Grebowsky (2000) modeled the effect of magnesium in the atmosphere of Mars and predicted a persistent layer of Mg^+ in the order of 10^4 cm^{-3} at around 80 km, which is a factor of around 20 times lower than the main ionospheric peak placed at 130 km. Molina-Cuberos et al. (2003) developed daytime and nighttime models for iron and magnesium produced by meteoric ablation, and the effect of solar activity, and seasonal variations was also explored. They found a meteoric layer formed by Fe^+ and Mg^+ with a magnitude of the order of 10^4 cm^{-3} at noon and decreases by two orders of magnitude during the night. The agreement between the model and the daytime measurements taken by Mars Express (Pätzold et al. 2005) some years later is quite good. Figure 7 shows (solid lines) the electron concentration predicted and a pair of experimental profiles that match better with the theoretical model.

Fig. 7 Theoretical predictions (Molina-Cuberos et al. 2003) (lines) and experimental determination (Pätzold et al. 2005) of meteoroid layers in the Martian ionosphere

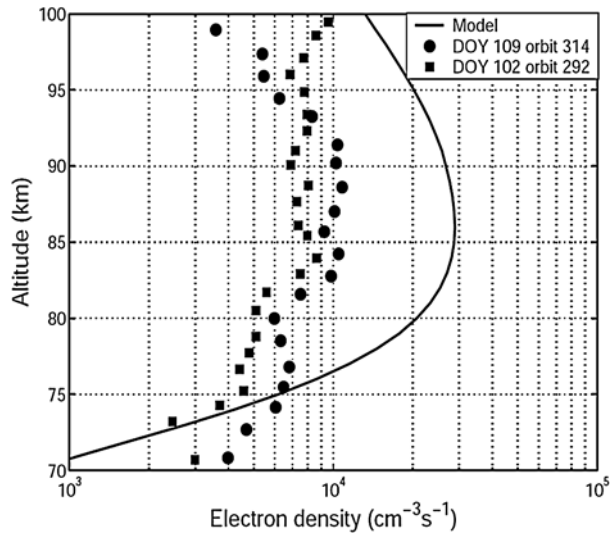
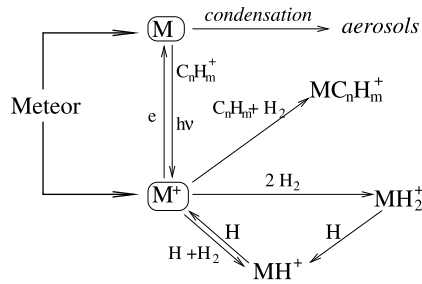


Fig. 8 Schematic diagram of the chemistry of metals in the ionosphere of giant planets, where M represents Mg^+ or Fe^+ . Electron recombination of metals have not been included for higher clarity, adapted from Kim et al. (2001)



4.3 Giant Planets

The hydrogen-hydrocarbons atmosphere of the giant planets is more difficult to model than the Martian one, due to the lack of measurements of reaction rates between hydrocarbons and metallic ions.

The ionospheres of gaseous planets are usually dominated by H^+ in the upper part and H_3^+ prevails below (Kim et al. 2001; Moses and Bass 2000). If metals are not considered in the modeling, then hydrocarbons are the major ionic species in the lower ionosphere (Kim and Fox 1994). However the predicted densities are much lower than the observed ones. The inclusion of meteoric ablation into the atmosphere that produces metal ions that take the place of hydrocarbons as the major ionic species in the lower part of the ionosphere (Kim et al. 2001; Moses and Bass 2000; Lyons 1995). Figure 8 shows a short scheme of the chemistry of metals in the atmosphere of giant planets.

The high gravitational focus of giant planets and the long distance to the Sun results in metallic ions that are mainly produced by charge exchange with atmospheric ions and direct ionization from ablation; the other source, photoionization of metallic neutrals, is much less important. Metallic ions are removed by three-body reactions with hydrocarbons and hydrogen, and the metallic atoms are lost by condensing onto existing aerosols or dust particles. The depletion of atoms is parameterized by assuming a constant lifetime, which depends on the number of dust particles and their size; typical values in the range of (1–6)

$\times 10^5$ s have been used in the modeling of Jupiter (Kim et al. 2001), Saturn (Moses and Bass 2000), and Neptune (Lyons 1995). Three-body reaction between metallic ions (Fe^+ , Mg^+) with the most abundant atmospheric neutral (H_2) has not been confirmed in the laboratory. Even by assuming that the reaction occurs, it results in little net loss of ions, due to the adduct ($\text{Fe}^+\cdot\text{H}_2$ or $\text{Mg}^+\cdot\text{H}_2$) undergoing a sequence of reactions with H atoms which restore metallic ions in the Jovian atmosphere (Kim et al. 2001).

The electron concentration profile of outer planets shows frequent gaps, specially at the dawn terminators, see Fig 3. The layering that is frequently observed in the lower ionospheres of the outer planets could be caused by long-lived atomic ions being moved by horizontal winds that possess vertical shear, such as might occur with atmospheric tides and gravity waves interacting with meteoric ions.

The ablation of meteoroid in the Jovian atmosphere and the chemistry of meteoric ions (O^+ , C^+ , Si^+ , Fe^+ , Mg^+ , Na^+ and S^+) was modeled by Kim et al. (2001). They found a layer of meteoric ions in the altitude region of 350–450 km, above the 1-bar level, with a peak total ion concentration of several times 10^4 cm^{-3} , which are comparable with the observed one. Moses and Bass (2000) studied the effects of external material on the chemistry and structure of Saturn's ionosphere, they found that the ionospheric structure is dominated by two major peak, with H^+ creating the high-altitude peak and Mg^+ , representing the metallic ions, the low-altitude peak with a magnitude in the order of 10^4 cm^{-3} , similar to the upper one. Lyons (1995) also considered metal ions to model the lower ionosphere of Neptune and calculated a concentration of Mg^+ around 1000 cm^{-3} , one order of magnitude lower than the one observed by Voyager 2. The differences are reduced if the magnesium ions are compressed into sharp layers by a sinusoidal vertical wind.

4.4 Titan

The Voyager 1 fly-by of Titan was able to determine an ionospheric peak of $2400 \pm 1100 \text{ cm}^{-3}$ at $1180 \pm 150 \text{ km}$ (Bird et al. 1997), and did not provide any information from below the peak. Therefore, Voyager did not show any evidence for meteoroids effects in its atmosphere. From 2004, Cassini has been orbiting Saturn and several opportunities to sound the ionosphere by radio-occultation and even to determine ionic mass by INMS (Ion and Neutral Mass Spectrometer) will be available in the near future. In spite of the lack of experimental observations of the effects of meteoroids at Titan, some theoretical models have investigated the effects of meteoroids in the composition of neutral (English et al. 1996) and ion (Molina-Cuberos et al. 2001) species. Molina-Cuberos et al. (2001) investigated the ablation of meteoroids and found that long-lived metallic ions considerably change the predictions of the electron number density due to models which only consider solar radiation and electrons trapped in the magnetosphere of Saturn. By using a simple model where metallic ions are lost by termolecular associations with neutral molecules, they concluded that an ionospheric layer could be present at around 700 km, with an electron concentration peak similar in magnitude to the one produced by solar radiation. Petrie (2004) has developed theoretical calculations of Mg^+ reactions with the atmospheric compounds of Titan (N_2 , CH_4 and some nitrogenated hydrocarbons), which allowed the magnesium chemistry to be modelled in more detail and predicted the radical MgNC as the final product.

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