

Discussion of Derivability of Local Residual Stress Level from Magnetic Stray Field Measurement

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Received: 8 April 2015 / Accepted: 10 June 2015 / Published online: 1 July 2015
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Abstract The NDT procedure dubbed ‘metal magnetic memory’ method and the related ISO 24497 standard has found wide industrial acceptance in some countries, mainly in Russia and China. The method has been claimed by some researchers (Roskosz and Bieniek in *NDT&E Int* 45:55–62, 2012; Wilson et al. in *Sens Actuators A* 135:381–387, 2007) as having potential for quantitative determination of local residual stress state in engineering structures, at least for some steel grades. This work presents a critical reexamination of a previous important study by Roskosz and Bieniek, who claimed to have found a direct relationship between local residual equivalent stress levels ranging from 0 to 50 MPa, and the stray field gradients in T/P24 steel sample placed in the Earth’s ambient magnetic field. We reconstruct their experiment in a magnetic finite element simulation, computing stray magnetic field and its tangential gradients along the axis of the sample. Different combinations of remanent induction and relative magnetic permeability levels have been modeled, and the influence of geometrical discontinuity is quantified. In order to validate magnetic finite element methodology, a new experiment is presented, along with its numerical counterpart. The magnetic finite element method allowed to obtain a good quantitative correlation with well-controlled stray field measurements. It is demonstrated, that the residual stress level of order of 50 MPa is not the only factor, on which the stray field measurement depends. The geometrical discontinuity and the remanent induction contribute to a higher extent to the field amplitudes. Consequently we prove, that a bidirectional correlation between

the magnetic field gradient and local stress levels cannot be determined because of at least three concurrent inseparable factors on which the measured stray field and its spatial gradient depends.

Keywords Metal magnetic memory · Stray magnetic field · Finite element method · Residual stress evaluation

1 Introduction

The NDT procedure based on notions of the metal magnetic memory (MMM) or residual magnetic field (RMF) method has been promoted by its inventor, A. Dubov, since 1994. In 2007 the procedure obtained an international standardization (ISO24997, [1]), and has been taken up commercially, mainly in Russia and China. It is worth noting, that the mentioned ISO standards were compiled by a direct translation of a Russian GOST standard.

The MMM procedure draws attention for at least two reasons:

- it has been applied in various safety-critical components [2–4]
- it has been claimed as capable of solving the essential problem of structural NDT, namely determining a local stress level from a fast and easily interpretable series of measurements in-situ [5–7].

In the literature, strong promotional aspects can be found, especially in the works by Dubov et al [3, 6, 7]. Dubov insists on a principal difference between MMM and classic stray field measurements (e.g. MFL) [6], to the advantage of the former. Actually, authors presenting a MMM-based qualitative defectoscopy [8, 9], exploit the well-known principles of

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flux leakage detection. MMM appears as a mixture of MFL (in its qualitative aspect) and magneto-elastic (Villari) effect, when stress concentrations are looked for. While a potential advantage of MMM consists in the ‘passive mode’ (i.e. no need to use an external magnetizing device), there is also a serious disadvantage of low signal-to-noise ratio and consequently much lower sensitivity of MMM to small defects. The low signal-to-noise ratio in MMM stems from the relatively low level of the geomagnetic field. The accuracy of the method remains unknown. In particular, Dubov failed to present reliable data on false positive and false negative rates of detection in comparison with other reference techniques.

A serious doubt arises when quantitative applications of MMM are claimed. In the fundamental paper [10] written by Wlasow and Dubov there is an interesting formula, which defines the local stress maximum as a simple function of K_{max} , K_{ave} and σ_m :

$$\sigma_{max} \sim \sqrt{\frac{K_{max}}{K_{ave}}} \sigma_m \quad (1)$$

where: K_{max} —the maximum value of the gradient of the RMF components in the area of stress concentration, K_{ave} —the average value of the gradient of the RMF components in the area under examination, σ_m —the ultimate strength of the material. These parameters, except for σ_m , are also defined in the ISO24997 standard [1]. The location of maximum field gradient is presented as a position of the highest “stress intensity factor”, not only due to an applied or residual stress. A defect or another geometrical change in test specimen could as well produce a stress intensity factor under loading.

If the fundamental formula (1) was correct and widely applicable, as claimed by Dubov et al, it would entail a fundamental break-through in NDT. Consequently, it requires a detailed evidence-based examination.

The quantitative use of any magnetic stray-field NDT (either active or passive) relies in the magneto-mechanical phenomena, i.e. a change of local magnetic properties of a solid as a function of stress. The magnetic properties in play are: magnetization $M(\sigma)$, magnetic permeability $\mu(\sigma)$, coercive field $H_c(\sigma)$, remanent induction $B_R(\sigma)$. All these variables form a general $B - H(\sigma)$ characteristic, defining the dependence of major and minor hysteresis loops on the local stress.

Several authors [11–15] have studied the direct magneto-mechanical problem on idealized flat samples, reproducing Villari’s experiment for different steel grades. Some empirical relationships can also be found in the pioneering work by Bozorth [16].

The foundation for analytical description of the $B - H(\sigma)$ was laid by Jiles [17], further extended and sometimes referred to as Jiles–Atherton–Sablík model. Although most of MMM researchers refer to Jiles’ model, only few of them

[18] emphasize the complex and multi-factorial relationships which require the use of partial differential equations.

Roskosz with his co-authors systematically studies the qualitative and putative quantitative potential of MMM. He notes the complexity of $B - H(\sigma)$ relationship and presents rather complicated patterns of measured magnetic stray field over a sample with a round hole [19]. However, in [2], he makes an unequivocal statement: ‘the magnetic metal memory is a nondestructive method with a great potential and is perfect for many applications.’ Another pair of works by Roskosz et al contain two contradicting conclusions:

- (a) [20] An attempt was made to strictly follow the MMM procedure as described in the ISO24497 standard, and assess local stress level in a dog-bone sample. Inconsistent results were found, and Roskosz states: ‘The only way to account for this is to state that either a mistake was made in the research or the methodology of the method is not fully understood, or the methodology itself is wrong. (...) Quantifying the level of stress concentration or the value of stress at the present stage of the development of the method of MMM testing is a disputable and doubt-raising issue.’
- (b) [5] Dog-bone samples of S235 and T/P24 steel grades are sequentially loaded beyond the yield limit. The residual stresses in the notch region are calculated using structural FEM. An inverse function: $dH(\sigma)/dx$ is plotted and presented as useful for quantitative determination of local stresses.

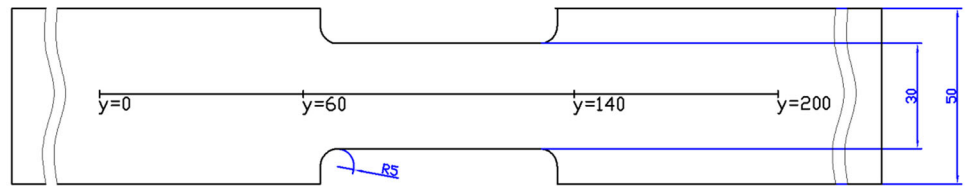
Facing these contradictory statements, we have formulated the following questions: Could the field gradients measured by Roskosz and Bieniek over S235 and T/P24 steel samples be due to the residual stress alone? Could a reliable inverse function be defined, allowing for deduction of a local residual stress level from the local stray field gradient?

The answer will be based on available literature, new original experimental data and magnetic finite element method (FEM) modeling. It is interesting to note, that the magnetic FEM was rarely if ever used in the context of quantitative MMM-related research. The authors know only of simulations concerning qualitative magnetic defectoscopy [9,21–23]. The only work aiming at defining a quantitative stress-field with magnetic FEM and experiment [24] did not attain the objectives, as declared by its author.

2 Experimental Set-Ups and Samples

Two experiments are discussed in this paper. The first, further referred to as ‘E1’, was performed by Roskosz and Bieniek and described in their paper [5]. Roskosz and Bieniek interpreted their results as a proof of a direct bidirectional

Fig. 1 Geometry of the sample used in experiment ‘E1’. Uniform thickness $t = 2$ mm



correlation between the local magnetic field gradient and the local stress level in the sample. The second experiment, referred to as ‘E2’ was performed by the authors of this article, in order to reexamine conclusions put forward by the aforementioned researchers.

2.1 Reinterpreted Experiment ‘E1’ by Roskosz and Bieniek

Roskosz and Bieniek measured two components of static magnetic field along the surface of a flat dog bone-shaped sample. The samples were made of S235 or T/P24 steel grades which were subjected to tensile stress exceeding yield limit and then unloaded. Figure 1 presents the samples’ dimensions and the adopted reference coordinate system.

A monotonic and thus potentially useful relationship between maximum magnetic field gradient and local stresses was postulated by those authors. The maximum of dH/dx was shown to be spatially coincident with the maximum of residual stress, produced by tensile loading beyond elastic limit and unloading of the sample. The residual stress was calculated by structural finite element analysis. The most important result was the claimed inverse relationship $\sigma_{eqv} \sim dH/dx$ plotted for T/P24 steel grade within the stress range (0–50 MPa) (Fig. 2).

The experiment by Roskosz and Bieniek had two important features, which influenced both the qualitative and quantitative character of the measurements:

- (a) Standard ferromagnetic clamps of a tensile test machine were used. In literature [8,25] it was shown, that the presence of magnetized clamps may produce an increase of stray field by an order of magnitude, and produce an uncontrollable remanent induction within the sample.
- (b) Unlike other authors [8,25], Roskosz and Bieniek extend their measurements over the notched portion of the sample, focusing on the peak found over the geometrical discontinuity. The potential geometrical effect is neglected in their interpretation. However, it can be deduced from the Gauss theorem, and has recently been confirmed experimentally [26,27], that the change of sample width is by itself an important source of stray field, and an enhancement of the local stress intensity factor resulting also in higher residual stress which Roskosz has computed by Finite Element Analysis.

2.2 New experiment ‘E2’

In order to resolve concerns about the interpretation of ‘E1’ results, authors of this paper conducted their own measurements, ‘E2’. The samples made of the same material and very similar in shape to those used in experiment ‘E1’, have been placed in the Earth’s magnetic field or alternatively in the controlled field produced by Helmholtz coils. Dimensions of sample ‘E2’ are shown in Fig. 3. Set-up used in the experiment, which in general enables measurements of stray magnetic field for a variable external magnetic field and loading force, is presented in Fig. 4. In contrast to the experiment ‘E1’, samples in ‘E2’ have not been loaded.

The new experiment ‘E2’ had two objectives. One consisted in assessing qualitatively and quantitatively the stray magnetic field around a sample with a geometric discontinuity, while providing well-controlled initial conditions. The second purpose was to validate the magnetic finite element methodology, comparing simulated results against experimental one.

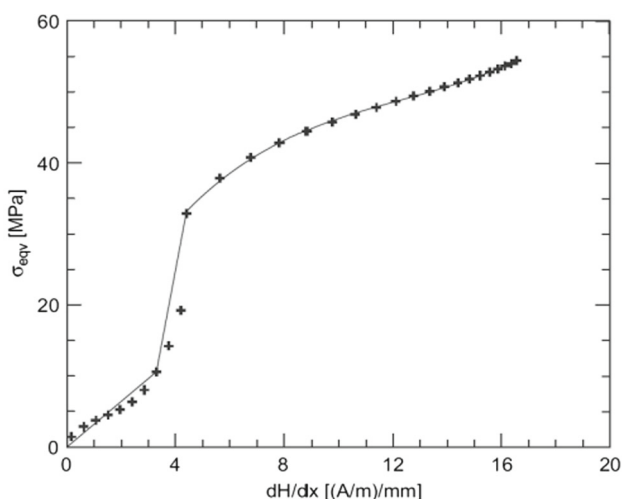


Fig. 2 Inverse relationship $\sigma_{eqv} \sim dH/dx$, steel T/P24 (reprinted from [5])

3 Finite Element Modeling and Its Validation

In order to perform an adequate interpretation of experimental results, FEM magnetostatic modeling in ANSYS software was used. A representative model (corresponding to experiment ‘E1’) is presented in Fig. 5.

Fig. 3 Geometry of the sample used in experiment 'E2'. Uniform thickness $t = 5$ mm

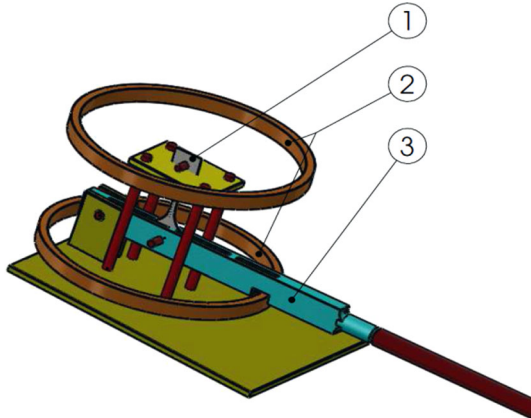
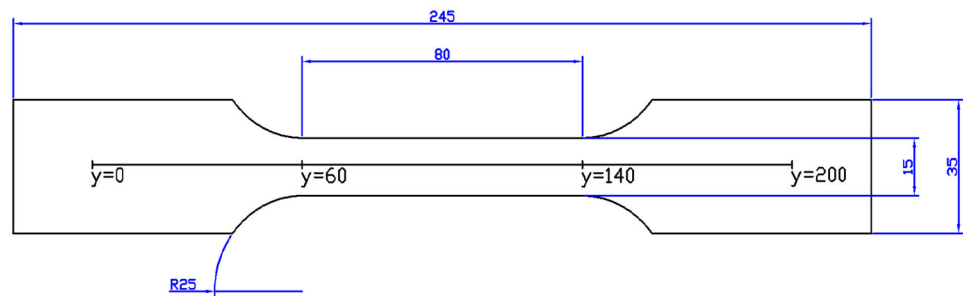


Fig. 4 Diagram of apparatus used in experiment 'E2'. 1—sample, 2—Helmholtz coils, 3—tensile machine made of austenitic steel (not used in the experiment)

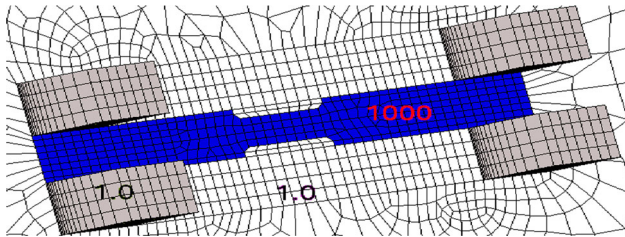


Fig. 5 FEM model of the reference sample in experimental series 'E1'. Numbers indicate typical values of the relative magnetic permeability of each material. Four blocks represent the clamps of the tensile test machine

Dimensions defined in the 3D FEM model are the same as in the experiment 'E1' by Roskosz and Bieniek. Modeled sample is surrounded by solid finite elements representing air. Ambient magnetic flux density in the reference case is equal to $40 \mu\text{T}$ (or ~ 35 A/m in terms of magnetic field strength), a value typical for the tangential component of the Earth's magnetic field. Presence of a ferromagnetic material locally modifies this value according to the Maxwell equations, namely the Gauss's law and Ampere's law. These relationships are compiled by the program into the matrix form and solved using sparse direct algorithm. Although there are four clamps in the model, in presented simulations they are inactive and have properties similar to the

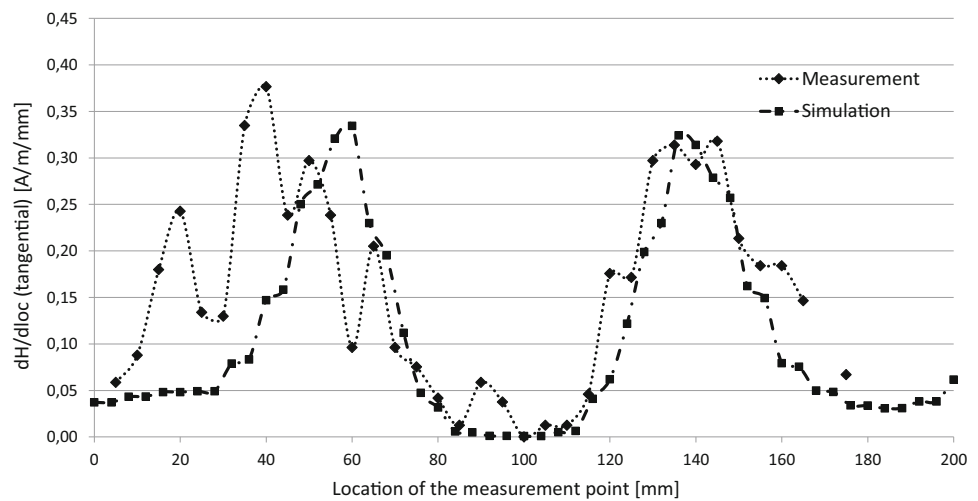
surrounding air. Magnetic parameters of the sample, namely its remanence and relative permeability, are variables. The critical region of the model, i.e. a 10 mm thick air space just above the sample, features regular 3D mesh arranged in layers of progressively increasing thickness.

Remanence modeling in ANSYS is possible in a simplified manner, as described below. The algorithm requires entry of magnetization level 'MG' [A/m] which has physical meaning of coercive force. Introduction of 'MG' parameter entails the displacement of the $B - H$ curve along the horizontal i.e. H axis, while its shape is preserved. Consequently, any (B, H) point located on the linear portion of the hysteresis loop can be accurately modeled, provided that the magnetization is monotonic.

It is relevant to demonstrate, that the adopted modeling scheme produces realistic results. A simulation was thus performed on a model representing the 'E2' sample, placed inside the Earth's field. The axial variation of the absolute gradient of H is shown in Fig. 6. We decided to present the gradient, and not the smooth $H(X)$, because the spatial gradient is the critical parameter claim we discuss, opposing the claims of Dubov and others. The gradient is calculated numerically, in an Excel sheet, taking two consecutive points into account (as defined in ISO 24497-2:2007, par. 7.2). Both the experimental and simulated $H(x)$ functions are relatively smooth, but they contain some noise which becomes apparent on the plot of gradient. In case of simulations, some noise is due to the relative coarse mesh sizing and some asymmetry in element shapes between the left and right side of the sample. The measurement is more noisy than the simulation, especially around the left curvature of the sample. These disturbances may be due to the manual finishing of the curved surfaces, and some non-uniform initial state of the material. Both these effects may come about in a real object studied with MFL.

Although there is some discrepancy between measured and computed function around the left peak, a good agreement has been obtained in terms of the maximum of $\text{abs.grad}(H)$, on which the "MMM" procedure relies. This result, together with those from previous works, [28–32] confirms that the magnetic FEM is capable of reproducing space and time field distribution in 3D steel objects of arbitrary shape.

Fig. 6 Measurement-based (new experiment—‘E2’) versus computed absolute values of tangential field strength gradient; sample initially demagnetized, unstressed, inside the Earth’s field



4 Numerical Sensitivity Analysis of the Rinterpreted Experiment (‘E1’)

A sensitivity analysis was performed using magnetostatic FEM. The sample had dimensions reproducing the experiment ‘E1’ by Roskosz and Bieniek. A linear magnetic permeability of 1000 and a negligible remanent induction (B_R) were assumed as a starting point. The ambient Earth’s field strength of 35 A/m was defined, in parallel to the sample’s axis. Both tangential and normal components of the stray magnetic field were calculated and stored. It was found that the behavior of $dH(\sigma)/dx$ parameter (both axial variation and dependency on studied factors) was similar for H_{TAN} and H_{NORM} , consistently with experimental observations by Roskosz and Bieniek. Consequently, for the sake of succinctness, only sensitivity analysis based on H_{TAN} is presented. The magnetic field strength H variation is shown in Fig 7a–c, while the curves in the Fig. 8a–c represent the gradient absolute value of H , for the same parameters made variable.

4.1 Sensitivity Analysis of Field Strength Axial Distribution $H(x)$

The axial variation of the stray magnetic field (H_{TAN}) depends in a different way on each studied factor (B_R , A , μ_r).

The influence of uniform remanent induction B_R is nearly linear. Realistic values of B_R (0.2–0.4 T) produce stray field levels exceeding by an order of magnitude the Earth’s ambient field. H_{TAN} changes its sign along with the change of sign of B_R , and the axial variation of H_{TAN} turns from convex to concave.

The influence of the geometrical discontinuity (A) is slightly nonlinear, yet monotonic. Lack of curvature ($A = B$) produces a nearly constant H_{TAN} along the sample’s axis. Increase of $A:B$ ratio results in deepening of the middle min-

imum of the curve. However, even the extreme H does not exceed the value of the applied external, ambient Earth’s field (–35 A/m).

The influence of μ is notably non-linear and non-monotonic. There is an extremum of H_{TAN} observed for $\mu \sim 100$. H_{TAN} decreases to nearly zero for $\mu \sim 10,000$, adopting some intermediate levels for $\mu \sim 1000$, which is the typical order of magnitude of initial magnetization of structural steels. Very low μ makes the sample similar to the surrounding air, and consequently H_{TAN} becomes close to the ambient Earth’s field strength H_E .

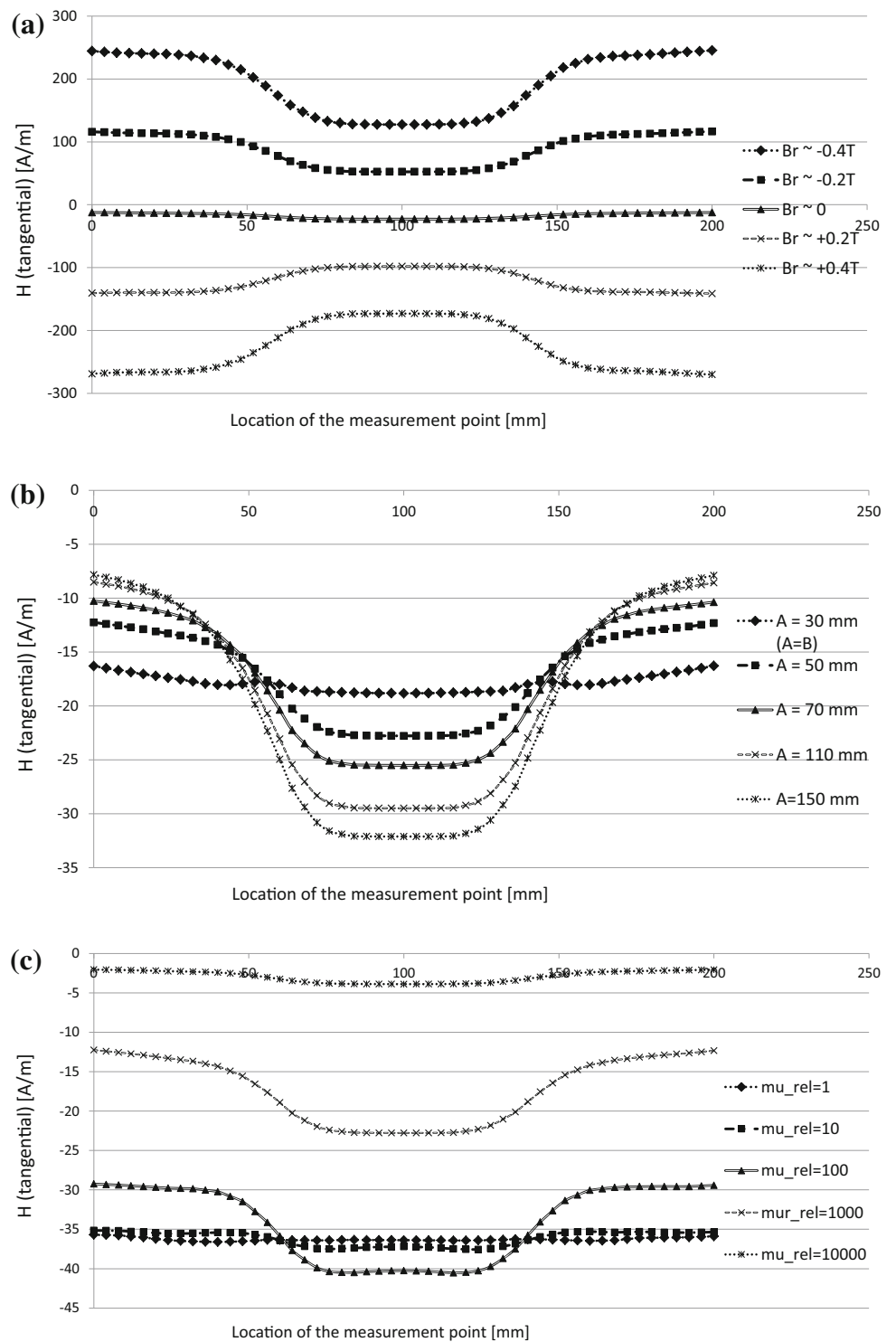
4.2 Sensitivity Analysis of Absolute Value of Field Gradient Axial Distribution $|dH/dx|$

The absolute value of gradient of H_{TAN} ($|dH_{TAN}/dx|$) adopts a nearly identical, two-peak variation regardless the studied factor, except for the case of $A = B$ (no geometrical discontinuity). The maxima of $|dH_{TAN}/dx|$ are located over the discontinuities. The highest value of $|dH_{TAN}/dx|$ (~ 4 (A/m)/mm) occurs at $B_R = -0.4$ T. Without the remanent induction, the maxima of $|dH_{TAN}/dx|$ reach values about an order of magnitude lower. The largest geometrical discontinuity ($A = 150$ mm) produces the maximum of $|dH_{TAN}/dx|$ equal to 0.7 (A/m)/mm, and when the initial permeability is variable, then the maximum of $|dH_{TAN}/dx|$ (~ 0.45 (A/m)/mm) is found for $\mu = 100$.

5 Discussion of Results

The essential question is: could the field gradients measured over S235 and T/P24 steel samples be due to the residual stress alone, as postulated by Roskosz and Bieniek? If that statement was true, then an inverse relationship $\sigma(|dH_{TAN}/dx|)$ could indeed be defined, and quantitative,

Fig. 7 Tangential magnetic field strength along the sample axis for varying: **a** uniform remanent induction, **b** width of external segments, **c** uniform magnetic permeability

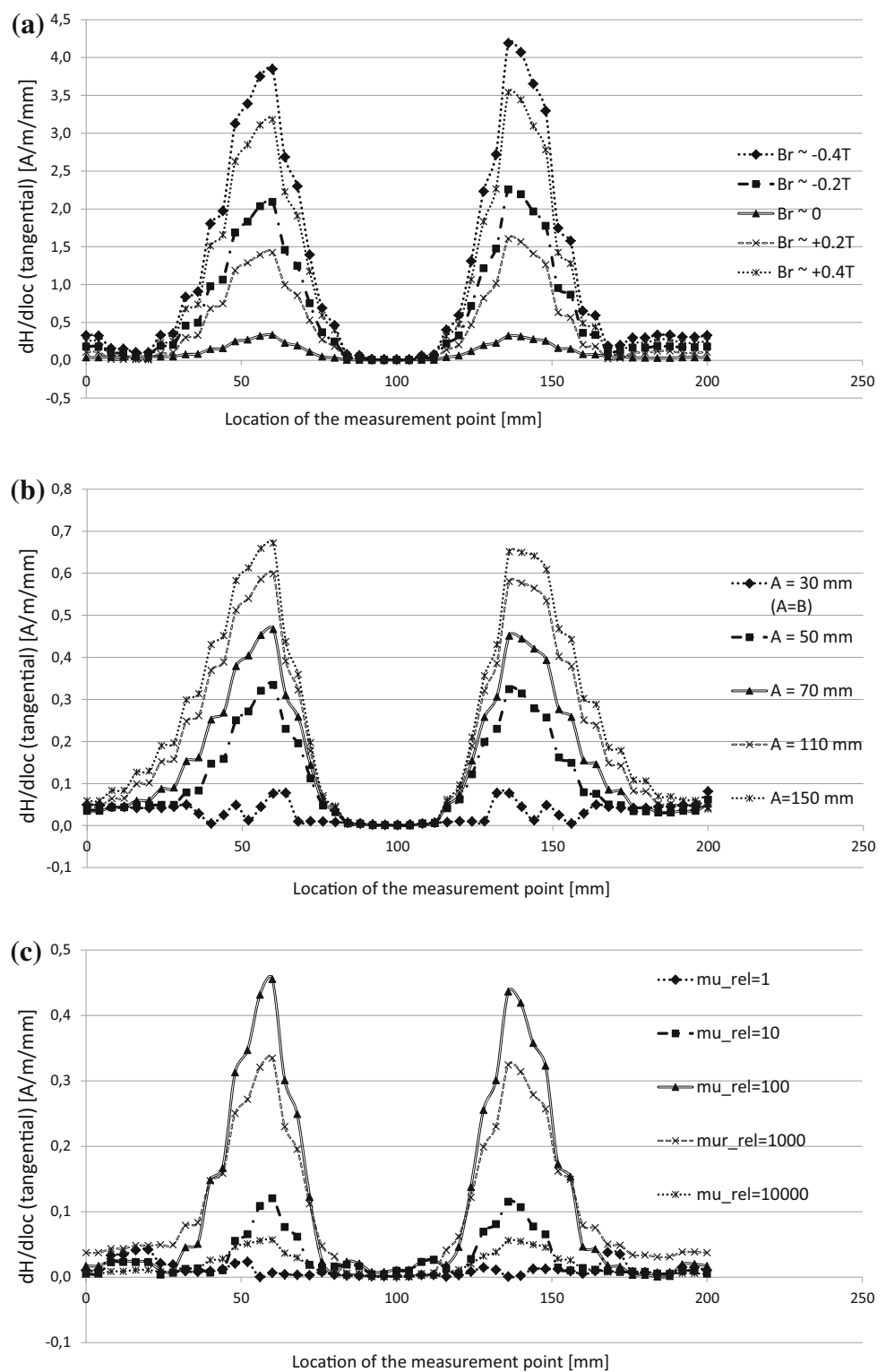


passive static stray field based stress determination would be possible.

Our FEM-based sensitivity analysis along with a consideration of dependencies between structural and magnetic parameters leads to a firm negative answer to the posed question. The main arguments are listed below.

(a) *The relationship between the material's structural state and the magnetic stray field are non-linear and non-local. It is well-known, that the stray field around a ferromagnetic object does not depend in a straightforward manner on the structural stress. Firstly, there is no local-local dependence. Solving Maxwell's equations,*

Fig. 8 Gradient absolute value of tangential magnetic field strength along the sample axis for varying: **a** uniform remanent induction, **b** width of external segments, **c** uniform magnetic permeability



one has to consider an entire magnetic system, composed of various segments of the sample, the ambient air, and possibly other components, such as e.g. the clamps of the tensile test machine. Any change to one of these components, changes the characteristics of the

magnetic circuit, thus provoking modification of stray field level and distribution. Secondly, the stress influences more than one local magnetic characteristic of the material, influencing the effective magnetic permeability, remanent induction, and coercivity. Moreover, these

influences are history-dependent, because a cyclically-magnetized ferromagnetic material may switch between different hysteresis loops, ranging from the anhysteretic behavior (initial magnetization curve), through minor loops up to the major hysteresis loop [33]. Finally, the residual stress depends on applied load and stress intensity factor as well as the yield strength, which adds further complexity to the problem.

- (b) *There can be no inverse function of a dependency on several variables.* Roskosz and Bieniek draw a relationship $|dH/dx| = f(\sigma_{RES})$, with measured magnetic field gradient and calculated local residual equivalent stress. Further they postulate and define an inverse function, namely $\sigma_{RES} = f^{-1}(|dH/dx|)$, claimed to be monotonic and valid for T/P24 steel grade. However, as quantified in our sensitivity analysis, the local stray field itself and its gradient are a function of at least three variables, namely the remanent induction B_R , geometric discontinuity factor A , and magnetic permeability μ :

$$H_{LOCAL} = g_1(B_R) + g_2(A) + g_3(\mu_r) \quad (2)$$

Furthermore, each of these functions can be separated into a stress-dependent, and a non-stress-dependent contribution, and the presented simple arithmetic sum is only a symbolic representation of a much more complex interdependence between the variables.

- (c) *Localized stresses of order of 50 MPa in a demagnetized sample cannot produce measurable variation of the stray field.* This fact was reported by [8], who gave an explanation derived from Jiles' theory. Our sensitivity analysis corroborates that statement. Based on studies by Anglada and Zurek [11,24] we estimated that the uniform static stress of ~ 50 MPa along with plastic strain of $\sim 30\%$ can produce ca 50% decrease in effective magnetic permeability and $\sim 50\%$ increase of remanent induction. The effect of localized stress and yield should be even smaller. This variation is not sufficient to produce stray fields and field gradients reported by [5]. Stray field levels measured by Roskosz and Bieniek are of order of 100–200 A/m, and absolute gradients of H reach 5 (A/m)/mm (S235 steel grade) and 20 (A/m)/mm (T/P24 steel grade). Comparing these values to those obtained in FEM-based sensitivity analysis one finds, that only an elevated residual induction exceeding 0.4 T could account for so high magnetic field strength. Indeed, the clamps of the tensile test machine were ferromagnetic, and their remanent induction was neither eliminated nor evaluated during the experiments by Roskosz and Bieniek. During sequential loading, the sample could acquire high, uncontrolled residual magnetization, being at origin of H levels exceeding 100 A/m.

- (d) *The same value of local $|dH/dx|$ may correspond to different local/global stress states.* The FEM-based sensitivity analysis shows, that using the $|dH/dx|$ parameter to determine localized stress levels inevitably leads to ambiguous stress assessment. In the dog-bone sample studied, the axial variation of $|dH/dx|$ is qualitatively identical for all the varied cases, except for a sample without any geometrical discontinuity. Quantification of the $|dH/dx|$ peak does not help either. We found that distinct cases, with possibly various stress levels, resulted in the same maximum of $|dH/dx|$.

6 Conclusion

It was found that the stray field gradients reported in the work by Roskosz and Bieniek are quantitatively due to uncontrolled remanent induction. Quantitatively their characteristic two peaks are primarily due to the geometrical discontinuity of the sample.

Several authors present an experimental proof that increasing local stress level or inducing plastic strain makes the strain field varies. However, this does not justify any claim of existence of a single inverse function. The inverse function cannot be defined, mainly because the original magnetic field dependence is a function of several variables. The calculated sensitivity analysis demonstrated the ambiguity of static stray field gradient measurements.

Taking into account our overall experience with electromagnetic NDT, and in view of recently published papers [26,27], the quantitative in-situ NDT based on passive measurement of magnetic stray-field is impossible, despite the attractiveness of its idea.

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