



Higher order asymptotic expansions to the solutions for a nonlinear damped wave equation

Tatsuki Kawakami and Hiroshi Takeda

Abstract. We study the Cauchy problem for a nonlinear damped wave equation. Under suitable assumptions for the nonlinearity and the initial data, we obtain the global solution which satisfies weighted L^1 and L^∞ estimates. Furthermore, we establish the higher order asymptotic expansion of the solution. This means that we construct the nonlinear approximation of the global solution with respect to the weight of the data. Our proof is based on the approximation formula of the linear solution, which is given by Takeda (Asymptot Anal 94:1–31, 2015), and the nonlinear approximation theory for a nonlinear parabolic equation developed by Ishige et al. (J Evol Equ 14:749–777, 2014).

Mathematics Subject Classification. 35L15, 35L71, 35B40.

Keywords. Asymptotic expansion, Large time behavior, Nonlinear damped wave equations, Nonlinear approximation.

1. Introduction

We consider the Cauchy problem for a nonlinear damped wave equation,

$$\begin{cases} \partial_t^2 u - \Delta u + \partial_t u = F(u) & \text{in } \mathbb{R}^N \times (0, \infty), \\ u(x, 0) = u_0(x), \quad \partial_t u(x, 0) = u_1(x) & \text{in } \mathbb{R}^N, \end{cases} \quad (1.1)$$

where $N = 1, 2, 3$, and $\partial_t = \partial/\partial t$. We assume that the nonlinear term $F \in C(\mathbb{R})$ satisfies

$$|F(\xi)| \leq C|\xi|^p, \quad \xi \in \mathbb{R}, \quad (1.2)$$

$$|F(\xi) - F(\eta)| \leq C(|\xi|^{p-1} + |\eta|^{p-1})|\xi - \eta|, \quad \xi, \eta \in \mathbb{R}, \quad (1.3)$$

for some constants $p > 1$ and $C > 0$, which are independent of ξ and η . The typical examples of our nonlinear terms are given by

$$F(\xi) = \pm|\xi|^{p-1}\xi \quad \text{or} \quad \pm|\xi|^p \quad (p > 1). \quad (1.4)$$

The aim of this paper is to study the large time behavior of the solution to (1.1). More precisely, we show the nonlinear approximation of the solution to (1.1) with respect to the order of the moment for the initial data. This point of view is shared with Ishige and the first author of the paper [17] studying a nonlinear heat equation.

In the classical paper [28], Matsumura considered the Cauchy problem of nonlinear wave equations with dissipation terms. His main tools in the proof are the estimates for the solutions to a linear damped wave equation

$$\begin{cases} \partial_t^2 v - \Delta v + \partial_t v = 0 & \text{in } \mathbb{R}^N \times (0, \infty), \\ v(x, 0) = v_0(x), \quad \partial_t v(x, 0) = v_1(x) & \text{in } \mathbb{R}^N. \end{cases} \quad (1.5)$$

Especially, he proved the decay estimates for the solution to (1.5) by the Fourier splitting method. Beginning this paper, many authors showed the large time behavior of the solution to (1.5) (see, e.g., [2, 23, 43]). In [33], Orive, Zuazua and Pozato obtained the higher order asymptotic expansion of the solution to (1.5) for the variable coefficient setting in L^2 base framework. Furthermore, the decomposition of the solution into the solutions of heat equations and wave equations are proposed (see, e.g., [27] for $N = 1$, [10] for $N = 2$, [30] for $N = 3$ and [29] for $N \geq 4$).

On the other hand, the Cauchy problem (1.1) with (1.4) has been studied by many mathematicians from various points of view. Especially, for the focusing case $F(\xi) = |\xi|^p$, it is well-known that the growth order $p = 1 + 2/N$ is critical situation for the existence of the global solution to (1.1). In [38], Todorova and Yordanov proved that, if $p > 1 + 2/N$, then there exists a unique global solution for the small compactly supported data, and if $1 < p < 1 + 2/N$, then the solution blows up in a finite time for small data (see also [26] for blow up results). Zhang [44] showed the small data blow-up results including $p = 1 + 2/N$. Ikehata and Tanizawa [14] obtained the global existence results for $p > 1 + 2/N$ under the non-compact supported data assumption (see, e.g., [12] for the case $F(\xi) = |\xi|^{p-1}\xi$). For the defocusing case $F(\xi) = -|\xi|^{p-1}\xi$ or $F(\xi) = -|\xi|^p$, it is also well-investigated, and it is well known that there exists a unique global solution with decay property for all $p > 1$ if data has sufficient regularity (see, e.g., [9, 13, 22]).

For the asymptotic profiles of the solution to damped wave equations, so-called diffusion phenomena is shown by many authors. Among others, Gallay and Raugel [7] proved that global solutions of nonlinear damped wave equation behaves like those of nonlinear heat equations with suitable data, including more general nonlinearity for $N = 1$. In [20], Karch proved the approximation of the solution to (1.1) by the heat kernel for $p \geq 1 + 4/N$. After that, Nishihara [30] proved it for $p > 1 + 2/N$ when $N = 3$. (See also [27] for $N = 1$, [10] for $N = 2$, [29] for $N = 4, 5$ and [8] for $N \geq 1$). In [21], Ueda and the first author of this paper obtained the second order nonlinear approximation of the solution to (1.1) for $p > 1 + 2/N$. Recently, the second author of this paper [37] proved the $(K + 1)$ th order expansion of the solution by the series of the heat kernel when $p > 1 + (K + 2)/N$.

We should also state the topic on the recent progress of the diffusion phenomena in damped wave equations. One of our motivation here, the precise description of the large time behavior of the solution, is shared with the related results for the following equations:

- The variable coefficient damping case (see, e.g., [24, 25, 31, 32, 39–42])

$$\partial_t^2 v - \Delta v + b(t, x)\partial_t v = 0 \quad \text{in } \mathbb{R}^N \times (0, \infty);$$

- The structural damped wave equations (see, e.g., [11, 15, 19, 34, 35] for $\sigma = 1$ and [4–6] for $0 < \sigma \leq 1$)

$$\partial_t^2 v - \Delta v + (-\Delta)^\sigma \partial_t v = 0 \quad \text{in } \mathbb{R}^N \times (0, \infty).$$

We remark that based on the linear estimates, the authors also treated the nonlinear perturbation by the various methods to have the diffusion phenomena.

As we have seen the above, there are many results for the diffusion phenomena of the dissipative type wave equations. On the other hand, there are few studies concerning the higher order asymptotic expansion of the global solutions of (1.1) except for the results [7] (1-d case), [21] (up to second order expansion) and [37] (under the regularity for the nonlinearity). Roughly speaking, the difficulty stems from the construction of the approximation functions as $t \rightarrow \infty$ and the singularity from the high frequency part. To avoid the difficulty, the previous results need to restrict the order of the expansion or to only treat the smooth nonlinearity. In this paper, we established the asymptotic expansion of the solution to (1.1) for (1.4) with $p > 1 + 2/N$ ($N = 1, 2$), $p \geq 2$ ($N = 3$) up to the suitable order, which depends on the order of the moment for the initial data. Our proof is based on the higher order asymptotic expansion formula of the linear solution by the solution of the heat equation, which is shown by the second author of this paper in [36], and nonlinear approximation technique for nonlinear parabolic equations developed by Ishige, Kobayashi and the first author of the paper in [18]. Our new ingredient here is to show the weighted L^∞ estimates for the solution to (1.5), which are important to apply the iteration scheme proposed in [18] to our problem (1.1). Furthermore, our results imply not only the detailed profile of the solution to (1.1) for large t , but also sharp decay estimates in each expansion order.

This paper is organized as follows. In Sect. 2, we prepare notations which used throughout this paper and we state the main results in this paper. Furthermore we mention important remarks on main results. Section 3 presents some preliminaries. Section 4 is devoted to the study of the weighted L^1 and L^∞ estimates for the solutions to (1.5). In Sects. 5 and 6, our main results are proved.

2. Main results

2.1. Notation

To state our results precisely, we summarize notion and notation. For $k \geq 0$, we denote $[k]$ the integer satisfying $k - 1 < [k] \leq k$. Let $\mathbb{N}_0 = \mathbb{N} \cup \{0\}$, $\mathbb{M} = \mathbb{N}_0^N$, and G be the N -dimensional heat kernel, that is,

$$G(x, t) := (4\pi t)^{-N/2} \exp\left(-\frac{|x|^2}{4t}\right). \tag{2.1}$$

Furthermore, for any $\phi \in L^q(\mathbb{R}^N)$ with $q \in [1, \infty]$, we denote by $e^{t\Delta}\phi$ the unique bounded solution for the Cauchy problem of the heat equation with the initial datum ϕ , that is,

$$(e^{t\Delta}\phi)(x) := \int_{\mathbb{R}^N} G(x - y, t)\phi(y)dy. \tag{2.2}$$

For any $\alpha = (\alpha_1, \dots, \alpha_N) \in \mathbb{M}$, we put

$$|\alpha| := \sum_{i=1}^N |\alpha_i|, \quad \alpha! := \prod_{i=1}^N \alpha_i!, \quad x^\alpha := \prod_{i=1}^N x_i^{\alpha_i}, \quad \partial_x^\alpha := \frac{\partial^{|\alpha|}}{\partial x_1^{\alpha_1} \dots \partial x_N^{\alpha_N}},$$

$$g(x, t) := G(x, 1 + t), \quad g_\alpha(x, t) := \frac{(-1)^{|\alpha|}}{\alpha!} \partial_x^\alpha g(x, t).$$

Let $\mathbb{M}_k := \{\alpha \in \mathbb{M} : |\alpha| \leq k\}$ for $k \geq 0$. For any $\alpha = (\alpha_1, \dots, \alpha_N), \beta = (\beta_1, \dots, \beta_N) \in \mathbb{M}$, we say

$$\alpha \leq \beta$$

if $\alpha_i \leq \beta_i$ for all $i \in \{1, \dots, N\}$. For ℓ, m and $n \in \mathbb{N}_0$, we put

$$\phi_\ell(r) := \left(\frac{2}{1 + \sqrt{1 - 4r}}\right)^{2\ell}, \quad \psi(r) := \frac{1}{\sqrt{1 - 4r}},$$

and

$$\Phi_{\ell, m} := \frac{1}{\ell! m!} \frac{d^m}{dr^m} \phi_\ell(r) \Big|_{r=0}, \quad \Psi_n := \frac{1}{n!} \frac{d^n}{dr^n} \psi(r) \Big|_{r=0}. \tag{2.3}$$

For any $q \in [1, \infty]$, we denote by L^q and $\|\cdot\|_{L^q}$ the usual $L^q(\mathbb{R}^N)$ space and its norm, respectively. Let $\ell \in \mathbb{N}_0$. Then $W^{\ell, q}$ denotes the Sobolev space of L^q functions, equipped with the norm

$$\|\phi\|_{W^{\ell, q}} := \left(\sum_{|\alpha| \leq \ell} \|\partial_x^\alpha \phi\|_{L^q}^q\right)^{1/q}.$$

For any $k \geq 0$ and $q \in \{1, \infty\}$, we denote by L_k^q the functional space $L^q(\mathbb{R}^N, (1 + |x|)^k dx)$, and put

$$\|\phi\|_{L_k^q} = \|\phi\|_{L^q((1+|x|)^k dx)}, \quad \|\phi\|_{W_k^{\ell, q}} = \sum_{|\alpha| \leq \ell} \|\partial_x^\alpha \phi\|_{L_k^q}.$$

Here we often identify $W_k^{0,1} = L_k^1$. Throughout the present paper, C denotes a various generic positive constant.

Let us give the definition of the solution for the Cauchy problem (1.1).

Definition 2.1. *Let $u \in C([0, \infty) : L^1) \cap L^\infty(0, \infty : L^\infty)$ and $F \in C(\mathbb{R})$. Then the function u is said to be a solution for the Cauchy problem (1.1) if there holds*

$$u(x, t) = (K_0(t)u_0)(x) + \left(K_1(t) \left(\frac{1}{2}u_0 + u_1 \right) \right) (x) + \int_0^t (K_1(t-s)F(u(s)))(x) ds \tag{2.4}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$, where the evolution operators $K_0(t)$ and $K_1(t)$ of the linear damped wave equation are given by

$$\begin{aligned} (K_0(t)\phi)(x) &:= \mathcal{F}^{-1} \left[e^{-\frac{t}{2}} \cos \left(t\sqrt{|\xi|^2 - \frac{1}{4}} \right) \mathcal{F}[\phi] \right] (x), \\ (K_1(t)\phi)(x) &:= \mathcal{F}^{-1} \left[e^{-\frac{t}{2}} \frac{\sin \left(t\sqrt{|\xi|^2 - \frac{1}{4}} \right)}{\sqrt{|\xi|^2 - \frac{1}{4}}} \mathcal{F}[\phi] \right] (x). \end{aligned} \tag{2.5}$$

Here we denote the Fourier and Fourier inverse transform by \mathcal{F} and \mathcal{F}^{-1} , respectively.

Let $M_\alpha(f, t)$ be the constant defined inductively (in α) by

$$\begin{aligned} M_0(f, t) &:= \int_{\mathbb{R}^N} f(x) dx, \quad \text{if } \alpha = 0, \\ M_\alpha(f, t) &:= \int_{\mathbb{R}^N} x^\alpha f(x) dx - \sum_{\beta \leq \alpha, \beta \neq \alpha} M_\beta(f, t) \int_{\mathbb{R}^N} x^\alpha g_\beta(x, t) dx \quad \text{if } \alpha \neq 0. \end{aligned} \tag{2.6}$$

By (2.3) and (2.6) we introduce the function $U_{lin}(t) = U_{lin}(x, t)$ by

$$\begin{aligned} U_{lin}(x, t) &= \sum_{\ell=0}^{[K/2]} \sum_{m=0}^{[K/2]-\ell} \Phi_{\ell, m} \left\{ \frac{1}{2} \sum_{|\alpha| \leq K-2(\ell+m)} M_\alpha(u_0, 0) (-t)^\ell (-\Delta)^{2\ell+m} g_\alpha(x, t) \right. \\ &+ \sum_{n=0}^{[K/2]-\ell-m} \sum_{|\alpha| \leq K-2(\ell+m+n)} \Psi_n \left(\frac{1}{2} M_\alpha(u_0, 0) + M_\alpha(u_1, 0) \right) \\ &\left. \times (-t)^\ell (-\Delta)^{2\ell+m+n} g_\alpha(x, t) \right\}. \end{aligned} \tag{2.7}$$

Following the notation of [16], we also introduce the linear operator P_K on L^1_K by

$$[P_K(t)f](x) := f(x) - \sum_{|\alpha| \leq K} M_\alpha(f, t) g_\alpha(x, t), \tag{2.8}$$

for $K \geq 0$ and $t \geq 0$ and $f \in L^1_K$. Then the operator $P_K(t)$ has the following property,

$$\int_{\mathbb{R}^N} x^\alpha [P_K(t)f](x) dx = 0, \quad t \geq 0. \tag{2.9}$$

for any $\alpha \in \mathbb{M}_K$ (see ‘‘Appendix’’). This property plays an important role of deriving our main results (for the detail, see also [16, 17, 21]).

With the notation and the notion above, we can define the sequence of functions $U_j = U_j(x, t)$ inductively: Let $K \geq 0$ and $(u_0, u_1) \in W_K^{[N/2],1} \cap W^{[N/2],\infty} \times L_K^1 \cap L^\infty$. For any $j \in \mathbb{N}$,

$$\begin{aligned}
 U_0(t) &:= U_{\text{lin}}(t) + \sum_{|\alpha| \leq K} \int_0^t K_1(t-s) M_\alpha(F(u(s)), s) g_\alpha(s) ds, \\
 U_j(t) &:= U_0(t) + \int_0^t K_1(t-s) P_K(s) F_{j-1}(s) ds,
 \end{aligned}
 \tag{2.10}$$

where $F_j(x, t) := F(U_j(x, t))$. Here $K_1(t)$ and U_{lin} are given in (2.5) and (2.7), respectively.

2.2. Main theorems

Now we are ready to treat our main results. Our first result gives the sufficient condition for the existence of global solution satisfying the weighted L^1 and L^∞ estimates with suitable decay property.

Theorem 2.1. *Let $K \geq 0$ and $(u_0, u_1) \in W_K^{[N/2],1} \cap W_K^{[N/2],\infty} \times L_K^1 \cap L_K^\infty$. Assume that $F \in C(\mathbb{R})$ satisfies (1.2) and (1.3) with*

$$p > p_* := 1 + \frac{2}{N}.
 \tag{2.11}$$

Then there exists a positive constant ε such that if $E_0 := \|u_0\|_{W^{[N/2],\infty}} + \|u_0\|_{W^{[N/2],1}} + \|u_1\|_{L^\infty} + \|u_1\|_{L^1} \leq \varepsilon$, then the Cauchy problem (1.1) admits a unique global solution u of (1.1) in the class

$$C([0, \infty) : L^1) \cap L^\infty(0, \infty : L^\infty)$$

satisfying

$$\|u(t)\|_{L^q} \leq CE_0(1+t)^{-\frac{N}{2}(1-\frac{1}{q})}, \quad t \geq 0,
 \tag{2.12}$$

for any $q \in [1, \infty]$. Moreover,

$$u \in C([0, \infty) : L_K^1) \cap L^\infty(0, \infty : L_K^\infty)$$

and

$$\|u(t)\|_{L_k^1} \leq CE_K(1+t)^{\frac{k}{2}}, \quad t \geq 0,
 \tag{2.13}$$

$$\|u(t)\|_{L_k^\infty} \leq CE_K(1+t)^{\frac{k-N}{2}}, \quad t \geq 0,
 \tag{2.14}$$

for any $k \in [0, K]$, where

$$E_K := \max_{k_i \in [0, K], i \in \{1, 2, 3, 4\}} \left\{ \|u_0\|_{W_{k_1}^{[N/2],\infty}} + \|u_0\|_{W_{k_2}^{[N/2],1}} + \|u_1\|_{L_{k_3}^\infty} + \|u_1\|_{L_{k_4}^1} \right\}.
 \tag{2.15}$$

Remark 2.1. Here we note that Theorem 2.1 states that to obtain the suitable estimates (2.13) and (2.14), we only assume that the smallness of E_0 , not E_K .

Our second result is the nonlinear approximation of the solutions to (1.1). In other words, the functions U_j defined by (2.10) are nonlinear approximation of the global solution to (1.1). (See also Remark 3.2.)

Theorem 2.2. *Let $j \in \mathbb{N}_0$, and let U_j be the functions given in (2.10). Put*

$$A := \frac{N}{2}(p-1) - 1 = \frac{N}{2}(p-p_*) > 0. \quad (2.16)$$

Assume that

$$p > p_* \quad (N = 1, 2), \quad p \geq 2 \quad (N = 3), \quad (2.17)$$

and there exists a unique global solution to (1.1) satisfying (2.12), (2.13) and (2.14). Then, for any $q \in [1, \infty]$, $k \in [0, K]$ and $\gamma \in \{1, \infty\}$,

$$\sup_{t>0} (1+t)^{\frac{N}{2}(1-\frac{1}{q})} \|U_j(t)\|_{L^q} + \sup_{t>0} (1+t)^{-\frac{k}{2} + \frac{N}{2}(1-\frac{1}{\gamma})} \|U_j(t)\|_{L_k^\gamma} < \infty \quad (2.18)$$

and

$$t^{\frac{N}{2}(1-\frac{1}{q})} \|u(t) - U_j(t)\|_{L^q} = \begin{cases} o(t^{-\frac{K}{2}}) + O(t^{-(j+1)A}) & \text{if } (j+1)A \neq K/2, \\ O(t^{-\frac{K}{2}} \log t) & \text{if } (j+1)A = K/2, \end{cases} \quad (2.19)$$

as $t \rightarrow \infty$.

Remark 2.2. Theorem 2.2 shows that once we have a unique solution to (1.1) satisfying the estimates (2.12), (2.13) and (2.14), without smallness of the data, then we see that the sequence of the functions U_j is well-defined in the sense of (2.18) and we have the asymptotic behavior (2.19) of the solutions as $t \rightarrow \infty$. We also note that the estimate (2.19) is easily rephrased as follows: For any fixed $p > p_*$ ($N = 1, 2$), $p \geq 2$ ($N = 3$) and $K \geq 0$, there exists a number $n_* \in \mathbb{N}_0$ such that

$$t^{\frac{N}{2}(1-\frac{1}{q})} \|u(t) - U_{n_*}(t)\|_{L^q} = o(t^{-\frac{K}{2}}) \quad \text{as } t \rightarrow \infty,$$

for any $q \in [1, \infty]$, which may be a more distinct description of the large time behavior of the global solution u . This convergence rate is same as the case $F \equiv 0$. See Corollary 3.1.

3. Preliminaries

3.1. Solutions of the heat equation

In this subsection we recall some preliminary results on the behavior of solutions for the heat equation and the operator $P_k(t)$.

Let $\alpha \in \mathbb{M}$ and G be the function given in (2.1). Then we have

$$|\partial_x^\alpha G(x, t)| \leq C t^{-\frac{N+|\alpha|}{2}} \left[1 + \left(\frac{|x|}{t^{1/2}} \right)^{|\alpha|} \right] \exp\left(-\frac{|x|^2}{4t}\right) \quad (3.1)$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. This inequality yields the inequalities

$$\|g_\alpha(t)\|_{L^q} \leq C(1+t)^{-\frac{N}{2}(1-\frac{1}{q}) - \frac{|\alpha|}{2}}, \quad t > 0, \quad (3.2)$$

$$\|g_\alpha(t)\|_{L_k^\gamma} \leq C(1+t)^{\frac{k-|\alpha|}{2} - \frac{N}{2}(1-\frac{1}{\gamma})}, \quad t > 0, \quad (3.3)$$

for any $q \in [1, \infty]$, $k \geq 0$ and $\gamma \in \{1, \infty\}$. Furthermore, applying the Young inequality to (2.2) with the aid of (3.1), for any $\alpha \in \mathbb{M}$ and $1 \leq r \leq q \leq \infty$, we have

$$\|\partial_x^\alpha e^{t\Delta} \phi\|_{L^r} \leq C t^{-\frac{N}{2}(\frac{1}{r} - \frac{1}{q}) - \frac{|\alpha|}{2}} \|\phi\|_{L^r}, \quad t > 0.$$

In particular, we obtain

$$\|e^{t\Delta} \phi\|_{L^r} \leq \|\phi\|_{L^r}, \quad t > 0.$$

We recall the following lemma, which is useful in our study for the asymptotic expansion of solutions.

Lemma 3.1. [18, Proposition 3.1] *Let $K \geq 0$. Then the following holds.*

(i) *For any $k \in [0, K]$,*

$$\int_{\mathbb{R}^N} |x|^k |e^{t\Delta} \phi(x)| \, dx \leq C t^{-\frac{K-k}{2}} \int_{\mathbb{R}^N} |x|^K |\phi(x)| \, dx, \quad t > 0,$$

for all $\phi \in L^1_K$ satisfying

$$\int_{\mathbb{R}^N} x^\alpha \phi(x) \, dx = 0, \quad \alpha \in \mathbb{M}_K. \tag{3.4}$$

(ii) *For any $k \in [0, K]$,*

$$\lim_{t \rightarrow \infty} t^{\frac{K-k}{2}} \int_{\mathbb{R}^N} |x|^k |e^{t\Delta} \phi(x)| \, dx = 0$$

for $\phi \in L^1_K$ satisfying (3.4).

Let us mention the important fact to show the proof of Theorem 2.2. The point of the following Proposition 3.1 is in the assertion that $e^{t\Delta} \phi$ is well-approximated by the sequence of $g_\alpha(t)$, not $\partial_x^\alpha G(x, t)$.

Proposition 3.1. [18, Theoreme 1.1] *Let $K \geq 0$ and $\phi \in L^1_K$. Then, for any $j \in \mathbb{N}_0$ and $q \in [1, \infty]$,*

$$t^{\frac{N}{2}(1-\frac{1}{q})+\frac{j}{2}} \left\| \nabla^j \left[e^{t\Delta} \phi - \sum_{|\alpha| \leq K} M_\alpha(\phi, 0) g_\alpha(t) \right] \right\|_{L^q} = o(t^{-\frac{K}{2}})$$

as $t \rightarrow \infty$.

At the end of this subsection we give a lemma on the estimate of the function $P_K(t)f(t)$.

Lemma 3.2. *Let $K \geq 0$ and $1 \leq q \leq \infty$. Let f be a measurable function in $\mathbb{R}^N \times (0, \infty)$ such that*

$$E_{K,q}[f](t) := (1+t)^{\frac{K}{2}} \left[(1+t)^{\frac{N}{2}(1-\frac{1}{q})} \|f(t)\|_{L^q} + \|f(t)\|_{L^1} \right] + \|f(t)\|_{L^1_K} + (1+t)^{\frac{N}{2}} \|f(t)\|_{L^\infty} \in L^\infty(0, T) \tag{3.5}$$

for any $T > 0$. Then, for any $0 \leq k \leq K$ and $\gamma \in \{1, \infty\}$,

$$|M_\alpha(f(t), t)| \leq C(1+t)^{-\frac{K-|\alpha|}{2}} E_{K,q}[f](t), \quad \alpha \in \mathbb{M}_K, \tag{3.6}$$

and

$$\begin{aligned} & (1+t)^{\frac{N}{2}(1-\frac{1}{q})} \|P_K(t)f(t)\|_{L^q} + (1+t)^{-\frac{k}{2} + \frac{N}{2}(1-\frac{1}{\gamma})} \|P_K(t)f(t)\|_{L_k^\gamma} \\ & \leq C(1+t)^{-\frac{K}{2}} E_{K,q}[f](t) \end{aligned} \tag{3.7}$$

for almost all $t > 0$.

Proof. Applying the same arguments as in the proof of [18, Lemma 2.2], we can easily prove this lemma. So we omit the proof. \square

3.2. Solutions of the damped wave equation

In this subsection, we recall some preliminary results on the properties of the fundamental solutions for the linearized damped equation (1.5). We first give the well-known representation formulas of evolution operators $K_1(t)\phi$ and $K_0(t)\phi$ (see, e.g., [3, 21, 30, 36]). For $N = 1, 2, 3$, we have

$$(K_1(t)\phi)(x) = \begin{cases} \frac{e^{-\frac{t}{2}}}{2} \int_{|z|\leq t} I_0\left(\sqrt{t^2 - |z|^2}/2\right) \phi(x+z) dz & (N = 1), \\ \frac{e^{-\frac{t}{2}}}{2\pi} \int_{|z|\leq t} \frac{\cosh\left(\sqrt{t^2 - |z|^2}/2\right)}{\sqrt{t^2 - |z|^2}} \phi(x+z) dz & (N = 2), \\ \frac{e^{-\frac{t}{2}}}{4\pi t} \partial_t \int_{|z|\leq t} I_0\left(\sqrt{t^2 - |z|^2}/2\right) \phi(x+z) dz & (N = 3), \end{cases}$$

where I_ν is the modified Bessel function of order ν defined by

$$I_\nu(y) = \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m + \nu + 1)} \left(\frac{y}{2}\right)^{2m+\nu}. \tag{3.8}$$

By changing the variable $z = ty$ we see that

$$(K_1(t)\phi)(x) = \begin{cases} \frac{e^{-\frac{t}{2}}}{2} \int_{|y|\leq 1} I_0\left(t\sqrt{1 - y^2}/2\right) \phi(x+ty) dy & (N = 1), \\ \frac{e^{-\frac{t}{2}}}{2\pi} \int_{|y|\leq 1} \frac{\cosh\left(t\sqrt{1 - |y|^2}/2\right)}{\sqrt{1 - |y|^2}} \phi(x+ty) dy & (N = 2), \\ \frac{e^{-\frac{t}{2}} t^2}{8\pi} \int_{|y|\leq 1} I_1\left(t\sqrt{1 - |y|^2}/2\right) \frac{\phi(x+ty)}{\sqrt{1 - |y|^2}} dy, \\ \quad + \frac{e^{-\frac{t}{2}} t}{4\pi} \int_{\mathbb{S}^2} \phi(x+t\omega) d\omega & (N = 3), \end{cases} \tag{3.9}$$

where $S^{N-1} := \{\omega \in \mathbb{R}^N, |\omega| = 1\}$ with its surface element $d\omega$. According to [30], for $N = 3$, we denote $J_1(t)\phi$ and $W_1(t)\phi$ by

$$(K_1(t)\phi)(x) = (J_1(t)\phi)(x) + (W_1(t)\phi)(x), \tag{3.10}$$

$$(J_1(t)\phi)(x) := \frac{e^{-\frac{t}{2}} t^2}{8\pi} \int_{|y|\leq 1} I_1\left(t\sqrt{1 - |y|^2}/2\right) \frac{\phi(x+ty)}{\sqrt{1 - |y|^2}} dy, \tag{3.11}$$

$$(W_1(t)\phi)(x) := \frac{e^{-\frac{t}{2}} t}{4\pi} \int_{\mathbb{S}^2} \phi(x+t\omega) d\omega. \tag{3.12}$$

Furthermore, it follows from the straightforward calculation (see, e.g., [27, 36]) that the representation formula for $K_0(t)\phi$ is given by

$$\begin{aligned}
 (K_0(t)\phi)(x) &= (\partial_t K_1(t)\phi)(x) + \frac{1}{2}(K_1(t)\phi)(x) \\
 &= \begin{cases} \frac{1}{2}(K_1(t)\phi)(x) + \frac{e^{-\frac{t}{2}}}{4} \int_{|y|\leq 1} I_1\left(t\sqrt{1-y^2}/2\right) \frac{\phi(x+ty)}{\sqrt{1-y^2}} dy, \\ \quad + e^{-\frac{t}{2}}(\phi(x-t) + \phi(x+t)) \quad (N=1), \\ \\ t^{-1}(K_1(t)\phi)(x) + \frac{e^{-\frac{t}{2}}t}{2\pi} \int_{|y|\leq 1} \sinh\left(t\sqrt{1-|y|^2}/2\right) \phi(x+ty) dy \\ \quad + \frac{e^{-\frac{t}{2}}t}{2\pi} \int_{|y|\leq 1} \frac{\cosh\left(t\sqrt{1-|y|^2}/2\right)}{\sqrt{1-|y|^2}} \nabla\phi(x+ty) \cdot y dy \quad (N=2), \\ \\ = 2t^{-1}(J_1(t)\phi)(x) + (\partial_t W_1(t)\phi)(x) \\ \quad + \frac{e^{-\frac{t}{2}}t^2}{8\pi} \int_{|y|\leq 1} \partial_t \left(I_1(t\sqrt{1-|y|^2}/2) \frac{\phi(x+ty)}{\sqrt{1-|y|^2}} \right) dy \quad (N=3). \end{cases} \tag{3.13}
 \end{aligned}$$

Next we begin with mentioning the L^q - L^r estimates for the evolution operators $K_0(t)$ and $K_1(t)$.

Lemma 3.3. [10, 27, 29, 30] *Let $1 \leq r \leq q \leq \infty$ and $0 < \delta < 1/2$. Assume that $(\phi, \psi) \in L^r \cap W^{[N/2],q} \times L^r \cap L^q$. Then it holds that*

$$\begin{aligned}
 \|K_0(t)\phi\|_{L^q} &\leq C(1+t)^{-\frac{N}{2}(\frac{1}{r}-\frac{1}{q})} \|\phi\|_{L^r} + Ce^{-\delta t} \|\phi\|_{W^{[N/2],q}}, \quad t \geq 0, \\
 \|K_1(t)\psi\|_{L^q} &\leq C(1+t)^{-\frac{N}{2}(\frac{1}{r}-\frac{1}{q})} \|\psi\|_{L^r} + Ce^{-\delta t} \|\psi\|_{L^q}, \quad t \geq 0. \tag{3.14}
 \end{aligned}$$

The L^1 weighted estimates for the operators $K_0(t)$ and $K_1(t)$ are well-known.

Lemma 3.4. [36, Proposition 5.2] *Let $k \geq 0$ and $0 < \delta < 1/8$. Then they hold that*

$$\|K_0(t)\phi\|_{L^1_k} \leq C\|\phi\|_{W^{[N/2],1}} + Ct^{\frac{k}{2}}\|\phi\|_{L^1}, \tag{3.15}$$

$$\|K_1(t)\phi\|_{L^1_k} \leq C\|\phi\|_{L^1_k} + Ct^{\frac{k}{2}}\|\phi\|_{L^1}, \tag{3.16}$$

for all $t \geq 0$.

Remark 3.1. The estimates (3.15) and (3.16) are already obtained in [21] by another method.

From now on, let us introduce the sequences of the heat semi-group, $V_0(t)$ and $V_1(t)$ as

$$V_0(t) := \frac{1}{2} \sum_{\ell=0}^{[K/2]} \sum_{m=0}^{[K/2]-\ell} \Phi_{\ell,m}(-t)^\ell (-\Delta)^{2\ell+m} e^{t\Delta}, \tag{3.17}$$

$$V_1(t) := \sum_{\ell=0}^{[K/2]} \sum_{m=0}^{[K/2]-\ell} \sum_{n=0}^{[K/2]-\ell-m} \Phi_{\ell,m} \Psi_n(-t)^\ell (-\Delta)^{2\ell+m+n} e^{t\Delta}, \tag{3.18}$$

for $K \geq 0$, where $(\Phi_{\ell,m}, \Psi_n)$ is given in (2.3). The following lemma states that the operators $K_0(t)$ and $K_1(t)$ are approximated by $V_0(t)$ and $V_1(t)$, respectively. This fact plays a crucial role in the proof of Theorem 2.2.

Lemma 3.5. [36, Proposition 4.1] *Let $K \geq 0$, $1 \leq r \leq q \leq \infty$ and $0 < \delta < 1/2$. Assume that $(\phi, \psi) \in L^r \cap W^{[N/2],q} \times L^r \cap L^q$. Then it holds that*

$$\begin{aligned} \|(K_0(t) - V_0(t))\phi\|_{L^q} &\leq Ct^{-\frac{N}{2}(\frac{1}{r}-\frac{1}{q})-(\lfloor\frac{K}{2}\rfloor+1)}\|\phi\|_{L^r} + Ce^{-\delta t}\|\phi\|_{W^{[N/2],q}}, \quad t > 0, \\ \|(K_1(t) - V_1(t))\phi\|_{L^q} &\leq Ct^{-\frac{N}{2}(\frac{1}{r}-\frac{1}{q})-(\lfloor\frac{K}{2}\rfloor+1)}\|\phi\|_{L^r} + Ce^{-\delta t}\|\phi\|_{L^q}, \quad t > 0, \end{aligned} \tag{3.19}$$

where $V_0(t)$ and $V_1(t)$ are the operators given in (3.17) and (3.18), respectively.

We note that using the notation $K_0(t)$ and $K_1(t)$, the solution of (1.5) is expressed as

$$v(x, t) = (K_0(t)v_0)(x) + (K_1(t)\left(\frac{1}{2}v_0 + v_1\right))(x) \tag{3.20}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. As a easy consequence of Lemma 3.5 with (3.20), we have the following:

Proposition 3.2. [36, Theorem 1.1] *Let $K \geq 0$ and $(v_0, v_1) \in W_K^{[N/2],1} \cap W^{[N/2],\infty} \times L^1_K \cap L^\infty$. Then there exists a unique solution v of (1.5) satisfying*

$$v(t) \in C([0, \infty) : L^1_K) \cap L^\infty(0, \infty : L^\infty).$$

Furthermore, for any $q \in [1, \infty]$,

$$t^{\frac{N}{2}(1-\frac{1}{q})} \left\| v(t) - V_0(t)v_0 - V_1(t)\left(\frac{1}{2}v_0 + v_1\right) \right\|_{L^q} = o(t^{-\frac{K}{2}})$$

as $t \rightarrow \infty$.

Combining Propositions 3.1 and 3.2, we can easily see the following.

Corollary 3.1. *Assume the same assumptions as in Proposition 3.2. Then, for any $q \in [1, \infty]$,*

$$t^{\frac{N}{2}(1-\frac{1}{q})} \|v(t) - U_{\text{lin}}(t)\|_{L^q} = o(t^{-\frac{K}{2}}) \tag{3.21}$$

as $t \rightarrow \infty$, where $U_{\text{lin}}(x, t)$ is defined by (2.7).

Remark 3.2. By this corollary we see that the function U_{lin} is an asymptotic expansion to the solution of (1.5), i.e. U_{lin} is a linear approximation of the linear part for the solution u for (1.1).

3.3. Useful formula

We prepare the point-wise estimate for the modified Bessel function.

Lemma 3.6. [1] *For any $\nu \in \mathbb{N}_0$, it holds that*

$$|I_\nu(x)| \leq \begin{cases} Cx^\nu, & 0 < x \leq 1, \\ Cx^{-\frac{1}{2}}e^x, & x \geq 1, \end{cases} \tag{3.22}$$

where I_ν is given in (3.8).

The following lemma is useful to obtain decay estimates.

Lemma 3.7. [36, Lemma 2.5] *Let $k \geq 0$. Then, for any $c > 0$ and $m = 1, 2, 3$, there exists a constant $C > 0$ such that*

$$\int_{|y| \leq 1} \frac{|y|^k e^{-ct|y|^2}}{(1 - |y|^2)^{\frac{m}{4}}} dy \leq Ct^{-\frac{N}{2} - \frac{k}{2}}, \quad t > 0. \tag{3.23}$$

4. Weighted estimates for the linear solutions

In this section we give weighted L^1 and L^∞ estimates for the solution v of (2.12). In other words, our new ingredient for the proof of the main result is the following weighted L^∞ estimates for the evolution operators $K_0(t)$ and $K_1(t)$, which are useful for us to estimate the nonlinear term in the nonlinear approximation.

Our purpose here is to obtain the following proposition on the weighted estimates for the solution v of the linear damped wave equation (1.5).

Proposition 4.1. *Let $K \geq 0$. Assume $(v_0, v_1) \in W_K^{[N/2], 1} \cap W^{[N/2]_{K}, \infty} \times L_K^1 \cap L_K^\infty$. Let v be a unique solution of (1.5). Then, for any $k \in [0, K]$,*

$$\|v(t)\|_{L_k^1} \leq CE_K(1 + t)^{\frac{k}{2}}, \quad t \geq 0, \tag{4.1}$$

$$\|v(t)\|_{L_k^\infty} \leq CE_K(1 + t)^{\frac{k-N}{2}}, \quad t \geq 0, \tag{4.2}$$

where E_K is given in (2.15) with (u_0, u_1) replaced by (v_0, v_1) , respectively.

To this end, we show the following lemmas, which are the weighted estimates for the evolution operators $K_0(t)$ and $K_1(t)$.

Lemma 4.1. *Assume the same conditions as in Lemma 3.4. Then it holds that*

$$\|K_0(t)\phi\|_{L_k^\infty} \leq C\|\phi\|_{W_k^{[N/2], \infty}} + Ct^{\frac{k}{2}}\|\phi\|_{L^\infty}, \tag{4.3}$$

$$\|K_1(t)\phi\|_{L_k^\infty} \leq C\|\phi\|_{L_k^\infty} + Ct^{\frac{k}{2}}\|\phi\|_{L^\infty}, \tag{4.4}$$

for all $t \geq 0$. Furthermore it holds that

$$\|K_0(t)\phi\|_{L_k^\infty} \leq Ct^{-\frac{N}{2}}\|\phi\|_{W_k^{[N/2], 1}} + Ct^{\frac{k-N}{2}}\|\phi\|_{L^1} + Ce^{-\delta t}\|\phi\|_{L_k^\infty}, \tag{4.5}$$

$$\|K_1(t)\phi\|_{L_k^\infty} \leq Ct^{-\frac{N}{2}}\|\phi\|_{L_k^1} + Ct^{\frac{k-N}{2}}\|\phi\|_{L^1} + Ce^{-\delta t}\|\phi\|_{L_k^\infty}, \tag{4.6}$$

for all $t \geq 1$.

By Lemma 4.1, for the weighted L^∞ estimates, we can remove the singularity at $t = 0$, and we obtain the following.

Corollary 4.1. *Assume the same assumptions as in Lemma 3.4. Then it holds that*

$$\begin{aligned} \|K_0(t)\phi\|_{L_k^\infty} &\leq C(1+t)^{-\frac{N}{2}} \left(\|\phi\|_{W_k^{[N/2],1}} + \|\phi\|_{W_k^{[N/2],\infty}} \right) \\ &\quad + C(1+t)^{\frac{k-N}{2}} \left(\|\phi\|_{W^{[N/2],1}} + \|\phi\|_{W^{[N/2],\infty}} \right), \end{aligned} \tag{4.7}$$

$$\begin{aligned} \|K_1(t)\phi\|_{L_k^\infty} &\leq C(1+t)^{-\frac{N}{2}} \left(\|\phi\|_{L_k^1} + \|\phi\|_{L_k^\infty} \right) \\ &\quad + C(1+t)^{\frac{k-N}{2}} \left(\|\phi\|_{L^1} + \|\phi\|_{L^\infty} \right), \end{aligned} \tag{4.8}$$

for all $t \geq 0$.

By Lemma 3.4 and Corollary 4.1 with the aid of (3.20) we can prove Proposition 4.1 immediately. So it suffices to prove Lemma 4.1.

Here, for the simplicity of the notation, we introduce the auxiliary functional $\mathcal{K}^{(m)}[\phi]$ as follows:

$$\mathcal{K}^{(m)}[\phi](x, t) := t^{\frac{N}{2}} \int_{|y| \leq 1} \frac{e^{-\frac{t}{2}|y|^2}}{(1-|y|^2)^{\frac{m}{4}}} |\phi(x+ty)| dy, \quad (x, t) \in \mathbb{R}^N \times (0, \infty), \tag{4.9}$$

where $m = 0, 1, 2, 3$. Then we have the weighted L^∞ estimates of $\mathcal{K}^{(m)}[\phi](x, t)$.

Lemma 4.2. *Assume the same conditions as in Lemma 3.4. Then, for any $m = 0, 1, 2, 3$,*

$$\|\mathcal{K}^{(m)}[\phi](t)\|_{L_k^\infty} \leq C\|\phi\|_{L_k^\infty} + Ct^{\frac{k}{2}}\|\phi\|_{L^\infty}, \quad t \geq 0, \tag{4.10}$$

and

$$\|\mathcal{K}^{(m)}[\phi](t)\|_{L_k^\infty} \leq Ct^{-\frac{N}{2}}\|\phi\|_{L_k^1} + Ct^{\frac{k-N}{2}}\|\phi\|_{L^1} + Ce^{-\delta t}\|\phi\|_{L_k^\infty}, \quad t > 0. \tag{4.11}$$

Proof. By (4.9) we see that

$$\begin{aligned} |x|^k |\mathcal{K}^{(m)}[\phi](x, t)| &\leq Ct^{\frac{N}{2}} \int_{|y| \leq 1} (|x+ty|^k + |ty|^k) \frac{e^{-\frac{t}{2}|y|^2}}{(1-|y|^2)^{\frac{m}{4}}} |\phi(x+ty)| dy \\ &=: CI_1(x, t) + CI_2(x, t) \end{aligned} \tag{4.12}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$.

We first show the estimate (4.10). Applying the Minkowski inequality with (3.23) we have

$$\|I_1(t)\|_{L^\infty} \leq Ct^{\frac{N}{2}} \int_{|y| \leq 1} \frac{e^{-\frac{t}{2}|y|^2}}{(1-|y|^2)^{\frac{m}{4}}} \|\cdot + ty\|^k \phi(\cdot + ty)\|_{L^\infty} dy \leq C\|\phi\|_{L_k^\infty}, \tag{4.13}$$

$$\|I_2(t)\|_{L^\infty} \leq Ct^{\frac{N}{2}+k} \int_{|y| \leq 1} \frac{e^{-\frac{t}{2}|y|^2} |y|^k}{(1-|y|^2)^{\frac{m}{4}}} \|\phi(\cdot + ty)\|_{L^\infty} dy \leq Ct^{\frac{k}{2}}\|\phi\|_{L^\infty}. \tag{4.14}$$

Then combining the above estimates (4.12)-(4.14), we obtain the estimate (4.10).

Next we prove the estimate (4.11). Observing the estimate (4.12), we estimate I_1 and I_2 , respectively. For the term I_1 , we decompose into the two parts:

$$\begin{aligned} I_1(x, t) &\leq Ct^{\frac{N}{2}} \int_{|y| \leq 1/2} |x + ty|^k e^{-\frac{t}{2}|y|^2} |\phi(x + ty)| dy \\ &\quad + Ct^{\frac{N}{2}} e^{-\frac{t}{8}} \int_{1/2 \leq |y| \leq 1} (1 - |y|^2)^{-\frac{m}{4}} |x + ty|^k |\phi(x + ty)| dy \\ &=: I_{11}(x, t) + I_{12}(x, t) \end{aligned}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. By changing the integral variable $z = x + ty$ we see that

$$I_{11}(x, t) = Ct^{-\frac{N}{2}} \int_{|z-x| \leq t/2} |z|^k e^{-\frac{|z-x|^2}{2t}} |\phi(z)| dz \leq Ct^{-\frac{N}{2}} \|\phi\|_{L^k_k} \quad (4.15)$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. For $I_{12}(x, t)$, since

$$\int_{1/2 \leq |y| \leq 1} (1 - |y|^2)^{-\frac{m}{4}} dy \leq C \quad (4.16)$$

for any $m = 0, 1, 2, 3$, we have

$$\begin{aligned} I_{12}(x, t) &\leq Ct^{\frac{N}{2}} e^{-\frac{t}{8}} \|\phi\|_{L^\infty_k} \int_{1/2 \leq |y| \leq 1} (1 - |y|^2)^{-\frac{m}{4}} dy \\ &\leq Ct^{\frac{N}{2}} e^{-\frac{t}{8}} \|\phi\|_{L^\infty_k} \end{aligned} \quad (4.17)$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. $I_2(x, t)$ is estimated by the similar way. Indeed, we again decompose $I_2(x, t)$ into two parts:

$$\begin{aligned} I_2(x, t) &\leq Ct^{\frac{N}{2}} \int_{|y| \leq 1/2} |ty|^k e^{-\frac{t}{2}|y|^2} |\phi(x + ty)| dy \\ &\quad + Ct^{\frac{N}{2}} e^{-\frac{t}{8}} \int_{1/2 \leq |y| \leq 1} (1 - |y|^2)^{-\frac{m}{4}} |ty|^k |\phi(x + ty)| dy \\ &=: I_{21}(x, t) + I_{22}(x, t) \end{aligned}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. By the well-known estimate

$$|y|^k e^{-\frac{t}{2}|y|^2} \leq Ct^{-\frac{k}{2}}$$

and changing the integral variable $z = x + ty$, we arrive at the estimate

$$\begin{aligned} I_{21}(x, t) &= Ct^{\frac{N}{2}+k} \int_{|y| \leq 1/2} |y|^k e^{-\frac{t}{2}|y|^2} |\phi(x + ty)| dy \\ &\leq Ct^{\frac{N+k}{2}} \int_{|y| \leq 1/2} |\phi(x + ty)| dy \\ &\leq Ct^{\frac{k-N}{2}} \int_{\mathbb{R}^N} |\phi(z)| dz \leq Ct^{\frac{k-N}{2}} \|\phi\|_{L^1} \end{aligned} \quad (4.18)$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. $I_{22}(x, t)$ is estimated easily by the estimate (4.16) and $|y| \leq 1$:

$$I_{22}(x, t) \leq Ct^{\frac{N}{2}} e^{-\frac{t}{8}} \|\phi\|_{L^\infty} \int_{1/2 \leq |y| \leq 1} (1 - |y|^2)^{-\frac{m}{4}} dy \leq Ct^{\frac{N}{2}} e^{-\frac{t}{8}} \|\phi\|_{L^\infty} \tag{4.19}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. Then, summing up the estimates (4.12), (4.15), (4.17), (4.18) and (4.19), we have

$$|x|^k |\mathcal{K}^{(m)}[\phi](x, t)| \leq Ct^{-\frac{N}{2}} \|\phi\|_{L^1_k} + Ct^{\frac{k-N}{2}} \|\phi\|_{L^1} + Ct^{\frac{N}{2}} e^{-\frac{t}{8}} \|\phi\|_{L^\infty_k},$$

which implies the estimate (4.11). Thus the proof of Lemma 4.2 is complete. □

As a next step, we show the point-wise estimates for the evolution operators $K_1(t)$ and $K_0(t)$ by the auxiliary functional $\mathcal{K}^{(m)}[\phi](x, t)$.

Lemma 4.3. *Let $\mathcal{K}^m[\phi]$ be functions given in (4.9). Then, for the case $N = 1, 2$,*

$$|(K_1(t)\phi)(x)| \leq C\mathcal{K}^{(N)}[\phi](x, t), \quad (x, t) \in \mathbb{R}^N \times (0, \infty), \tag{4.20}$$

and, for the case $N = 3$,

$$|(J_1(t)\phi)(x)| \leq C\mathcal{K}^{(N)}[\phi](x, t), \quad (x, t) \in \mathbb{R}^N \times (0, \infty). \tag{4.21}$$

Proof. The proof is an easy consequence of Lemma 3.6. By (3.22) we see that

$$|I_\nu(x)| \leq Cx^{-\frac{1}{2}}e^x \tag{4.22}$$

for $\nu \in \mathbb{N}_0$ and $x > 0$. Furthermore, for any $y \in \{y \in \mathbb{R}^N : |y| < 1\}$, we have

$$-\frac{1}{2} + \frac{\sqrt{1 - |y|^2}}{2} = \frac{-|y|^2}{2(1 + \sqrt{1 - |y|^2})} \leq -\frac{|y|^2}{2}. \tag{4.23}$$

We first prove (4.20). For $N = 1$, applying the estimate (4.22) to the expression (3.9) and (4.23), we have

$$\begin{aligned} |(K_1(t)\phi)(x)| &\leq Ce^{-\frac{t}{2}} \int_{|y| \leq 1} \frac{e^{\frac{t\sqrt{1-y^2}}{2}}}{t^{\frac{1}{2}}(1-y^2)^{\frac{1}{4}}} |\phi(x+ty)| dy \\ &\leq Ct^{\frac{1}{2}} \int_{|y| \leq 1} \frac{e^{-\frac{t}{2}|y|^2}}{(1-y^2)^{\frac{1}{4}}} |\phi(x+ty)| dy \end{aligned}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. Therefore we obtain the desired estimate (4.20) for $N = 1$. Furthermore, by the definition of $\cosh y$, we see that

$$0 \leq \cosh y \leq e^y, \quad y \in \mathbb{R}. \tag{4.24}$$

Thus, again using the estimate (4.23), we can easily have the desired estimate (4.20) for $N = 2$. Next we prove the estimate (4.21). Applying the estimates (4.22) and (4.23) to the expression (3.11) we have

$$\begin{aligned} |(J_1(t)\phi)(x)| &\leq C e^{-\frac{t}{2}} t^2 \int_{|y|\leq 1} \frac{e^{\frac{t\sqrt{1-y^2}}{2}}}{t^{\frac{1}{2}}(1-y^2)^{\frac{3}{4}}} |\phi(x+ty)| dy \\ &\leq C t^{\frac{3}{2}} \int_{|y|\leq 1} \frac{e^{-\frac{1}{2}t|y|^2}}{(1-y^2)^{\frac{3}{4}}} |\phi(x+ty)| dy \end{aligned}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$, which is the desired estimate. Then Lemma 4.3 follows. □

Lemma 4.4. *Assume the same assumption as in Lemma 4.3. Then the following holds:*

(i) For $N = 1$,

$$|(K_0(t)\phi)(x)| \leq C(|(K_1(t)\phi)(x)| + C e^{-\frac{t}{2}}(|\phi(x-t)| + |\phi(x+t)|) + CK^{(3)}[\phi](x, t)) \tag{4.25}$$

for all $(x, t) \in \mathbb{R} \times (0, \infty)$;

(ii) For $N = 2$,

$$|(K_0(t)\phi)(x)| \leq C t^{-1} |(K_1(t)\phi)(x)| + CK^{(0)}[\phi](x, t) + CK^{(2)}[|\nabla\phi|](x, t) \tag{4.26}$$

for all $(x, t) \in \mathbb{R}^2 \times (0, \infty)$;

(iii) For $N = 3$,

$$\begin{aligned} |(K_0(t)\phi)(x)| &\leq C(t^{-1}|(J_1(t)\phi)(x)| + \mathcal{K}^{(1)}(x, t) + |(K_1(t)|\nabla\phi|)(x)| \\ &\quad + |(W_1(t)\phi)(x)| + |(W_1(t)|\nabla\phi|)(x)| \end{aligned} \tag{4.27}$$

for all $(x, t) \in \mathbb{R}^3 \times (0, \infty)$.

Proof. We first prove for the case $N = 1$. Observing the estimates (4.10), we use the estimates (3.22) and (4.23) to see that

$$\begin{aligned} |(K_0(t)\phi)(x)| &\leq C(|(K_1(t)\phi)(x)| + e^{-\frac{t}{2}}(|\phi(x-t)| + |\phi(x+t)|)) \\ &\quad + C e^{-\frac{t}{2}} t^{\frac{1}{2}} \int_{|y|\leq 1} \frac{e^{\frac{t\sqrt{1-y^2}}{2}}}{(1-y^2)^{\frac{3}{4}}} |\phi(x+ty)| dy \end{aligned} \tag{4.28}$$

for all $(x, t) \in \mathbb{R} \times (0, \infty)$. In addition, by (4.22) and (4.23), it is easy to see that

$$C e^{-\frac{t}{2}} t^{\frac{1}{2}} \int_{|y|\leq 1} \frac{e^{\frac{t\sqrt{1-y^2}}{2}}}{(1-y^2)^{\frac{3}{4}}} |\phi(x+ty)| dy \leq CK^{(3)}(x, t)$$

for all $(x, t) \in \mathbb{R} \times (0, \infty)$. This together with (4.28) yields the desired estimate (4.25).

Next we show the point-wise estimate for $K_0(t)g$ with $N = 2$. We recall the estimate (4.24) and

$$0 \leq \sinh y \leq e^y$$

for $y \geq 0$. Then we use (3.13) to have

$$\begin{aligned} |(K_0(t)\phi)(x)| &\leq t^{-1}|(K_1(t)\phi)(x)| + Ce^{-\frac{t}{2}} \int_{|y|\leq 1} e^{t\frac{\sqrt{1-|y|^2}}{2}} |\phi(x+ty)| dy \\ &\quad + Ce^{-\frac{t}{2}} \int_{|y|\leq 1} \frac{e^{t\frac{\sqrt{1-|y|^2}}{2}}}{\sqrt{1-|y|^2}} |\nabla\phi(x+ty)||y| dy \\ &\leq t^{-1}|(K_1(t)\phi)(x)| + C\mathcal{K}^{(0)}[\phi](x,t) + C\mathcal{K}^{(2)}[|\nabla\phi|](x,t), \end{aligned}$$

which is the desired estimate (4.26).

Finally we show the case $N = 3$. To this end, we use the point-wise estimate, which is given in [36, (5.13)], as follows:

$$\begin{aligned} |(K_0(t)\phi)(x)| &\leq Ct^{-1}|(J_1(t)\phi)(x)| + Ct^{\frac{3}{2}} \int_{|y|\leq 1} \frac{e^{\frac{-t|y|^2}{2(1+\sqrt{1-|y|^2})}}}{(\sqrt{1-|y|^2})^{\frac{1}{2}}} |\phi(x+ty)| dy \\ &\quad + |(K_1(t)|\nabla\phi)(x)| + C|(W_1(t)\phi)(x)| + C|(W_1(t)|\nabla\phi)(x)| \end{aligned}$$

Then by the definition of $\mathcal{K}^{(1)}$, we obtain the estimate (4.27). Thus the proof of Lemma 4.4 is complete. \square

Finally we show Lemma 4.1.

Proof of Lemma 4.1. We first prove the case $N = 1$. By Lemma 4.2 with (4.20) we have (4.4). Since

$$\begin{aligned} &|x|^k(|\phi(x-t)| + |\phi(x+t)|) \\ &\leq C(|x-t|^k|\phi(x-t)| + t^k|\phi(x-t)| + |x+t|^k|\phi(x+t)| + t^k|\phi(x+t)|), \end{aligned}$$

we see that

$$\|(|\phi(\cdot-t)| + |\phi(\cdot+t)|)\|_{L_k^\infty} \leq C(\|\phi\|_{L_k^\infty} + t^k\|\phi\|_{L^\infty}), \quad t \geq 0. \tag{4.29}$$

This together with (4.4), (4.10) and (4.25) yields the estimate (4.3). Furthermore, combining the estimates (4.11), (4.20), (4.25) and (4.29), we have the estimates (4.5) and (4.6).

Next we show the case $N = 2$. Similarly to the case $N = 1$, by Lemma 4.2 with (4.20) we have (4.4). By (4.23) and (4.24) we observe that

$$\begin{aligned} |(K_1(t)\phi)(x)| &\leq Cte^{-\frac{t}{2}} \int_{|y|\leq 1} \frac{\cosh\left(t\frac{\sqrt{1-|y|^2}}{2}\right)}{\sqrt{1-|y|^2}} |\phi(x+ty)| dy \\ &\leq Ct \int_{|y|\leq 1} \frac{|\phi(x+ty)|}{\sqrt{1-|y|^2}} dy \end{aligned} \tag{4.30}$$

for all $(x,t) \in \mathbb{R}^2 \times (0,\infty)$. Then, by (3.23) and applying the Minkowski inequality to (4.30), we see that

$$\begin{aligned} |x|^k|(K_1(t)\phi)(x)| &\leq Ct \int_{|y|\leq 1} \frac{|x+ty|^k|\phi(x+ty)| + |ty|^k|\phi(x+ty)|}{\sqrt{1-|y|^2}} dy \\ &\leq Ct\|\phi\|_{L_k^\infty} + Ct^{\frac{k}{2}}\|\phi\|_{L^\infty} \end{aligned}$$

for all $(x, t) \in \mathbb{R}^2 \times (0, \infty)$. This together with (4.10) and (4.26) implies (4.3). Furthermore, combining the estimates (4.11), (4.20) and (4.26), we have the estimates (4.5) and (4.6).

Finally we prove the case $N = 3$. For the term $W_1(t)g$, by (3.12) we have

$$\begin{aligned} |x|^k |(W_1(t)\phi)(x)| &\leq C e^{-\frac{t}{2}} \int_{\mathbb{S}^2} (|x + t\omega|^k + t^k) |\phi(x + t\omega)| d\omega \\ &\leq C e^{-\frac{t}{2}} t (\|\phi\|_{L_k^\infty} + t^k \|\phi\|_{L^\infty}) \end{aligned} \tag{4.31}$$

for all $(x, t) \in \mathbb{R}^3 \times (0, \infty)$. Then, noting (3.10), (4.10), (4.21) and (4.31), we arrive at the estimate

$$\begin{aligned} \|K_1(t)\phi\|_{L_k^\infty} &\leq C \|\mathcal{K}^{(3)}[\phi](t)\|_{L_k^\infty} + C e^{-\frac{t}{2}} t (\|\phi\|_{L_k^\infty} + t^k \|\phi\|_{L^\infty}) \\ &\leq C \|\phi\|_{L_k^\infty} + C t^{\frac{k}{2}} \|\phi\|_{L^\infty}, \quad t \geq 0, \end{aligned}$$

which is the estimate (4.4). To show the estimate (4.3), we recall that the estimate (3.22) implies

$$|I_\nu(x)| \leq C e^x, \quad x > 0,$$

for $\nu \in \mathbb{N}_0$. Then, by (3.23) we obtain

$$\begin{aligned} |x|^k |(J_1(t)\phi)(x)| &\leq C e^{-\frac{t}{2}} \int_{|y| \leq 1} e^{\frac{t\sqrt{1-|y|^2}}{2}} \frac{|x|^k |\phi(x + ty)|}{\sqrt{1-|y|^2}} dy \\ &\leq C t \int_{|y| \leq 1} \frac{(|x + ty|^k + |ty|^k) |\phi(x + ty)|}{\sqrt{1-|y|^2}} dy \\ &\leq C t (\|\phi\|_{L_k^\infty} + t^{\frac{k}{2}} \|\phi\|_{L^\infty}) \end{aligned} \tag{4.32}$$

for all $(x, t) \in \mathbb{R}^3 \times (0, \infty)$. Summing up the estimates (4.10), (4.20), (4.27), (4.31) and (4.32), we obtain the estimate (4.3). By (3.10), (4.11), (4.21) and (4.31) we reached at the estimate (4.6). Furthermore, by (4.6), (4.11), (4.20) and (4.31) we can easily obtain the estimate (4.5). Thus Lemma 4.1 follows. \square

5. Proof of Theorem 2.1

This section is devoted to the proof of Theorem 2.1. We note that the our strategy for the proof of the estimates (2.13) and (2.14) are the combination of the argument similar to the proof of [21, Lemma 4.1] and the weighted L^∞ estimates for the linearized solution which is developed in Section 4.

First we give the following weighted L^1 estimate for the solution to (1.1).

Lemma 5.1. *Assume the same conditions as in Theorem 2.1. Then there exists a positive constant ε_0 such that, if $E_0 \leq \varepsilon_0$, then the estimate (2.13) holds.*

Proof. For any $t \geq 0$, we put

$$E(t) := \sup_{0 \leq s \leq t} (1 + s)^{-\frac{k}{2}} \|u(s)\|_{L_k^1}.$$

Then, by (2.4), (3.16) and (4.1) we obtain

$$\|u(t)\|_{L_k^1} \leq CE_K(1+t)^{\frac{k}{2}} + C \int_0^t (\|F(u(s))\|_{L_k^1} + (t-s)^{\frac{k}{2}} \|F(u(s))\|_{L^1}) ds \quad (5.1)$$

for all $t \geq 0$. Here, by (1.2) and (1.3) we have

$$\|F(u(t))\|_{L^1} \leq C\|u(t)\|_{L^p}^p \leq CE_0^p(1+t)^{-A}, \quad (5.2)$$

$$\|F(u(t))\|_{L_k^1} \leq C\|u(t)\|_{L^\infty}^{p-1} \|u(t)\|_{L_k^1} \leq CE_0^{p-1}(1+t)^{-A} \|u(t)\|_{L_k^1}, \quad (5.3)$$

for all $t \geq 0$, where A is defined by (2.16). Since $A > 1$, substituting the above estimates into (5.1), this yields

$$\begin{aligned} \|u(t)\|_{L_k^1} &\leq CE_K(1+t)^{\frac{k}{2}} + CE_0^{p-1}E(t) \int_0^t (1+s)^{\frac{k}{2}-A} ds \\ &\quad + CE_0^p \int_0^t (t-s)^{\frac{k}{2}} (1+s)^{-A} ds \\ &\leq C(E_K + E_0^p + E_0^{p-1}E(t))(1+t)^{\frac{k}{2}}, \quad t \geq 0. \end{aligned}$$

Eventually, we obtain

$$E(t) \leq C_0(E_K + E_0^p + E_0^{p-1}E(t)), \quad t \geq 0,$$

where C_0 is a positive constant. Choosing ε_0 such that

$$0 < \varepsilon_0^{p-1} \leq \min\{1/(2C_0), E_K^{(p-1)/p}\},$$

and letting E_0 sufficiently small such that $E_0 \leq \varepsilon_0$, we obtain $E(t) \leq 4C_0E_K$ and the desired estimate (2.13). Thus Lemma 5.1 follows. \square

Secondly, we give the following weighted L^∞ estimate for the solution u of (1.1).

Lemma 5.2. *Assume the same conditions as in Theorem 2.1. Then there exists a positive constant ε_1 such that, if $E_0 \leq \varepsilon_1$, then the estimate (2.14) holds.*

Proof. For any $t \geq 0$, we put

$$\tilde{E}(t) := \sup_{0 \leq s \leq t} (1+s)^{\frac{N-k}{2}} \|u(s)\|_{L_k^\infty}.$$

Then, by (2.4) and (4.2) we obtain

$$\|u(t)\|_{L_k^\infty} \leq CE_K(1+t)^{\frac{k-N}{2}} + \left(\int_0^{t/2} + \int_{t/2}^t \right) \|K_1(t-s)F(u(s))\|_{L_k^\infty} ds \quad (5.4)$$

for all $t \geq 0$. Here, by (1.2) and (1.3) we have

$$\|F(u(t))\|_{L^\infty} \leq C\|u(t)\|_{L^\infty}^p \leq CE_0^p(1+t)^{-A-1-\frac{N}{2}}, \quad (5.5)$$

$$\|F(u(t))\|_{L_k^\infty} \leq C\|u(t)\|_{L^\infty}^{p-1} \|u(t)\|_{L_k^\infty} \leq CE_0^{p-1}(1+t)^{-A-1} \|u(t)\|_{L_k^\infty}, \quad (5.6)$$

for all $t \geq 0$, where A is defined by (2.16). Then, since $A > 0$, by (4.4), (5.5) and (5.6) we obtain

$$\begin{aligned} & \int_{t/2}^t \|K_1(t-s)F(u(s))\|_{L_k^\infty} ds \\ & \leq C \int_{t/2}^t (\|F(u(s))\|_{L_k^\infty} + (t-s)^{\frac{k}{2}} \|F(u(s))\|_{L^\infty}) ds \\ & \leq CE_0^{p-1} \tilde{E}(t) \int_{t/2}^t (1+s)^{-A-1-\frac{N-k}{2}} ds + CE_0^p \int_{t/2}^t (t-s)^{\frac{k}{2}} (1+s)^{-A-1-\frac{N}{2}} ds \\ & \leq CE_0^{p-1} (\tilde{E}(t) + E_0) (1+t)^{-\frac{N-k}{2}} \end{aligned} \tag{5.7}$$

for all $t \geq 0$. Furthermore, since $A > 0$ again, by (2.13), (4.8), (5.2), (5.3), (5.5) and (5.6) we have

$$\begin{aligned} & \int_0^{t/2} \|K_1(t-s)F(u(s))\|_{L_k^\infty} ds \\ & \leq C \int_0^{t/2} (1+t-s)^{-\frac{N}{2}} (\|F(u(s))\|_{L_k^1} + \|F(u(s))\|_{L_k^\infty}) ds \\ & \quad + C \int_0^{t/2} (1+t-s)^{-\frac{N-k}{2}} (\|F(u(s))\|_{L^1} + \|F(u(s))\|_{L^\infty}) ds \\ & \leq CE_0^{p-1} (1+t)^{-\frac{N}{2}} \int_0^{t/2} (1+s)^{-A-1} (\|u(s)\|_{L_k^1} + \|u(s)\|_{L_k^\infty}) ds \\ & \quad + CE_0^p (1+t)^{-\frac{N-k}{2}} \int_0^{t/2} (1+s)^{-A-1} (1+(1+s)^{-\frac{N}{2}}) ds \\ & \leq CE_0^{p-1} E_K (1+t)^{-\frac{N}{2}} \int_0^{t/2} (1+s)^{-A-1+\frac{k}{2}} ds \\ & \quad + CE_0^{p-1} \tilde{E}(t) (1+t)^{-\frac{N}{2}} \int_0^{t/2} (1+s)^{-A-1-\frac{N-k}{2}} ds + CE_0^p (1+t)^{-\frac{N-k}{2}} \\ & \leq CE_0^{p-1} (E_K + \tilde{E}(t) + E_0) (1+t)^{-\frac{N-k}{2}} \end{aligned}$$

for all $t \geq 0$. This together with (5.4) and (5.7) implies that

$$\tilde{E}(t) \leq C_1 (E_K + E_0^p + E_0^{p-1} E_K + E_0^{p-1} \tilde{E}(t)), \quad t \geq 0.$$

where C_1 is a positive constant. Choosing ε_1 such that

$$0 < \varepsilon_1^{p-1} \leq \min\{1, 1/(2C_1), E_K^{p-1}\}$$

and letting E_0 sufficiently small such that $E_0 \leq \varepsilon_1$, we obtain $\tilde{E}(t) \leq 6C_1 E_K$ and the desired estimate (2.14), and the proof of Lemma 5.1 is complete. \square

Proof of Theorem 2.1 We note that under the assumptions for the nonlinearity (1.2), (1.3) and (2.11) and the smallness of the initial data with regularity, we easily have the existence of global solution to (1.1) satisfying (2.12) (see e.g. [10], [27], [30]). Therefore, to show Theorem 2.1, it suffices to prove the estimates (2.13) and (2.14). Moreover, by Lemmas 5.1 and 5.2, we obtain

the estimates (2.13) and (2.14) with putting $\varepsilon := \min\{\varepsilon_0, \varepsilon_1\}$. The proof of Theorem 2.1 is completed.

6. Proof of Theorem 2.2

In this section, applying the argument as in [18], we prove Theorem 2.2. The proof consists of two parts, the derivation of the estimates (2.18) and (2.19).

Proof of (2.18). The proof is by induction. Applying estimates (3.2) and (3.3) to the definition of U_{lin} , that is (2.7), we can easily obtain

$$\sup_{t>0} (1+t)^{\frac{N}{2}(1-\frac{1}{q})} \|U_{\text{lin}}(t)\|_{L^q} + \sup_{t>0} (1+t)^{-\frac{k}{2}+\frac{N}{2}(1-\frac{1}{\gamma})} \|U_{\text{lin}}(t)\|_{L_k^\gamma} < \infty \quad (6.1)$$

for any $q \in [1, \infty]$, $k \in [0, K]$ and $\gamma \in \{1, \infty\}$. On the other hand, by (1.2), (2.12), (2.13), (2.14) and (3.5) we have

$$E_{K,q}[F(u)](t) \leq C \|u(t)\|_{L^\infty}^{p-1} E_{K,q}[u](t) \leq C(1+t)^{\frac{K}{2}-A-1}, \quad t > 0, \quad (6.2)$$

where A is given in (2.16). By (3.6) and (6.2) we obtain

$$\begin{aligned} & \left\| \int_0^{t/2} K_1(t-s) M_\alpha(F(u(s)), s) g_\alpha(s) ds \right\|_{L_k^\gamma} \\ & \leq \int_0^{t/2} (\|M_\alpha F(u(s)), s\| \|K_1(t-s) g_\alpha(s)\|_{L_k^\gamma}) ds \\ & \leq C \int_0^{t/2} (1+s)^{-\frac{K-|\alpha|}{2}} E_{K,q}[F(u)](s) \|K_1(t-s) g_\alpha(s)\|_{L_k^\gamma} ds \\ & \leq C \int_0^{t/2} (1+s)^{-A-1+\frac{|\alpha|}{2}} \|K_1(t-s) g_\alpha(s)\|_{L_k^\gamma} ds, \quad t > 0. \end{aligned} \quad (6.3)$$

Here, by (3.2), (3.3), (3.16) and (4.8) we have

$$\begin{aligned} \|K_1(t-s) g_\alpha(s)\|_{L_k^1} & \leq C \|g_\alpha(s)\|_{L_k^1} + C(t-s)^{\frac{k}{2}} \|g_\alpha(s)\|_{L^1} \\ & \leq C(1+t)^{\frac{k-|\alpha|}{2}} + C(t-s)^{\frac{k}{2}} (1+t)^{-\frac{|\alpha|}{2}}, \quad t > 0, \end{aligned} \quad (6.4)$$

and

$$\begin{aligned} \|K_1(t-s) g_\alpha(s)\|_{L_k^\infty} & \leq C(1+t-s)^{-\frac{N}{2}} (\|g_\alpha(s)\|_{L_k^1} + \|g_\alpha(s)\|_{L_k^\infty}) \\ & \quad + C(1+t-s)^{\frac{k-N}{2}} (\|g_\alpha(s)\|_{L^1} + \|g_\alpha(s)\|_{L^\infty}) \\ & \leq C(1+t-s)^{-\frac{N}{2}} (1+s)^{\frac{k-|\alpha|}{2}} \\ & \quad + C(1+t-s)^{\frac{k-N}{2}} (1+s)^{-\frac{|\alpha|}{2}}, \quad t > 0. \end{aligned}$$

These together with (6.3) and $A > 0$ yield

$$\begin{aligned} & \left\| \int_0^{t/2} K_1(t-s)M_\alpha(F(u(s)),s)g_\alpha(s) ds \right\|_{L_k^\gamma} \\ & \leq C \int_0^{t/2} \left((1+t-s)^{-\frac{N}{2}(1-\frac{1}{\gamma})}(1+s)^{\frac{k}{2}-A-1} \right. \\ & \quad \left. + (1+t-s)^{\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})}(1+s)^{-A-1} \right) ds \\ & \leq C(1+t)^{\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})} \int_0^\infty (1+s)^{-A-1} ds \leq C(1+t)^{\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})}, \quad t > 0. \end{aligned} \tag{6.5}$$

Furthermore, since it follows from (3.2), (3.3) and (4.4) that

$$\begin{aligned} \|K_1(t-s)g_\alpha(s)\|_{L_k^\infty} & \leq C\|g_\alpha(s)\|_{L_k^\infty} + C(t-s)^{\frac{k}{2}}\|g_\alpha(s)\|_{L^\infty} \\ & \leq C(1+s)^{\frac{k-N-|\alpha|}{2}} + C(t-s)^{\frac{k}{2}}(1+s)^{-\frac{N+|\alpha|}{2}}, \quad t > 0, \end{aligned} \tag{6.6}$$

by (6.3), (6.4), (6.6) and $A > 0$ we have

$$\begin{aligned} & \left\| \int_{t/2}^t K_1(t-s)M_\alpha(F(u(s)),s)g_\alpha(s) ds \right\|_{L_k^\gamma} \\ & \leq C \int_{t/2}^t (1+s)^{-A+\frac{|\alpha|}{2}} \|K_1(t-s)g_\alpha(s)\|_{L_k^\gamma} ds \\ & \leq C \int_{t/2}^t \left((1+s)^{-A-1+\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})} + (t-s)^{\frac{k}{2}}(1+s)^{-A-1-\frac{N}{2}(1-\frac{1}{\gamma})} \right) ds \\ & \leq C(1+t)^{\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})} \int_0^\infty ((1+s)^{-A-1} ds \leq C(1+t)^{\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})} \quad t > 0. \end{aligned}$$

This together with (2.10), (6.1) and (6.5) yields (2.18) for $j = 0$.

Assume that (2.18) holds for some $j = m \in \mathbb{N}_0$. Similarly to (6.2) with (2.18) for the case $j = m$, we have

$$E_{K,q}[F_m](t) \leq C\|U_m(t)\|_{L^\infty}^{p-1} E_{K,q}[U_m](t) \leq C(1+t)^{\frac{K}{2}-A-1}, \quad t > 0. \tag{6.7}$$

By (2.10) and (2.18) with $j = 0$ we see that

$$\begin{aligned} \|U_{m+1}(t)\|_{L_k^\gamma} & \leq C(1+t)^{\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})} + \left\| \left(\int_0^{t/2} + \int_{t/2}^t \right) K_1(t-s)P_K(s)F_m(s) ds \right\|_{L_k^\gamma} \\ & =: C(1+t)^{\frac{k}{2}-\frac{N}{2}(1-\frac{1}{\gamma})} + J_1(t) + J_2(t), \quad t > 0. \end{aligned} \tag{6.8}$$

For the case $\gamma = 1$, by (3.7) and (3.16) we have

$$\begin{aligned}
 J_1(t) &\leq \int_0^{t/2} \|K_1(t-s)P_K(s)F_m(s)\|_{L_k^1} ds \\
 &\leq C \int_0^{t/2} (\|P_K(s)F_m(s)\|_{L_k^1} + (t-s)^{\frac{k}{2}} \|P_K(s)F_m(s)\|_{L^1}) ds \\
 &\leq C \int_0^{t/2} ((1+s)^{-\frac{K-k}{2}} + t^{\frac{k}{2}}(1+s)^{-\frac{K}{2}}) E_{K,q}[F_m](s) ds \\
 &\leq C(1+t)^{\frac{k}{2}} \int_0^\infty (1+s)^{-\frac{K}{2}} E_{K,q}[F_m](s) ds, \quad t > 0.
 \end{aligned} \tag{6.9}$$

Furthermore, for the case $\gamma = \infty$, by (3.7) and (4.8) we obtain

$$\begin{aligned}
 J_1(t) &\leq \int_0^{t/2} \|K_1(t-s)P_K(s)F_m(s)\|_{L_k^\infty} ds \\
 &\leq C \int_0^{t/2} ((1+t-s)^{-\frac{N}{2}} (\|P_K(s)F_m(s)\|_{L_k^1} + \|P_K(s)F_m(s)\|_{L_k^\infty}) ds \\
 &\quad + C \int_0^{t/2} (1+t-s)^{\frac{k-N}{2}} (\|P_K(s)F_m(s)\|_{L^1} + \|P_K(s)F_m(s)\|_{L^\infty}) ds \\
 &\leq C(1+t)^{-\frac{N}{2}} \int_0^{t/2} (1+s)^{-\frac{K-k}{2}} E_{K,q}[F_m](s) ds \\
 &\quad + C(1+t)^{\frac{k-N}{2}} \int_0^{t/2} (1+s)^{-\frac{K}{2}} E_{K,q}[F_m](s) ds \\
 &\leq C(1+t)^{\frac{k-N}{2}} \int_0^\infty (1+s)^{-\frac{K}{2}} E_{K,q}[F_m](s) ds, \quad t > 0.
 \end{aligned} \tag{6.10}$$

Moreover, by (3.7), (3.16) and (4.4) we have

$$\begin{aligned}
 J_2(t) &\leq \int_{t/2}^t \|K_1(t-s)P_K(s)F_m(s)\|_{L_k^\gamma} ds \\
 &\leq C \int_{t/2}^t (\|P_K(s)F_m(s)\|_{L_k^\gamma} + (t-s)^{\frac{k}{2}} \|P_K(s)F_m(s)\|_{L^\gamma}) ds \\
 &\leq C \int_{t/2}^t ((1+s)^{-\frac{K-k}{2} - \frac{N}{2}(1-\frac{1}{\gamma})} + t^{\frac{k}{2}}(1+s)^{-\frac{K}{2} - \frac{N}{2}(1-\frac{1}{\gamma})}) E_{K,q}[F_m](s) ds \\
 &\leq C(1+t)^{\frac{k}{2} - \frac{N}{2}(1-\frac{1}{\gamma})} \int_0^\infty (1+s)^{-\frac{K}{2}} E_{K,q}[F_m](s) ds, \quad t > 0.
 \end{aligned}$$

This together with (6.7), (6.9), (6.10) and $A > 0$ implies that

$$J_1(t) + J_2(t) \leq C(1+t)^{\frac{k}{2} - \frac{N}{2}(1-\frac{1}{\gamma})} \int_0^\infty (1+s)^{-A-1} ds \leq C(1+t)^{\frac{k}{2} - \frac{N}{2}(1-\frac{1}{\gamma})}, \quad t > 0. \tag{6.11}$$

Therefore, by (6.8) and (6.11) we obtain (2.18) with $j = m + 1$. Hence, by induction we see that (2.18) holds for all $j \in \mathbb{N}_0$.

Proof of (2.19). The proof is also by induction. We put

$$u_{\text{lin}}(x, t) := (K_0(t)u_0)(x) + \left(K_1 \left(\frac{1}{2}u_0 + u_1 \right) \right) (x), \tag{6.12}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$. Let $j \in \mathbb{N}_0$ and $F_{-1} \equiv 0$. Since it follows from (2.4), (2.10) and (6.12) that

$$\begin{aligned} u(x, t) - U_j(x, t) &= u_{\text{lin}}(x, t) + \int_0^t (K_1(t-s)F(u(s)))(x) ds \\ &\quad - U_0(x, t) - \int_0^t (K_1(t-s)P_K(s)F_{j-1}(s))(x) ds \\ &= u_{\text{lin}}(x, t) - U_{\text{lin}}(x, t) + \int_0^t (K_1(t-s)F(u(s)))(x) ds \\ &\quad - \sum_{|\alpha| \leq K} \int_0^t (K_1(t-s)M_\alpha(F(u(s)), s)g_\alpha(s))(x) ds \\ &\quad - \int_0^t (K_1(t-s)P_K(s)F_{j-1}(s))(x) ds \\ &= u_{\text{lin}}(x, t) - U_{\text{lin}}(x, t) + \int_0^t (K_1(t-s)P_K(s) \\ &\quad \times (F(u(s)) - F_{j-1}(s)))(x) ds \end{aligned}$$

for all $(x, t) \in \mathbb{R}^N \times (0, \infty)$, for any $q \in [1, \infty]$, we have

$$\begin{aligned} \|u(t) - U_j(t)\|_{L^q} &\leq \|u_{\text{lin}}(t) - U_{\text{lin}}(t)\|_{L^q} + \left\| \int_0^t K_1(t-s)P_K\tilde{F}_{j-1}(s) ds \right\|_{L^q} \\ &\leq \|u_{\text{lin}}(t) - U_{\text{lin}}(t)\|_{L^q} + \left\| \int_{t/2}^t K_1(t-s)P_K\tilde{F}_{j-1}(s) ds \right\|_{L^q} \\ &\quad + \left\| \int_0^{t/2} (K_1(t-s) - V_1(t-s)) P_K\tilde{F}_{j-1}(s) ds \right\|_{L^q} \\ &\quad + \left\| \int_0^{t/2} V_1(t-s)P_K\tilde{F}_{j-1}(s) ds \right\|_{L^q} \\ &=: \tilde{J}_1(t) + \tilde{J}_2(t) + \tilde{J}_3(t) + \tilde{J}_4(t), \quad t > 0, \end{aligned} \tag{6.13}$$

where

$$\tilde{F}_{j-1}(x, t) := F(u(x, t)) - F_{j-1}(x, t). \tag{6.14}$$

By (3.21) we obtain

$$t^{\frac{N}{2}(1-\frac{1}{q})}\tilde{J}_1(t) = o(t^{-\frac{K}{2}}) \quad \text{as } t \rightarrow \infty. \tag{6.15}$$

So it suffices to consider the terms $\tilde{J}_n(t)$ with $n = 2, 3, 4$.

We first consider the case $j = 0$. By (6.2) and (6.14) we have

$$E_{K,q}[\tilde{F}_{-1}](t) = E_{K,q}[F(u)](t) \leq C(1+t)^{\frac{K}{2}-A-1}, \quad t > 0. \tag{6.16}$$

Since $A > 0$, it follows from (3.7), (3.14) and (6.16) that

$$\begin{aligned} t^{\frac{N}{2}(1-\frac{1}{q})} \tilde{J}_2(t) &\leq C t^{\frac{N}{2}(1-\frac{1}{q})} \int_{t/2}^t (1 + e^{-\delta(t-s)}) \|P_K(s) \tilde{F}_{-1}(s)\|_{L^q} ds \\ &\leq C t^{\frac{N}{2}(1-\frac{1}{q})} \int_{t/2}^t (1+s)^{-\frac{K}{2}-\frac{N}{2}(1-\frac{1}{q})} E_{K,q}[\tilde{F}_{-1}](s) ds \\ &\leq C \int_{t/2}^{\infty} (1+s)^{-A-1} ds \leq C(1+t)^{-A}, \quad t > 0. \end{aligned} \quad (6.17)$$

Furthermore, since $A > 0$, by (3.7), (3.19) and (6.16) we obtain

$$\begin{aligned} t^{\frac{N}{2}(1-\frac{1}{q})} \tilde{J}_3(t) &\leq C t^{\frac{N}{2}(1-\frac{1}{q})} \int_0^{t/2} \left((t-s)^{-\frac{N}{2}(1-\frac{1}{q}) - ([\frac{K}{2}] + 1)} \|P_K \tilde{F}_{-1}(s)\|_{L^1} \right. \\ &\quad \left. + e^{-\delta(t-s)} \|P_K \tilde{F}_{-1}(s)\|_{L^q} \right) ds \\ &\leq C t^{-([\frac{K}{2}] + 1)} \int_0^{t/2} (1+s)^{-\frac{K}{2}} E_{K,q}[\tilde{F}_{-1}](s) ds \\ &\leq C t^{-([\frac{K}{2}] + 1)} \int_0^{\infty} (1+s)^{-A-1} ds \leq C t^{-([\frac{K}{2}] + 1)}, \quad t > 0. \end{aligned} \quad (6.18)$$

On the other hand, it follows from (3.18) that

$$\begin{aligned} \|V_1(t-s) P_K \tilde{F}_{-1}(s)\|_{L^q} &\leq C \left\| V \left(\frac{t-s}{2} \right) e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s) \right\|_{L^q} \\ &\leq C (t-s)^{-\frac{N}{2}(1-\frac{1}{q})} \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} \end{aligned}$$

for all $t \geq s + \delta > 0$. This implies that

$$\begin{aligned} t^{\frac{N}{2}(1-\frac{1}{q})} \tilde{J}_4(t) &\leq t^{\frac{N}{2}(1-\frac{1}{q})} \int_0^{t/2} \|V_1(t-s) P_K \tilde{F}_{-1}(s)\|_{L^q} ds \\ &\leq C t^{\frac{N}{2}(1-\frac{1}{q})} \int_0^{t/2} (t-s)^{-\frac{N}{2}(1-\frac{1}{q})} \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} ds \\ &\leq C \int_0^{t/2} \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} ds, \quad t > 0. \end{aligned} \quad (6.19)$$

For any $T > 0$, applying Lemma 3.1 (ii) with (2.9), we have

$$\lim_{t \rightarrow \infty} (t-s)^{\frac{K}{2}} \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} = 0$$

for any $s \in (0, T)$. Then, by the Lebesgue dominated convergence theorem we see that

$$\begin{aligned} \limsup_{t \rightarrow \infty} t^{\frac{K}{2}} \int_0^T \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} ds \\ \leq \limsup_{t \rightarrow \infty} \int_0^T (t-s)^{\frac{K}{2}} \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} ds = 0. \end{aligned} \quad (6.20)$$

Furthermore, applying Lemma 3.1 (i) with (2.9), for any $\delta > 0$, we deduce from (3.7) that

$$\begin{aligned} & \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} \leq C(t-s)^{-\frac{K}{2}} \|P_K \tilde{F}_{-1}(s)\|_{L^1_K} \\ & \leq C(t-s)^{-\frac{K}{2}} E_{K,q}[\tilde{F}_{-1}](s) \end{aligned} \tag{6.21}$$

for all $t \geq s + \delta > 0$. Then, for any $T_0 > 0$, by (6.16) and (6.21) we see that

$$\begin{aligned} \int_T^{t/2} \|e^{(\frac{t-s}{2})\Delta} P_K \tilde{F}_{-1}(s)\|_{L^1} ds & \leq C \int_T^{t/2} (t-s)^{-\frac{K}{2}} E_{K,q}[\tilde{F}_{-1}](s) ds \\ & \leq Ct^{-\frac{K}{2}} \int_T^t (1+s)^{\frac{K}{2}-A-1} ds \end{aligned} \tag{6.22}$$

for all $t \geq 2T$ and $T \geq T_0$. Therefore, by (6.19), (6.20) and (6.22), for any $d > 0$ and $T \geq T_0$, we have

$$t^{\frac{N}{2}(1-\frac{1}{q})} \tilde{J}_4(t) \leq dt^{-\frac{K}{2}} + C_* t^{-\frac{K}{2}} \int_T^t (1+s)^{\frac{K}{2}-A-1} ds \tag{6.23}$$

for all sufficiently large t , where C_* is a constant independent of $T \in [T_0, \infty)$ and $d > 0$. Hence, by (6.17), (6.18) and (6.23) we obtain (2.19) with $j = 0$.

Next we assume that (2.19) holds for some $j = m \in \mathbb{N}_0$. Then, for any $q \in [1, \infty]$, by (1.3), (2.12), (2.18) with $j = m$ and (6.14) we have

$$\begin{aligned} (1+t)^{\frac{N}{2}(1-\frac{1}{q})} \|\tilde{F}_m(t)\|_{L^q} & \leq C(1+t)^{\frac{N}{2}(1-\frac{1}{q})} \|u(t) - U_m(t)\|_{L^q} \\ & \quad \times \max \left\{ \|u(t)\|_{L^\infty}^{p-1}, \|U_m(t)\|_{L^\infty}^{p-1} \right\} \\ & \leq C(1+t)^{\frac{N}{2}(1-\frac{1}{q})-A-1} \|u(t) - U_m(t)\|_{L^q}, \quad t > 0 \end{aligned} \tag{6.24}$$

Similarly to (6.24) with (2.17), (2.13) and (2.14), for any $\gamma \in \{1, \infty\}$, we obtain

$$\begin{aligned} (1+t)^{\frac{N}{2}(1-\frac{1}{\gamma})} \|\tilde{F}_m(t)\|_{L^\gamma_K} & \leq C(1+t)^{\frac{N}{2}(1-\frac{1}{\gamma})} \|u(t) - U_m(t)\|_{L^\infty} \\ & \quad \times \max \left\{ \|u(t)\|_{L^\infty}^{p-2} \|u(t)\|_{L^\gamma_K}, \|U_m(t)\|_{L^\infty}^{p-2} \|U_m(t)\|_{L^\gamma_K} \right\} \\ & \leq C(1+t)^{\frac{N}{2}(1-\frac{1}{\gamma})+\frac{K}{2}-A-1} \|u(t) - U_m(t)\|_{L^\infty}, \quad t > 0. \end{aligned} \tag{6.25}$$

By (2.19) with $j = m$, (3.5), (6.24) and (6.25) we see that

$$\begin{aligned} & E_{K,q}[\tilde{F}_m](t) \\ & \leq C(1+t)^{\frac{K}{2}-A-1} \left((1+t)^{\frac{N}{2}(1-\frac{1}{q})} \|u(t) - U_m(t)\|_{L^q} + \|u(t) - U_m(t)\|_{L^\infty} \right) \\ & = \begin{cases} o(t^{-A-1}) + O(t^{\frac{K}{2}-(m+2)A-1}) & \text{if } (m+1)A \neq K/2, \\ O(t^{-A-1} \log t) & \text{if } (m+1)A = K/2, \end{cases} \end{aligned}$$

as $t \rightarrow \infty$. Therefore, applying same arguments as in the proof of the case $j = 0$, we can easily show that

$$t^{\frac{N}{2}(1-\frac{1}{q})} \left(\tilde{J}_2(t) + \tilde{J}_3(t) + \tilde{J}_4(t) \right) = \begin{cases} o(t^{-\frac{K}{2}}) + O(t^{-(m+2)A}) & \text{if } (m+2)A \neq K/2, \\ O(t^{-\frac{K}{2}} \log t) & \text{if } (m+2)A = K/2, \end{cases}$$

as $t \rightarrow \infty$. This together with (6.13) and (6.15) implies (2.19) with $j = m + 1$. Hence, by induction we see that (2.19) holds for all $j \in \mathbb{N}_0$, and the proof of Theorem 2.2 is complete. \square

Appendix

In this Appendix, for the convenience of the reader, we give the proof of (2.9). This proof is completely same as in the one of assertion (ii) of [18, Lemma 2.1].

Proof of (2.9). Let $K \geq 0$. For any $f \in L_K^1$ and $\alpha \in \mathbb{M}_K$, since

$$\int_{\mathbb{R}^N} x^\alpha g_\beta(x, t) dx = 0 \quad \text{if not } \beta \leq \alpha,$$

by (2.6) and (2.8) we have

$$\begin{aligned} \int_{\mathbb{R}^N} x^\alpha [P_K(t)f](x) dx &= \int_{\mathbb{R}^N} x^\alpha f(x) dx - \sum_{\beta \leq \alpha} M_\beta(f, t) \int_{\mathbb{R}^N} x^\alpha g_\beta(x, t) dx \\ &= \int_{\mathbb{R}^N} x^\alpha f(x) dx - M_\alpha(f, t) - \sum_{\beta \leq \alpha, \alpha \neq \beta} M_\beta(f, t) \int_{\mathbb{R}^N} x^\alpha g_\beta(x, t) dx = 0 \end{aligned}$$

for $t \geq 0$. Thus (2.9) follows. \square

Acknowledgments

The authors would like to thank referee for useful comments for the original manuscript. The work of the first author (T. Kawakami) was supported in part by Grant-in-Aid for Young Scientists (B) (No. 24740107) and (No. 16K17629) of JSPS (Japan Society for the Promotion of Science) and by the JSPS Program for Advancing Strategic International Networks to Accelerate the Circulation of Talented Researchers “Mathematical Science of Symmetry, Topology and Moduli, Evolution of International Research Network based on OCAMI”, and the second author (H. Takeda) by Grant-in-Aid for Young Scientists (B) (No. 15K17581) of JSPS.

References

- [1] Abramowitz, M., Stegun, I.A.: Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables. Dover Publications Inc., New York (1992) (reprint of the 1972 edition)
- [2] Bellout, H., Friedman, A.: Blow-up estimates for a nonlinear hyperbolic heat equation. SIAM J. Math. Anal. **20**, 354–366 (1989)
- [3] Courant, R., Hilbert, D.: Methods of Mathematical Physics II. Wiley, New York (1989)

- [4] D'Abbicco, M., Ebert, M.R.: An application of L^p - L^q decay estimates to the semi-linear wave equation with parabolic-like structural damping. *Nonlinear Anal.* **99**, 16–34 (2014)
- [5] D'Abbicco, M., Ebert, M.R.: Diffusion phenomena for the wave equation with structural damping in the L^p - L^q framework. *J. Differ. Equ.* **256**, 2307–2336 (2014)
- [6] D'Abbicco, M., Reissig, M.: Semilinear structural damped waves. *Math. Methods Appl. Sci.* **37**, 1570–1592 (2014)
- [7] Gallay, T., Raugel, G.: Scaling variables and asymptotic expansions in damped wave equations. *J. Differ. Equ.* **150**, 42–97 (1998)
- [8] Hayashi, N., Kaikina, E.I., Naumkin, P.I.: Damped wave equation with super critical nonlinearities. *Differ. Integr. Equ.* **17**, 637–652 (2004)
- [9] Hayashi, N., Kaikina, E.I., Naumkin, P.I.: On the critical nonlinear damped wave equation with large initial data. *J. Math. Anal. Appl.* **334**, 1400–1425 (2007)
- [10] Hosono, T., Ogawa, T.: Large time behavior and L^p - L^q estimate of 2-dimensional nonlinear damped wave equations. *J. Differ. Equ.* **203**, 82–118 (2004)
- [11] Ikehata, R.: Asymptotic profiles for wave equations with strong damping. *J. Differ. Equ.* **257**, 2159–2177 (2014)
- [12] Ikehata, R., Miyaoka, Y., Nakatake, T.: Decay estimates of solutions for dissipative wave equations in \mathbb{R}^N with lower power nonlinearities. *J. Math. Soc. Japan* **56**, 365–373 (2004)
- [13] Ikehata, R., Nishihara, K., Zhao, H.: Global asymptotics of solutions to the Cauchy problem for the damped wave equation with absorption. *J. Differ. Equ.* **226**, 1–29 (2006)
- [14] Ikehata, R., Tanizawa, K.: Global existence of solutions for semilinear damped wave equations in \mathbb{R}^N with noncompactly supported initial data. *Nonlinear Anal.* **61**, 1189–1208 (2005)
- [15] Ikehata, R., Todorova, G., Yordanov, B.: Wave equations with strong damping in Hilbert spaces. *J. Differ. Equ.* **254**, 3352–3368 (2013)
- [16] Ishige, K., Ishiwata, M., Kawakami, T.: The decay of the solutions for the heat equation with a potential. *Indiana Univ. Math. J.* **58**, 2673–2708 (2009)
- [17] Ishige, K., Kawakami, T.: Refined asymptotic profiles for a semilinear heat equation. *Math. Ann.* **353**, 161–192 (2012)
- [18] Ishige, K., Kawakami, T., Kobayashi, K.: Asymptotics for a nonlinear integral equation with a generalized heat kernel. *J. Evol. Equ.* **14**, 749–777 (2014)
- [19] Jonov, B., Sideris, T.C.: Global and almost global existence of small solutions to a dissipative wave equation in 3D with nearly null nonlinear terms. *Commun. Pure Appl. Anal.* **14**, 1407–1442 (2015)

- [20] Karch, G.: Selfsimilar profiles in large time asymptotics of solutions to damped wave equations. *Studia Math.* **143**, 175–197 (2000)
- [21] Kawakami, T., Ueda, Y.: Asymptotic profiles to the solutions for a nonlinear damped wave equations. *Differ. Integr. Equ.* **26**, 781–814 (2013)
- [22] Kawashima, S., Nakao, M., Ono, K.: On the decay property of solutions to the Cauchy problem of the semilinear wave equation with a dissipative term. *J. Math. Soc. Japan* **47**, 617–653 (1995)
- [23] Li, T.-T.: Nonlinear heat conduction with finite speed of propagation. In: *China-Japan Symposium on Reaction-Diffusion Equations and their Applications and Computational Aspects* (Shanghai, 1994), pp. 81–91. World Sci. Publ., River Edge (1997)
- [24] Lin, J., Nishihara, K., Zhai, J.: Decay property of solutions for damped wave equations with space-time dependent damping term. *J. Math. Anal. Appl.* **374**, 602–614 (2011)
- [25] Lin, J., Nishihara, K., Zhai, J.: Critical exponent for the semilinear wave equation with time-dependent damping. *Discrete Contin. Dyn. Syst.* **32**, 4307–4320 (2012)
- [26] Li, T.T., Zhou, Y.: Breakdown of solutions to $\square u + u_t = |u|^{1+\alpha}$. *Discrete Contin. Dyn. Syst.* **1**, 503–520 (1995)
- [27] Marcati, P., Nishihara, K.: The L^p - L^q estimates of solutions to one-dimensional damped wave equations and their application to compressible flow through porous media. *J. Differ. Equ.* **191**, 445–469 (2003)
- [28] Matsumura, A.: On the asymptotic behavior of solutions of semi-linear wave equations. *Publ. Res. Inst. Math. Sci.* **12**, 169–189 (1976)
- [29] Narazaki, T.: L^p - L^q estimates for the damped wave equations and their applications to semi-linear problem. *J. Math. Soc. Japan* **56**, 585–626 (2004)
- [30] Nishihara, K.: L^p - L^q estimates of solutions to the damped wave equation in 3-dimensional space and their application. *Math. Z.* **244**, 631–649 (2003)
- [31] Nishihara, K.: Decay properties for the damped wave equation with space dependent potential and absorbed semilinear term. *Commun. Partial Differ. Equ.* **35**, 1402–1418 (2010)
- [32] Nishihara, K.: Asymptotic behavior of solutions to the semilinear wave equation with time-dependent damping. *Tokyo J. Math.* **34**, 327–343 (2011)
- [33] Orive, R., Zuazua, E., Pazoto, A.F.: Asymptotic expansion for damped wave equations with periodic coefficients. *Math. Methods Appl. Sci.* **11**, 1285–1310 (2001)
- [34] Ponce, G.: Global existence of small solutions to a class of nonlinear evolution equations. *Nonlinear Anal.* **9**, 399–418 (1985)

- [35] Shibata, Y.: On the rate of decay of solutions to linear viscoelastic equation. *Math. Methods Appl. Sci.* **23**, 203–226 (2000)
- [36] Takeda, H.: Higher-order expansion of solutions for a damped wave equation. *Asymptot. Anal.* **94**, 1–31 (2015)
- [37] Takeda, H.: Large time behavior of solutions for a nonlinear damped wave equation. *Commun. Pure Appl. Anal.* **15**, 41–55 (2016)
- [38] Todorova, G., Yordanov, B.: Critical exponent for a nonlinear wave equation with damping. *J. Differ. Equ.* **174**, 464–489 (2001)
- [39] Wakasugi, Y.: Scaling variables and asymptotic profiles of solutions to the semi-linear damped wave equation with variable coefficients. [arXiv:1508.05778](https://arxiv.org/abs/1508.05778)
- [40] Wirth, J.: Wave equations with time-dependent dissipation. I. Non-effective dissipation. *J. Differ. Equ.* **222**, 487–514 (2006)
- [41] Wirth, J.: Wave equations with time-dependent dissipation. II. Effective dissipation. *J. Differ. Equ.* **232**, 74–103 (2007)
- [42] Yamazaki, T.: Asymptotic behavior for abstract wave equations with decaying dissipation. *Adv. Differ. Equ.* **11**, 419–456 (2006)
- [43] Yang, H., Milani, A.: On the diffusion phenomenon of quasilinear hyperbolic waves. *Bull. Sci. Math.* **124**, 415–433 (2000)
- [44] Zhang, Q.S.: A blow-up result for a nonlinear wave equation with damping; The critical case, *C.R. Acad. Sci. Paris, Série I*, **333** (2001), 109–114.

Tatsuki Kawakami
Department of Mathematical Sciences
Osaka Prefecture University
Sakai 599-8531, Japan
e-mail: kawakami@ms.osakafu-u.ac.jp

Hiroshi Takeda
Faculty of Engineering
Fukuoka Institute of Technology
Fukuoka 811-0295, Japan
e-mail: h-takeda@fit.ac.jp

Received: 10 May 2016.

Accepted: 6 September 2016.