

Breakings of the neutrino μ - τ reflection symmetry

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ABSTRACT: The neutrino μ - τ reflection symmetry has been attracting a lot of attention as it predicts the interesting results $\theta_{23} = \pi/4$ and $\delta = \pm\pi/2$. But it is reasonable to consider breakings of such a symmetry either from the theoretical considerations or on the basis of experimental results. We thus perform a systematic study for the possible symmetry-breaking patterns and their implications for the mixing parameters. The general treatment is applied to some specific symmetry breaking arising from the renormalization group effects for illustration.

KEYWORDS: Discrete Symmetries, Neutrino Physics, CP violation

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1 Introduction

The discovery of neutrino oscillations indicates that neutrinos are massive and bear flavor mixing. The neutrino mixing arises from the mismatch between their mass and flavor eigenstates, and is described by a 3×3 unitary matrix $U = U_l^\dagger U_\nu$ (with U_l and U_ν being respectively the unitary matrix for diagonalizing the charged-lepton mass matrix $M_l M_l^\dagger$ and neutrino mass matrix M_ν). In the standard parametrization, U reads

$$U = P_\phi \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} P_\nu, \quad (1.1)$$

where θ_{ij} (for $ij = 12, 13, 23$) are the mixing angles (with $c_{ij} = \cos \theta_{ij}$ and $s_{ij} = \sin \theta_{ij}$) and δ is the Dirac CP phase. $P_\nu = \text{Diag}(e^{i\rho}, e^{i\sigma}, 1)$ contains two Majorana CP phases ρ and σ , while $P_\phi = \text{Diag}(e^{i\phi_1}, e^{i\phi_2}, e^{i\phi_3})$ consists of three unphysical phases $\phi_{1,2,3}$ that can be removed via the charged-lepton field rephasing. In addition, neutrino oscillations are also controlled by two mass-squared differences $\Delta m_{ij}^2 = m_i^2 - m_j^2$ (for $ij = 21, 31$). Thanks to various neutrino-oscillation experiments [1], the neutrino mixing parameters have been measured to a good accuracy. A global-fit result [2] for them is given by

$$\begin{aligned} \sin^2 \theta_{12} &= 0.308 \pm 0.017, & \Delta m_{21}^2 &= (7.54 \pm 0.24) \times 10^{-5} \text{ eV}^2, \\ \sin^2 \theta_{13} &= 0.0234 \pm 0.0020, & |\Delta m_{31}^2| &= (2.47 \pm 0.06) \times 10^{-3} \text{ eV}^2. \end{aligned} \quad (1.2)$$

Note that the sign of Δm_{31}^2 remains undetermined, allowing for two possible neutrino mass orderings $m_1 < m_2 < m_3$ (referred to as the normal hierarchy and NH for short) or $m_3 < m_1 < m_2$ (the inverted hierarchy and IH for short). The absolute neutrino mass scale is not known either, but subject to the constraint $m_1 + m_2 + m_3 < 0.23 \text{ eV}$ from cosmological observations [3]. Particularly noteworthy, a recent result from the NOvA experiment ($\theta_{23} = 39.5^\circ \pm 1.7^\circ$ or $52.1^\circ \pm 1.7^\circ$ in the NH case) disfavors the popular maximal mixing scenario $\theta_{23} = 45^\circ$ with 2.6σ significance [4]. On the other hand, it is

interesting to find that the best-fit result for δ is around 270° ($261^\circ \pm 55^\circ$ for NH and $277^\circ \pm 43^\circ$ for IH) [5].

How to understand the neutrino mixing pattern poses an interesting question. As symmetries (e.g., the $SU(3)_q$ quark flavor symmetry) have been serving as a guideline for understanding the particle physics, they may play a similar role in addressing the flavor issues. Along this line, many discrete groups have been proposed as the lepton flavor symmetry [6, 7]. A simplest example is the μ - τ permutation symmetry [8–12]: in the basis of M_l being diagonal, M_ν should keep unchanged with respect to the transformation $\nu_\mu \leftrightarrow \nu_\tau$ and thus feature $M_{e\mu} = M_{e\tau}$ and $M_{\mu\mu} = M_{\tau\tau}$ (with $M_{\alpha\beta}$ for $\alpha, \beta = e, \mu, \tau$ being the matrix elements of M_ν). Such a symmetry (which results in $\theta_{23} = \pi/4$ and $\theta_{13} = 0$) was historically motivated by the experimental facts that θ_{23} takes a value close to $\pi/4$ while θ_{13} was only constrained by $\sin^2 2\theta_{13} < 0.18$ [13] (and thus might be negligibly small) at the time. However, the relatively large $\theta_{13} \simeq 0.15$ observed recently [14] requires a significant breaking of this symmetry unless neutrinos are quasi-degenerate in masses [15]. Hence we need to go beyond this simple possibility to accommodate the experimental results in a better way. In this connection, the μ - τ reflection symmetry [12, 16] may serve as a unique alternative: when M_l is diagonal, M_ν should remain invariant under the transformation¹

$$\nu_e \leftrightarrow \nu_e^c, \quad \nu_\mu \leftrightarrow \nu_\tau^c, \quad \nu_\tau \leftrightarrow \nu_\mu^c, \quad (1.3)$$

and thus be characterized by

$$M_{e\mu} = M_{e\tau}^*, \quad M_{\mu\mu} = M_{\tau\tau}^*, \quad M_{ee} = M_{ee}^*, \quad M_{\mu\tau} = M_{\mu\tau}^*. \quad (1.4)$$

In addition to allowing for an arbitrary θ_{13} , this symmetry predicts $\theta_{23} = \pi/4$ and $\delta = \pm\pi/2$ [19] which are close to the present data, thereby having been attracting a lot of interest [20–34]. Moreover, ρ and σ are required to take the trivial values 0 or $\pi/2$.

Nevertheless, it is hard to believe that the μ - τ reflection symmetry can remain as an exact one. On the experimental side, the aforementioned results seem to hint towards $\theta_{23} \neq \pi/4$ (and possibly $\delta \neq \pm\pi/2$). On the theoretical side, flavor symmetries are generally implemented at a superhigh energy scale and so the renormalization group (RG) running effect may provide a source for the symmetry breaking as we will see. In view of these considerations, it is worthwhile to consider the breaking of this symmetry. In the next section, we perform a systematic study of the possible symmetry-breaking patterns and their implications for the mixing parameters. First of all, we establish an equation set relating the symmetry-breaking parameters in an M_ν of approximate μ - τ reflection symmetry and the deviations of mixing parameters from their special values taken in the symmetry context. While the numerical results for these equations are analyzed in section 2.1, some analytical approximations will be derived in section 2.2 to explain the corresponding numerical results. In section 2.3 the general treatment is applied to some specific symmetry breaking arising from the RG running effect. Finally, we summarize our main results in section 3.

¹This operation is a combination of the μ - τ exchange and CP conjugate transformations — a typical kind of the generalized CP transformations [17, 18].

2 Breaking of the μ - τ reflection symmetry

Above all, we define the following parameters to characterize breaking of μ - τ reflection symmetry

$$\epsilon_1 = \frac{M_{e\mu} - M_{e\tau}^*}{M_{e\mu} + M_{e\tau}^*}, \quad \epsilon_2 = \frac{M_{\mu\mu} - M_{\tau\tau}^*}{M_{\mu\mu} + M_{\tau\tau}^*}, \quad \epsilon_3 = \frac{\text{Im}(M_{ee})}{\text{Re}(M_{ee})}, \quad \epsilon_4 = \frac{\text{Im}(M_{\mu\tau})}{\text{Re}(M_{\mu\tau})}, \quad (2.1)$$

by analogy with the discussions about breaking of μ - τ permutation symmetry in ref. [35]. Note that they correspond to the four symmetry conditions in eq. (1.4) one by one. These parameters have to be small in magnitude (say $|\epsilon_i| \leq 0.1$ for $i = 1, 2, 3, 4$) so as to keep the μ - τ reflection symmetry as an approximate one. In terms of them, the most general neutrino mass matrix of an approximate μ - τ reflection symmetry can always be parameterized in a manner as follows: suppose, at the symmetry level, there is a neutrino mass matrix of the form

$$M_\nu^{(0)} = \begin{pmatrix} A_0 & B_0 & B_0^* \\ B_0 & C_0 & D_0 \\ B_0^* & D_0 & C_0^* \end{pmatrix}, \quad (2.2)$$

in which A_0 and D_0 are real. This neutrino mass matrix can be diagonalized by a unitary matrix $U^{(0)}$ (an analogue of U) with its parameters satisfying the requirements

$$\phi_1^{(0)} = \phi_2^{(0)} + \phi_3^{(0)} = 0, \quad \theta_{23}^{(0)} = \pi/4, \quad \delta^{(0)} = \pm\pi/2, \quad \rho^{(0)}, \sigma^{(0)} = 0 \text{ or } \pi/2. \quad (2.3)$$

After the symmetry is softly broken, $M_\nu^{(0)}$ may receive a general perturbation as given by

$$M_\nu^{(1)} = \begin{pmatrix} \delta_{ee} & \delta_{e\mu} & \delta_{e\tau} \\ \delta_{e\mu} & \delta_{\mu\mu} & \delta_{\mu\tau} \\ \delta_{e\tau} & \delta_{\mu\tau} & \delta_{\tau\tau} \end{pmatrix}, \quad (2.4)$$

which can be decomposed into two parts as

$$M_\nu^{(1)} = \frac{1}{2} \begin{pmatrix} 2\text{Re}(\delta_{ee}) & \delta_{e\mu} + \delta_{e\tau}^* & \delta_{e\mu}^* + \delta_{e\tau} \\ \delta_{e\mu} + \delta_{e\tau}^* & \delta_{\mu\mu} + \delta_{\tau\tau}^* & 2\text{Re}(\delta_{\mu\tau}) \\ \delta_{e\mu}^* + \delta_{e\tau} & 2\text{Re}(\delta_{\mu\tau}) & \delta_{\mu\mu}^* + \delta_{\tau\tau} \end{pmatrix} + \frac{1}{2} \begin{pmatrix} 2i\text{Im}(\delta_{ee}) & \delta_{e\mu} - \delta_{e\tau}^* & \delta_{e\tau} - \delta_{e\mu}^* \\ \delta_{e\mu} - \delta_{e\tau}^* & \delta_{\mu\mu} - \delta_{\tau\tau}^* & 2i\text{Im}(\delta_{\mu\tau}) \\ \delta_{e\tau} - \delta_{e\mu}^* & 2i\text{Im}(\delta_{\mu\tau}) & \delta_{\tau\tau} - \delta_{\mu\mu}^* \end{pmatrix}. \quad (2.5)$$

Consequently, the complete neutrino mass matrix $M_\nu = M_\nu^{(0)} + M_\nu^{(1)}$ can be parameterized as

$$M_\nu = \begin{pmatrix} A(1 + i\epsilon_3) & B(1 + \epsilon_1) & B^*(1 - \epsilon_1^*) \\ B(1 + \epsilon_1) & C(1 + \epsilon_2) & D(1 + i\epsilon_4) \\ B^*(1 - \epsilon_1^*) & D(1 + i\epsilon_4) & C^*(1 - \epsilon_2^*) \end{pmatrix}, \quad (2.6)$$

with

$$\begin{aligned} A &= A_0 + \text{Re}(\delta_{ee}), & B &= B_0 + \frac{\delta_{e\mu} + \delta_{e\tau}^*}{2}, \\ D &= D_0 + \text{Re}(\delta_{\mu\tau}), & C &= C_0 + \frac{\delta_{\mu\mu} + \delta_{\tau\tau}^*}{2}, \end{aligned} \quad (2.7)$$

and

$$\epsilon_1 = \frac{\delta_{e\mu} - \delta_{e\tau}^*}{2B}, \quad \epsilon_2 = \frac{\delta_{\mu\mu} - \delta_{\tau\tau}^*}{2C}, \quad \epsilon_3 = \frac{\text{Im}(\delta_{ee})}{A}, \quad \epsilon_4 = \frac{\text{Im}(\delta_{\mu\tau})}{D}. \quad (2.8)$$

It should be noted that $\text{Im}(\epsilon_{1,2})$ and $\epsilon_{3,4}$ will transform in a way as

$$\begin{aligned} \text{Im}(\epsilon_1) &\rightarrow \text{Im}(\epsilon_1) + \varphi_1 + (\varphi_2 + \varphi_3)/2, & \epsilon_3 &\rightarrow \epsilon_3 + 2\varphi_1, \\ \text{Im}(\epsilon_2) &\rightarrow \text{Im}(\epsilon_2) + \varphi_2 + \varphi_3, & \epsilon_4 &\rightarrow \epsilon_4 + \varphi_2 + \varphi_3, \end{aligned} \quad (2.9)$$

under the neutrino-field rephasing

$$\begin{aligned} \nu_e &\rightarrow e^{i\varphi_1}\nu_e \simeq (1 + i\varphi_1)\nu_e, & \nu_\mu &\rightarrow e^{i\varphi_2}\nu_\mu \simeq (1 + i\varphi_2)\nu_\mu, \\ \nu_\tau &\rightarrow e^{i\varphi_3}\nu_\tau \simeq (1 + i\varphi_3)\nu_\tau, \end{aligned} \quad (2.10)$$

with $\varphi_{1,2,3}$ being some small parameters comparable to ϵ_i . Taking advantage of such a freedom, one can always achieve $\epsilon_{3,4} = 0$ from the general case given by eq. (2.6). In the following discussions, we therefore concentrate on this particular case without loss of generality.

Starting from an M_ν of the form in eq. (2.6) but with $\epsilon_{3,4} = 0$, we study the dependence of mixing parameters on $\epsilon_{1,2}$. To this end, we diagonalize such an M_ν with one unitary matrix in a straightforward way

$$U^\dagger M_\nu U^* = \text{Diag}(m_1, m_2, m_3). \quad (2.11)$$

The mixing parameters in U are supposed to lie around those special values in eq. (2.3) and the corresponding deviations

$$\begin{aligned} \Delta\phi_1 &= \phi_1 - 0, & \Delta\phi &= (\phi_2 + \phi_3)/2 - 0, & \Delta\theta &= \theta_{23} - \pi/4, \\ \Delta\delta &= \delta - \delta^{(0)}, & \Delta\rho &= \rho - \rho^{(0)}, & \Delta\sigma &= \sigma - \sigma^{(0)}, \end{aligned} \quad (2.12)$$

are some small quantities. By making perturbation expansions for these small quantities in eq. (2.11), one reaches the following relations connecting the mixing-parameter deviations with $\epsilon_{1,2}$

$$\begin{aligned} m_3 s_{13}^2 \Delta\delta + \bar{m}_1 c_{12}^2 \Delta\rho + \bar{m}_2 s_{12}^2 \Delta\sigma &= (m_3 s_{13}^2 - m_{11}) \Delta\phi_1, \\ 2m_{12} \bar{s}_{13} \Delta\theta - m_{11} s_{13}^2 \Delta\delta - \bar{m}_1 s_{12}^2 \Delta\rho - \bar{m}_2 c_{12}^2 \Delta\sigma &= (m_{22} - m_3) \Delta\phi, \\ [m_{12} + i(m_{11} + m_3) \bar{s}_{13}] \Delta\theta - (m_{11} - m_3) \bar{s}_{13} \Delta\delta \\ - 2\bar{m}_1 c_{12} (i s_{12} + c_{12} \bar{s}_{13}) \Delta\rho + 2\bar{m}_2 s_{12} (i c_{12} - s_{12} \bar{s}_{13}) \Delta\sigma &= \\ [m_{12} - i(m_{11} + m_3) \bar{s}_{13}] (i \Delta\phi_1 + i \Delta\phi - \epsilon_1), \\ 2(m_{22} - m_3) \Delta\theta - 2(m_{12} - i m_{11} \bar{s}_{13}) \bar{s}_{13} \Delta\delta \\ - 2\bar{m}_1 s_{12} (i s_{12} + 2c_{12} \bar{s}_{13}) \Delta\rho - 2\bar{m}_2 c_{12} (i c_{12} - 2s_{12} \bar{s}_{13}) \Delta\sigma &= \\ (m_{22} + m_3 - 2i m_{12} \bar{s}_{13}) (2i \Delta\phi - \epsilon_2). \end{aligned} \quad (2.13)$$

In order to make the expressions compact, the definitions

$$\begin{aligned}
 m_{11} &= \bar{m}_1 c_{12}^2 + \bar{m}_2 s_{12}^2, & m_{12} &= (\bar{m}_1 - \bar{m}_2) c_{12} s_{12}, & m_{22} &= \bar{m}_1 s_{12}^2 + \bar{m}_2 c_{12}^2, \\
 \bar{m}_1 &= m_1 \exp \left[2i\rho^{(0)} \right], & \bar{m}_2 &= m_2 \exp \left[2i\sigma^{(0)} \right], & \bar{s}_{13} &= -is_{13} \exp \left[i\delta^{(0)} \right],
 \end{aligned} \tag{2.14}$$

have been taken.

2.1 Numerical results

By solving eq. (2.13), one can learn how much mixing-parameter deviations are predicted given some small values of symmetry-breaking parameters. As we will see, the resulting mixing-parameter deviations depend on the neutrino mass spectrum and the combinations of $\rho^{(0)}$ and $\sigma^{(0)}$. So they are to be evaluated as functions of the lightest neutrino mass (m_1 for NH or m_3 for IH) for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ (i.e., $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$, $[\pi/2, 0]$, $[0, \pi/2]$ and $[\pi/2, \pi/2]$). As typical examples, the mixing-parameter deviations respectively arising from $R_1 \equiv \text{Re}(\epsilon_1) = 0.1$ (in figure 1), $I_1 \equiv \text{Im}(\epsilon_1) = 0.1$ (in figure 2), $R_2 \equiv \text{Re}(\epsilon_2) = 0.1$ (in figure 3) and $I_2 \equiv \text{Im}(\epsilon_2) = 0.1$ (in figure 4) are presented in such a way. Here and in the following, the black, red, green and blue colors are assigned to $\Delta\theta$, $\Delta\delta$, $\Delta\rho$ and $\Delta\sigma$, respectively. To save space, the absolute value of a mixing-parameter deviation will be represented by a dashed line when it is negative. By contrast, the full line will be used when the mixing-parameter deviations are positive. In consideration of the experimental sensitivity, the region where the mixing-parameter deviations have magnitudes smaller than 0.01 has not been shown. In doing the calculations we have specified $\delta^{(0)} = -\pi/2$. If it takes the opposite value $\pi/2$, the mixing-parameter deviations produced by R_1 or R_2 will change in the way

$$\Delta\theta \rightarrow \Delta\theta, \quad \Delta\delta \rightarrow -\Delta\delta, \quad \Delta\rho \rightarrow -\Delta\rho, \quad \Delta\sigma \rightarrow -\Delta\sigma, \tag{2.15}$$

while those produced by I_1 or I_2 will change in the way

$$\Delta\theta \rightarrow -\Delta\theta, \quad \Delta\delta \rightarrow \Delta\delta, \quad \Delta\rho \rightarrow \Delta\rho, \quad \Delta\sigma \rightarrow \Delta\sigma. \tag{2.16}$$

The point is that eq. (2.13) stays invariant with respect to the transformations

$$\begin{aligned}
 \Delta\theta &\rightarrow \Delta\theta, & \Delta\delta &\rightarrow -\Delta\delta, & \Delta\rho &\rightarrow -\Delta\rho, & \Delta\sigma &\rightarrow -\Delta\sigma, \\
 R_1 &\rightarrow R_1, & I_1 &\rightarrow -I_1, & R_2 &\rightarrow R_2, & I_2 &\rightarrow -I_2,
 \end{aligned} \tag{2.17}$$

combined with $\bar{s}_{13} \rightarrow -\bar{s}_{13}$ as well as $\Delta\phi_1 \rightarrow -\Delta\phi_1$ and $\Delta\phi \rightarrow -\Delta\phi$. For the convenience of readers, from the results in figures 1–4 we take some representative values of $\Delta\theta$ (in table 1), $\Delta\delta$ (in table 2), $\Delta\rho$ (in table 3) and $\Delta\sigma$ (in table 4) at $m_1(m_3) = 0.001, 0.01$ and 0.1 eV in the NH (IH) case. When a value of the mixing-parameter deviations falls in the range $[-0.01, 0]$ or $[0, 0.01]$, it will be reported as 0.00 or -0.00 .

There are two immediate remarks for the above results. For the convenience of discussions, we use the contribution of R_1 to $\Delta\delta$ as an illustration. But the following discussions apply to the contribution of each symmetry-breaking parameter to each mixing-parameter deviation. (1) The resulting $\Delta\delta$ still can be inferred from the results in figure 1 when

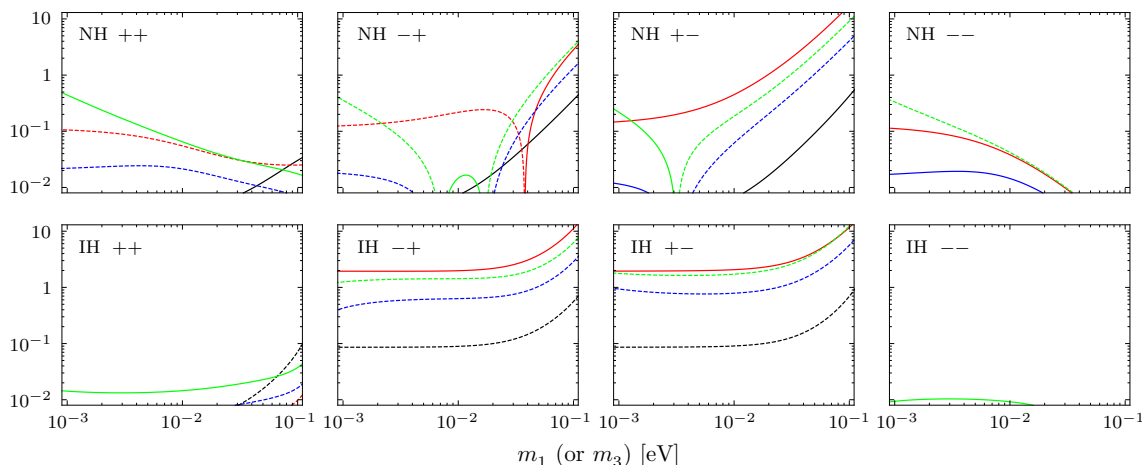


Figure 1. The mixing-parameter deviations induced by $R_1 = 0.1$ ($\Delta\theta$ in black, $\Delta\delta$ in red, $\Delta\rho$ in green and $\Delta\sigma$ in blue) against the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$. The signs ++, --, +- and -- respectively stand for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0], [\pi/2, 0], [0, \pi/2]$ and $[\pi/2, \pi/2]$.

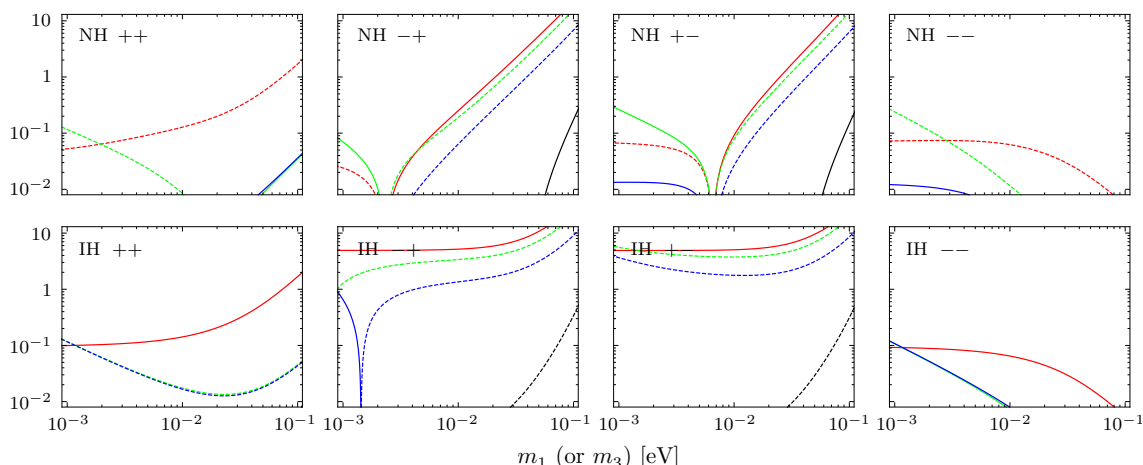


Figure 2. The mixing-parameter deviations induced by $I_1 = 0.1$ ($\Delta\theta$ in black, $\Delta\delta$ in red, $\Delta\rho$ in green and $\Delta\sigma$ in blue) against the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$. The signs ++, --, +- and -- respectively stand for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0], [\pi/2, 0], [0, \pi/2]$ and $[\pi/2, \pi/2]$.

R_1 takes a value other than 0.1.² Since eq. (2.13) is a system of linear equations, $\Delta\delta$ is expected to have a linear dependence on R_1 , implying that the ratio $\Delta\delta/R_1$ must be a quantity independent of the value of R_1 . The size of this ratio measures the sensitivity of $\Delta\delta$ to R_1 . When it is of $\leq \mathcal{O}(0.1)$, $\mathcal{O}(1)$ or $\geq \mathcal{O}(10)$, we say that $\Delta\delta$ may receive a small, modest or large contribution from R_1 . For $R_1 = 0.1$, the “small”, “modest” and “large” contributions are respectively of $\leq \mathcal{O}(0.01)$, $\mathcal{O}(0.1)$ and $\mathcal{O}(1)$. Provided that the

²Note that in some scenarios (e.g., the RG-induced symmetry breaking as will be discussed in section 2.3) the symmetry-breaking parameters have no chance to acquire a size up to 0.1.

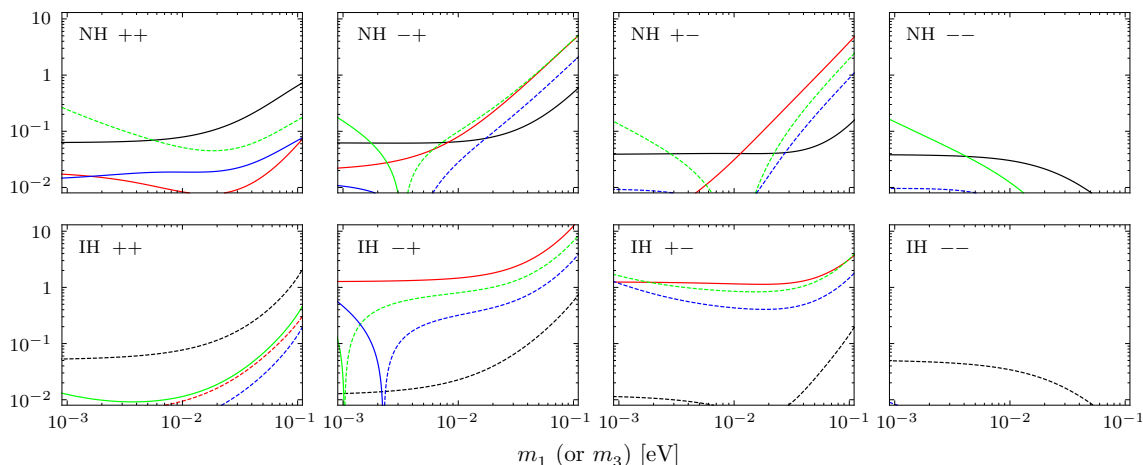


Figure 3. The mixing-parameter deviations induced by $R_2 = 0.1$ ($\Delta\theta$ in black, $\Delta\delta$ in red, $\Delta\rho$ in green and $\Delta\sigma$ in blue) against the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$. The signs ++, --, +- and -- respectively stand for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0], [\pi/2, 0], [0, \pi/2]$ and $[\pi/2, \pi/2]$.

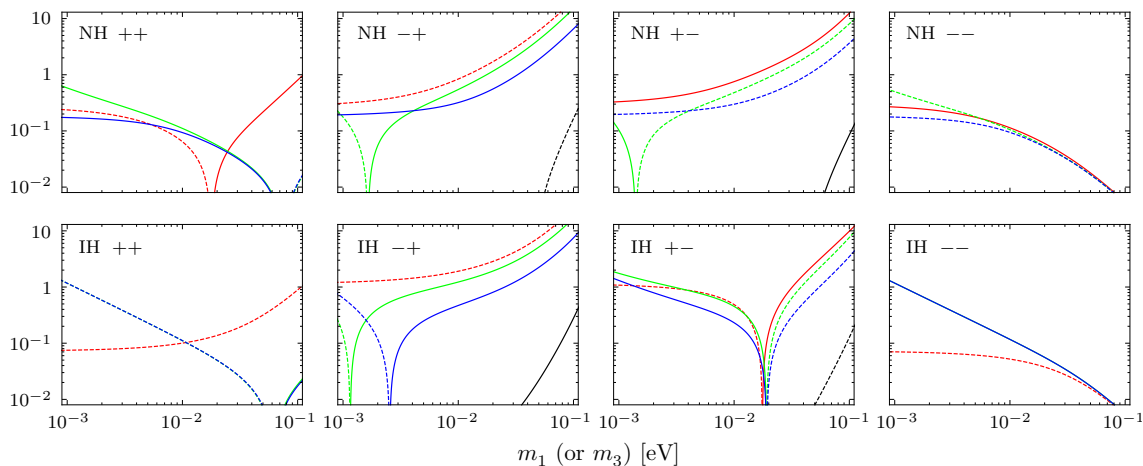


Figure 4. The mixing-parameter deviations induced by $I_2 = 0.1$ ($\Delta\theta$ in black, $\Delta\delta$ in red, $\Delta\rho$ in green and $\Delta\sigma$ in blue) against the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$. The signs ++, --, +- and -- respectively stand for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0], [\pi/2, 0], [0, \pi/2]$ and $[\pi/2, \pi/2]$.

$\Delta\delta'$ arising from a given value R'_1 of R_1 (e.g., 0.1) has already been known (e.g., the red lines in figure 1), the $\Delta\delta''$ generated by another value R''_1 of R_1 can be directly obtained as

$$\Delta\delta'' = \frac{\Delta\delta'}{R'_1} R''_1 = \frac{R''_1}{R'_1} \Delta\delta'. \quad (2.18)$$

For example, the $\Delta\delta$ produced by $R_1 = 0.01$ will be a tenth of the result given by figure 1. (2) But there is one thing to take care of: in some cases (in particular the case of quasi-degenerate neutrino mass spectrum combined with $\rho^{(0)} \neq \sigma^{(0)}$) $\Delta\delta$ seems to have gained a contribution of $\geq \mathcal{O}(1)$ from $R_1 = 0.1$. For instance, as shown by the red line of the

$[\rho^{(0)}, \sigma^{(0)}]$ \diagdown $m_1[m_3]$		0.001 eV	0.01 eV	0.1 eV
[0, 0]	$R_1 = 0.1$	0.00 [-0.00]	0.00 [-0.00]	0.03 [-0.07]
$[\pi/2, 0]$		0.00 [-0.09]	0.01 [-0.09]	0.36 [-0.56]
$[0, \pi/2]$		0.00 [-0.09]	0.01 [-0.09]	0.43 [-0.70]
$[\pi/2, \pi/2]$		0.00 [-0.00]	0.00 [-0.00]	0.00 [-0.00]
[0, 0]	$I_1 = 0.1$	0.00 [0.00]	0.00 [0.00]	0.00 [0.01]
$[\pi/2, 0]$		0.00 [-0.00]	-0.00 [-0.00]	0.17 [-0.34]
$[0, \pi/2]$		0.00 [-0.00]	-0.00 [-0.00]	0.15 [-0.31]
$[\pi/2, \pi/2]$		0.00 [0.00]	0.00 [0.00]	0.00 [0.00]
[0, 0]	$R_2 = 0.1$	0.06 [-0.05]	0.08 [-0.08]	0.65 [-1.5]
$[\pi/2, 0]$		0.06 [-0.01]	0.07 [-0.02]	0.49 [-0.57]
$[0, \pi/2]$		0.04 [-0.01]	0.04 [-0.01]	0.13 [-0.14]
$[\pi/2, \pi/2]$		0.04 [-0.05]	0.03 [-0.03]	0.00 [-0.00]
[0, 0]	$I_2 = 0.1$	0.00 [-0.00]	0.00 [-0.00]	-0.00 [-0.00]
$[\pi/2, 0]$		0.00 [0.00]	0.00 [0.00]	-0.17 [0.28]
$[0, \pi/2]$		-0.00 [0.00]	-0.00 [0.00]	0.09 [-0.13]
$[\pi/2, \pi/2]$		-0.00 [-0.00]	-0.00 [-0.00]	-0.00 [-0.00]

Table 1. Some representative values of $\Delta\theta$ (induced by $R_1 = 0.1$, $I_1 = 0.1$, $R_2 = 0.1$ and $I_2 = 0.1$, respectively) at $m_1(m_3) = 0.001, 0.01$ and 0.1 eV in the NH (IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$.

sub-figure labelled by “NH +-” in figure 1, $\Delta\delta$ takes a value of 10 at $m_1 \simeq 0.077$ eV in the case of NH combined with $[\rho^{(0)}, \sigma^{(0)}] = [0, \pi/2]$. This would never be a realistic result but some signal for the breakdown of our approximation method used to derive eq. (2.13). Nevertheless, the ratio $\Delta\delta/R_1$ which takes a value of 100 in this specific case (implying that even a small $R_1 = 0.001$ can lead to some sizable $\Delta\delta \simeq 0.1 = 6^\circ$) is still useful. To keep $\Delta\delta$ within an acceptably small range in such a case (remember that it is right the possible closeness of θ_{23} and δ to $\pi/4$ and $-\pi/2$ that motivates us to study the μ - τ reflection symmetry), R_1 should be of $\leq \mathcal{O}(0.001)$. If R_1 were unfortunately much greater than 0.001, the resulting $\Delta\delta$ would be of $\mathcal{O}(1)$. And one has to invoke an exact method (instead of the approximation one adopted here) to obtain the precise value of $\Delta\delta$. But this is beyond the scope of our interest, because it would be insignificant to discuss the μ - τ reflection symmetry any more if δ turned out to be far from $\pm\pi/2$.

Given small values of symmetry-breaking parameters (i.e., $R_1 = 0.1$, $I_1 = 0.1$, $R_2 = 0.1$ and $I_2 = 0.1$), the resulting $\Delta\theta$ appears as:

1. The $|\Delta\theta|$ from $R_1 = 0.1$ is: negligibly small in the case of $m_1 \ll m_2 \ll m_3$ (e.g., the particular case of $m_1 = 0.001$ eV in table 1); negligibly small or 0.1 in the case of $m_3 \ll m_1 \simeq m_2$ (e.g., the particular case of $m_3 = 0.001$ eV in table 1) combined with

$[\rho^{(0)}, \sigma^{(0)}] \backslash m_1[m_3]$		0.001 eV	0.01 eV	0.1 eV
$[0, 0]$	$R_1 = 0.1$	-0.11 [-0.00]	-0.06 [-0.00]	-0.03 [-0.01]
$[\pi/2, 0]$		-0.13 [1.94]	-0.22 [1.98]	2.83 [11.2]
$[0, \pi/2]$		0.15 [1.95]	0.45 [2.07]	17.3 [11.7]
$[\pi/2, \pi/2]$		0.11 [-0.00]	0.04 [-0.00]	0.00 [-0.00]
$[0, 0]$	$I_1 = 0.1$	-0.05 [0.10]	-0.13 [0.14]	-1.70 [1.68]
$[\pi/2, 0]$		-0.02[4.92]	0.25 [5.18]	22.0 [31.5]
$[0, \pi/2]$		-0.07 [4.90]	0.09 [5.05]	21.6 [24.3]
$[\pi/2, \pi/2]$		-0.07 [0.09]	-0.06 [0.07]	-0.01 [0.01]
$[0, 0]$	$R_2 = 0.1$	0.02 [-0.01]	0.01 [-0.01]	0.06 [-0.23]
$[\pi/2, 0]$		0.02[1.27]	0.08 [1.47]	4.20 [12.4]
$[0, \pi/2]$		0.00 [1.24]	0.03 [1.16]	3.86 [3.16]
$[\pi/2, \pi/2]$		0.00 [-0.01]	-0.00 [-0.00]	0.00 [-0.00]
$[0, 0]$	$I_2 = 0.1$	-0.24 [-0.08]	-0.06 [-0.10]	0.80 [-0.89]
$[\pi/2, 0]$		-0.31[-1.22]	-0.84 [-1.91]	-21.1 [-25.9]
$[0, \pi/2]$		0.33 [-1.08]	0.75 [-0.49]	12.5 [9.86]
$[\pi/2, \pi/2]$		0.27 [-0.07]	0.11 [-0.05]	0.01 [-0.01]

Table 2. Some representative values of $\Delta\delta$ (induced by $R_1 = 0.1$, $I_1 = 0.1$, $R_2 = 0.1$ and $I_2 = 0.1$, respectively) at $m_1(m_3) = 0.001, 0.01$ and 0.1 eV in the NH (IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$.

$\rho^{(0)} = \sigma^{(0)}$ or $\rho^{(0)} \neq \sigma^{(0)}$; 0.03 or $\mathcal{O}(1)$ or negligibly small in the case of $m_1 \simeq m_2 \simeq m_3$ (e.g., the particular case of $m_1 = 0.1$ eV in table 1) combined with $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$ or $[\pi/2, 0]$ (and $[0, \pi/2]$) or $[\pi/2, \pi/2]$.

2. The $|\Delta\theta|$ from $I_{1,2} = 0.1$ is: around 0.1 in the case of $m_1 \simeq m_2 \simeq m_3$ combined with $\rho^{(0)} \neq \sigma^{(0)}$; negligibly small in other cases.
3. The $|\Delta\theta|$ from $R_2 = 0.1$ is: 0.05 in the case of $m_1 \ll m_2 \ll m_3$; 0.05 or negligibly small in the case of $m_3 \ll m_1 \simeq m_2$ combined with $\rho^{(0)} = \sigma^{(0)}$ or $\rho^{(0)} \neq \sigma^{(0)}$; $\mathcal{O}(1)$ or 0.1 or negligibly small in the case of $m_1 \simeq m_2 \simeq m_3$ combined with $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$ (and $[\pi/2, 0]$) or $[0, \pi/2]$ or $[\pi/2, \pi/2]$.

On the other hand, the resulting $\Delta\delta$ turns out to be:

1. The $|\Delta\delta|$ from $R_1 = 0.1$ is: around 0.1 in the case of $m_1 \ll m_2 \ll m_3$; negligibly small or $\mathcal{O}(1)$ in the case of $m_3 \ll m_1 \simeq m_2$ combined with $\rho^{(0)} = \sigma^{(0)}$ or $\rho^{(0)} \neq \sigma^{(0)}$; 0.03 or $\mathcal{O}(1)$ or negligibly small in the case of $m_1 \simeq m_2 \simeq m_3$ combined with $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$ or $[\pi/2, 0]$ (and $[0, \pi/2]$) or $[\pi/2, \pi/2]$.

$[\rho^{(0)}, \sigma^{(0)}]$ \diagdown $m_1[m_3]$		0.001 eV	0.01 eV	0.1 eV
[0, 0]	$R_1 = 0.1$	0.44 [0.01]	0.07 [0.01]	0.02 [0.04]
$[\pi/2, 0]$		-0.36 [-1.25]	0.02 [-1.43]	-3.24 [-6.25]
$[0, \pi/2]$		0.22 [-1.77]	-0.19 [-1.72]	-9.14 [-12.4]
$[\pi/2, \pi/2]$		-0.33 [0.01]	-0.04 [0.01]	-0.00 [0.00]
[0, 0]	$I_1 = 0.1$	-0.12 [-0.12]	-0.01 [-0.02]	0.04 [-0.05]
$[\pi/2, 0]$		0.07 [-1.27]	-0.19 [-3.38]	-15.7 [-20.7]
$[0, \pi/2]$		0.26 [-5.57]	-0.08 [-3.75]	-14.1 [-19.0]
$[\pi/2, \pi/2]$		-0.24 [0.11]	-0.01 [0.01]	-0.00 [0.00]
[0, 0]	$R_2 = 0.1$	-0.24 [0.01]	-0.05 [0.01]	-0.16 [0.34]
$[\pi/2, 0]$		0.15 [0.02]	-0.10 [-0.80]	-4.27 [-6.70]
$[0, \pi/2]$		-0.13 [-1.61]	-0.00 [-0.85]	-1.96 [-3.21]
$[\pi/2, \pi/2]$		0.15 [-0.00]	0.01 [0.00]	0.00 [0.00]
[0, 0]	$I_2 = 0.1$	0.58 [-1.17]	0.11 [-0.11]	-0.01 [0.02]
$[\pi/2, 0]$		-0.18 [-0.14]	0.55 [1.23]	15.0 [17.0]
$[0, \pi/2]$		0.10 [1.75]	-0.48 [0.45]	-8.07 [-7.62]
$[\pi/2, \pi/2]$		-0.50 [1.17]	-0.11 [0.12]	-0.01 [0.01]

Table 3. Some representative values of $\Delta\rho$ (induced by $R_1 = 0.1$, $I_1 = 0.1$, $R_2 = 0.1$ and $I_2 = 0.1$, respectively) at $m_1(m_3) = 0.001, 0.01$ and 0.1 eV in the NH (IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$.

2. The $|\Delta\delta|$ from $I_{1,2} = 0.1$ is: $\mathcal{O}(0.1)$ in the case of $m_1 \ll m_2 \ll m_3$; 0.1 or $\mathcal{O}(1)$ in the case of $m_3 \ll m_1 \simeq m_2$ combined with $\rho^{(0)} = \sigma^{(0)}$ or $\rho^{(0)} \neq \sigma^{(0)}$; $\mathcal{O}(1)$ in the case of $m_1 \simeq m_2 \simeq m_3$ but negligibly small for $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$.
3. The $|\Delta\delta|$ from $R_2 = 0.1$ is: negligibly small in the case of $m_1 \ll m_2 \ll m_3$; negligibly small or $\mathcal{O}(1)$ in the case of $m_3 \ll m_1 \simeq m_2$ combined with $\rho^{(0)} = \sigma^{(0)}$ or $\rho^{(0)} \neq \sigma^{(0)}$; 0.06 or $\mathcal{O}(1)$ or negligibly small in the case of $m_1 \simeq m_2 \simeq m_3$ combined with $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$ or $[\pi/2, 0]$ (and $[0, \pi/2]$) or $[\pi/2, \pi/2]$.

Note that in the cases where some mixing-parameter deviation acquires a size of $\mathcal{O}(1)$ from $R_1 = 0.1$, $I_1 = 0.1$, $R_2 = 0.1$ or $I_2 = 0.1$, the corresponding symmetry-breaking parameter should be of $\leq \mathcal{O}(0.01)$ to keep this mixing-parameter deviation within an acceptably small range. From the above results one can draw the following observations: (1) The mixing-parameter deviations tend to get remarkably enhanced (or suppressed) in the case of $m_1 \simeq m_2 \simeq m_3$ combined with $\rho^{(0)} \neq \sigma^{(0)}$ (or $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$). (2) When the absolute neutrino mass scale is small and $\rho^{(0)} \neq \sigma^{(0)}$, the mixing-parameter deviations generally have a much greater magnitude in the IH case than in the NH case. (3) $\Delta\theta$ is

$[\rho^{(0)}, \sigma^{(0)}]$ \diagdown $m_1[m_3]$		0.001 eV	0.01 eV	0.1 eV
[0, 0]	$R_1 = 0.1$	-0.02 [-0.00]	-0.02 [-0.01]	-0.01 [-0.02]
$[\pi/2, 0]$		-0.02 [-0.43]	0.01 [-0.63]	-1.33 [-2.88]
$[0, \pi/2]$		0.01 [-0.94]	-0.06 [-0.79]	-4.07 [-5.63]
$[\pi/2, \pi/2]$		0.02 [-0.01]	0.01 [-0.00]	0.00 [-0.00]
[0, 0]	$I_1 = 0.1$	0.01 [-0.12]	0.00 [-0.02]	0.04 [-0.04]
$[\pi/2, 0]$		0.00 [0.65]	-0.06 [-1.35]	-6.83 [-9.09]
$[0, \pi/2]$		0.01 [-3.65]	-0.03 [-1.77]	-6.32 [-8.58]
$[\pi/2, \pi/2]$		0.01 [0.11]	0.00 [0.01]	-0.00 [0.00]
[0, 0]	$R_2 = 0.1$	0.02 [0.00]	0.02 [-0.00]	0.07 [-0.14]
$[\pi/2, 0]$		0.01 [0.46]	-0.03 [-0.32]	-1.74 [-3.08]
$[0, \pi/2]$		-0.01 [-1.18]	-0.00 [-0.43]	-0.88 [-1.47]
$[\pi/2, \pi/2]$		-0.01 [-0.01]	-0.01 [-0.00]	-0.00 [-0.00]
[0, 0]	$I_2 = 0.1$	0.17 [-1.17]	0.09 [-0.11]	-0.01 [0.02]
$[\pi/2, 0]$		0.19 [-0.64]	0.32 [0.47]	6.57 [7.50]
$[0, \pi/2]$		-0.20 [1.31]	-0.30 [0.23]	-3.66 [-3.48]
$[\pi/2, \pi/2]$		-0.18 [1.17]	-0.09 [0.12]	-0.01 [0.01]

Table 4. Some representative values of $\Delta\sigma$ (induced by $R_1 = 0.1$, $I_1 = 0.1$, $R_2 = 0.1$ and $I_2 = 0.1$, respectively) at $m_1(m_3) = 0.001, 0.01$ and 0.1 eV in the NH (IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$.

most sensitive to R_2 while $\Delta\delta$ to all the symmetry-breaking parameters. In magnitude, $\Delta\delta$ is generally much greater than $\Delta\theta$.

For illustration, we give a toy example to show how to make use of the above results. In this connection, we discuss how the global-fit results $\theta_{23}/^\circ = 41.6_{-1.2}^{+1.5}$ and $\delta/^\circ = 261_{-59}^{+51}$ in the NH case [5] may arise from an approximate μ - τ reflection symmetry. (For simplicity, only the best-fit results will be used.) An interesting possibility is that a single R_2 ³ is responsible for the two measured mixing-parameter deviations $\Delta\theta \simeq -3.4^\circ$ and $\Delta\delta \simeq -9^\circ$. This can be achieved when the $\Delta\delta$ and $\Delta\theta$ generated by R_2 happen to fulfill the condition $\Delta\delta/\Delta\theta = (-9^\circ)/(-3.4^\circ) \simeq 2.65$. With the aid of the sub-figure labelled by “NH $-+$ ” in figure 3, it turns out that this requirement can be satisfied in the case of $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, 0]$: at $m_1 = 0.019$ eV, the values of $\Delta\theta/R_2$ and $\Delta\delta/R_2$ (as well as $\Delta\rho/R_2$ and $\Delta\sigma/R_2$) respectively read 7.8 and 20.7 (as well as -24.5 and -9). Accordingly, $R_2 \simeq -0.076$ is capable of producing $\theta_{23} \simeq 41.6^\circ$ and $\delta = 261^\circ$ (as well as $\rho \simeq 100.7^\circ$ and $\sigma \simeq 3.9^\circ$).

³Of course, in a realistic context, the mixing-parameter deviations may receive contributions from not merely one symmetry-breaking parameter.

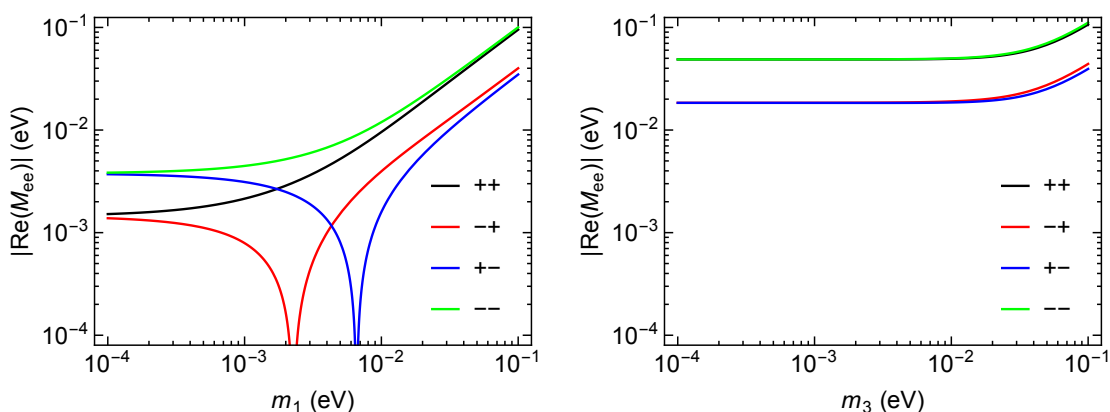


Figure 5. The possible values of $|\text{Re}(M_{ee})|$ against the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$. The signs $++$, $-+$, $+-$ and $--$ respectively stand for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$, $[\pi/2, 0]$, $[0, \pi/2]$ and $[\pi/2, \pi/2]$.

Finally, we discuss the consequences of breaking of μ - τ reflection symmetry on the allowed range of effective Majorana neutrino mass $|M_{ee}|$ which directly controls the rates of neutrinoless double-beta decays [36–39]. For this purpose, one obtains

$$\begin{aligned} \text{Re}(M_{ee}) &\simeq \bar{m}_1 c_{12}^2 c_{13}^2 + \bar{m}_2 s_{12}^2 c_{13}^2 - m_3 s_{13}^2, \\ \text{Im}(M_{ee}) &\simeq \bar{m}_1 c_{12}^2 \Delta\rho + \bar{m}_2 s_{12}^2 \Delta\sigma + m_3 s_{13}^2 \Delta\delta + (m_{11} - m_3 s_{13}^2) \Delta\phi_1. \end{aligned} \quad (2.19)$$

Because the symmetry-breaking parameter $\epsilon_3 = \text{Im}(M_{ee})/\text{Re}(M_{ee})$ should be a small quantity (e.g., $|\epsilon_3| \leq 0.1$) if we want to maintain the μ - τ reflection symmetry as an approximate one, the value of

$$|M_{ee}| = \sqrt{[\text{Re}(M_{ee})]^2 + [\text{Im}(M_{ee})]^2} \simeq |\text{Re}(M_{ee})| \left(1 + \frac{\epsilon_3^2}{2}\right), \quad (2.20)$$

can be well approximated by that of $|\text{Re}(M_{ee})|$. It is thus fair to say that the consequences of breaking of μ - τ reflection symmetry on the allowed range of $|M_{ee}|$ are negligibly small. In figure 5 we present the possible values of $|\text{Re}(M_{ee})|$ as a function of the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ [23]. (1) In the NH case, the three components of $|\text{Re}(M_{ee})|$ add constructively to a maximal level for $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$. By contrast, the three components will cancel each other out (i.e., $|\text{Re}(M_{ee})| \simeq 0$) at $m_1 \simeq 0.002$ eV (or 0.007 eV) for $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, 0]$ (or $[0, \pi/2]$). (2) In the IH case, the value of $|\text{Re}(M_{ee})|$ is mainly determined by the first two components as the third one is highly suppressed. Because of $m_1 \simeq m_2$ in the IH case, $|\text{Re}(M_{ee})|$ approximates to m_1 (or $m_1(c_{12}^2 - s_{12}^2)$) for $\rho^{(0)} = \sigma^{(0)}$ (or $\rho^{(0)} \neq \sigma^{(0)}$).

2.2 Analytical approximations

In this section, we give the analytical expressions of $\Delta\delta$ and $\Delta\theta$ (which are of more practical interest than $\Delta\rho$ and $\Delta\sigma$ since the Majorana phases cannot be pinned down in a foreseeable

future) to explain the numerical results. After a lengthy but straightforward calculation, one obtains the approximation results

$$\begin{aligned}
 \Delta\delta &\simeq [2\mathbb{T}(\bar{m}_1 + \bar{m}_2)(\bar{m}_1 - m_3)(\bar{m}_2 - m_3)\bar{s}_{13}]^{-1} \{ [2\mathbb{T}m_{12}(m_{22} - m_3)(\bar{m}_1 + \bar{m}_2 - 2m_3s_{13}^2) \\
 &\quad - 4m_{12}^2(m_{11} + m_3 + \mathbb{T}m_{12})s_{13}^2 + 4(\bar{m}_1 - m_3)(\bar{m}_2 - m_3)(m_{11} + m_3)s_{13}^2] \mathbb{R}_1 - [4m_{12} \\
 &\quad \times (\bar{m}_1 - m_3)(\bar{m}_2 - m_3) - 2\mathbb{T}(\bar{m}_1 + \bar{m}_2 - 2m_3s_{13}^2)(m_{11} + m_3)(m_{22} - m_3) - 4\mathbb{T}\Omega \\
 &\quad \times m_{12}^2m_3s_{13}^2] \bar{s}_{13}\mathbb{I}_1 - [\mathbb{T}m_{12}(m_{22} + m_3)(\bar{m}_1 + \bar{m}_2 - 2m_3s_{13}^2) - 2\mathbb{T}m_{12}(\bar{m}_1 - m_3) \\
 &\quad \times (\bar{m}_2 - m_3)s_{13}^2 - 2(m_{11} - m_3)(m_{22} + m_3)(m_{11} + m_3 + \mathbb{T}m_{12})s_{13}^2] \mathbb{R}_2 \\
 &\quad + \mathbb{T}(m_{22} + m_3)[m_{12}^2 - (m_{22} - m_3)(2m_{11} + m_{22} - m_3)] \bar{s}_{13}\mathbb{I}_2 \} , \\
 \Delta\theta &\simeq [2\mathbb{T}(\bar{m}_1 + \bar{m}_2)(\bar{m}_1 - m_3)(\bar{m}_2 - m_3)]^{-1} \{ 2\mathbb{T}(\bar{m}_1 + \bar{m}_2)[m_{12}^2 - (m_{11}^2 - m_3^2)s_{13}^2] \mathbb{R}_1 \\
 &\quad + 4m_{11}m_{12}[\mathbb{T}(\bar{m}_1 + \bar{m}_2) + 2m_{12}s_{13}^2] \bar{s}_{13}\mathbb{I}_1 - \mathbb{T}(\bar{m}_1 + \bar{m}_2)(m_{11} - m_3)(m_{22} + m_3)\mathbb{R}_2 \\
 &\quad - \mathbb{T}m_{12}[(\bar{m}_1 + \bar{m}_2)(2m_{11} + m_{22} - m_3) - 2m_{11}(m_{22} + m_3)s_{13}^2] \bar{s}_{13}\mathbb{I}_2 \} , \quad (2.21)
 \end{aligned}$$

with $\mathbb{T} = \tan 2\theta_{12}$ and $\Omega = (m_{11} + m_3)/(m_{11} - m_3)$.

For illustration, we discuss the possible values of $\Delta\delta$ and $\Delta\theta$ in several typical cases.

(1) In the case of $m_1 \ll m_2 \simeq \sqrt{\Delta m_{21}^2} \ll m_3 \simeq \sqrt{|\Delta m_{31}^2|}$, $\Delta\delta$ and $\Delta\theta$ approximate to

$$\begin{aligned}
 \Delta\delta &\simeq \frac{m_3}{2\bar{m}_2\bar{s}_{13}} \left[\left(2rc_{12}s_{12} + \frac{4}{\mathbb{T}}s_{13}^2 \right) \mathbb{R}_1 - \bar{s}_{13}\mathbb{I}_2 + \left(rc_{12}s_{12} - \frac{2}{\mathbb{T}}s_{13}^2 \right) \mathbb{R}_2 \right] - 2s_{12}^2\mathbb{I}_1 , \\
 \Delta\theta &\simeq s_{13}^2\mathbb{R}_1 - 2|r|c_{12}s_{12}^3\bar{s}_{13}\mathbb{I}_1 + \frac{1}{2}\mathbb{R}_2 - \frac{\bar{m}_2}{2m_3}c_{12}s_{12}\bar{s}_{13}\mathbb{I}_2 , \quad (2.22)
 \end{aligned}$$

where $r = \Delta m_{21}^2/\Delta m_{31}^2$ takes a value of ± 0.03 for NH or IH. One finds that $\Delta\delta$ ($\Delta\theta$) can receive modest contributions from \mathbb{R}_1 and $\mathbb{I}_{1,2}$ (\mathbb{R}_2). (2) For $m_1 \simeq m_2 \simeq \sqrt{|\Delta m_{31}^2|} \gg m_3$, in which case one has $m_2 - m_1 \simeq \Delta m_{21}^2/(2\sqrt{|\Delta m_{31}^2|})$, the results are strongly dependent on the combinations of $\rho^{(0)}$ and $\sigma^{(0)}$. If $\rho^{(0)}$ and $\sigma^{(0)}$ take the same value, $\Delta\delta$ and $\Delta\theta$ are simplified to

$$\begin{aligned}
 \Delta\delta &\simeq \mathbb{I}_1 - \frac{3}{4}\mathbb{I}_2 + \frac{1}{4\mathbb{T}\bar{s}_{13}} [(4s_{13}^2 - 2\mathbb{T}|r|c_{12}s_{12}) \mathbb{R}_1 + (2s_{13}^2 + \mathbb{T}|r|c_{12}s_{12}) \mathbb{R}_2] , \\
 \Delta\theta &\simeq -s_{13}^2\mathbb{R}_1 - |r|c_{12}s_{12}\bar{s}_{13} \left(\mathbb{I}_1 - \frac{3}{4}\mathbb{I}_2 \right) - \frac{1}{2}\mathbb{R}_2 . \quad (2.23)
 \end{aligned}$$

It turns out that $\Delta\delta$ ($\Delta\theta$) can acquire modest contributions from \mathbb{I}_1 and \mathbb{I}_2 (\mathbb{R}_2). When $\rho^{(0)}$ and $\sigma^{(0)}$ differ from each other, one will have

$$\begin{aligned}
 \Delta\delta &\simeq \frac{1}{|r|} [-32c_{12}^3s_{12}^3\bar{s}_{13}\mathbb{R}_1 + 4\cos 2\theta_{12}\mathbb{I}_1 - 4c_{12}s_{12}\bar{s}_{13}\mathbb{R}_2 - 4c_{12}^2s_{12}^2\cos 2\theta_{12}\mathbb{I}_2] , \\
 \Delta\theta &\simeq -4c_{12}^2s_{12}^2\mathbb{R}_1 - c_{12}s_{12}\cos 2\theta_{12} \left[1 - \frac{4}{|r|}\cos 2\theta_{12}s_{13}^2 \right] \bar{s}_{13}(4\mathbb{I}_1 - \mathbb{I}_2) - \frac{1}{2}\cos^2 2\theta_{12}\mathbb{R}_2 . \quad (2.24)
 \end{aligned}$$

Enhanced by the factor $1/|r|$, $\Delta\delta$ is easy to gain a large contribution. On the other hand, $\Delta\theta$ can only obtain a modest contribution from \mathbb{R}_1 . (3) When it comes to $m_1 \simeq m_2 \simeq m_3 \simeq m_0$, in which case $m_2 - m_1 \simeq \Delta m_{21}^2/(2m_0)$ and $m_3 - m_1 \simeq \Delta m_{31}^2/(2m_0)$, the mixing-parameter

deviations may get remarkably magnified in some cases. For $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$, $\Delta\delta$ and $\Delta\theta$ are approximately given by

$$\begin{aligned}\Delta\delta &\simeq \frac{1}{T\bar{s}_{13}} \left[(\text{Tr}c_{12}s_{12} + 2s_{13}^2) R_1 + \frac{2m_0^2}{\Delta m_{31}^2} (\text{Tr}c_{12}s_{12} - 2s_{13}^2) R_2 \right] - \frac{2m_0^2}{\Delta m_{31}^2} (2I_1 - I_2) , \\ \Delta\theta &\simeq \frac{2m_0^2}{\Delta m_{31}^2} [2s_{13}^2 R_1 - rc_{12}s_{12}\bar{s}_{13} (2I_1 - I_2) + R_2] .\end{aligned}\quad (2.25)$$

Because of the enhancement factor $m_0^2/\Delta m_{31}^2$, $\Delta\delta$ ($\Delta\theta$) can get large contributions from I_1 and I_2 (R_2). In the case of $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, 0]$, $\Delta\delta$ and $\Delta\theta$ appear as

$$\begin{aligned}\Delta\delta &\simeq \frac{2m_0^2}{\Delta m_{21}^2} \left\{ \left[\frac{4}{T}s_{12}^2 - 2c_{12}s_{12} (1 + 2s_{12}^2) \right] \bar{s}_{13}R_1 + 2\cos 2\theta_{12}I_1 - 2c_{12}s_{12}\bar{s}_{13}R_2 - c_{12}^2I_2 \right\} , \\ \Delta\theta &\simeq \frac{2m_0^2}{\Delta m_{31}^2} \left[2c_{12}^2s_{12}^2R_1 - \frac{4m_0^2}{\Delta m_{21}^2} c_{12}s_{12} \cos 2\theta_{12}\bar{s}_{13}^3 (2\cos 2\theta_{12}I_1 - c_{12}^2I_2) + c_{12}^4R_2 \right] .\end{aligned}\quad (2.26)$$

When $[\rho^{(0)}, \sigma^{(0)}] = [0, \pi/2]$, the results become

$$\begin{aligned}\Delta\delta &\simeq \frac{2m_0^2}{\Delta m_{21}^2} \left\{ - \left[\frac{4}{T}c_{12}^2 + 2c_{12}s_{12} (1 + 2c_{12}^2) \right] \bar{s}_{13}R_1 + 2\cos 2\theta_{12}I_1 - 2c_{12}s_{12}\bar{s}_{13}R_2 + s_{12}^2I_2 \right\} , \\ \Delta\theta &\simeq \frac{2m_0^2}{\Delta m_{31}^2} \left[2c_{12}^2s_{12}^2R_1 - \frac{4m_0^2}{\Delta m_{21}^2} c_{12}s_{12} \cos 2\theta_{12}\bar{s}_{13}^3 (2\cos 2\theta_{12}I_1 + s_{12}^2I_2) + s_{12}^4R_2 \right] .\end{aligned}\quad (2.27)$$

In these two cases $\Delta\delta$ is greatly enhanced by the factor $m_0^2/\Delta m_{21}^2$, while the enhancement factor $m_0^2/\Delta m_{31}^2$ for $\Delta\theta$ is also significant but not so huge. Finally, $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$ will lead us to

$$\begin{aligned}\Delta\delta &\simeq -\frac{\Delta m_{31}^2}{4m_0^2\bar{s}_{13}} \left[\left(rc_{12}s_{12} + \frac{2}{T}s_{13}^2 \right) R_1 + \bar{s}_{13} (I_1 - I_2) \right] , \\ \Delta\theta &\simeq \frac{\Delta m_{31}^2}{8m_0^2} [2s_{13}^2 R_1 - 2rc_{12}s_{12}\bar{s}_{13} (I_1 - I_2) + R_2] .\end{aligned}\quad (2.28)$$

It is easy to see that all the mixing-parameter deviations are extremely suppressed in this case. One will find that all the above analytical results agree well with the corresponding numerical results.

2.3 RG induced symmetry breaking

This section is devoted to the RG-induced breaking of μ - τ reflection symmetry. A flavor symmetry [6, 7] together with the associated new fields is usually introduced at an energy scale Λ_{FS} much higher than the electroweak (EW) one Λ_{EW} . In this case one must consider the RG running effect when confronting the flavor-symmetry model with the low-energy data [40]. During the RG evolution process the significant difference between m_μ and m_τ can serve as a unique source for the breaking of μ - τ reflection symmetry. The general symmetry breaking studied in the above thus finds an interesting application in such a specific situation [41–44]. The energy dependence of neutrino mass matrix is described by its RG equation, which at the one-loop level appears as [45–49]

$$16\pi^2 \frac{dM_\nu}{dt} = C \left(Y_l^\dagger Y_l \right)^T M_\nu + C M_\nu \left(Y_l^\dagger Y_l \right) + \alpha M_\nu .\quad (2.29)$$

Here t is defined as $\ln(\mu/\mu_0)$ with μ denoting the renormalization scale, whereas C and α read

$$\begin{aligned} C &= -\frac{3}{2}, & \alpha &\simeq -3g_2^2 + 6y_t^2 + \lambda, & \text{in the SM;} \\ C &= 1, & \alpha &\simeq -\frac{6}{5}g_1^2 - 6g_2^2 + 6y_t^2, & \text{in the MSSM.} \end{aligned} \quad (2.30)$$

In eq. (2.29) the α -term is flavor universal and therefore just contributes an overall rescaling factor (which will be referred to as I_α), while the other two terms may modify the structure of M_ν . In the basis under study, the Yukawa coupling matrix of three charged leptons is given by $Y_l = \text{Diag}(y_e, y_\mu, y_\tau)$. In light of $y_e \ll y_\mu \ll y_\tau$, it is reasonable to neglect the contributions of y_e and y_μ . Integration of the RG equation enables us to connect the neutrino mass matrix $M_\nu(\Lambda_{\text{FS}})$ at Λ_{FS} with the corresponding one at Λ_{EW} in a manner as [50, 51]

$$M_\nu(\Lambda_{\text{EW}}) = I_\alpha I_\tau^\dagger M_\nu(\Lambda_{\text{FS}}) I_\tau^*, \quad (2.31)$$

where $I_\tau \simeq \text{Diag}\{1, 1, 1 - \Delta_\tau\}$ and

$$I_\alpha = \exp\left(\frac{1}{16\pi^2} \int_{\ln \Lambda_{\text{FS}}}^{\ln \Lambda_{\text{EW}}} \alpha dt\right), \quad \Delta_\tau = \frac{C}{16\pi^2} \int_{\ln \Lambda_{\text{EW}}}^{\ln \Lambda_{\text{FS}}} y_\tau^2 dt. \quad (2.32)$$

In the SM case, the RG running effect is negligible due to the smallness of $y_\tau \simeq 0.01$ (which renders $\Delta_\tau \simeq \mathcal{O}(10^{-5})$). By contrast, $y_\tau^2 = (1 + \tan^2 \beta)m_\tau^2/v^2$ can be enhanced by a large $\tan \beta$ in the MSSM case. Given $\Lambda_{\text{FS}} \simeq 10^{13}$ GeV, for example, the value of Δ_τ depends on $\tan \beta$ in a way as

$$\Delta_\tau \simeq 0.042 \left(\frac{\tan \beta}{50}\right)^2. \quad (2.33)$$

With the help of eq. (2.31), one will get the RG-corrected neutrino mass matrix at Λ_{EW}

$$M_\nu(\Lambda_{\text{EW}}) \simeq I_\alpha \left[M_\nu(\Lambda_{\text{FS}}) - \Delta_\tau \begin{pmatrix} 0 & 0 & M_{e\tau} \\ 0 & 0 & M_{\mu\tau} \\ M_{e\tau} & M_{\mu\tau} & 2M_{\tau\tau} \end{pmatrix} \right], \quad (2.34)$$

from a neutrino mass matrix respecting the μ - τ reflection symmetry at Λ_{FS} . By means of the above-mentioned treatment, one may arrange $M_\nu(\Lambda_{\text{EW}})$ in a form as given by eq. (2.6) with $\epsilon_2 = 2\epsilon_1 = \Delta_\tau$ and $\epsilon_{3,4} = 0$. The relations between the mixing-parameter deviations and Δ_τ can therefore be obtained by simply taking $R_2 = 2R_1 = \Delta_\tau$ and $I_{1,2} = 0$ in eq. (2.13). As before, the mixing-parameter deviations have a linear dependence on Δ_τ and thus a square dependence on $\tan \beta$. In figures 6–7 we display the mixing-parameter deviations respectively arising from $\tan \beta = 50$ and $\tan \beta = 30$ against the absolute neutrino mass scale for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$. By analogy with eq. (2.18) (with a square dependence in place of the linear dependence), the mixing-parameter deviations generated by other values of $\tan \beta$ can also be inferred from these results. Qualitatively, the mixing-parameter deviations produced by Δ_τ closely resemble those resulting from $R_{1,2}$ in a few aspects: (1) Their magnitudes tend to grow with the absolute neutrino mass

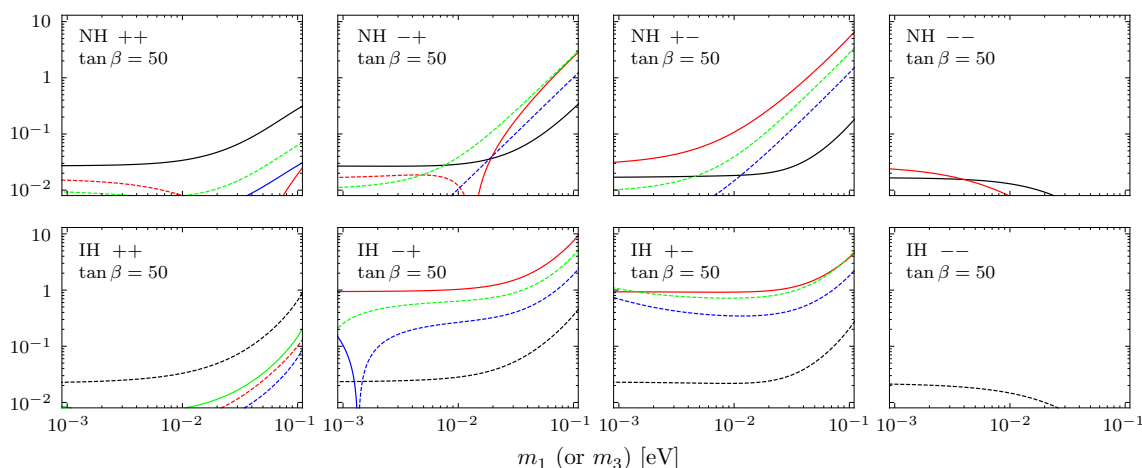


Figure 6. The mixing-parameter deviations in the MSSM case with $\tan \beta = 50$ ($\Delta\theta$ in black, $\Delta\delta$ in red, $\Delta\rho$ in green and $\Delta\sigma$ in blue) against the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$. The signs ++, -+, +- and -- respectively stand for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0], [\pi/2, 0], [0, \pi/2]$ and $[\pi/2, \pi/2]$.

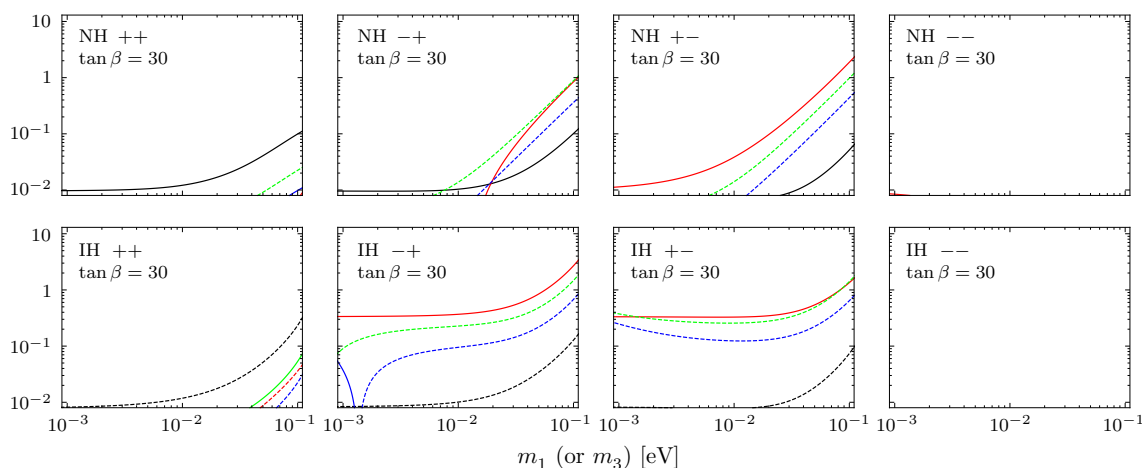


Figure 7. The mixing-parameter deviations in the MSSM case with $\tan \beta = 30$ ($\Delta\theta$ in black, $\Delta\delta$ in red, $\Delta\rho$ in green and $\Delta\sigma$ in blue) against the lightest neutrino mass m_1 (or m_3) in the NH (or IH) case for various combinations of $\rho^{(0)}$ and $\sigma^{(0)}$ with $\delta^{(0)} = -\pi/2$. The signs ++, -+, +- and -- respectively stand for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0], [\pi/2, 0], [0, \pi/2]$ and $[\pi/2, \pi/2]$.

scale (except in the case of $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$). (2) $\Delta\theta$ generally receives a modest contribution from Δ_τ in most cases, while $\Delta\delta$, $\Delta\rho$ and $\Delta\sigma$ may get remarkably magnified for $\rho^{(0)} \neq \sigma^{(0)}$. (3) $|\Delta\delta|$ and $|\Delta\rho|$ are comparable to each other, while $|\Delta\sigma|$ is somewhat smaller. (4) $\Delta\theta$ is always positive (negative) in the NH (IH) case. Quantitatively, the allowed $\Delta\theta$ and $\Delta\delta$ from $\Delta_\tau \leq 0.04$ (for $\tan \beta \leq 50$) are as follows: (1) In the case of $m_1 \ll m_2 \ll m_3$ or $m_3 \ll m_1 \simeq m_2$, $|\Delta\theta|$ is smaller than 0.02. When the neutrino masses are quasi-degenerate (except in the case of $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$), $|\Delta\theta|$ may receive a contribution up to 0.3, implying that $\Delta_\tau \simeq 0.02$ (corresponding to $\tan \beta \simeq 35$) is large

enough to generate $|\Delta\theta| \simeq 0.1$. (2) $|\Delta\delta|$ is also smaller than 0.02 for $m_1 \ll m_2 \ll m_3$. In the case of $m_3 \ll m_1 \simeq m_2$ or $m_1 \simeq m_2 \simeq m_3$, $|\Delta\delta|$ will be greatly enhanced to $\mathcal{O}(1)$ for $\rho^{(0)} \neq \sigma^{(0)}$ (even Δ_τ of $\mathcal{O}(0.001)$ that corresponds to $\tan\beta \simeq 10$ is sufficient for producing $|\Delta\delta|$ of $\mathcal{O}(0.1)$) but negligibly small for $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$.

On the other hand, the analytical expressions for $\Delta\theta$, $\Delta\rho$ and Δ are found to be

$$\begin{aligned}
 \Delta\theta &\simeq [2(\bar{m}_1 - m_3)(\bar{m}_2 - m_3)]^{-1} [m_{12}^2 - (m_{11} - m_3)(m_{22} + m_3)] \Delta_\tau, \\
 \Delta\delta &\simeq -[\text{T}(\bar{m}_1 + \bar{m}_2)(\bar{m}_1 - m_3)(\bar{m}_2 - m_3)\bar{s}_{13}]^{-1} [\text{T}m_{12}m_3(\bar{m}_1 + \bar{m}_2 - 2m_3s_{13}^2) \\
 &\quad - 2(m_{11} + \text{T}m_{12} + m_3)(\bar{m}_1\bar{m}_2 - m_{22}m_3)s_{13}^2] \Delta_\tau, \\
 \Delta\rho &\simeq -[2\bar{m}_1m_3(\bar{m}_1 + \bar{m}_2)(\bar{m}_1 - m_3)(\bar{m}_2 - m_3)\bar{s}_{13}t_{12}]^{-1} \{ \text{T}m_{12}m_3^2(\bar{m}_1 + \bar{m}_2)(m_{11} \\
 &\quad - m_3s_{13}^2 + 2m_{22}t_{12}^2s_{13}^2) - m_{11}[\text{T}m_{12}m_3(\bar{m}_1 + \bar{m}_2 - 2m_3s_{13}^2) - 2(m_{11} + \text{T}m_{12} \\
 &\quad + m_3)(\bar{m}_1\bar{m}_2 - m_{22}m_3)s_{13}^2] [m_3 + m_{11}\Omega s_{13}^2 - m_{22}t_{12}^2s_{13}^2] \} \Delta_\tau, \tag{2.35}
 \end{aligned}$$

with $t_{12} = \tan\theta_{12}$, while $\Delta\sigma$ can be obtained from $\bar{m}_1\Delta\rho/\bar{m}_2$ by making the replacement $t_{12} \rightarrow -1/t_{12}$. These results can help us understand the numerical results: (1) For $m_1 \ll m_2 \ll m_3$, eq. (2.35) is simplified to

$$\Delta\theta \simeq \frac{1}{2}\Delta_\tau, \quad \Delta\delta \simeq \frac{\bar{m}_2c_{12}s_{12}}{m_3\bar{s}_{13}}\Delta_\tau, \quad \Delta\rho \simeq \frac{c_{12}\bar{s}_{13}}{s_{12}}\Delta_\tau. \tag{2.36}$$

(2) In the case of $m_1 \simeq m_2 \gg m_3$, one will have

$$\Delta\delta \simeq \frac{1}{\text{T}}\bar{s}_{13}\Delta_\tau, \quad \Delta\rho \simeq -\frac{[m_3 + \bar{m}_1(1 - t_{12}^2)s_{13}^2]\bar{s}_{13}}{2m_3t_{12}}\Delta_\tau, \tag{2.37}$$

for $\rho^{(0)} = \sigma^{(0)}$, or

$$\Delta\delta \simeq -\frac{8}{|r|}c_{12}s_{12}\bar{s}_{13}\Delta_\tau, \quad \Delta\rho \simeq \frac{8[m_3 + \bar{m}_1(c_{12}^2 - s_{12}^2)(1 + t_{12}^2)s_{13}^2]c_{12}^3s_{12}\bar{s}_{13}}{m_3|r|}\Delta_\tau, \tag{2.38}$$

for $\rho^{(0)} \neq \sigma^{(0)}$ together with $\Delta\theta \simeq -1/2\Delta_\tau$. (3) When the case of $m_1 \simeq m_2 \simeq m_3 \simeq m_0$ is considered, the results become

$$\begin{aligned}
 \Delta\theta &\simeq \frac{2m_0^2}{\Delta m_{31}^2}\Delta_\tau, & \Delta\delta &\simeq \frac{2m_0^2[\text{T}rc_{12}s_{12} - 2s_{13}^2]}{\text{T}\Delta m_{31}^2\bar{s}_{13}}\Delta_\tau, & \Delta\rho &\simeq \frac{\text{T}m_0^2rc_{12}^2}{\Delta m_{31}^2\bar{s}_{13}}\Delta_\tau; \\
 \Delta\theta &\simeq \frac{2m_0^2c_{12}^2}{\Delta m_{31}^2}\Delta_\tau, & \Delta\delta &\simeq \frac{2m_0^2(r - 4s_{13}^2)c_{12}s_{12}}{\Delta m_{21}^2\bar{s}_{13}}\Delta_\tau, & \Delta\rho &\simeq \frac{8m_0^2c_{12}^3s_{12}\bar{s}_{13}}{\Delta m_{21}^2}\Delta_\tau; \\
 \Delta\theta &\simeq \frac{2m_0^2s_{12}^2}{\Delta m_{31}^2}\Delta_\tau, & \Delta\delta &\simeq -\frac{2m_0^2(r + 4s_{13}^2)c_{12}s_{12}}{\Delta m_{21}^2\bar{s}_{13}}\Delta_\tau, & \Delta\rho &\simeq \frac{8m_0^2c_{12}^3s_{12}\bar{s}_{13}}{\Delta m_{21}^2}\Delta_\tau; \\
 \Delta\theta &\simeq \frac{\Delta m_{31}^2}{8m_0^2}\Delta_\tau, & \Delta\delta &\simeq -\frac{\text{T}\Delta m_{21}^2c_{12}s_{12} + 2\Delta m_{31}^2s_{13}^2}{8\text{T}m_0^2\bar{s}_{13}}\Delta_\tau, & \Delta\rho &\simeq \frac{\Delta m_{31}^2\bar{s}_{13}}{8m_0^2t_{12}}\Delta_\tau, \tag{2.39}
 \end{aligned}$$

for $[\rho^{(0)}, \sigma^{(0)}] = [0, 0]$, $[\pi/2, 0]$, $[0, \pi/2]$ and $[\pi/2, \pi/2]$. One can see that these approximation results agree well with the corresponding numerical results.

3 Summary

To summarize, the μ - τ reflection symmetry deserves particular attention as it leads to the interesting results $\theta_{23} = \pi/4$ and $\delta = \pm\pi/2$ (which are close to the current experimental data) as well as trivial Majorana phases. Nevertheless, it is reasonable for us to consider the breaking of such a symmetry either from the theoretical considerations (e.g., the RG running effect may provide a source for the symmetry breaking) or on the basis of experimental results (e.g., the newly-reported NOvA result disfavors the maximal mixing scenario at a 2.6σ level). Consequently, we have performed a systematic study for the possible symmetry-breaking patterns and their implications for the mixing parameters.

We first define some parameters measuring the symmetry-breaking strengths and then derive an equation set relating them with the deviations of mixing parameters from the special values taken in the symmetry context. By solving these equations in both a numerical and analytical way, the sensitivity of mixing-parameter deviations to the symmetry-breaking parameters for various neutrino mass schemes and the Majorana-phase combinations are investigated in some detail. It turns out that $\Delta\theta$ is most sensitive to R_2 while $\Delta\delta$, $\Delta\rho$ and $\Delta\sigma$ to all the symmetry-breaking parameters. And the magnitudes of $\Delta\delta$, $\Delta\rho$ and $\Delta\sigma$ are generally much greater than that of $\Delta\theta$. This means that the symmetry-breaking pattern with a sizable R_2 will be favored if $\Delta\theta$ is notable, but any symmetry-breaking pattern may cause significant $\Delta\delta$, $\Delta\rho$ and $\Delta\sigma$. The mixing-parameter deviations tend to get remarkably magnified (a symmetry-breaking parameter of $\leq \mathcal{O}(0.01)$ may give rise to some mixing-parameter deviations of $\geq \mathcal{O}(0.1)$) in the case of $m_1 \simeq m_2 \simeq m_3$ combined with $\rho^{(0)} \neq \sigma^{(0)}$. But they will be highly suppressed and inconsiderable in the case of $m_1 \simeq m_2 \simeq m_3$ combined with $[\rho^{(0)}, \sigma^{(0)}] = [\pi/2, \pi/2]$. With these general results as guide, one may easily find an appropriate specific way to break the μ - τ reflection symmetry so as to generate the required mixing-parameter deviations when necessary. Finally, as a unique illustration, the general treatment is applied to the specific symmetry breaking induced by the RG running effect.

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