# Properties of holomorphic Wiener functions skeleton, contraction, and local Taylor expansion 

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Summary. We show that each holomorphic Wiener function has a skeleton which is intrinsic from several viewpoints. In particular, we study the topological aspects of the skeletons by using the local Taylor expansion for holomorphic Wiener functions.

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## 1 Introduction

Holomorphic functions of infinitely many complex variables have been discussed by many authors, among others by Shigekawa [7], who first constructed them on the complex Wiener space using the techniques of the Malliavin calculus. He noticed several properties of his holomorphic functions (let us call them holomorphic Wiener functions), such as the Itô-Wiener expansion in the $L^{p}$-sense. In this paper, we will investigate some other properties of holomorphic Wiener functions.

We use Shigekawa's framework and notation [7]. Let $(B, H, \mu, J)$ be an almost complex abstract Wiener space, i.e., $B$ is a real separable Banach space, $H$ is a real separable Hilbert space continuously and densely imbedded in $B$, $\mu$ is a Gaussian measure satisfying

$$
\int_{B} \exp (\sqrt{-1}\langle\varphi, z\rangle) \mu(d z)=\exp \left(-\frac{1}{4}\|\varphi\|_{H^{*}}^{2}\right), \quad \varphi \in B^{*} \subset H^{*},
$$

and $J: B \rightarrow B$, the almost complex structure, is an isometric mapping such that $J^{2}=-$ id and $\left.J\right|_{H}: H \rightarrow H$ is also isometric (see [7]).

[^0]In this paper, we will first show that from several viewpoints, we can endow each holomorphic function on ( $B, H, \mu, J$ ) with an intrinsic and unique skeleton, i.e., a function defined on $H$ which is considered to be the restriction of the original Wiener function. Note that, of course, since $\mu(H)=0$, it makes no measure-theoretical sense to endow general Wiener functions with skeletons.

Let $F: B \rightarrow \mathbf{C}$ be $L^{1+}$-holomorphic (for precise definition, see Definition 2.1 below). Note however that $F$ may not be continuous. Taking into account the mean value theorem for usual holomorphic functions on $\mathbf{C}^{n}$ and the rotation invariance of the Gaussian measure $\mu$, we may guess that the skeleton of $F$ should be

$$
F(h)=\int_{B} F(z+h) \mu(d z), \quad h \in H
$$

On the other hand, if we can give an intrinsic meaning to a function $F(\alpha z)$ for $\alpha>0$, which will be done for $0<\alpha<1$ and will be called the contraction operation, then the skeleton of $F$ should be

$$
F(h)=\lim _{\alpha \rightarrow 0} F(\alpha z+h) \text { in probability, } \quad h \in H .
$$

As expected, we can show that these two are consistent (Theorem 2.8).
Skeletons are expected to appear in the theory of large deviations as rate functions [1]. In infinite dimensional spaces, the theory of large deviations necessarily involves the topology of the spaces. Consequently, to investigate skeletons of Wiener functions, it is important to look at their fluctuations in small balls centered at each $h \in H$. In this context, we will show that

$$
\begin{equation*}
\lim _{r \rightarrow 0} \frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}|F(z+h)-F(h)|^{2} \mu(d z)=0, \quad h \in H \tag{1}
\end{equation*}
$$

for each $L^{2+}$-holomorphic function $F$, where $B_{r}$ denotes the centered $\|\cdot\|_{B^{-}}$ ball with radius $r>0$ (Theorem 4.1). But there, we must require the norm $\|\cdot\|_{B}$ of $B$ to have some good property which is fitted to the almost complex structure $J$.

In proving our theorems, we will make use of the rotation invariance of the Gaussian measure $\mu$. In fact, some results, such as the theorem of local Taylor expansion for holomorphic Wiener functions (Theorem 3.6), will be proved only by means of the rotation invariance.

## 2 Mean value theorem and contraction operation

First we define holomorphic functions on the almost complex abstract Wiener space ( $B, H, \mu, J$ ). We will review [7] briefly.

Let $B^{*}$ be the topological dual space of $B$ and let $B^{* C}$ be its complexification, i.e., $B^{* \mathrm{C}}:=B^{*} \oplus \sqrt{-1} B^{*}$. Defining

$$
\begin{aligned}
& B^{*(1,0)}=\left\{\varphi \in B^{* \mathrm{C}} \mid J^{*} \varphi=\sqrt{-1} \varphi\right\} \\
& B^{*(0,1)}:=\left\{\varphi \in B^{* \mathrm{C}} \mid J^{*} \varphi=-\sqrt{-1} \varphi\right\}
\end{aligned}
$$

we see that

$$
B^{* \mathrm{C}}=B^{*(1,0)} \oplus B^{*(0,1)} .
$$

The Hilbert spaces $H^{* \mathrm{C}}, H^{*(1,0)}$ and $H^{*(0,1)}$ are defined similarly.
A function $F: B \rightarrow \mathbf{C}$ is called a holomorphic polynomial, if it is expressed in the form

$$
F(z)=f\left(\left\langle\varphi_{1}, z\right\rangle, \ldots,\left\langle\varphi_{n}, z\right\rangle\right)
$$

where $n \in \mathbf{N}, f: \mathbf{C}^{n} \rightarrow \mathbf{C}$ is a polynomial with complex coefficients, and $\varphi_{1}, \ldots, \varphi_{n} \in B^{*(1,0)}$. The class of holomorphic polynomials is denoted by $\mathscr{P}_{h}$.

Definition 2.1 We define the space $\mathscr{H}^{p}(B, \mu)$ of $L^{p}$-holomorphic Wiener functions as the $L^{p}(B, \mu)$-closure of $\mathscr{P}_{h}$ :

$$
\mathscr{H}^{p}(B, \mu):=\overline{\mathscr{P}}_{h}^{I_{p}^{p}(B, \mu)}, \quad 1 \leqq p<\infty
$$

Further we define the auxiliary spaces by

$$
\mathscr{H}^{p+}(B, \mu):=\bigcup_{p<p^{\prime}} \mathscr{H}^{P^{\prime}}(B, \mu), \quad 1 \leqq p<\infty
$$

Note that the above definition of $\mathscr{H}^{p}(B, \mu)$ is equivalent to Shigekawa's definition (see [7, Proposition 4.2).

Since the Cameron-Martin density (see Definition 2.5 below) has all moments, it is easy to see that, for each $F \in \mathscr{H}^{p+}(B, \mu)$, the translated function $F(z+h)$ with $h \in H$ also belongs to $\mathscr{H}^{p^{+}}(B, \mu)$. Note that there are many $L^{p}$-holomorphic functions which are not continuous with respect to any measurable norm (see [10]; on measurable norms, see [2, 3]).

Since we are assuming the almost complex structure $J$, the measurable norm $\|\cdot\|_{B}$ is naturally required to have the following property.

Assumption 2.2 (Rotation invariance of the norm) We assume that, for arbitrary $a, b \in \mathbf{R}$,

$$
\|(a+b J) z\|_{B}=|a+\sqrt{-1} b|\|z\|_{B}, \quad z \in B .
$$

In particular, we assume

$$
\left\|e^{J \theta} z\right\|_{B}=\|z\|_{B}, \quad \theta \in \mathbf{R}
$$

where $e^{J \theta} z$ is an abbreviation for $(\cos \theta+(\sin \theta) J) z$.
Example 2.3 Let $\left(W\left(\mathbf{R}^{2}\right), H\left(\mathbf{R}^{2}\right), P^{W}\right)$ be the two-dimensional Wiener space, i.e.,

$$
\begin{aligned}
W\left(\mathbf{R}^{2}\right):= & \left\{w=\left(w^{1}, w^{2}\right) \in C\left([0,1] \rightarrow \mathbf{R}^{2}\right) \mid w(0)=0 \in \mathbf{R}^{2}\right\} \\
H\left(\mathbf{R}^{2}\right):=\{h= & \left(h^{1}, h^{2}\right) \in W \mid h(t) \text { is absolutely continuous and } \\
& \left.d h / d t \in L^{2}\left([0,1] \rightarrow \mathbf{R}^{2}, d t\right)\right\}
\end{aligned}
$$

$p^{W}:=$ the standard two-dimensional Wiener measure.

Then the abstract Wiener space $\left(W\left(\mathbf{R}^{2}\right), H\left(\mathbf{R}^{2}\right), P^{W}\right)$ is naturally equipped with an almost complex structure $J$, by the identification $\mathbf{R}^{2} \cong \mathbf{C}$. Namely, we define $J: W\left(\mathbf{R}^{2}\right) \rightarrow W\left(\mathbf{R}^{2}\right)$ by

$$
J w:=\left(-w^{2}, w^{1}\right), \quad w=\left(w^{1}, w^{2}\right) \in W\left(\mathbf{R}^{2}\right)
$$

Under this almost complex structure $J$, the following measurable norms are typical examples of rotation invariant ones:

$$
\begin{gathered}
\|w\|_{\infty}:=\max _{0 \leqq t \leqq 1}|w(t)|_{\mathbf{R}^{2}}, \\
\|w\|_{p}:=\left(\int_{0}^{1}|w(t)|_{\mathbb{R}^{2}}^{p} \mathrm{~d} t\right)^{1 / p}, \quad 1 \leqq p<\infty, \\
\|w\|_{(\alpha)}:=\sup _{0 \leqq s<t \leqq 1} \frac{|w(t)-w(s)|_{\mathbf{R}^{2}}}{|t-s|^{\alpha}}, \quad 0<\alpha<\frac{1}{2},
\end{gathered}
$$

where $|\cdot|_{\mathbf{R}^{2}}$ denotes the Euclid norm in $\mathbf{R}^{2} ;|(x, y)|_{\mathbf{R}^{2}}:=\sqrt{x^{2}+y^{2}}$.
Lemma 2.4 (Mean value theorem) Let $B_{r}$ be the centered $\|\cdot\|_{B}$-ball with radius $r>0$,

$$
B_{r}:=\left\{z \in B \mid\|z\|_{\boldsymbol{B}}<r\right\} .
$$

Then for each $F \in \mathscr{K}^{1}(B, \mu)$ and each $r>0$,

$$
\frac{1}{\mu\left(B_{r}\right)} \int_{B^{r}} F(z) \mu(d z)=\int_{B} F(z) \mu(d z)
$$

Proof. It is shown by Shigekawa [7, Theorem 4.1] that

$$
\int_{B} F(z) \mu(d z)=\frac{1}{2 \pi} \int_{0}^{2 \pi}\left(U_{\theta} F\right)(z) \mu(d \theta), \quad \mu \text {-a.e. } z \in B
$$

where $\left(U_{\theta} F\right)(z):=F\left(e^{J \theta} z\right)$, which is well-defined on account of the rotation invariance of $\mu$. The distribution of $U_{\theta} F$ is therefore just the same as that of $F$ for each $\theta \in \mathbf{R}$. It then follows from Fubini's theorem and the rotation invariance of $B_{r}$ that

$$
\begin{aligned}
\mu\left(B_{r}\right) \times \int_{B} F(z) \mu(d z) & =\int_{B_{r}} \mu(d z) \cdot \frac{1}{2 \pi} \int_{0}^{2 \pi}\left(U_{\theta} F\right)(z) d \theta \\
& =\frac{1}{2 \pi} \int_{0}^{2 \pi} d \theta \int_{B_{r}}\left(U_{\theta} F\right)(z) \mu(d z) \\
& =\frac{1}{2 \pi} \int_{0}^{2 \pi} d \theta \int_{B_{r}} F(z) \mu(d z) \\
& =\int_{B_{r}} F(z) \mu(d z)
\end{aligned}
$$

which completes the proof.

When $F \in \mathscr{H}^{1}(B, \mu)$ is continuous in $\|\cdot\|_{B}$, we have

$$
F(0)=\lim _{r \rightarrow 0} \frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}} F(z) \mu(d z)=\int_{B} F(z) \mu(d z) .
$$

Hence we present the following definition, which will be justified in the sequel.
Definition 2.5 For each $F \in \mathscr{H}^{1+}(B, \mu)$, we define its skeleton by

$$
F(h):=\int_{B} F(z+h) \mu(d z)=\int_{B} e^{2\langle h, z\rangle-|h|_{\bar{A}}^{2}} F(z) \mu(d z), \quad h \in H .
$$

Here note that the multiplier $e^{2\langle h, z\rangle-|h|_{n}^{2}}$ is the Cameron-Martin density.
Lemma 2.6 Let $\alpha \in \mathbf{R}$ be such that $0<\alpha<1$, and $F \in \mathscr{H}^{p}(B, \mu)$. If a sequence $\left\{F_{n}\right\}, F_{n} \in \mathscr{P}_{h}$, converges to $F$ in $\mathscr{H}^{p}(B, \mu)$, then $\left\{F_{n}(\alpha z)\right\}$ is also convergent in $\mathscr{H}^{p}(B, \mu)$. The limit does not depend on the choice of approximating sequence $\left\{F_{n}\right\}$ of holomorphic polynomials.
Proof. Since each holomorphic Fourier-Hermite function is in fact a monomial ( $[4,7]$ and expression (3) in the next section), we note that the Ornstein-Uhlenbeck semigroup $\left\{T_{t}\right\}_{t \geqq 0}$ (for details see [9] ${ }^{1}$ ) operates on $G \in \mathscr{P}_{h}$ as a family of contractions. Namely,

$$
\left(T_{t} G\right)(z)=G\left(e^{-t / 2} z\right), \quad z \in B \text { and } t \geqq 0
$$

Consequently, $F_{n}(\alpha z)=T_{-2 \log \alpha} F_{n}(z)$. Thus it is clear that $\left\{F_{n}(\alpha z)\right\}$ is convergent in $L^{p}(B, \mu)$, hence in $\mathscr{H}^{p}(B, \mu)$, and that the limit is $T_{-2 \log \alpha} F(z)$, which does not depend on the choice of approximating sequence.

By the above lemma, we may present the following definition of the contraction operation.

Definition 2.7 Let $\alpha \in \mathbf{R}$ be such that $0<\alpha<1$, and $F \in \mathscr{H}^{p+}(B, \mu)$. Then we define

$$
F(\alpha z):=T_{-2 \log \alpha} F(z) \in \mathscr{H}^{P}(B, \mu) .
$$

Let us show first that Definitions 2.5 and 2.7 are consistent.
Theorem 2.8 For any $F \in \mathscr{H}^{p+}(B, \mu)$,

$$
\lim _{\alpha \rightarrow 0} F(\alpha z+h)=F(h) \quad \text { in } \mathscr{H}^{p+}(B, \mu) \text { for each } h \in H
$$

Proof. For any $G \in L^{p}(B, \mu)$, we have

$$
\lim _{z \rightarrow \infty}\left\|T_{t} G-\int_{B} G(z) \mu(d z)\right\|_{L^{p}}=0
$$

which proves the claim.

[^1]
## 3 Local Taylor expansion

In this section, we will consider a localization of holomorphic Wiener functions.

We recall the Itô-Wiener expansion for $L^{2}$-holomorphic Wiener functions. According to [4, 7], the Hilbert space $\mathscr{H}^{2}(B, \mu)$ is decomposed into an orthogonal infinite direct sum as follows:

$$
\begin{equation*}
\mathscr{H}^{2}(B, \mu)=\bigoplus_{n=0}^{\infty} C_{(n, 0)} \tag{2}
\end{equation*}
$$

where $C_{(n, 0)}$ is the space of holomorphic $n$-fold Wiener integrals, i.e., the $L^{2}$-closure of the space of all the finite linear combinations of holomorphic Fourier-Hermite functions of degree $n$.

Recall that each holomorphic Fourier-Hermite function of degree $n$ is expressed in the form

$$
\begin{equation*}
\prod_{k=1}^{K}\left\langle\varphi_{k}, z\right\rangle^{m_{k}} \tag{3}
\end{equation*}
$$

where $\left\{\varphi_{k}\right\}_{k=1}^{K}$ is an orthonormal system of $H^{*(1,0)}$, while $m_{1}, \ldots, m_{K}$ are positive integers such that $\sum_{k=1}^{K} m_{k}=n$. Therefore we may call it a holomorphic monomial, and we may call the Itô-Wiener expansion (2) the global Taylor expansion in the $L^{2}$-sense.

In this section, we will present a local version of this expansion. As before, let $B_{r}$ be the centered $\|\cdot\|_{B}$-ball with radius $r>0$.

Lemma 3.1 Let $G_{1}$ and $G_{2}$ be two holomorphic monomials of different degrees. Then $\left.G_{1}\right|_{B_{r}}$ and $\left.G_{2}\right|_{B_{r}}$ are mutually orthogonal in $L^{2}\left(B_{r}, \mu\right)$.

Proof. Assume $G_{1}$ is of degree $m$ and $G_{2}$ is of degree $n$, where $m \neq n$. Then we obviously have

$$
\begin{aligned}
& U_{\theta} G_{1}=e^{\sqrt{-1} m \theta} G_{1}, \\
& U_{\theta} G_{2}=e^{\sqrt{-1} n \theta} G_{2},
\end{aligned}
$$

for $\theta \in \mathbf{R}$. Since $U_{\theta} B_{r}=B_{r}$ and the measure $\mu$ is $U_{\theta}$-invariant, we have

$$
\begin{aligned}
\int_{B_{r}} G_{1} \bar{G}_{2} d \mu & =\int_{B_{r}} U_{\theta}\left(G_{1} \bar{G}_{2}\right) d \mu \\
& =\int_{B_{r}} e^{\sqrt{-1} m \theta} e^{-\sqrt{-1} n \theta} G_{1} \bar{G}_{2} d \mu \\
& =e^{\sqrt{-1}(m-n) \theta} \int_{B_{r}} G_{1} \bar{G}_{2} d \mu,
\end{aligned}
$$

for arbitrary $\theta \in \mathbf{R}$. Thus we see that

$$
\int_{B_{r}} G_{1} \bar{G}_{2} d \mu=0
$$

Remark 3.2 Since $\mu$ is rotation invariant while the real part and the imaginary part of each holomorphic function are harmonic, Lemma 3.1 is essentially a consequence of the following fact: Two homogeneous harmonic polynomials on $\mathbf{R}^{n}$ of different degrees are mutually orthogonal with respect to the uniform measure on the sphere $S^{n-1}$ [8, p. 69].

In $\mathscr{H}^{2}(B, \mu)$, two holomorphic monomials of a same degree are mutually orthogonal, if they differ in one of the exponents $m_{k}$ in the expression (3). Now what about $\mathscr{H}^{2}\left(B_{r}, \mu\right)$ ? To answer this question, we need a class of good measurable norms.

First we introduce the notion of splitting of our almost complex abstract Wiener space (see, for example, [7] for details). Suppose that $\varphi \in H^{*(1,0)}$. Then decompose $\varphi=\psi+\sqrt{-1} \psi^{\prime}$ where $\psi, \psi^{\prime} \in H^{*} \cong H$. Let $H_{\varphi}$ be a two-dimensional subspace of $H$ spanned by $\psi$ and $\psi^{\prime}$, let $H_{\varphi}^{\perp}$ be its orthogonal complement, and let $B_{\varphi}^{\perp}$ be the closure of $H_{\varphi}^{\perp}$ in $B$. Accordingly, $\mu$ is split into $\mu=\mu_{\varphi} \otimes \mu_{\varphi}^{\perp}$, where $\mu_{\varphi}$ is a Gaussian measure on $H_{\varphi} \cong \mathbf{C}$, while $\mu_{\varphi}^{\perp}$ is also a Gaussian measure on $B_{\varphi}^{\perp}$. Now we have two almost complex abstract Wiener spaces $\left(H_{\varphi}, H_{\varphi}, \mu_{\varphi},\left.J\right|_{H_{\varphi}}\right)$ and ( $B_{\varphi}^{\perp}, H_{\varphi}^{\perp}, \mu_{\varphi}^{\perp},\left.J\right|_{\mathcal{B}_{\varphi}^{\perp}}$ ), and the following natural identification:

$$
\left(H_{\varphi} \oplus B_{\varphi}^{\perp}, H_{\varphi} \oplus H_{\varphi}^{\perp}, \mu_{\varphi} \otimes \mu_{\varphi}^{\perp},\left.\left.J\right|_{H_{\varphi}} \oplus J\right|_{B_{\varphi}^{\perp}}\right) \cong(B, H, \mu, J)
$$

Using this notation, we give a definition of the best fitted norms to the complex structure.

Definition 3.3 The norm $\|\cdot\|_{B}$ of $B$ is said to be completely rotation invariant, if it is rotation invariant and there exists a complete orthonormal system $\left\{\varphi_{k}\right\}$ of $H^{*(1,0)}$ such that for any $\varphi:=\varphi_{k}$,

$$
\left\|\left(e^{J \theta} z_{1}\right) \oplus z_{2}\right\|_{B}=\left\|z_{1} \oplus z_{2}\right\|_{B}, \quad z_{1} \in H_{\varphi}, \quad z_{2} \in B_{\varphi}^{\perp}, \quad \theta \in \mathbf{R}
$$

Example 3.4 Let $\left\{\varphi_{k}\right\}$ be a complete orthonormal system of $H^{*(1,0)}$. Put

$$
\begin{equation*}
\|z\|:=\left(\sum_{k} a_{k}\left|\left\langle\varphi_{k}, z\right\rangle\right|^{2}\right)^{1 / 2} \tag{4}
\end{equation*}
$$

where $a_{k}>0, \sum_{k} a_{k}<\infty$. Then $\|\cdot\|$ is measurable and completely rotation invariant. In the case of the two-dimensional Wiener space with almost complex structure $J$ defined in Example 2.3, the $L^{2}$-norm

$$
\|w\|_{2}=\left(\int_{0}^{1}|w(t)|_{\mathbf{R}^{2}}^{2} d t\right)^{1 / 2}
$$

has the above expression (4) and hence is completely rotation invariant. (In proving this fact, use the trigonometric series expansion for each $w$.)

Lemma 3.5 Let $\|\cdot\|_{B}$ be completely rotation invariant. Then two holomorphic monomials of a same degree are mutually orthogonal in $L^{2}\left(B_{r}, \mu\right)$, if they differ in one of the exponents $m_{k}$ in the expression (3).

Proof. Let $G_{1}$ and $G_{2}$ be holomorphic monomials of degree $n$, expressed as follows:

$$
\begin{array}{lll}
G_{1}(z)=\prod_{k=1}^{K}\left\langle\varphi_{k}, z\right\rangle^{m_{k}}, & m_{k} \geqq 0, & \sum_{k=1}^{K} m_{k}=n, \\
G_{2}(z)=\prod_{k=1}^{K}\left\langle\varphi_{k}, z\right\rangle^{m_{k}^{\prime}}, & m_{k}^{\prime} \geqq 0, & \sum_{k=1}^{K} m_{k}^{\prime}=n,
\end{array}
$$

where $m_{1} \neq m_{1}^{\prime}$. Then

$$
G_{1}(z) \overline{G_{2}(z)}=F_{1}(z) F_{2}(z)
$$

where $F_{1}$ and $F_{2}$ are given by

$$
\begin{aligned}
& F_{1}(z):=\left\langle\varphi_{1}, z\right\rangle^{m_{1}}{\overline{\left\langle\varphi_{1}, z\right\rangle^{m_{1}^{\prime}}}}^{F_{2}(z):=\prod_{k=2}^{K}\left\langle\varphi_{k}, z\right\rangle^{m_{k}}{\overline{\left\langle\varphi_{k}, z\right\rangle^{m_{k}^{\prime}}}}^{\text {. }}} .
\end{aligned}
$$

Putting $\varphi:=\varphi_{1}$, we have

$$
\begin{equation*}
\int_{B_{r}} G_{1}(z) \overline{G_{2}(z)} \mu(d z)=\int_{H_{\varphi}} \mu_{\varphi}\left(d z_{1}\right) F_{1}\left(z_{1}\right) \int_{B_{r}\left(z_{1}\right)} F_{2}\left(z_{2}\right) \mu_{\varphi}^{\perp}\left(d z_{2}\right), \tag{5}
\end{equation*}
$$

where $B_{r}\left(z_{1}\right):=\left\{z_{2} \in B_{\varphi}^{\perp} \mid z_{1} \oplus z_{2} \in B_{r}\right\}, z_{1} \in H_{\varphi}$. For $\theta \in \mathbf{R}$, note that

$$
F_{1}\left(e^{\jmath \theta} z_{1}\right)=e^{\sqrt{-1}\left(m_{1}-m_{1}^{\prime}\right) \theta} F_{1}\left(z_{1}\right), \quad z_{1} \in H_{\varphi}
$$

Hereafter we identify $H_{\varphi}$ with $\mathbf{C}$. Then since the complete rotation invariance of the norm implies $B_{r}\left(e^{J \theta} z_{1}\right)=B_{r}\left(z_{1}\right)$, we have

$$
\begin{aligned}
& \int_{\arg z_{1} \in[\theta, \theta+\eta)} \mu_{\varphi}\left(d z_{1}\right) F_{1}\left(z_{1}\right) \int_{B_{r}\left(z_{1}\right)} F_{2}\left(z_{2}\right) \mu_{\varphi}^{\perp}\left(d z_{2}\right) \\
= & \int_{\arg z_{1} \in[0, \eta)} \mu_{\varphi}\left(d z_{1}\right) F_{1}\left(e^{J \theta} z_{1}\right) \\
= & \int_{B_{r}\left(e^{J z_{1}}\right)} F_{2}\left(z_{2}\right) \mu_{\varphi}^{\perp}\left(d z_{2}\right) \\
= & \int_{\arg z_{1} \in[0, \eta)} \mu_{\varphi}\left(d z_{1}\right) e^{\sqrt{-1}\left(m_{1}-m_{1}^{\prime}\right) \theta} F_{1}\left(z_{1}\right) \int_{B_{r}\left(z_{1}\right)} F_{2}\left(z_{2}\right) \mu_{\varphi}^{\perp}\left(d z_{2}\right) \\
= & \int_{\arg z_{1} \in[0, \eta)} \mu_{\varphi}\left(d z_{1}\right) F_{1}\left(z_{1}\right) \int_{B_{r}\left(z_{1}\right)} F_{2}\left(z_{2}\right) \mu_{\varphi}^{\perp}\left(d z_{2}\right) .
\end{aligned}
$$

Consequently, putting $\eta=2 \pi$, we get

$$
\int_{B_{r}} G_{1}(z) \overline{G_{2}(z)} \mu(d z)=e^{\sqrt{-1}\left(m_{1}-m_{i}^{\prime}\right) \theta} \int_{B_{r}} G_{1}(z) \overline{G_{2}(z)} \mu(d z)
$$

for every $\theta$, which shows this value should be equal to zero.
We define the space of $L^{p}$-holomorphic functions defined on $B_{r}$ as follows.
Definition 3.6 By $\mathscr{H}^{p}\left(B_{r}, \mu\right)$, we denote the $L^{p}\left(B_{r}, \mu\right)$-closure of $\left.\mathscr{P}_{h}\right|_{B_{r}}$ :

$$
\mathscr{H}^{p}\left(B_{r}, \mu\right):={\overline{\left.\mathscr{P}_{h}\right|_{B r}}}^{L^{p}\left(\boldsymbol{B}_{r}, \mu\right)}, \quad 1 \leqq p<\infty,
$$

and we define the auxiliary spaces by

$$
\mathscr{H}^{p^{+}}\left(B_{r}, \mu\right):=\bigcup_{p<p^{\prime}} \mathscr{H}^{p^{\prime}}\left(B_{r}, \mu\right), \quad 1 \leqq p<\infty .
$$

We have thus defined $L^{p}$-local holomorphic functions for each $p$ ( $1 \leqq p<\infty$ ), but we will exclusively deal with the case $p=2$. The next theorem readily follows from Lemmas 3.1 and 3.5.

Theorem 3.7 The Hilbert space $\mathscr{H}^{2}\left(B_{r}, \mu\right)$ is decomposed into an infinite orthogonal direct sum

$$
\begin{equation*}
\mathscr{H}^{2}\left(B_{r}, \mu\right)=\oplus_{n=0}^{\infty} C_{(n, 0)}\left(B_{r}\right), \tag{6}
\end{equation*}
$$

where $C_{(n, 0)}\left(B_{r}\right)$ is the $L^{2}\left(B_{r}, \mu\right)$-closure of all the finite linear combinations of holomorphic monomials of degree n. If, in addition, the given measurable norm is completely rotation invariant, each $F \in C_{(n, 0)}\left(B_{r}\right)$ is orthogonally decomposed into an infinite linear combination of holomorphic monomials of degree $n$.

By Shigekawa's uniqueness theorem [7, Theorem 4.3], a mapping

$$
\left.\mathscr{H}^{2}(B, \mu) \ni F \mapsto F\right|_{B_{r}} \in \mathscr{H}^{2}\left(B_{r}, \mu\right)
$$

is injective, so $\mathscr{H}^{2}(B, \mu)$ is continuously imbedded in $\mathscr{H}^{2}\left(B_{r}, \mu\right)$. Of course, there exist many functions in $\mathscr{H}^{2}\left(B_{r}, \mu\right)$ which cannot be extended to the whole space $B$ as elements of $\mathscr{H}^{2}(B, \mu)$. Since (6) is considered as a local Taylor expansion, we may say that $\mathscr{H}^{2}\left(B_{r}, \mu\right)$ is the space of local $L^{2}$ holomorphic functions with radius of convergence at least $r$.

To conclude this section, we present a corollary which is an immediate consequence of Theorem 3.7.

Corollary 3.8 (i) For $F \in \mathscr{H}^{2}(B, \mu)$,

$$
\int_{B_{r}}|F(z)|^{2} \mu(d z)=\sum_{n=0}^{\infty} \int_{B_{r}}\left|J_{(n, 0)} F(z)\right|^{2} \mu(d z),
$$

where $J_{(n, 0)}$ is the orthogonal projection from $\mathscr{H}^{2}(B, \mu)$ onto $C_{(n, 0)}$.
(ii) Suppose that the norm $\|\cdot\|_{B}$ is completely rotation invariant with respect to a CONS $\left\{\varphi_{k}\right\} \subset H^{*(1,0)}$ and that $F \in C_{(n, 0)}$ has an expression

$$
F=\sum_{k=0}^{\infty} a_{k} G_{k},
$$

where $\left\{a_{k}\right\} \in l^{2}$, while $\left\{G_{k}\right\}$ is a sequence of holomorphic monomials of degree $n$ which are based on $\left\{\varphi_{k}\right\}$, normalized and mutually orthogonal in $\mathscr{H}^{2}(B, \mu)$. Then we have

$$
\int_{B_{r}}|F(z)|^{2} \mu(d z)=\sum_{k=0}^{\infty}\left|a_{k}\right|^{2} \int_{B_{r}}\left|G_{k}(z)\right|^{2} \mu(d z) .
$$

Remark 3.9 It is clearly possible to give a local version of Corollary 3.8 as well as of the theorem in the next section.
Remark 3.10 In [5], another type of localization of holomorphic Wiener functions is given in connection with SDEs with holomorphic coefficients.

## 4 Topological aspects of skeletons

In this section, we will consider the topological aspects of the skeletons of holomorphic Wiener functions.

Let $F \in \mathscr{H}^{1+}(B, \mu)$. As we mentioned in Definition 2.5, the skeleton of $F$ is the mean of $F(z+h)$, which is equal to the local mean restricted to the centered ball $B_{r}$ of arbitrary radius $r>0$ (Lemma 2.4). Here we will show that the fluctuation of $F(z+h)$ around the skeleton $F(h)$ in $B_{r}$ decreases as the radius $r$ tends to zero. Namely, we will show the following theorem.

Theorem 4.1 Let $\|\cdot\|_{B}$ be completely rotation invariant and let $B_{r}$ be the centered $\|\cdot\|_{B^{-}}$-ball with radius $r>0$. Then for each $F \in \mathscr{H}^{2+}(B, \mu)$,

$$
\begin{equation*}
\lim _{r \rightarrow 0} \frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}|F(z+h)-F(h)|^{2} \mu(d z)=0, \quad h \in H \tag{7}
\end{equation*}
$$

In particular,

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \mu\left(|F(z)-F(h)|>\varepsilon \mid\|z-h\|_{B}<\delta\right)=0, \quad \varepsilon>0 \tag{8}
\end{equation*}
$$

Lemma 4.2 Let $\|\cdot\|_{B}$ be completely rotation invariant and let $B_{r}$ be the centered $\|\cdot\|_{B^{-}}$-ball with radius $r>0$. Then for each $F \in \mathscr{H}^{2}(B, \mu)$,

$$
\frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}|F(z)|^{2} \mu(d z) \leqq \int_{B}|F(z)|^{2} \mu(d z) .
$$

Proof. By Corollary 3.8 (i), we have

$$
\frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}|F(z)|^{2} \mu(d z)=\sum_{n=0}^{\infty} \frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}\left|J_{(n, 0)} F(z)\right|^{2} \mu(d z),
$$

and hence it is sufficient to show that, for each $n=0,1, \ldots$,

$$
\frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}\left|J_{(n, 0)} F(z)\right|^{2} \mu(d z) \leqq \int_{B}\left|J_{(n, 0)} F(z)\right|^{2} \mu(d z) .
$$

By Corollary $3.8($ ii), we have only to show that

$$
\begin{equation*}
\frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}\left|G_{k}(z)\right|^{2} \mu(d z) \leqq \int_{B}\left|G_{k}(z)\right|^{2} \mu(d z) \tag{9}
\end{equation*}
$$

for each holomorphic monomial $G_{k}$ which appeared in Corollary 3.8 (ii). Since $G_{k}$ is a monomial, $\left|G_{k}(z)\right|^{2}$ increases as $\|z\|_{B}$ increases, so (9) is intuitively obvious. For a rigorous proof, see Appendix.

Before proving Theorem 4.1, it is important to notice that, if the norm is not rotation invariant, the assertion of Theorem 4.1 may be false. Indeed, by using the same method as in [11], we can construct a holomorphic Wiener function $F$ and a measurable norm $\|\cdot\|_{B}$ without rotation invariance for which (8) does not hold. However we do not know whether "complete rotation invariance" is necessary or not. Indeed, if Lemma 4.2 holds for $\|\cdot\|_{B}$ which is not completely rotation invariant but only rotation invariant, then Theorem 4.1 would hold without complete rotation invariance of the norm.

Proof of Theorem 4.1 We will show (7) for $h=0$ only. For arbitrary $h \in H,(7)$ and (8) are proved by using the Cameron-Martin density (cf. [11]).
Take an arbitrary $\varepsilon>0$. Since

$$
\int_{B}\{F(z)-F(0)\} \mu(d z)=0,
$$

there exists a continuous holomorphic polynomial $G$ such that

$$
\begin{gathered}
\int_{B} G(z) \mu(d z)=0, \\
\int_{B}|F(z)-F(0)-G(z)|^{2} \mu(d z)<\frac{1}{4} \varepsilon .
\end{gathered}
$$

Applying Lemma 4.2, we see that

$$
\begin{aligned}
& \frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}|F(z)-F(0)|^{2} \mu(d z) \\
& \quad \leqq \frac{2}{\mu\left(B_{r}\right)_{B_{r}}} \int_{B_{r}}|G(z)|^{2} \mu(d z)+\frac{2}{\mu\left(B_{r}\right)_{B_{r}}} \int_{|F(z)-F(0)-G(z)|^{2} \mu(d z)} \quad \leqq \frac{2}{\mu\left(B_{r}\right)_{B_{r}}} \int|G(z)|^{2} \mu(d z)+2 \int_{B}|F(z)-F(0)-G(z)|^{2} \mu(d z) \\
& \quad \leqq \frac{2}{\mu\left(B_{r}\right)} \int_{B_{r}}|G(z)|^{2} \mu(d z)+\frac{1}{2} \varepsilon .
\end{aligned}
$$

Since $G$ is continuous and

$$
G(0)=\int_{B} G(z) \mu(d z)=0,
$$

there exists an $r_{0}>0$ such that for $0<r<r_{0}$,

$$
\frac{2}{\mu\left(B_{r}\right)} \int_{B_{r}}|G(z)|^{2} \mu(d z)<\frac{1}{2} \varepsilon,
$$

which implies

$$
\frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}|F(z)-F(0)|^{2} \mu(d z)<\varepsilon .
$$

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## Appendix: Proof of Lemma 4.2

Here we will present a rigorous proof for (9) to complete the proof of Lemma 4.2.

Let $\|\cdot\|_{B}$ be completely rotation invariant, and let $\left\{\varphi_{k}\right\}$ be the corresponding complete orthonormal system of $H^{*(1,0)}$, as in Definition 3.2.

Lemma A1 Let $\varphi$ be one of the $\varphi_{k}$ 's. Then

$$
\left\|\left(a z_{1}\right) \oplus z_{2}\right\|_{B}, \quad z_{1} \in H_{\varphi}, \quad z_{2} \in B_{\varphi}^{\frac{1}{\varphi}},
$$

is non-decreasing as $a>0$ increases.
Proof. It is sufficient to show that

$$
\left\|\left(a z_{1}\right) \oplus z_{2}\right\|_{B} \geqq\left\|z_{1} \oplus z_{2}\right\|_{B}=: r, \quad a>1
$$

Assume that this is not true for some $a>1$. Then $\left(a z_{1}\right) \oplus z_{2}$ must be an interior point of a ball $B_{r}$. Note that $0 \oplus z_{2} \in B_{r}$, because $\left(-z_{1}\right) \oplus z_{2} \in B_{r}$ by the complete rotation invariance of $\|\cdot\|_{B}$ and $0 \oplus z_{2}=(1 / 2)\left\{z_{1} \oplus z_{2}+\left(-z_{1}\right) \oplus\right.$ $\left.z_{2}\right\}$. Since $z_{1} \oplus z_{2}$ can be expressed as a convex combination of $0 \oplus z_{2} \in B_{r}$ and $\left(a z_{1}\right) \oplus z_{2}$, which is an interior point of $B_{r}, z_{1} \oplus z_{2}$ would be an interior point of $B_{r}$. This is a contradiction.

For the moment, assume $B$ to be of finite dimension; $B \cong \mathbf{C}^{N}$, and suppose that $G$ has an expression

$$
G\left(z_{1}, \ldots, z_{K}\right)=\prod_{k=1}^{K} z_{k}^{m_{k}}, \quad\left(z_{1}, \ldots, z_{K}\right) \in \mathbf{C}^{K},
$$

where $K \leqq N$. Letting $\mu\left(d z_{1} \ldots d z_{N}\right)$ be the standard Gaussian measure on $\mathbf{C}^{N}$, we may ask whether
$\frac{1}{\mu\left(B_{r}\right)} \int_{B_{r}}\left|G\left(z_{1}, \ldots, z_{K}\right)\right|^{2} \mu\left(d z_{1} \ldots d z_{N}\right) \leqq \int_{C^{N}}\left|G\left(z_{1}, \ldots, z_{K}\right)\right|^{2} \mu\left(d z_{1} \ldots d z_{N}\right)$.
Here $B_{r}$ denotes a centered ball of radius $r>0$ with respect to a completely rotation invariant norm.

In the following lemma, we will state an assertion under the identification $\mathbf{C}^{N} \cong \mathbf{R}^{2 N}$, but using the same symbols.

Lemma A2 Let $F: \mathbf{R}_{+}^{2 N}:=[0, \infty)^{2 N} \rightarrow \mathbf{R}$ be non-decreasing in the following sense:

$$
F\left(x_{1}, \ldots, x_{2 N}\right) \leqq F\left(y_{1}, \ldots, y_{2 N}\right), \quad 0 \leqq x_{k} \leqq y_{k}, \quad k=1, \ldots, 2 N
$$

Then putting $B_{r}^{+}:=B_{r} \cap \mathbf{R}_{+}^{2 N}, r>0$, we have

$$
\frac{1}{\mu\left(B_{r}^{+}\right)} \int_{B_{r}^{+}} F\left(x_{1}, \ldots, x_{2 N}\right) \mu\left(d x_{1} \ldots d x_{2 N}\right) \leqq \int_{\mathbf{R}_{+}^{2 N}} F\left(x_{1}, \ldots, x_{2 N}\right) \mu\left(d x_{1} \ldots d x_{2 N}\right)
$$

Proof. Define a function $\rho_{r}(x), x \in \mathbf{R}_{+}^{2 N}, 0<r \leqq \infty$, by

$$
\begin{aligned}
& \rho_{r}(x):=\frac{1}{\mu\left(B_{r}^{+}\right)} \mathbf{1}_{B_{r}^{+}}(x) g(x), \quad 0<r<\infty, \\
& \rho_{\infty}(x):=g(x)
\end{aligned}
$$

where

$$
g(x):=(2 \pi)^{-N} \exp \left(-\sum_{k=1}^{2 N} x_{k}^{2} / 2\right), \quad x=\left(x_{1}, \ldots, x_{2 N}\right) .
$$

Let $0<r$, and define $x \vee y$ and $x \wedge y, x, y \in \mathbf{R}_{+}^{2 N}$, by

$$
\begin{aligned}
& x \vee y:=\left(x_{1} \vee y_{1}, \ldots, x_{2 N} \vee y_{2 N}\right), \\
& x \wedge y:=\left(x_{1} \wedge y_{1}, \ldots, x_{2 N} \wedge y_{2 N}\right) .
\end{aligned}
$$

Thus the $x_{k}$ and $y_{k}$ are the components of $x$ and $y$, respectively. Lemma A1 implies that the indicator function $\mathbf{1}_{B_{r}^{+}}$has the property

$$
\mathbf{1}_{B_{r}^{+}}(x \wedge y) \geqq \mathbf{1}_{B_{r}^{*}}(y) .
$$

On the other hand, we readily see that

$$
g(x \vee y) g(x \wedge y)=g(x) g(y) .
$$

Hence it is easy to see that

$$
\rho_{\infty}(x \vee y) \rho_{r}(x \wedge y) \geqq \rho_{\infty}(x) \rho_{r}(y) .
$$

Therefore it follows from a version of the FKG-inequality due to Preston [6, Theorem 3] that

$$
\int_{\mathbf{R}_{+}^{2 N}} F(x) \rho_{r}(x) \mu(d x) \leqq \int_{\mathbf{R}_{+}^{2 N}} F(x) \rho_{\infty}(x) \mu(d x),
$$

which completes the proof.
By Lemma A2, it follows easily that (10) holds. Finally, letting the dimension $N$ tend to infinity, we obtain (9). Thus the proof of Lemma 4.2 is complete.

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[^1]:    ${ }^{1}$ Our $T_{t}$ here corresponds to $T_{t / 2}$ in [9]

